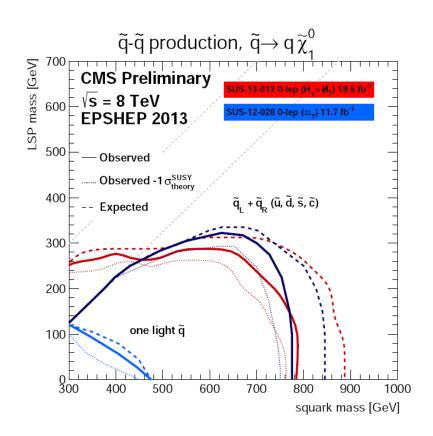
MITP workshop, Mainz July 17th 2014

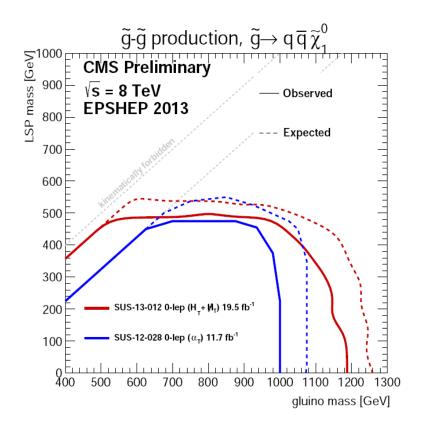
Probing light Neutralino Dark Matter at the LHC

Lorenzo Calibbi
ULB

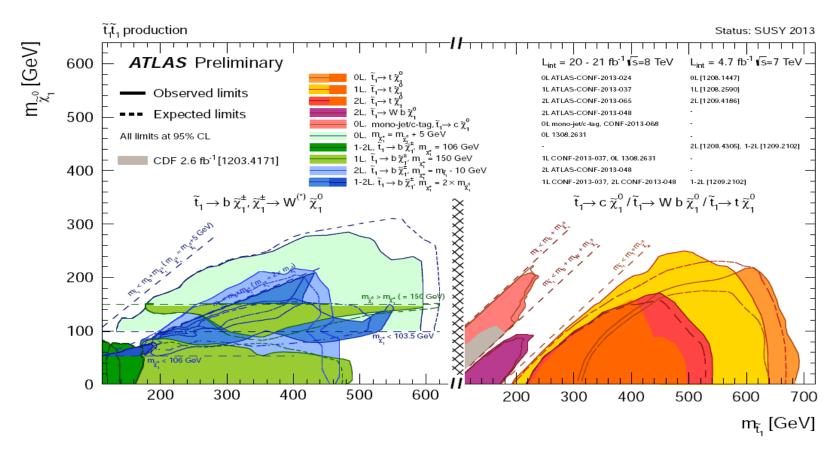


based on collaborations with J. Lindert, T. Ota, Y. Takanishi

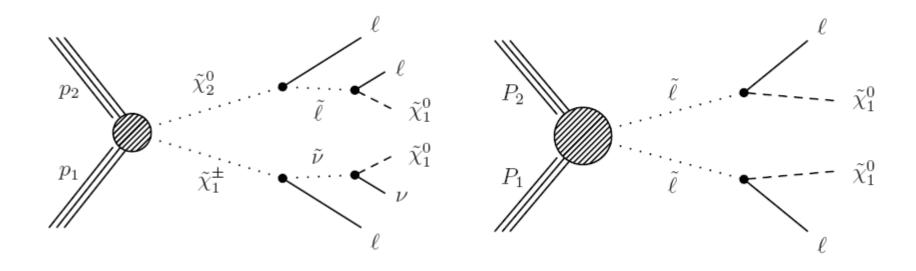




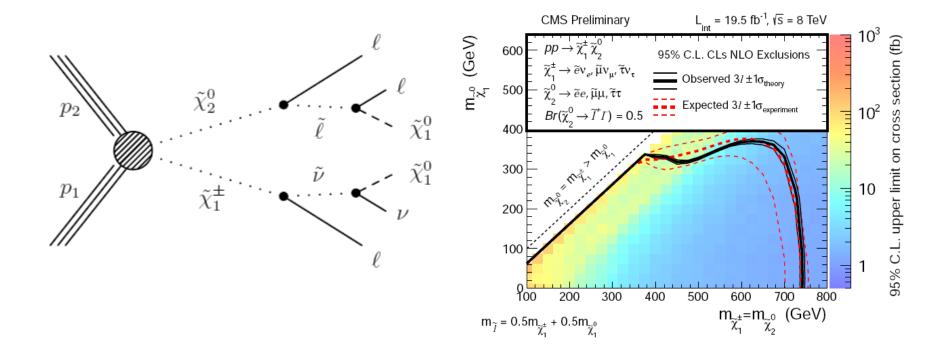
Strongly-interacting SUSY particles (gluinos, 1st generation squarks) typically excluded by the LHC above ~1 TeV



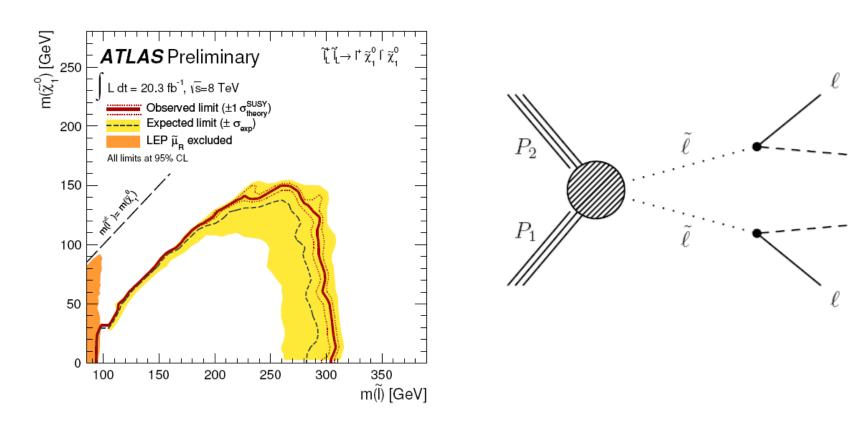
Strongly-interacting SUSY particles (gluinos, 1st generation squarks) typically excluded by the LHC above ~1 TeV



Direct bounds on EW-iteracting particles relatively weaker (smaller production) EW-searches at the LHC started to go considerably beyond the limits set by LEP



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The EW SUSY sector might be lighter than the strong sector but LHC can test it!

We discuss an example of the potential of EW searches at the LHC: probing light Neutralino Dark Matter

Assumptions:

- Only MSSM superfields
- R-parity
- Dark Matter thermal relic, standard history of the universe
- Neutralino DM candidate, $\Omega_{\rm DM}h^2 \leq 0.124$

How light the neutralino is allowed to be after LHC searches at 8 TeV?

Light Neutralino Dark Matter in the MSSM

MSSM neutralinos:

$$\left(\tilde{B}, \tilde{W}_3, \tilde{H}_d^0, \tilde{H}_u^0\right)$$

MSSM charginos: $(\tilde{W}^{\pm}, \tilde{H}_{u}^{+}, \tilde{H}_{J}^{+})$

$$(\tilde{W}^{\pm},\,\tilde{H}_u^+,\,\tilde{H}_d^+)$$

$$\mathbf{M}_{\widetilde{N}} = \begin{pmatrix} M_1 & 0 & -c_{\beta} s_W m_Z & s_{\beta} s_W m_Z \\ 0 & M_2 & c_{\beta} c_W m_Z & -s_{\beta} c_W m_Z \\ -c_{\beta} s_W m_Z & c_{\beta} c_W m_Z & 0 & -\mu \\ s_{\beta} s_W m_Z & -s_{\beta} c_W m_Z & -\mu & 0 \end{pmatrix} \quad \mathcal{M}_{\pm} = \begin{pmatrix} M_2 & \sqrt{2} M_W \sin\beta \\ \sqrt{2} M_W \cos\beta & \mu \end{pmatrix}$$

Mass eigenstates:

LEP chargino searches: $m_{\widetilde{\chi}^{\pm}} > 103.5 \; \mathrm{GeV}$

$$\widetilde{\chi}_{i}^{0} = N_{i1}\widetilde{B} + N_{i2}\widetilde{W} + N_{i3}\widetilde{H}_{d} + N_{i4}\widetilde{H}_{u}$$
 $M_{2}, \ \mu \gtrsim 100 \text{ GeV}$

$$m_{\tilde{\chi}_1^0} \lesssim 30 \text{ GeV} \implies M_1 \ll M_2, \mu \iff \tilde{\chi}_1^0 \approx \tilde{B}$$

No gaugino mass unification at the GUT scale ($M_1 \simeq 0.5 \times M_2$)

Light Neutralino Dark Matter in the MSSM

MSSM neutralinos:

$$\left(\tilde{B}, \tilde{W}_3, \tilde{H}_d^0, \tilde{H}_u^0\right)$$

MSSM charginos: $(\tilde{W}^{\pm}, \tilde{H}_{u}^{+}, \tilde{H}_{d}^{+})$

$$(\tilde{W}^\pm,\,\tilde{H}_u^+,\,\tilde{H}_d^+)$$

$$\mathbf{M}_{\widetilde{N}} = \begin{pmatrix} M_1 & 0 & -c_{\beta} s_W m_Z & s_{\beta} s_W m_Z \\ 0 & M_2 & c_{\beta} c_W m_Z & -s_{\beta} c_W m_Z \\ -c_{\beta} s_W m_Z & c_{\beta} c_W m_Z & 0 & -\mu \\ s_{\beta} s_W m_Z & -s_{\beta} c_W m_Z & -\mu & 0 \end{pmatrix} \quad \mathcal{M}_{\pm} = \begin{pmatrix} M_2 & \sqrt{2} M_W \sin\beta \\ \sqrt{2} M_W \cos\beta & \mu \end{pmatrix}$$

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Relic density (WMAP, Planck):
$$\Omega_{\rm Y} h^2 \sim 0.1$$

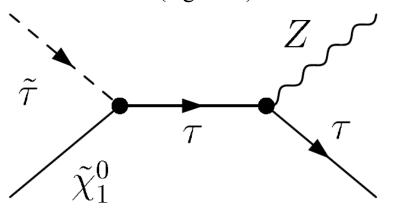
Efficient annihilation required

Annihilation mechanisms

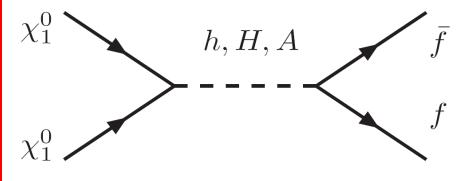
t-channel exchange (e.g. sfermions):

 χ_1 \tilde{f} χ_1^0

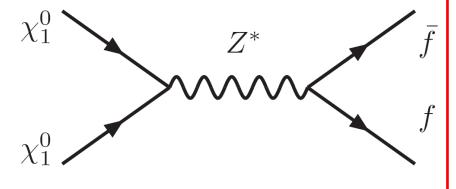
Coannihilations (e.g. stau):



Higgs mediation:



Z mediation:

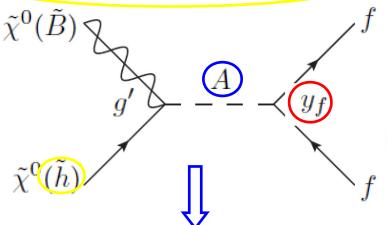


 \Longrightarrow

different regions of the SUSY parameter space are selected

Annihilation mechanisms: Higgs exchange

Higgs mediated annihilation:



$$\propto \frac{m_{\tilde{\chi}_1^0}^2}{m_A^2} \frac{m_f}{v} (N_{11}N_{13,14}) \tan \beta$$

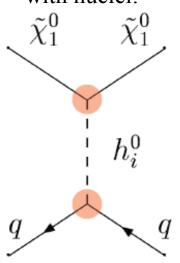
Drees Nojiri '92

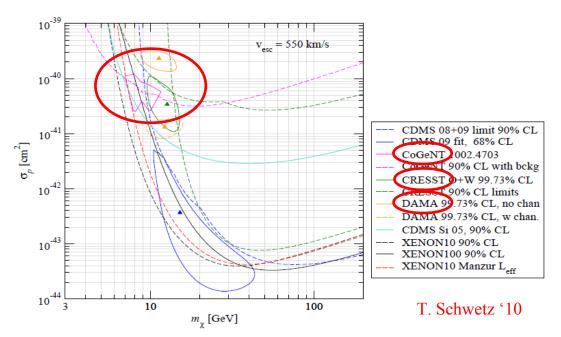
Small $m_{\tilde{\chi}_1^0} \iff \text{large } \tan \beta, \text{ small } m_A \text{ (and } \mu)$

 $m_A \approx 100 \div 120 \,\mathrm{GeV}$

Bottino et al. '02-'10

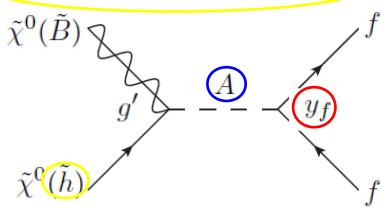
Large scattering cross-section with nuclei:





Annihilation mechanisms: Higgs exchange

Higgs mediated annihilation:



$$\propto \frac{m_{\tilde{\chi}_1^0}^2}{m_A^2} \frac{m_f}{v} (N_{11}N_{13,14}) \tan \beta$$

Drees Nojiri '92

Small $m_{\tilde{\chi}_1^0} \iff \text{large } \tan \beta, \text{ small } m_A \text{ (and } \mu)$

 $m_A \approx 100 \div 120 \, \mathrm{GeV}$

Bottino et al. '02- '10

This parameter space for light neutralinos is now excluded by:

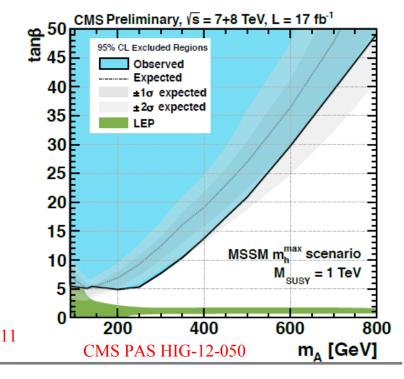
•
$$B_s \to \mu^+ \mu^-$$

• Searches for extra Higgses at the LHC, $pp \rightarrow X \Phi \rightarrow \tau \tau$

Higgs mass and couplings

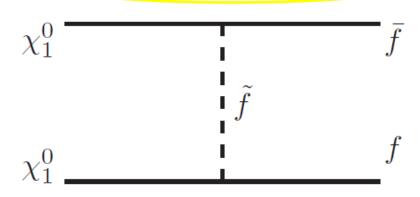
MSSM incompatible with DAMA, CoGeNT

LC Ota Takanishi '11



Annihilation mechanisms: sfermion exchange

Sfermion mediated annihilation:



LEP limits: $e^+e^- \to \tilde{f}\tilde{f}^* \to f\bar{f}\tilde{\chi}_1^0\tilde{\chi}_1^0$ Can one evade it with compressed spectra?

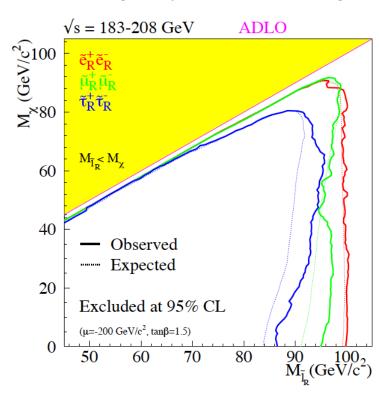
If $m_{\tilde{\chi}_1^0} \lesssim 30 \text{ GeV}$ this is not possible because Z width bounds:

$$Z \to \tilde{f}\tilde{f}^* \implies m_{\tilde{f}} \gtrsim 40 \,\mathrm{GeV}$$

(Possible exception: light sbottom with tuned LR mixing) Arbey Battaglia Mahmoudi '13

$$\propto \frac{m_{\tilde{\chi}_1^0} m_f}{m_{\tilde{f}}^2}$$

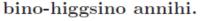
Challenged by LEP bounds, e.g.:

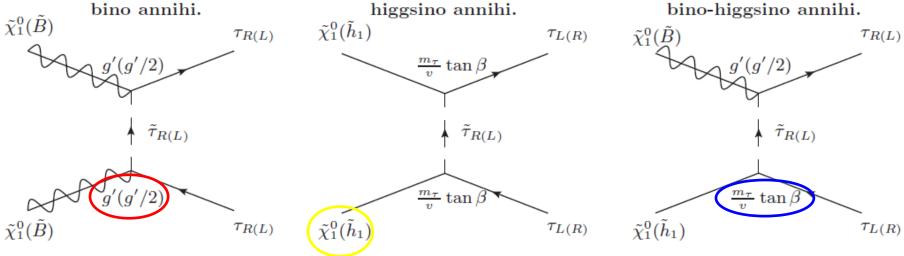


Light neutralino DM possible in presence of a light stau

Albornoz Bélanger Boehm '11

Grothaus Lindner Takanishi '12





- RH stau much more efficient (cross-section 16x larger than LH one)
- Sizeable higgsino component: small μ
- Yukawa interactions: large tanβ

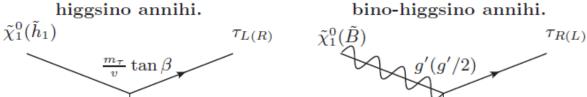
Relic density essentially controlled by 4 parameters only:

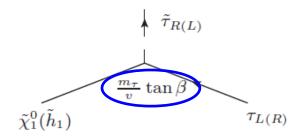
$$M_1, m_{\tilde{\tau}_R}, \mu, \tan \beta$$

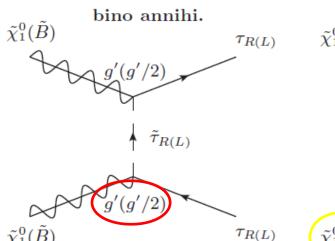
Light neutralino DM possible in presence of a light stau

Albornoz Bélanger Boehm '11

Grothaus Lindner Takanishi '12







 $\widetilde{\chi}_{1}^{0}(\widetilde{h}_{1}) \qquad \qquad \tau_{L(R)}$

 $\tilde{\tau}_{R(L)}$

- RH stau much more efficient (cross-section 16x larger than LH one)
- Sizeable higgsino component: small μ
- Yukawa interactions: large tanβ

Light neutralino DM necessarily implies:

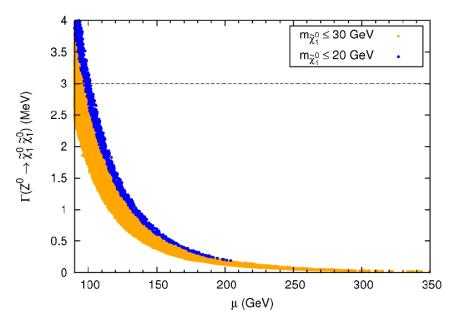
light stau, light higgsino-like neutralinos and charginos

Parameters scan and constraints

SuSpect, micrOMEGAs Djouadi et al. '02 Belanger et al. '06

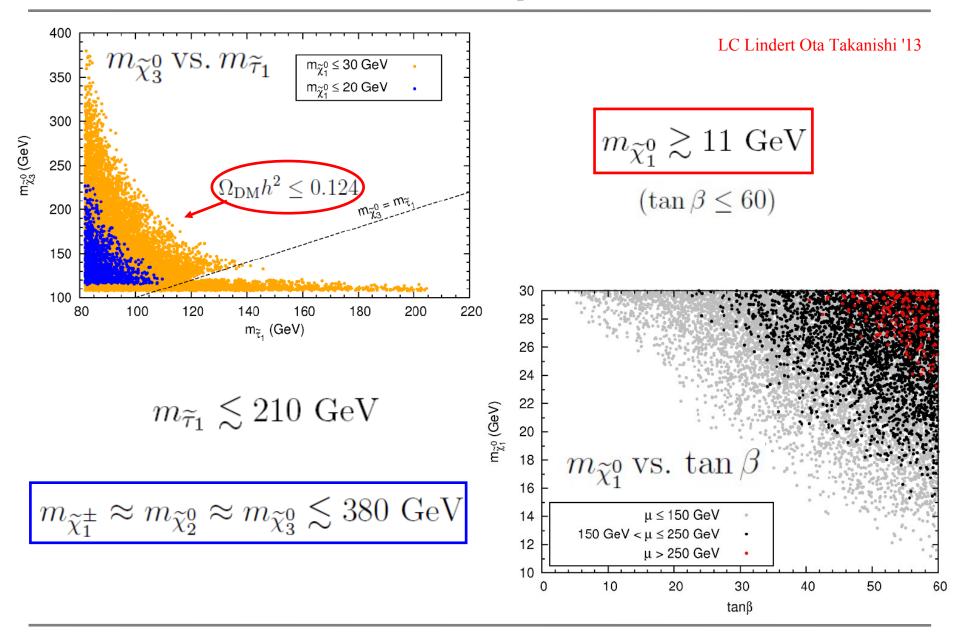
- $m_{\widetilde{\chi}_1^0} \leq 30 \text{ GeV}$
- CMB, Planck (3 σ): $\Omega_{\rm DM} h^2 \leq 0.124$
- LEP2: $m_{\tilde{\tau}_R} \ge 81.9 \,\text{GeV}, \quad m_{\tilde{\chi}_1^{\pm}} \ge 103.5 \,\text{GeV}$
- LHC: limits on charginos depend on smuon/selectron masses and can be evaded
- ullet Flavour: $\Omega_{\rm DM}$ does not depend on heavy Higgs/squark masses, flavour observables do not constrain the DM parameter space
- Z invisible width, LEP: $\Delta\Gamma_Z^{\rm inv} < 3~{
 m MeV}$

LC Lindert Ota Takanishi '13

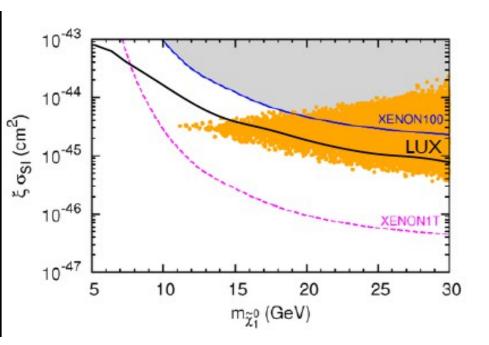


$$\Gamma(Z \to \widetilde{\chi}_1^0 \widetilde{\chi}_1^0) = \frac{G_F}{\sqrt{2}} \frac{M_Z^3}{12\pi} \left(1 - \frac{4M_{\widetilde{\chi}_1^0}^2}{M_Z^2} \right)^{\frac{3}{2}} \left[N_{13}^2 - N_{14}^2 \right]^2$$

Parameter space



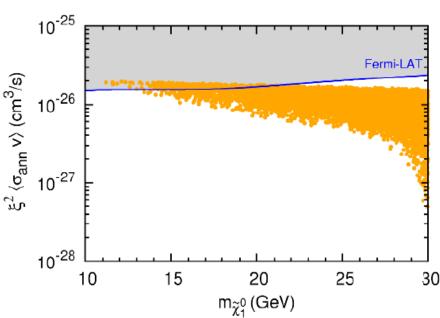
Spin-independent scattering cross section



Th. prediction can be lowered by 2x by uncert. in light quark masses/form factors

Heavy Higgs slow decoupling (possible additional ~3x suppression)

Annihilation cross-section



Gamma rays from satellite galaxies

LHC phenomenology: production and decays

Relic density constr. imply that we have at least 4 states at O(100) GeV:

$$\widetilde{\tau}_1, \ \widetilde{\chi}_2^0, \ \widetilde{\chi}_3^0, \ \widetilde{\chi}_1^{\pm}$$

The rest of the spectrum *can* be decoupled. Still sizeable EW production:

$$pp \to \widetilde{\tau}_1^+ \widetilde{\tau}_1^- + X, \quad pp \to \widetilde{\chi}_i^0 \widetilde{\chi}_j^0 + X, \quad pp \to \widetilde{\chi}_i^0 \widetilde{\chi}_1^{\pm} + X, \quad pp \to \widetilde{\chi}_1^+ \widetilde{\chi}_1^- + X$$

Drell-Yan, up to O(1) pb at LHC8

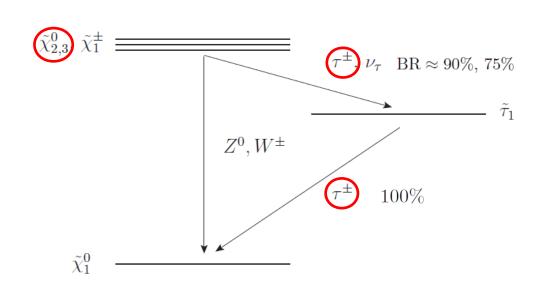
Decays:

$$m_{\widetilde{\chi}_{1}^{\pm}} \simeq m_{\widetilde{\chi}_{2,3}^{0}} > m_{\widetilde{\tau}_{1}} > m_{\widetilde{\chi}_{1}^{0}}$$

$$pp \longrightarrow \widetilde{\chi}_{2,3}^{0} \, \widetilde{\chi}_{2,3}^{0} \longrightarrow 4\tau + \cancel{E}_{T}$$

$$pp \longrightarrow \widetilde{\chi}_{2,3}^{0} \, \widetilde{\chi}_{1}^{\pm} \longrightarrow 3\tau + \cancel{E}_{T}$$

multi-tau signals!



LHC phenomenology: production and decays

Relic density constr. imply that we have at least 4 states at O(100) GeV:

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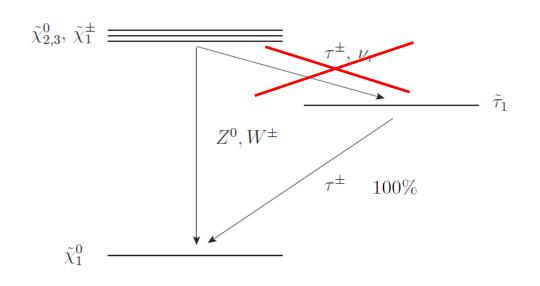
Drell-Yan, up to O(1) pb at LHC8

Decays:

$$m_{\widetilde{\tau}_1} > m_{\widetilde{\chi}_{2,3}^0} > m_{\widetilde{\chi}_1^0}$$

3-body decays it might be more difficult

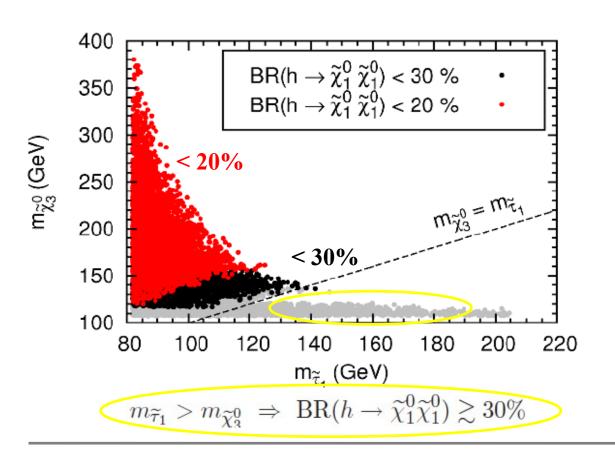
However...



LHC phenomenology: invisible Higgs decay

$$\Gamma(h \to \tilde{\chi}_{1}^{0} \tilde{\chi}_{1}^{0}) = \frac{G_{F} M_{W}^{2} m_{h}}{2\sqrt{2}\pi} \left(1 - 4m_{\tilde{\chi}_{1}^{0}}^{2}/m_{h}^{2}\right)^{3/2} \left|C_{h\tilde{\chi}_{1}^{0}\tilde{\chi}_{1}^{0}}\right|^{2}$$

$$C_{h\tilde{\chi}_{1}^{0} \tilde{\chi}_{1}^{0}} = \left(N_{12} - \tan \theta_{W} N_{11}\right) \left(\sin \beta N_{14} - \cos \beta N_{13}\right)$$



Last fits to Higgs data:

$$\mathrm{BR}_h^\mathrm{inv} < 16\% \quad (95\% \, \mathrm{CL})$$
Falkowski et al. '13
 $\mathrm{BR}_h^\mathrm{inv} < 19\% \quad (95\% \, \mathrm{CL})$
Giardino et al. '13

but with 20% theo unc.:

$${
m BR}_h^{
m inv} < 52\% \quad (68\%~{
m CL})$$

Djouadi Moreau '13

$$m_{\widetilde{ au}_1} > m_{\widetilde{\chi}^0_{2,3}} > m_{\widetilde{\chi}^0_1}$$
 strongly disfavoured!





ATLAS-CONF-2013-028

March 10, 2013



Search for electroweak production of supersymmetric particles in final states with at least two hadronically decaying taus and missing transverse momentum with the ATLAS detector in proton-proton collisions at $\sqrt{s} = 8 \text{ TeV}$

The ATLAS Collaboration

News from Atlas

ATLAS-CONF-2013-028

Signal region	requirements
OS $m_{\rm T2}$	at least 1 OS tau pair
	jet veto
	Z-veto
	$E_{\rm T}^{\rm miss} > 40~{\rm GeV}$
	$m_{\rm T2} > 90 {\rm GeV}$
OS m _{T2} -nobjet	at least 1 OS tau pair
	b-jet veto
	Z-veto
	$E_{\rm T}^{\rm miss} > 40~{\rm GeV}$
	$m_{\rm T2} > 100 {\rm GeV}$

№ 450 ATLAS Preliminary SR combined L dt = 20.7 fb⁻¹, \sqrt{s} =8 TeV μ [GeV]

M₁=50 GeV, tan β=50, 2 taus

Table 1: Definition of the signal regions.

SM process	SR OS m_{T2}	SR OS $m_{\rm T2}$ -nobjet
top	$0.2 \pm 0.5 \pm 0.1$	$1.6 \pm 0.8 \pm 1.2$
Z+jets	$0.28 \pm 0.26 \pm 0.23$	$0.4 \pm 0.3 \pm 0.3$
diboson	$2.2 \pm 0.5 \pm 0.5$	$2.5 \pm 0.5 \pm 0.9$
multi-jet & W+jets	$8.4 \pm 2.6 \pm 1.4$	$12 \pm 3 \pm 3$
SM total	$11.0 \pm 2.7 \pm 1.5$	$17 \pm 4 \pm 3$
data	6	14

$$m_{\tilde{\tau}_1} = 95 \,\text{GeV}, \quad m_{\tilde{\chi}_1^0} = 50 \,\text{GeV}$$



Simulation

Herwig++ (event samples)

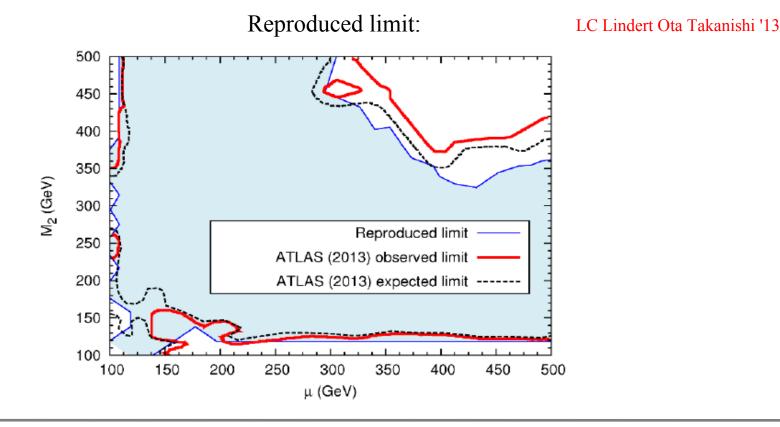
Cuts of two Atlas SR applied

Prospino 2 (NLO K-factors)

95% CL limits on the number of events:

Delphes 3 (fast detector simul.)

 $S_{
m SR1}^{95} < 5.6 \quad S_{
m SR2}^{95} < 10.4$



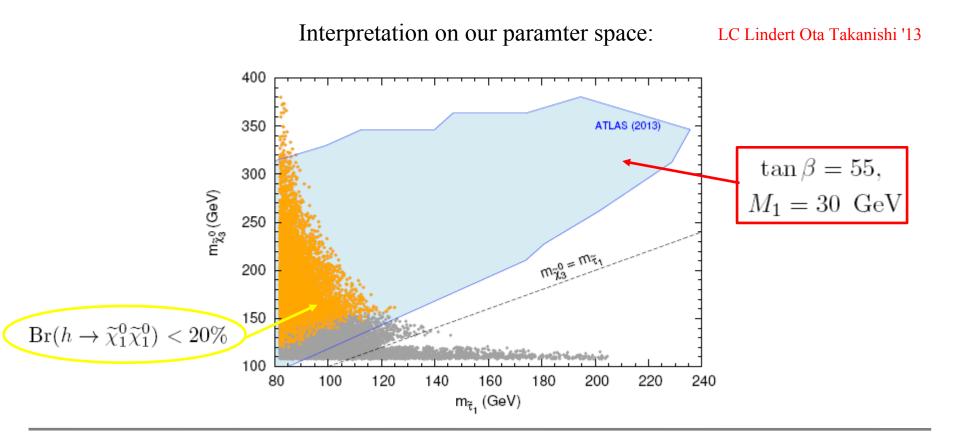
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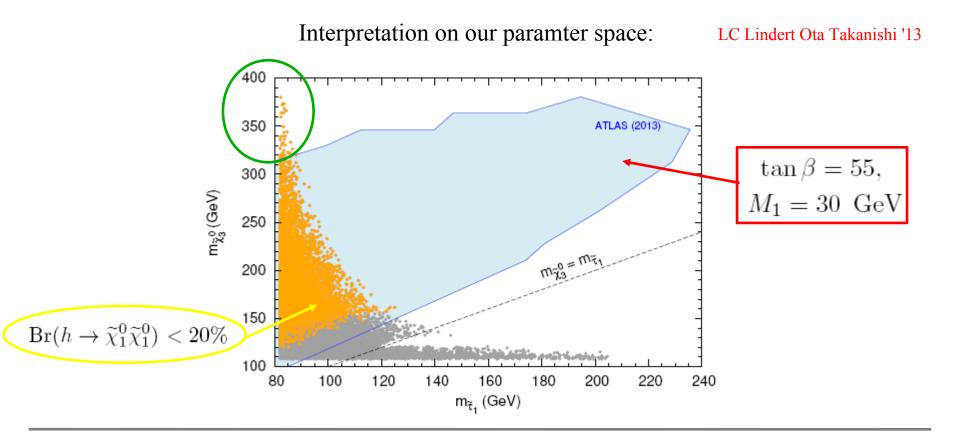
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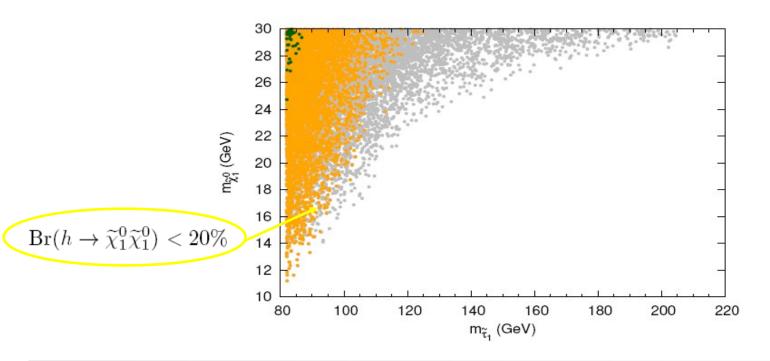
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Interpretation on the stau-neutralino mass plane:



Cuts of two Atlas SR applied

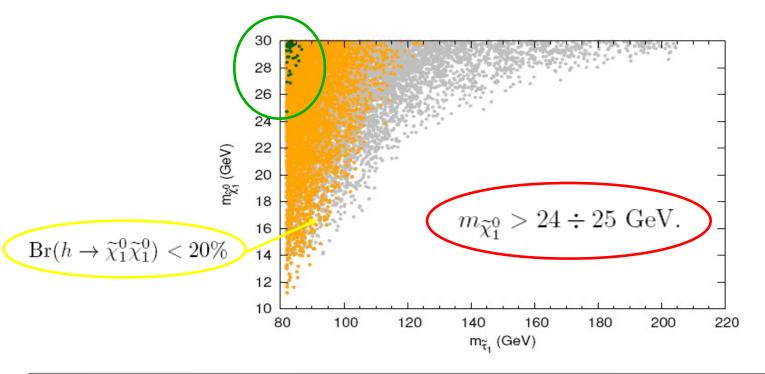
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Interpretation on the stau-neutralino mass plane:



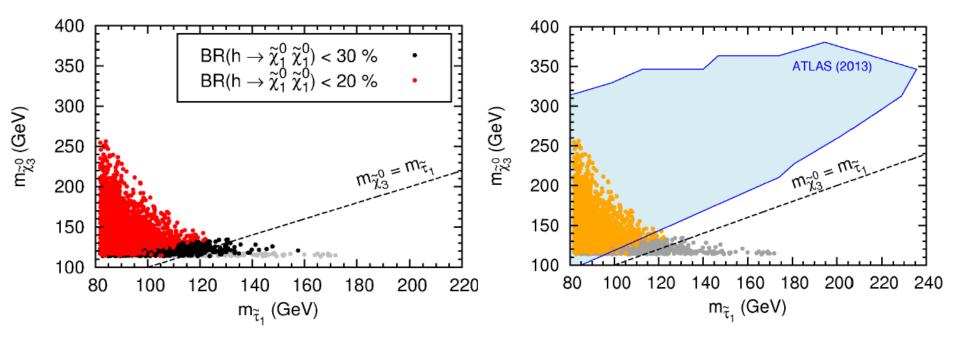
Flipping the sign of μ :

LC Lindert Ota Takanishi '14

Destructive interference between two annihilation modes: $\tilde{B}\tilde{B}$ and $\tilde{B}\tilde{H}_d$

Partial cancellation in $\tilde{\chi}_1^0 \tilde{\chi}_1^0 h \implies \text{lower } BR(h \to \tilde{\chi}_1^0 \tilde{\chi}_1^0)$

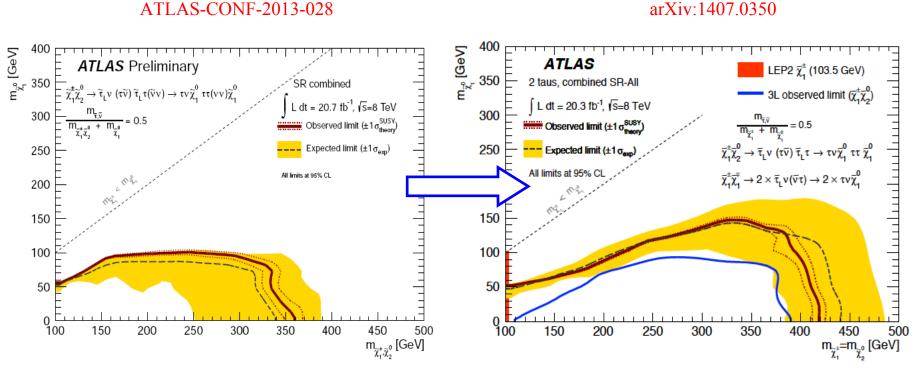
$$N_{14} = -\frac{M_Z s_W}{\mu} \left[c_\beta + s_\beta \frac{M_1}{\mu} \right]$$



Bound from the invisible decay milder, but DM heavier because annihilation less efficient \Rightarrow limit on the neutralino mass similar to the case $\mu>0$

Search for the direct production of charginos, neutralinos and staus in final states with at least two hadronically decaying taus and missing transverse momentum in pp collisions at \sqrt{s} = 8 TeV with the ATLAS detector

ATLAS-CONF-2013-028



Conclusions

Neutralino Dark Matter lighter than ~ 30 GeV requires light staus and higgsinos (relic density constraints)

Few parameters involved: manageable simplified model

Generic prediction: multi-tau + missing energy signal at the LHC

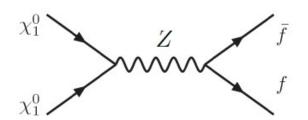
Reinterpreation of a ATLAS search sets strong constraint on light Neutralino Dark Matter (in combination with Higgs -> invisible)

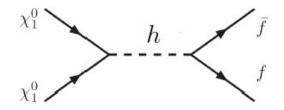
Limit stronger than direct/indirect DM searches

Nice interplay among different experimental info (collider and not)

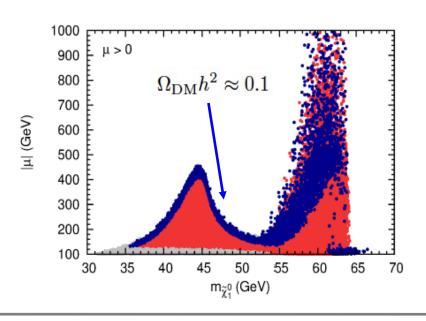
What if there are no scalars below few hundreds GeV?

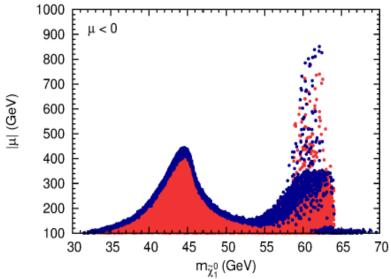
Dark Matter might hide close to the Z or h "resonances"



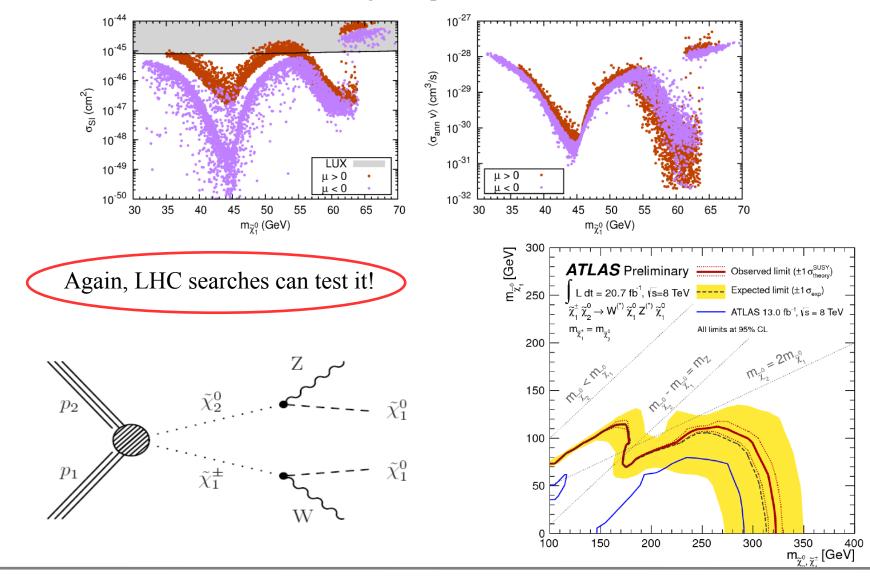


Only higgsino/gaugino parameters are involved: M_1 , M_2 , μ , $\tan \beta$





Close to the resonances, the scenario might be problematic for direct/indirect DM searches:



Experiments interpret the search in terms of Wino production and BR($\tilde{\chi}_2^0 \to Z\tilde{\chi}_1^0$)=100% Again, we have to recast the exclusion. Preliminary results (using Checkmate):

