# Spectroscopy of exotic states (Experimental aspects) 

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SFB 1044 Summer School
29.08.-01.09.2017

Boppard, Germany

## OUTLINE

- Lecture 1: The Past and Present
- How it all began (on a parking lot)
- The $X$, the $Y$ and the $Z$
- What is „beam constraint"?
- Lecture 2: The Future
- Dalitz formalism: how to „see" quantum numbers
- (Bottomonium: recoil mass, kinematic reflections $\rightarrow$ BACKUP)
- Belle II (ready to take data ?)
- Panda, an X factory! (,,detailed balance")



## Discovery of Charmonium (J/ $\Psi$ )

- SLAC (Stanford)

Mark I, Richter et al.
$\mathrm{e}+\mathrm{e}-\rightarrow$ hadrons, $\mathrm{e}+\mathrm{e}-, \mu+\mu-$

- BNL (Brookhaven) E598, Ting et al.
$p+A \rightarrow[e+e-] X$
- new, very narrow state (decay to light mesons blocked by OZI) width $\sim 100 \mathrm{keV}$
 (dith 100 kV )

- experimental proof for existance of $4^{\text {th }}$ quark

MARK I group reacted quickly
$\rightarrow$ it was feasible to modify the accelerator, so that the beam energies could be changed to $\leq 1 \mathrm{MeV}$ every minute.

## Discovery of $\psi^{\prime}$

- first exited ( $\mathrm{n}=2$ ) state of $\mathrm{J} / \psi$
" only 3 weeks after $J / \psi$
" beginning of charmonium spectroscopy
- Decay:
$\Psi^{\prime} \rightarrow J / \Psi \pi+\pi-$




## Heavy Quarkonium



## Charmonium vs. Positronium

Decays to light quarks suppressed $\rightarrow$ narrow widths


Charmonium


Positronium


## Cornell-Potential

- Coulomb-Potential
+ Confinement-Term
$k=0.5 \mathrm{GeV} / \mathrm{fm}$
$V(r)=-\frac{4}{3} \frac{\alpha_{s}}{r}+k r$
spin-spin $+\frac{32 \pi \alpha_{s}}{9 m_{c}^{2}} \delta_{r} \vec{S}_{c} \vec{S}_{\bar{c}}$
spin-orbit $\quad+\frac{1}{m_{c}^{2}}\left(\frac{2 \alpha_{s}}{r^{3}}-\frac{k}{2 r}\right) \vec{L} \vec{S}$
tensor $+\frac{1}{m_{c}^{2}} \frac{4 \alpha_{s}}{r^{3}}\left(\frac{3 \vec{S}_{c} \vec{r} \cdot \vec{S}_{c} \vec{r}}{r^{2}}-\vec{S}_{c} \vec{S}_{\bar{c}}\right)$
- solve Schrödinger equation (quark mass heavy $\rightarrow$ non-relativistic)

$$
\Psi(r, \theta, \phi)=R_{n l}(r) Y_{l m}(\theta, \phi)
$$

$$
\left[-\frac{1}{m_{q}}\left(\frac{\partial^{2}}{\partial r^{2}}+\frac{2}{r} \frac{\partial}{\partial r}+\frac{l(l+1)}{m_{q} r^{2}}+V(r)\right)\right] R_{n l}(r)=E_{n l} R_{n l}(r) \quad \mathrm{JPC}
$$

## Quarkonium Excited States






Charmonium





Radial Wavefunctions




The agreement between prediction by the potential model and experimental observation is and was encouraging (level $\sim 10^{-3}, 2-3 \mathrm{MeV}$ compared to mass of $3-10 \mathrm{GeV}$ ) and heavy quark hadron physicists were living happily.

About 30 years passed.

## Belle and KEKB, Japan


$\rightarrow$ extend decay length symm. of $B$ mesons

| Accelerator Laboratory | CESR <br> Cornell | KEKB KEK | $\begin{aligned} & \text { PEP-II } \\ & \text { SLAC } \end{aligned}$ | SuperKEKB KEK |
| :---: | :---: | :---: | :---: | :---: |
| Detector | CLEO III | Belle | BaBar | Belle II |
|  | (achieved) | (achieved) | (achieved) | (planned) |
| Circumference (km) | 0.768 | 3.0 | 2.2 | 3.0 |
| Energy $e^{-} / e^{+}(\mathrm{GeV})$ | 5.3/5.3 | 8.0/3.5 | 9.0/3.1 | 7.0/4.0 |
| Lorentz boost $\beta \gamma$ | 0 | 0.43 | 0.56 | 0.28 |
| Beam current $e^{-} / e^{+}$(A) | 0.5/0.5 | 1.6/1.2 ${ }^{\dagger}$ | 3.2/2.1 | 3.6/2.6 |
| Number of bunches | 45 | 5120 | 1732 | 2500 |
| Crossing angle (mrad) | $\pm 2.3$ | $\pm 11$ | 0 | 83 |
| Luminosity ( $10^{33} / \mathrm{cm}^{2} \mathrm{~s}$ ) | 1.55 | 21.08 | 12.07 | 800 |
| $\sigma_{x}(\mu \mathrm{~m})$ | n.a. | 103-116 | 120 | 7.2-8.9 |
| $\sigma_{y}(\mu \mathrm{~m})$ | n.a. | 0.94 | 4 | $36 \times 10^{-3}$ |
| $\sigma_{z}(\mathrm{~mm})$ | n.a. | 6 | 11 | 5 |

## The BELLE Detector


A. The $\Upsilon(n S)$ states


## Luminosity

```
Belle Run Summary(v1.84) - Exp 47 Run 529
Start Time: 2005 Dec 7, 22:58:37 took 78 sec to start
Stop Time: 2005 Dec 7, 23:14:01 took 924 sec
Stop Reason: FATAL from [EFARM1] E1TRK CDC timeout (1 sec rx0 stat=28 len=64/0/11448 ev=193658,
\begin{tabular}{ll} 
Expert shift: & S. Lange \\
Non-Expert: & K. Kinoshita \\
BCG shift: & Ishikawa (4862) \\
& \\
Run Mode: & Luminosity Run
\end{tabular}
```



Luminosity:
at start
at stop peak/fill

```
ECL
155.80 e 32
156.55 e 32
157.78 e 32
EFC KEKB
140.07e32 108.39e32
139.11e32 109.06e32
L>1 x 1034 s-1 cm
```

$1 \times 10^{34} \mathrm{~s}^{-1} \mathrm{~cm}^{-2}$

1 barn $=10^{-24} \mathrm{~cm}^{2}$

24 hours $\times 60 \mathrm{~min} \times 60$ seconds $\times 1 / 10^{10}$ barn $^{-1}$
$=86.400 \times 1 / 10^{10}$ barn $^{-1}$
$\sim 1 / 10^{15}$ barn $^{-1}$
$=1$ inverse femtobarn per 1 day

1 nb cross section
$\times 1.000 .000 \mathrm{nb}^{-1}$ per 1 day
$=1.000 .000 \mathrm{Y}(4 \mathrm{~S})$ per 1 day !

Almost every $\mathrm{Y}(4 \mathrm{~S})$ decays into a B anti- B meson pair.
$\rightarrow$ we call it a „B meson factory"

## Example Events with Charmonium

$$
J / \psi \rightarrow e+e-
$$

$$
\mathrm{J} / \psi \rightarrow \mu^{+} \mu^{-}
$$



## $J / \psi \rightarrow \mathrm{e}+\mathrm{e}-, \mu+\mu-$ Invariant Mass




## $\mathrm{K}_{\text {long }}$ Detection, also with the muon detector



## A difference between BaBar and Belle: the Cerenkov detectors

## Belle: Aerogel Cerenkov Detector



## Particle Indentification with Aerogel ",threshold mode" (Belle)


pulse height / number of photoelectrons

## BaBar DIRC (Detection of internally reflected Cerenkov light)


$4 \times 1.225 \mathrm{~m}$
Synthetic Fused Silica
Bars glued end-to-end

## Belle Silicon Vertex Detector



FIG. 28. Belle SVD 1.4 (left) and SVD 2.0 (right).
about $10^{5}$ readout channels vertex resolution in beam direction $\Delta z \sim 50 \mu \mathrm{~m}$

## Particle Detection with Silicon Vertex Detectors

single sided


Schematic of a single-sided silicon-strip detector
G.J.Barker, b Quark Physics at LEP



Reconstruction of higher charmonium states by invariant mass

$$
m=\sqrt{\left(\sum_{i} E_{i}^{2}\right)-\left(\sum_{i}{\vec{p}_{i}^{2}}^{2}\right)}
$$

## often <br> DOUBLE GAUSSIAN




Why we often use double Gaussian to fit a signal:
a.) for real data
any effect, which is not in MC,
is parametrized as an additional Gaussian
b.) if a peak is constructed from two (or more) daughter particles
each daughter particle contributes as a Gaussian
(in particular, when combining a neutral and a charged particle)

Example: $\phi \rightarrow \mathrm{K}+\mathrm{K}-$


Single Gaussian
sigma 3.12 MeV

double Gaussian sigma 2.00 MeV

## input value in generator: width 4.26 MeV

Note: narrow and wide Gaussian are forced to same mean.

## Breit Wigner Shape

in particular for particles with a natural width comparable to detector resolution

$$
p(m)=\frac{1}{\left(m^{2}-m_{0}^{2}\right)^{2}+m_{0}^{2} \Gamma^{2}(m)}
$$

width may change
from resonance theory („T matrix"):
with m!!! the pole of a resonance is a complex number:
a real part (m) and an imaginary part $(\Gamma)$

## Voigtian shape Breit-Wigner and Gauss, folded

- Width of BW fixed to 4.26 MeV (input width in EvtGen)
- Result: Gaussian adds 0.9 MeV (= resolution)




## Attention!

FIG. 109. Difference of the line shape of a Breit-Wigner probability distribution and a Gaussian probability distribution R. K. Bock, W. Krischer, The Data Analysis BriefBook

## Beam Constrained Mass m $m_{B C}$

or
energy substituted mass $\mathrm{m}_{\mathrm{ES}}$

$$
\begin{gathered}
m_{b c}=\sqrt{\frac{\left(E_{b e a m}^{c m s}\right)^{2}}{4}-\left(p_{B}^{c m s}\right)^{2}} \\
=\sqrt{\frac{E_{B C}^{2}}{4}-\left(\left[\sum_{i} p_{i x}\right]^{2}+\left[\sum_{i} p_{i y}\right]^{2}+\left[\sum_{i} p_{i z}\right]^{2}\right)}
\end{gathered}
$$

$$
\mathrm{B} \rightarrow \mathrm{~K}\left[\mathrm{~J} / \psi \pi_{1} \pi_{2}\right]
$$

$$
\begin{aligned}
& \mathrm{px} \_\mathrm{Bc}=\text { jpsi_vector. } \mathrm{x}()+\mathrm{pcms}[\text { kaon }] . \mathrm{x}()+\mathrm{pcms}[\text { pion1]. } \times()+\text { pcms[pion2].x(); } \\
& \text { py_B }=\text { jpsi_vector. } \mathrm{y}()+\text { pcms }[\text { kaon }] \cdot \mathrm{y}()+\mathrm{pcms}[\text { pion1] } \mathrm{y}()+\text { pcms }[\text { pion2 } 2 \cdot \mathrm{y}() \text {; } \\
& \mathrm{pz} \_\mathrm{B}=\text { jpsi_vector. } \mathrm{z}()+\text { pcms }[\text { kaon }] . \mathrm{z}()+\mathrm{pcms}[\text { pion1].z( })+\text { pcms[pion2].z(); } \\
& \text { esum_B }=\text { jpsi_vector.e() }+ \text { pcms[kaon].e() }+ \text { pcms[pion1].e() }+ \text { pcms[pion2].e(); } \\
& \text { deltaE }=\mathrm{ECM} / 2 \text { - esum_B; } \\
& \text { mass_B = sqrt( esum_B*esum_B - (px_B*px_B + py_B*py_B + pz_B*pz_B) ); } \\
& \text { mass_BC }=\operatorname{sqrt}\left(E C M * E C M / 4 .-\left(p x \_B^{*} p x \_B+p y \_B^{*} p y \_B+p z \_B^{*} p z \_B\right)\right. \text { ); }
\end{aligned}
$$

ECM comes from accelerator measurement.

B meson
invariant mass

Beam constrained mass $\mathrm{m}_{\mathrm{BC}}$
or
energy substituted mass
mes


## Statistical Significance



Fit $S+B G$
$\chi^{2}=42.08$
Significance $=\sqrt{ }(-42.08+73.08)=5.6$, sigma" if „likelihood fit", then $\chi^{2} \rightarrow 2 \ln$ (likelihood)
Upper limit $(<3 \sigma)$, evidence $(>3 \sigma)$, observation $(>5 \sigma)$

## $\chi^{2}$ fit vs. log likelihood fit

$\mathrm{S}=\sqrt{ }\left(\chi^{2}(\mathrm{~B})-\chi^{2}(\mathrm{~S}+\mathrm{B})\right) \quad P_{\lambda}(k)=\frac{\lambda^{k}}{k!} \mathrm{e}^{-\lambda}$,
If number of events is small $\rightarrow$ Poisson statistics $\rightarrow \exp ()$ Term $\rightarrow \log$ likelihood is better than $\chi^{2}$ (in other words, removing the shape of $\exp ()$ and make it flat, when searching for the global minimum)
$\chi^{2}=-2 \ln L+$ constant
(don't forget the minus sign!
minimum $\chi^{2}$, but maximum log likelihood)

$$
S=\sqrt{ }(2(-\ln L(B)--\ln L(S+B)))
$$

## HOWTO scale the significance

- for a known resonance (control channel) mass and width are known from PDG only yield is floating $\rightarrow$ ndof $=1$
- for a new resonance (maybe exotic) nothing is known mass, width, and yield are floating $\rightarrow$ ndof $=3$
rule: we must scale the significance
chi2 $=52.58$ for a fit with ndof $=3$
root [4] TMath::Prob(42.08,3)
(Double_t) $3.85831510574144829 e-09$
search for new chi2 with same $p$-value, but ndof $=1$
root [68] TMath::Prob(34.694,1) // corrected chi2 (Double_t) $3.85818878297546610 \mathrm{e}-09$


## X(3872)

## A molecular state?

$$
B^{ \pm} \rightarrow K^{ \pm} \underbrace{J / \psi \pi^{+} \pi^{-}}_{\text {resonant state? }}
$$



Product branching fraction small
$B(B$ decay $) \times B(X$ decay $) \simeq 10^{-5}$

## X(3872)



X(3872)



This is $B$ meson decays
This is not simply invariant mass, but fitted $m_{B C}$ yield (for a given mass bin) many fits!

Trick: subtract the $\mathrm{J} / \psi$ mass (so the experimental resolution of the $\mathrm{J} / \psi$ is „taken out")

This is „inclusive" production (not B decays)
This is invariant mass.
$J / \psi$ mass is not subtracted.

## $X(3872 \rightarrow J / \psi \pi \pi$ event - can you „see" the $\psi$ shape ?



## What do we know about the $\mathrm{X}(3872)$ ?

- Observed by 7 experiments
- Observed in 5 decay channels
- very near to $\mathrm{D}^{*}$ threshold $\mathrm{E}_{\mathrm{B}}=0.01 \pm 0.18 \mathrm{MeV}$
but $\mathrm{D} \overline{\mathrm{D}}^{*}$ decays dominant (factor $\sim 10$ )
- $\Gamma \leq 1.2 \mathrm{MeV} \rightarrow$ very narrow (Belle, by 3-dim overconstrained fit)
- JPC=1++

Cornell-potential: $\chi_{\mathrm{c} 1}{ }^{\prime}$


Barnes et al., Phys. Rev. D72(2005)054026
$\rightarrow$ predicted mass $\geq 50 \mathrm{MeV}$ higher
$\rightarrow$ predicted width factor $\geq 100$ larger

- isospin violating decays

Is the $X(3872)$ exotic ?

## TETRAQUARK



$$
[q Q]_{8}[\overline{q Q}]_{8}
$$

Diquarks
are colored

Maiani, Riquer, Piccinini, Polosa, Burns;
Ebert, Faustov, Galkin; Chiu, Hsieh;
Ali, Hambrock, Wang

THRESHOLD CUSP


## MOLECULE

Intriguing Analogon


Tornqvist; Swanson; Braaten, Kusonoki, Wong; Voloshin; Close, Page Guo, Hanhart, Meissner

Bugg; Swanson

## Subresonant structure of $X(3872) \rightarrow J / \psi\left[\pi^{+} \pi^{-}\right]$


almost no non-resonant phasespace component ! dominated by $\rho^{0}(\sim 100 \%)$ ! ISOSPIN VIOLATING 2 particle decay (back-to-back), not a 3-particle decay !

## Isospin violating charmonium transistions

Only 2 decays for charmonium measured in PDG

## Decays into $J / \psi(1 S)$ and anything

|  | $J / \psi(1 S)$ anything | $(61.0 \pm 0.6) \%$ |
| :--- | :--- | :---: |
|  | $J / \psi(1 S)$ neutrals | $(25.11 \pm 0.33) \%$ |
| $J / \psi(1 S) \pi^{+} \pi^{-}$ | $(34.46 \pm 0.30) \%$ |  |
| $\boldsymbol{\psi ( 2 S )}$ | $\left(\psi(1 S) \pi^{0} \pi^{0}\right.$ | $(3.14 \pm 0.31) \%$ |
| $J / \psi(1 S) \eta$ | $(1.268 \pm 0.05) \%$ |  |
|  | $J / \psi(1 S) \pi^{0}$ | Hadronic decays |
|  | $(8.6 \pm 1.3) \times 10^{-3}$ |  |
| $\pi^{0} h_{c}(1 P)$ |  |  |

but branching fraction of
$\mathcal{B}(X(3872) \rightarrow J / \psi \rho)$ is order of $\sim 5-10 \%$ factor $\sim 10^{2}$ too large

Is there isospin inside the $X(3872)$ ?


No evidence. Significant $\rho / \omega$ interference can explain lineshape.
(proposed by Terasaki, Prog. Theor. Phys. 122(2010)1285)

## What important knowledge is missing? <br> $\rightarrow$ Width of $\mathrm{X}(3872)$

upper limit on width (Belle I), $\Gamma<1.2 \mathrm{MeV}$
for pure $\chi_{\mathrm{c} 1}{ }^{\prime}$ charmonium state,
prediction $\Gamma=40 \mathrm{keV}$
G. Y. Chen, J. P. Ma, arXiv:0802.2982[hep-ph], Phys. Rev. D77(2008]097501.
if molecule

- must be larger than width of $\mathrm{D}^{*}$
$\Gamma>82.3 \pm 1.2 \pm 1.4 \mathrm{keV}$
E. Braaten, arXiv:0711.1854 [hep-ph], Phys. Rev. D77(2008)034019.
- long-range molecular components in the wavefunction?
$\rightarrow$ measure the width of the $X(3872)$
in the sub-MeV regime


## X(3872) Width Measurement at Belle I

$M_{\mathrm{bc}} \equiv \sqrt{\left(E_{\mathrm{beam}}^{\mathrm{cms}}\right)^{2}-\left(p_{B}^{\mathrm{cms}}\right)^{2}}$

$\mathrm{M}_{\mathrm{BC}} / \mathrm{GeV}$


$M\left(\mathrm{~J} / \psi \pi^{+} \pi^{-}\right) / \mathrm{GeV}$

$$
\Delta E \equiv E_{B}^{\mathrm{cms}}-E_{\text {beam }}^{\mathrm{cms}}
$$


$\Delta \mathrm{E} / \mathrm{GeV}$

3-dim fit $\rightarrow$ kinematical over-constraint provides access to observables smaller than detector resolution

Belle, Phys. Rev. D84(2011)052004

Reference Analysis: $\mathrm{B} \rightarrow \mathrm{K} \psi^{\prime}, \psi^{\prime} \rightarrow \mathrm{J} / \psi \pi^{+} \pi^{-}$
$M_{\mathrm{bc}} \equiv \sqrt{\left(E_{\mathrm{beam}}^{\mathrm{cms}}\right)^{2}-\left(p_{B}^{\mathrm{cms}}\right)^{2}}$



$$
\Delta E \equiv E_{B}^{\mathrm{cms}}-E_{\underline{\mathrm{beam}}}^{\mathrm{cms}}
$$




factor $\sim 10$ more statistics than $\mathrm{X}(3872) \rightarrow$ use as reference signal
$\rightarrow$ fix resolution parameters
$\rightarrow$ fix absolute mass scale (MC/data shift $+0.92 \pm 0.06 \mathrm{MeV}$ )

## Measurement of width of $X(3872)$

- Correlation function from MC
$\Gamma$ (output) $=\mathrm{f}(\Gamma$ (input) $)$
- 3-dim fits validated with $\psi^{\prime}$ width $\Gamma_{\psi}=0.52 \pm 0.11 \mathrm{MeV}$ (PDG $0.304 \pm 0.009 \mathrm{MeV}$ )
$\rightarrow$ bias $0.23 \pm 0.11 \mathrm{MeV}$
- procedure for upper limit: width in 3-dim fit fixed

$\mathrm{n}_{\text {signal }}$ and $\mathrm{n}_{\mathrm{BG}}$ floating
$\rightarrow$ calculate likelihood
- $\quad \Gamma_{\times(3872)}<\underbrace{0.95 \mathrm{MeV}+\text { bias }}$


### 1.2 MeV

- implication: width of $X(3872)$ can be measured at Belle II



## $X(3872)$ in $B$ decays



## Exercise:

$$
\begin{aligned}
& \mathrm{B} \rightarrow \mathrm{~K} \times(3872) \\
& 0 \rightarrow 0-1+ \\
& \text { parity }(-1) \rightarrow \text { parity }(-1) \times(+1) \times(-1)^{\text {L }}
\end{aligned}
$$

## Exercise:

$$
\begin{gathered}
B \rightarrow K \times(3872) \\
0 \rightarrow 0-1+
\end{gathered}
$$

$$
\text { parity }(-1) \rightarrow \text { parity }(-1) \times(+1) \times(-1)^{\llcorner }
$$

We need $\mathrm{L}=1$ to create $\mathrm{J}=1$, but this violates parity.
The $X(3872)$ is created in a parity violating weak decay.

## Y(4260) <br> a gluonic hybrid state?

Note: recent notation
by PDG as X(4260)

## Y(4260)

- initial state radiation events
$\mathrm{e}^{+\mathrm{e}-\rightarrow \gamma \mathrm{J} / \psi \pi+\pi-}$ (undetected $\gamma$ parallel to beam axis)
- mass $>4 \mathrm{GeV}$
far above $\mathrm{D}\left({ }^{*}\right) \mathrm{D}\left(^{*}\right)$ thresholds decay to $J / \psi \pi+\pi-$ should be suppressed
- width $<100 \mathrm{MeV}$ quite narrow for such a high state
- quantum numbers must be (based upon production mechanism)
${ }^{P} P C=1--$
initial state radiation




## Y(4260) Parameters

|  | BaBar [1] | CLEO-c [2] | Belle [3] | Belle [4] | BaBar [5] | BaBar [6] |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\mathcal{L}$ | $211 \mathrm{fb}^{-1}$ | $13.3 \mathrm{fb}^{-1}$ | $553 \mathrm{fb}^{-1}$ | $548 \mathrm{fb}^{-1}$ | $454 \mathrm{fb}^{-1}$ | $454 \mathrm{fb}^{-1}$ |
| N | $125 \pm 23$ | $14.1_{-4.2}^{+5.2}$ | $165 \pm 24$ | $324 \pm 21$ | $344 \pm 39$ | - |
| Significance | $\simeq 8 \sigma$ | $\simeq 4.9 \sigma$ | $\geq 7 \sigma$ | $\geq 15 \sigma$ | - | - |
| $m / \mathrm{MeV}$ | $4259 \pm 8_{-6}^{+2}$ | $4283_{-16}^{+17} \pm 4$ | $4295 \pm 10_{-3}^{+10}$ | $4247 \pm 12_{-32}^{+17}$ | $4252 \pm 6_{-3}^{+2}$ | $4244 \pm 5 \pm 4$ |
| $\Gamma / \mathrm{MeV}$ | $88 \pm 23_{-4}^{+6}$ | $70_{-25}^{+40}$ | $133 \pm 26_{-6}^{+13}$ | $108 \pm 19 \pm 10$ | $105 \pm 18_{-6}^{+4}$ | $114_{-15}^{+16} \pm$ 7 |

[1] BaBar Collaboration, arXiv:hep-ex/0506081, Phys. Rev. Lett. 95(2005)142001.
[2] CLEO-c Collaboration, arXiv:hep-ex/0611021, Phys. Rev. D74(2006)091104.
[3] Belle Collaboration, arXiv:hep-ex/0612006.
[4] Belle Collaboration, arXiv:0707.2541[hep-ex], Phys. Rev. Lett. 99(2007)182004.
[5] BaBar Collaboration, arXiv:0808.1543[hep-ex].
[6] BaBar Collaboration, arXiv:1204.2158[hep-ex], Phys. Rev. D86(2012)0511け2.

## BESIII Experiment at BEPC II (symmetric !)



Superconducting solenoid
$\mathrm{B}=1 \mathrm{~T}$
no silicon vertex detector, because only D mesons (no separation of B mesons and D mesons required)
no Cerenkov detector (time-of-flight sufficient for $\mathrm{K} / \pi$ separation, low momentum)

## $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \mathrm{Y}(4260) \rightarrow \mathrm{J} / \psi \pi^{+} \pi^{-}$at BESIII

BESIII, PRL110(2013)252001


## Is the $Y(4260)$ exotic?

## TETRAQUARK

HADRO-CHARMONIUM $\left[J / \psi \mathrm{f}_{0}(980)\right]$


Voloshin, Li
(Guo, Hanhart, Meissner)


Zhu; Kou, Pene; Close, Page;
Lattice QCD, Bernard et al.; Mei, Luo
$\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \gamma_{\text {ISR }} \mathrm{J} / \psi\left(\psi^{\prime}\right) \pi^{+} \pi^{-}: Y(4008,4260,4350,4660)$


Sören Lange | Exotic States (Experiment) | SFB 1044 Summer School, Boppard, 08/2017

## What is the tail around 4.7 GeV ?

- Threshold $m(D)+m\left(D^{* *}\right)=4326 \mathrm{MeV}$

Lineshape distorted? No.

- Non-corrected radiative effects? No.

Radiative lower mass tail in $\mathrm{J} / \psi \rightarrow \mathrm{e}^{+} \mathrm{e}^{-}$ might generate higher mass tail in $\mathrm{m}\left(\mathrm{J} / \psi\right.$-with-wrong-mass $\left.\pi^{+} \pi^{-}\right)$.

- Fit funtion: Breit Wigner x Phasespace x Efficiency Efficiency $\mathrm{a}\left(\mathrm{m}-\mathrm{m}_{0}\right)+\mathrm{b}$ with $\mathrm{a}=7.4 \pm 1.3 \mathrm{GeV}-1, \mathrm{~b}=9.31 \pm 0.07$ (Belle) changes factor $\sim 2$ over peak




## Overpopulation of $1^{--}$States



All same quantum number


No more $\left[J / \psi \pi^{+} \pi^{-}\right]$state up to 7 GeV

Note: radiative transitions between the states forbidden by parity
but apparently

- no mixing with other $\psi$ states
- no mixing among them $Y(4260)$ seems not decay to $\psi^{\prime} \pi^{+} \pi^{-}$ $\mathrm{Y}(4350)$ seems not decay to $\mathrm{J} / \psi \pi^{+} \pi^{-}$


## Y(4260): Comparison Belle and BaBar

- BaBar collisions head-on, dipole magnet close to IR
- Belle: $\pm 11 \mathrm{mrad}$
- slightly higher background at BaBar (also seen as MRad SVD radiation dose)
- backward acceptance for $0 \sim 180$ • limited

hep-ex/0612006, 553/fb



CMS polar angle of $\mathrm{Y}(4260)$ to $\mathrm{e}-$ beam

What is the $Y(4260)$ ?
A hybrid? $[Q Q]_{8} g$

## Does the $\mathrm{Y}(4260)$ decay to $\mathrm{e}+\mathrm{e}-$ ?

- very small coupling to $\mathrm{e}+\mathrm{e}-$
(although JP=1--)
$\operatorname{BR}\left(\mathrm{J} / \psi \pi^{+} \pi^{-}\right) \times \Gamma\left(\mathrm{e}^{+} \mathrm{e}^{-}\right)=$
(7.5 $\pm 0.9 \pm 0.8) \mathrm{eV}$

BaBar, arXiv:0808.1543

- This is a partial width of the order ,,eV" of a state which is $\sim 100 \mathrm{MeV}$ total width!
 $\rightarrow$ factor $10^{8}$ suppressed

What is blocking these decays?
(maybe the gluonic string ?)

## Z(4430)+

$$
\begin{array}{r}
\mathrm{B}^{0} \rightarrow \mathrm{~K}^{+} \psi^{\prime} \pi^{-} \\
\psi^{\prime} \rightarrow \mathrm{J} / \psi \pi^{+} \pi^{-}
\end{array}
$$





# A charged state can never be a charmonium state. 

$$
\mathrm{Z}(4430)^{+}
$$




Belle
Phys. Rev. D80(2009)031104 605/fb

 $m=4433 \pm 4 \pm 2 \mathrm{MeV}, \Gamma=45_{-13}^{+18+30} \mathrm{MeV}$ Belle, Phys. Rev D80(2009)031104


## Long discussion between Belle and BaBar (2007-2014)

$$
\text { Is the } \mathrm{Z}+(4430) \text { a kinematical effect ? }
$$

- $\cos \theta_{K}$, the normalized dot-product between the $K \pi$ three-momentum vector in the parent- $B$ rest frame and the kaon three-momentum vector after a Lorentz transformation from the $B$ meson rest frame to the $K \pi$ rest frame
- $\cos \theta_{\psi}$, the normalized dot-product of the $\psi^{\prime} \pi^{\mp}$ three-momentum vector in the parent $B$ meson rest frame and the $\psi^{\prime}$ three-momentum vector in the $\psi^{\prime} \pi^{\mp}$ rest frame.
$\cos \theta_{K}$ is correlated with $m\left(K^{ \pm} \pi^{\mp}\right), \cos \theta_{\psi}$ is correlabed with $m\left(\psi \pi^{\mp}\right), \cos \theta_{K}$ is correlated with $\cos \theta_{\psi}$


## TRUE!



MC with angular correlations can describe data well.
No $Z^{+}$states in red line (MC) required! BaBar, arXiv:1111.5919, Phys.Rev. D85(2012)052003

