

An Improved Limit for Γ_{ee} of $X(3872)$ and Γ_{ee} Measurement of $\psi(3686)$

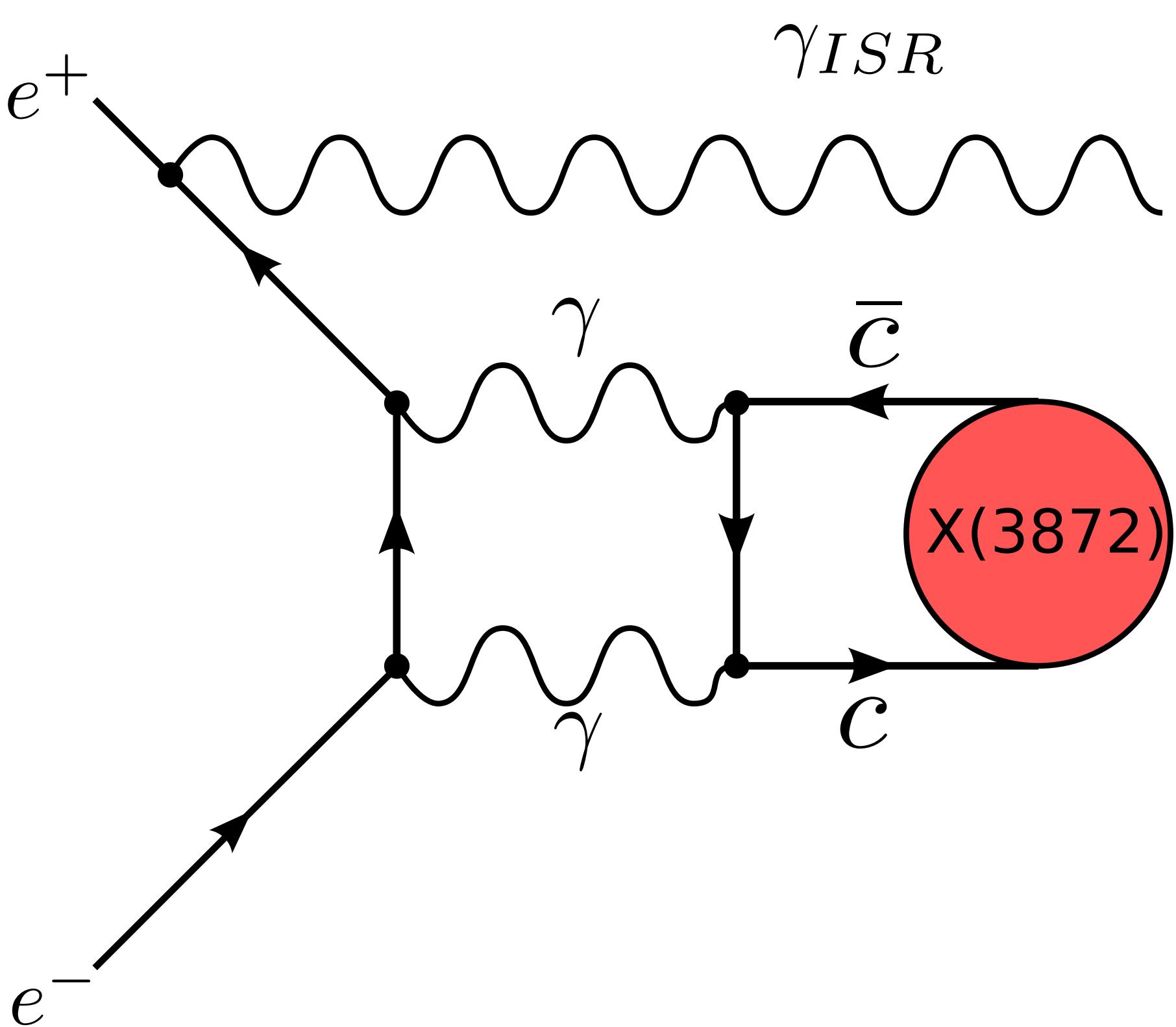
PLB 749 (2015) 414-420

Martin Ripka on behalf of the BESIII Collaboration

Institute for Nuclear Physics, Johannes-Gutenberg University Mainz

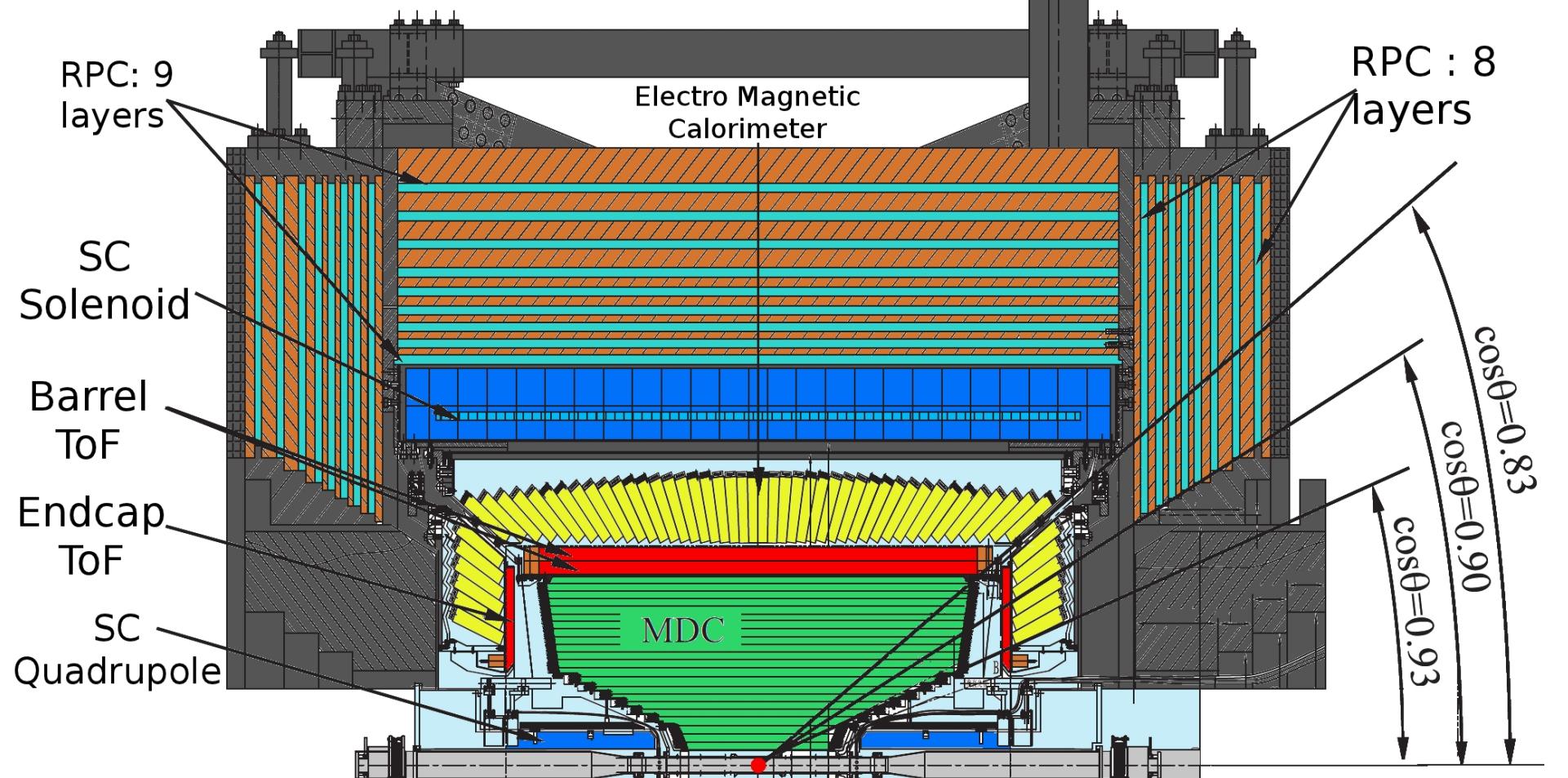
INTRODUCTION

- $X(3872)$ First observed in 2003 by Belle [PRL 91, 262001 (2003)]
- $J^{PC} = 1^{++}$ [PRL 110, 222001 (2013)]
- Close to $D^0\bar{D}^{*0}$ threshold
→ meson molecule? [PLB 725, 127 (2013)]
- Large decay rate $X(3872) \rightarrow \gamma\psi(3686)$ compared to $X(3872) \rightarrow \gamma J/\psi$
→ tetraquark? [Nuc.Ph. B 886, 665 (2014)] [PRL 112, 092001 (2014)]



- Decay $Y(4260) \rightarrow \gamma X(3872)$ recently observed at BESIII [PRL 112, 092001 (2014)]
- Theoretical calculation predicts $\Gamma_{ee}^{X(3872)} \approx 0.03$ eV [PLB 736, 221 (2014)]
- The current upper limit for $\Gamma_{ee}^{X(3872)}$ is at the $\mathcal{O}(10^2)$ eV level [PDG (2014)]
- 1^{++} state never observed directly in e^+e^- annihilation
- Process may occur via a two-photon box diagram

BESIII DETECTOR, DATA, MC



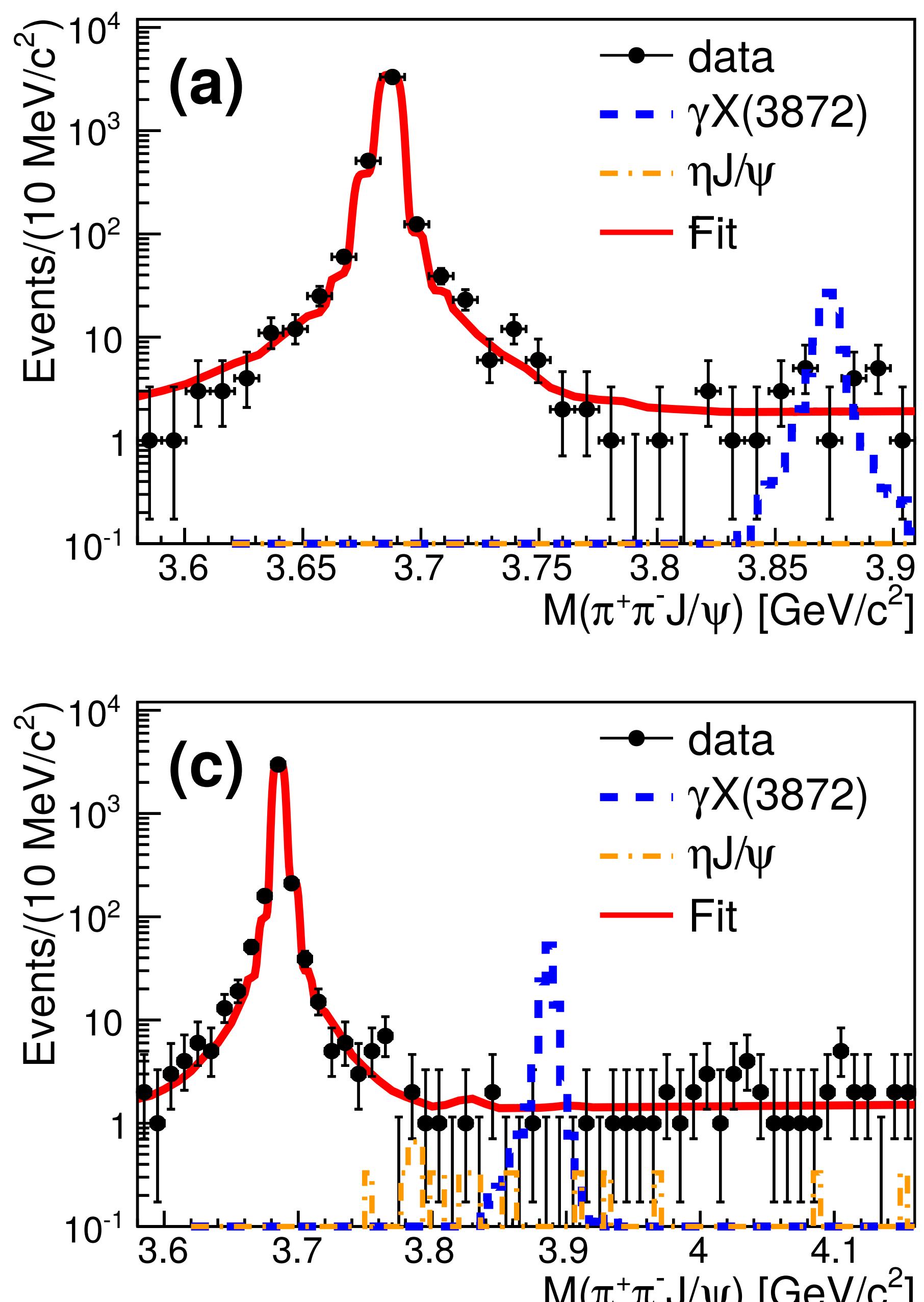
- Operating at BEPC II in Beijing/China
- Symmetric e^+e^- collider [$2-4.6$ GeV/c 2]

482 pb^{-1} @ 4.009 GeV, 1092 pb^{-1} @ 4.230 GeV,
 826 pb^{-1} @ 4.260 GeV, 540 pb^{-1} @ 4.360 GeV

ANALYSIS STRATEGY

- Using Initial State Radiation (ISR) technique to access $X(3872)$ resonantly
- $e^+e^- \rightarrow \gamma_{\text{ISR}} X(3872)$
- $X(3872) \rightarrow \pi^+\pi^- J/\psi$ ($\mathcal{B} > 3.8\%$)
- $J/\psi \rightarrow \ell^+\ell^-$, ($\ell = \mu, e$) ($\mathcal{B} = 11.96\%$)
- Untagged ISR photon: $|\cos \theta_{\text{ISR}}| > 0.95$

$\pi^+\pi^- J/\psi$ MASS SPECTRUM



- (a) 4.009 GeV, (b) 4.230 GeV, (c) 4.260 GeV and (d) 4.360 GeV
- No significant $X(3872)$ peak observed at any of the four c.m. energies

CALCULATION OF Γ_{ee}

- Event count rate $\frac{dN_A}{dx}$ ($A = \psi(3686), X(3872)$) obtained by unbinned maximum likelihood fits
- Fit PDF: $\psi(3686)$ -MC-shape \otimes Gaussian + Polynomial + $X(3872)$ -MC-shape
- Relation to non-radiative cross section: $\frac{dN_A}{dx} = W(s, x)\varepsilon_A \mathcal{L} \sigma(e^+e^- \rightarrow A)\mathcal{B}(A \rightarrow \pi^+\pi^- J/\psi)$

$$\Gamma_{ee}^A \mathcal{B}(A \rightarrow \pi^+\pi^- J/\psi) = \frac{N_A}{\varepsilon_A \mathcal{L} I_A \mathcal{B}(J/\psi \rightarrow \ell^+\ell^-)}$$

- with $I_A = 12\pi \Gamma_{\text{tot}} \int_{x_1}^{x_2} dx \frac{W(s, x)}{(s' - M_A^2)^2 + \Gamma_{\text{tot}}^2 M_A^2}$ and $x = \sqrt{1 - M(\pi^+\pi^- J/\psi)^2/s}$
- Set an upper limit at the 90% confidence level (C.L.)
- Four likelihood curves from fit: $L_i(\gamma)$, $i = 1 \dots 4$, and $\gamma = \Gamma_{ee}^{X(3872)} \mathcal{B}(X(3872) \rightarrow \pi^+\pi^- J/\psi)$
- Look for γ_i^{up} in $\int_0^{\gamma_i^{\text{up}}} d\gamma L_i(\gamma) = 0.9 \int_0^\infty d\gamma L_i(\gamma)$
- Combining the four measurements: Look for $\gamma_{\text{tot}}^{\text{up}}$ in the product of the single likelihood curves

$$\int_0^{\gamma_{\text{tot}}^{\text{up}}} d\gamma \prod_{i=1}^4 L_i(\gamma) = 0.9 \int_0^\infty d\gamma \prod_{i=1}^4 L_i(\gamma)$$

RESULTS

- $\Gamma_{ee}^{X(3872)} \mathcal{B}(X(3872) \rightarrow \pi\pi J/\psi) < 0.13$ eV at 90% C.L., improves recent results ≈ 60
- $\Gamma_{ee}^{\psi(3686)} = (2213 \pm 18_{\text{stat}} \pm 99_{\text{sys}})$ eV

SYSTEMATIC UNCERTAINTIES

Source	$\sigma_{\text{sys}}^{X(3872)} [\%]$	$\sigma_{\text{sys}}^{\psi(3686)} [\%]$
Luminosity	1.0	1.0
Tracking	4.0	4.0
J/ψ selection	0.2	0.2
Kinematic Fit	0.4	0.4
Integrals I_A	0.7	0.7
Branching ratio	0.5	1.4
$X(3872)$ width	2.7	-
ISR simulation	3.4	-
$\psi(3686)$ fit model	-	1.0
Total	6.1	4.5