

Recent Hadronic Cross Section Measurements from *BABAR*

Konrad Griessinger
on behalf of the *BABAR* Collaboration

Institute for Nuclear Physics
Mainz University

International Workshop on e^+e^- Collisions from Phi to Psi,
June 2017



BABAR

TM and © Nefvana, All Rights Reserved

JG|U

Outline

- 1 Introduction
- 2 Cross section $e^+e^- \rightarrow \pi^+\pi^-\pi^0\pi^0$
- 3 Cross section $e^+e^- \rightarrow \pi^+\pi^-\eta$
- 4 Cross sections $e^+e^- \rightarrow K_S K_L \pi^0, K_S K_L \pi^0 \pi^0, K_S K_L \eta$

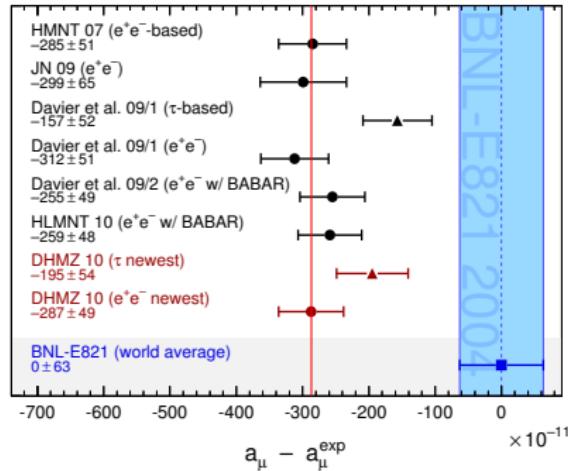
The contributions to a_μ and its uncertainty

$$\vec{\mu} = g \frac{e}{2m} \vec{s}$$

$$(g_\mu - 2)/2 =: a_\mu^{\text{SM}} = a_\mu^{\text{QED}} + a_\mu^{\text{weak}} + a_\mu^{\text{hadronic}}$$

Interaction	Contribution [$\cdot 10^{-11}$]	Uncertainty [$\cdot 10^{-11}$]
QED [1]	116 584 718.951	0.080
EW [9]	153.6	1
hadronic VP [6, 12]	6837	43
hadronic LbL [11, 2]	119	41
total theory	116 591 828	60
E821 experiment [15]	116 592 089	63
deviation exp-theo	261	87

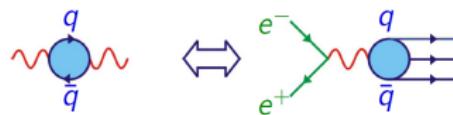
Discrepancy between SM prediction and direct measurement from Eur.Phys.J., C71:1515, 2011 [6].



Just a fluctuation?

3 σ effect, thus reduction of uncertainties necessary!

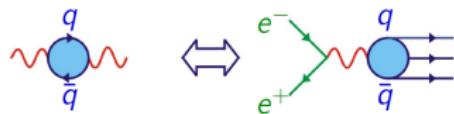
Connection between a_μ and σ_{had}



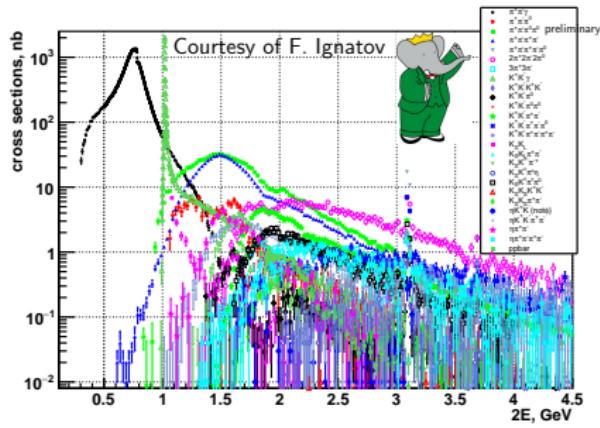
$$a_\mu^{\text{had}} \approx \frac{1}{4\pi^3} \int_{m_\pi^2}^{\infty} K_\mu(s) \cdot \sigma_{e^+e^- \rightarrow \text{had}}(s) ds$$

Kernel function cross section

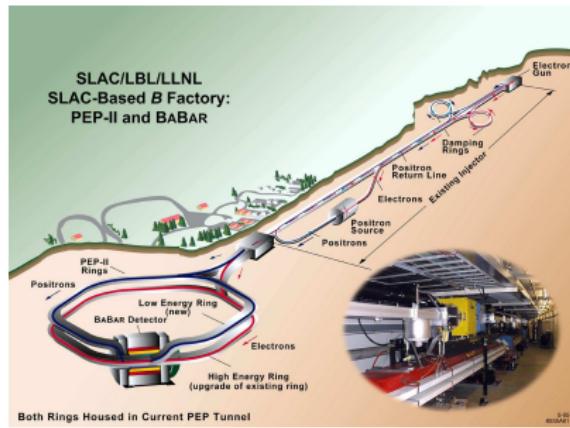
Connection between a_μ and σ_{had}



$$a_\mu^{\text{had}} \approx \frac{1}{4\pi^3} \int_{m_\pi^2}^\infty K_\mu(s) \cdot \sigma_{e^+ e^- \rightarrow \text{had}}(s) ds$$



The *BABAR* Experiment

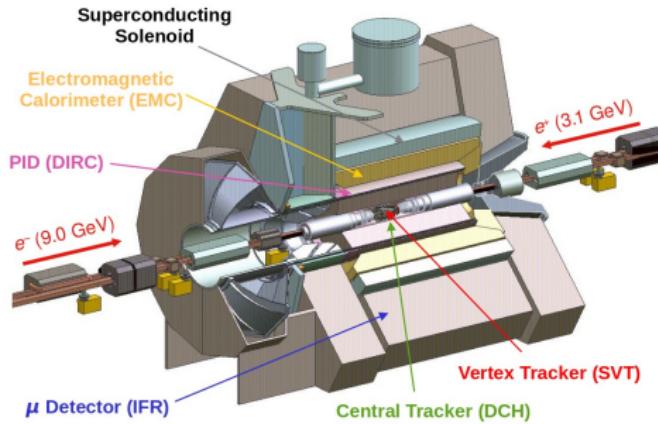


Experimental specifications

Energy: $\sqrt{s} \approx 10.58 \text{ GeV}$ ($E_{e^-} \approx 9.0 \text{ GeV}$, $E_{e^+} \approx 3.1 \text{ GeV}$),

Luminosity: $\mathcal{L} \approx 500 \text{ fb}^{-1}$ ($\Upsilon(4S)$)

The *BABAR* Experiment

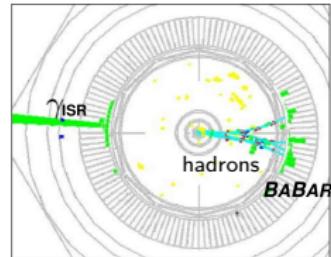
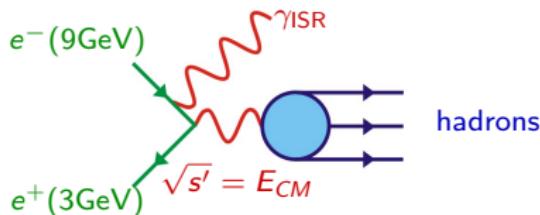


Experimental specifications

Energy: $\sqrt{s} \approx 10.58 \text{ GeV}$ ($E_{e^-} \approx 9.0 \text{ GeV}$, $E_{e^+} \approx 3.1 \text{ GeV}$),

Luminosity: $\mathcal{L} \approx 500 \text{ fb}^{-1}$ ($\Upsilon(4S)$)

Initial State Radiation (ISR) events at *BABAR*



ISR selection

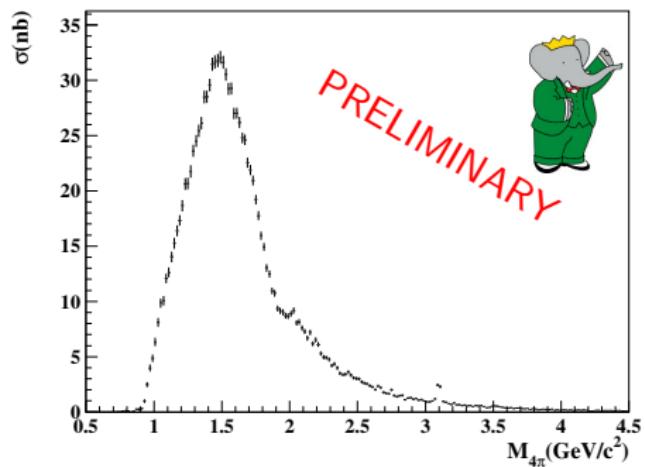
- Detected high energy photon: $E_\gamma > 3\text{GeV}$
→ defines E_{CM} & provides strong background rejection
- Event topology: γ_{ISR} back-to-back to hadrons
→ high acceptance
- Kinematic fit including γ_{ISR}
→ very good energy resolution (4 – 15MeV)
- Continuous measurement from threshold to $\sim 5\text{GeV}$
→ provides common, consistent systematic uncertainties

$$e^+ e^- \rightarrow \pi^+ \pi^- \pi^0 \pi^0$$

PRELIMINARY

Resulting cross section

$$e^+ e^- \rightarrow \pi^+ \pi^- \pi^0 \pi^0$$

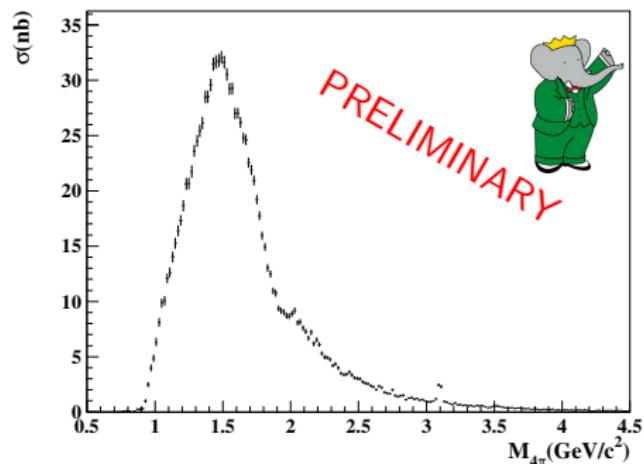


- dominant ISR-bkg $\pi^+ \pi^- 3\pi^0$ removed using data
- most precise measurement to date
- widest energy range $0.85 < E_{\text{CM}} < 4.5 \text{ GeV}$

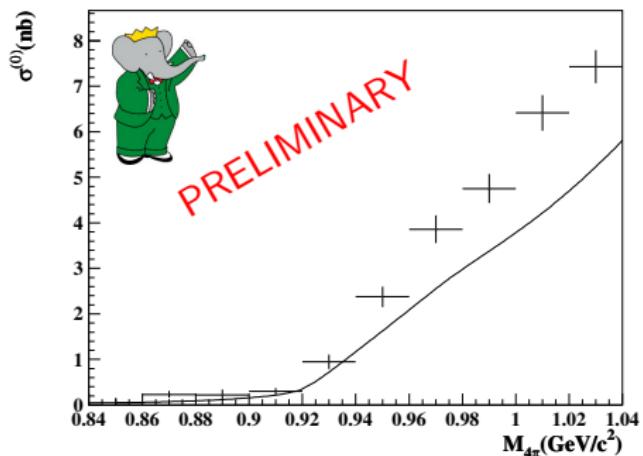
$E_{\text{CM}}(\text{GeV})$	Syst. unc.
1.2 – 2.7	3.1%
2.7 – 3.2	6.7%
> 3.2	7.1%

Resulting cross section

$$e^+ e^- \rightarrow \pi^+ \pi^- \pi^0 \pi^0$$



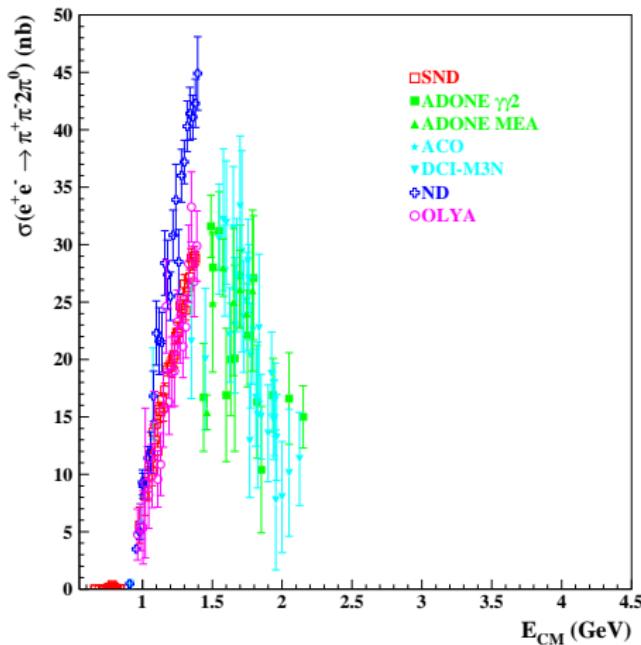
$E_{\text{CM}}(\text{GeV})$	Syst. unc.
1.2 – 2.7	3.1%
2.7 – 3.2	6.7%
> 3.2	7.1%



Comparison to Chiral Pert. Theo.

(Eur.Phys.J., C24:535–545, 2002 [8])

Contribution of $\pi^+ \pi^- 2\pi^0$ to $g_\mu - 2$

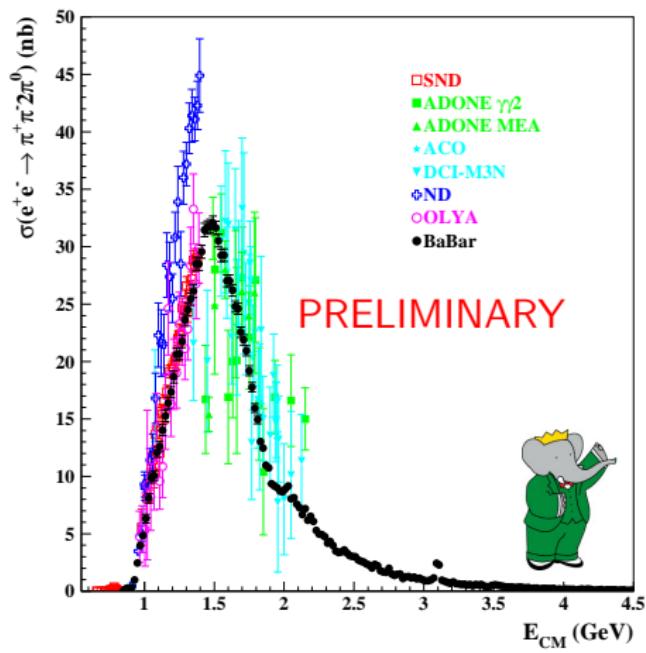


$$a_\mu^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} \frac{\sqrt{1 - \frac{4m_e^2}{s}}}{1 + \frac{2m_e^2}{s}} K_\mu(s) \sigma(s) ds$$

Before *BABAR* (Eur.Phys.J., C31:503,2003) [5]

$$a_\mu(1.02 < \sqrt{s} < 1.8 \text{ GeV}) = \\ (16.76 \pm 1.31 \pm 0.20_{\text{rad}}) \times 10^{-10}$$

Contribution of $\pi^+ \pi^- 2\pi^0$ to $g_\mu - 2$



$$a_\mu^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} \frac{\sqrt{1 - \frac{4m_e^2}{s}}}{1 + \frac{2m_e^2}{s}} K_\mu(s) \sigma(s) ds$$

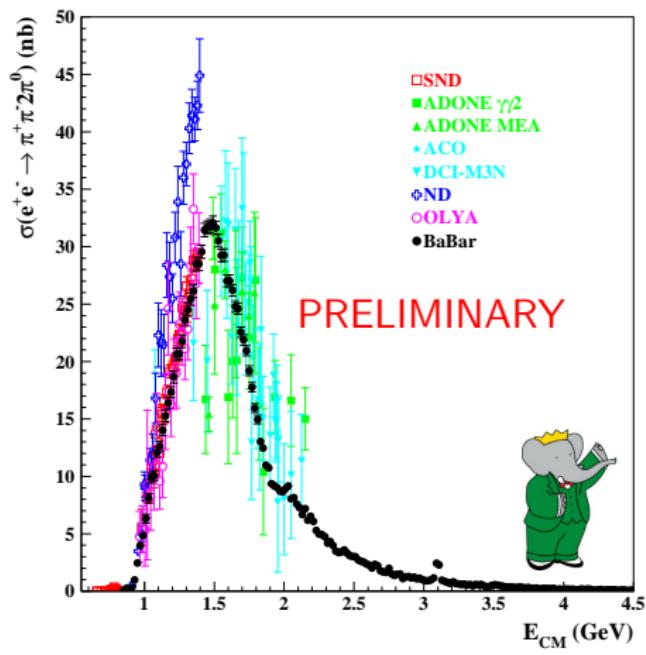
Before *BABAR* (Eur.Phys.J., C31:503,2003) [5]

$$a_\mu(1.02 < \sqrt{s} < 1.8 \text{ GeV}) = (16.76 \pm 1.31 \pm 0.20_{\text{rad}}) \times 10^{-10}$$

New result in the same energy range

$$a_\mu(1.02 < \sqrt{s} < 1.8 \text{ GeV}) = (17.4 \pm 0.1_{\text{stat}} \pm 0.6_{\text{syst}}) \times 10^{-10}$$

Contribution of $\pi^+ \pi^- 2\pi^0$ to $g_\mu - 2$



$$a_\mu^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} \frac{\sqrt{1 - \frac{4m_e^2}{s}}}{1 + \frac{2m_e^2}{s}} K_\mu(s) \sigma(s) ds$$

New result starting at lower limit

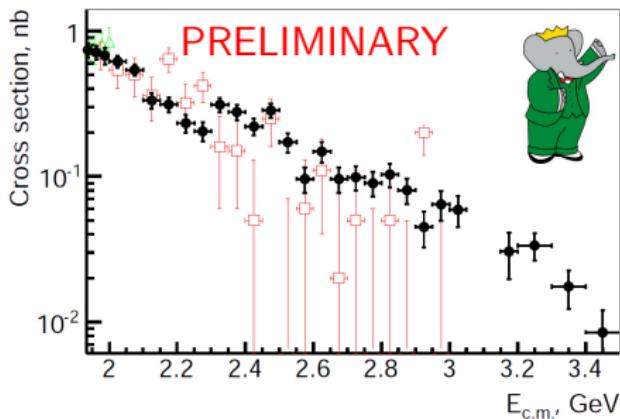
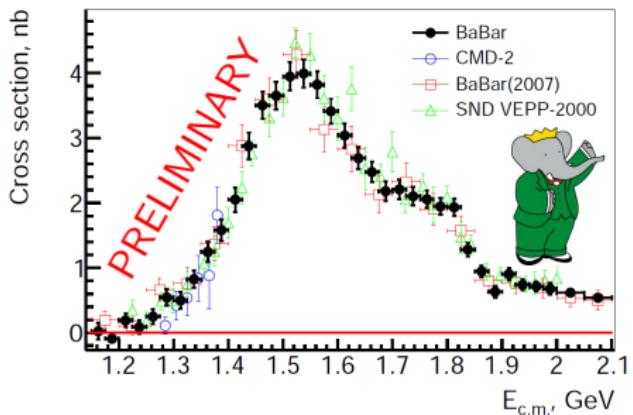
$$a_\mu(0.85 < \sqrt{s} < 1.8 \text{ GeV}) = (17.9 \pm 0.1_{\text{stat}} \pm 0.6_{\text{syst}}) \times 10^{-10}$$

New result in a wider energy range

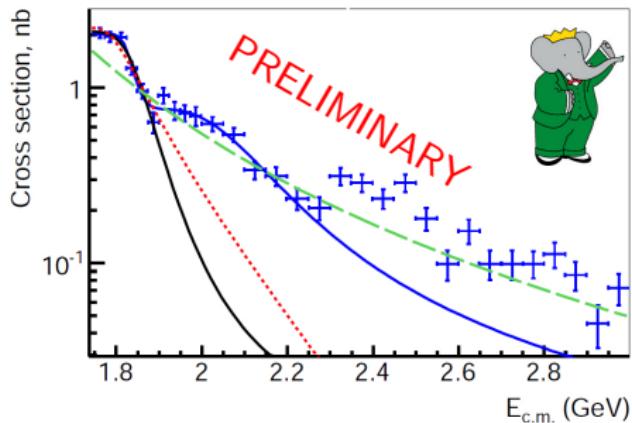
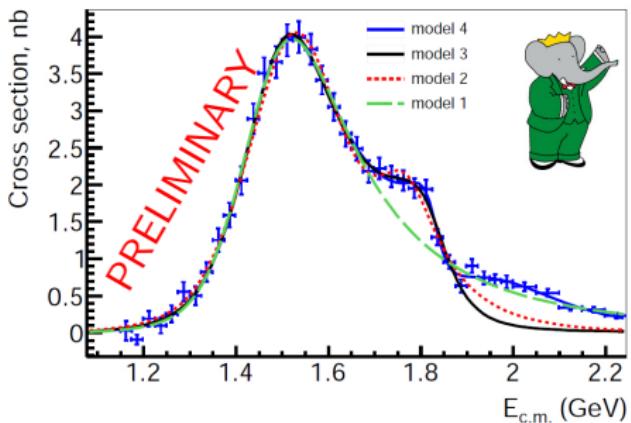
$$a_\mu(0.85 < \sqrt{s} < 3.0 \text{ GeV}) = (21.8 \pm 0.1_{\text{stat}} \pm 0.7_{\text{syst}}) \times 10^{-10}$$

$$e^+e^- \rightarrow \pi^+\pi^-\eta$$

PRELIMINARY

Cross section $e^+e^- \rightarrow \pi^+\pi^-\eta$ 

- Most accurate $\sigma(e^+e^- \rightarrow \pi^+\pi^-\eta)$ measurement to date
- First measurement up to 3.5 GeV
- Especially above 1.6 GeV more precise than previous data

Fits to the cross section $e^+e^- \rightarrow \pi^+\pi^-\eta$ 

- Model 1: $\rho(770) - \rho(1450)$, fit: $E_{\text{CM}} < 1.7$ GeV
- Model 2: $\rho(770) - \rho(1450) - \rho(1700)$, fit: $E_{\text{CM}} < 1.9$ GeV
- Model 3: $\rho(770) - \rho(1450) + \rho(1700)$, fit: $E_{\text{CM}} < 1.9$ GeV
- Model 4: $\rho(770) - \rho(1450) + \rho(1700) + \rho(2150)$, fit: $E_{\text{CM}} < 2.2$ GeV

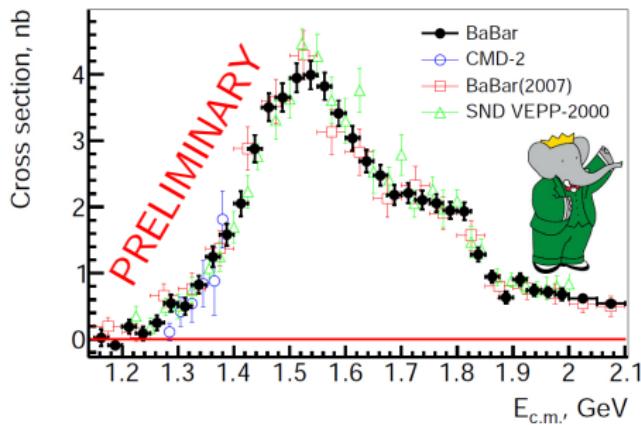
"+": relative phase 0° , "-": relative phase 180°

Contribution of $\pi^+ \pi^- \eta$ to $g_\mu - 2$

$$a_\mu^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} \frac{\sqrt{1 - \frac{4m_e^2}{s}}}{1 + \frac{2m_e^2}{s}} K_\mu(s) \sigma(s) ds$$

HLMNT 2011 [10]

$$a_\mu(\sqrt{s} < 1.8 \text{ GeV}) = (0.88 \pm 0.10) \times 10^{-10}$$



DHMZ 2011 [6]

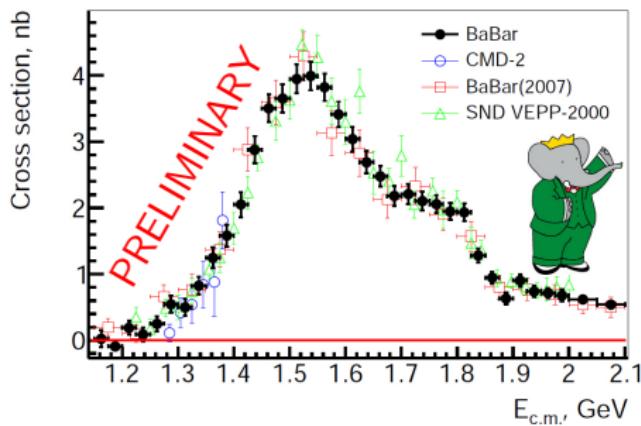
$$a_\mu(\sqrt{s} < 1.8 \text{ GeV}) = (1.15 \pm 0.06_{\text{stat}} \pm 0.08_{\text{syst}}) \times 10^{-10}$$

Contribution of $\pi^+ \pi^- \eta$ to $g_\mu - 2$

$$a_\mu^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} \frac{\sqrt{1 - \frac{4m_e^2}{s}}}{1 + \frac{2m_e^2}{s}} K_\mu(s) \sigma(s) ds$$

HLMNT 2011 [10]

$$a_\mu(\sqrt{s} < 1.8 \text{ GeV}) = (0.88 \pm 0.10) \times 10^{-10}$$



DHMZ 2011 [6]

$$a_\mu(\sqrt{s} < 1.8 \text{ GeV}) = (1.15 \pm 0.06_{\text{stat}} \pm 0.08_{\text{syst}}) \times 10^{-10}$$

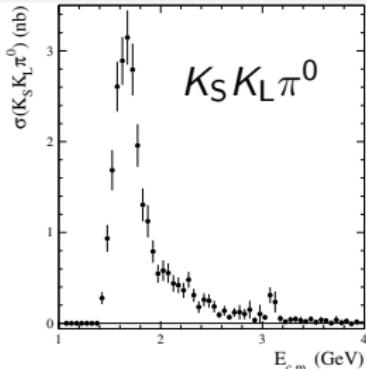
New result

$$a_\mu(\sqrt{s} < 1.8 \text{ GeV}) = (1.19 \pm 0.02_{\text{stat}} \pm 0.06_{\text{syst}}) \times 10^{-10}$$

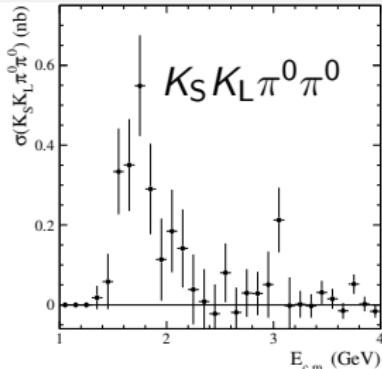
$$e^+e^- \rightarrow K_S K_L \pi^0, K_S K_L \pi^0 \pi^0, K_S K_L \eta$$
$$e^+e^- \rightarrow K_S K^\pm \pi^\mp \pi^0, K_S K^\pm \pi^\mp \eta$$

PRD 95 (2017), 052001
PRD 95 (2017), 092005

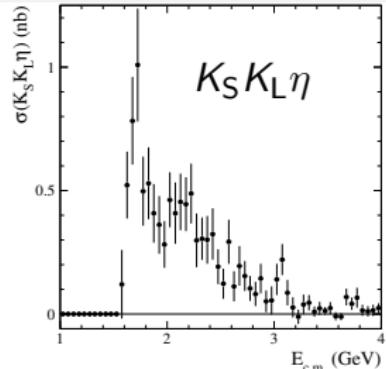
$$e^+ e^- \rightarrow K_S K_L \pi^0, K_S K_L \pi^0 \pi^0, K_S K_L \eta$$



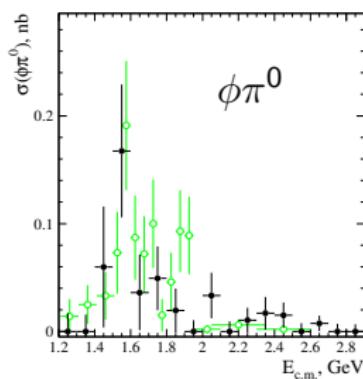
syst.: 10–30 %



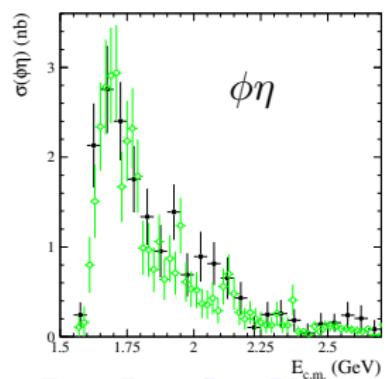
syst.: 25–60 %



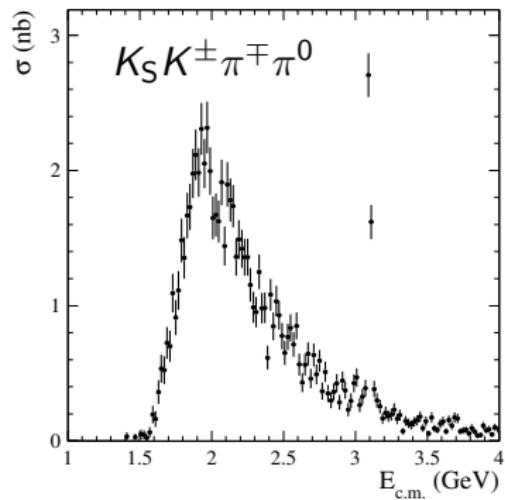
syst.: 15–30 %



Published in
PRD 95 (2017), 052001



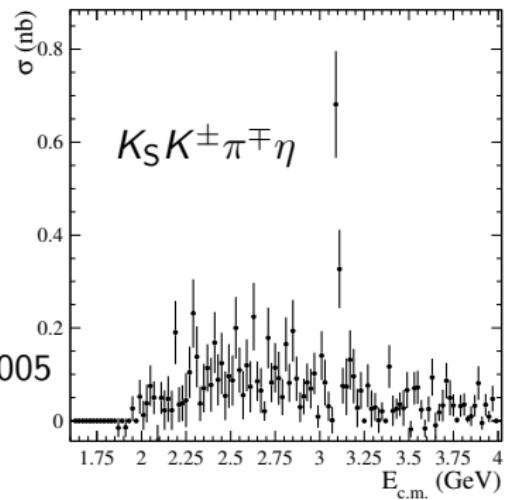
$$e^+e^- \rightarrow K_S K^\pm \pi^\mp \pi^0, K_S K^\pm \pi^\mp \eta$$



Published in
PRD 95 (2017), 092005

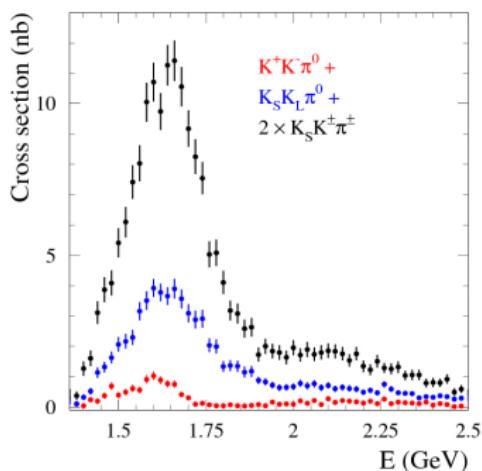


total unc.: 6–12 %

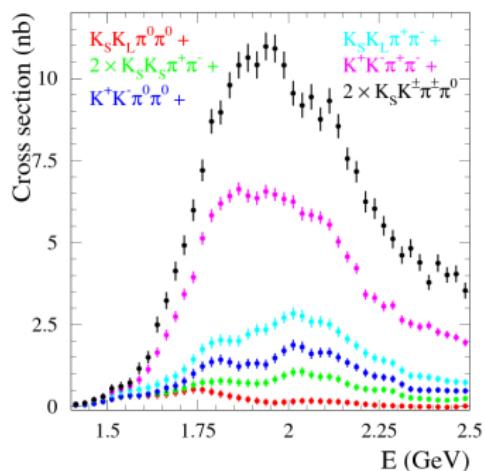


total unc.: 13–19 %

Total $e^+e^- \rightarrow KK\pi(\pi)$ cross sections



Courtesy of
V. Druzhinin [7]



- all $KK\pi(\pi)$ channels now measured by *BABAR*
 \Rightarrow isospin relations not necessary any more
- new contributions to $g_\mu - 2$:
 - ★ $a_\mu(KK\pi) = (2.45 \pm 0.15) \times 10^{-10}$ [4]
 - ★ $a_\mu(KK\pi\pi) = (0.85 \pm 0.05) \times 10^{-10}$

Summary

- ISR physics has proven to be a very productive field even years after the end of data taking at the B-factories
- Precision measurements of hadronic cross sections have greatly improved a_μ^{SM} & more hadronic final states in preparation
- New results from *BABAR*:
 - ★ $e^+e^- \rightarrow \pi^+\pi^-\pi^0\pi^0$
 - ★ $e^+e^- \rightarrow \pi^+\pi^-\eta$
 - ★ first measurement of
 $e^+e^- \rightarrow K_S K_L \pi^0, K_S K_L \pi^0\pi^0, K_S K_L \eta, K_S K^\pm \pi^\mp \pi^0, K_S K^\pm \pi^\mp \eta$



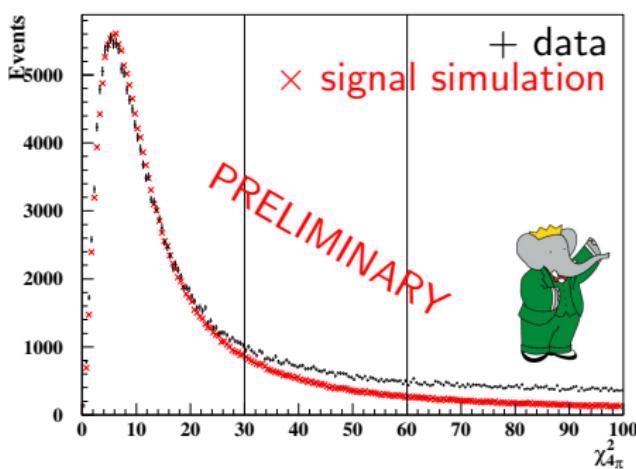
Thank you!
Any questions?

Backup

$\pi^+\pi^-2\pi^0\gamma$ Event Selection

Full $\Upsilon(4S)$ on peak data set of 454.4 fb^{-1}

$$e^+ e^- \rightarrow \pi^+\pi^-\pi^0\pi^0\gamma_{\text{ISR}}$$



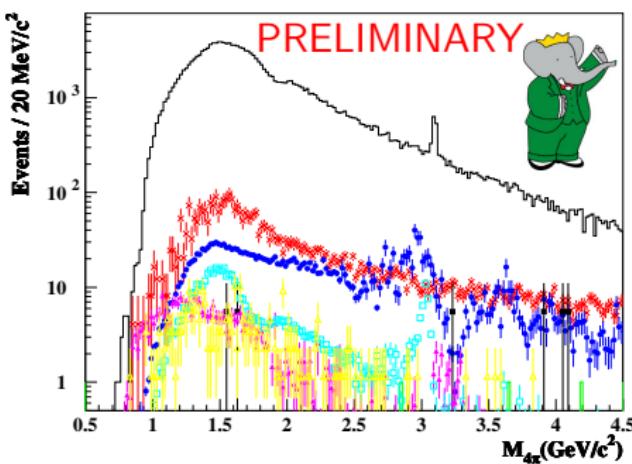
Main Selection Requirements

- exactly 2 charged tracks
- ≥ 5 photons
- $E_\gamma^{\text{lab}} > 0.05 \text{ GeV}$
- $|M_{\pi^0}^{\text{reco}} - M_{\pi^0}^{\text{PDG}}| < 0.03 \text{ GeV}$
- $E_{\gamma_{\text{ISR}}} > 3 \text{ GeV}$
- 6C kinematic fit: $\chi^2_{2\pi 2\pi^0\gamma} < 30$
- reject other hypotheses
- Muon and Kaon PID

Background subtraction

$$e^+ e^- \rightarrow \pi^+ \pi^- \pi^0 \pi^0 \gamma_{\text{ISR}}$$

Simulated background channels:

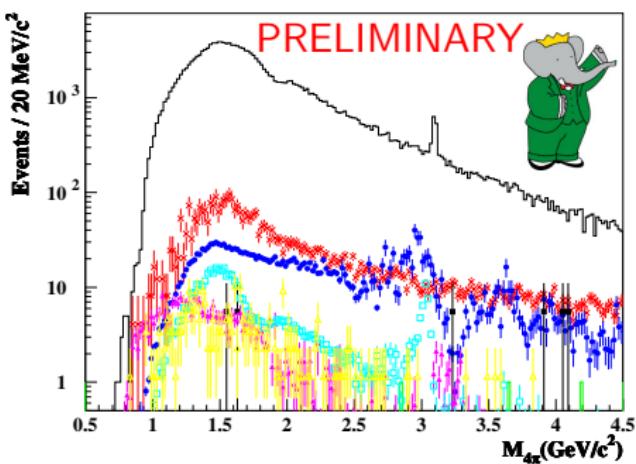


$q\bar{q}$, 3π , $4\pi 2\pi^0$, $K_s K\pi$, $K^+ K^- 2\pi^0$, $\tau\tau$,
 $\pi^+ \pi^- 3\pi^0$

Background subtraction

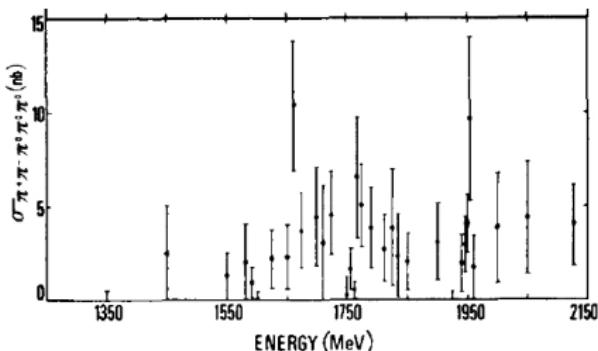
$$e^+ e^- \rightarrow \pi^+ \pi^- \pi^0 \pi^0 \gamma_{\text{ISR}}$$

Simulated background channels:



$q\bar{q}$, 3π , $4\pi 2\pi^0$, $K_s K\pi$, $K^+ K^- 2\pi^0$, $\tau\tau$,
 $\pi^+ \pi^- 3\pi^0$

Main issue: background from
 $e^+ e^- \rightarrow \pi^+ \pi^- 3\pi^0$

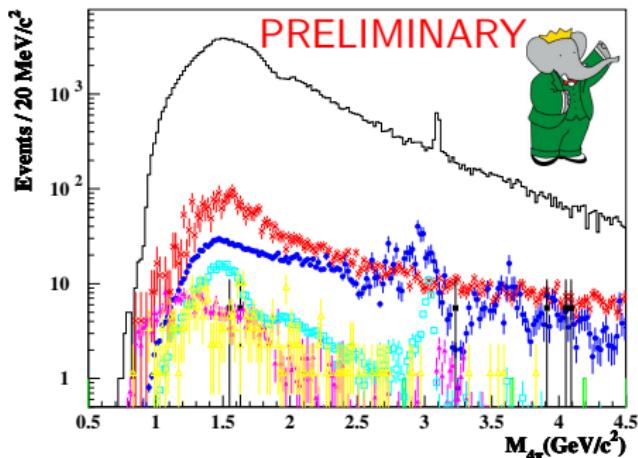


Only little data [3] and no full simulation available

Background subtraction

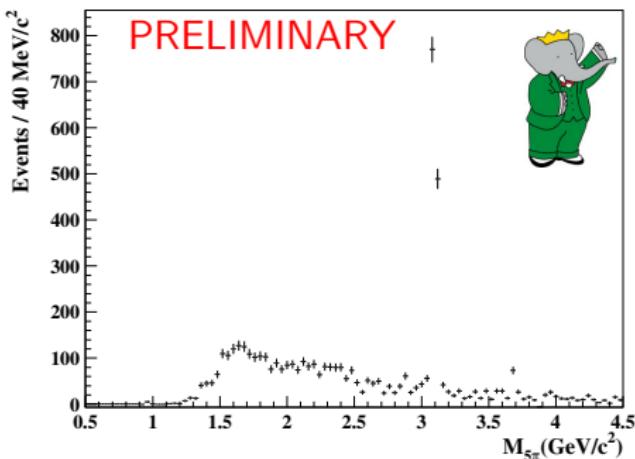
$$e^+ e^- \rightarrow \pi^+ \pi^- \pi^0 \pi^0 \gamma_{\text{ISR}}$$

Simulated background channels:



$q\bar{q}$, 3π , $4\pi 2\pi^0$, $K_s K\pi$, $K^+ K^- 2\pi^0$, $\tau\tau$,
 $\pi^+ \pi^- 3\pi^0$

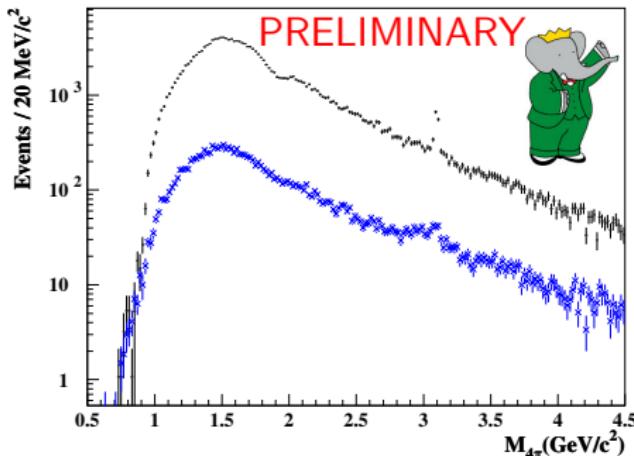
Main issue: background from
 $e^+ e^- \rightarrow \pi^+ \pi^- 3\pi^0$



\Rightarrow dedicated *BABAR* measurement
 \Rightarrow adjust simulation

Background subtraction: cross check

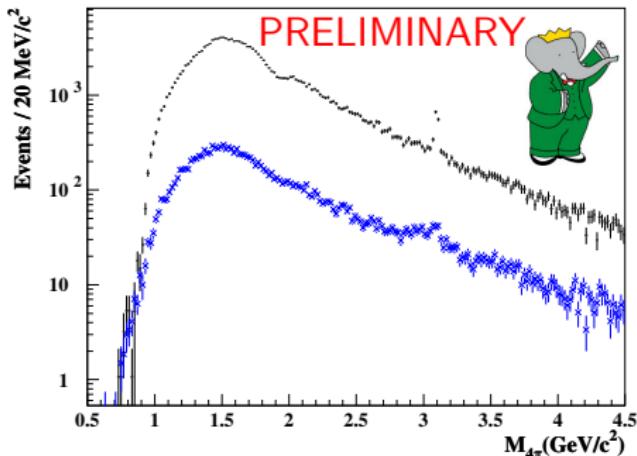
Sideband bkg subtraction



$$N_{1s} = \frac{\beta}{\beta-\alpha} \cdot N_1 - \frac{1}{\beta-\alpha} \cdot N_2$$
$$\alpha := \frac{N_{2s}}{N_{1s}}, \quad \beta := \frac{N_{2b}}{N_{1b}}$$

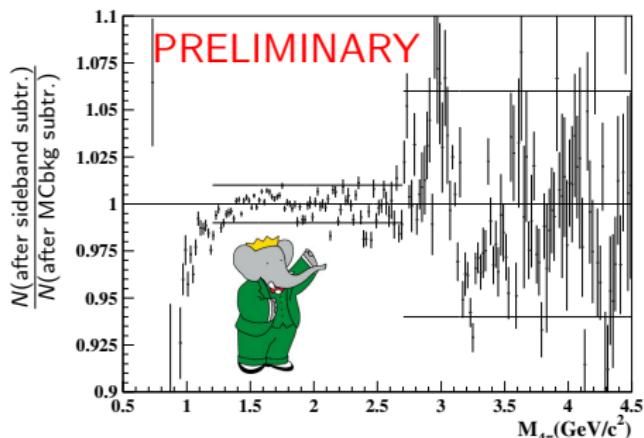
Background subtraction: cross check

Sideband bkg subtraction



$$N_{1s} = \frac{\beta}{\beta - \alpha} \cdot N_1 - \frac{1}{\beta - \alpha} \cdot N_2$$
$$\alpha := \frac{N_{2s}}{N_{1s}}, \quad \beta := \frac{N_{2b}}{N_{1b}}$$

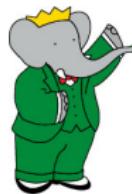
Comparison of both methods



Less than 1 % discrepancy in the peak region around 1.5 GeV/c²

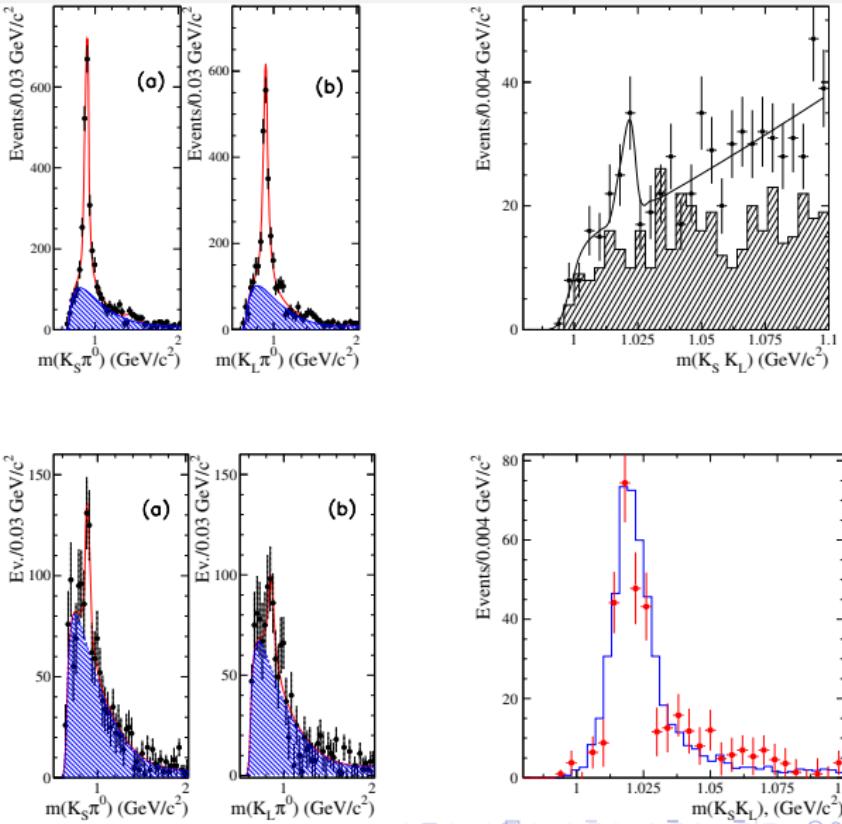
$e^+ e^- \rightarrow K_S K_L \pi^0, K_S K_L \eta$ intermediate st.: $K^*(892)$ and ϕ

$e^+ e^- \rightarrow K_S K_L \pi^0$

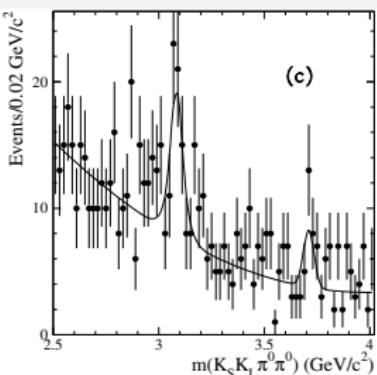
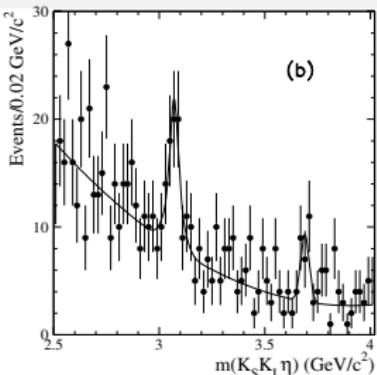
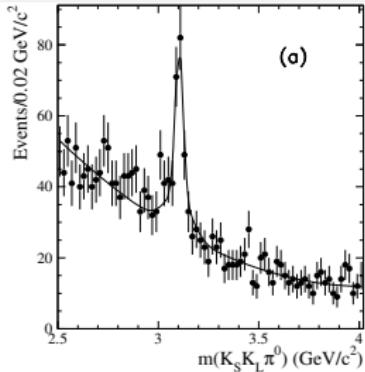


Published in
PRD 95 (2017), 052001

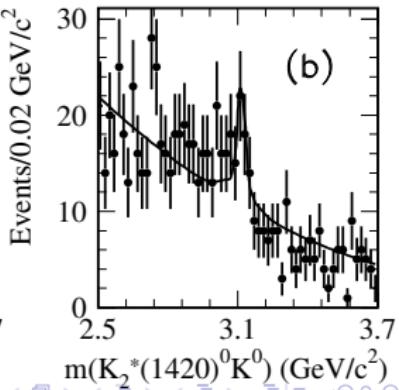
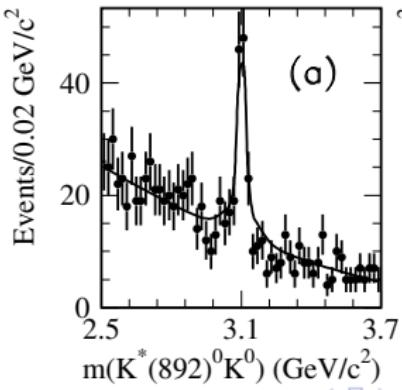
$e^+ e^- \rightarrow K_S K_L \eta$



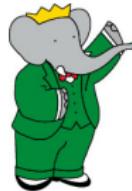
$e^+ e^- \rightarrow K_S K_L \pi^0, K_S K_L \eta, K_S K_L \pi^0 \pi^0$, int. st.: ψ and K^*



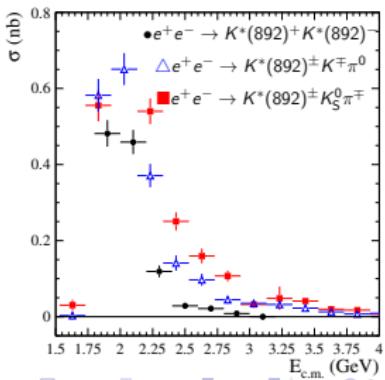
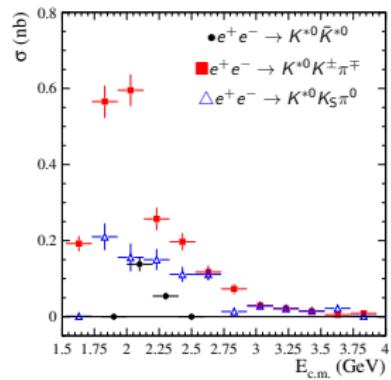
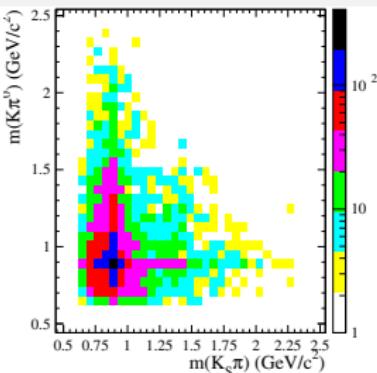
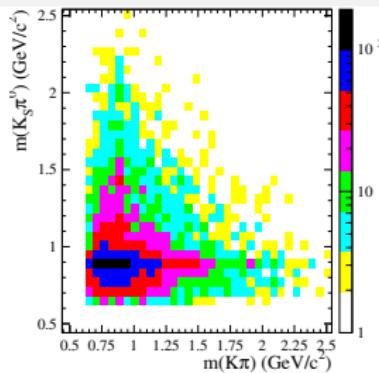
Published in
PRD 95 (2017), 052001



$e^+ e^- \rightarrow K_S K^\pm \pi^\mp \pi^0$ intermediate states



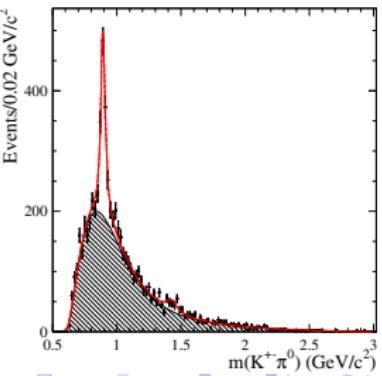
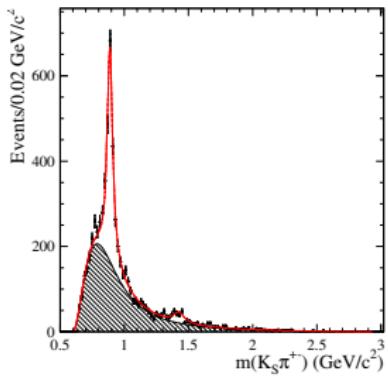
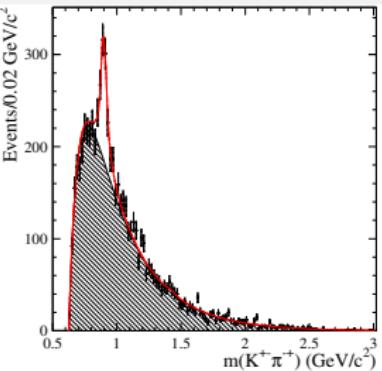
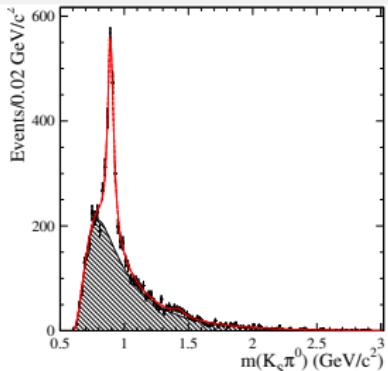
Published in
PRD 95 (2017), 092005



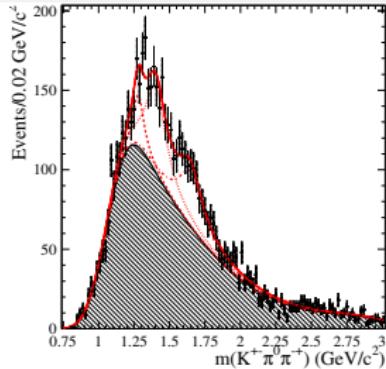
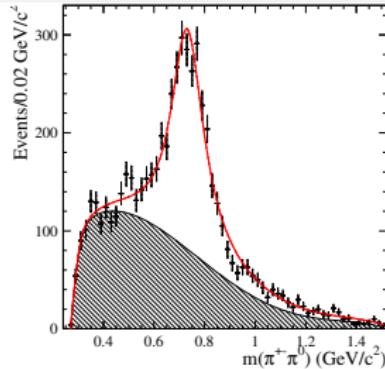
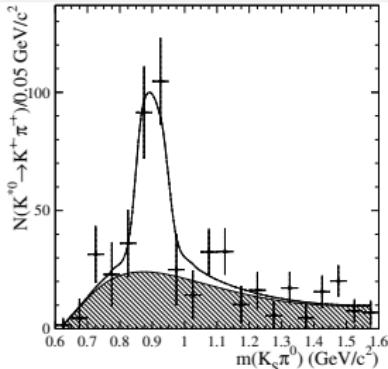
$e^+ e^- \rightarrow K_S K^\pm \pi^\mp \pi^0$ intermediate states



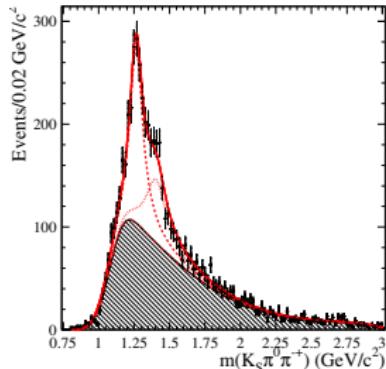
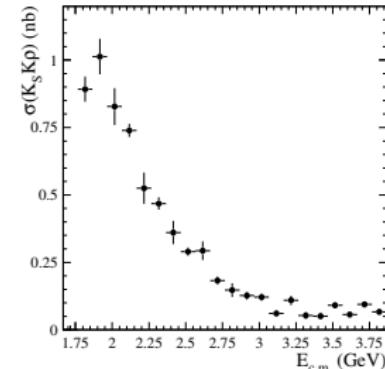
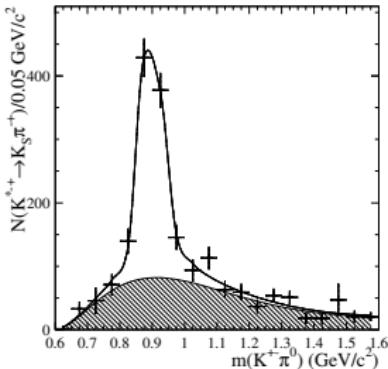
Published in
PRD 95 (2017), 092005



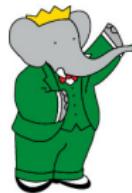
$e^+e^- \rightarrow K_S K^\pm \pi^\mp \pi^0$ intermediate states



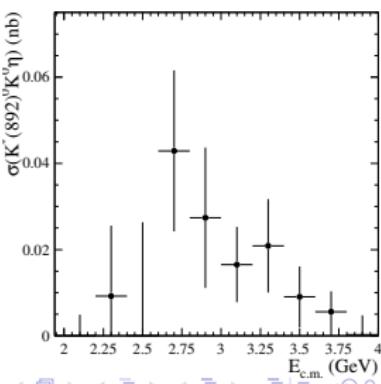
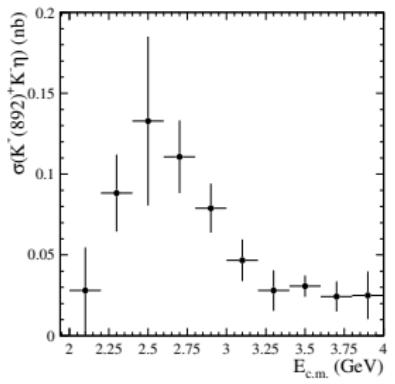
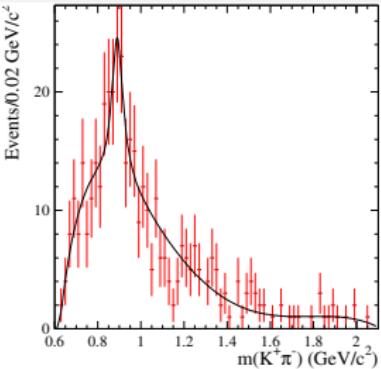
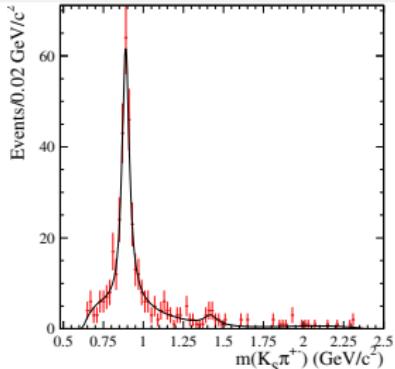
Published in PRD 95 (2017), 092005



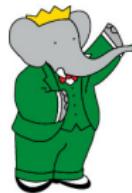
$e^+ e^- \rightarrow K_S K^\pm \pi^\mp \eta$ intermediate states



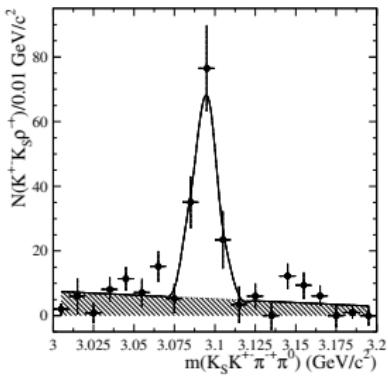
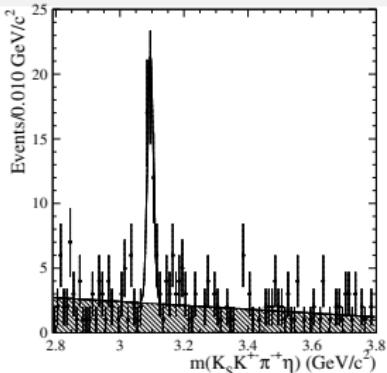
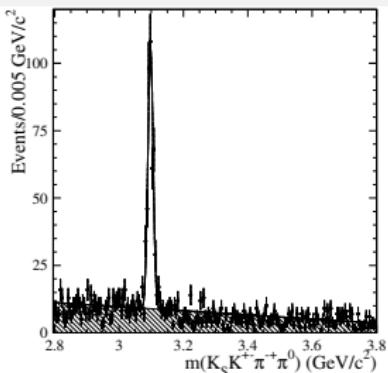
Published in
PRD 95 (2017), 092005



$e^+ e^- \rightarrow K_S K^\pm \pi^\mp \pi^0 / \eta$ intermediate states: J/ψ



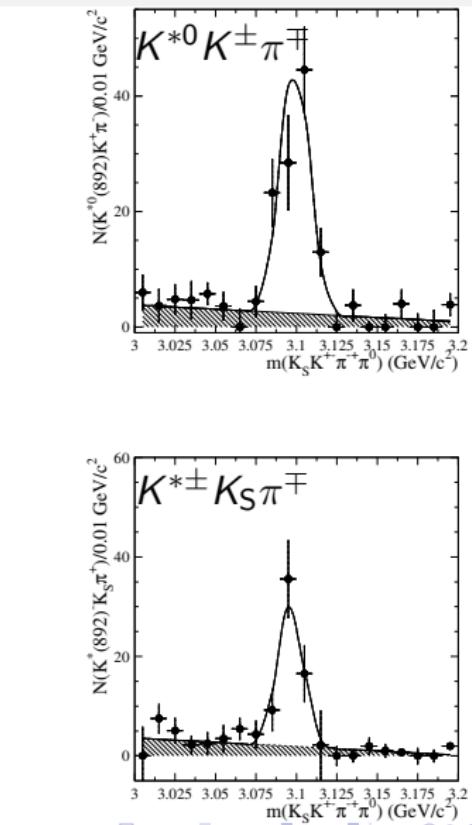
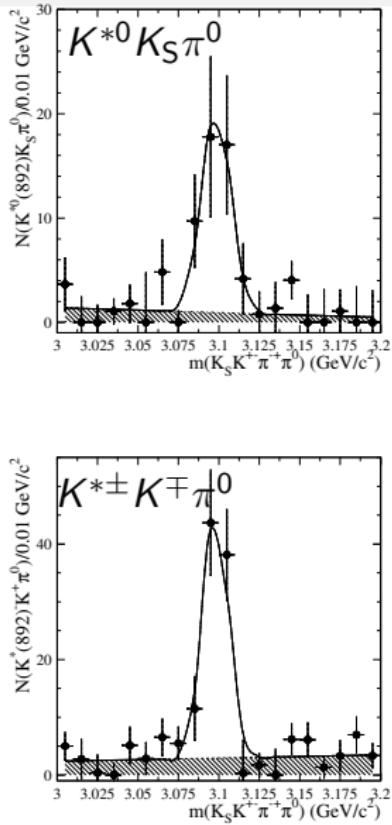
Published in
PRD 95 (2017), 092005



$e^+ e^- \rightarrow K_S K^\pm \pi^\mp \pi^0$ intermediate states: K^*



Published in
PRD 95 (2017), 092005



References

References I

- [1] T. Aoyama, M. Hayakawa, T. Kinoshita, and M. Nio.
Complete tenth-order qed contribution to the muon $g - 2$.
Phys. Rev. Lett., 109:111808, Sep 2012.
- [2] G. Colangelo, M. Hoferichter, A. Nyffeler, M. Passera, and P. Stoffer.
Remarks on higher-order hadronic corrections to the muon $g - 2$.
Phys.Lett., B735:90–91, 2014.
- [3] G. Cosme, B. Dudelzak, B. Grelaud, B. Jean-Marie, S. Jullian, et al.
Hadronic Cross-Sections Study in e^+e^- Collisions from 1.350GeV to 2.125GeV.
Nucl.Phys., B152:215, 1979.

References II

- [4] M. Davier.
Update of the hadronic vacuum polarization contribution to the muon $g - 2$.
Nuclear and Particle Physics Proceedings, 287-288:70 – 75, 2017.
The 14th International Workshop on Tau Lepton Physics.
- [5] M. Davier, S. Eidelman, A. Höcker, and Z. Zhang.
Updated estimate of the muon magnetic moment using revised results from e^+e^- annihilation.
Eur. Phys. J., C31:503–510, 2003.
- [6] M. Davier, A. Hoecker, B. Malaescu, and Z. Zhang.
Reevaluation of the Hadronic Contributions to the Muon $g - 2$ and to $\alpha(M_Z)$.
Eur.Phys.J., C71:1515, 2011.

References III

- [7] V. P. Druzhinin.
Study of e^+e^- annihilation to hadrons at low energies at BaBar.
EPJ Web Conf., 142:01013, 2017.
- [8] G. Ecker and R. Unterdorfer.
Four-Pion Production in e^+e^- Annihilation.
Eur. Phys. J., C24:535–545, 2002.
- [9] C. Gnendiger, D. Stöckinger, and H. Stöckinger-Kim.
The electroweak contributions to $(g - 2)_\mu$ after the Higgs boson mass measurement.
Phys. Rev., D88(5):053005, 2013.
- [10] K. Hagiwara, R. Liao, A. D. Martin, D. Nomura, and T. Teubner.
 $(g - 2)_\mu$ and $\alpha(M_Z^2)$ re-evaluated using new precise data.
J. Phys., G38:085003, 2011.

References IV

- [11] F. Jegerlehner and A. Nyffeler.
The Muon g-2.
Phys.Rept., 477:1–110, 2009.
- [12] A. Kurz, T. Liu, P. Marquard, and M. Steinhauser.
Hadronic contribution to the muon anomalous magnetic moment to next-to-next-to-leading order.
Phys.Lett., B734:144–147, 2014.
- [13] J. P. Lees et al.
Cross sections for the reactions $e^+e^- \rightarrow K_S^0 K_L^0 \pi^0$, $K_S^0 K_L^0 \eta$, and $K_S^0 K_L^0 \pi^0 \pi^0$ from events with initial-state radiation.
Phys. Rev., D95(5):052001, 2017.

References V

[14] J. P. Lees et al.

Measurement of the $e^+e^- \rightarrow K_s^0 K^\pm \pi^\mp \pi^0$ and $K_s^0 K^\pm \pi^\mp \eta$ cross sections using initial-state radiation.

Phys. Rev., D95(9):092005, 2017.

[15] K. Olive et al.

Review of Particle Physics.

Chin.Phys., C38:090001, 2014.