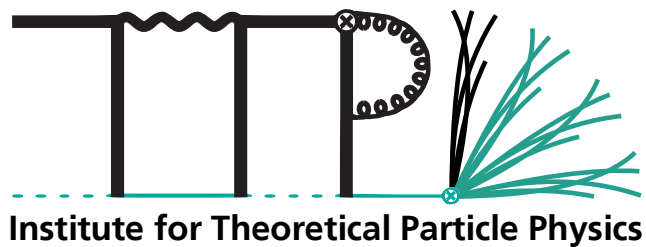


Closing the window on neutralino dark matter with a 100 TeV hadron collider



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Karlsruher Institut für Technologie (KIT).

R.Mahbubani, P. Schwaller, JZ, [arXiv 1609.abcde]

Effective Field Theories for Collider Physics, Flavor Phenomena
and Electroweak Symmetry Breaking, Eltville am Rhein, 13.09.2016

Outline

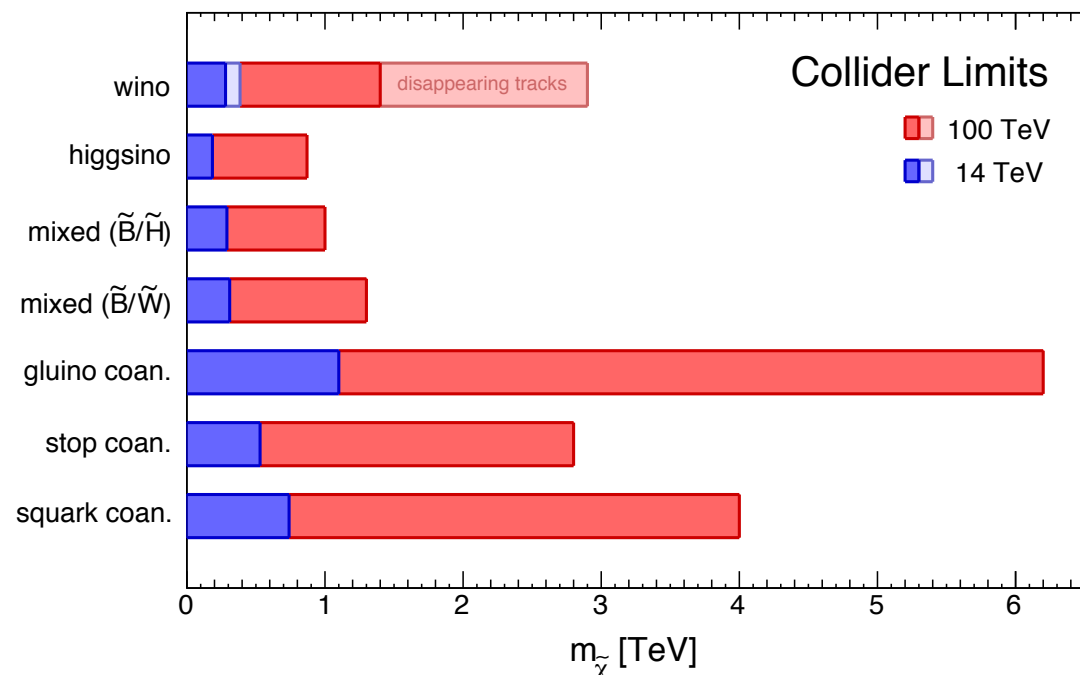
- Motivation
- Phenomenology of Bino/Higgsino DM
- FCC searches
- Conclusions

TeV Higgsinos: MSSM's last stand

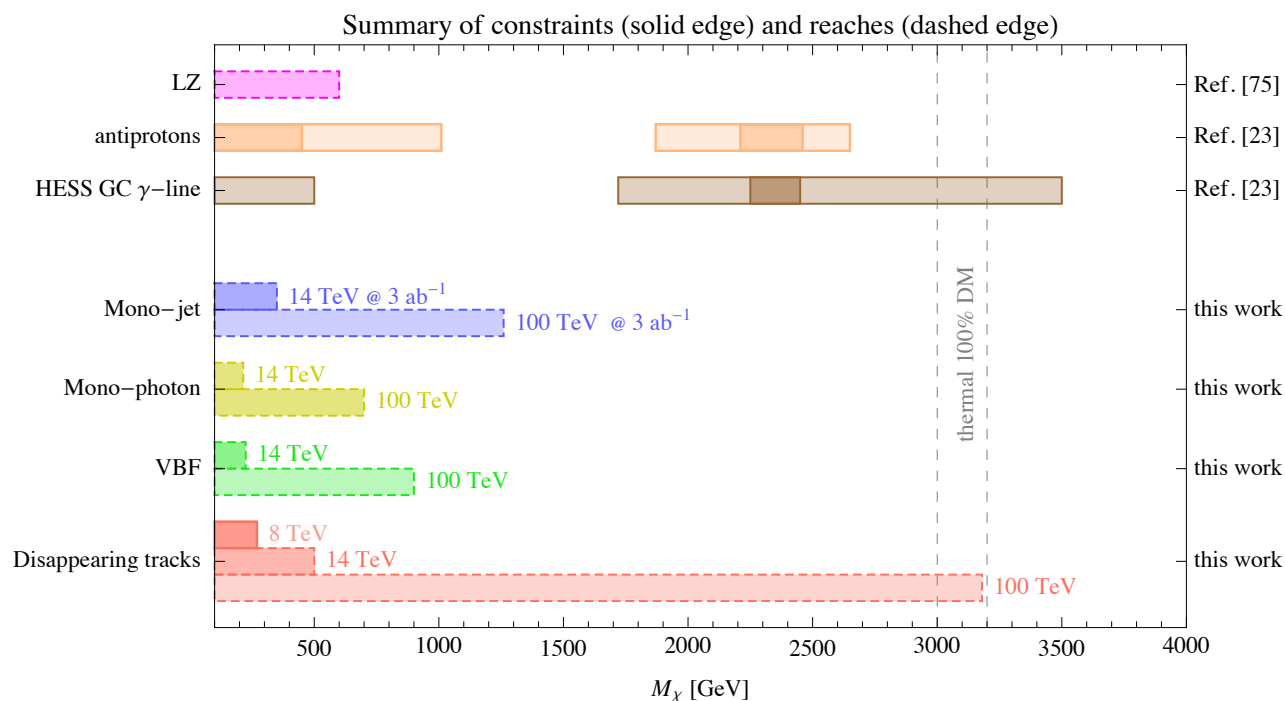
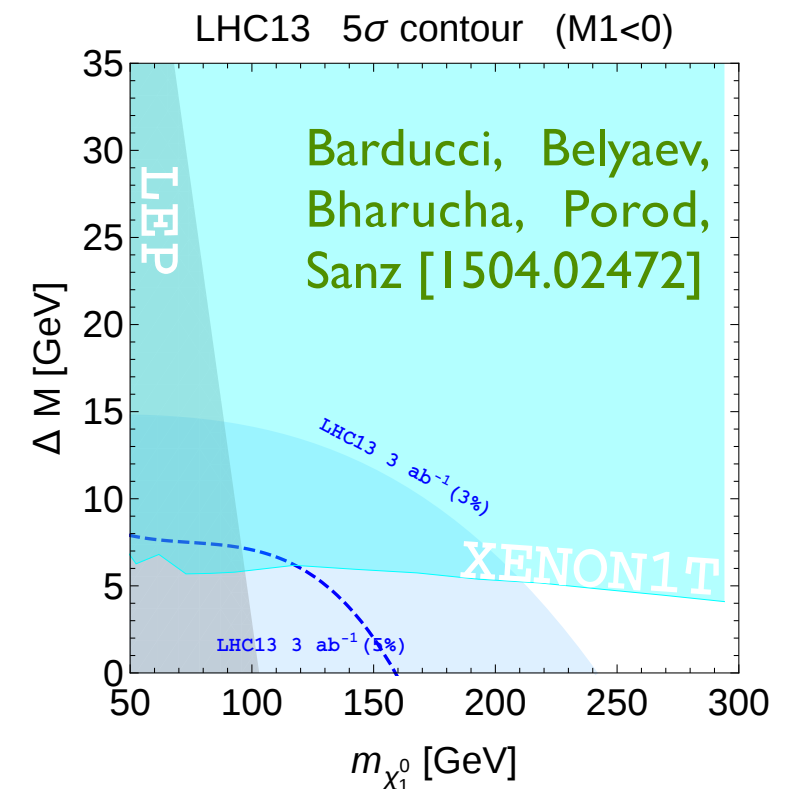
- The neutralino is the MSSM DM candidate, made out of Bino, Wino and Higgsino.
- Relic density sets mass for pure: \tilde{B} (100 GeV), \tilde{W} (2.7 TeV) and \tilde{H} (1.1 TeV).
- All pure cases and mixtures (wino/Higgsino, bino/Wino, etc) well studied in the literature.
- Only “uncoverable” spot is an “almost pure”^{*} Higgsino, mass ~ 1.1 TeV:
 - DM search @ LHC(FCC): mono-jet. Reaches 200 (900) GeV (Low, Wang:1404.0682).
 - Disappearing tracks: Alternative to monojets. Reach limited due to Higgsino lifetime.
- This talk:
 - a) disappearing tracks at FCC with a customised tracker for 1.1 TeV Higgsino
 - b) mono-Z study with highly boosted di-lepton final states (beats mono-jet!)

^{*} A pure Higgsino, EW doublet, is ruled out, because both neutral states are mass degenerate, and the Z - n_1 - n_2 coupling is actually Z - n_1 - n_1 . Z currents with weak couplings are excluded by direct detection experiments (XENON100, LUX, etc). Some additional Bino and/or Wino component is required.

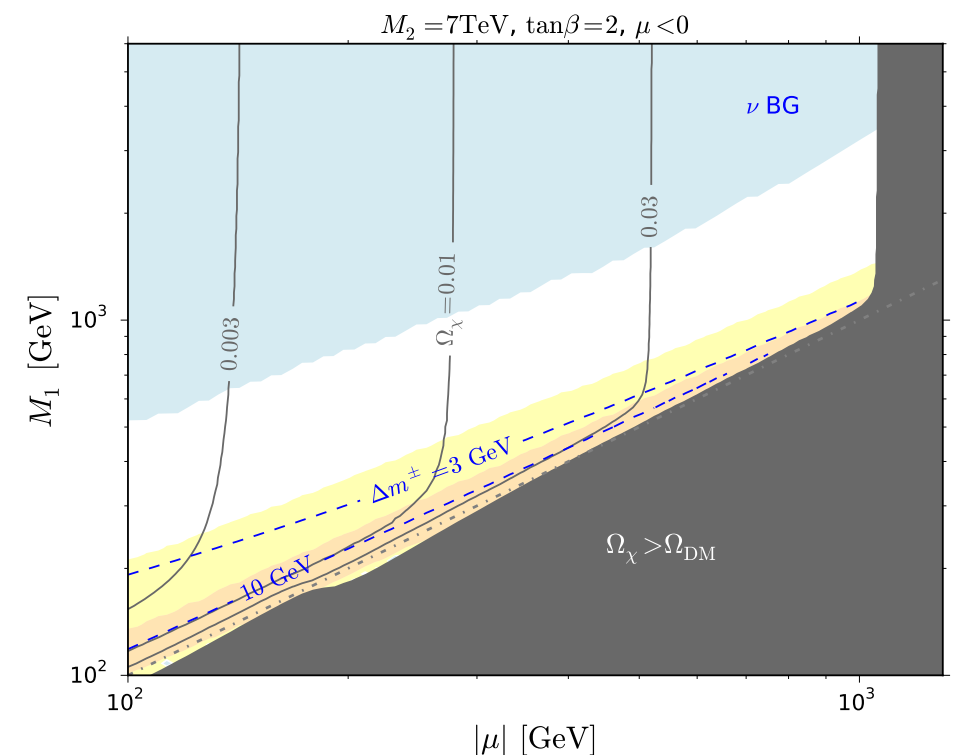
Recent Analyses



M. Low, L.T.Wang [1404.0682]



Cirelli, Sala, Taoso [1407.7058]



Badziak, Delgado, Olechowski, Pokorski, Sakurai [1506.07177]

Higgsino/Bino

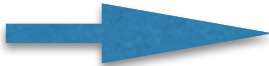
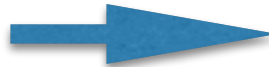
dark matter:

phenomenology

Electroweakino Sector

$$\begin{array}{ccc}
 & \begin{array}{c} \text{gauge} \\ \text{eigenstates} \end{array} & \begin{array}{c} \text{mass} \\ \text{eigenstates} \end{array} \\
 \mathcal{M}_\chi^\pm = \begin{pmatrix} M_2 & \sqrt{2} \sin \beta M_W \\ \sqrt{2} \cos \beta M_W & \mu \end{pmatrix} & \tilde{W}^+, \tilde{H}_u^+ \sim \chi_1^+, \chi_2^+ & \longrightarrow \text{Charged winos and Higgsinos mix into charginos (2).} \\
 \\
 \mathcal{M}_\chi^0 = \begin{pmatrix} M_1 & 0 & -c_\beta s_W M_Z & s_\beta s_W M_Z \\ 0 & M_2 & c_\beta c_W M_Z & -s_\beta c_W M_Z \\ -c_\beta s_W M_Z & c_\beta c_W M_Z & 0 & -\mu \\ s_\beta s_W M_Z & -s_\beta c_W M_Z & -\mu & 0 \end{pmatrix} & \begin{array}{c} \text{gauge} \\ \text{eigenstates} \end{array} \tilde{H}_u^0, \tilde{H}_d^0, \tilde{B}^0, \tilde{W}_3^0 \sim \begin{array}{c} \text{mass} \\ \text{eigenstates} \end{array} \chi_1^0, \dots, \chi_4^0 & \searrow \text{Neutral Wino, neutral Higgsinos and bino mix into neutralinos (4).}
 \end{array}$$


Too much hassle! Simplify by:

- 1) decoupling the Wino ($M_2 \rightarrow \infty$)  1 neutral and 1 charged state removed
- 2) taking a Higgsino/Bino hierarchy ($|\mu| \ll M_1$)  1 neutral state removed

Important!

Since EW symmetry is broken, in an EW multiplet neutral components correct their masses due to Z-loops, charged components also have W, γ -loops.

Thomas, Wells, hep-ph/9804359,
Cirelli, Formengo, Strumia, hep-ph/051209



$$\Delta_{1\text{-loop}} \begin{cases} \tilde{H} & 340 \text{ MeV} \\ \tilde{W} & 170 \text{ MeV} \end{cases}$$

Simplified Bino/Higgsino

$$M = \begin{pmatrix} M_1 & -mc_\beta & ms_\beta \\ -mc_\beta & 0 & \mp\mu \\ ms_\beta & \mp\mu & 0 \end{pmatrix} \quad m = m_Z s_W \approx 43.8 \text{ GeV}$$

Expanding in μ/M_1

340 MeV

$\Delta_{1-\text{loop}}$

Δ_+

$$\Delta_+ = \Delta_{1-\text{loop}} + \frac{96 \text{ MeV} (1 \mp s_{2\beta})}{(M_1/10 \text{ TeV})} + \mathcal{O}\left(\frac{|\mu|}{M_1}, \frac{m}{M_1}\right)$$

Δ_0

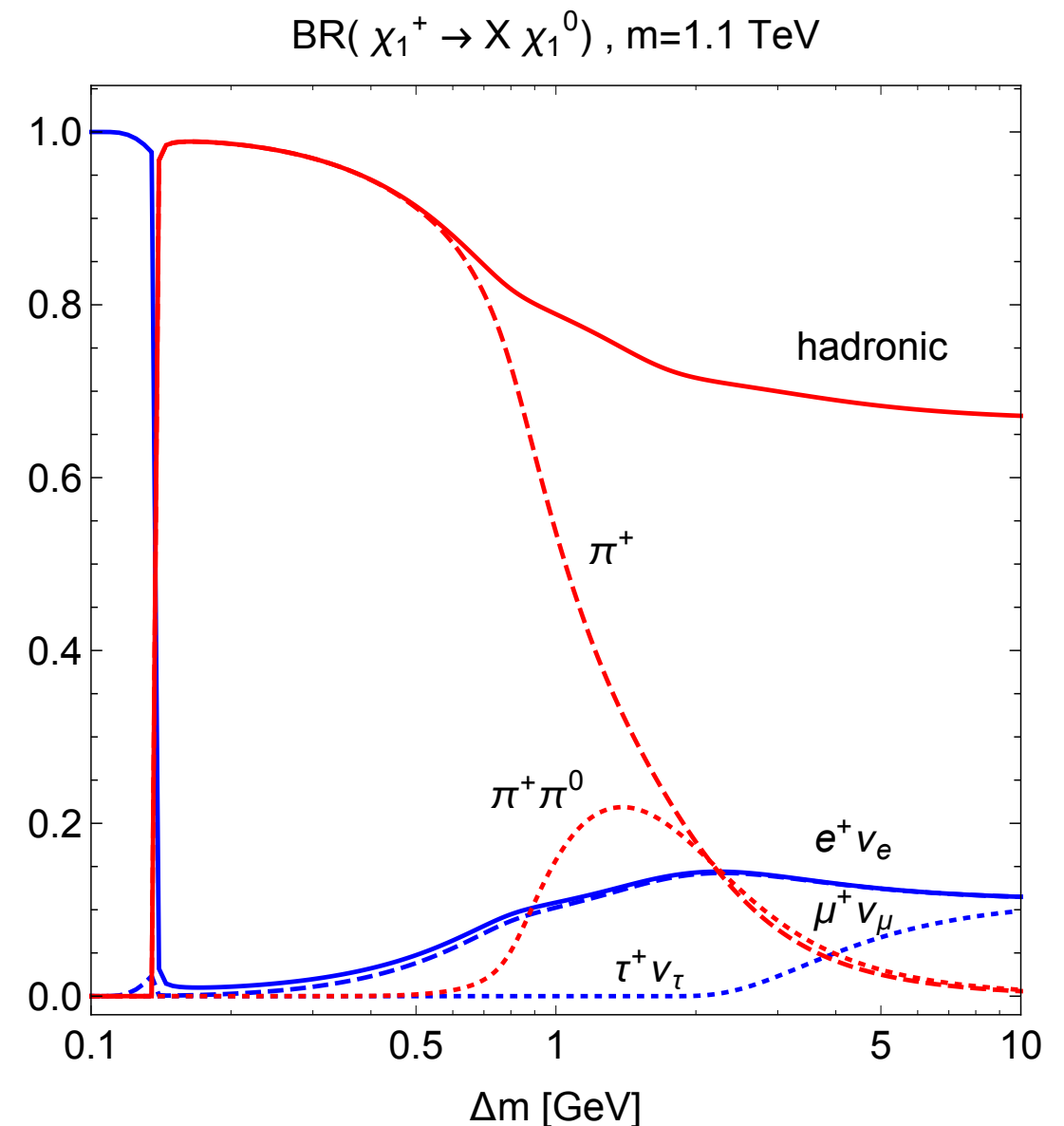
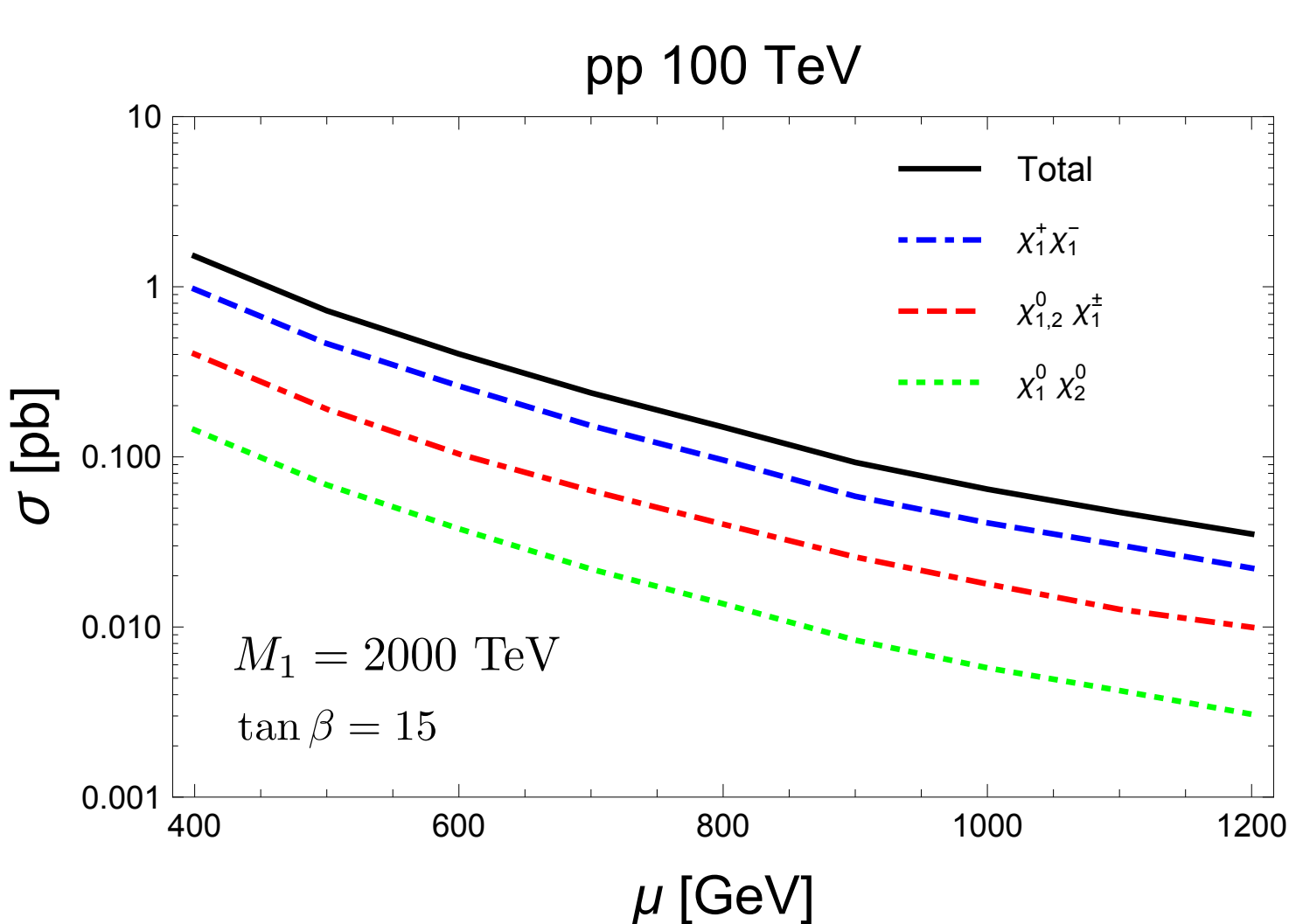
$$\Delta_0 = \frac{192 \text{ MeV}}{(M_1/10 \text{ TeV})} + \mathcal{O}\left(\frac{|\mu|}{M_1}, \frac{m}{M_1}\right)$$

Limiting cases:

1. $\Delta_0 \geq \Delta_+$: decay open only to first neutralino \longrightarrow only for $M_1 \lesssim 3|\mu|$.
2. $\Delta_0 = 0$, $\Delta_+ = 340 \text{ MeV}$: decays to both, lifetime reduced by half.

$\Delta_0 < 100 \text{ KeV}$ gives inelastic scattering @ DD $\longrightarrow M_1 < 20 \text{ PeV}$.

XS and branching ratios



$$\sigma(1.1 \text{ TeV})[\text{fb}] = 47.23 \text{ (39.05) NLO (LO)}.$$

PROSPINO

Beenakker, Klasen, Krämer, Plehn,
Spira, Zerwas, hep-ph/9906298

Decays formulae (mostly) from
Chen, Drees, Gunion: hep-ph/951230, 9607421, 9902302.

$$\Delta_+ = 340 \text{ MeV} \Rightarrow \text{BR}(\chi_1^\pm \rightarrow \pi^\pm \chi_0^{(1)}) \sim 97\%.$$

$$\Gamma(\chi_1^\pm \rightarrow \pi^\pm \chi_0^{(1)}) \sim (\Delta_+)^{-3} \Rightarrow \tau_{\tilde{W}} \sim 8 \tau_{\tilde{H}} \quad \text{Higgsino lifetime too short!}$$

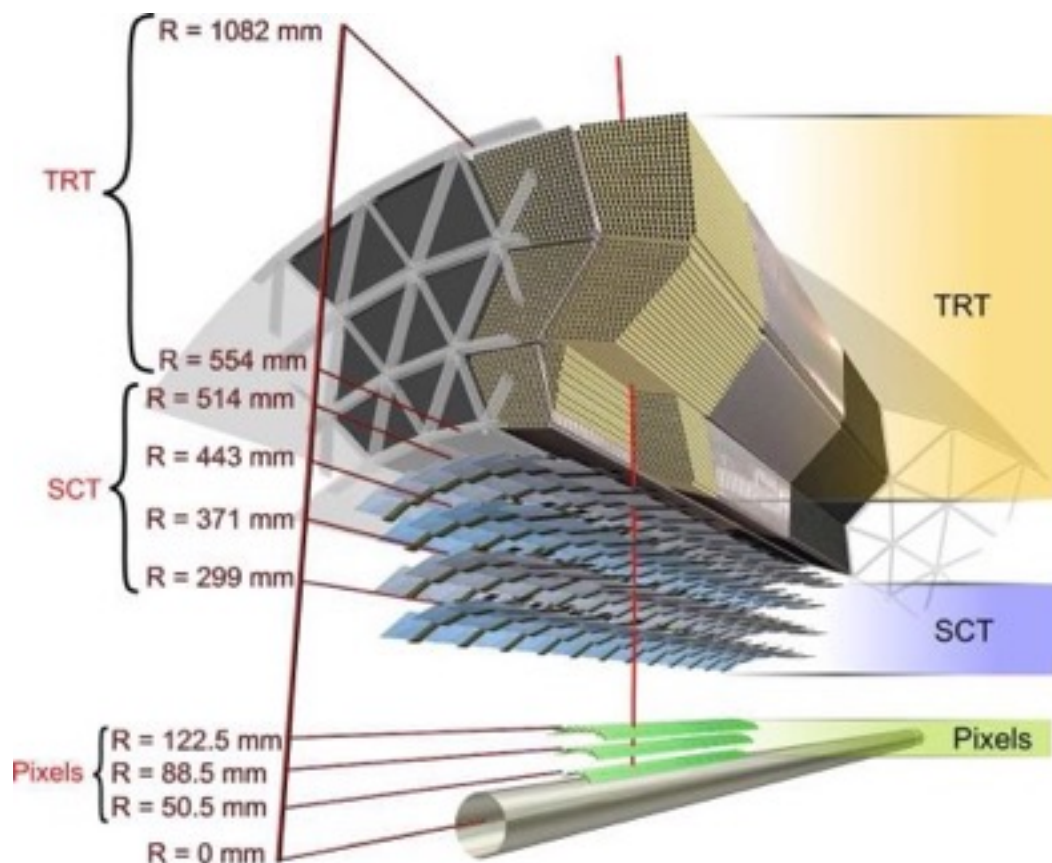
Disappearing tracks

@ 100 TeV FCC-hh

Disappearing tracks @ LHC

ATLAS: CERN-PH-EP-2013-155 [CMS: CERN-PH-EP-2013-037]

- Charged particle (track) decays into neutral + SM (unreconstructed): disappeared!!!
- Event selection requires:
 - 1 “good quality”* (isolated, well reconstructed) track with large p_T .
 - large missing transverse energy ($MET > O(100 \text{ GeV})$).
 - 1 hard jet, $p_T > 100 \text{ GeV}$ (from initial state radiation, to trigger the event).
 - $\Delta\Phi(\text{jet}, MET) > 1.0$ (0.5) @ ATLAS (CMS) : kills mismeasured QCD multijets.



* “Reinheits track”

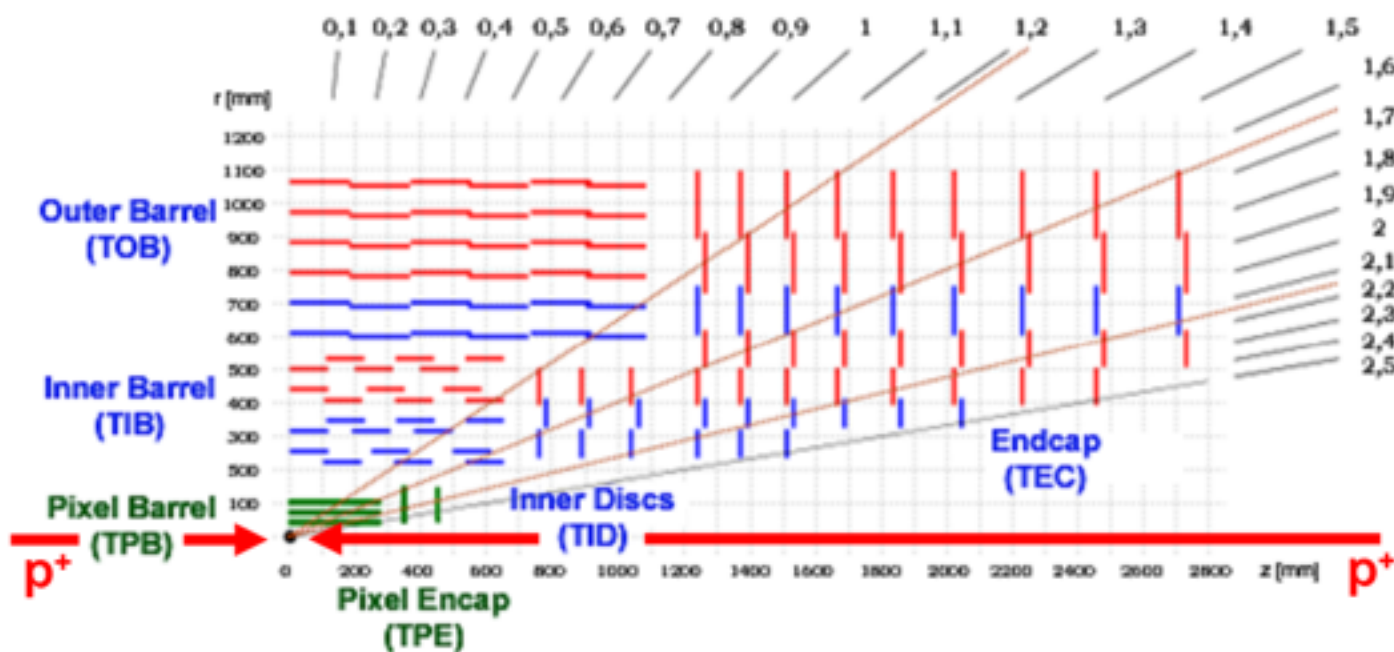
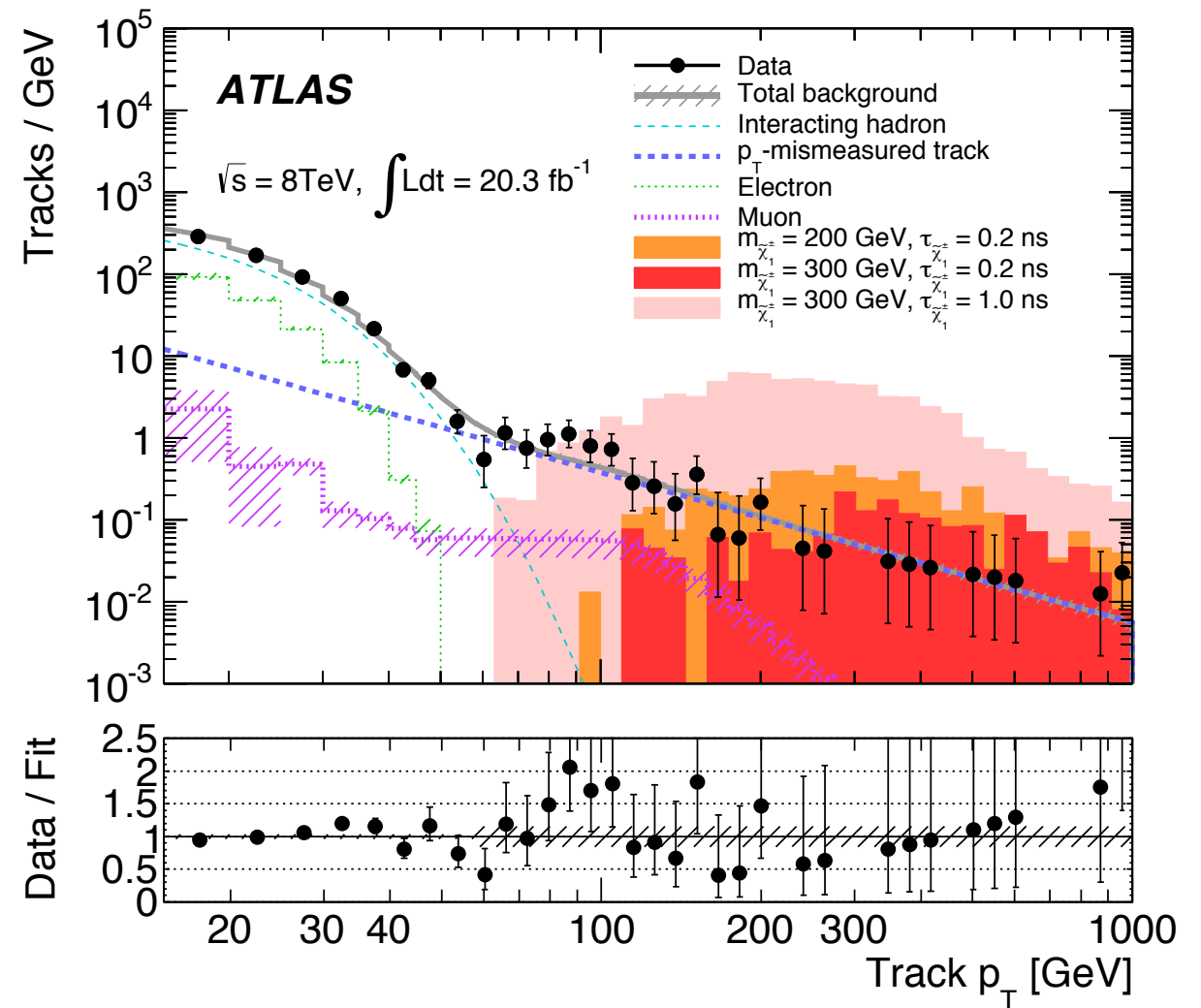
- At least 3 hits in pixel detectors.
 - At least 2 hits in the SCT.
 - Less than 5 hits in the TRT**
 - $p_T > 15 \text{ GeV}, 0.1 < |\eta| < 1.9$ (hard and central)
- $d_{min} \approx 30 \text{ cm}$

** SM particle leaves (on average) 32 hits in TRT

Don't forget about backgrounds!

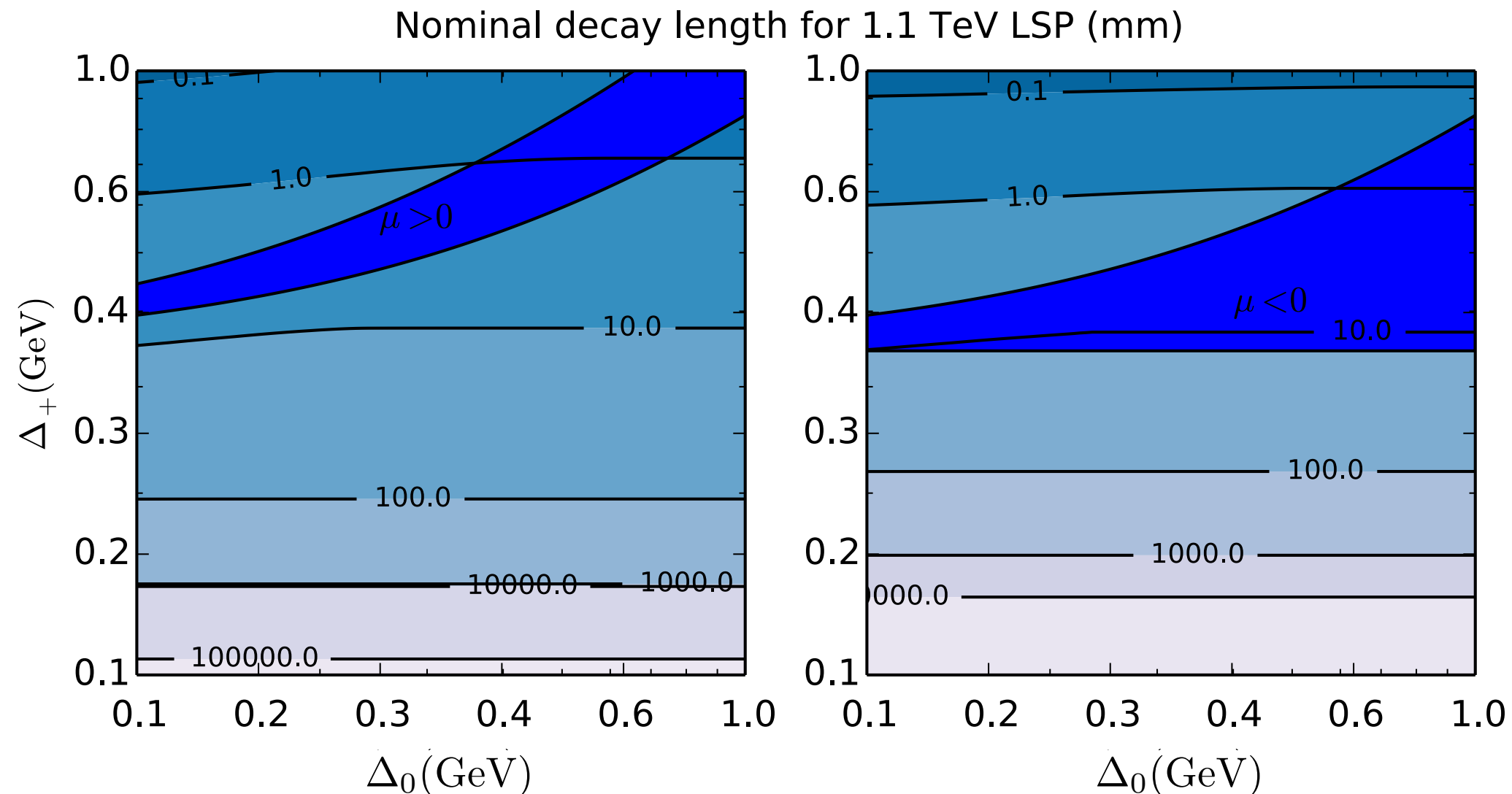
- Background sources:
 - Interacting hadron-tracks
 - Lepton tracks
 - pT-mismeasured tracks
(dominant if $p_T > 100$ GeV)

CMS: cuts on $E_{\text{calo}} < 10$ GeV



Same principle, but...
too messy tracker!!!

Decay lengths and parameter space

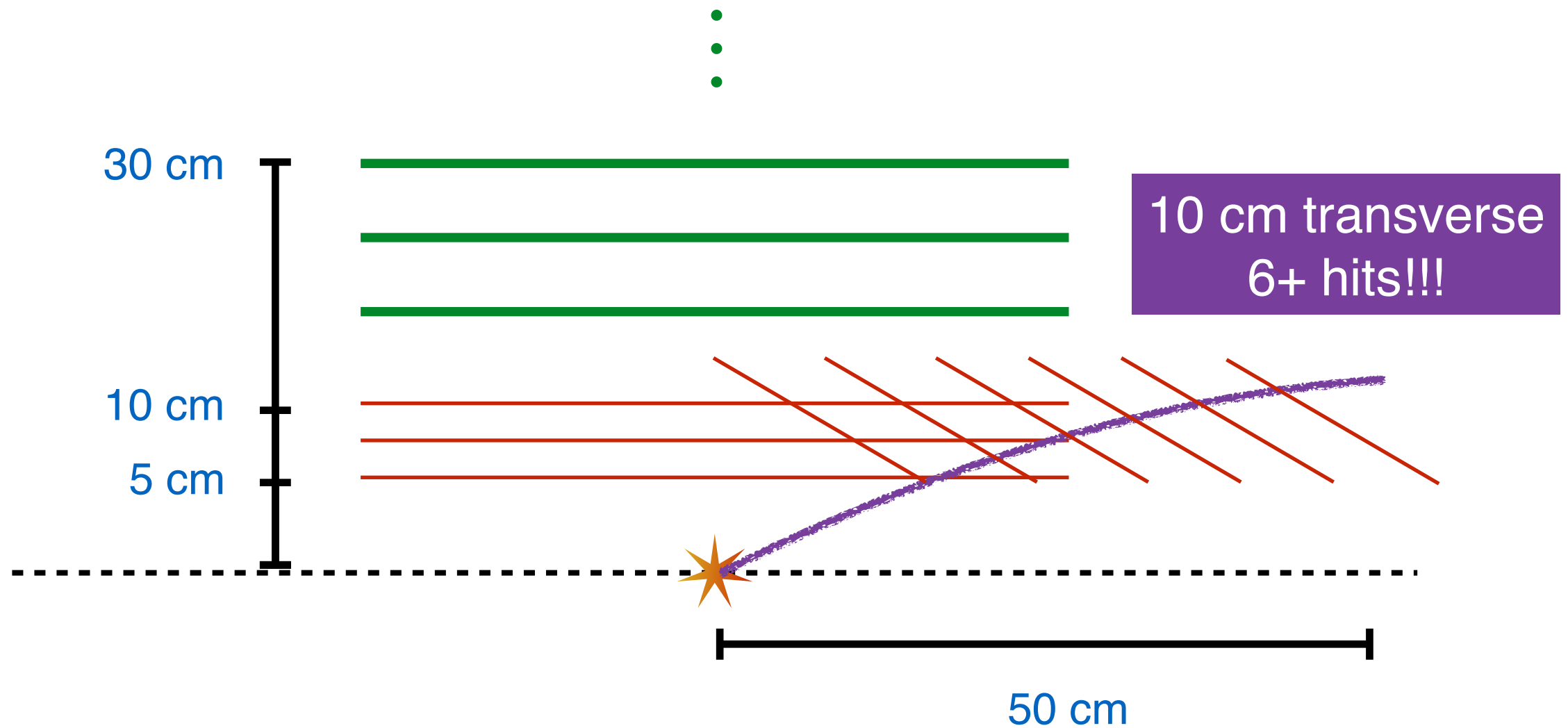


- Lifetimes ~ 1 mm, larger values (up to 10 mm) reachable for $\mu < 0$.

But $\begin{cases} 1 - \text{Lorentz boost} & d = (c\tau)\gamma\beta \\ 2 - \text{Poisson process} & P(d > d_0) = \exp(-m_{\chi^\pm} d_0 \Gamma / |\vec{p}|) \end{cases}$ \rightarrow A tale of two tails

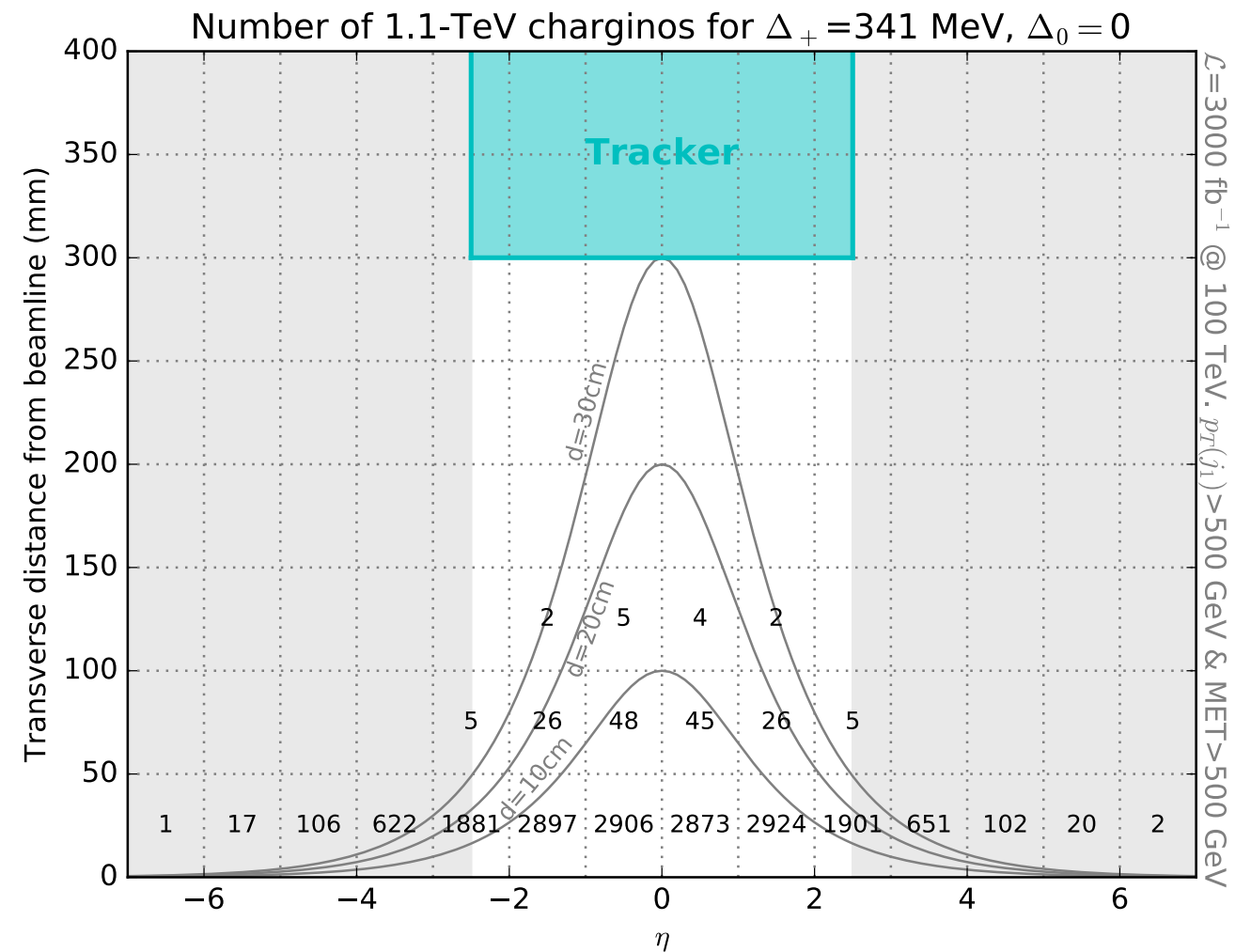
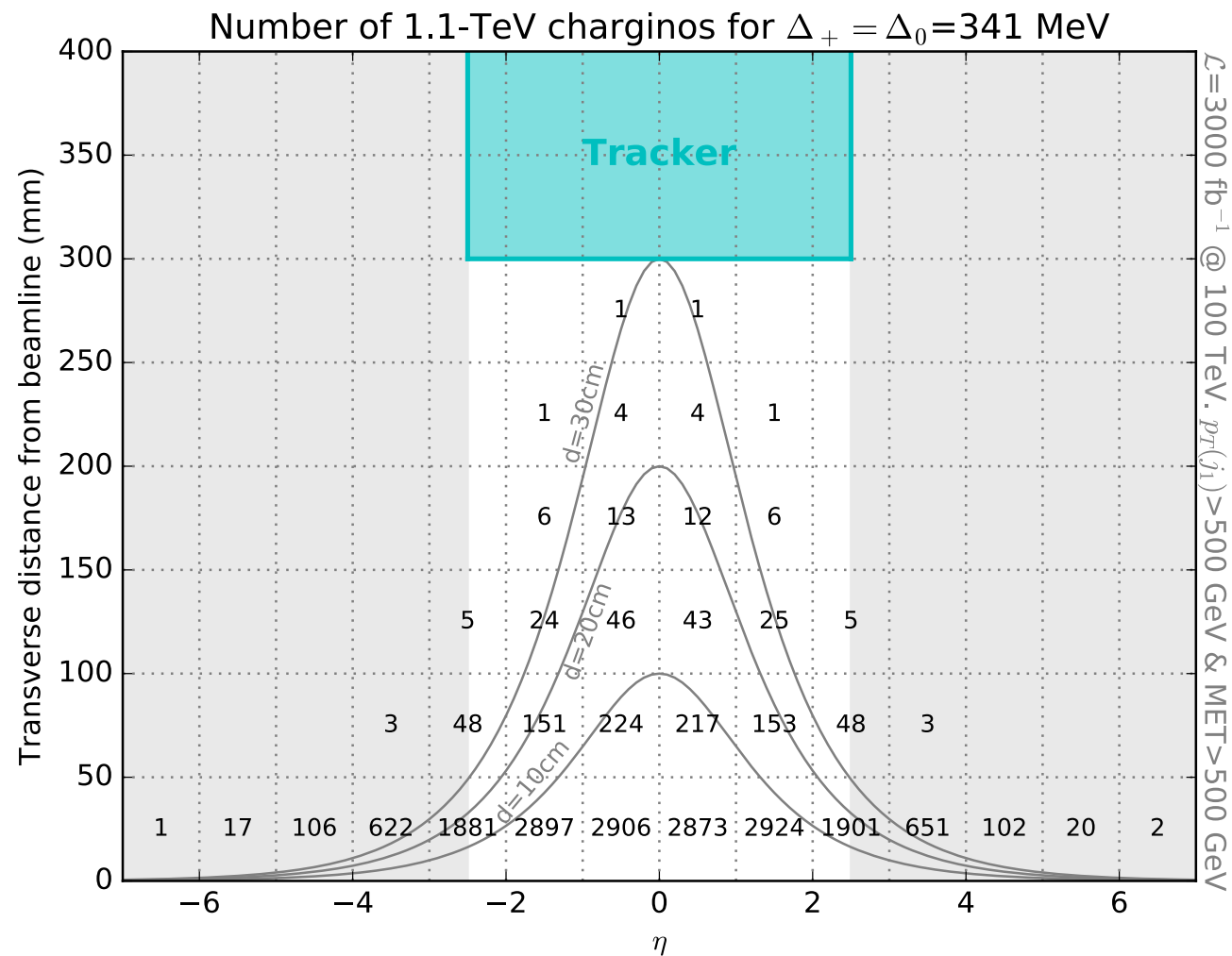
Moreover... who designed the LHC trackers??? why $d_{\min} = 30$ cm???

Detector proposal



Tracker design

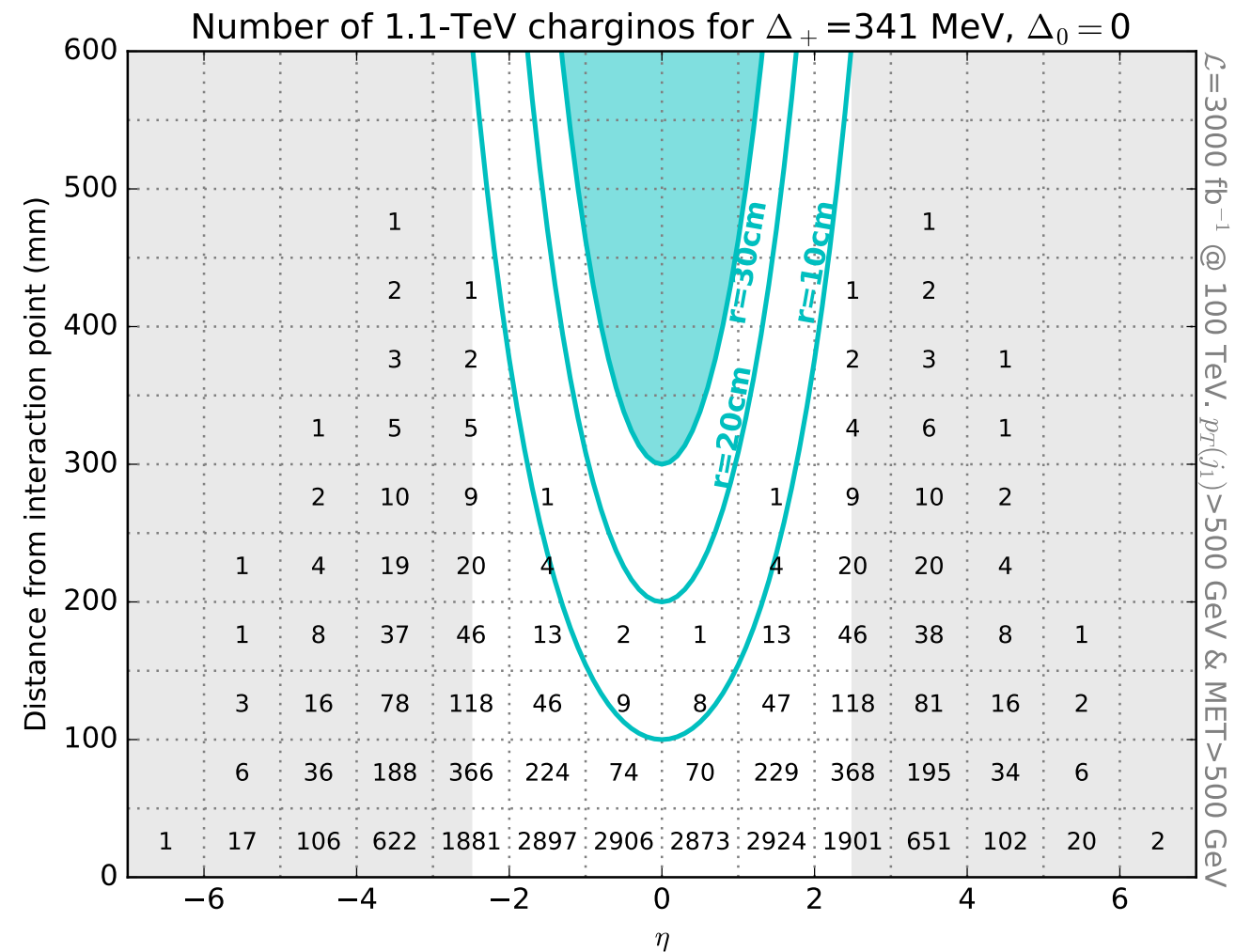
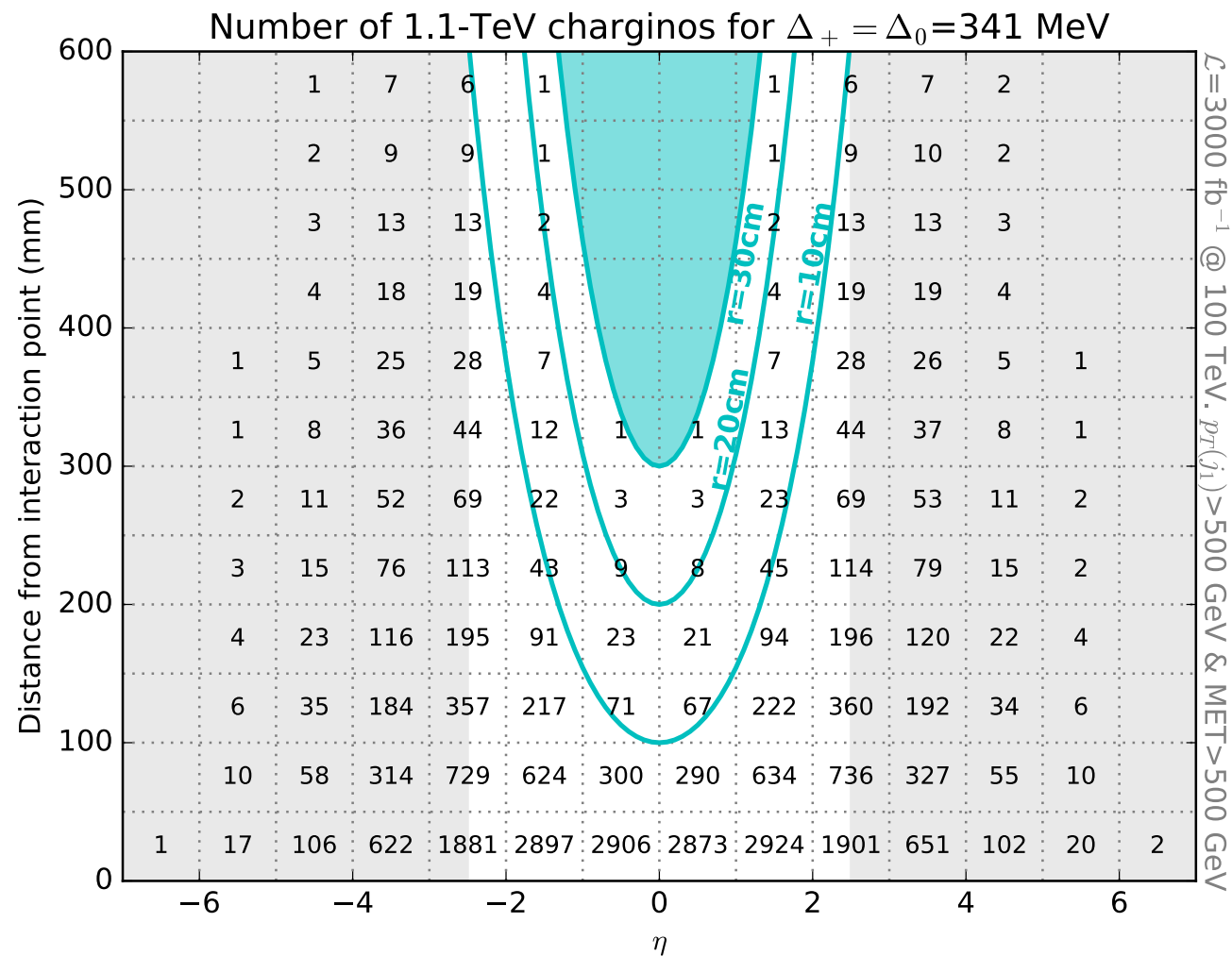
A: “Squeezing” current geometry



PRELIMINARY

Tracker design

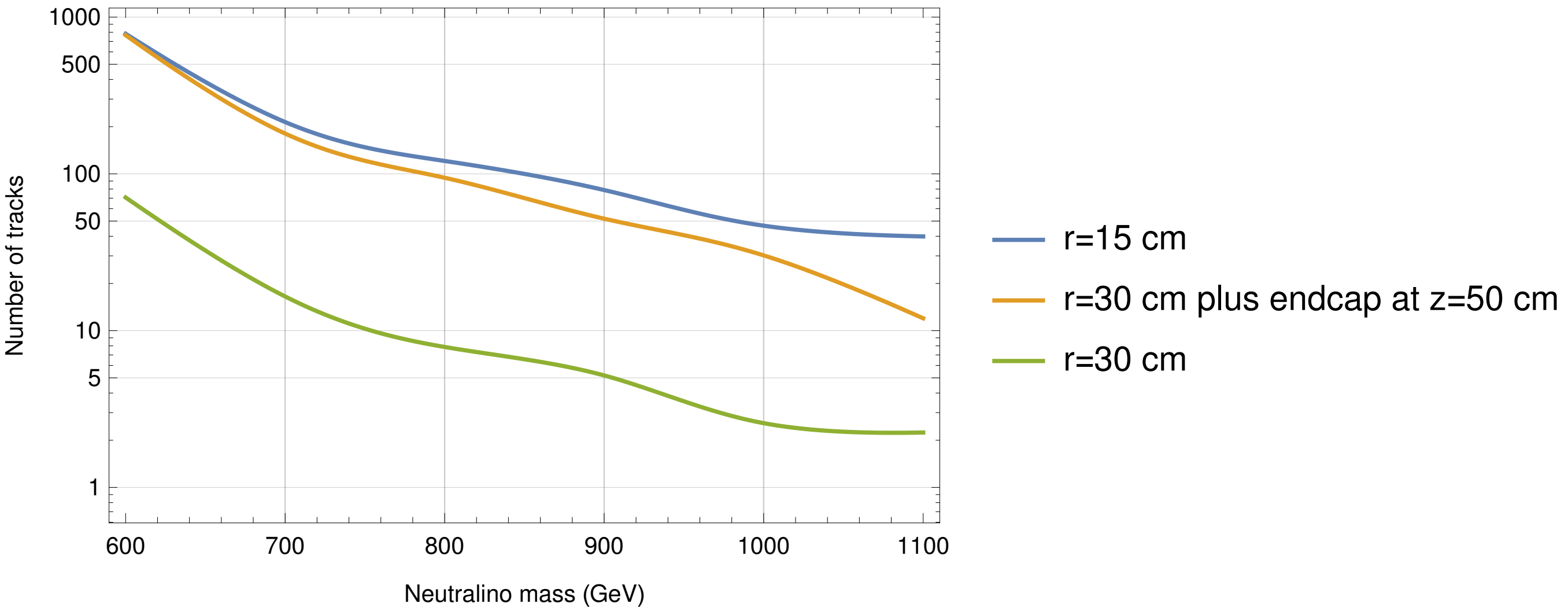
B: Spherical Tracker



PRELIMINARY

Tracker design

C: Forward (“end-cap”) tracker (very original!)



PRELIMINARY

Mono-Z analysis

@ 100 TeV FCC-hh

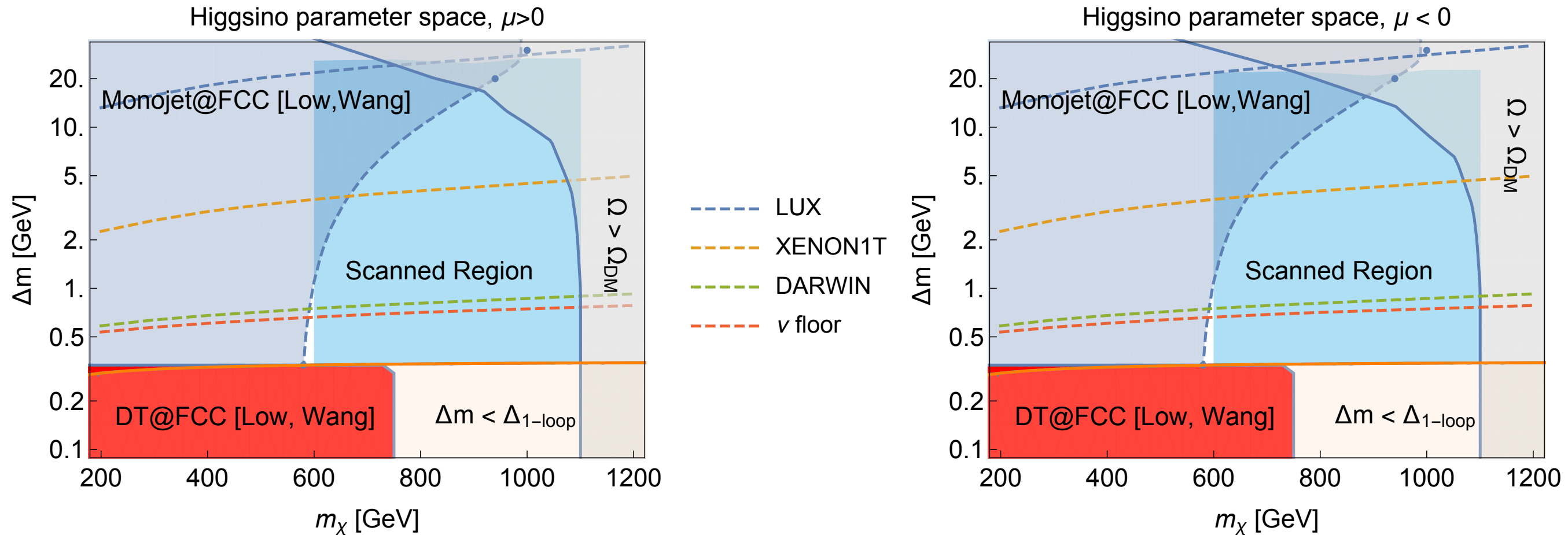
Mono-Z vs mono-jets

At the LHC, LHC, the mono-Z search for EWkinos ([Anandakrishnan, Carpenter, Raby, 1407.1833](#)) is much less sensitive than mono-jets, mono-jets plus soft-leptons ([Schwaller, JZ, 1312.7350](#); [Low, Wang, 1404.0682](#); [Barducci, Belyaev, Bharucha, Porod, Sanz 1504.02472](#) + [Badziak, Delgado, Olechowski, Pokorski, Sakurai \[1506.07177\]...](#)).

Potential advantages for mono-Z at FCC:

- Soft leptons might not be viable (depend on p_T thresholds).
- Weak coupling stronger at FCC energies.
- Weak effects in PDFs are important ([Rojo, 1605.08302](#))
- EW Sudakovs can have a large impact ([Becher, Garcia i Tormo, 1305.4202 1509.01961](#)).
- Very different systematics (crucial to estimate the sensitivity).

The parameter space



- Xenon I-T forces splittings below 2-5 GeV.
- LHC 95% C.L bounds give $m_\chi > 200$ GeV.
- FCC monojet bounds: $m_\chi > 600$ GeV for nominal splitting.
- Relic density forces $m_\chi < 1100$ GeV.
- Scanned region: $|\mu| = 600, 750, 900, 1000, 1100$; $\tau_\beta = 15$, M_1 scans Δ_+ .

Analysis pipeline

- MEs: MG4 and FR+MG5. PS: Pythia 6 (same results with Pythia 8).
Detector Simulation: Delphes with customised FCC card (loose ID, larger η)
- Backgrounds:
 - irreducible: $ZZ, WW \rightarrow l^+ l^- \nu \nu$ + fully leptonic $t\bar{t}$.
 - fake/lost leptons: W +jets, semi leptonic $t\bar{t}$ (matched up to 1 jet).
 - fake \cancel{E}_T : $Z (\rightarrow l^+ l^-) + \text{jets}$ (similarly ZW, ZZ).
- Parton level cuts: $p_T(l^+, l^-) > 400 \text{ GeV}$ or $H_T, \cancel{E}_T > 400 \text{ GeV}$.
- Event selection (basic cuts):
 - Tighter cut on $\cancel{E}_T > 500 \text{ GeV}$.
 - Two OS leptons satisfying $p_T(l) > 50 \text{ GeV}, |m(l^+, l^-) - m_Z| < 15 \text{ GeV}$.
 - Jets: Allow up to one additional hard jet ($p_T > 50 \text{ GeV}$), veto-b-jets.
 - Ignore any hard jet within $\Delta R < 0.5$ from the leptons.

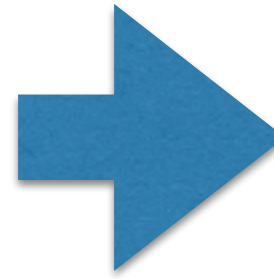
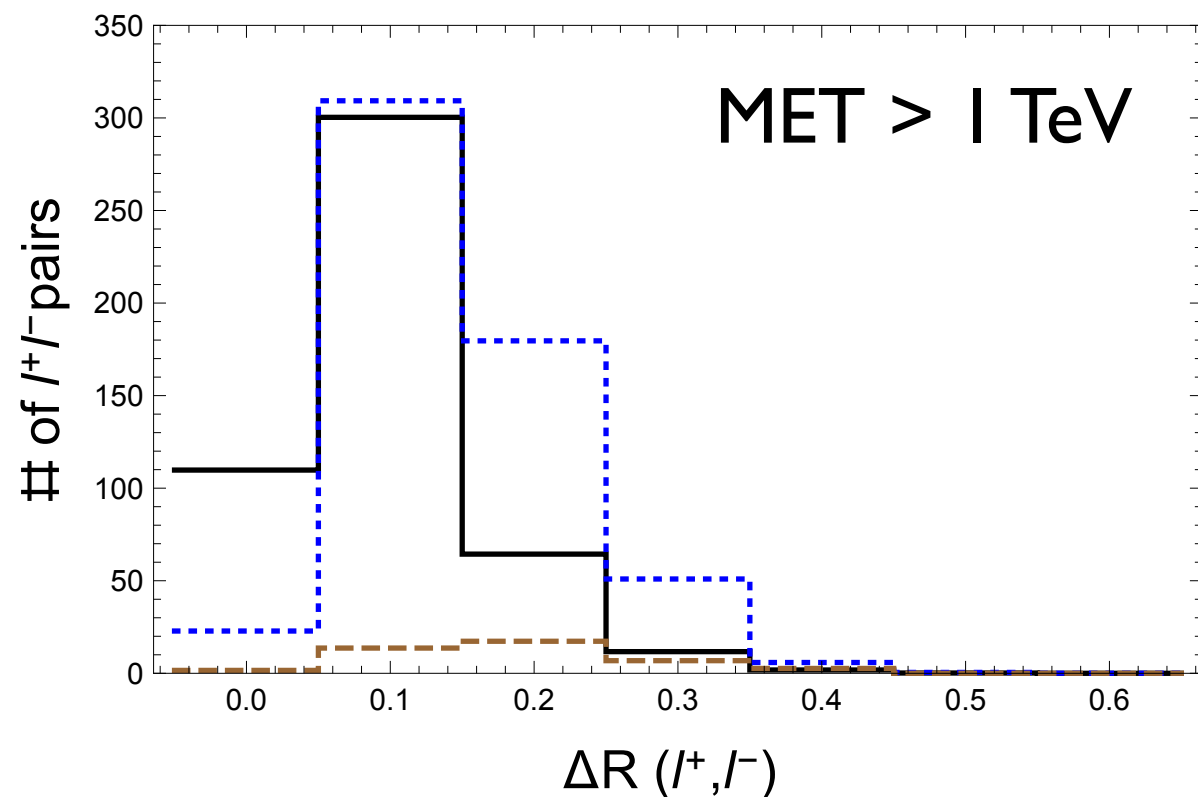
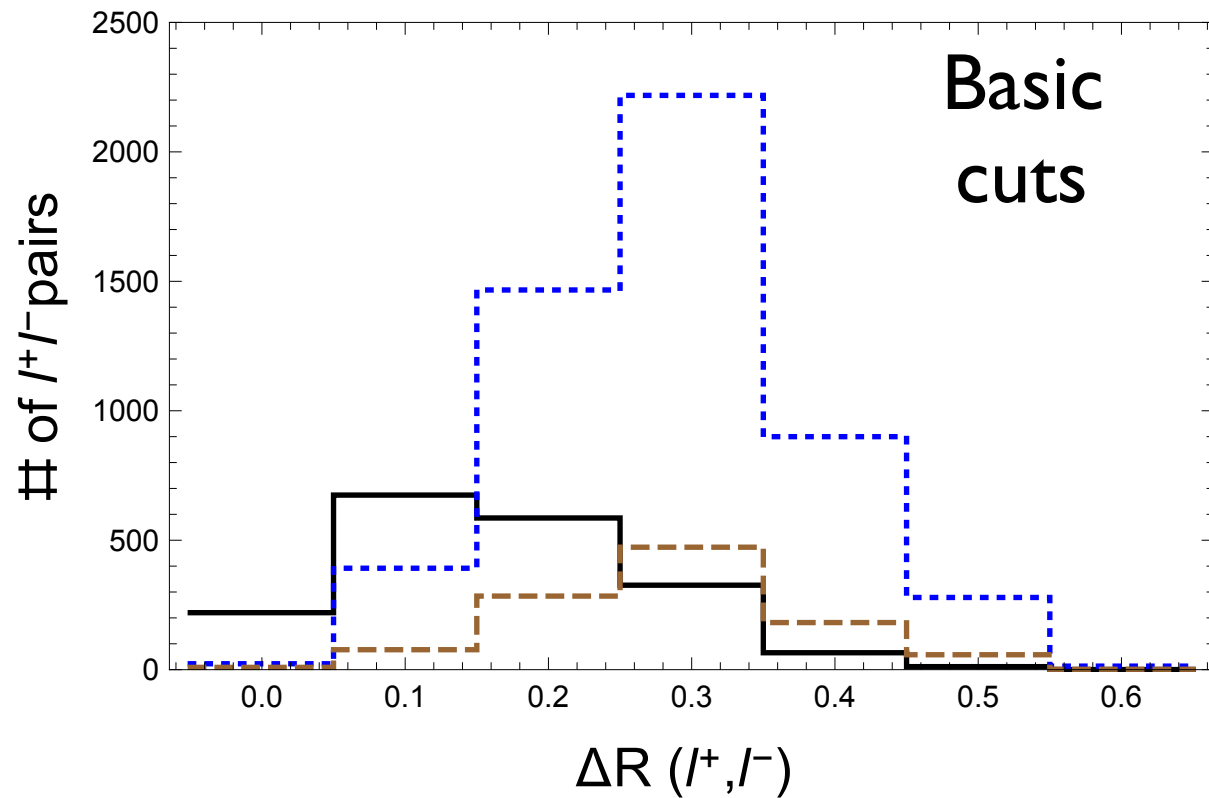
Angular separation

Benchmark J:

$$m_\chi = 1.1 \text{ TeV},$$

$$\Delta_+ = 340 \text{ MeV}$$

$$M_I = 20000 \text{ TeV}$$



This analysis heavily relies on the capability of tagging a highly boosted, leptonically decaying Z.

Optimisation and cut-flow

Optimal: $\Delta\phi(l^+l^-, \cancel{E}_T) > 0.7$, $\Delta\phi(j_1, \cancel{E}_T) > 0.1 + \cancel{E}_T > 900(+X)$ GeV.

X value chosen for the 0% systematics case

process	$\delta M \leq 15$	$N_j \leq 1$	$\Delta\phi(j, \cancel{E}_T) > 0.1$	$\Delta\phi(Z, \cancel{E}_T) > 0.7$	$\cancel{X}_T > 0.9$	$\cancel{X}_T > 1.5$
signal J	241.2	190.7	188.4	188.2	112.52	47.88
$ZZ \rightarrow l^+l^-\nu\nu$	6059.2	5346.1	5291.9	5291.9	831.6	118.3
$W^+W^- \rightarrow l^+\nu l^-\nu$	0.0134	0.0089	0.0089	0.0089	0.	0.
$tt \rightarrow l^+b\nu l^-\bar{b}\nu$	123.4	67.3	62.5	62.15	1.9	0.
$tt \rightarrow l\nu b\bar{b}jj$	255.9	95.27	94.97	8.21	1.76	0.0433
$(Z \rightarrow l^+l^-) + \text{jets}$	29342	9402.6	1370.7	1084.4	84.42	3.15
$(W \rightarrow l\nu) + \text{jets}$	336.4	115.9	115.5	10.2	0.366	0.
$ZW \rightarrow l^+l^-l\nu$	399.8	336.7	325.4	325.4	31.66	2.73
$ZZ \rightarrow l^+l^-jj$	68.50	35.86	3.36	2.47	0.0436	0.
$ZW \rightarrow l^+l^-jj$	58.12	29.09	2.92	2.2	0.	0.
100 S/B	0.658	1.23	2.59	2.77	11.8	38.5
Significance ($\beta = 0$)	1.26	1.54	2.21	2.28	3.65	4.30
Significance ($\beta = 0.1$)	0.07	0.12	0.26	0.28	1.12	2.87

Table 1: Cut flow for the backgrounds and for a signal with $\mu = 1100$ GeV, $\tan\beta = 15$ and $M_1 = \text{TeV}$. The numbers of events quoted correspond to a total integrated luminosity of 3000 fb^{-1} at a 100 TeV center-of-mass energy. We have defined $\delta M = |\frac{M_{ll} - m_Z}{\text{GeV}}|$ and $\cancel{X}_T = \frac{\cancel{E}_T}{\text{TeV}}$. The significance is computed assuming a) no systematic errors and b) 10 % systematic errors.

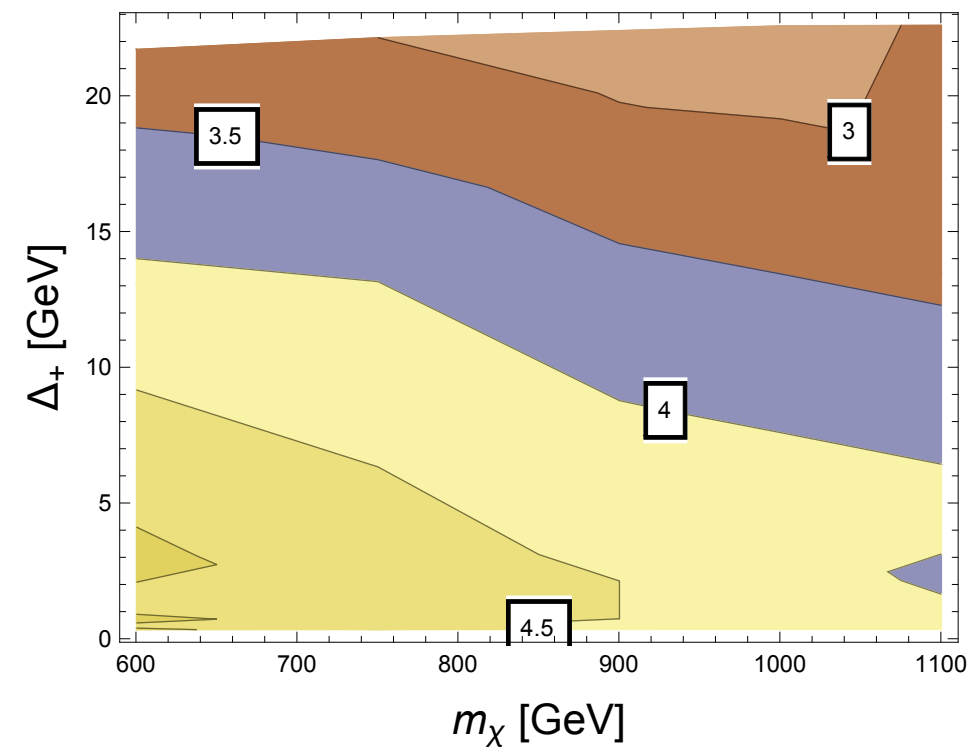
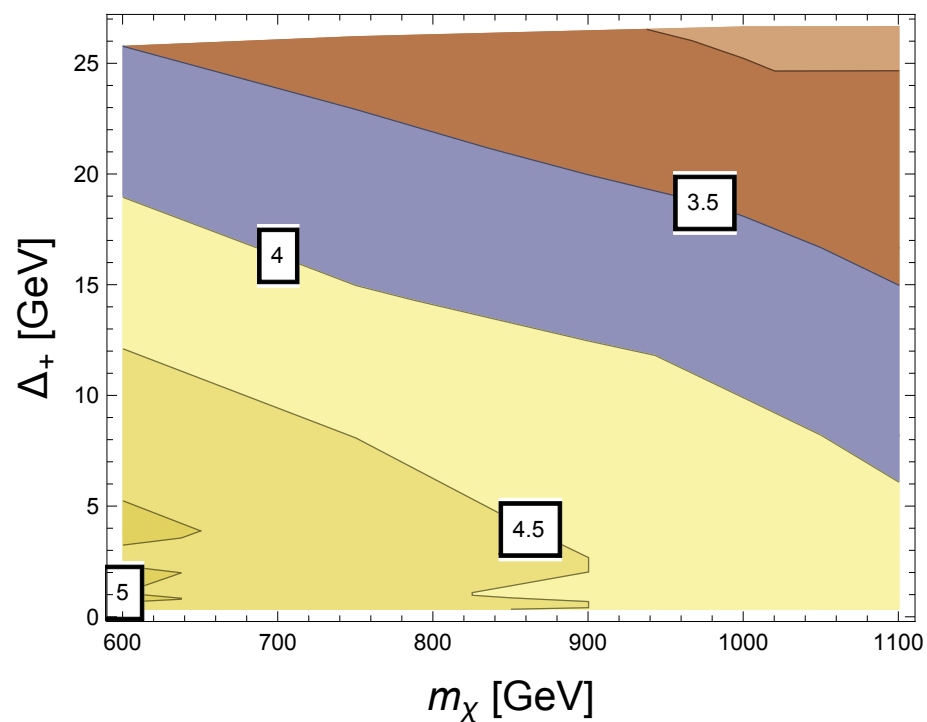
FCC reach

Lumi=3ab⁻¹

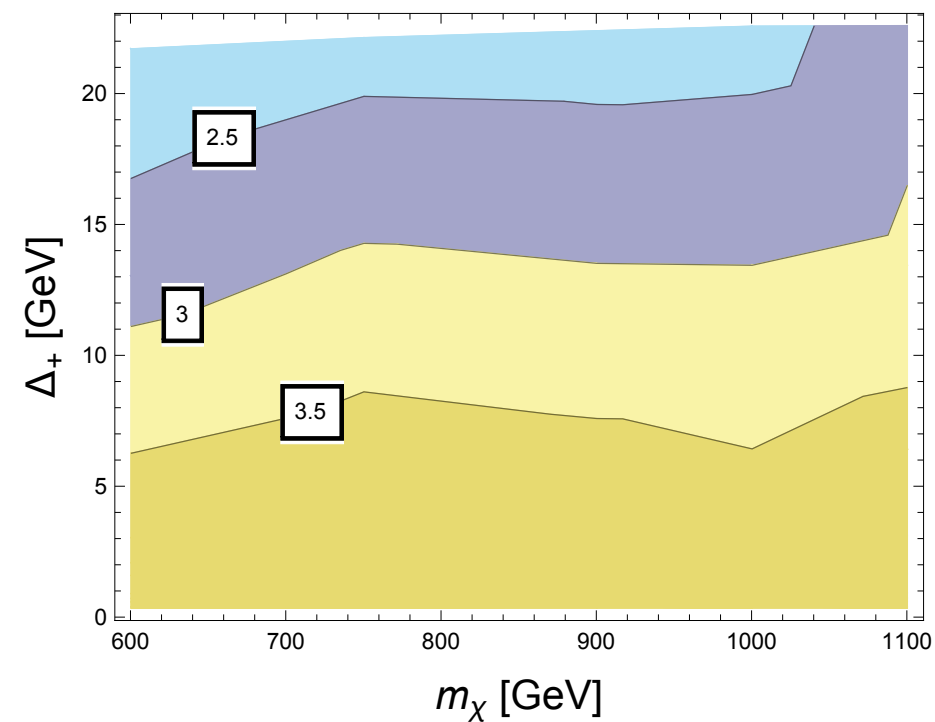
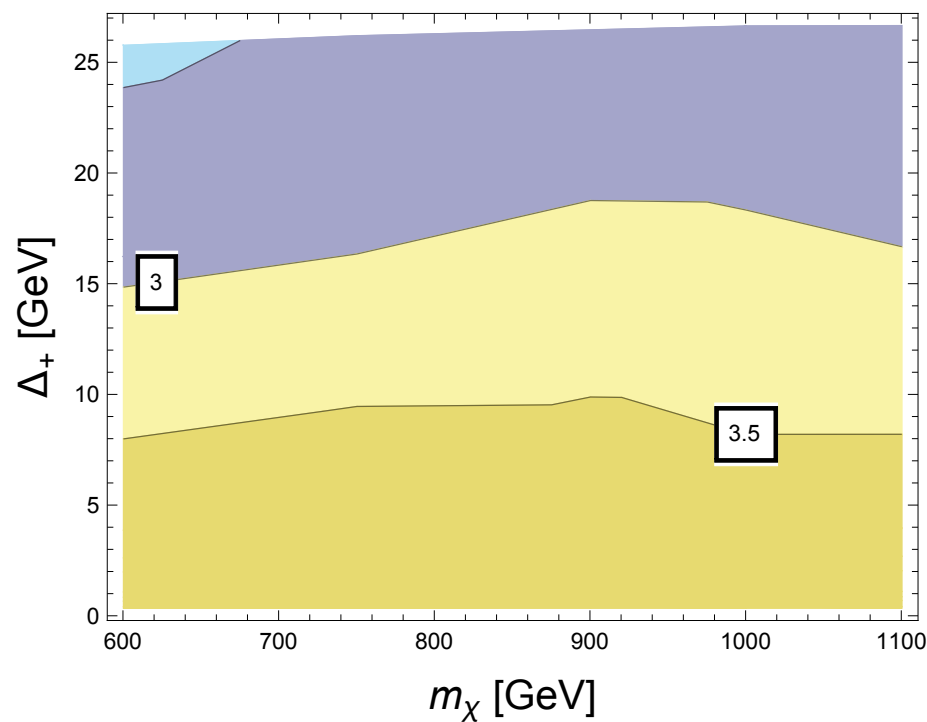
$\mu > 0$

$\mu < 0$

syst=0%

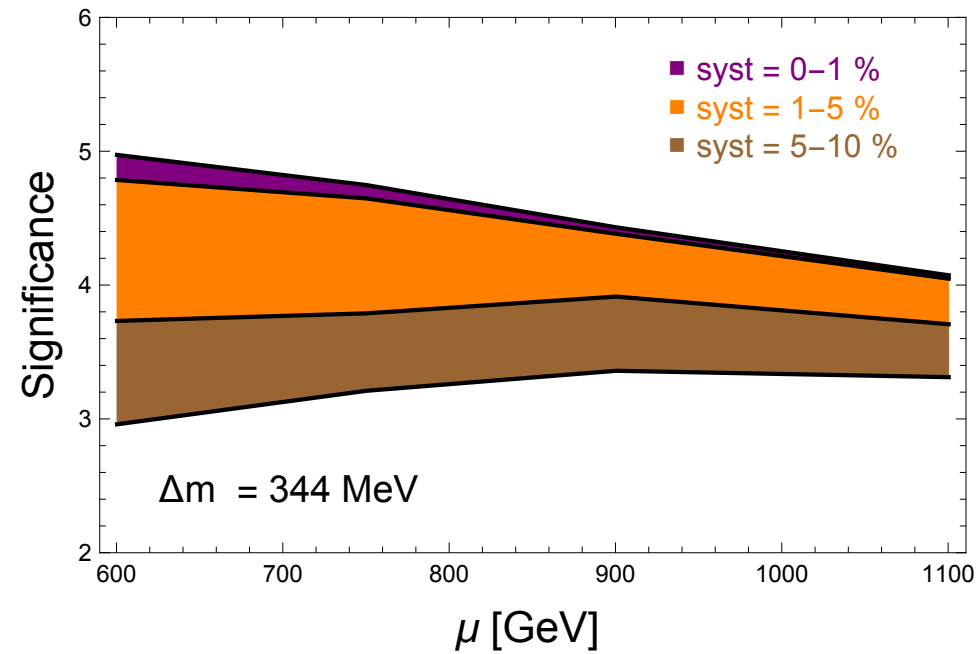


syst=5%

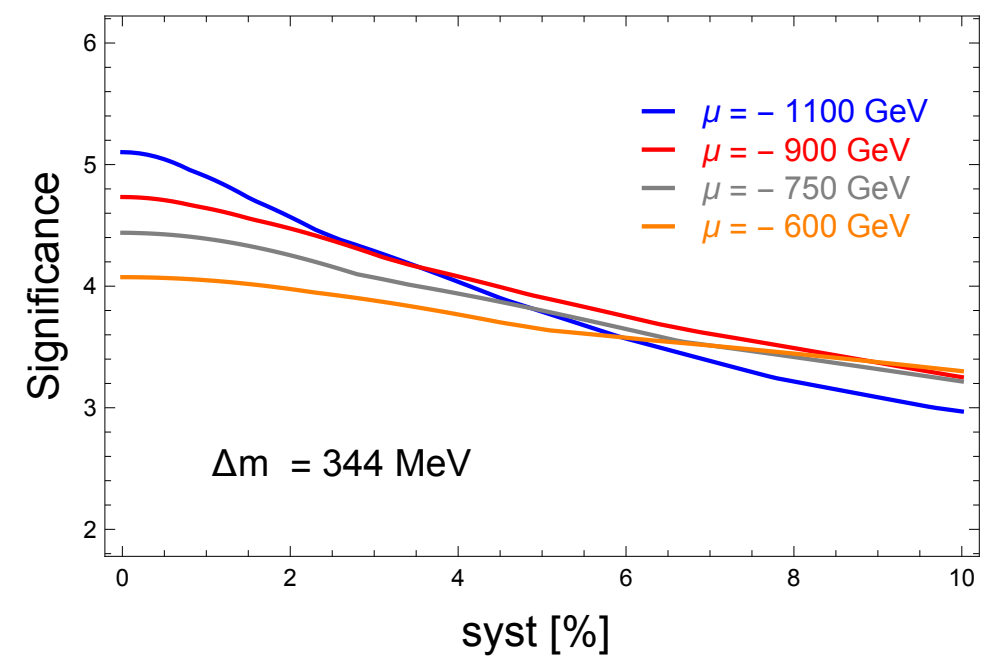
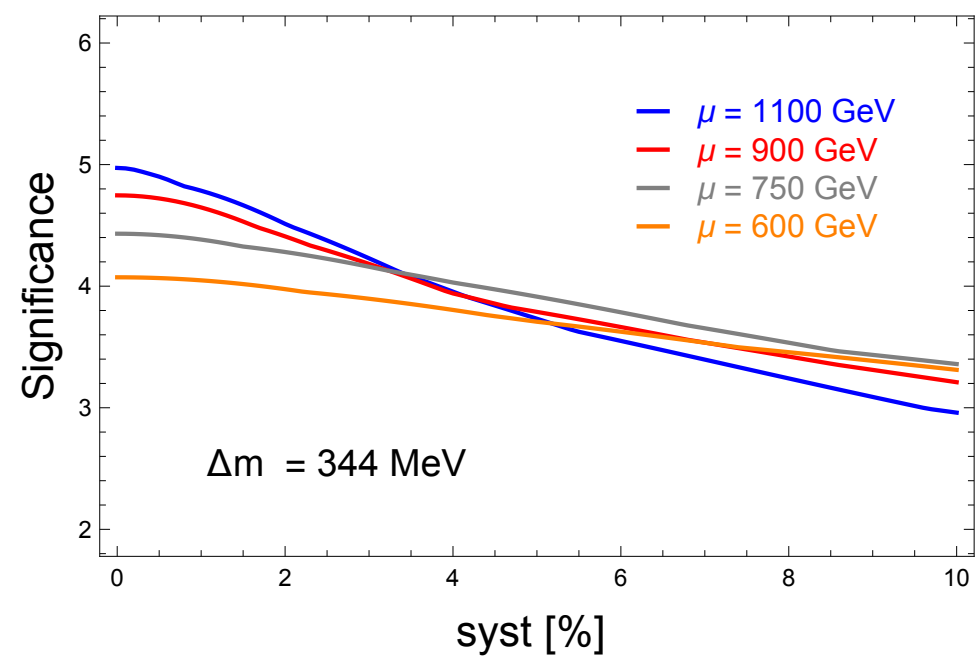
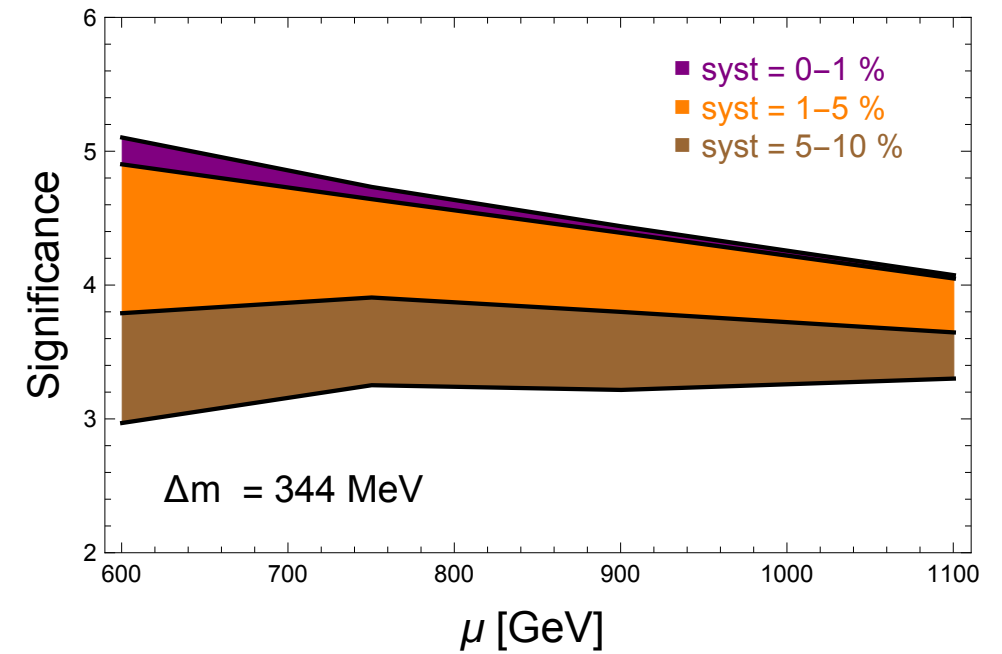


Role of systematics

$$\mu > 0$$



$$\mu < 0$$

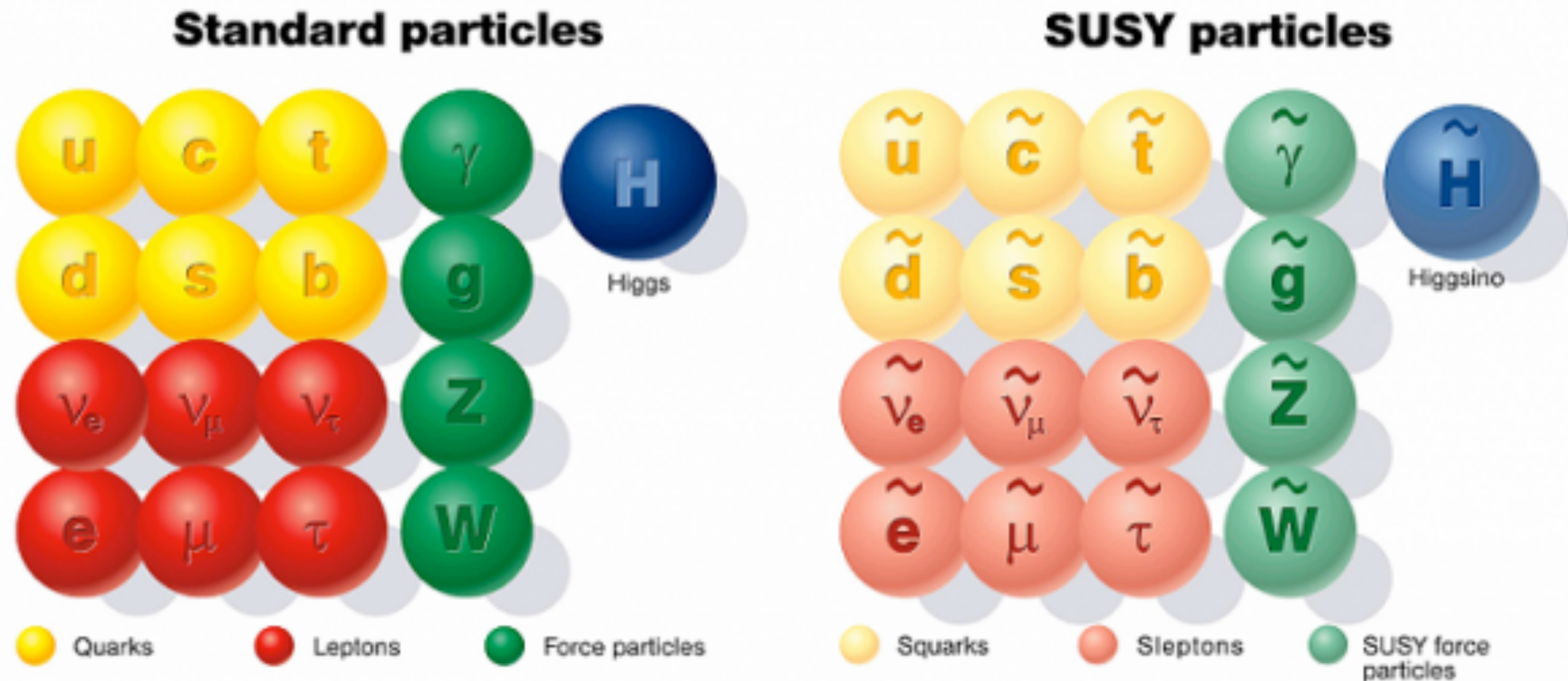


Conclusions (I)

- SUSY is tightly constrained, but still room for surprises/discoveries.
- “Leave no stones unturned” approach: full coverage of parameter space.
- Compressed spectra is “natural”: $O(\text{GeV})$ splittings among components of the same EW-multiplet. Favored by (lack of signal in) direct detection experiments.
- Discussed two strategies @ FCC: (a) disappearing tracks and (b) mono-Z.
 - (a) 100 TeV detectors yet to be designed: historical chance to write the TDR. Hard to estimate the significance, yet $O(10)$ signal events in “almost background-free” environment for 1.1 TeV Higgsinos.
 - (b) With 3000 fb^{-1} , 5% (0%) systematics one achieves 3.7 (4.3) σ for 1.1 TeV, and “full” coverage of parameter space at the 3 (4) σ level.
- Worth taking time to think what the FCC/next collider (not LHC-100!!!) can do for your favourite physics scenario. What would you need???

Outlook (II)

Supersymmetry is out there!!!



“Phenomenological studies are right or wrong based on if they can be reproduced by real experimentalists and real detectors or not.”