



Canada's national laboratory
for particle and nuclear physics
and accelerator-based science

Looking for New Physics with Pion Decays

Luca Doria
TRIUMF

Jan 2017



Outline

- General Context
- Motivation
- Experimental Technique
- Data Analysis
- Results
- Conclusions & Outlook



History and Physics Motivation

1935: H. Yukawa predicts a new particle

1936: Discovery of the Muon

1947: C. Powell and collaborators discover the Pion

M.Lattes, H.Muirhead, G.Occhialini, C.Powell:

Nature, 159:694-697 (1947)



1949: H.Yukawa awarded the Nobel Prize.

1950: C. Powell awarded the Nobel Prize



- Pion discovered with $\pi^+ \rightarrow \mu^+ \nu$
 $\quad \quad \quad \searrow$
 $\quad \quad \quad e^+ \nu \bar{\nu}$

- But: $m_e = 0.511 \text{ MeV}$

$$m_\mu = 105 \text{ MeV}$$

- Why don't we see $\pi^+ \rightarrow e^+ \nu$?

- 1950s: Many experimental indications that the weak interactions were violating parity. "V-A" structure:

$$H_w \sim \left(\frac{g^2 V_{ud}}{8m_W^2} \right) \bar{l} \gamma_\lambda (1 - \gamma_5) \bar{\nu}_l \bar{u} \gamma^\lambda \gamma_5 d$$

Note on the Decay of the π -Meson

M. RUDERMAN AND R. FINKELSTEIN

California Institute of Technology, Pasadena, California

(Received July 25, 1949)

Assuming the symmetric coupling scheme proposed by Wheeler and Tiomno, and others, we have calculated the ratio of the decay rate π -meson \rightarrow electron + neutrino to the decay rate of π -meson $\rightarrow \mu$ -meson + neutrino. The electron-neutrino decay proceeds faster, in disagreement with experiment, unless the π -meson is pseudoscalar and the β -decay coupling is pseudovector. Hence if the symmetric coupling scheme is correct and no other direct couplings are introduced, the π -meson must be pseudoscalar and β -decay must be at least partially pseudovector. If symmetric coupling is not assumed, no conclusion of this kind can be drawn.

Meson	Scalar	f	f	f	f	f
	P -scalar	f	5.1	f	1.0×10^{-4}	f
	Vector	f	f	4.0	f	2.4
	P -vector	f	f	f	f	2.4

PHYSICAL REVIEW

VOLUME 76, NUMBER 10

NOVEMBER 15, 1949

Note on the Decay of the π -Meson

M. RUDERMAN AND R. FINKELSTEIN
California Institute of Technology, Pasadena, California
 (Received July 25, 1949)

Assuming the symmetric coupling scheme proposed by Wheeler and Tiomno, and others, we have calculated the ratio of the decay rate π -meson \rightarrow electron + neutrino to the decay rate of π -meson $\rightarrow \mu$ -meson + neutrino. The electron-neutrino decay proceeds faster, in disagreement with experiment, unless the π -meson is pseudoscalar and the β -decay coupling is pseudovector. Hence if the symmetric coupling scheme is correct and no other direct couplings are introduced, the π -meson must be pseudoscalar and β -decay must be at least partially pseudovector. If symmetric coupling is not assumed, no conclusion of this kind can be drawn.

Meson	Scalar	f	f	f	f	f
	P -scalar	f	5.1	f	1.0×10^{-4}	f
	Vector	f	f	4.0	f	2.4
	P -vector	f	f	f	f	2.4

PHYSICAL REVIEW

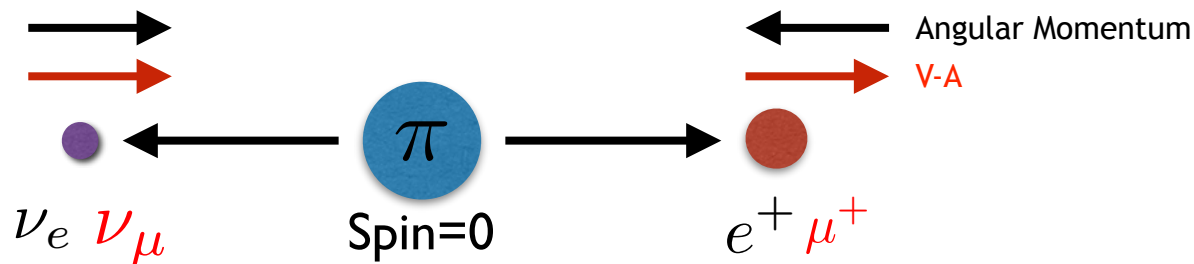
VOLUME 109, NUMBER 1

JANUARY 1, 1958

Theory of the Fermi Interaction

R. P. FEYNMAN AND M. GELL-MANN
California Institute of Technology, Pasadena, California
 (Received September 16, 1957)

Experimentally¹⁶ no $\pi \rightarrow e + \nu$ have been found, indicating that the ratio is less than 10^{-5} . This is a very serious discrepancy. The authors have no idea on how it can be resolved.

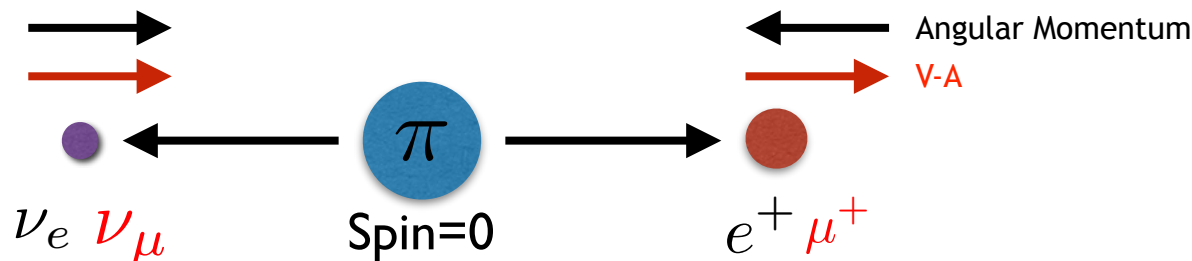


Neutrinos produced only by weak interactions:

Neutrinos: left-handed helicity

Antineutrinos: right-handed helicity

Weak interaction forces the electron into the “wrong” helicity state



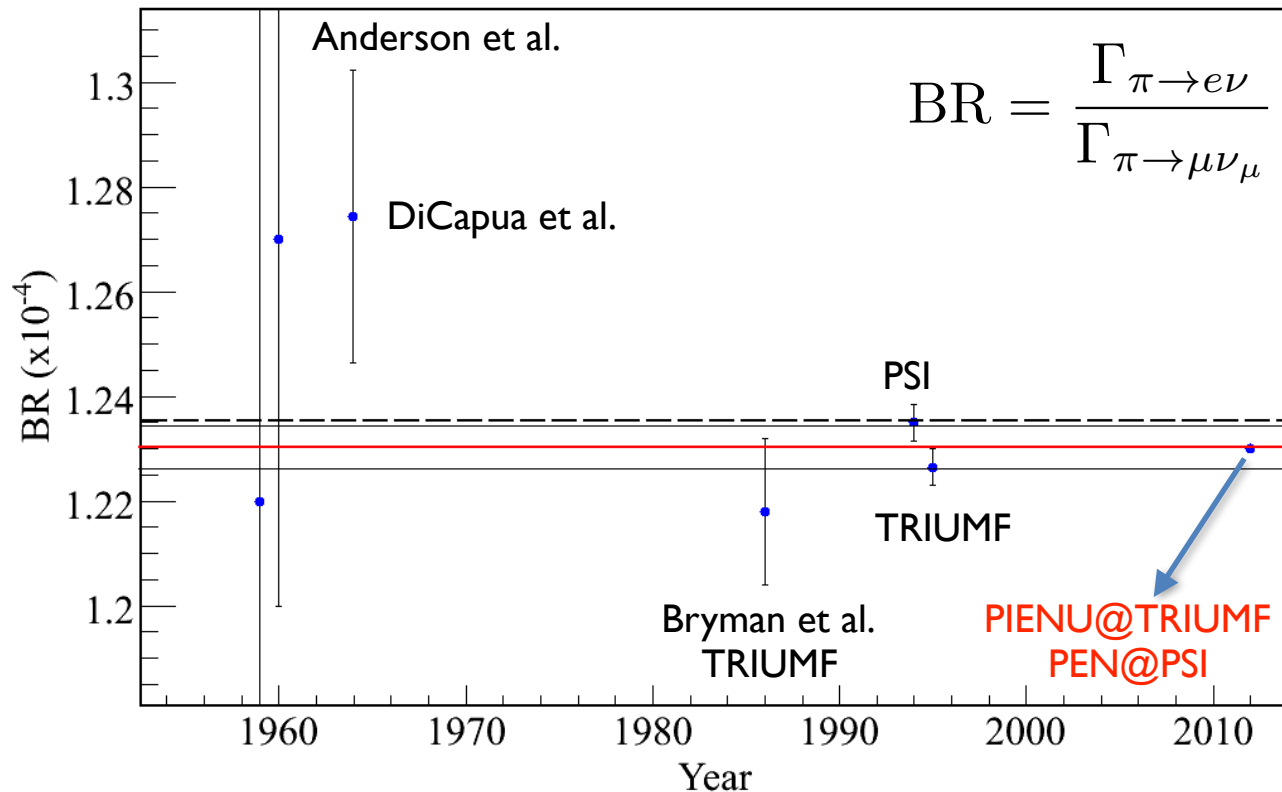
Neutrinos produced only by weak interactions:

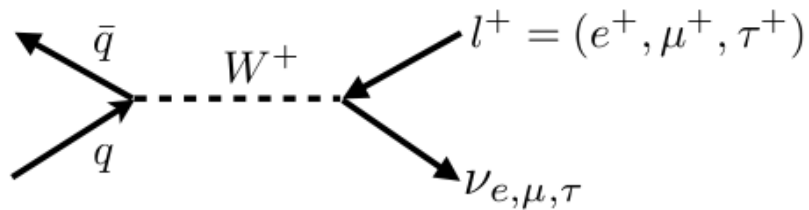
Neutrinos: left-handed helicity

Antineutrinos: right-handed helicity

Weak interaction forces the electron into the “wrong” helicity state

The V-A structure of the weak interactions explains why the muon decay mode is favoured!

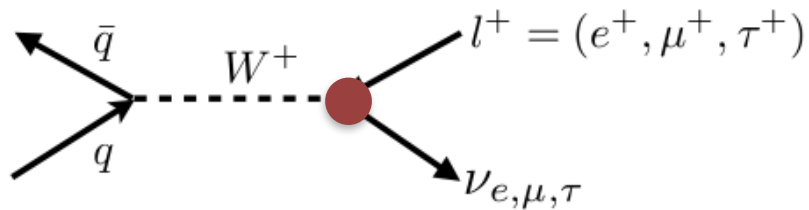




$$\Gamma_{\pi \rightarrow l \nu_l} = G_{e,\mu,\tau}^2 \frac{m_\pi f_\pi^2 m_l^2}{8\pi} \left(1 - \frac{m_l^2}{m_\pi^2}\right)^2$$

$$R_0 = \frac{\Gamma_{\pi \rightarrow e \nu_e}}{\Gamma_{\pi \rightarrow \mu \nu_\mu}} = \frac{m_e^2}{m_\mu^2} \left(\frac{m_\pi^2 - m_e^2}{m_\pi^2 - m_\mu^2}\right)^2 = 1.28336(2) \times 10^{-4}$$

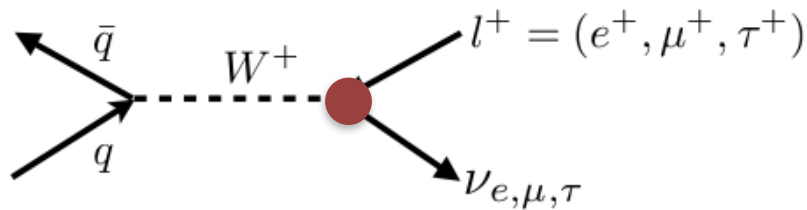
Helicity Suppression



$$\Gamma_{\pi \rightarrow l \nu_l} = G_{e,\mu,\tau}^2 \frac{m_\pi f_\pi^2 m_l^2}{8\pi} \left(1 - \frac{m_l^2}{m_\pi^2}\right)^2$$

$$R_0 = \frac{\Gamma_{\pi \rightarrow e \nu_e}}{\Gamma_{\pi \rightarrow \mu \nu_\mu}} = \frac{m_e^2}{m_\mu^2} \left(\frac{m_\pi^2 - m_e^2}{m_\pi^2 - m_\mu^2}\right)^2 = 1.28336(2) \times 10^{-4}$$

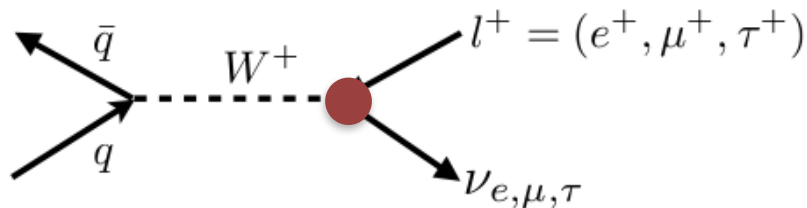
Helicity Suppression



$$\Gamma_{\pi \rightarrow l \nu_l} = \underbrace{G_{e,\mu,\tau}^2}_{\text{Helicity Suppression}} \frac{m_\pi f_\pi^2 m_l^2}{8\pi} \left(1 - \frac{m_l^2}{m_\pi^2}\right)^2$$

$$R_0 = \frac{\Gamma_{\pi \rightarrow e \nu_e}}{\Gamma_{\pi \rightarrow \mu \nu_\mu}} = \frac{m_e^2}{m_\mu^2} \left(\frac{m_\pi^2 - m_e^2}{m_\pi^2 - m_\mu^2}\right)^2 = 1.28336(2) \times 10^{-4}$$

Helicity Suppression



$$\Gamma_{\pi \rightarrow l \nu_l} = \frac{G_{e, \mu, \tau}^2 m_\pi f_\pi^2 m_l^2}{8\pi} \left(1 - \frac{m_l^2}{m_\pi^2}\right)^2$$

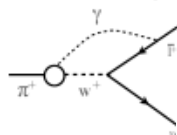
$$R_0 = \frac{\Gamma_{\pi \rightarrow e \nu_e}}{\Gamma_{\pi \rightarrow \mu \nu_\mu}} = \frac{m_e^2}{m_\mu^2} \left(\frac{m_\pi^2 - m_e^2}{m_\pi^2 - m_\mu^2}\right)^2 = 1.28336(2) \times 10^{-4}$$

Helicity Suppression

$$R = R_0 \times \left[1 + \frac{\alpha}{\pi} \left\{ F\left(\frac{m_e}{m_\pi}\right) - F\left(\frac{m_\mu}{m_\pi}\right) + \frac{m_\mu^2}{m_\rho^2} \left(c_2 \ln \frac{m_\rho^2}{m_\mu^2} + c_3 \right) + c_4 \frac{m_\pi^6}{m_e^2 m_\rho^4} \right\} + c_5 \left(\frac{\alpha}{\pi} \ln \frac{m_\mu}{m_e} \right)^2 + \dots \right] (+SD_\pi)$$

V. Cirigliano, I. Rosell: Phys. Rev. Lett. 99(23), 231801 (2007)

M. Terent'ev: Yad. Fiz. 18(870) (1973)

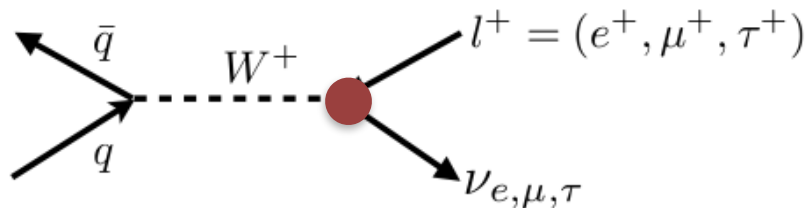


S. Berman: Phys. Rev. Lett. 1(12), 468 (1958)

T. Kinoshita: Phys. Rev. Lett. 2(11), 477 (1959)

T. Goldman, W. Wilson: Phys. Rev. D 14(9), 2428 (1976)

W. Marciano, A. Sirlin: Phys. Rev. Lett. 36(24), 1425 (1976)



$$\Gamma_{\pi \rightarrow l \nu_l} = \frac{G_{e,\mu,\tau}^2 m_\pi f_\pi^2 m_l^2}{8\pi} \left(1 - \frac{m_l^2}{m_\pi^2}\right)^2$$

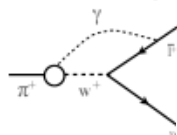
$$R_0 = \frac{\Gamma_{\pi \rightarrow e \nu_e}}{\Gamma_{\pi \rightarrow \mu \nu_\mu}} = \frac{m_e^2}{m_\mu^2} \left(\frac{m_\pi^2 - m_e^2}{m_\pi^2 - m_\mu^2}\right)^2 = 1.28336(2) \times 10^{-4}$$

Helicity Suppression

$$R = R_0 \times \left[1 + \frac{\alpha}{\pi} \left\{ F\left(\frac{m_e}{m_\pi}\right) - F\left(\frac{m_\mu}{m_\pi}\right) + \frac{m_\mu^2}{m_\rho^2} \left(c_2 \ln \frac{m_\rho^2}{m_\mu^2} + c_3 \right) + c_4 \frac{m_\pi^6}{m_e^2 m_\rho^4} \right\} + c_5 \left(\frac{\alpha}{\pi} \ln \frac{m_\mu}{m_e} \right)^2 + \dots \right] (+SD_\pi)$$

V. Cirigliano, I. Rosell: Phys. Rev. Lett. 99(23), 231801 (2007)

M. Terent'ev: Yad. Fiz. 18(870) (1973)



- S. Berman: Phys. Rev. Lett. 1(12), 468 (1958)
- T. Kinoshita: Phys. Rev. Lett. 2(11), 477 (1959)
- T. Goldman, W. Wilson: Phys. Rev. D 14(9), 2428 (1976)
- W. Marciano, A. Sirlin: Phys. Rev. Lett. 36(24), 1425 (1976)

Standard Model: $R^{SM} = 1.2353(1) \times 10^{-4}$

Experiments: $R^{exp} = 1.230 \pm 0.004 \times 10^{-4}$

TRIUMF: D. Britton et al. Phys. Rev. Lett. 68:3000-3003 (1992)

PSI: G. Czapke et al. Phys. Rev. Lett. 70:17-20 (1993)

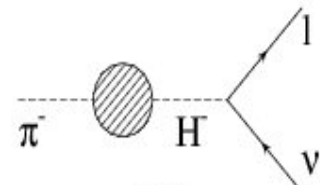
New pseudo-scalar interactions (no helicity suppression) B.Campbell, D. Maybury: Nucl.Phys. B709, 419 (2005)

$$1 - \frac{R^{exp}}{R^{SM}} \sim \mp \frac{\sqrt{2}\pi}{G_\mu} \frac{1}{\Lambda^2} \frac{m_\pi^2}{m_e(m_d + m_u)} \sim \left(\frac{1\text{TeV}}{\Lambda}\right)^2 \times 10^3 \Rightarrow 1000\text{TeV}$$

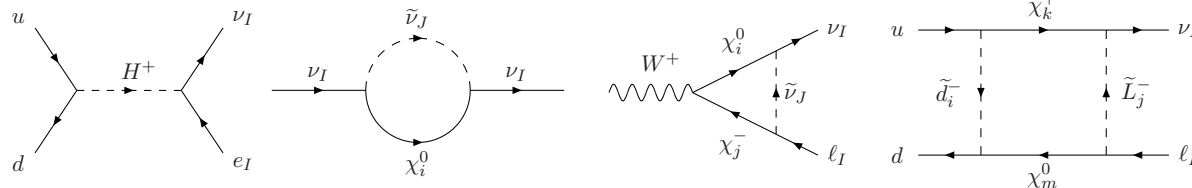
Charged Higgs (with non-SM couplings) O. Shanker: Nucl.Phys. B204(3), 375 (1982)

Relevant for SUSY models, Changes the BR through loop contributions

$$1 - \frac{R^{exp}}{R^{SM}} \sim \mp \frac{2m_\pi^2}{m_e(m_d + m_u)} \frac{m_W^2}{m_{H^\pm}} \lambda_{ud} (\lambda_{e\nu} - \frac{m_e}{m_\mu} \lambda_{\mu\nu}) \quad M_{H^\pm}^\pm \sim 400\text{GeV}$$



SUperSYmmetry M. Ramsey-Musolf, S.Su, S.Tulin: Phys.Rev. D76, 095017 (2007)



R-parity violating SUSY affects the BR already at tree-level!

And More: Leptoquarks, new scalar interactions, massive neutrinos,...

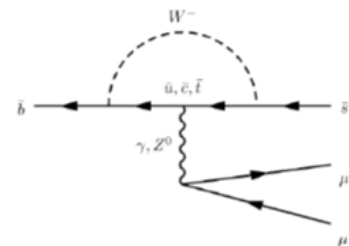
LHCb Collaboration,

R. Aaij et al., “Test of lepton universality using $B^+ \rightarrow K^+ l^+ l^-$ decays,” arXiv:1406.6482.

- Previously measured by Belle and BaBar at 20-50% precision level
- $R_K = 1$ expected from SM
- Theoretically clean observable with small corrections
- Analysis : $1 < q^2 < 6 \text{ GeV}^2/c^4$

$$R_K = \frac{B^+ \rightarrow K^+ \mu^+ \mu^-}{B^+ \rightarrow K^+ e^+ e^-} = 0.745_{-0.074}^{+0.090}(\text{stat}) \pm 0.036(\text{cyst})$$

2.6 σ deviation from the SM value



Feruglio, Paradisi, Patteri, Phys. Rev. Lett. 118, 011801 (2017)

$$R_{D^{(*)}}^{\tau/\ell} = \frac{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau \bar{\nu})_{\text{exp}} / \mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau \bar{\nu})_{\text{SM}}}{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell \bar{\nu})_{\text{exp}} / \mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell \bar{\nu})_{\text{SM}}},$$

$$R_D^{\tau/\ell} = 1.37 \pm 0.17, \quad R_{D^*}^{\tau/\ell} = 1.28 \pm 0.08.$$

3.9 σ deviation from the SM value

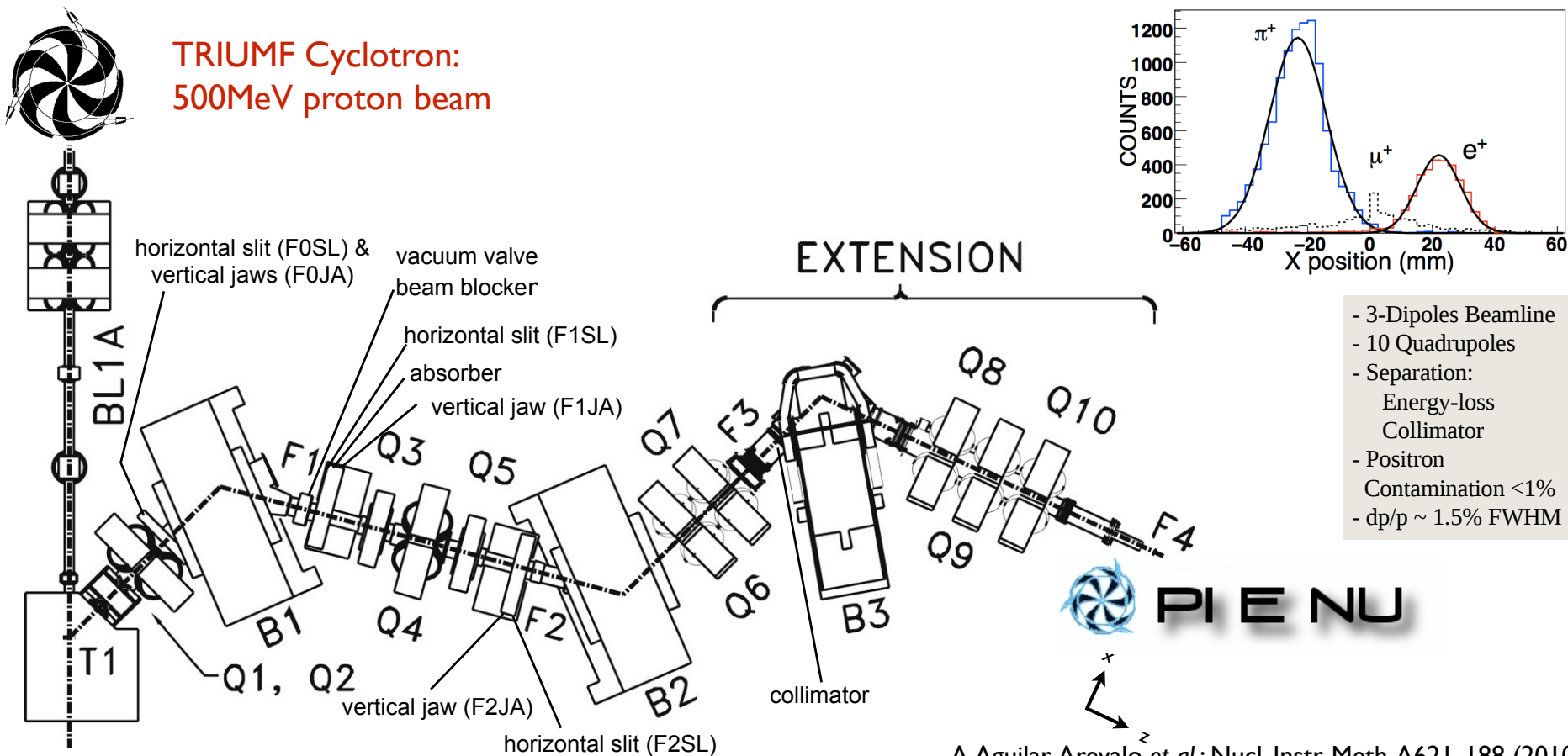
J. P. Lees *et al.* [BaBar Collaboration], Phys. Rev. D **88** (2013) 7, 072012 [[arXiv:1303.0571](https://arxiv.org/abs/1303.0571)].

M. Huschle *et al.* [Belle Collaboration], Phys. Rev. D **92** (2015) 7, 072014 [[arXiv:1507.03233](https://arxiv.org/abs/1507.03233)].

R. Aaij *et al.* [LHCb Collaboration], Phys. Rev. Lett. **115** (2015) 15, 159901 [[arXiv:1506.08614](https://arxiv.org/abs/1506.08614)].

The PIENU Experiment at TRIUMF

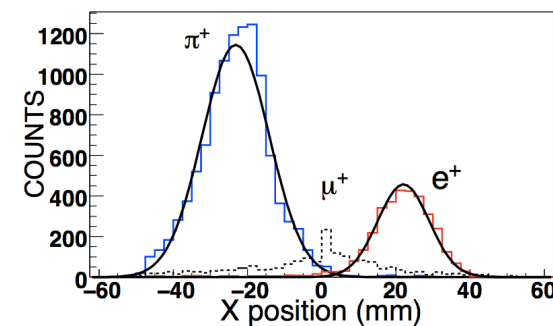
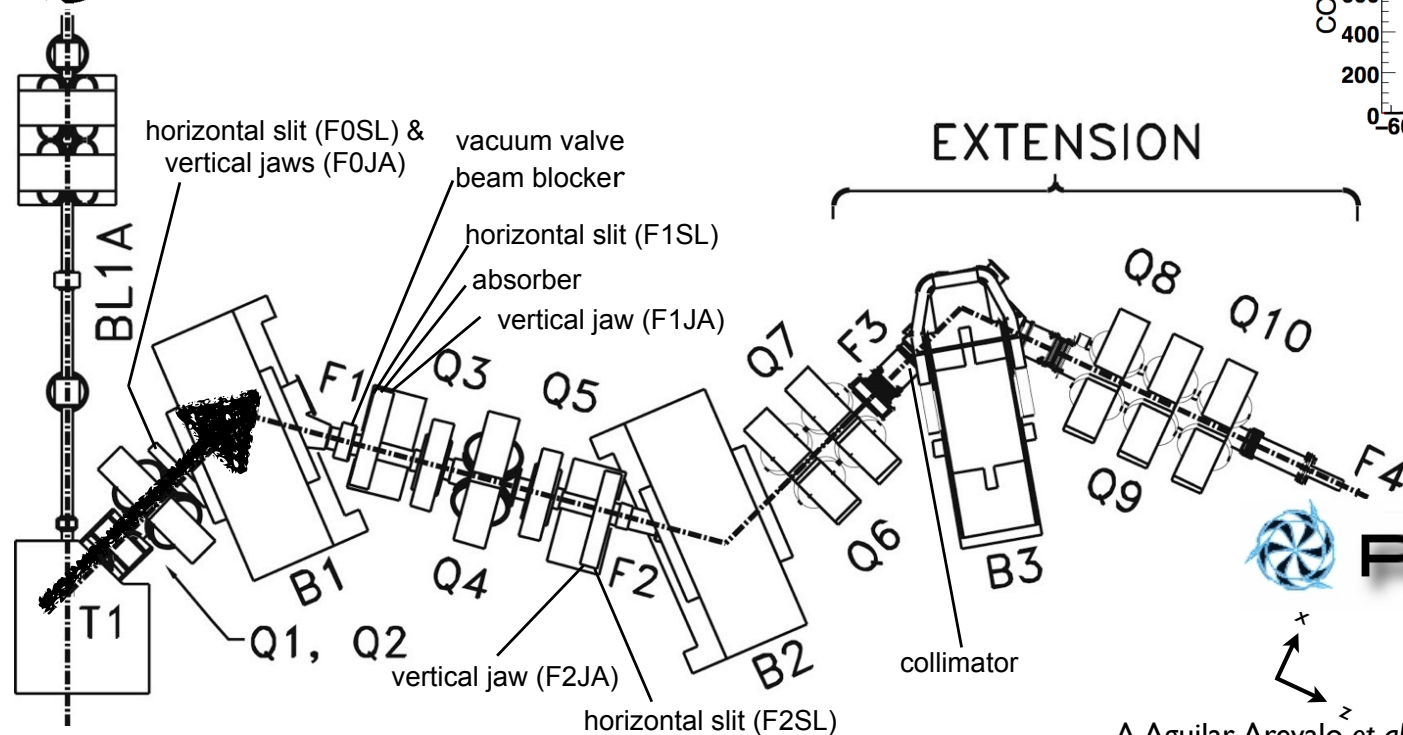
TRIUMF Cyclotron:
500MeV proton beam



A.Aguilar-Arevalo et al.: Nucl. Instr. Meth. A621, 188 (2010)



TRIUMF Cyclotron:
500MeV proton beam

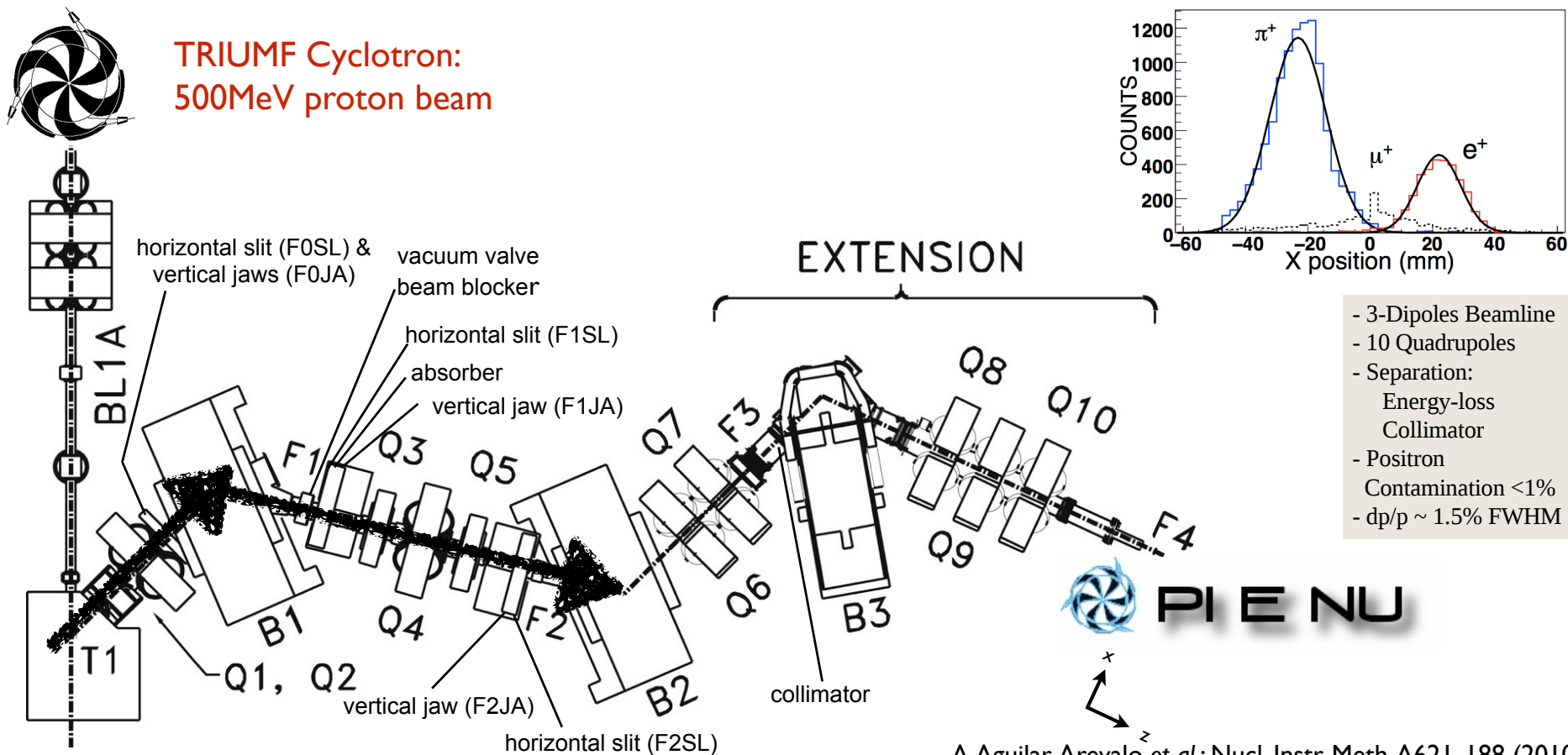


- 3-Dipoles Beamline
- 10 Quadrupoles
- Separation:
Energy-loss
Collimator
- Positron
Contamination <1%
- $dp/p \sim 1.5\%$ FWHM



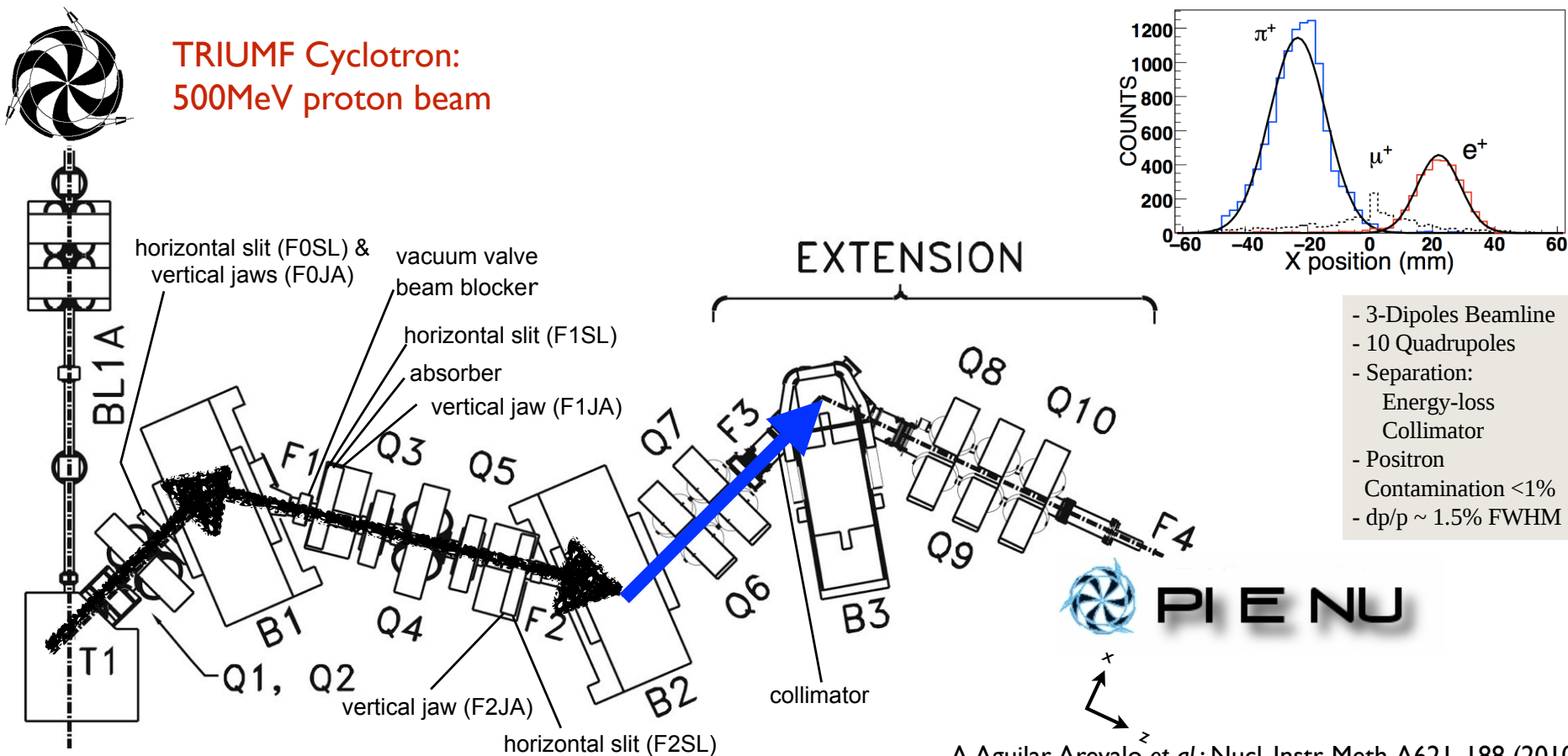
PIE NU

A.Aguilar-Arevalo et al.: Nucl. Instr. Meth. A621, 188 (2010)



A.Aguilar-Arevalo et al.: Nucl. Instr. Meth. A621, 188 (2010)

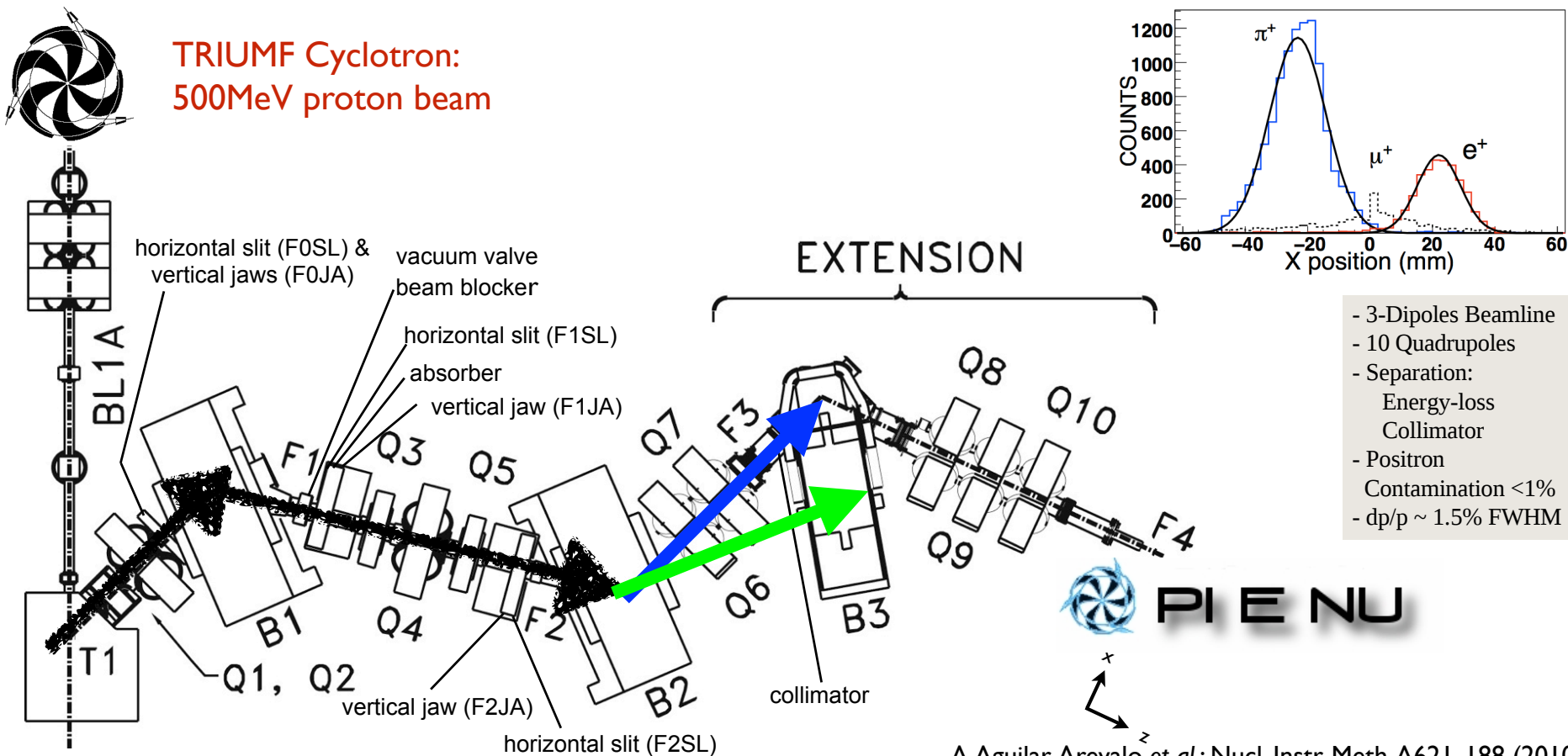
TRIUMF Cyclotron:
500MeV proton beam



- 3-Dipoles Beamline
- 10 Quadrupoles
- Separation:
 - Energy-loss
 - Collimator
- Positron
- Contamination <1%
- $dp/p \sim 1.5\%$ FWHM

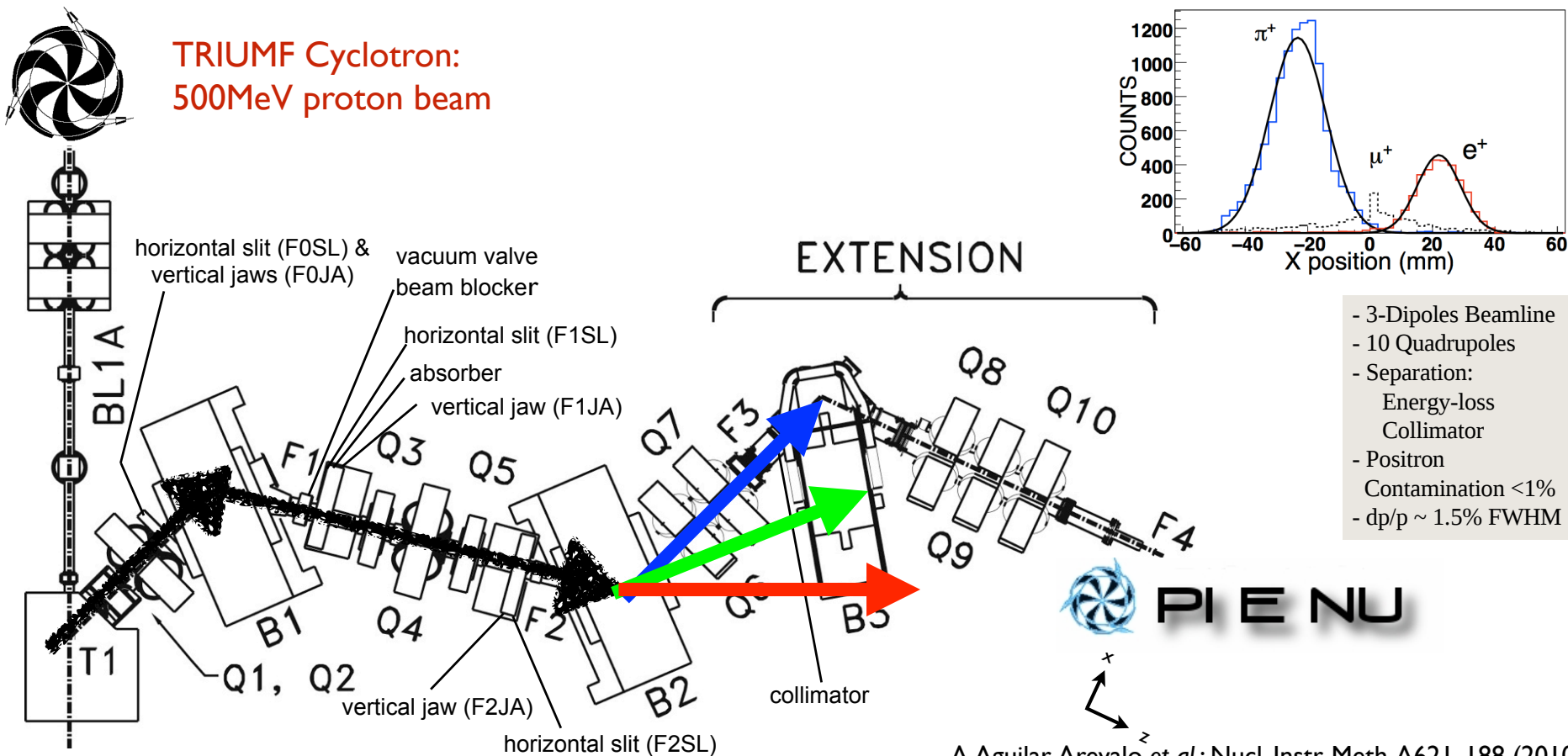
A.Aguilar-Arevalo et al.: Nucl. Instr. Meth. A621, 188 (2010)

TRIUMF Cyclotron:
500MeV proton beam



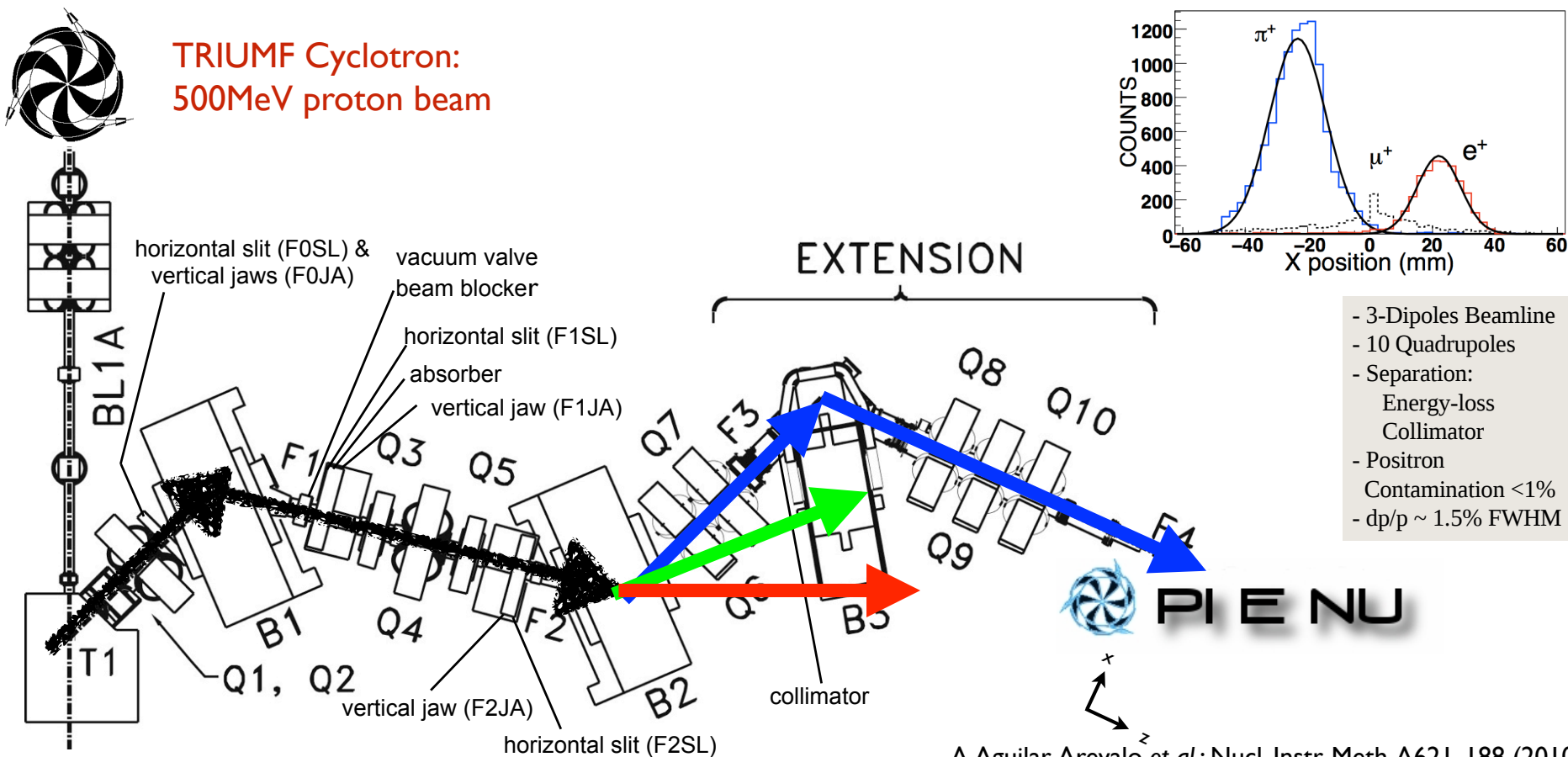
A.Aguilar-Arevalo et al.: Nucl. Instr. Meth. A621, 188 (2010)

TRIUMF Cyclotron:
500MeV proton beam

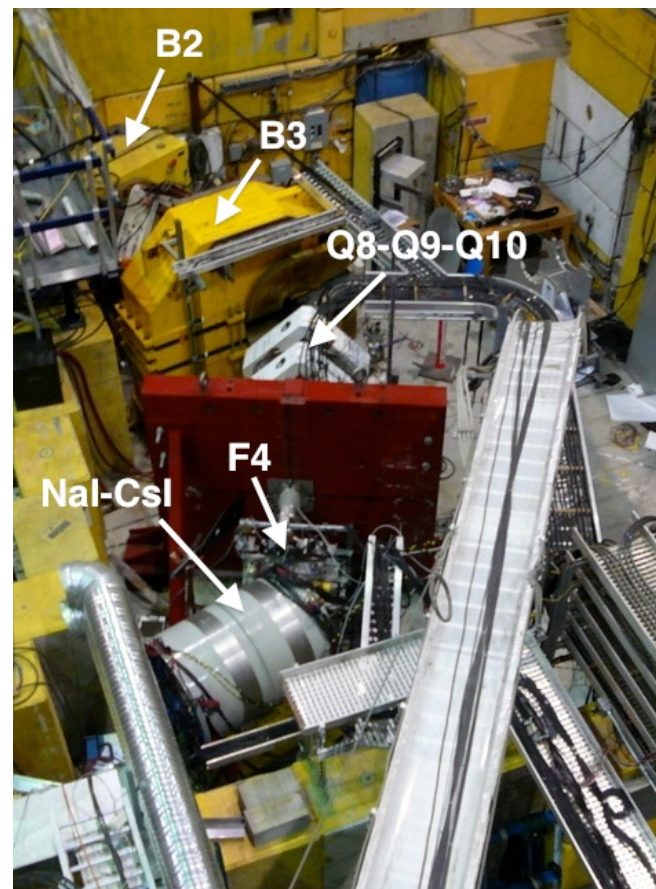
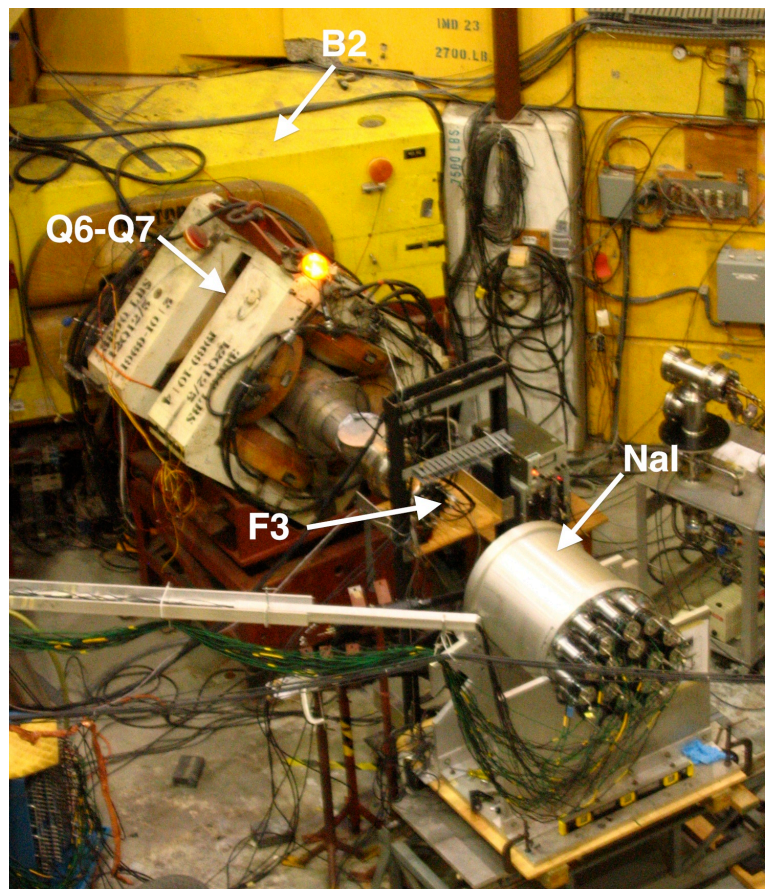


A.Aguilar-Arevalo et al.: Nucl. Instr. Meth. A621, 188 (2010)

TRIUMF Cyclotron:
500MeV proton beam



A.Aguilar-Arevalo et al.: Nucl. Instr. Meth. A621, 188 (2010)

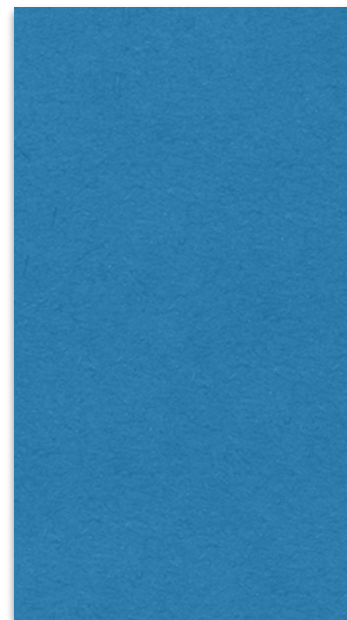
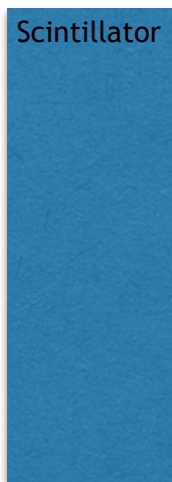


Idea: Stop the pions in an active target

Measure the decay positrons from the two decays

Advantages:

- Simultaneous measurement (**energy and time**)
- Same acceptance and conditions
- Systematic uncertainties cancel in the BR

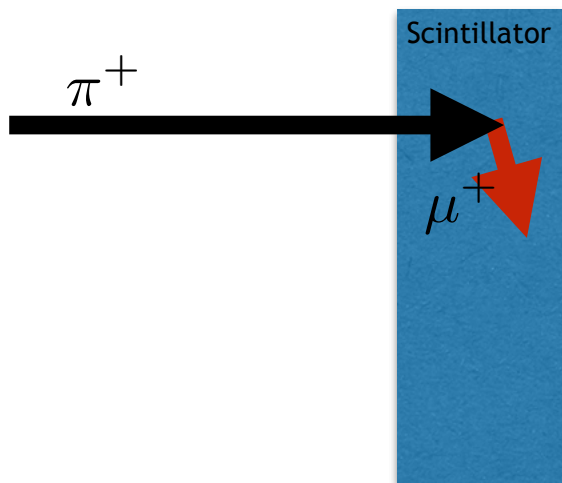


Idea: Stop the pions in an active target

Measure the decay positrons from the two decays

Advantages:

- Simultaneous measurement (**energy and time**)
- Same acceptance and conditions
- Systematic uncertainties cancel in the BR

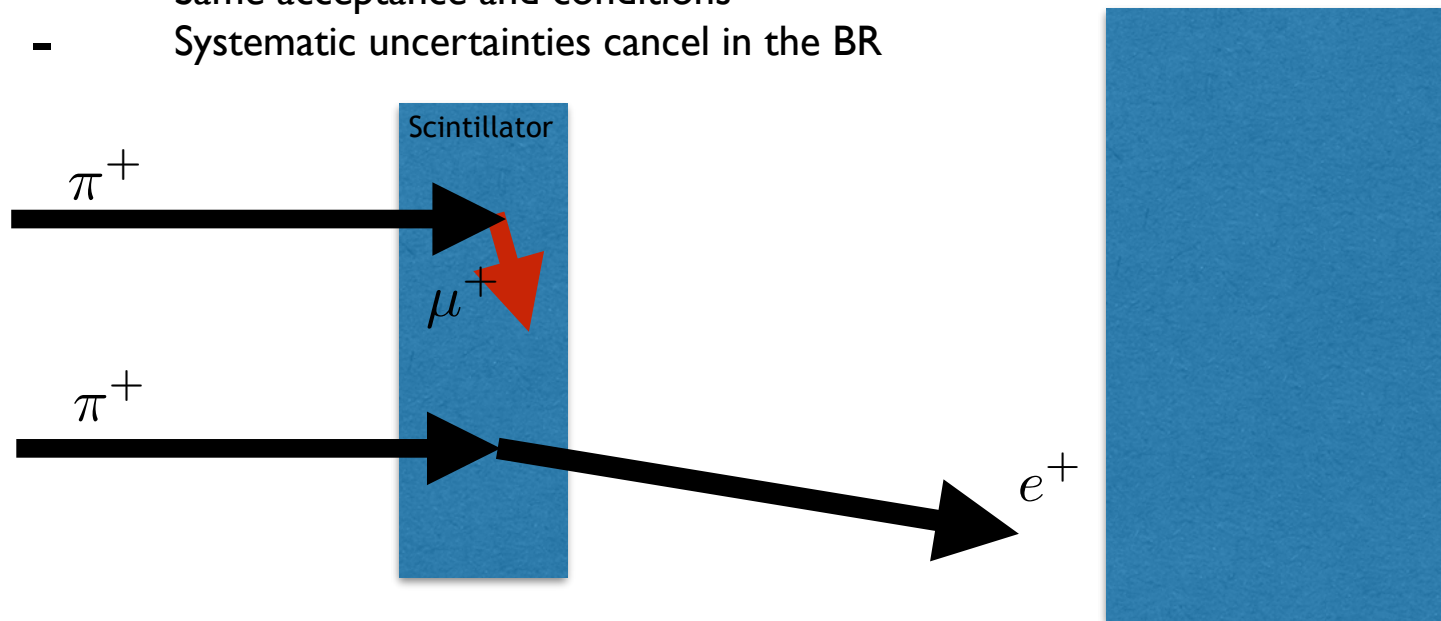


Idea: Stop the pions in an active target

Measure the decay positrons from the two decays

Advantages:

- Simultaneous measurement (**energy and time**)
- Same acceptance and conditions
- Systematic uncertainties cancel in the BR

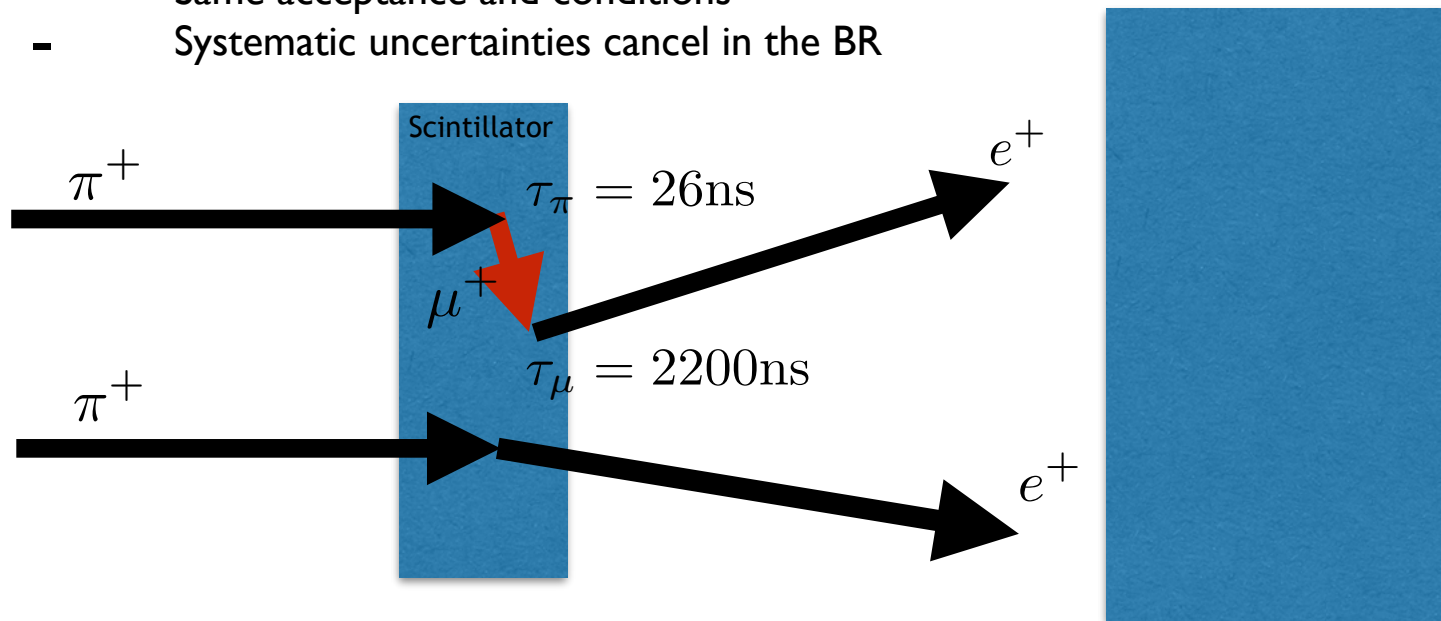


Idea: Stop the pions in an active target

Measure the decay positrons from the two decays

Advantages:

- Simultaneous measurement (**energy and time**)
- Same acceptance and conditions
- Systematic uncertainties cancel in the BR

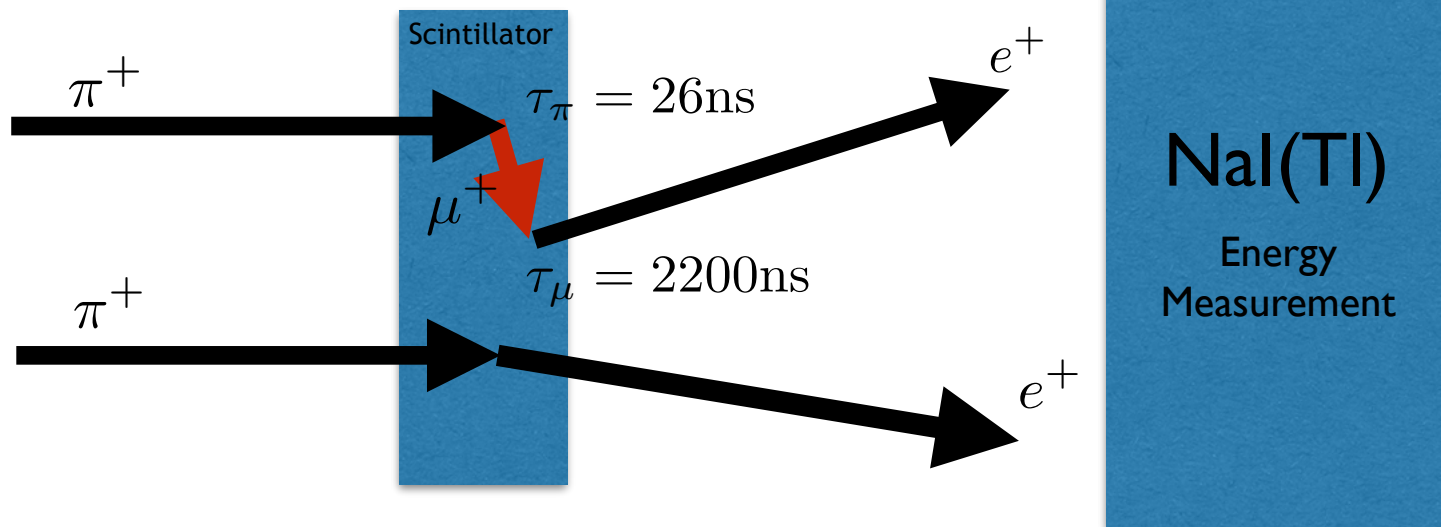


Idea: Stop the pions in an active target

Measure the decay positrons from the two decays

Advantages:

- Simultaneous measurement (**energy and time**)
- Same acceptance and conditions
- Systematic uncertainties cancel in the BR

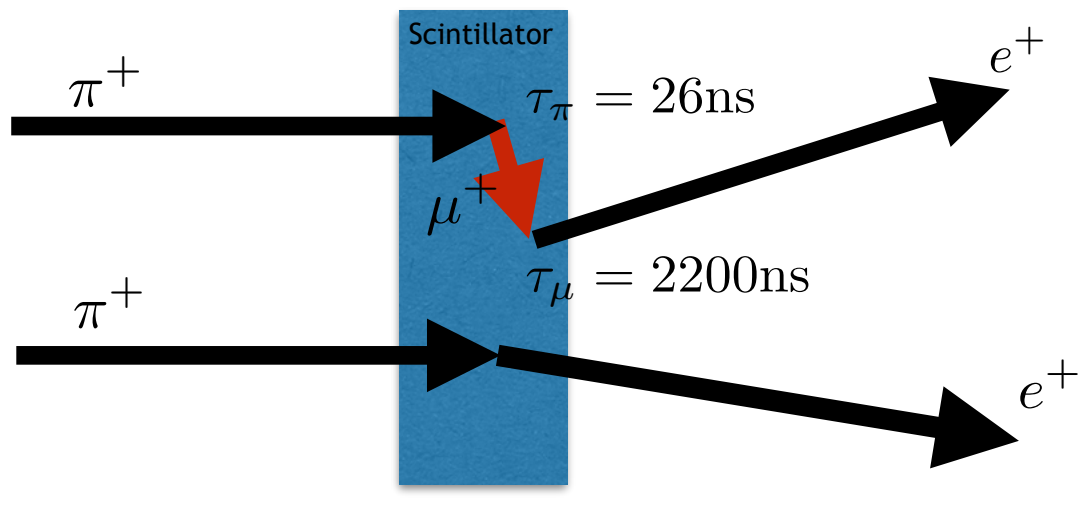
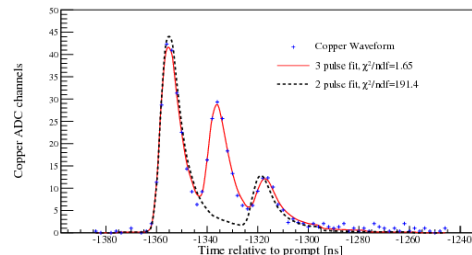


Idea: Stop the pions in an active target

Measure the decay positrons from the two decays

Advantages:

- Simultaneous measurement (**energy and time**)
- Same acceptance and conditions
- Systematic uncertainties cancel in the BR

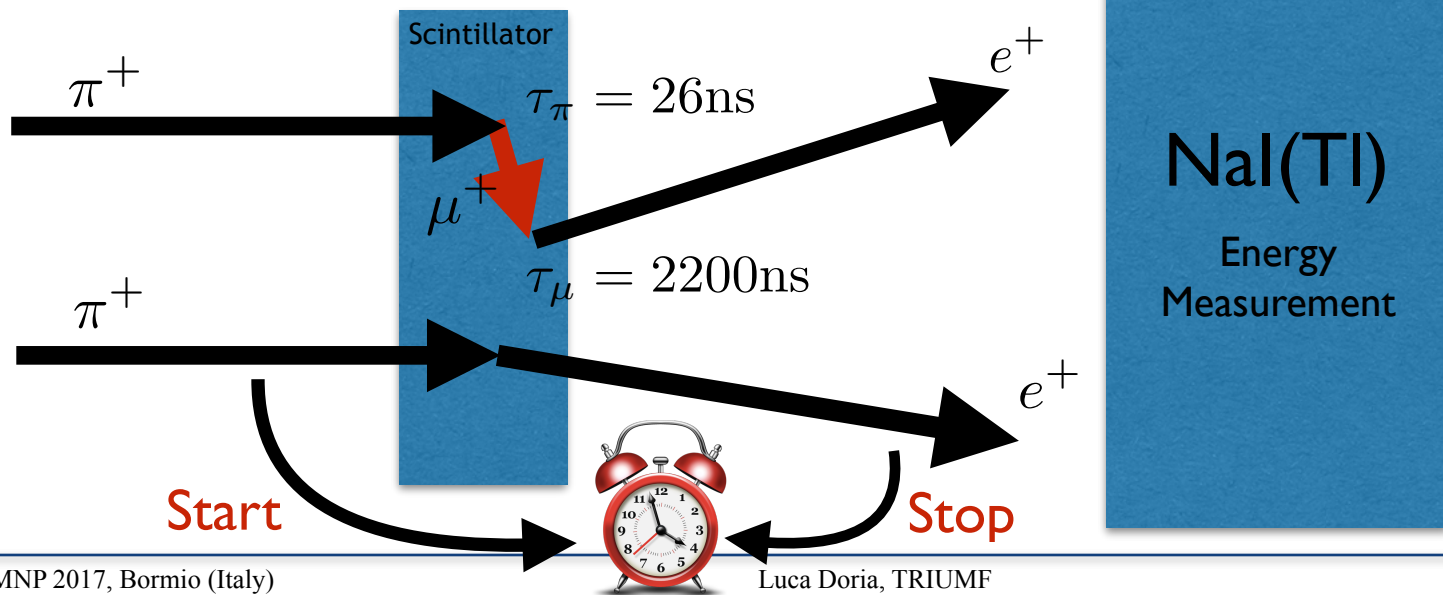
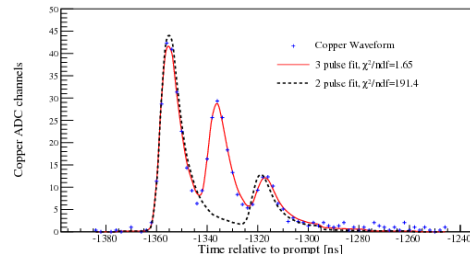


Idea: Stop the pions in an active target

Measure the decay positrons from the two decays

Advantages:

- Simultaneous measurement (**energy and time**)
- Same acceptance and conditions
- Systematic uncertainties cancel in the BR



Beam:

60kHz pions @ 75 MeV/c

$\pi : \mu : e = 85 : 14 : 1$

Detector:

Acceptance: 20%

Plastic Scintillators

NaI(Tl) + CsI Calorimeter

Wire Chambers

Silicon Strips

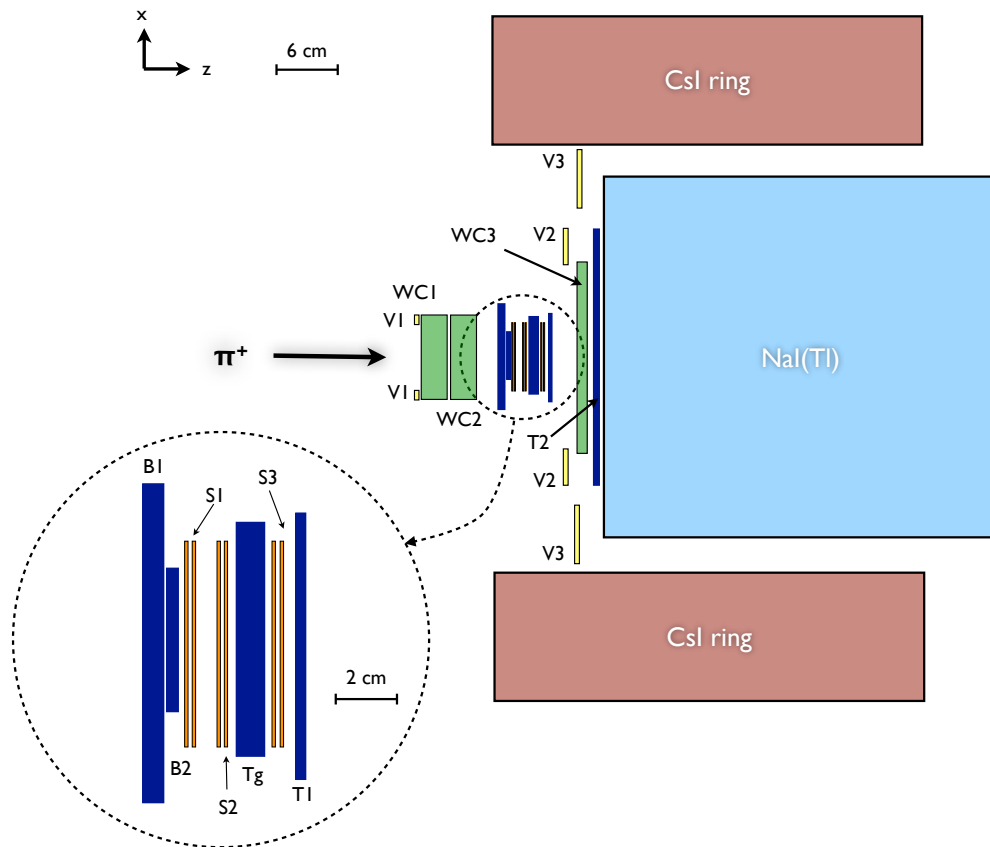
Energy resolution:

2.2% FWHM @ 70MeV

Temperature Stabilization

Data taking:

2009-2012



Beam:

60kHz pions @ 75 MeV/c

$\pi : \mu : e = 85 : 14 : 1$

Detector:

Acceptance: 20%

Plastic Scintillators

NaI(Tl) + CsI Calorimeter

Wire Chambers

Silicon Strips

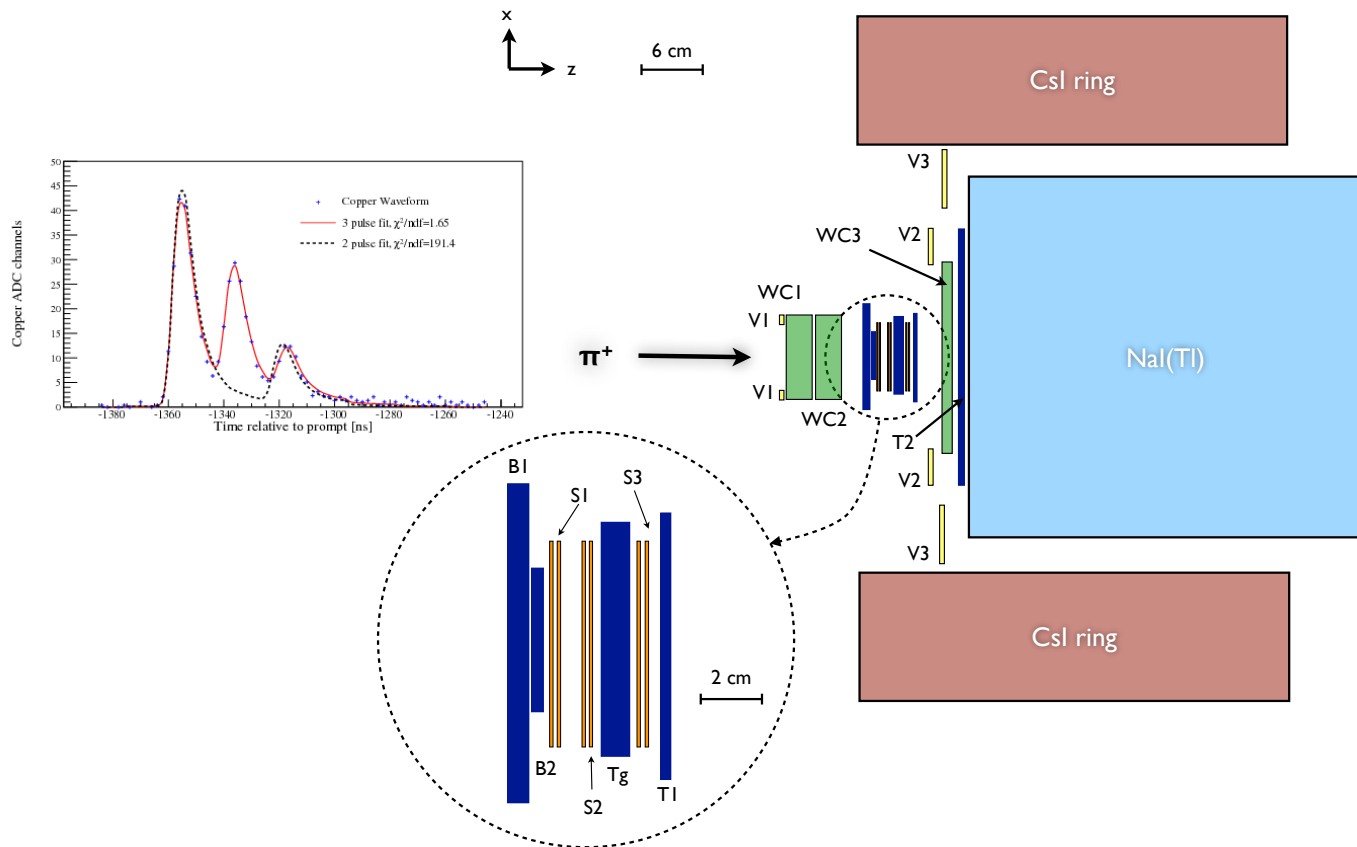
Energy resolution:

2.2% FWHM @ 70MeV

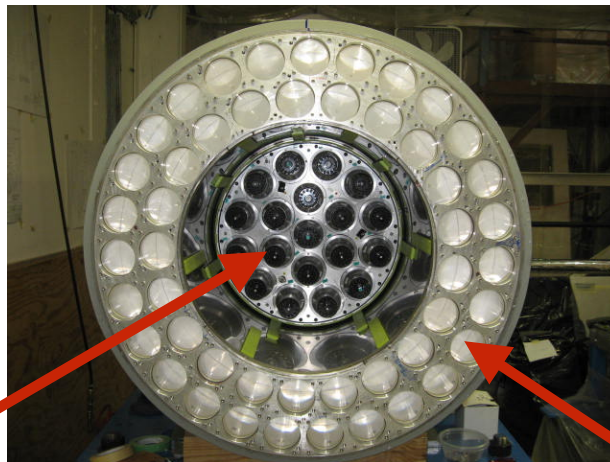
Temperature Stabilization

Data taking:

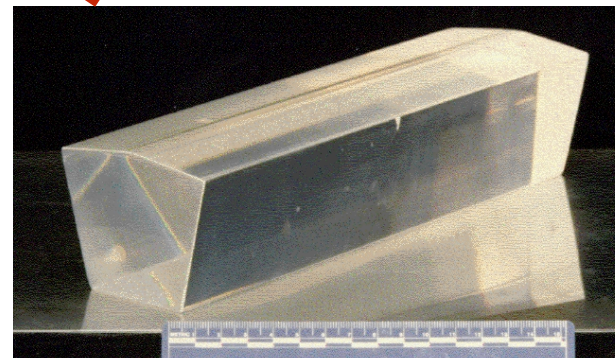
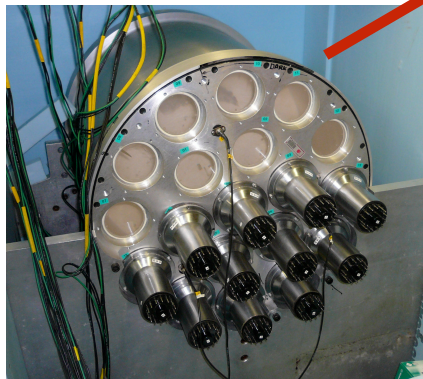
2009-2012

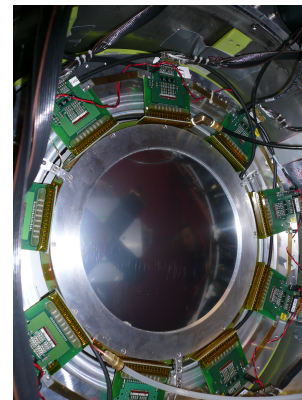
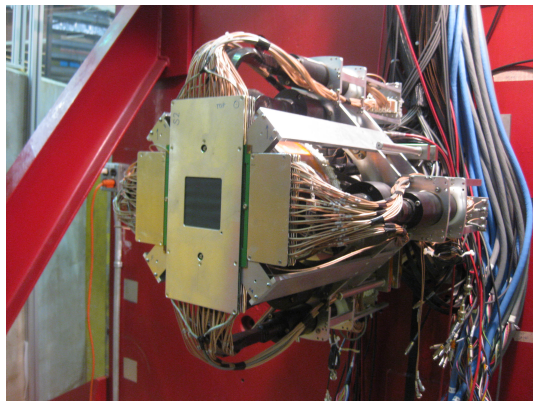
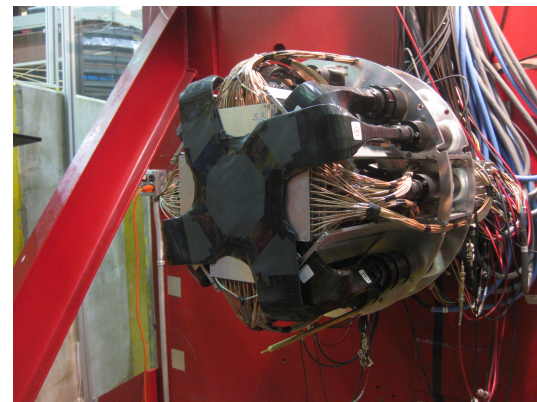
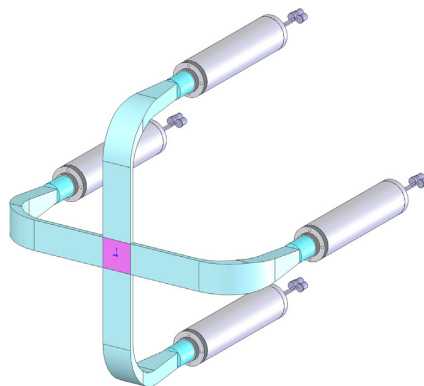
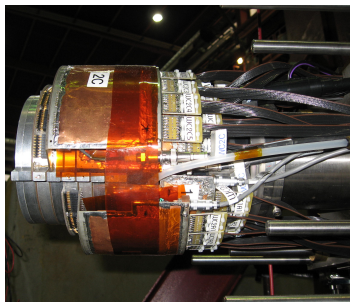
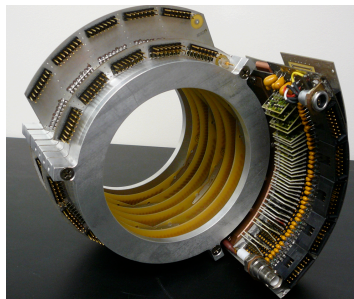


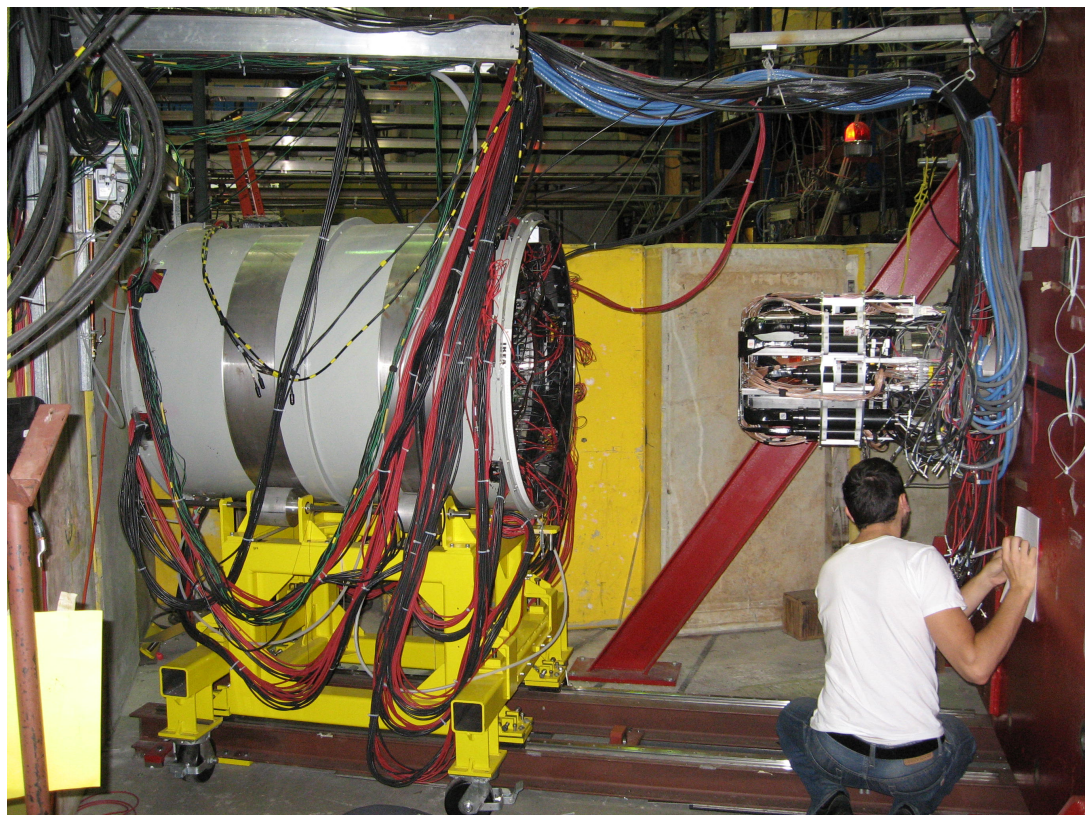
“BiNa”:
Monolithic 48x48cm
NaI(Tl) crystal
19-PMTs readout



97 pure CsI crystals
single PMT readout







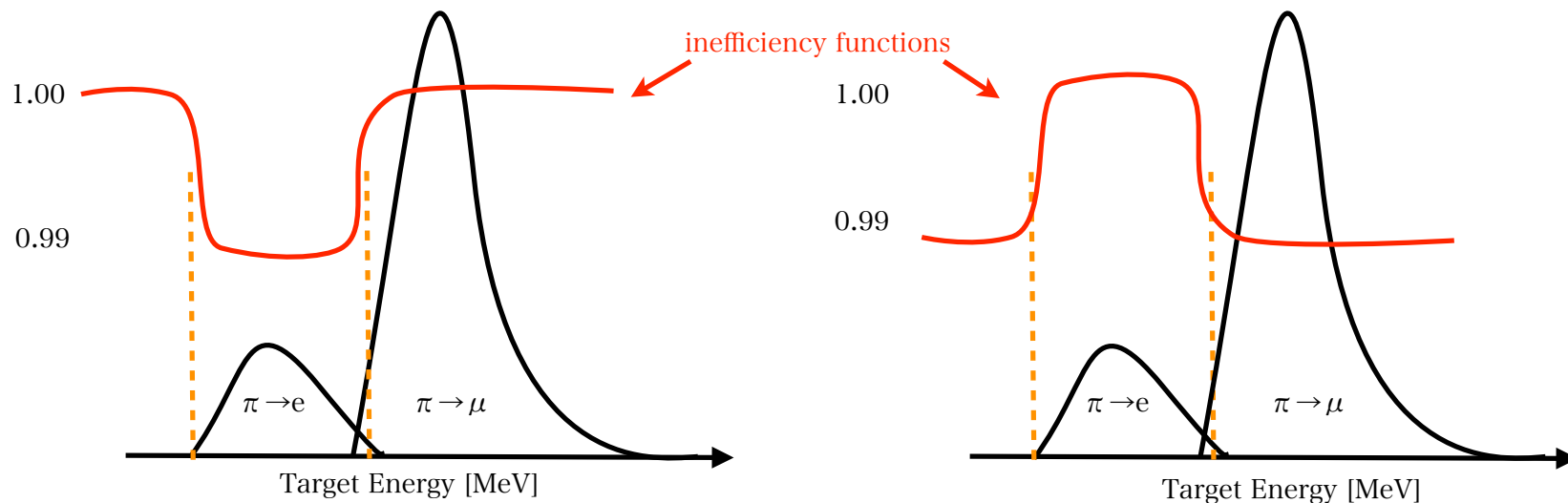
π^+ Beam

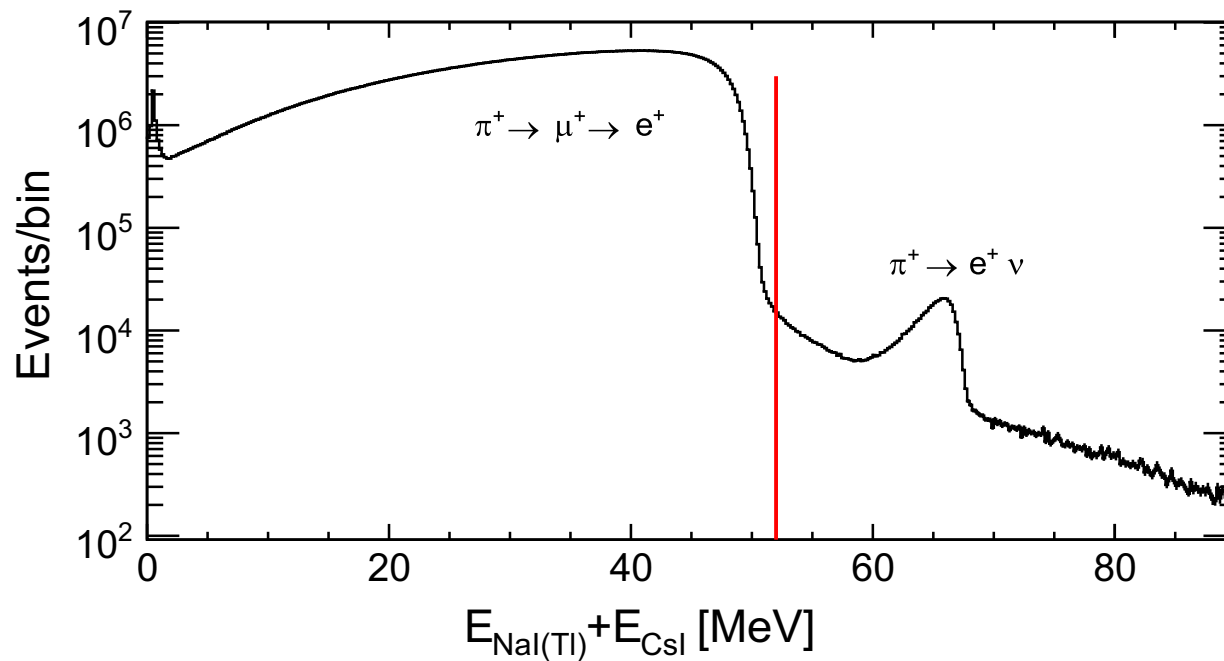


A.Aguilar-Arevalo et al: Nucl. Instr. Meth. A79, 38-46 (2015)

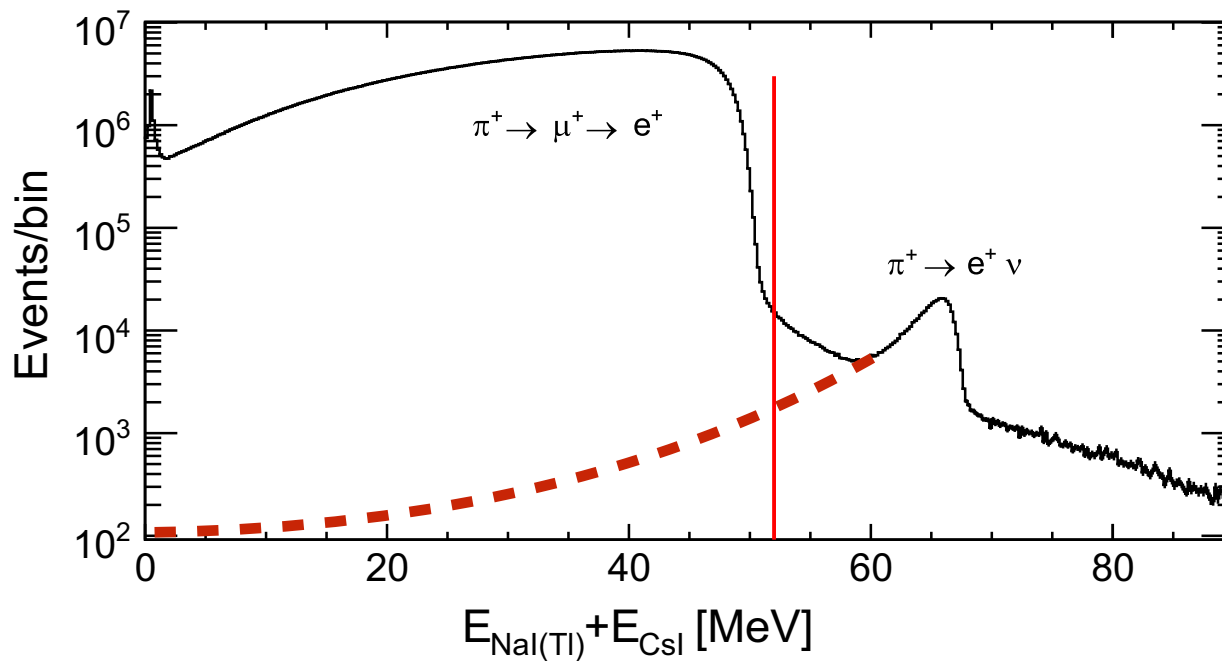
Data Analysis

- **Avoid biases in precision experiments!**
- Blinding procedure done before starting the analysis.
- One of the two decays is slightly suppressed: BR changes.
- Random and unknown inefficiency factor
- “Unblinding” only when the Collaboration agrees on the analysis procedure and systematic error estimates.

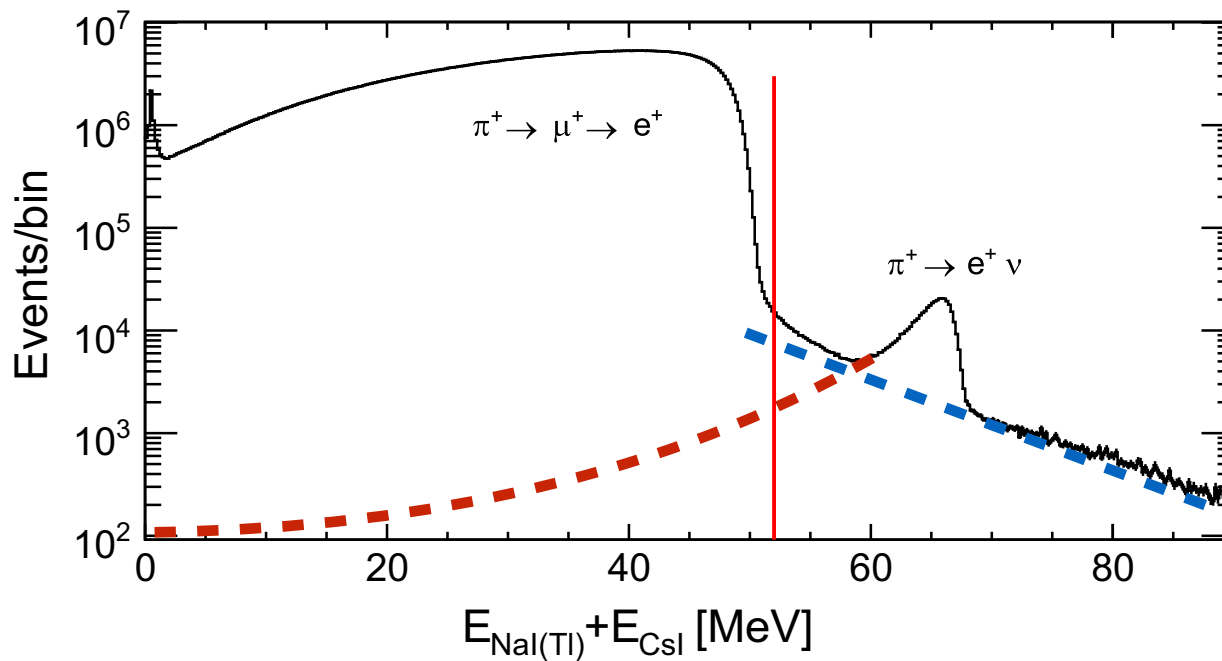




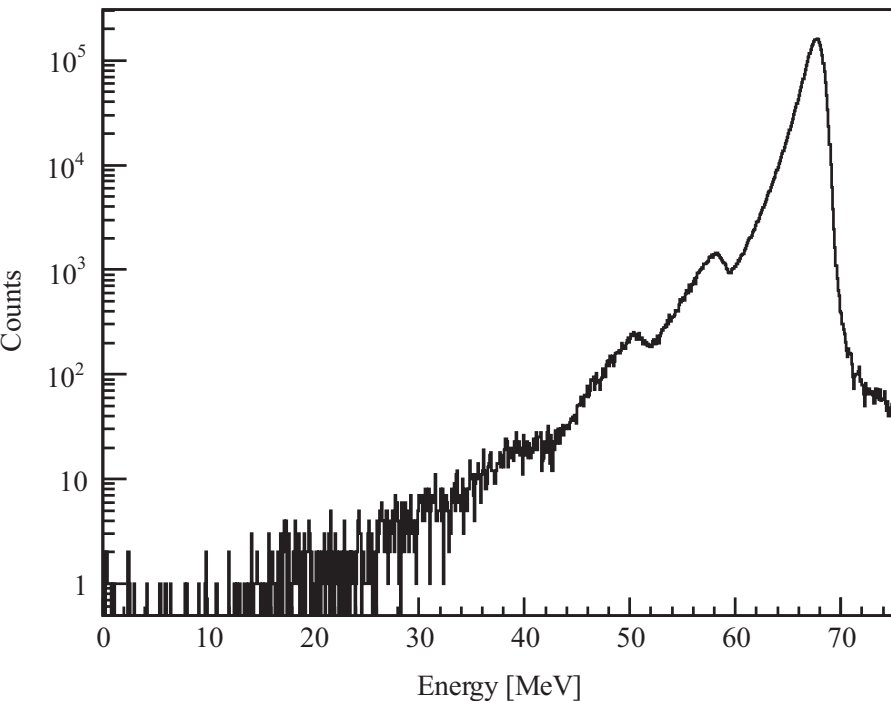
- Low Energy Tail
- Pileup



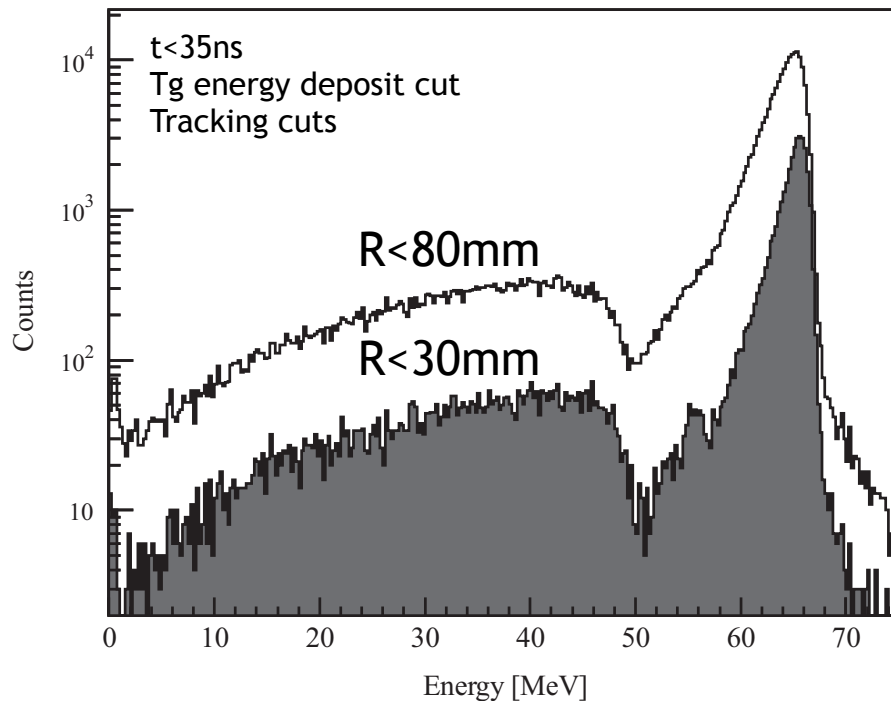
- Low Energy Tail
- Pileup



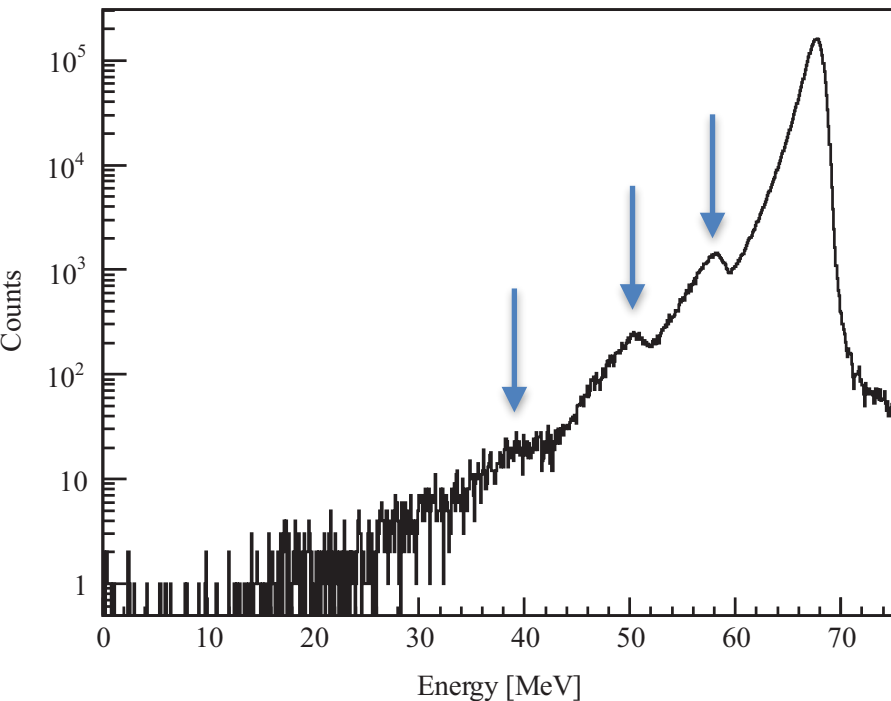
- Low Energy Tail
- Pileup



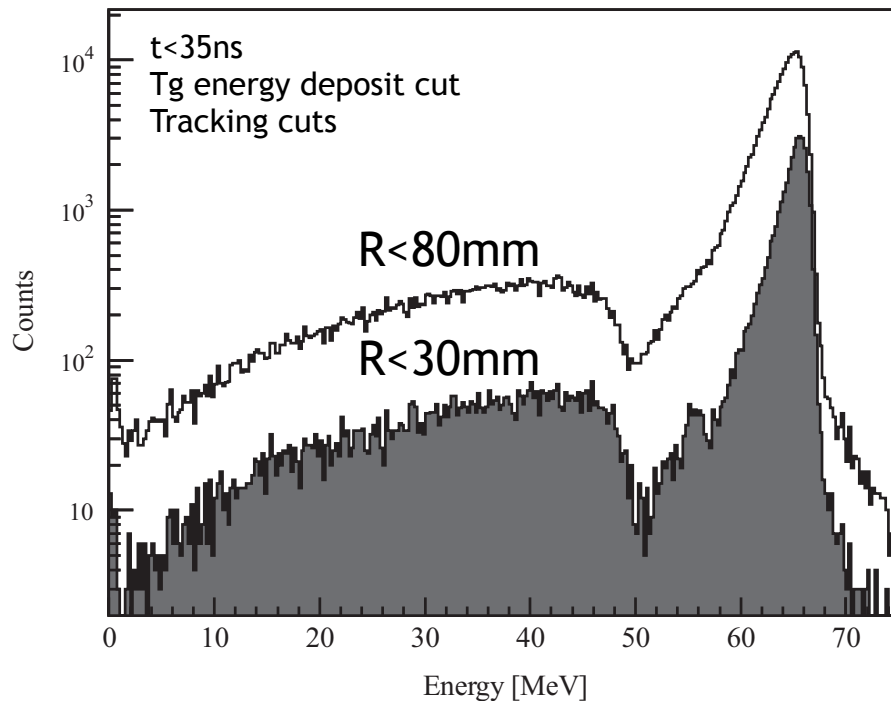
Positron Beam: Upper limit



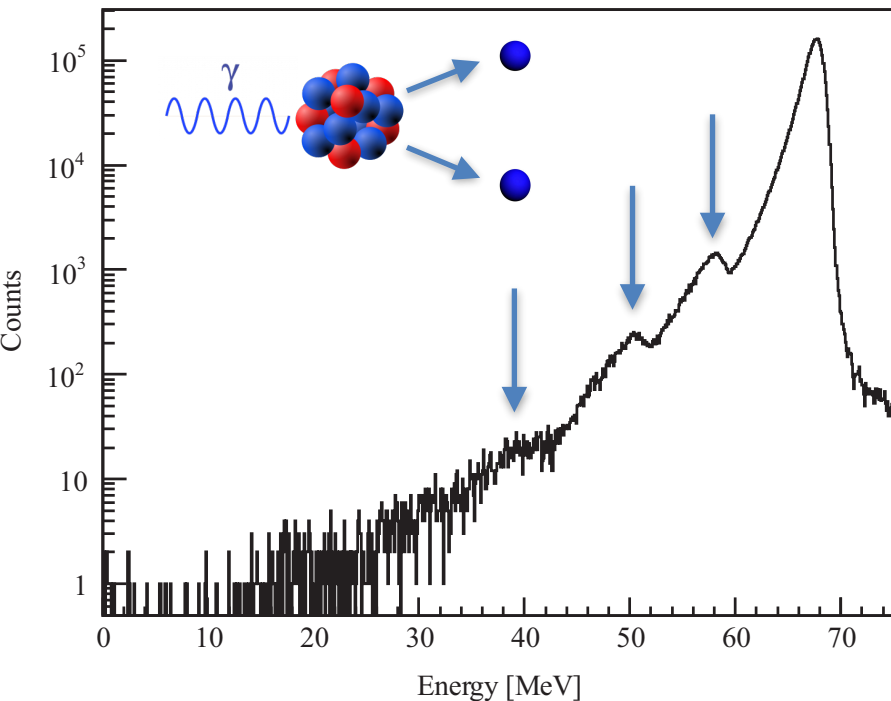
Suppressed Spectrum: Lower Limit



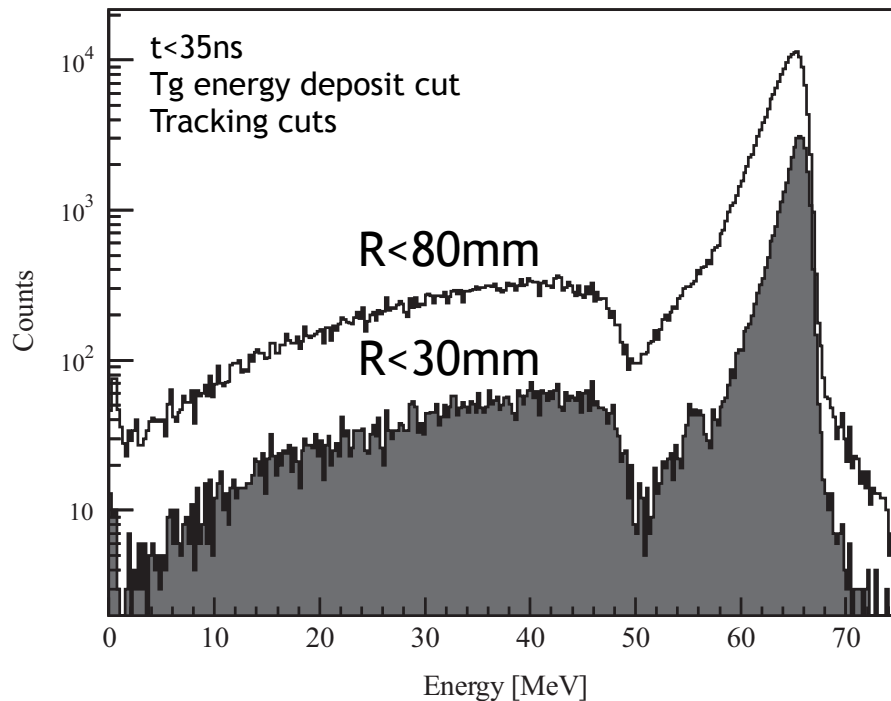
Positron Beam: Upper limit



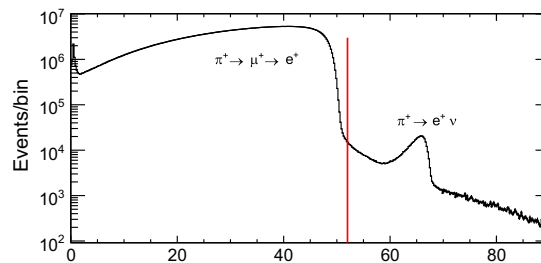
Suppressed Spectrum: Lower Limit



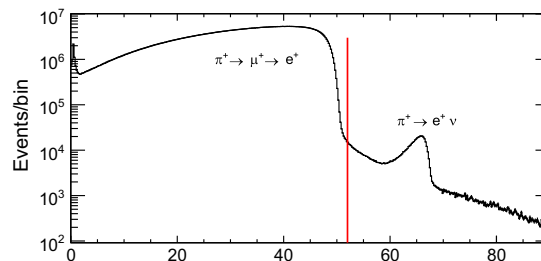
Positron Beam: Upper limit



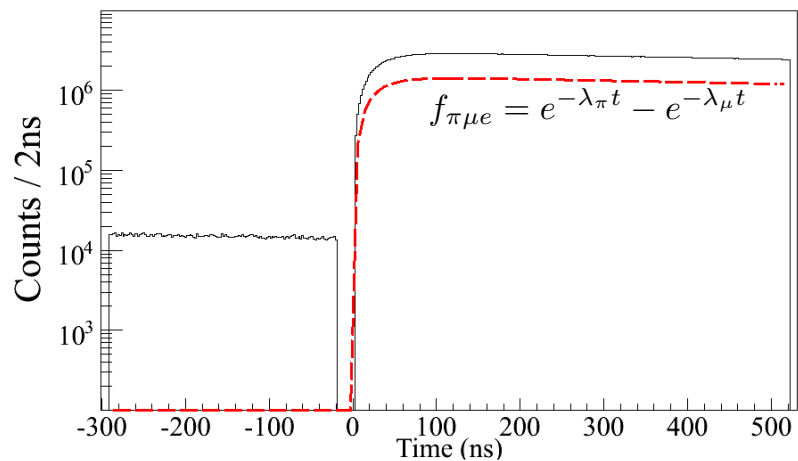
Suppressed Spectrum: Lower Limit



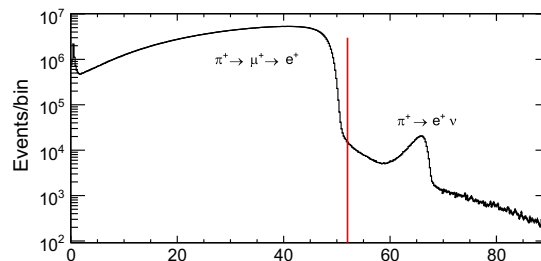
$E < 52 \text{ MeV}$



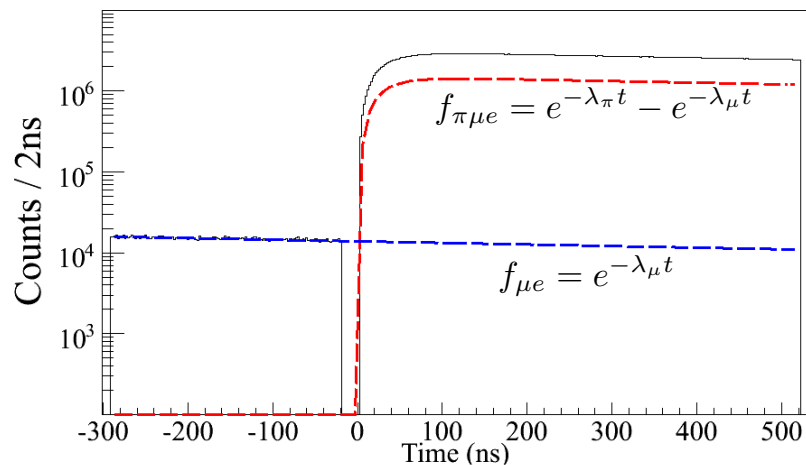
$E > 52 \text{ MeV}$



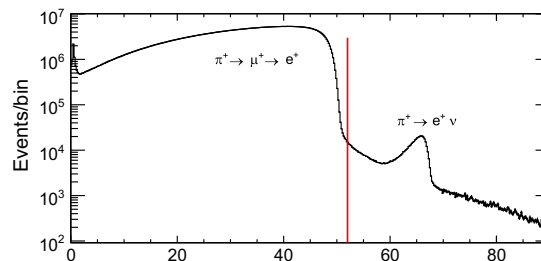
$E < 52 \text{ MeV}$



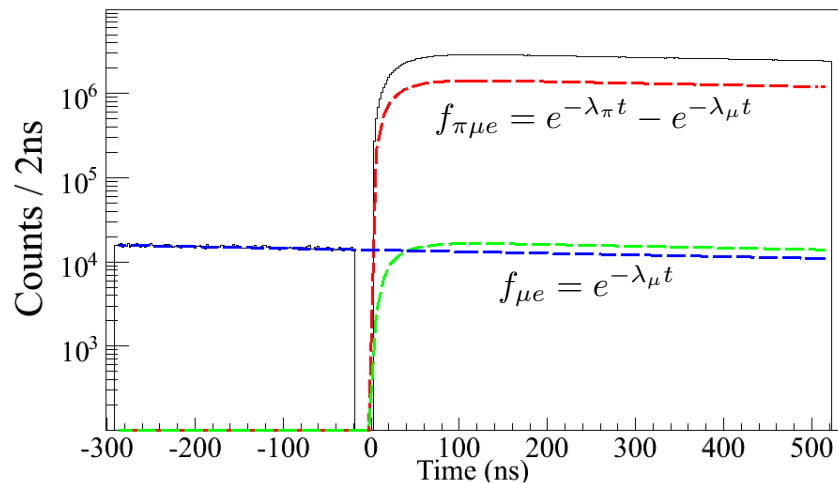
$E > 52 \text{ MeV}$



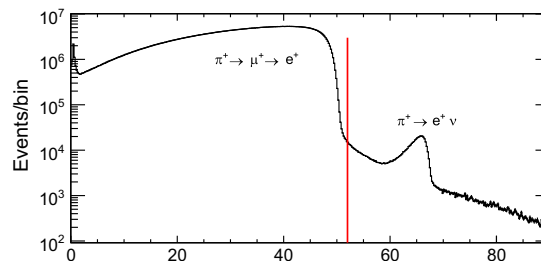
$E < 52 \text{ MeV}$



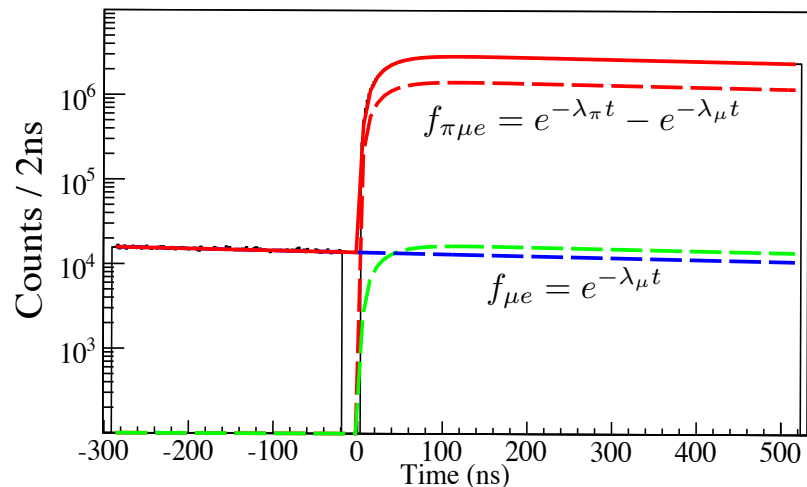
$E > 52 \text{ MeV}$



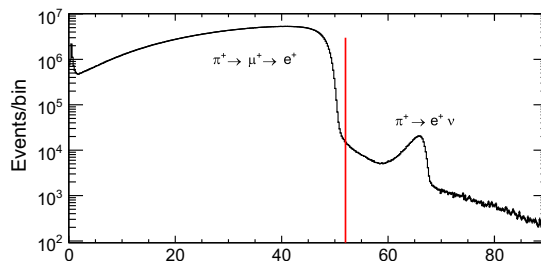
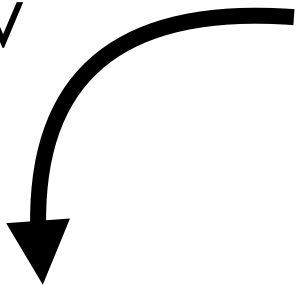
$E < 52 \text{ MeV}$



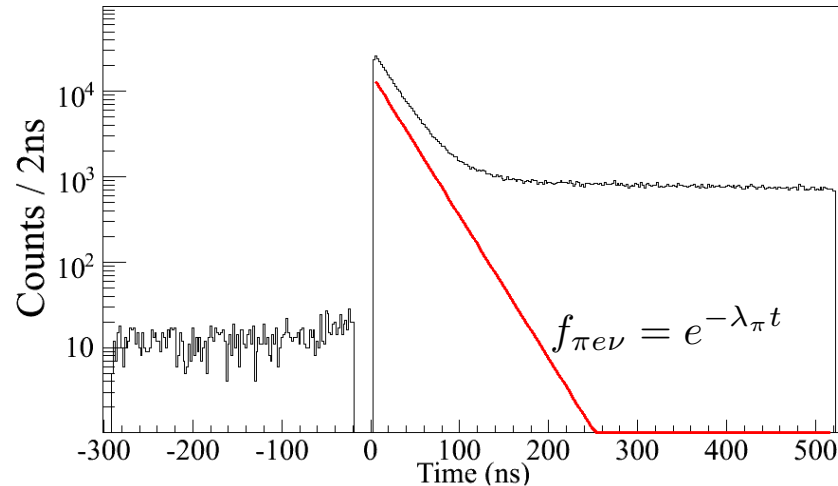
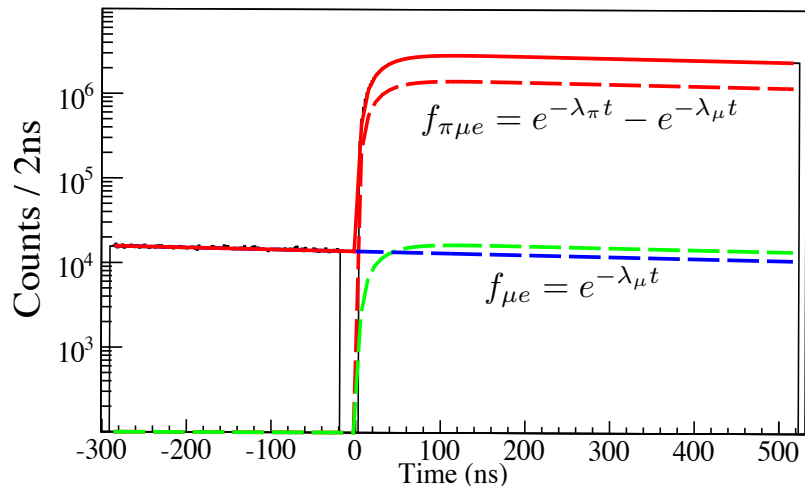
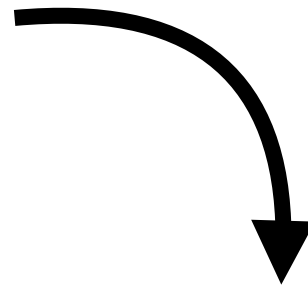
$E > 52 \text{ MeV}$



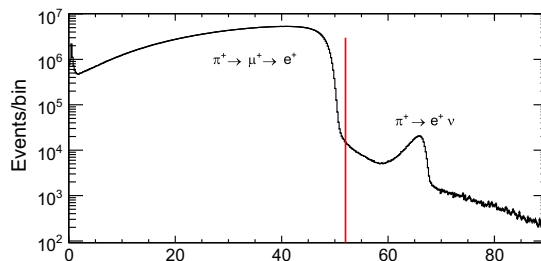
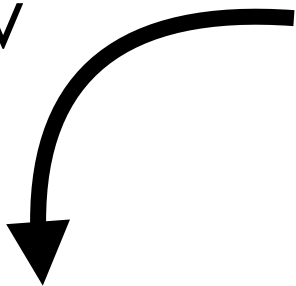
$E < 52 \text{ MeV}$



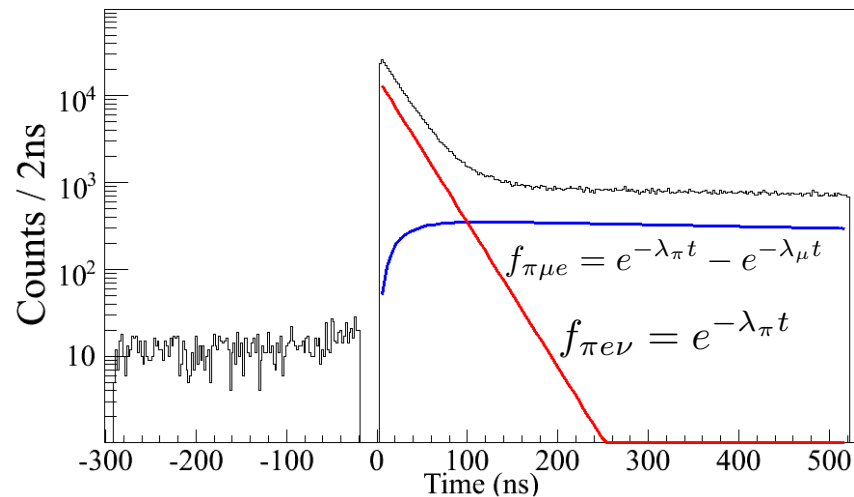
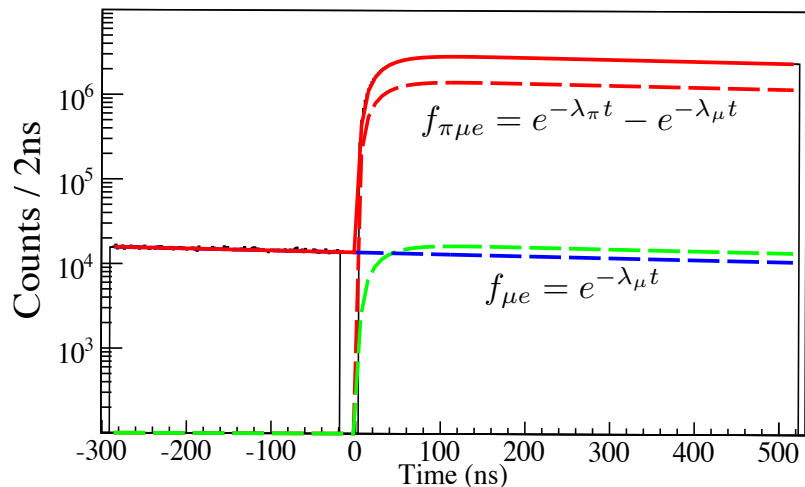
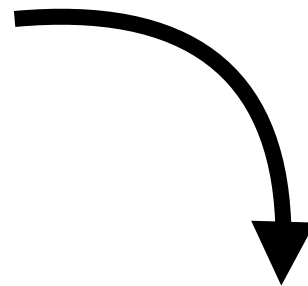
$E > 52 \text{ MeV}$



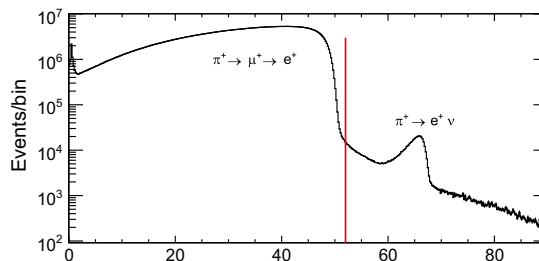
$E < 52 \text{ MeV}$



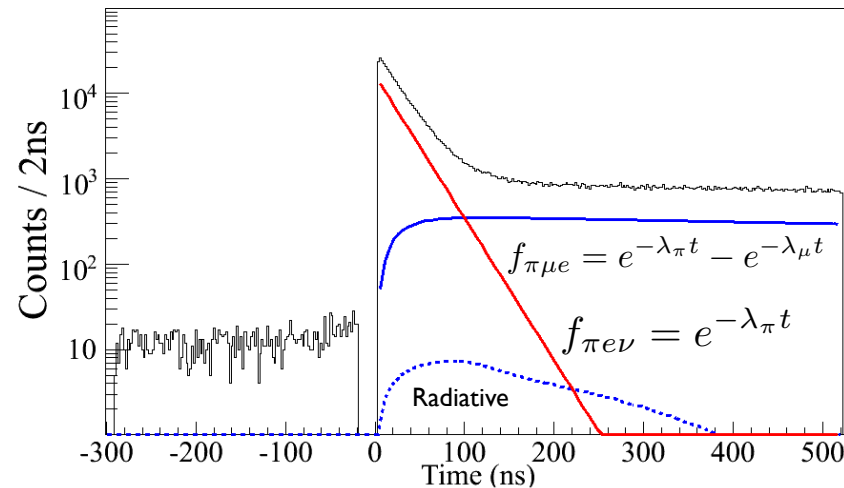
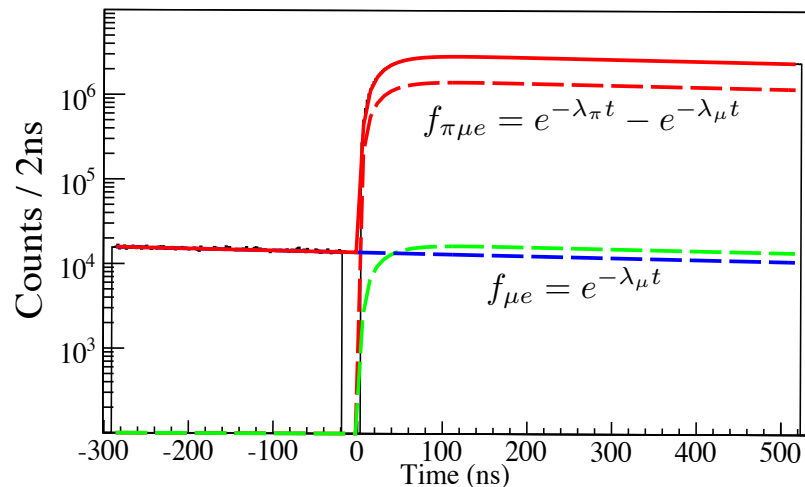
$E > 52 \text{ MeV}$

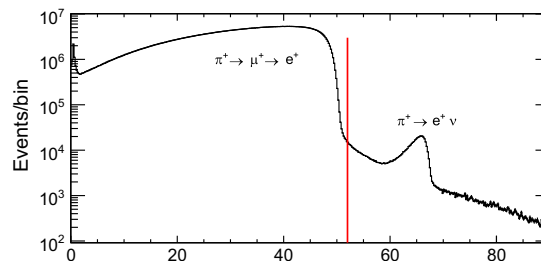
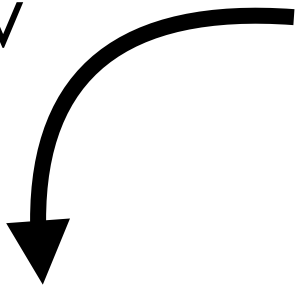
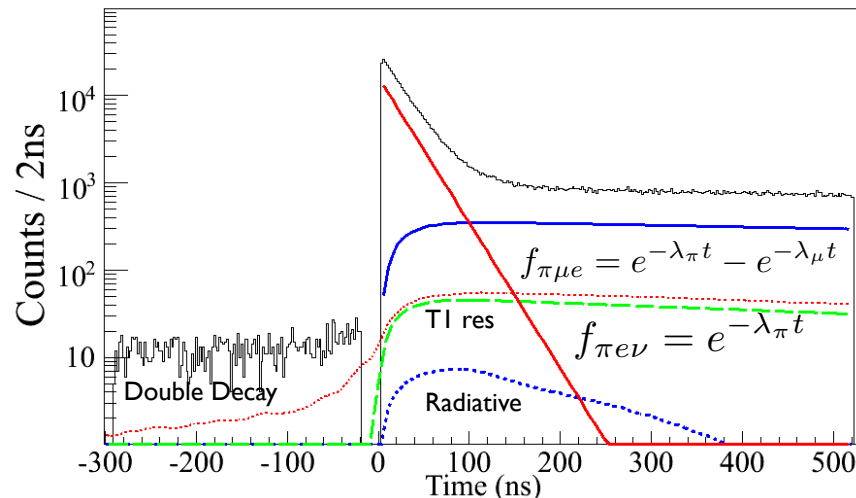
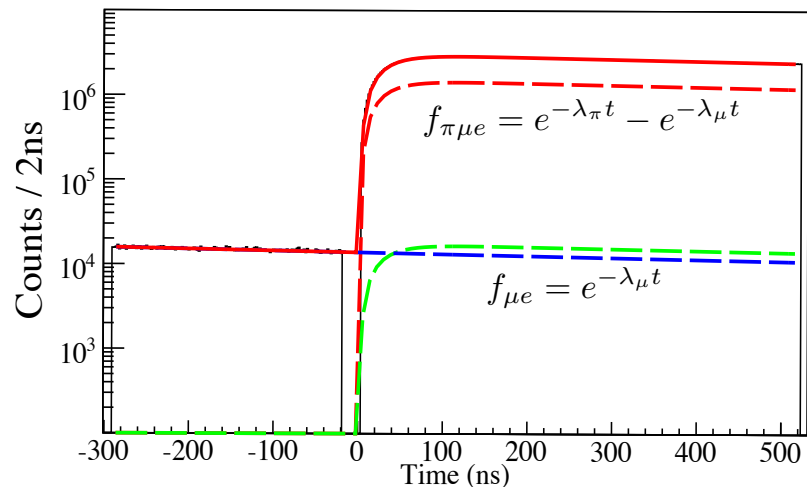
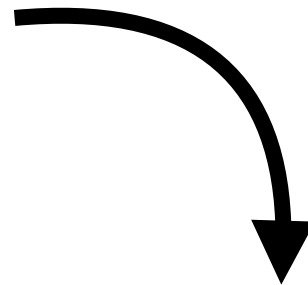


$E < 52 \text{ MeV}$

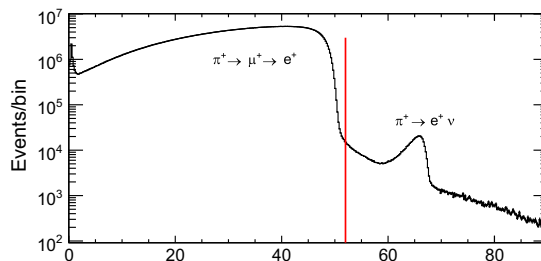


$E > 52 \text{ MeV}$

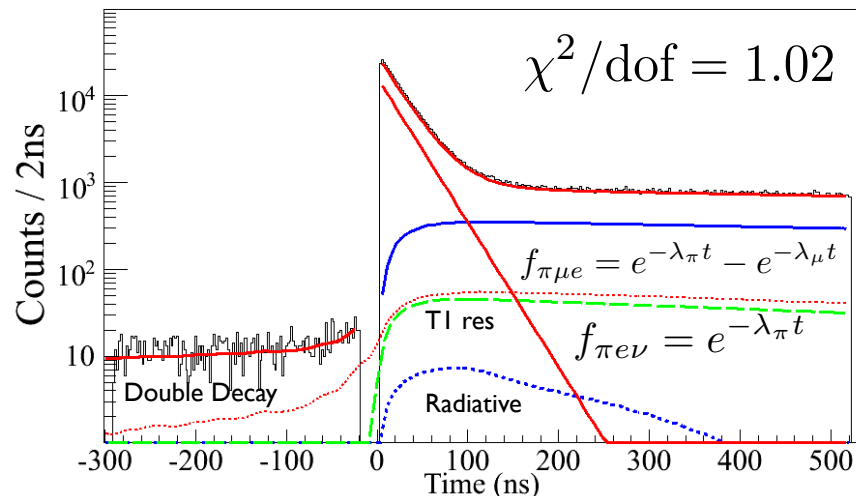
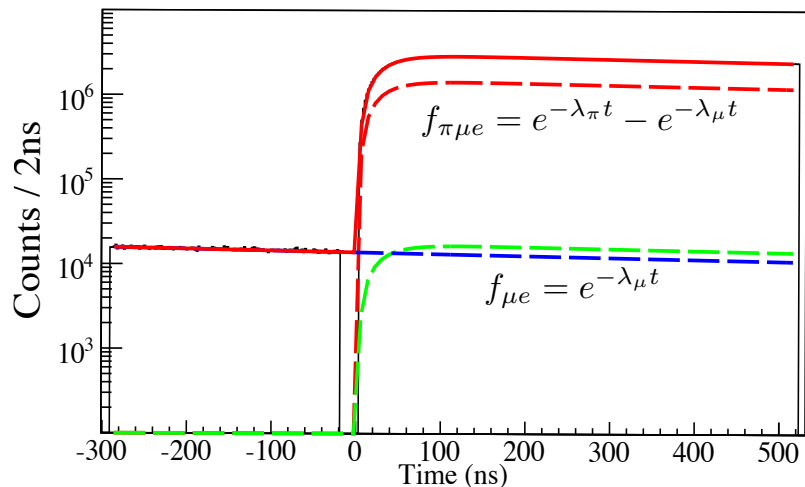


$E < 52 \text{ MeV}$

 $E > 52 \text{ MeV}$


$E < 52 \text{ MeV}$

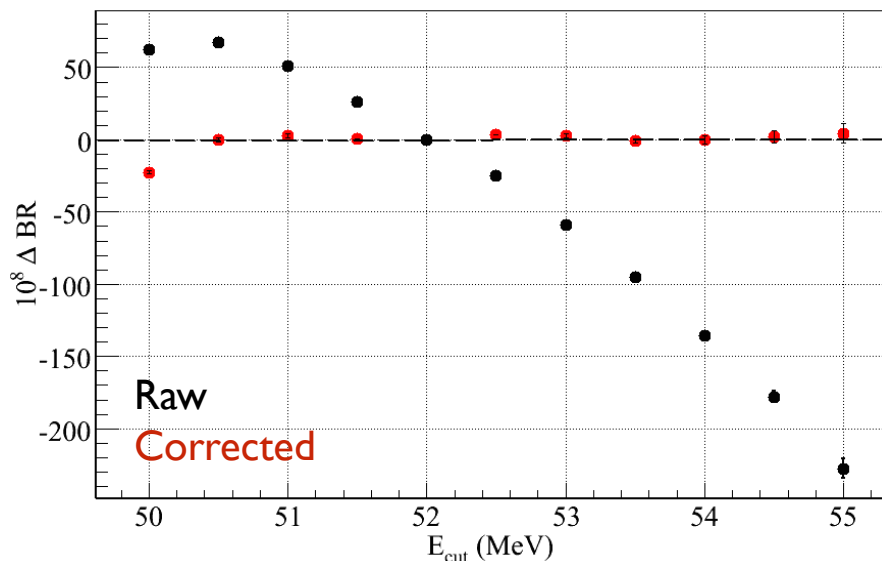


$E > 52 \text{ MeV}$



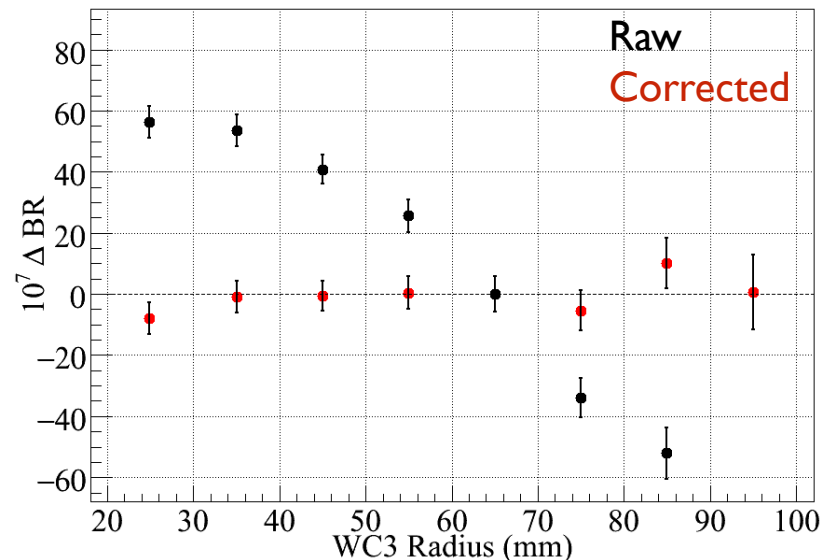
Acceptance Radius Dependence

- R= 60 mm
- Errors adjusted to statistics change
- Maximum R investigated with e^+ beam



Energy cut dependence

- Tail/muDIF corrections applied



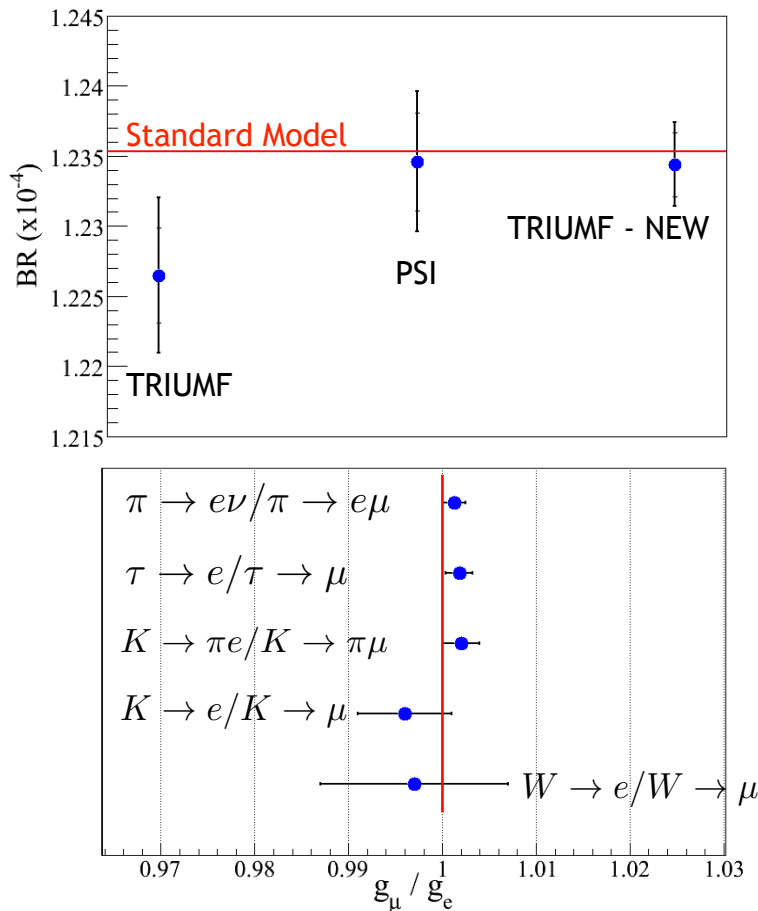
Results

Phys. Rev. Lett. 115, 071801 (2015)

	Values	Uncertainties	
		<i>Stat</i>	<i>Syst</i>
$R_{e/\mu}^{Raw} (10^{-4})$	1.1972	0.0022	0.0005
π, μ lifetimes			0.0001
other parameters			0.0003
excluded components			0.0005
Corrections			
Acceptance	0.9991		0.0003
Low energy tail	1.0316		0.0012
Other	1.0004		0.0008
$R_{e/\mu}^{Exp} (10^{-4})$	1.2344	0.0023	0.0019

$$R_{e/\mu}^{Th} (10^{-4}) = 1.2352(2)$$

$$e - \mu \text{ Universality: } g_e/g_\mu = 0.9996 \pm 0.0012$$



Dataset	BR	Status
2010	$1.2344 \pm 0.0023 \pm 0.0012$	Published
2011	$1.2XX \pm 0.0018 \pm 0.0013$	Completed, blind
2012	$1.2XX \pm 0.0009 \pm 0.00X$	In progress

Dataset	BR	Status
2010	$1.2344 \pm 0.0023 \pm 0.0012$	Published
2011	$1.2XX \pm 0.0018 \pm 0.0013$	Completed, blind
2012	$1.2XX \pm 0.0009 \pm 0.00X$	In progress

Final Goal: 0.1% precision

$\pi \rightarrow e\nu$ is a two-body decay

The pion decays at rest

→ Kinematics fully known if e^+ is measured:

$$m_\nu = \sqrt{m_\pi^2 + m_e^2 - 2m_\pi E_e}$$

If a massive neutrino can be produced, it will show up as a peak in the energy spectrum.

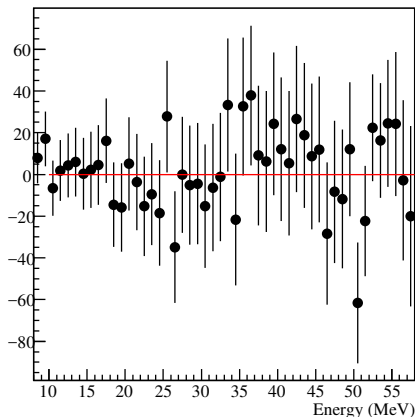
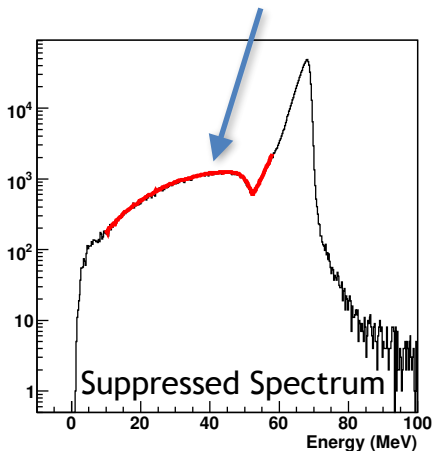
$\pi \rightarrow e\nu$ is a two-body decay

The pion decays at rest

→ Kinematics fully known if e^+ is measured:

$$m_\nu = \sqrt{m_\pi^2 + m_e^2 - 2m_\pi E_e}$$

If a massive neutrino can be produced, it will show up as a peak in the energy spectrum.



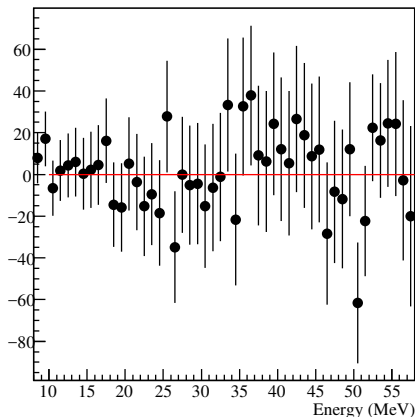
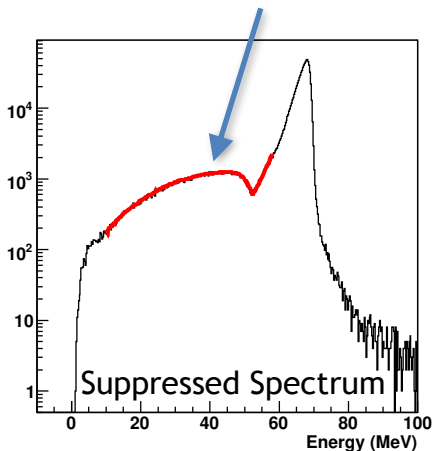
$\pi \rightarrow e\nu$ is a two-body decay

The pion decays at rest

→ Kinematics fully known if e^+ is measured:

$$m_\nu = \sqrt{m_\pi^2 + m_e^2 - 2m_\pi E_e}$$

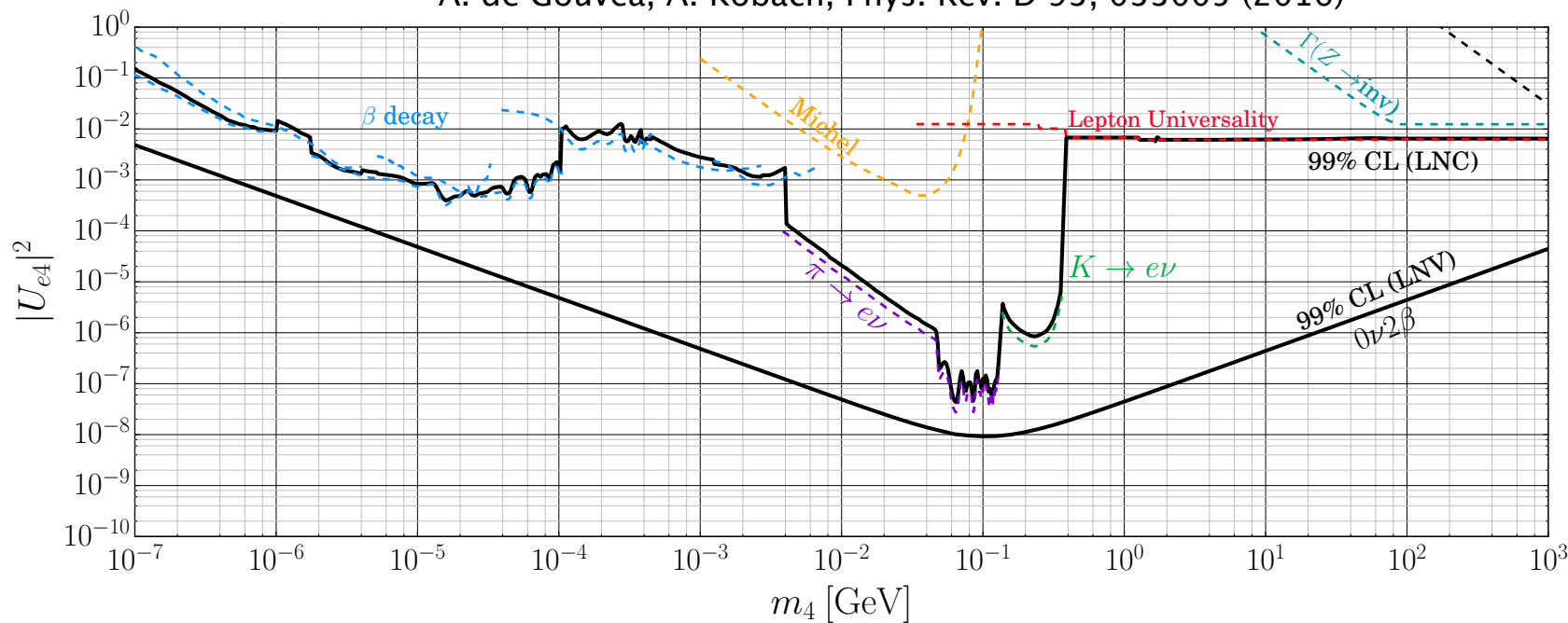
If a massive neutrino can be produced, it will show up as a peak in the energy spectrum.



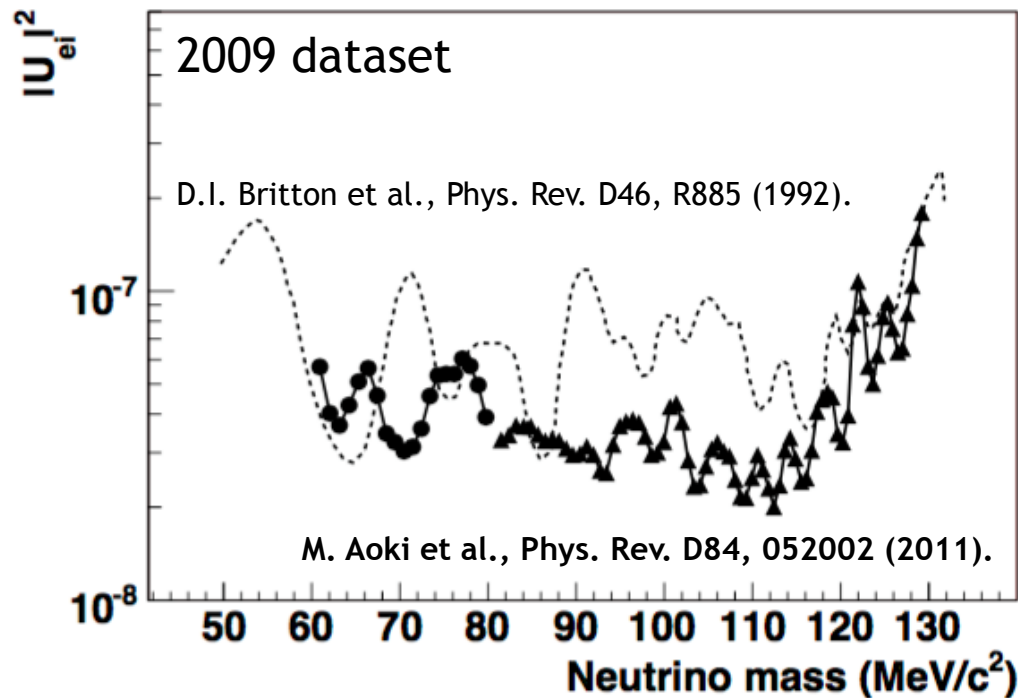
- 1) Consider the suppressed spectrum
- 2) Fit the spectrum with signal+bkg shapes:
 - $\pi \rightarrow \mu\nu \rightarrow e\nu\nu$ (data, $t > 150\text{ns}$)
 - Muon decays in flight (MC)
 - $\pi \rightarrow e\nu$ shape (MC)
- 3) Set upper limits to the BR for the pion decay to massive neutrinos.

$$\frac{N(\pi \rightarrow e\nu_i)_{UL}}{N(\pi \rightarrow e\nu_l)} = |U_{ei}|_{UL}^2 \rho_e$$

A. de Gouvêa, A. Kobach, Phys. Rev. D 93, 033005 (2016)



(a)



Full data analysis ongoing.

- Best limit on lepton universality violation established
- Work ongoing towards full dataset analysis
- Massive neutrino searches in the MeV range in pion and muon decays

- Best limit on lepton universality violation established
- Work ongoing towards full dataset analysis
- Massive neutrino searches in the MeV range in pion and muon decays

After >60 years the pion is still an important testing ground for the SM!

- Best limit on lepton universality violation established
- Work ongoing towards full dataset analysis
- Massive neutrino searches in the MeV range in pion and muon decays

After >60 years the pion is still an important testing ground for the SM!

Final results coming soon: stay tuned!



Canada's national laboratory
for particle and nuclear physics
and accelerator-based science

Thank you! Merci!

TRIUMF: Alberta | British Columbia | Calgary |
Carleton | Guelph | Manitoba | McGill | McMaster |
Montréal | Northern British Columbia | Queen's |
Regina | Saint Mary's | Simon Fraser | Toronto |
Victoria | Western | Winnipeg | York

Follow us at TRIUMFLab

