

Recent Results on Soft QCD Topics from ATLAS

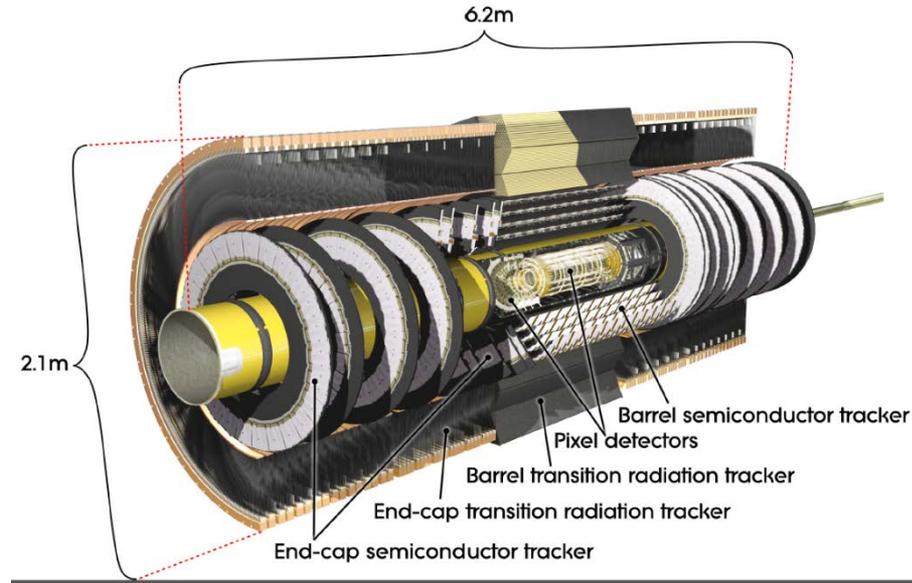
**Bormio, 55. International Winter
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on behalf of the ATLAS Collaboration**

Outline

- The ATLAS detector is designed to investigate very rare high p_T processes in the environment of high luminosities and large number (several tens) of interactions per one bunch crossing
- But ATLAS provides also a window onto important soft QCD processes which have large cross sections and the data for such studies are taken in special low luminosity runs
- These processes have intrinsic interest. They can not be predicted by a perturbative QCD. Various MC models, having a large number of adjustable tuning parameters, are used to describe them
- As well they are needed to understand pile-up and underlying event activity in *all* LHC measurements
- This presentation concerns the 4 following recent ATLAS soft QCD studies:
 - I. **Charged-particles distributions in $\sqrt{s} = 13$ TeV pp interactions**
 - Phys.Lett. B (2016) 758
 - II. **Measurement of charged-particle distributions sensitive to underlying event in $\sqrt{s} = 13$ TeV pp collisions**
 - arxiv:1701.05390
 - III. **Study of hard double-parton scattering in four jet events in pp collisions at $\sqrt{s} = 7$ TeV**
 - JHEP 11 (2016) 110
 - IV. **Measurement of the total cross section from elastic scattering in pp collisions at $\sqrt{s} = 8$ TeV**
 - Phys.Lett. B (2016) 158

I-II. Inner Detector

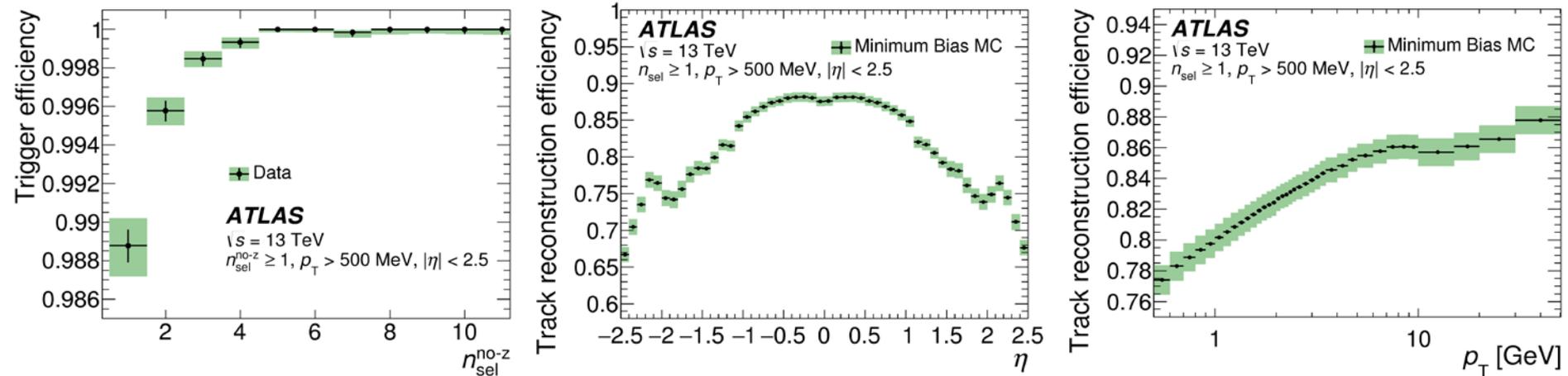


- **Studies I and II are based on the reconstructed in the Inner Detector (ID) (Figure above) charged tracks**
- **ID has full coverage in ϕ and covers the pseudorapidity range $|\eta| < 2.5$ and it is placed inside 2T axial magnetic field**
- **Its barrel (end-cap) part consists of**
 - four (three) pixel layers
 - four (nine) double-layers of single-sided silicon microstrips
 - 73 (160) layers of TRT (Transition Radiation Tracker) straws
- **A typical track in the barrel part has 4 pixel hits, 8 SCT hits and more than 30 TRT straw hits**
- **Minimum-bias trigger scintillators (MBTS), covering region $2.07 < |\eta| < 3.86$, are used for triggering. They are mounted at each end of the ID and segmented into 12 sectors each**

I-II. Data selection, corrections

- **Special low luminosity run with $\langle\mu\rangle = 0.005$, about 9M events**
- **Trigger: one or more fired MBTS counters on either side of the ID**
- **A number of corrections used and their uncertainties included to the corresponding systematic errors**
- **The corrections account for inefficiencies due to trigger selection, vertex and track reconstruction, background from the secondary interactions**
- **The trigger efficiency is rather high (next slide) as well as a vertex reconstruction efficiency**

I-II. Trigger and Track Reconstruction Efficiency

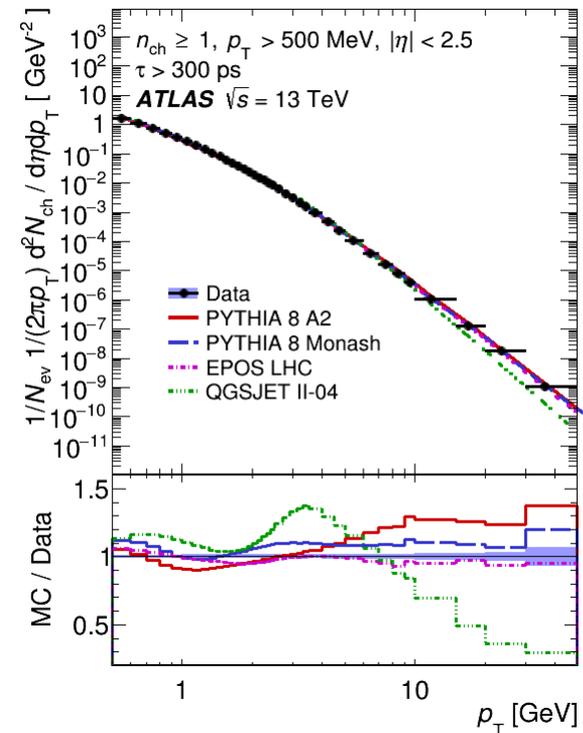
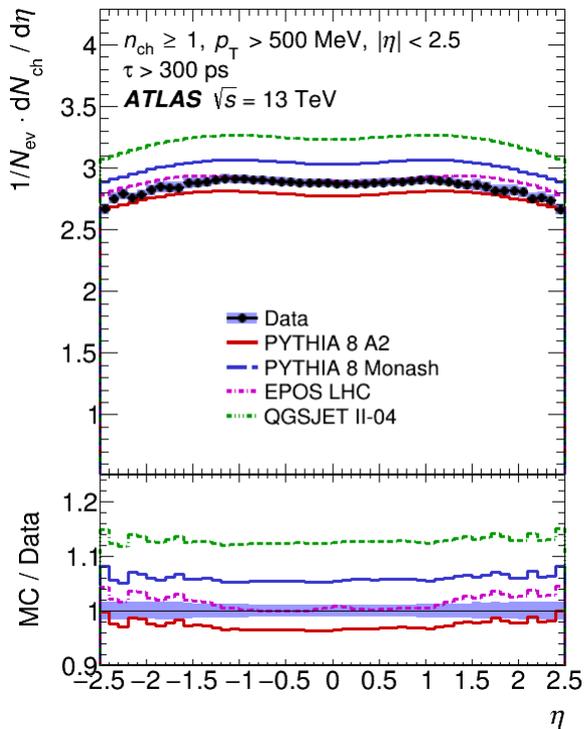


- The primary track reconstruction efficiency, $\varepsilon_{\text{trk}}(p_T, \eta)$, is determined from the simulation
- ε_{trk} depends on the amount of material in the detector, due to particle interactions that lead to efficiency losses
- The resulting reconstruction efficiency as a function of η and p_T is shown in the above Figures together with its uncertainty
- The efficiency is lower at $|\eta| > 1$ due to more material in that region
- The total uncertainty due to imperfect knowledge of the detector material is $\pm 0.7\%$ in the most central region and $\pm 1.5\%$ in the most forward region
- The Bayesian unfolding is applied to the p_T and multiplicity distributions to correct from the observed track distributions to those for primary charged particles

I-II. Monte Carlo generators used

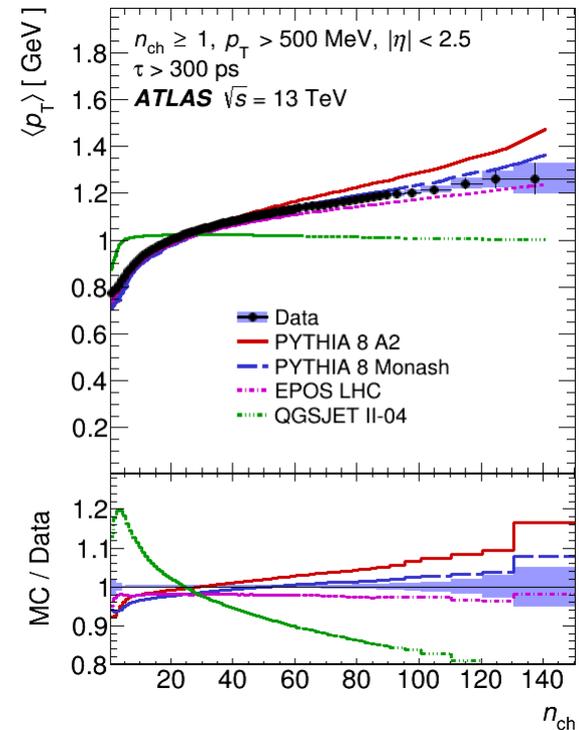
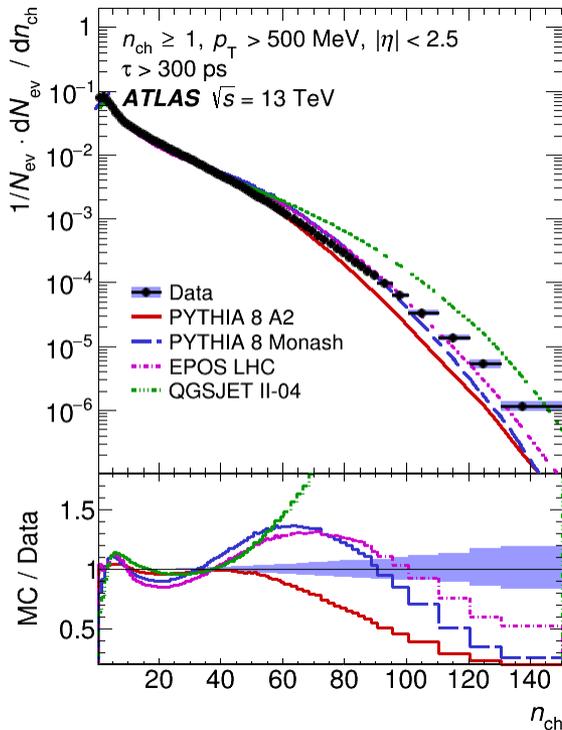
- **PYTHIA8:**
 - includes non-diffractive (ND) processes dominated by t-channel gluon exchange
 - and diffractive processes involving a color-singlet exchange
 - includes multi parton interactions (MPI)
 - for a hard scattering uses perturbative $2 \rightarrow 2$ QCD matrix element with leading logarithmic initial and final state parton showers
 - Lund string hadronization model
- **Herwig:**
 - for a hard scattering uses perturbative $2 \rightarrow 2$ QCD matrix element with leading logarithmic initial and final state parton showers
 - includes MPI (simpler parametrization than in PYTHIA)
 - a cluster-oriented color-disruption mechanism is used in hadronization
- **EPOS:**
 - a specialist softQCD/cosmic-ray air-shower MC generator based on an implementation of parton-based Gribov-Regge theory
 - a QCD inspired effective-field theory describing the hard and soft scattering simultaneously
- **QGSJET-II:**
 - Provides a phenomenological treatment of hadronic and nuclear interactions in the Reggeon field theory framework

I. Mean charged particle density as a function of η and p_T



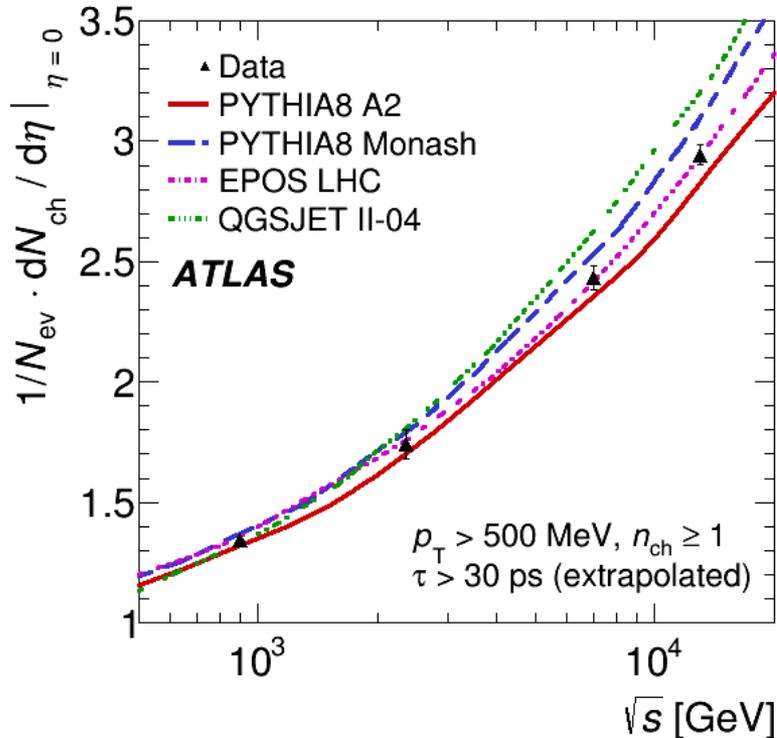
- **The mean particle density is approximately constant at 2.9 for $|\eta| < 1$ and decreases at higher $|\eta|$**
 - EPOS describes the data well for $|\eta| < 1$ but predicts slightly larger multiplicities for larger $|\eta|$
 - QGSJET-II and PYTHIA8 Monash predict too large multiplicities (by 15% and 5%)
 - PYTHIA8 A2 – by about 3% too low in the central region but is OK in the forward region
- **The distribution on p_T decreases by about 9 orders of value**
 - EPOS describes the data well over the entire p_T spectrum
 - The PYTHIA 8 tunes describe the data well but are slightly above at large p_T
 - QGSJET-II gives a poor prediction over the entire spectrum

I. The charged particle multiplicity distribution and $\langle p_T \rangle$ vs n_{ch}



- **The charged particle multiplicity distribution decreases by 5 orders. The high n_{ch} region has significant contribution from events with numerous MPI**
 - All the models describes the data satisfactory in the region $n_{ch} < 30-50$ but fail at higher n_{ch}
 - At high n_{ch} the strongest deviations are for QGSJET-II and PYTHIA8 A2
- **$\langle p_T \rangle$ rises with n_{ch} from 0.8 to 1.2 GeV. The increase is modelled by a color reconnection mechanism in PYTHIA8 or by the hydrodynamical evolution model used in EPOS**
 - EPOS predicts slightly lower $\langle p_T \rangle$, but describes the dependence on n_{ch} well
 - The PYTHIA8 tunes predict a steeper rise of $\langle p_T \rangle$ with n_{ch} than the data
 - QGSJET-II predicts $\langle p_T \rangle$ of ~ 1 GeV, with very little dependence on n_{ch}

I. Energy dependence of the charged particles density at $\eta = 0$

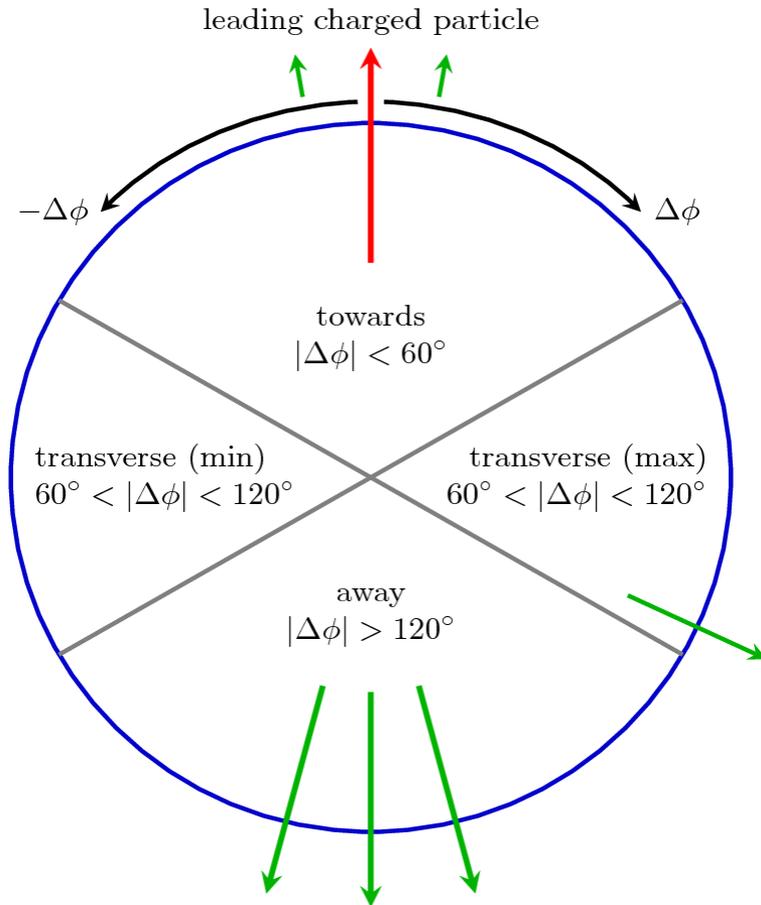


- The mean number of primary charged particles in the central region is $2.874 \pm 0.001 \text{ (stat.)} \pm 0.033 \text{ (syst.)}$
- The value increases by a factor of 2.2 when \sqrt{s} increases by a factor of about 14 from 0.9 TeV to 13 TeV
- EPOS and PYTHIA8 A2 describe the dependence very well
- PYTHIA8 Monash and QGSJET-II predict a steeper rise in multiplicity with c.m.s. energy

Conclusions

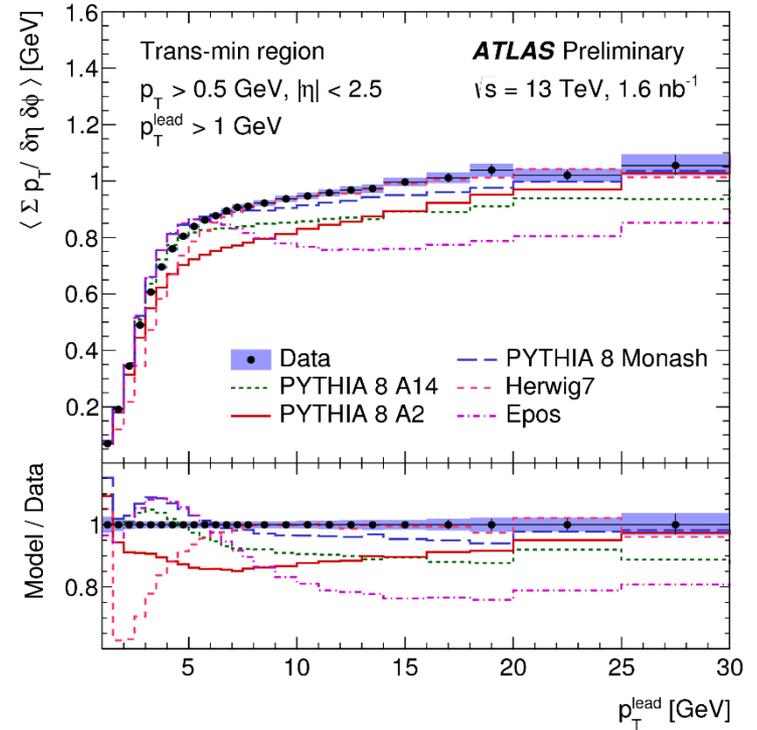
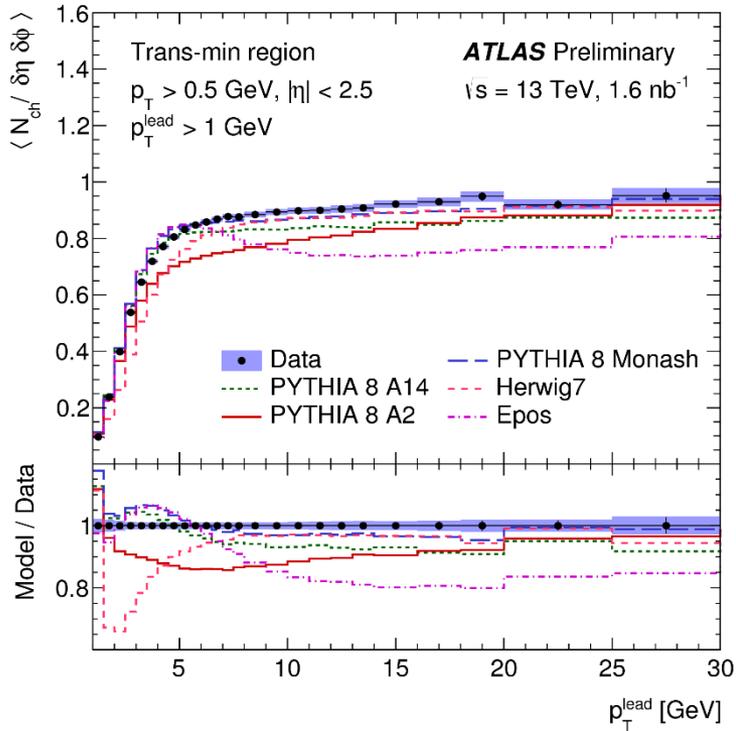
- Primary-charged-particle multiplicity measurements with the ATLAS detector using pp interactions at $\sqrt{s} = 13 \text{ TeV}$ are presented
- The results highlight clear difference between MC models and the measured distributions
- EPOS reproduces data the best, PYTHIA8 A2 and Monash give reasonable descriptions of the data and QGSJET-II provides the worst description

II. Underlying event with leading p_T charged particle



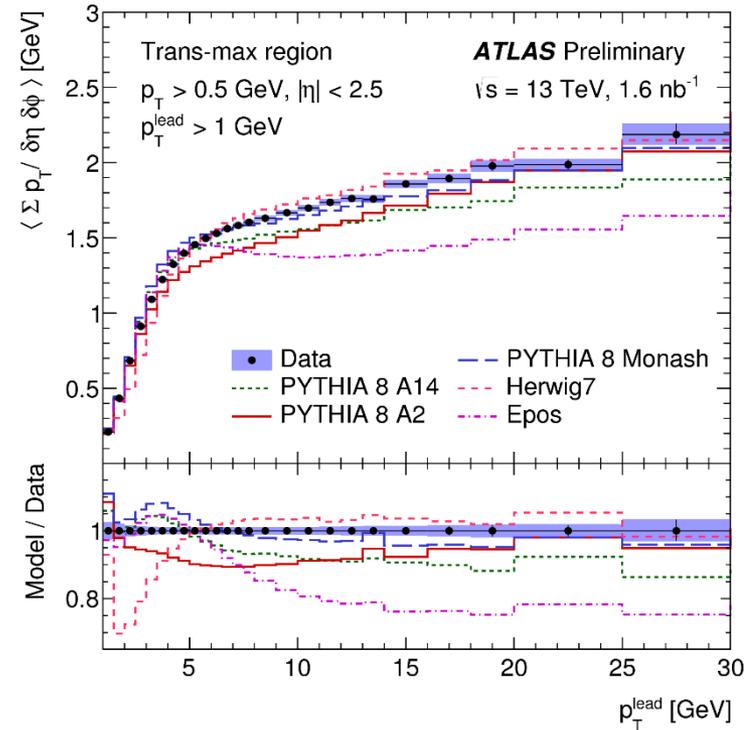
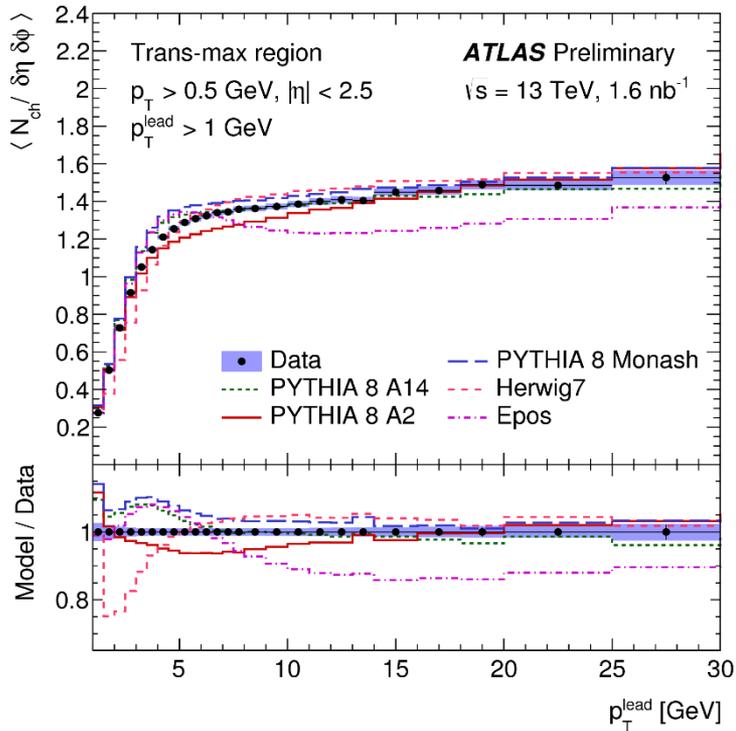
- Underlying event is defined as activity accompanying any hard scattering in a event
- It includes
 - Partons not participating in a hard scattering (beam remnants)
 - Multiple parton interactions
 - Initial state gluon radiation (ISR)
- It is impossible to separate the UE from the hard scattering process on an event by event base
- However distributions have been measured which are sensitive to the properties of the UE
- These are the distributions of particles in the transverse regions (Figure)
- The figure illustrates distribution of charged particles in the transverse to the beam plane with respect to the leading p_T charged particle
- The towards and away regions contain products of the hard scatter mostly

II. Trans-min region



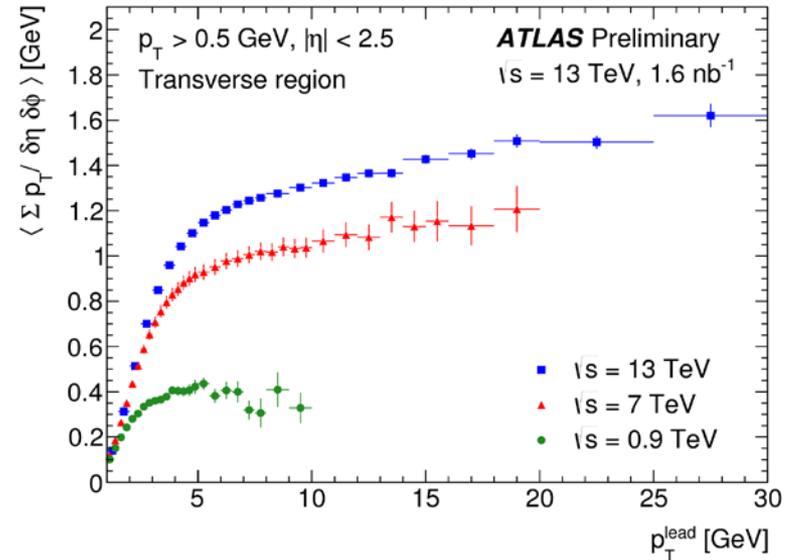
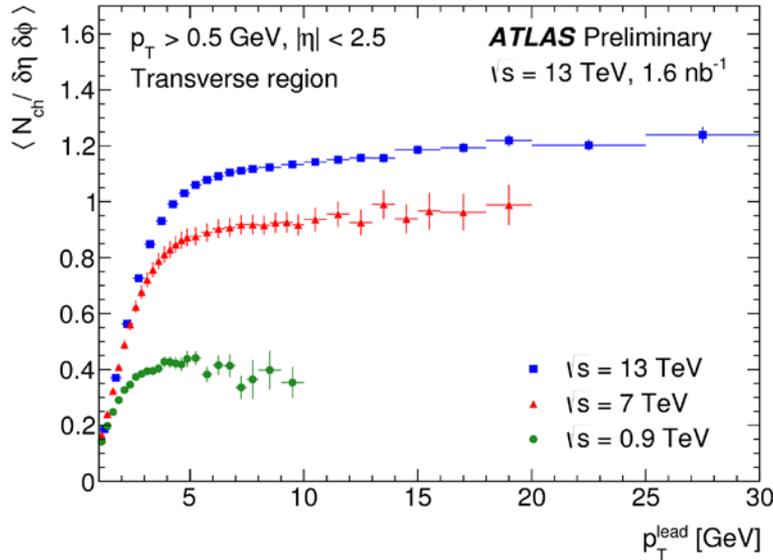
- **Trans-min region most sensitive to MPI and hard process contamination is small: the distributions are rather flat above $p_T = 5$ GeV**
- **EPOS and PYTHIA8 A2 are too low at the plateau region**
- **Herwig7 gives one of the best descriptions at the plateau but severe undershoots transition region below $p_T = 5$ GeV**

II. Trans-max region



- **The trans-max region has larger hard process contamination: the rise in the plateau region is stronger and its mean level is noticeably larger**
- **In general MC generators describes the data better than for the trans-min**
- **EPOS strongly undershoots data at $p_T > 10$ GeV**
- **Herwig7 again undershoots $p_T < 5$ GeV**

II. Conclusions



- A number of distributions sensitive to properties of the UE are presented
- Present measurement completes previous ATLAS measurements at lower energies (above figures)
- An increase of UE activity of approximately 20% is observed when going from 7 TeV to 13 TeV pp collisions
- Comparison against prediction from several MC generator tunes indicate that for most observables the models describe UE data to better than 5% accuracy
- EPOS gives the worst predictions
- The data can be used to improve MC tunes

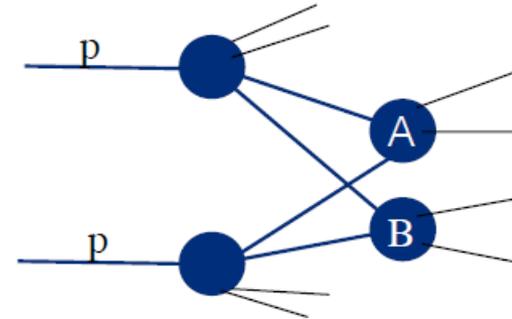
III. Hard double parton scattering in 4-jet events

$$\hat{\sigma}_{(A,B)}^{\text{DPS}} = \frac{1}{1+\delta_{AB}} \frac{\hat{\sigma}_A \hat{\sigma}_B}{\sigma_{\text{eff}}},$$

$$\Rightarrow \sigma_{\text{eff}} = \frac{1}{1+\delta_{AB}} \frac{\hat{\sigma}_A \hat{\sigma}_B}{f_{\text{DPS}} \cdot \hat{\sigma}_{(A,B)}^{\text{tot}}}$$

Kroneker $\delta = 1$, if A and B the same

fraction of DPS events



$$\sigma_{\text{eff}} = \frac{1}{1 + \delta_{AB}} \frac{1}{f_{\text{DPS}}} \frac{\sigma_{2j}^A \sigma_{2j}^B}{\sigma_{4j}}$$

- σ_{eff} is an effective cross section (to be found)
 - phenomenological parameter describing the effective overlap between the interacting hadrons
 - determines the overall size of DPS cross section
 - Assumed to be process and cut independent
- Data taken in pp collisions at $\sqrt{s} = 7$ TeV during 2010, $\langle \mu \rangle = 0.41$, integrated luminosity 37.3 pb^{-1}

III. Event selection, templates

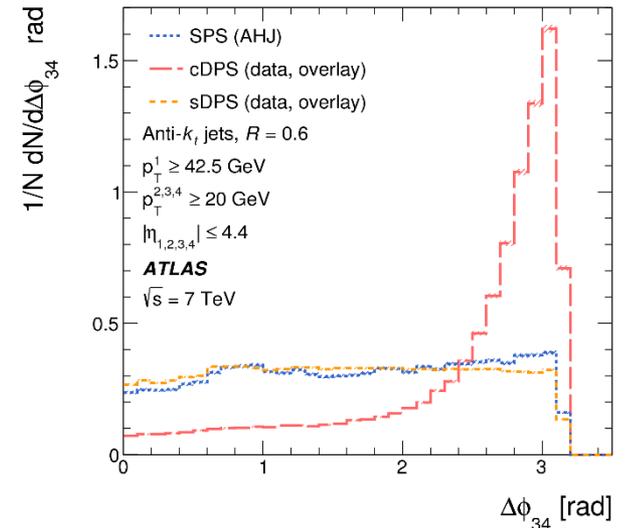
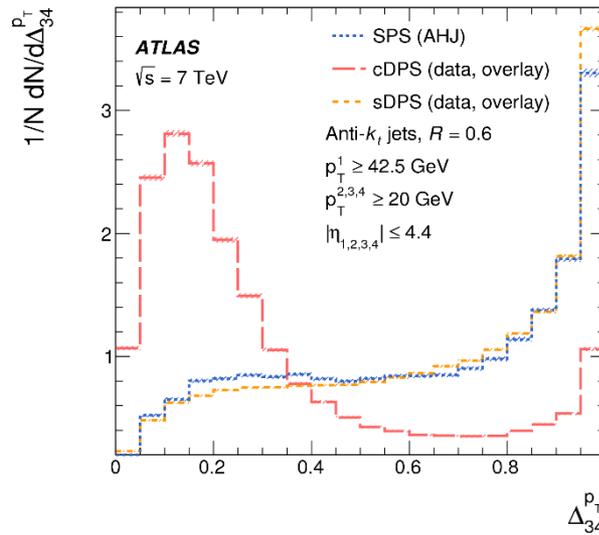
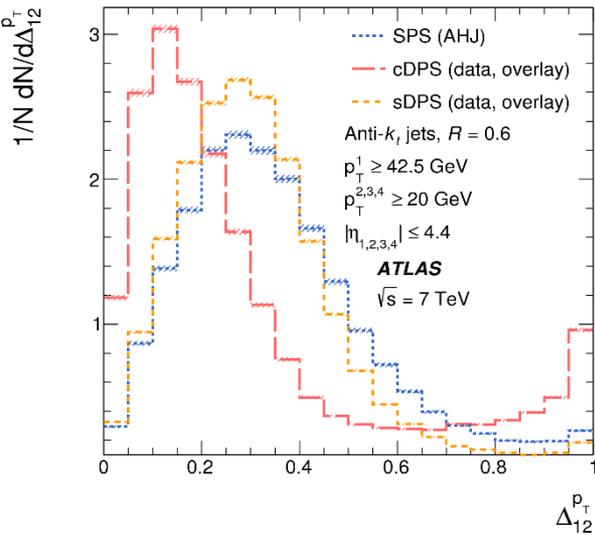
- **Single vertex events only**
- **Use anti- k_T jets with $R = 0.6$ found in calorimeters at $|\eta| < 4.4$**
- **4-jet event selection: $N_{\text{jet}} \geq 4$, $p_T^1 > 42.5$ GeV, $p_T^{2,3,4} > 20$ GeV**
- **Di-jet samples $N_{\text{jet}} \geq 2$, A: $p_T^{1,2} > 20$ GeV; B: $p_T^1 > 42.5$ GeV, $p_T^2 > 20$ GeV**
- **4-jet and B samples are subsamples of sample A**
- **AHJ MC: Alpgen ($2 \rightarrow n$, n up to 5) + Jimmi (MPI) + Herwig (hadronization)**
- **Event record of AHJ is used to distinguish between ($2 \rightarrow 4$) and (two $2 \rightarrow 2$) processes**
- **Sample SPS: all 4 jets from one hard scatter, based on AHJ MC**
- **Sample complete DPS (cDPS): 2 jets from 1st hard scatter, 2 jets from 2nd one. Obtained by overlaying two di-jet data events**
- **Sample semi DPS (sDPS): 3 jets from 1st hard scatter, 1 jet from 2nd one. Obtained by overlaying two di-jet data events**
- **Template fit is used to determine $f_{\text{DPS}} = f_{\text{cDPS}} + f_{\text{sDPS}}$**

III. Discriminating variables

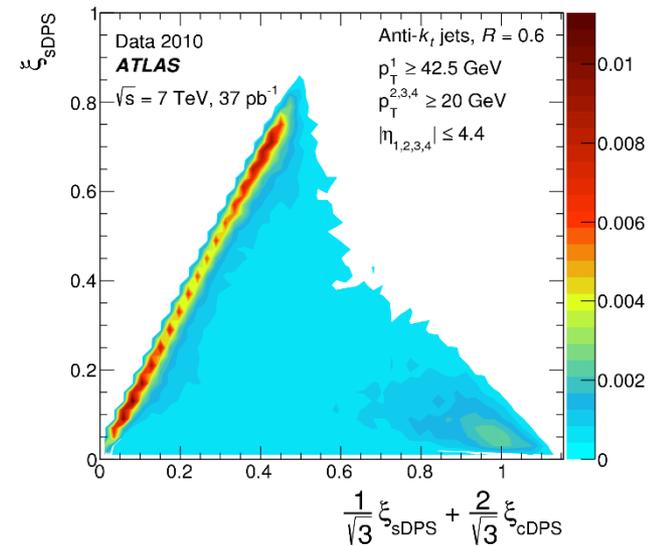
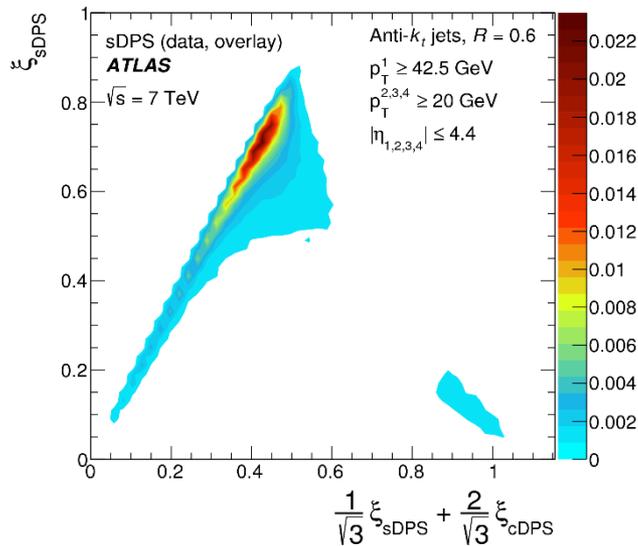
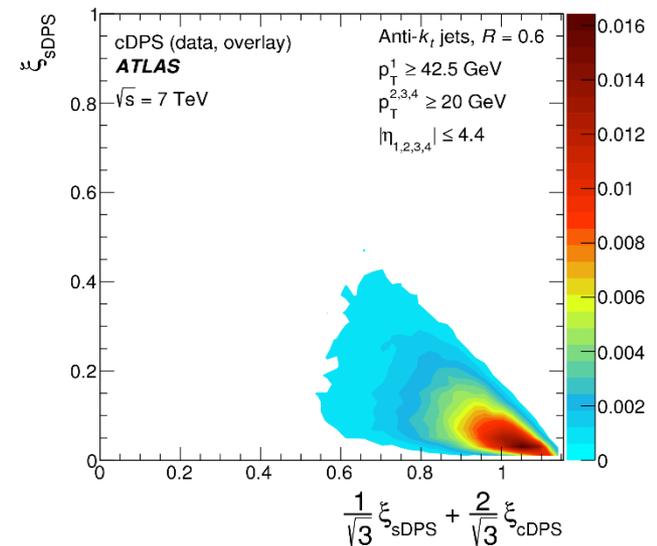
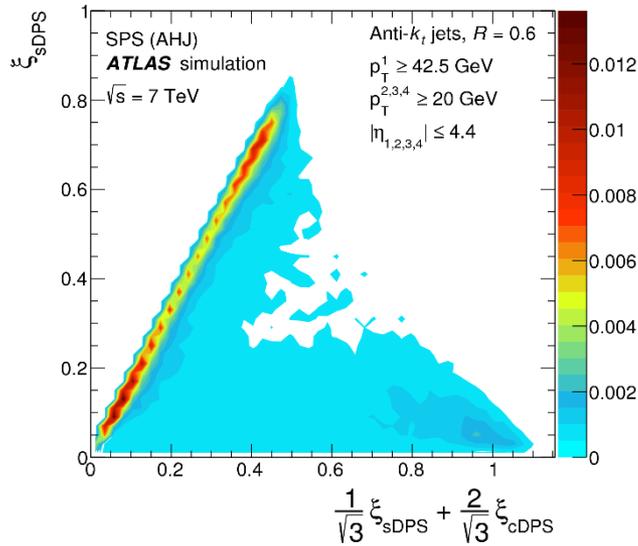
- 21 discriminating variables are used to classify events as belonging to SPS (ξ_{SPS}), cDPS (ξ_{cDPS}), or sDPS (ξ_{sDPS}): $\xi_{\text{SPS}} + \xi_{\text{cDPS}} + \xi_{\text{sDPS}} = 1$
- Neural network is used for event classification

$$\Delta_{ij}^{p_T} = \frac{|\vec{p}_T^i + \vec{p}_T^j|}{p_T^i + p_T^j}; \quad \Delta\phi_{ij} = |\phi_i - \phi_j|; \quad \Delta y_{ij} = |y_i - y_j|;$$

$$|\phi_{1+2} - \phi_{3+4}|; \quad |\phi_{1+3} - \phi_{2+4}|; \quad |\phi_{1+4} - \phi_{2+3}|;$$



III. NN output for test samples



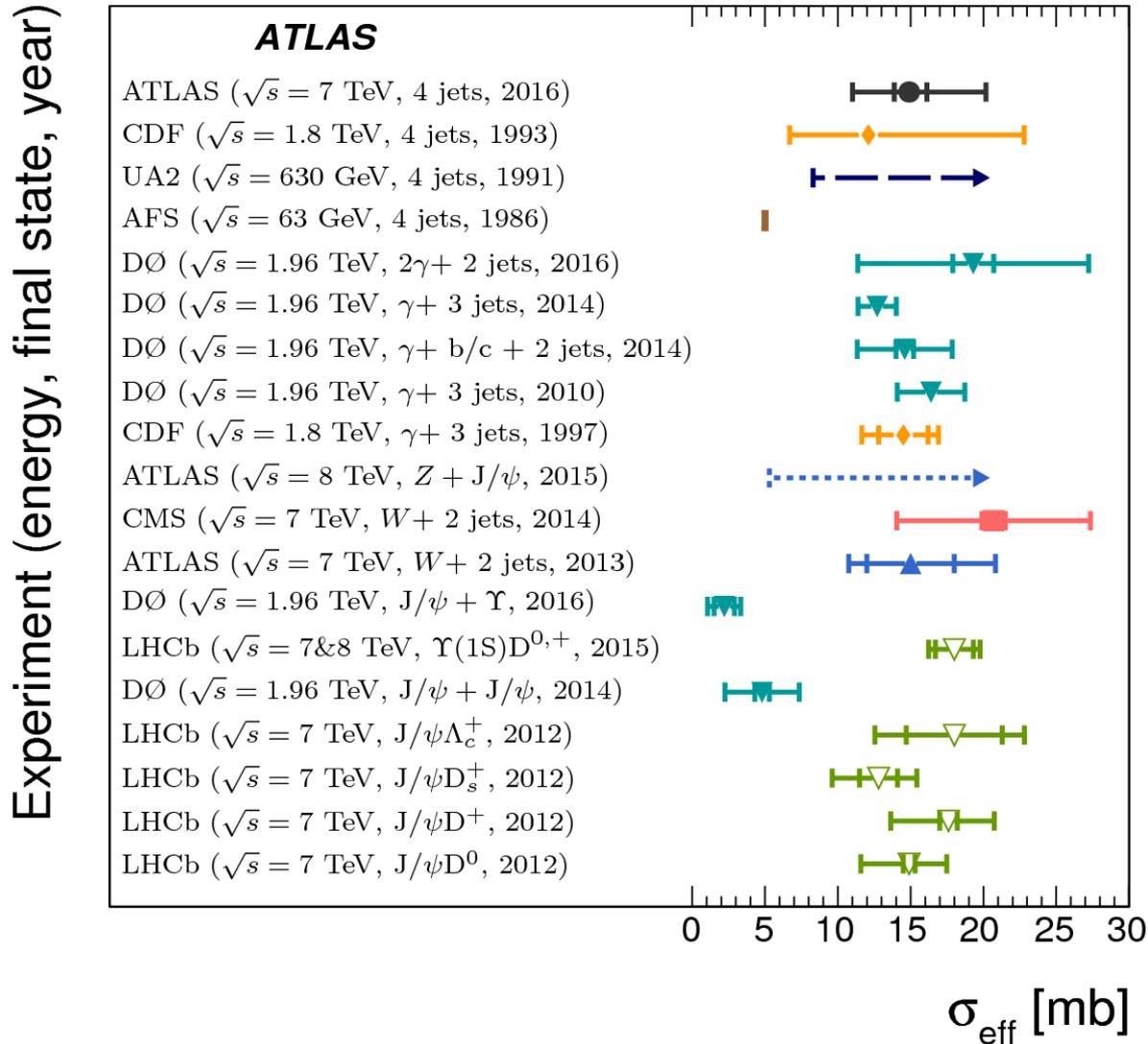
III. Results

- χ^2 minimization with Minuit to find f_{cDPS} and f_{sDPS}
- $\mathbf{D} = (1 - f_{\text{cDPS}} - f_{\text{sDPS}}) * \mathbf{M}_{\text{SPS}} + f_{\text{cDPS}} * \mathbf{M}_{\text{cDPS}} + f_{\text{sDPS}} * \mathbf{M}_{\text{sDPS}}$
- Here \mathbf{D} is ternary data distribution
- \mathbf{M}_{SPS} , \mathbf{M}_{cDPS} and \mathbf{M}_{sDPS} are ternary distributions in test samples normalized to measured 4-jet cross section
- Different sources of systematic uncertainties were taken in to account and propagated to the final results
- The dominating source is jet energy scale (JES) uncertainty: about 4.5% in the central region and rising to about 10% in the forward one. It gives about 30% uncertainty in the final results

$$f_{\text{DPS}} = 0.092^{+0.005}_{-0.011} \text{ (stat.) }^{+0.033}_{-0.037} \text{ (syst.)},$$

$$\sigma_{\text{eff}} = 14.9^{+1.2}_{-1.0} \text{ (stat.) }^{+5.1}_{-3.8} \text{ (syst.) mb}$$

III. Results overview



IV. Total and elastic cross sections

- **Optical theorem relates total pp cross section to the elastic-scattering amplitude extrapolated to the forward direction**

$$\sigma_{tot} = 4\pi * \text{Im}(f_{el})_{t \rightarrow 0}$$

- **ATLAS uses optical theorem and luminosity-dependent method of the total cross section extraction. With this method**

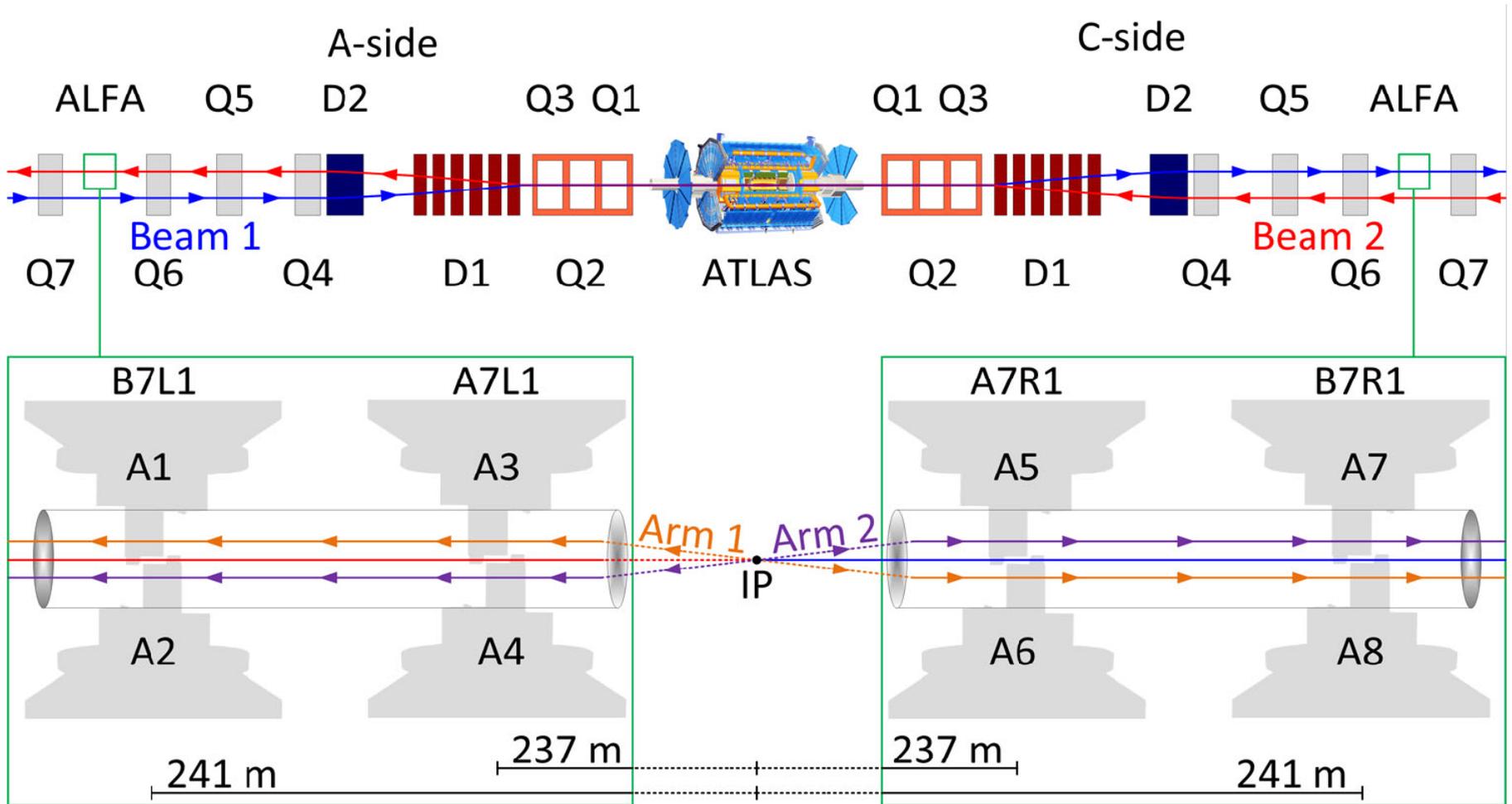
$$\sigma_{tot}^2 = \frac{16\pi(hc)^2}{1 + \rho^2} \left. \frac{d\sigma_{el}}{dt} \right|_{t \rightarrow 0}$$

- **Where $\rho = 0.1362$ represents a small correction arising from the ratio of the real to the imaginary part of the elastic-scattering amplitude**

IV. ALFA detector, data taking

- ALFA detector is used to record elastic-scattering data
- Consists of Roman Pot (RP) tracking detector stations placed at 237 m (inner) and 241 m (outer) on either side of the ATLAS IP
- Each station houses two scintillating fibre detectors with a spatial resolution of about 35 μm
- The detectors are supplemented with trigger counters consisting of plain scintillator tiles
- The data recorded in a single low luminosity run with special high β^* optics for pp interactions at $\sqrt{s} = 8 \text{ TeV}$
- 3.8 M selected elastic events
- Measure elastic track positions at ALFA to get the scattering angle Θ and thereby the t-spectrum $d\sigma/dt$
- To calculate Θ from the measured tracks the transport matrix elements of the beam optics are used

IV. ALFA detector



IV. Analysis of elastic data

- **Data-driven method to calculate the reconstruction efficiency of about 90%**
- **Tuning of the beam optics model with ALFA constrains → effective optics**
- **Trigger efficiency is very high ~99.9%; determined from data stream with looser conditions**
- **Dedicated luminosity determination resulting in a small uncertainty of 1.5%**

$$\left(\frac{d\sigma}{dt}\right)_i = \frac{1}{t_i} \cdot \frac{M^{-1}[N_i - B_i]}{A_i \cdot \epsilon^{reco} \cdot \epsilon^{trig} \cdot \epsilon^{DAQ} \cdot L_{int}}$$

- **Δt_i : the width of the bins in t**
- **A : acceptance(t)**
- **M^{-1} : symbolizes the unfolding procedure applied to the background subtracted number of events $N_i - B_i$**
- **ϵ^{reco} : event reconstruction efficiency**
- **ϵ^{trig} : trigger efficiency**
- **ϵ^{DAQ} : the dead time correction**
- **L_{int} : the integrated luminosity**

IV. Fitting formula for $d\sigma/dt$

- The theoretical prediction used to fit the elastic data consists of the Coulomb term, the Coulomb-Nuclear-Interference term and the dominant Nuclear term

$$\frac{d\sigma}{dt} = \frac{4\pi\alpha^2(\hbar c)^2}{|t|^2} \cdot G^4(t) \quad \text{Coulomb}$$

$$\text{CNI} \quad - \quad \sigma_{\text{tot}} \cdot \frac{\alpha G^2(t)}{|t|} [\sin(\alpha\phi(t)) + \rho \cos(\alpha\phi(t))] \cdot \exp\left(\frac{-B|t|}{2}\right)$$

$$\text{Nuc.} \quad + \quad \sigma_{\text{tot}}^2 \frac{1 + \rho^2}{16\pi(\hbar c)^2} \cdot \exp(-B|t|)$$

ρ	0.1362
Λ	0.71 GeV ²
ϕ_C	0.577

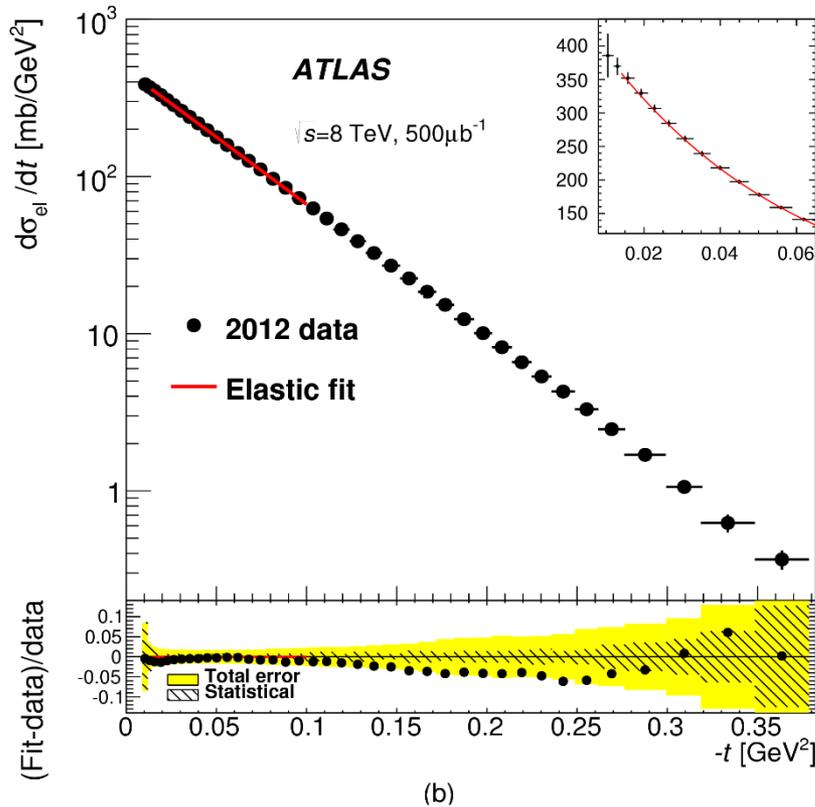
$$G(t) = \left(\frac{\Lambda}{\Lambda + |t|} \right)^2,$$

Proton dipole form factor

$$\phi(t) = -\ln \frac{B|t|}{2} - \phi_C,$$

Coulomb phase

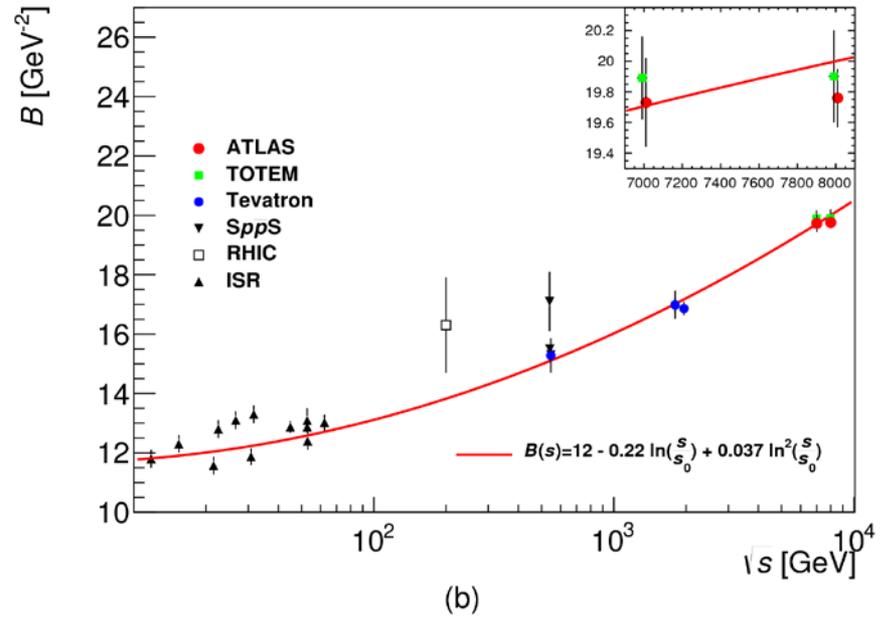
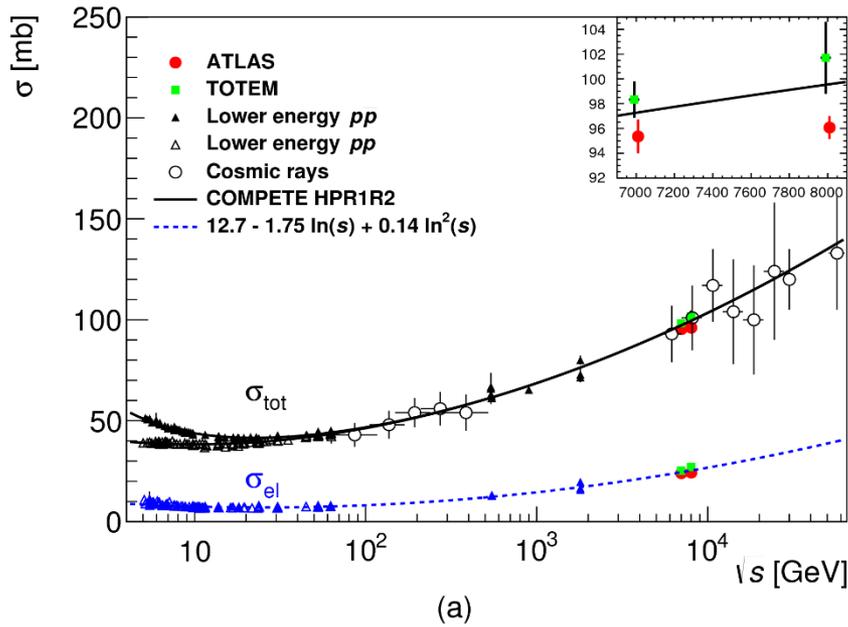
Fit results



- The fit includes experimental systematic uncertainties in the χ^2
- Main systematics
 - t-independent: luminosity $\pm 1.5\%$
 - t-dependent: beam energy $\pm 0.65\%$
- The fit range is set to $-t[0.014, 0.1] \text{ GeV}^2$, where possible deviations from exponential form of the nuclear amplitude are expected to be small
- The extrapolation uncertainty is evaluated by a variation of the fit range

- $\sigma_{\text{tot}} = 96.07 \pm 0.18(\text{stat.}) \pm 0.85(\text{exp.}) \pm 0.31(\text{extr.})$
- $B = 19.74 \pm 0.05(\text{stat.}) \pm 0.16(\text{exp.}) \pm 0.15(\text{extr.})$

IV. Energy evolution



- Comparison with COMPETE model [Chin.Phys. C, 38, 090001 \(2014\)](#) for the evolution of the total cross section. High accuracy of the ATLAS data due to precise luminosity measurement

- Comparison with a model from Schegelsky and Ryskin [Phys.Rev. D85, 094024 \(2012\)](#) for the evolution of the nuclear slope

Summary

- **Primary-charged-particle multiplicity measurements with the ATLAS detector using pp interactions at $\sqrt{s} = 13$ TeV were done**
- **The results highlight clear difference between MC models and the measured distributions**
- **EPOS reproduces data the best, PYTHIA8 A2 and Monash give reasonable descriptions of the data and QGSJET-II provides the worst description**
- **A number of distributions sensitive to properties of the UE are presented for pp interactions at $\sqrt{s} = 13$ TeV**
- **EPOS fails to describe the data properly contrary to the minimum bias event description, for which EPOS is the best**
- **Comparison against prediction from several other MC generator tunes indicate that for most observables the models describe UE data to better than 5% accuracy**
- **Hard double parton scattering in 4-jet events is investigated in pp interactions at $\sqrt{s} = 7$ TeV**
- **The phenomenological DPI parameter σ_{eff} is extracted**
- **Distribution $d\sigma/dt$ for proton-proton elastic scattering at $\sqrt{s} = 8$ TeV is measured and used to estimate $\sigma_{\text{tot}} = 96.07 \pm 0.18(\text{stat.}) \pm 0.85(\text{exp.}) \pm 0.31(\text{extr.})$**