

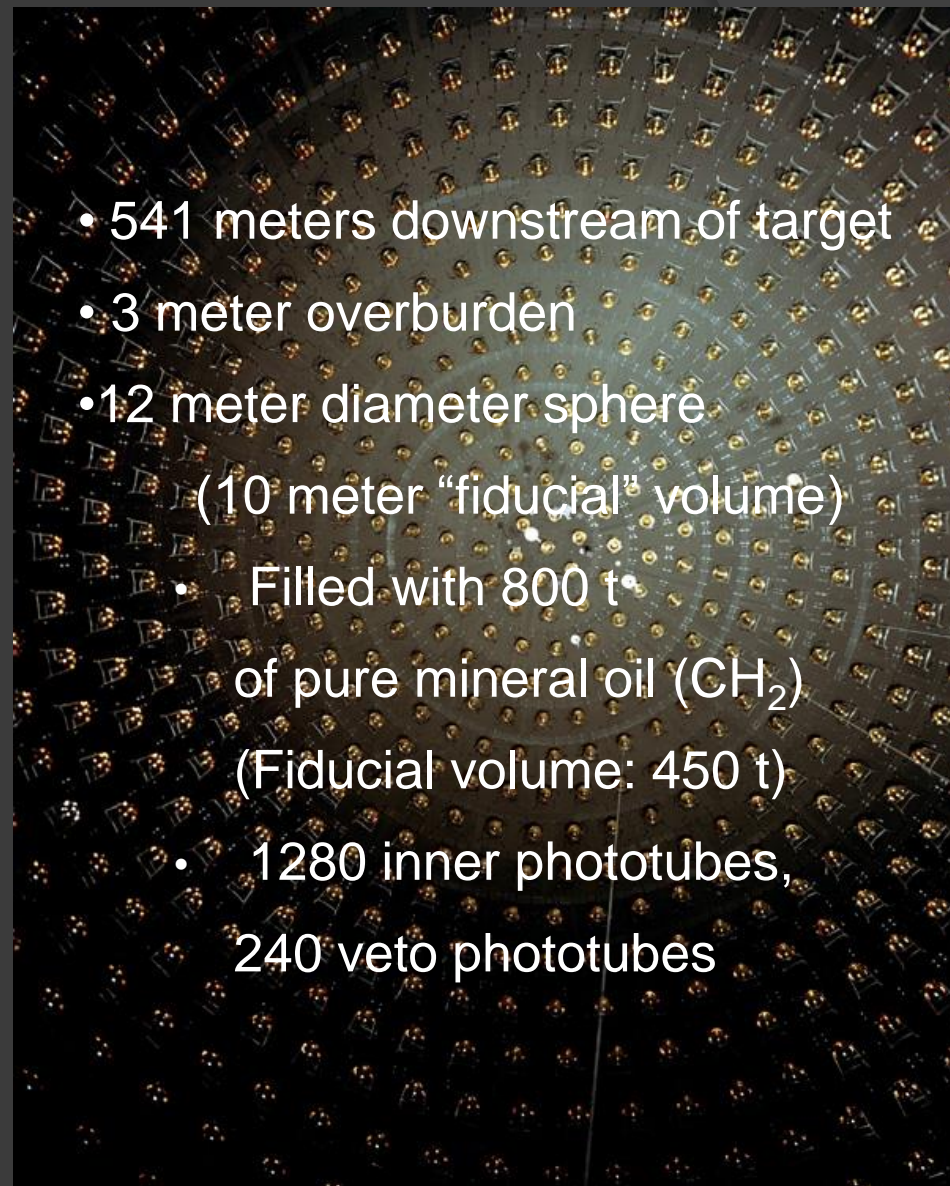
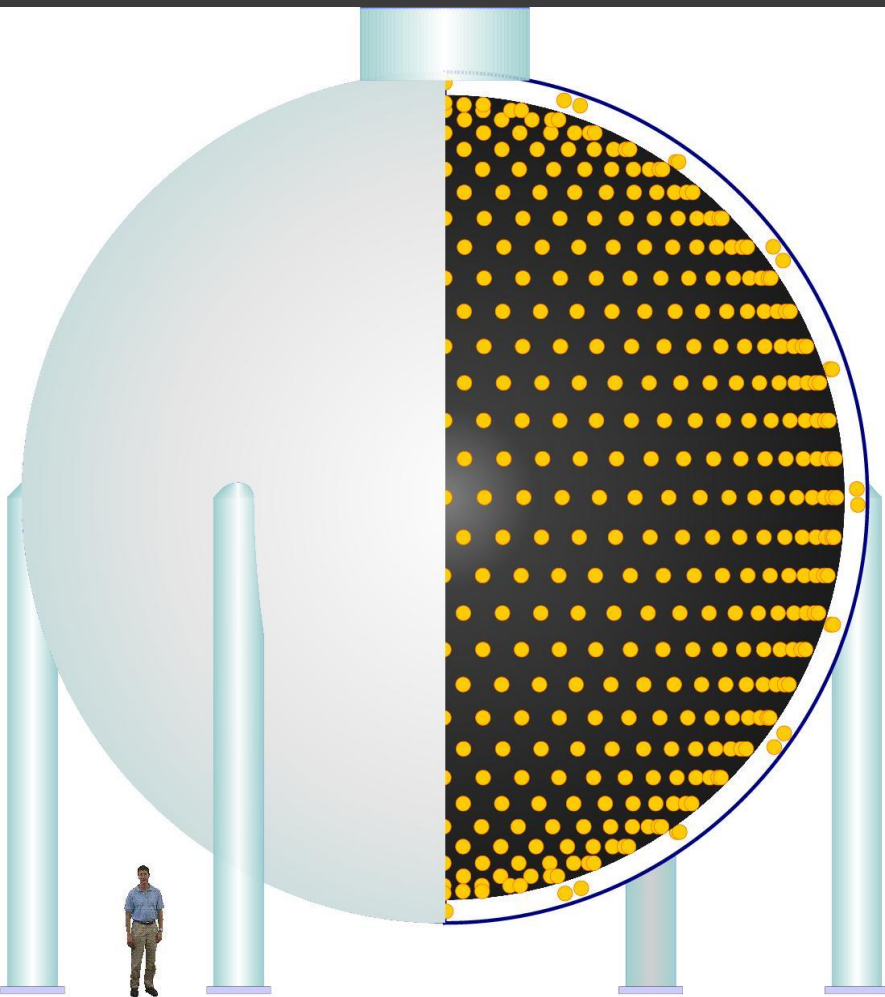
# MiniBooNE Reconstruction Overview

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*Louisiana State University*

FroST-II Workshop, Mainz, Germany  
October 22-24, 2016

# The MiniBooNE Detector



- 541 meters downstream of target
- 3 meter overburden
- 12 meter diameter sphere  
(10 meter “fiducial” volume)
- Filled with 800 t of pure mineral oil ( $\text{CH}_2$ )  
(Fiducial volume: 450 t)
- 1280 inner phototubes,  
240 veto phototubes

*A. A. Aguilar-Arevalo et al., NIM A599, 28 (2009)*

# Subevents

A 19.2  $\mu\text{s}$  beam trigger window

- encompasses the 1.6  $\mu\text{s}$  spill
- starts 4  $\mu\text{s}$  before the beam

Subevent:

Multiple hits within a  $\sim 100$  ns window form “subevents”

Most events are from

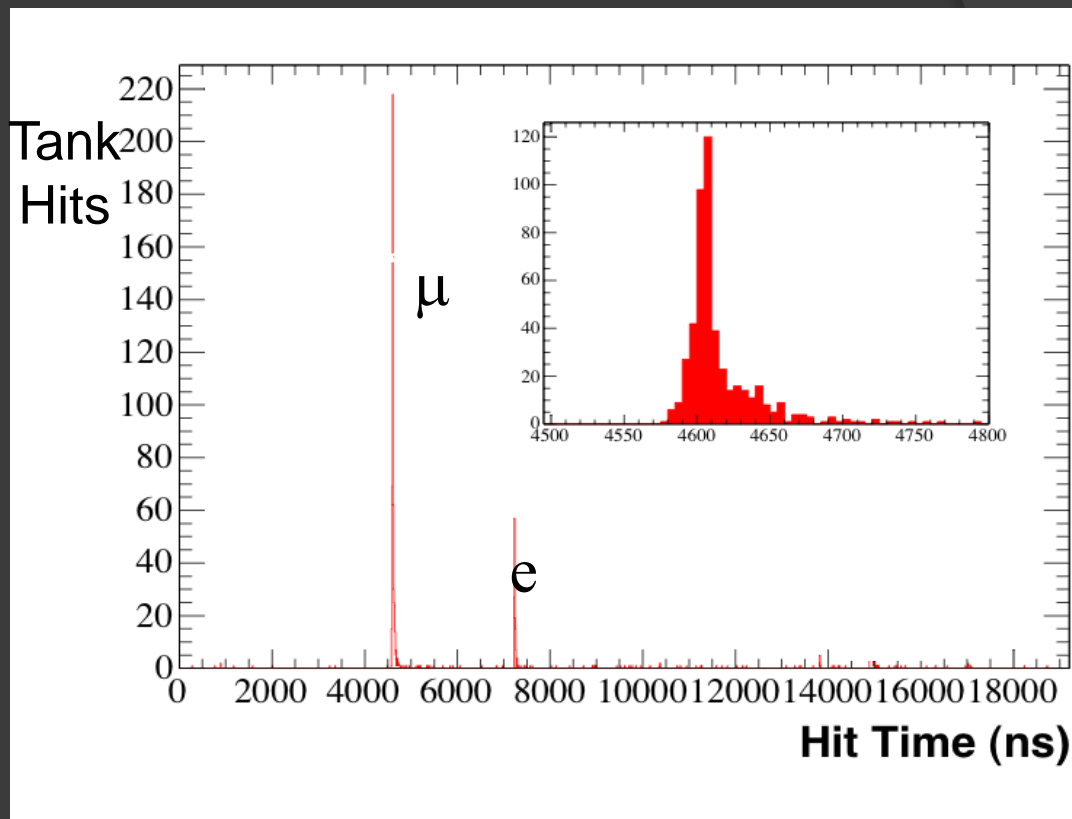
$\nu_\mu$  CC interactions

( $\nu+n \rightarrow \mu+p$ )

with characteristic two

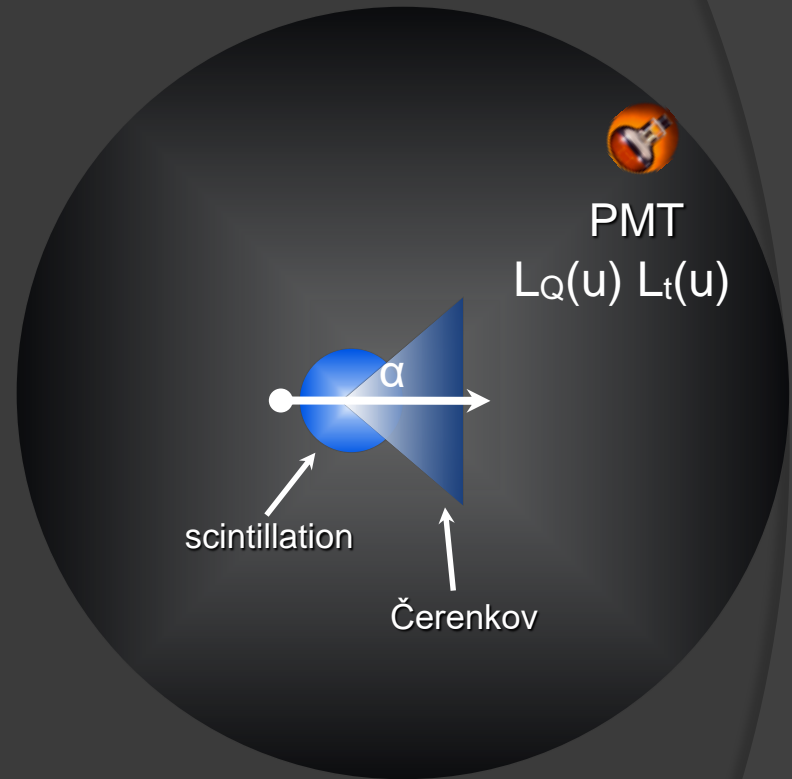
“subevent” structure from

stopped  $\mu \rightarrow \nu_\mu \nu_e e$



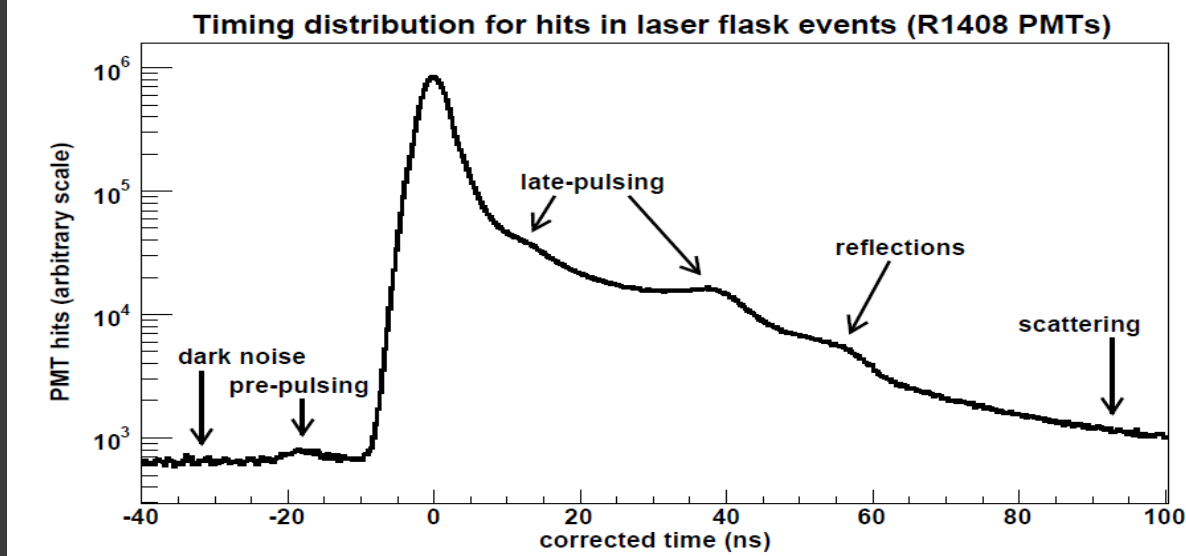
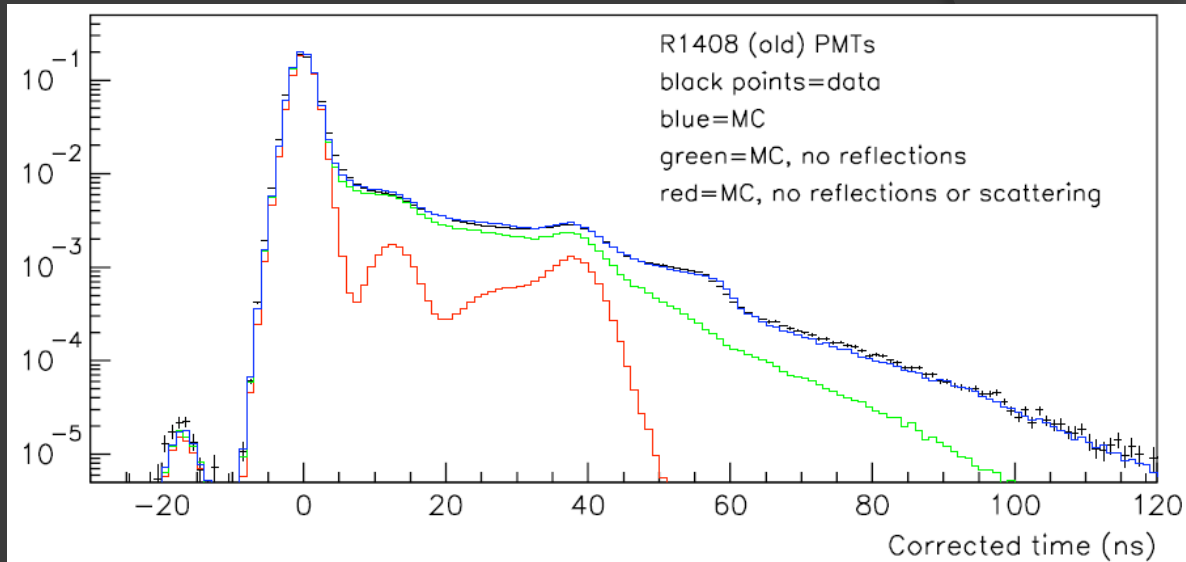
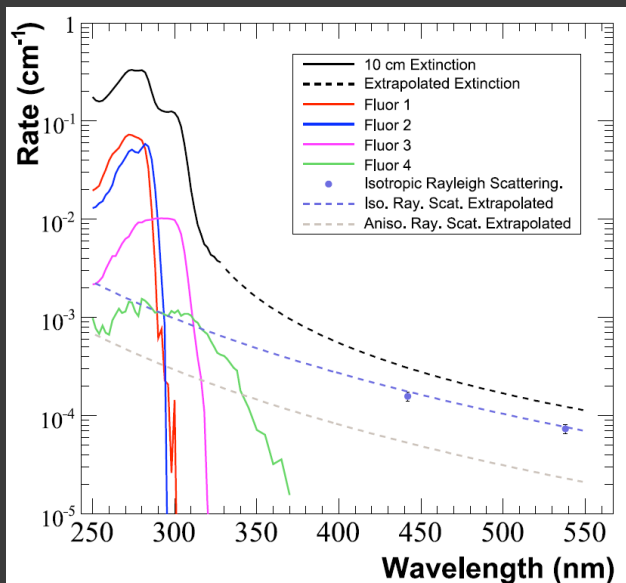
# Track Reconstruction

- A particle is parametrized as a “track” in the oil.
  - Vertex:  $(x,y,z)$
  - Time:  $(t)$
  - Direction:  $(\theta,\varphi)$
  - Kinetic energy:  $(E)$
- At each point of the track scintillation and Čerenkov light is produced. This depends on the type of particle.
- This light propagates through the mineral oil to the PMTs.



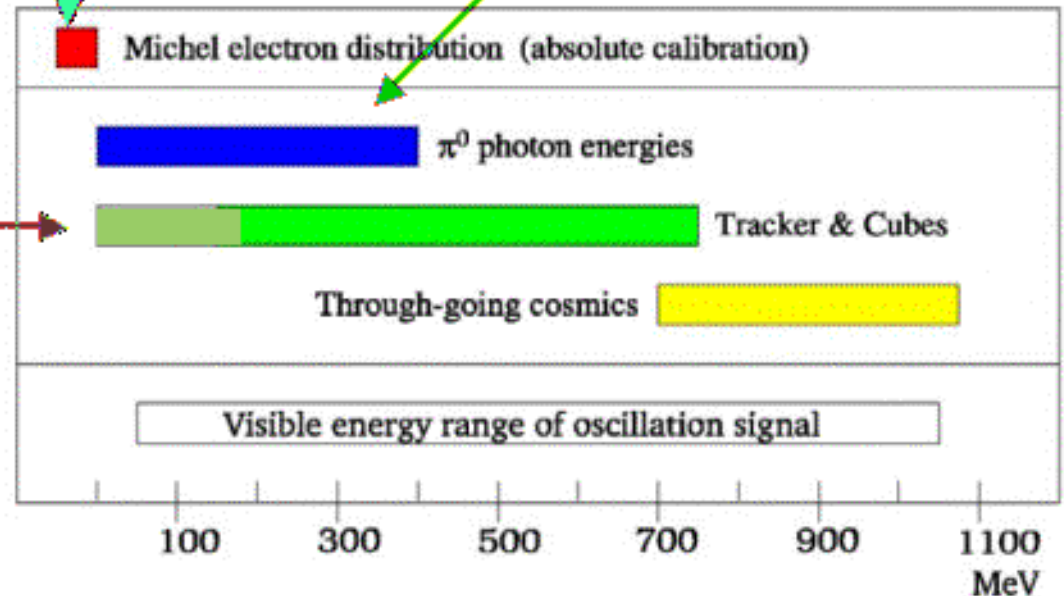
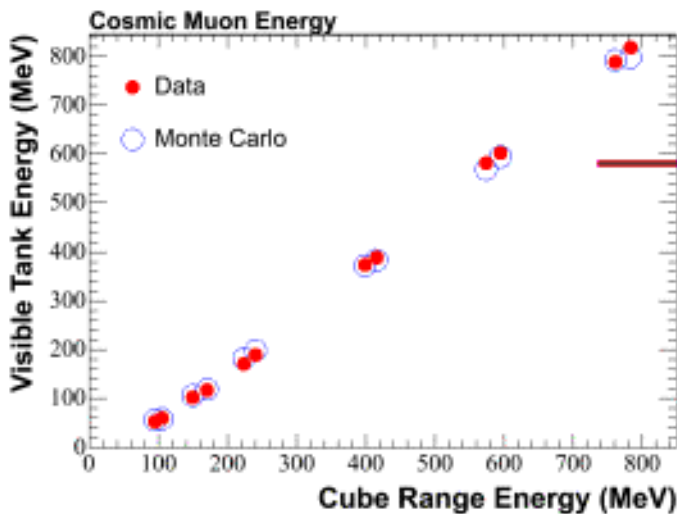
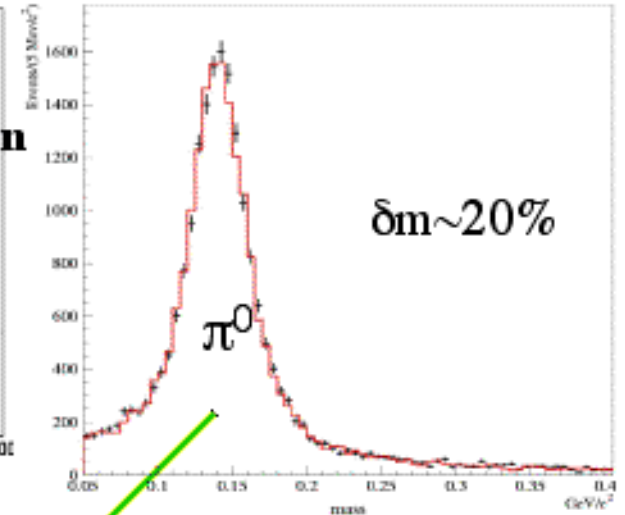
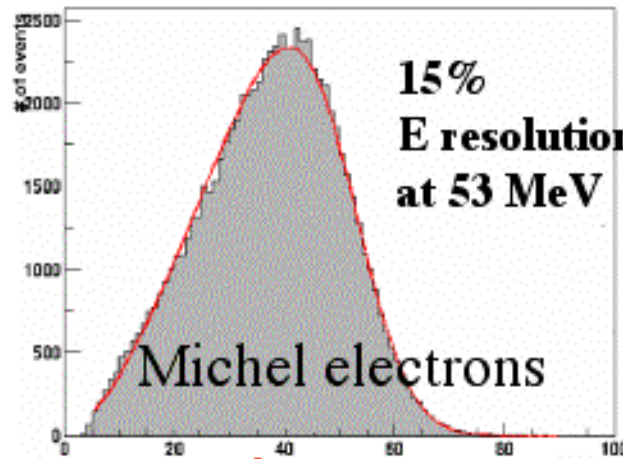
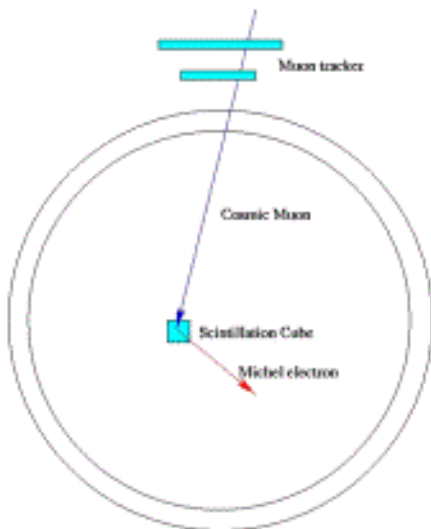
# Optical Model

For the first 2-3 years MB was mainly an experiment in optics.



# Detector Calibration

## Tracker system





# Track Fitting - Likelihood

$$\mathcal{L}(\mathbf{x}) = \prod_{i=1}^{N_{unhit}} \mathcal{P}_i(\text{unhit}; \mathbf{x}) \prod_{j=1}^{N_{hit}} \mathcal{P}_j(\text{hit}; \mathbf{x}) f(q_j; \mathbf{x}) f(t_j; \mathbf{x})$$

$$-\log(\mathcal{L})(\mathbf{x}) = F_q(\mathbf{x}) + F_t(\mathbf{x})$$

$$F_q(\mathbf{x}) = - \sum_{i=1}^{N_{unhit}} \log(\mathcal{P}_i(\text{unhit}; \mathbf{x})) - \sum_{j=1}^{N_{hit}} \log(\mathcal{P}_j(\text{hit}; \mathbf{x}) f(q_j; \mathbf{x})),$$

$$F_t(\mathbf{x}) = - \sum_{j=1}^{N_{hit}} \log(f(t_j; \mathbf{x})).$$

$$\mathcal{P}(\text{hit}; \mu(\mathbf{x})) = 1 - \mathcal{P}(\text{unhit}; \mu(\mathbf{x})) = 1 - e^{-\mu}.$$

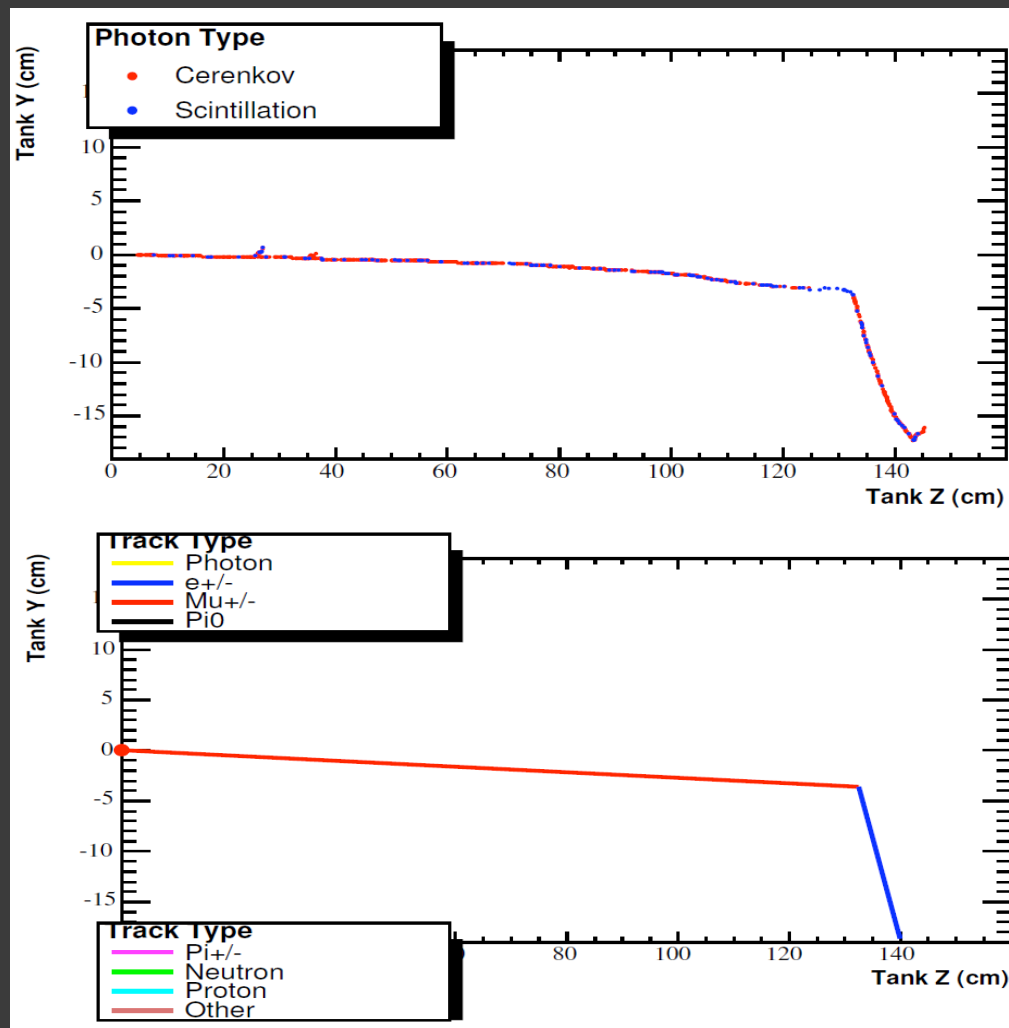
L – likelihood.

$\mathcal{P}_i(\text{unhit}; \mathbf{x})$  – probability a tube i to be unhit given  $\mathbf{x}$ .

$f(q_j; \mathbf{x})$  - charge PDF for PMT j.

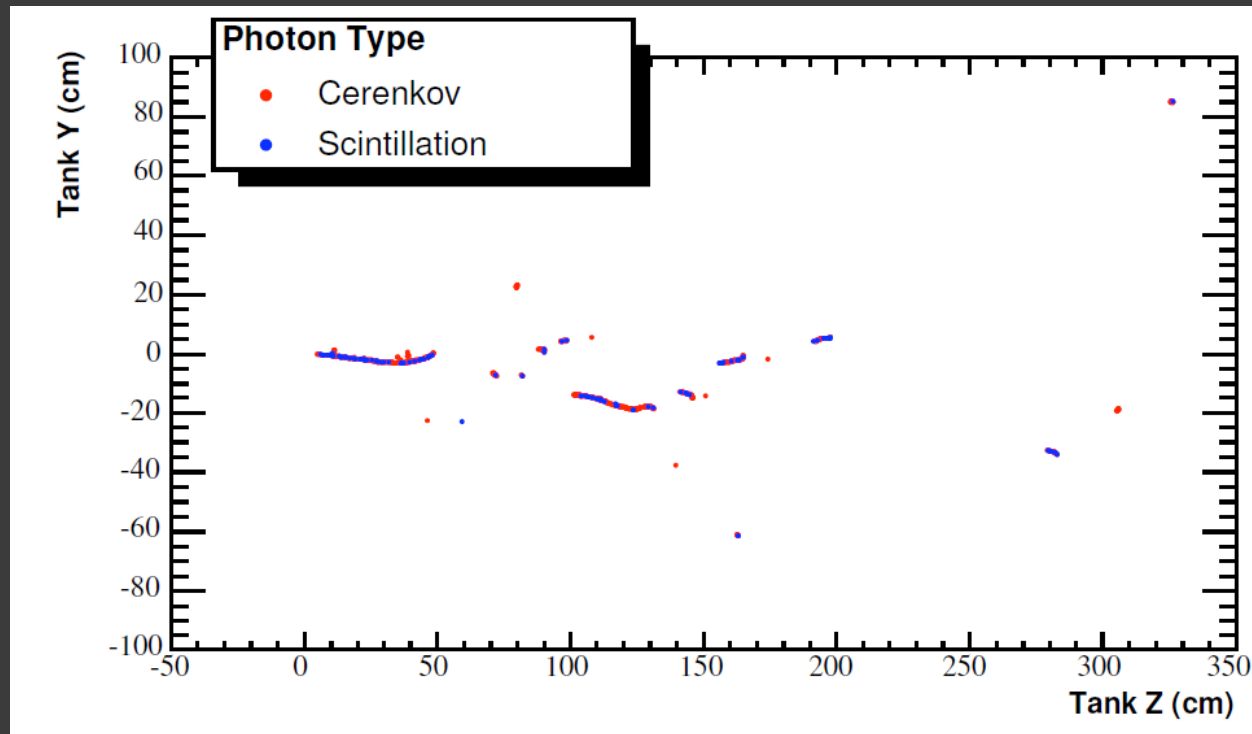
$f(t_j; \mathbf{x})$  - time PDF for PMT j.

# Photon Emission Along the Track - Muon





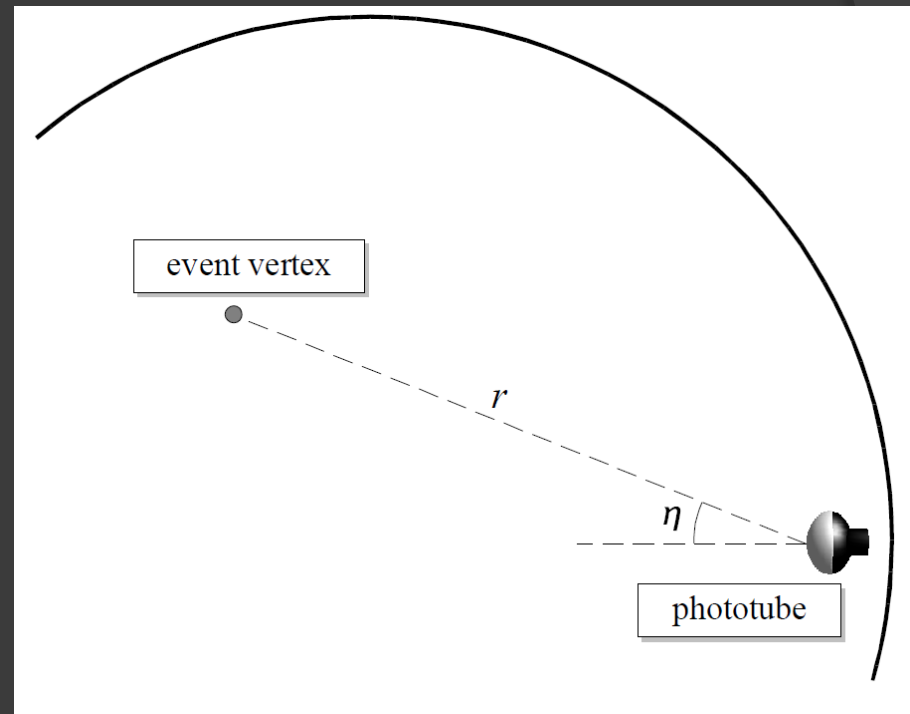
# Photon Emission Along the Track - Electron



# Track Fitting – Predicted Charge Point Source

$$\mu = \Phi \Omega(r) T(r) \epsilon(\eta)$$

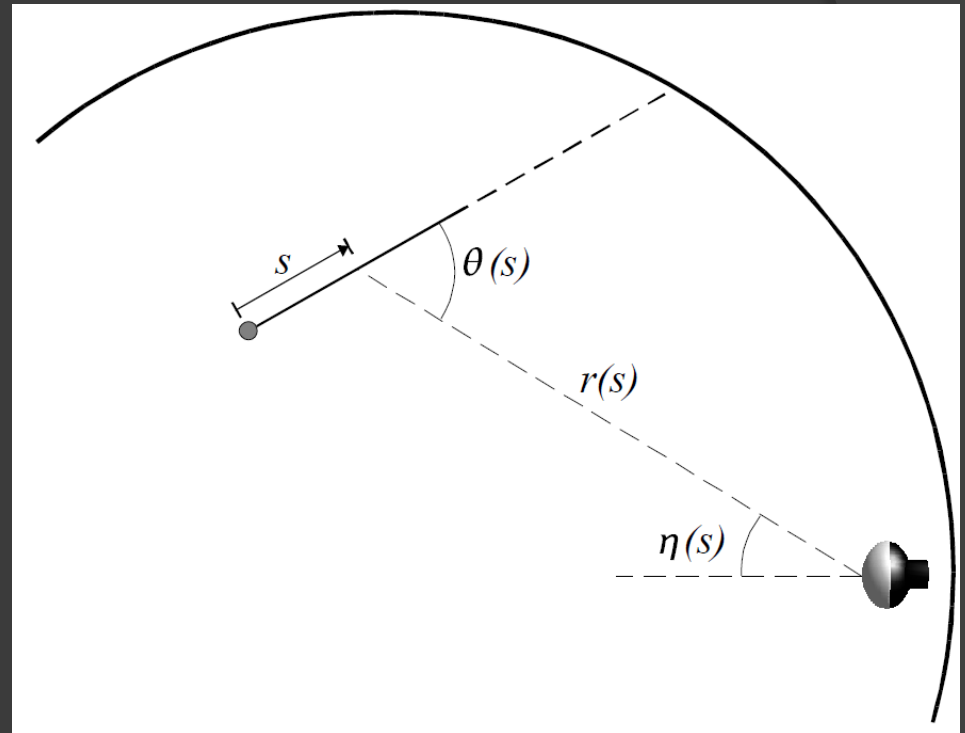
$\mu$  - predicted charge  
 $\Phi$  - light yield  
 $\Omega$  - solid angle PMT  
 $T$  - transition  
 $\epsilon$  - acceptance



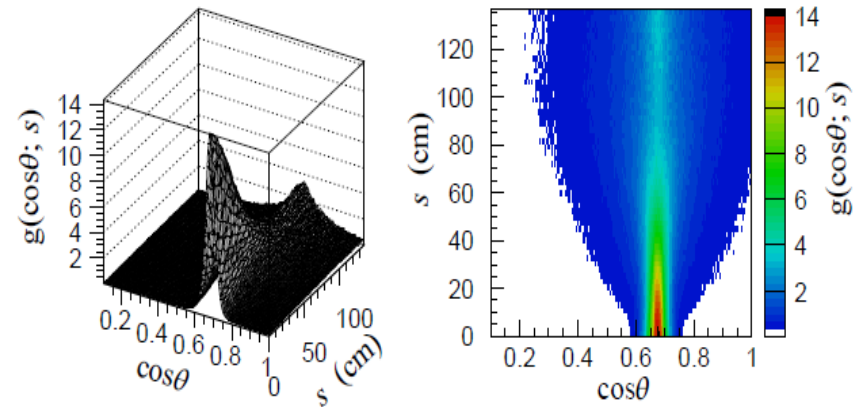
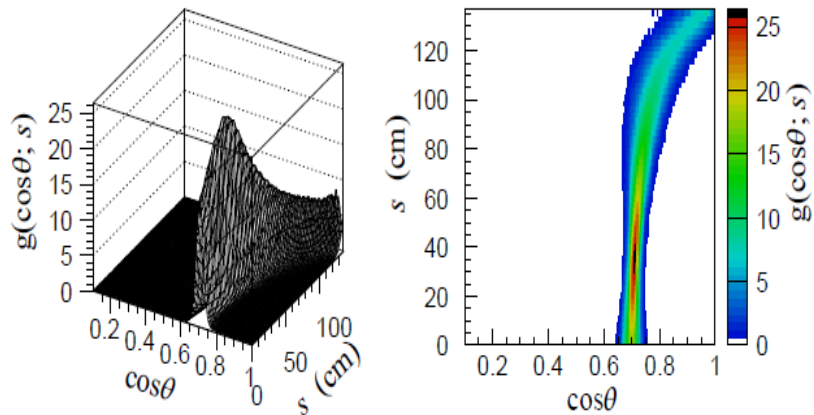
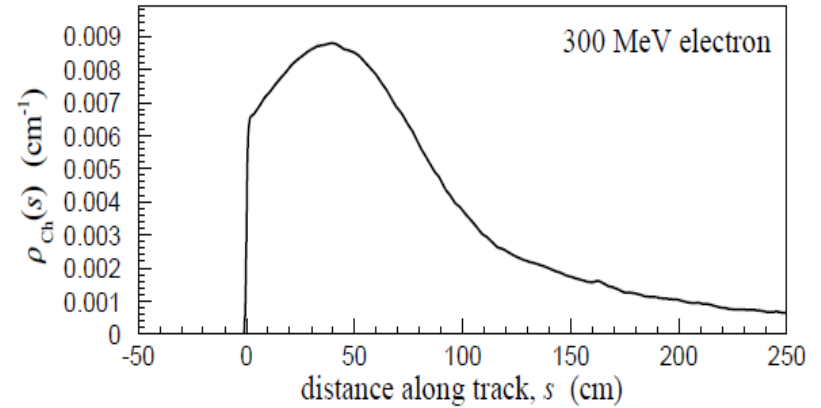
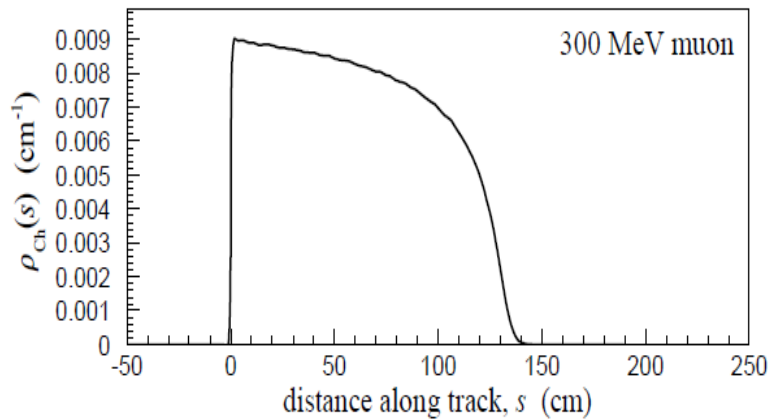
# Track Fitting – Predicted Charge Extended Track Directional

$$\mu_{\text{Ch}} = \Phi_{\text{Ch}} \int_{-\infty}^{\infty} ds \rho_{\text{Ch}}(s) \Omega(s) T_{\text{Ch}}(s) \epsilon(s) g(\cos \theta(s); s)$$

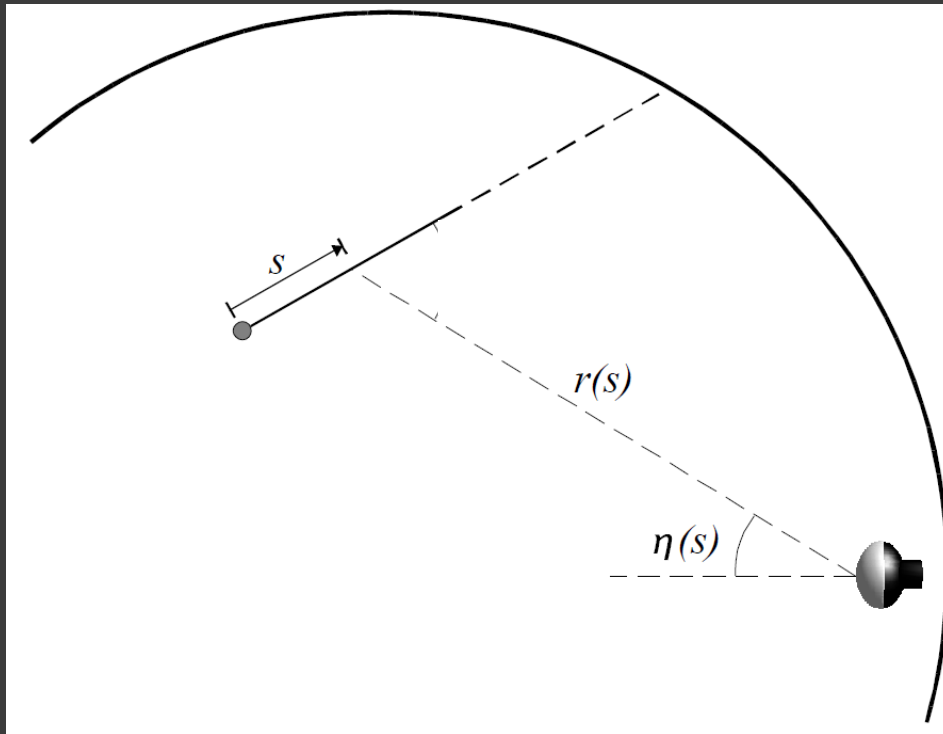
$g(\cos \theta(s); s)$  – angular emission profile



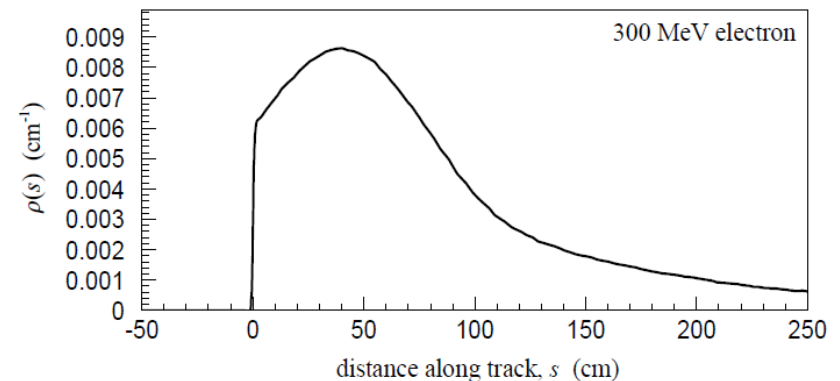
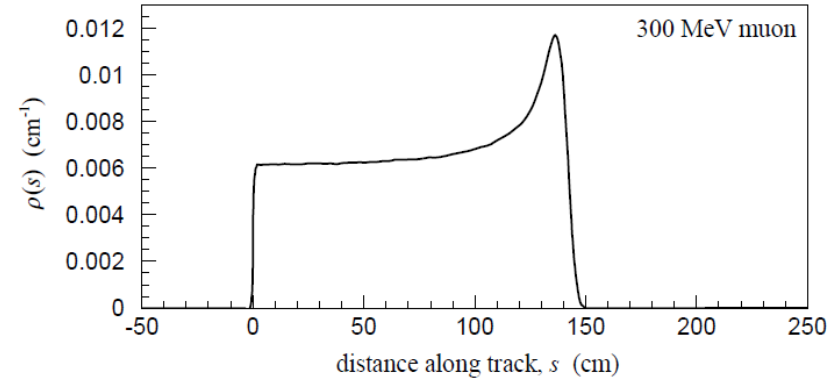
# Emission Profile - Cherenkov



# Track Fitting – Predicted Charge Extended Track



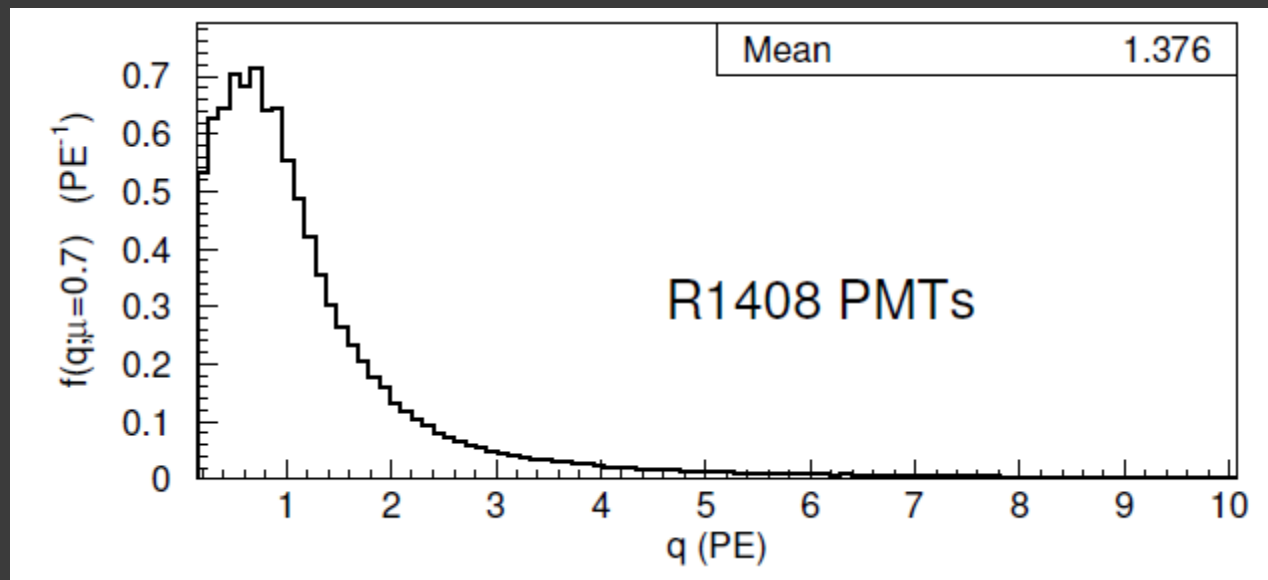
$\rho(s)$  normalized to 1



$$\mu_{\text{sci}} = \Phi_{\text{sci}} \int_{-\infty}^{\infty} ds \rho_{\text{sci}}(s) \Omega(s) T_{\text{sci}}(s) \epsilon(s)$$

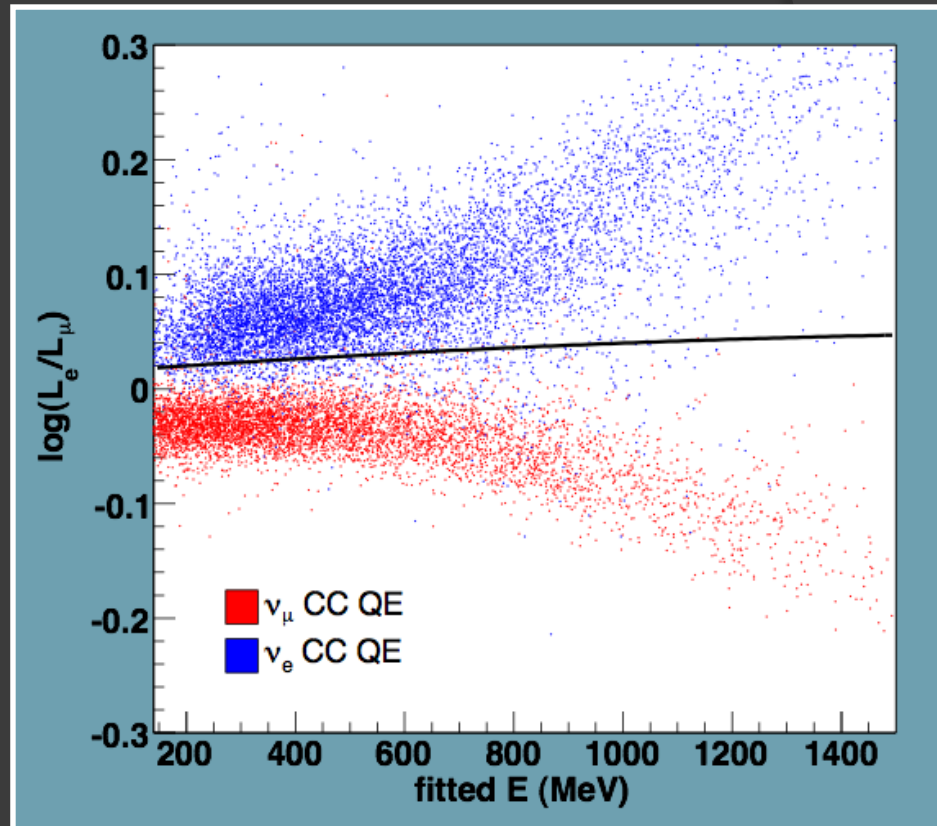
# Charge PDF

Measured by *in-situ* laser with control light output.



# Track-Based Analysis Rejecting Muon-like Events

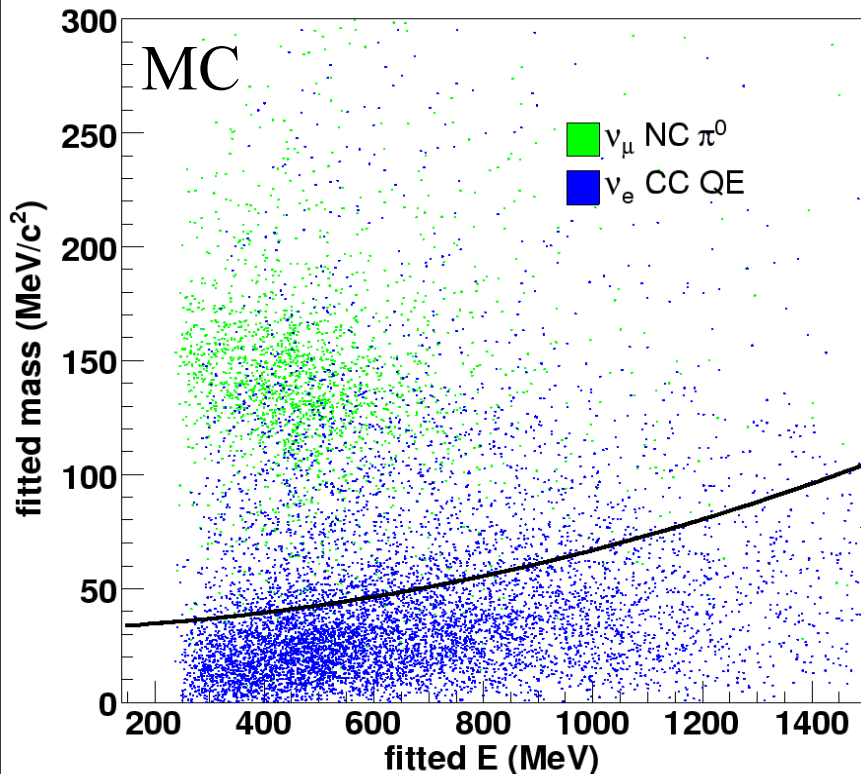
- Single track fit to muon and electron hypothesis
- $\log(L_e/L_\mu) > 0$  selects electron hypothesis.
- The cut is a quadratic function with energy, optimizing oscillation sensitivity.
- Separation is clean at high energies where muon-like events are long.



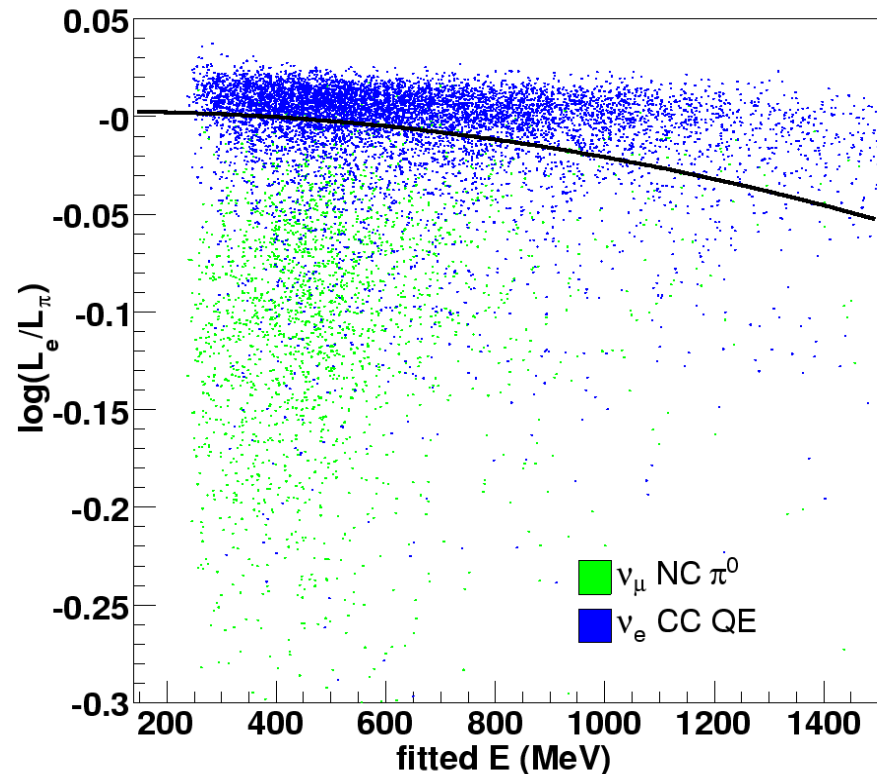


# Track-based Analysis Rejecting $\pi^0$ Events

Two track fit no mass constraint



Two track fit with  $\pi^0$  mass



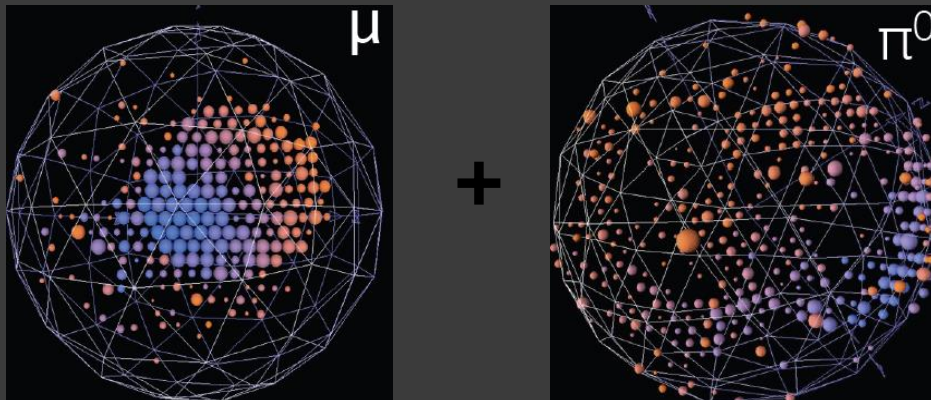
- Cuts are quadratic functions chosen to maximize  $\nu_\mu \rightarrow \nu_e$  sensitivity.  
 $\log(L_e/L_\pi) > 0$  – electron hypothesis fits better.

# $\nu_{\mu}$ CC $\pi^0$ Challenges

- CC  $\pi^0$  is tagged by one stopped muon decay electron (also CCQE signature).
- CC  $\pi^0$  is a small fraction (6%) in sample dominated by CCQE events 63%

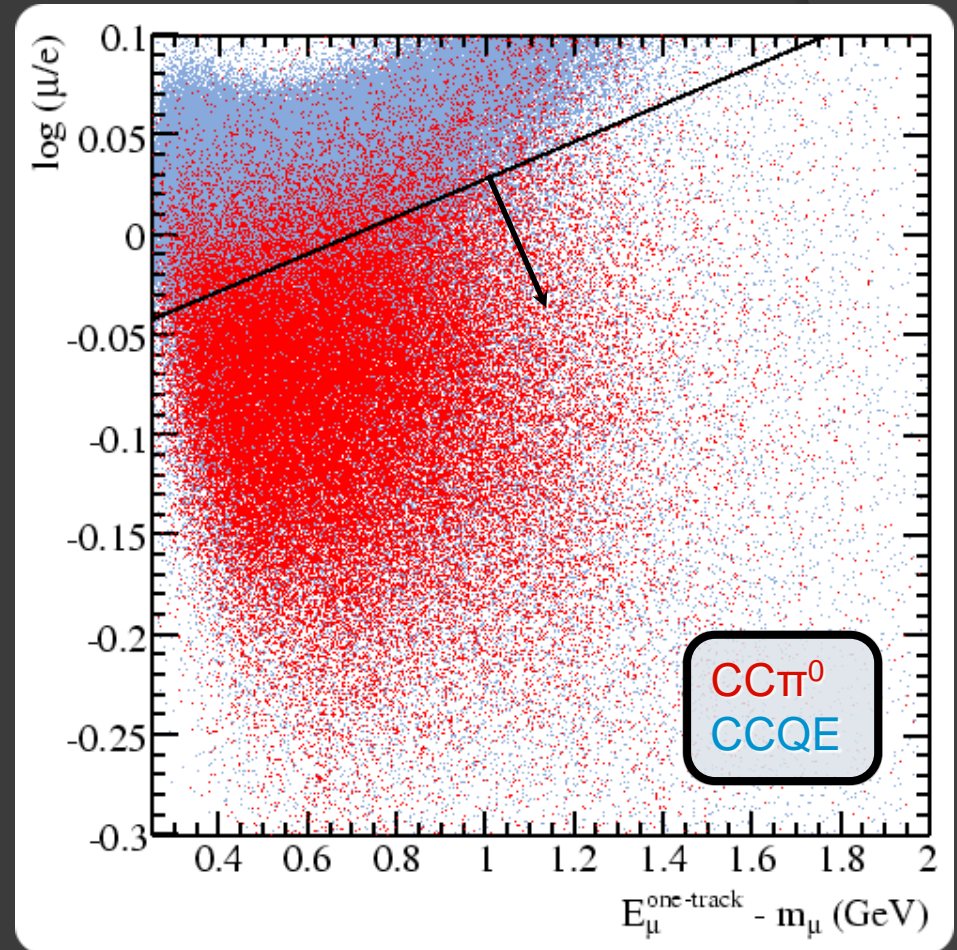
Sample	Events	Fraction
total MC	267007	100%
CCQE	168723	63%
CC $\pi^0$	16504	6%
CC $\pi^+$	66268	25%

- Overlapping rings make reconstruction more difficult.



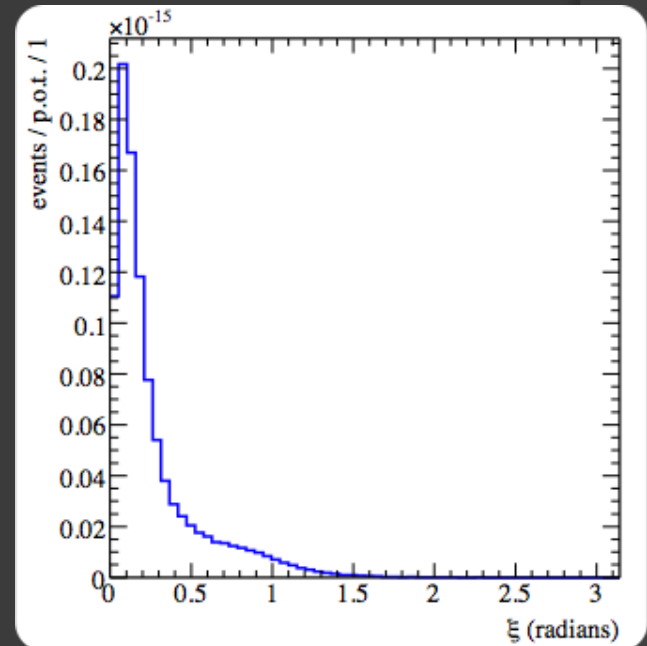
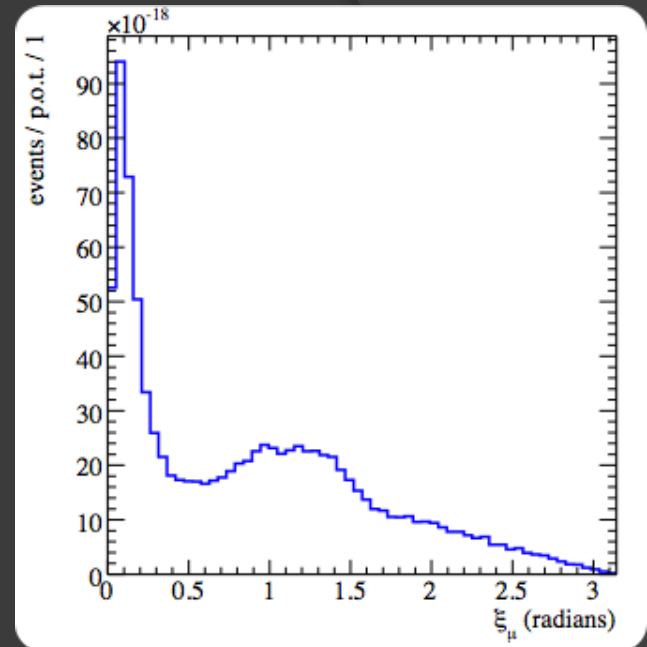
# Pre-filtering before the fit

- 2 subevents.
- Tank hits  $> 200$  (1<sup>st</sup> subevent)  
Tank hits  $< 200$  (2<sup>nd</sup> subevent)  
Veto hits  $< 6$  (both subevents)
- We need to reduce the two-subevent sample down to something more manageable before the fitter is run.
- A one-track likelihood ratio cut vs one-track energy reduces CCQE events by 98% while keeping 86% of  $CC\pi^0$  events.



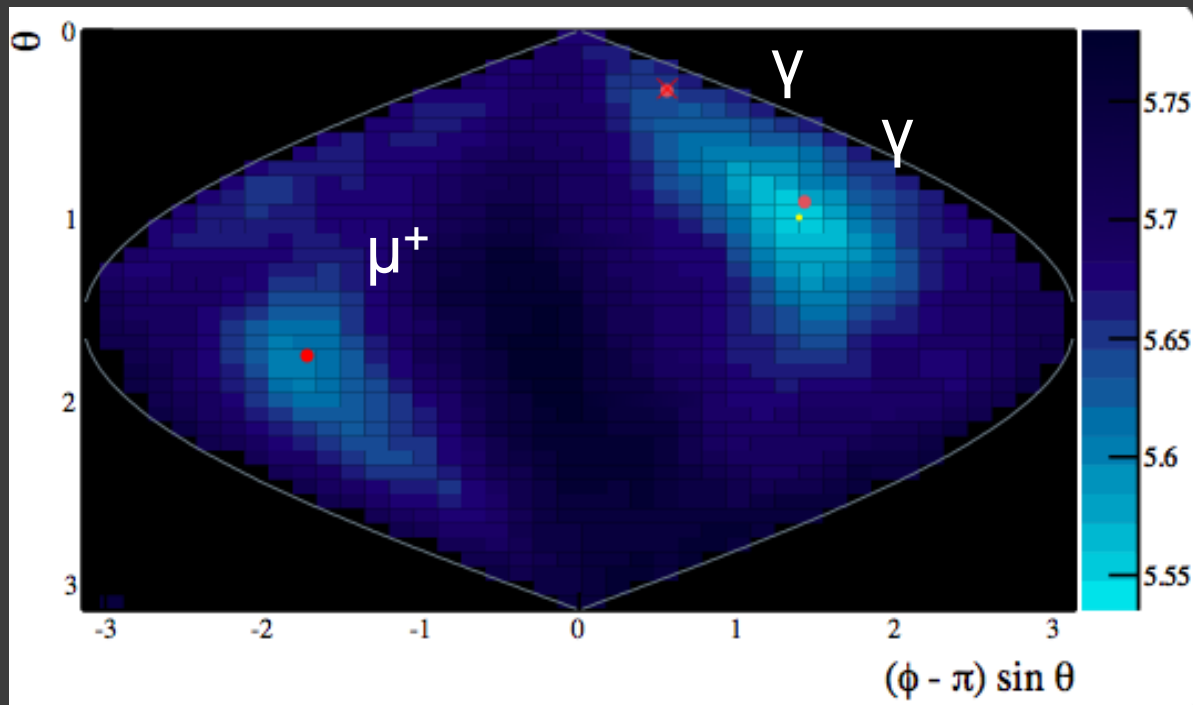
# Three-track fitting

- ❖ We start with a muon hypothesis.
- ❖ Measure the angle vs the true muon.
- ❖ That fit only finds the true muon  $\sim 1/3^{\text{rd}}$  of the time.
- ❖ However, it does a good job of finding one of the three rings.



# Reconstructing $CC\pi^0$ events

- ✦ Fixing the one-track muon fit in the likelihood function, we scan (in solid angle) for a second track.

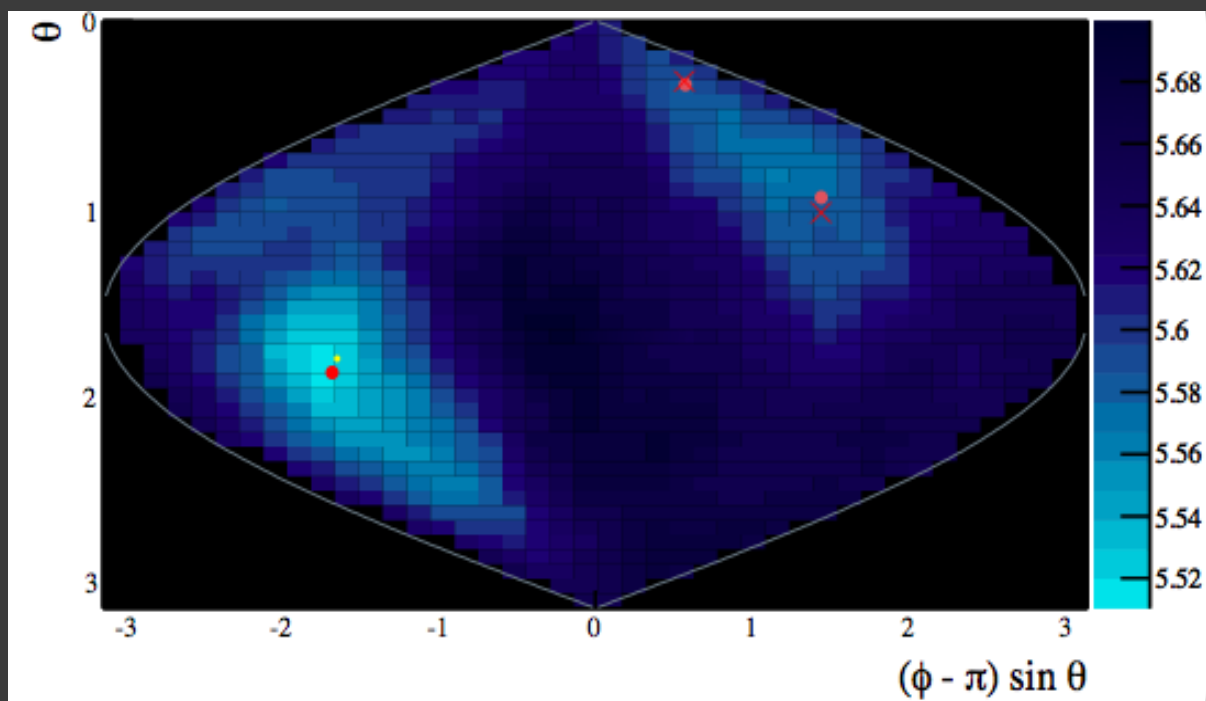


- : true track
- x : fit track
- : max likelihood

- ✦ The one track fit found one of the photons in this event.
- ✦ The scan found the second photon.
- ✦ After this scan, both tracks are allowed to float in a two-track fit.

# Reconstructing $CC\pi^0$ events

- Both tracks are fixed in the likelihood function. A third track is scanned for in all directions of solid angle.

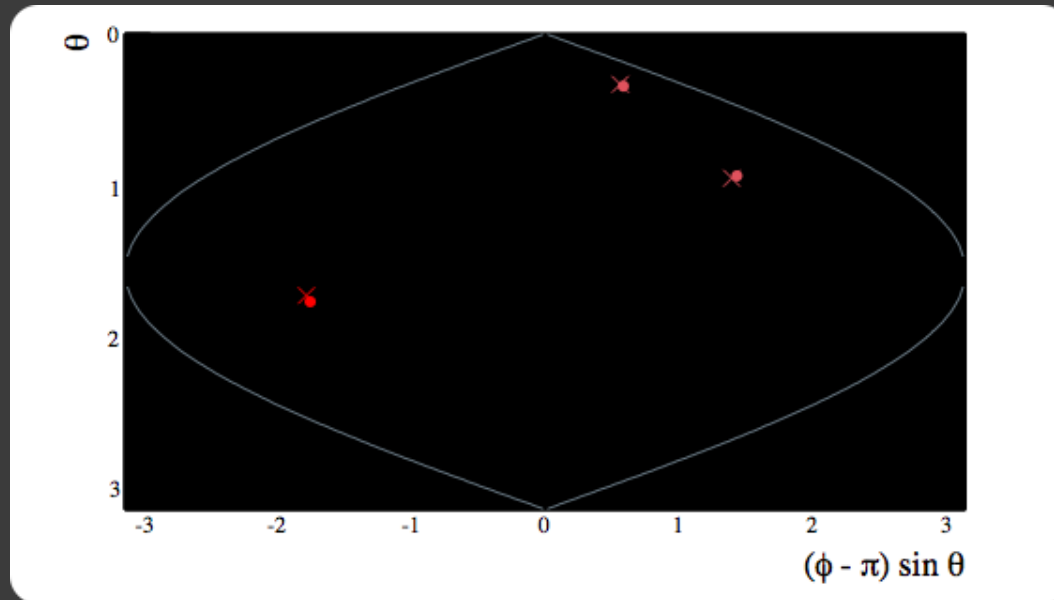


● : true track  
x : fit track  
● : max likelihood

- The two-track fit dimmed likelihood around the second photon, and brightened the likelihood around the muon.
- The scan found the muon in this event.

# Reconstructing $CC\pi^0$ events

- For all three possible particle configurations, additional three-track fits are performed. Swapping out two of the tracks for photons.

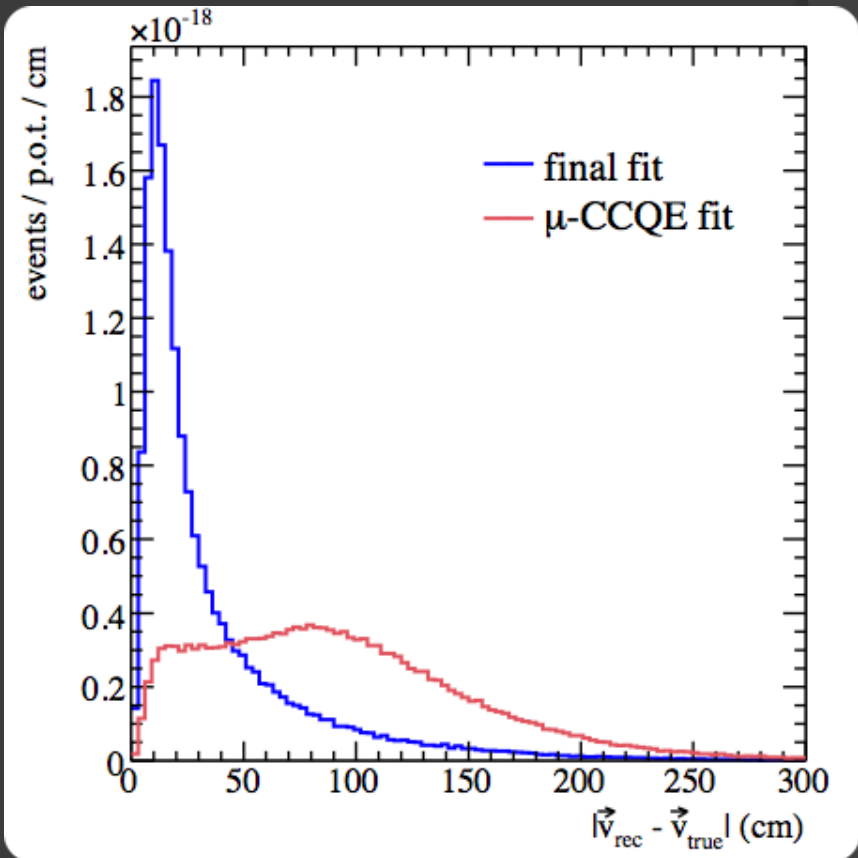
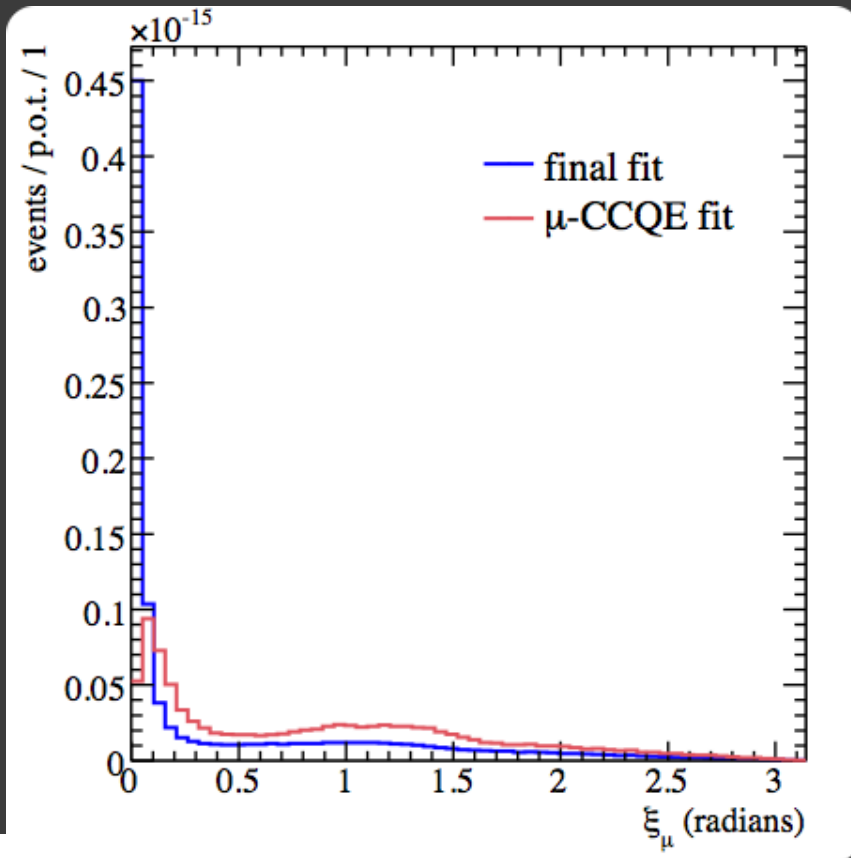


- Particle ID is performed by combining the fit likelihood and the direction to the 2<sup>nd</sup> subevent vertex (muon decay) vs the assumed muon in the fit as an additional likelihood.
- The three-track fit has identified all three particles ( $\mu, \gamma, \gamma$ ) in this event.



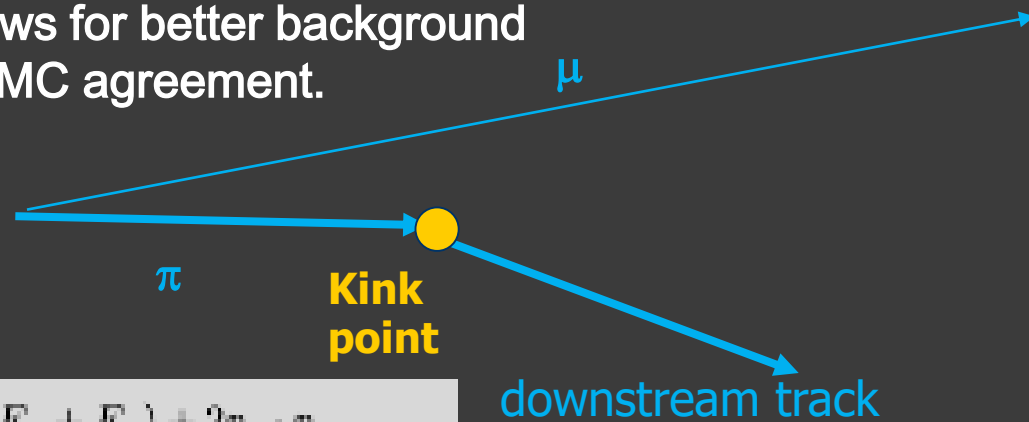
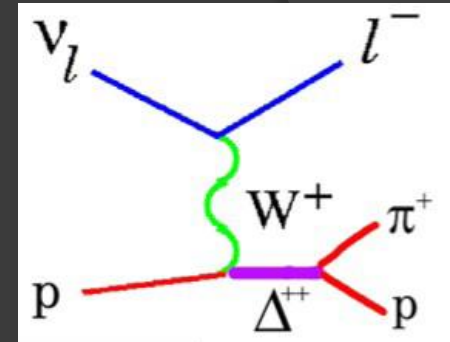
# Muon angle and event vertex

- ❖ The fitter has significantly improved the muon angular reconstruction.
- ❖ The event vertex has significantly improved.



# $\nu_\mu$ CC $\pi^+$ Reconstruction- a Step Further

- We can reconstruct the whole event if we reconstruct both  $\mu$  and  $\pi$  kinematics (assuming neutrino direction and target nucleon at rest).
- Need to reconstruct the pion – **kinked fitter**.
- Developed CC $\pi^+$  dedicated reconstruction .
- Better reconstruction allows for better background rejection and better data MC agreement.



$$E_\nu = \frac{m_\mu^2 + m_\pi^2 - 2m_N(E_\mu + E_\pi) + 2\mathbf{p}_\mu \cdot \mathbf{p}_\pi}{2(E_\mu + E_\pi - |\mathbf{p}_\mu| \cos \theta_{\nu,\mu} - |\mathbf{p}_\pi| \cos \theta_{\nu,\pi} - m_N)}$$

# CC Inclusive Event Reconstruction

New event reconstruction for MiniBooNE

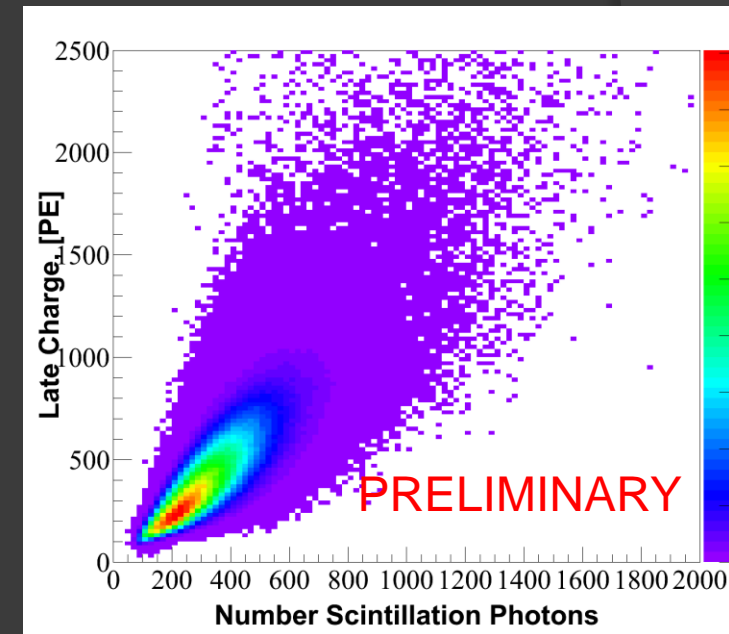
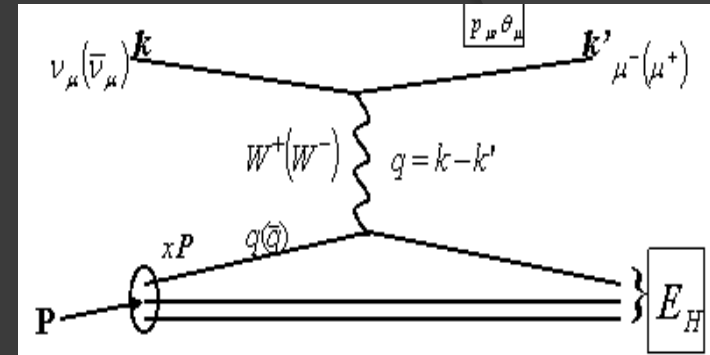
- Muon kinematics from 2-track likelihood fit:

Second ring of the fit absorbs the bias due to second most prominent ring.

- Neutrino energy – MiniBooNE detector as calorimeter.

Small scintillation light component produces late hits in the event. The charge of the late hits is used as a measure of the neutrino energy.

**Fully reconstruct the lepton vertex – no assumptions for the target!!!**

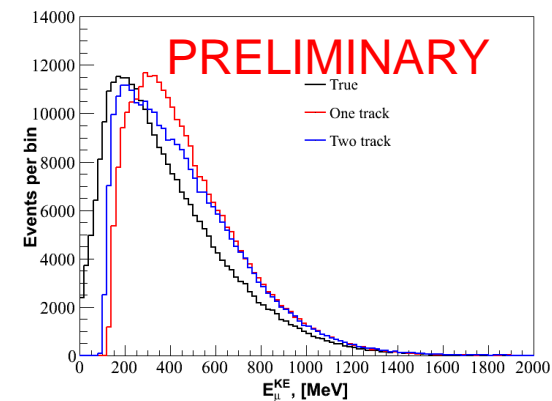
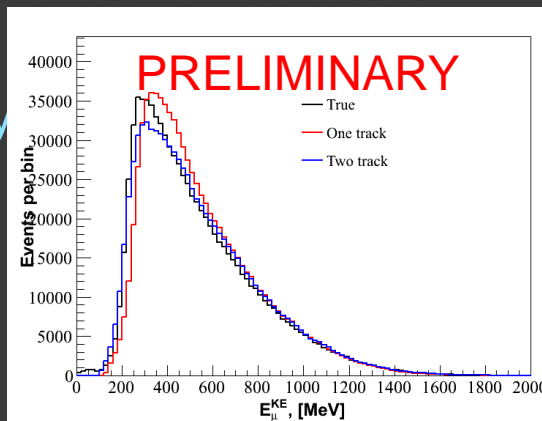


Plots are from MC

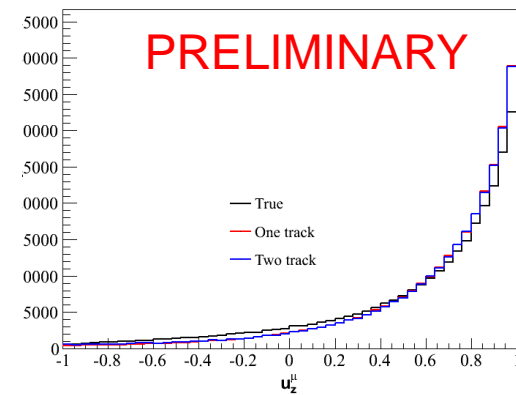
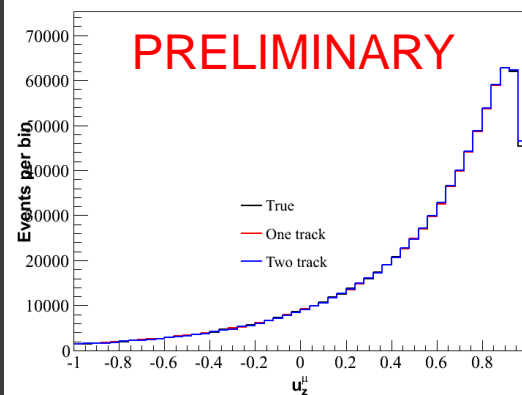
# Muon Kinematics Reconstruction Performance

2-track fit improves significantly reconstruction of the muon kinetic energy compared to one track fit. Muon kinetic energy resolution is about 5%.

$T_\mu$



$u_z^\mu$



No significant improvement for the muon angle. Muon angle resolution is better than 1°.

CCQE

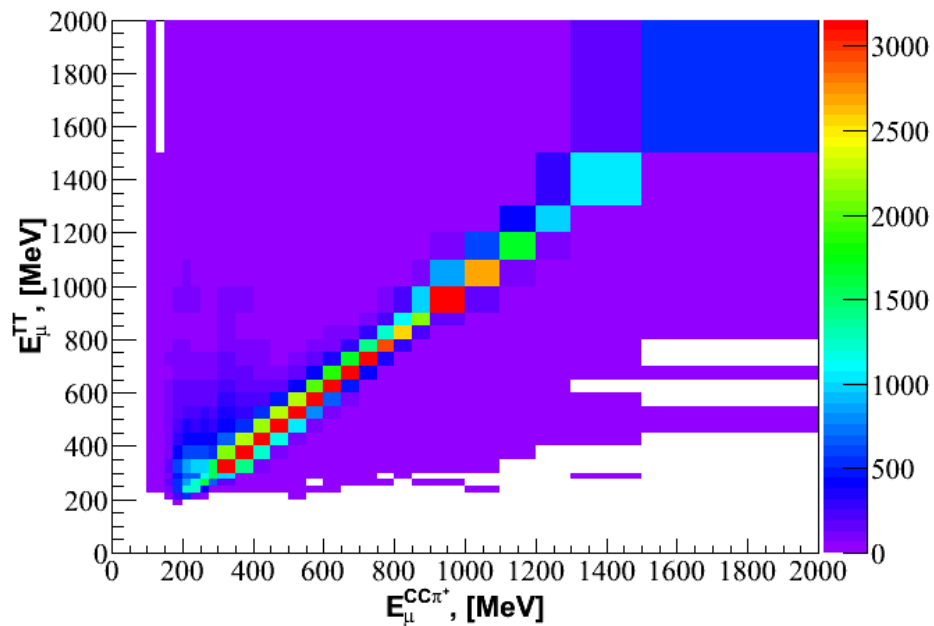
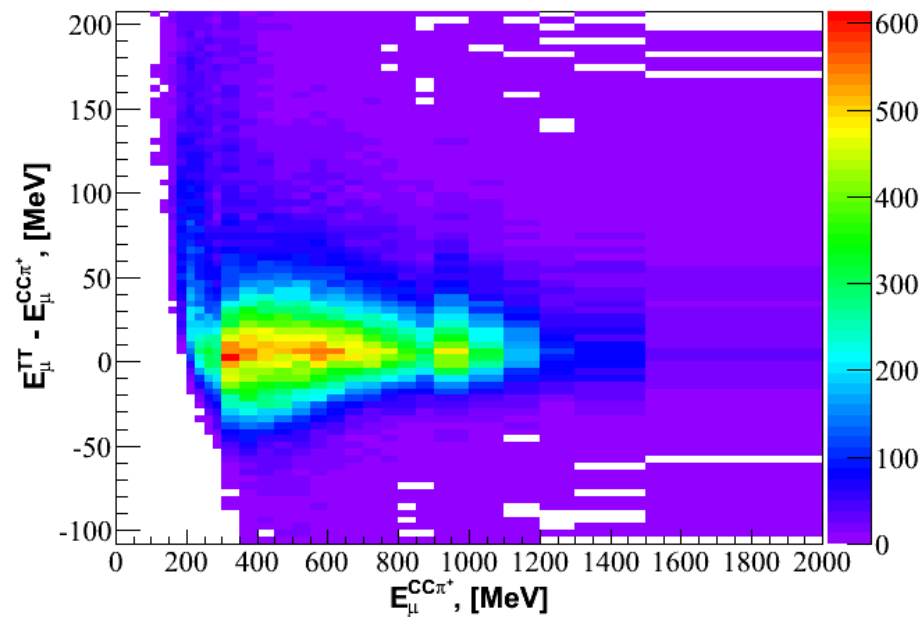
CCπ+

Plots are from MC.

# CC Inclusive Reconstruction – $\text{CC}\pi^+$ Sample

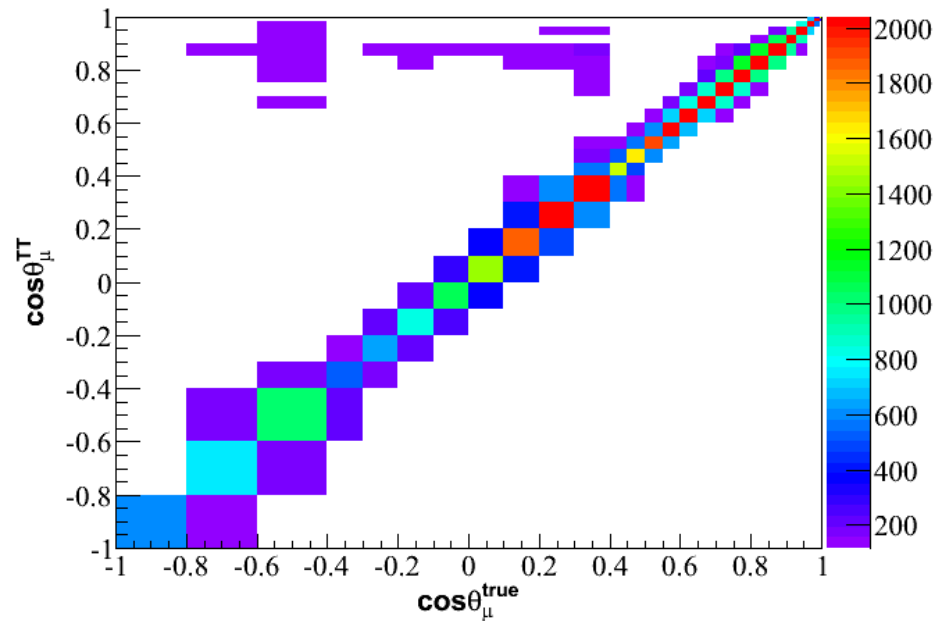
Event-by-event difference

TT vs  $\text{CC}\pi$  fitter

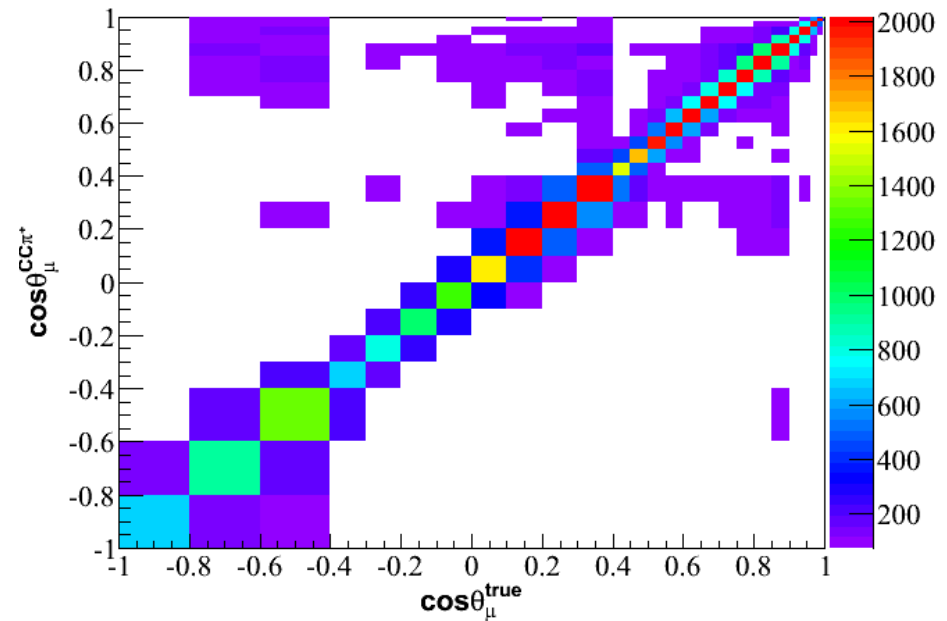


# CC Inclusive Reconstruction – $CC\pi^+$ Sample

## Uz Rec vs Uz true



Two Track Fit



$CC\pi^+$  Fit

# Neutrino Energy Reconstruction Performance

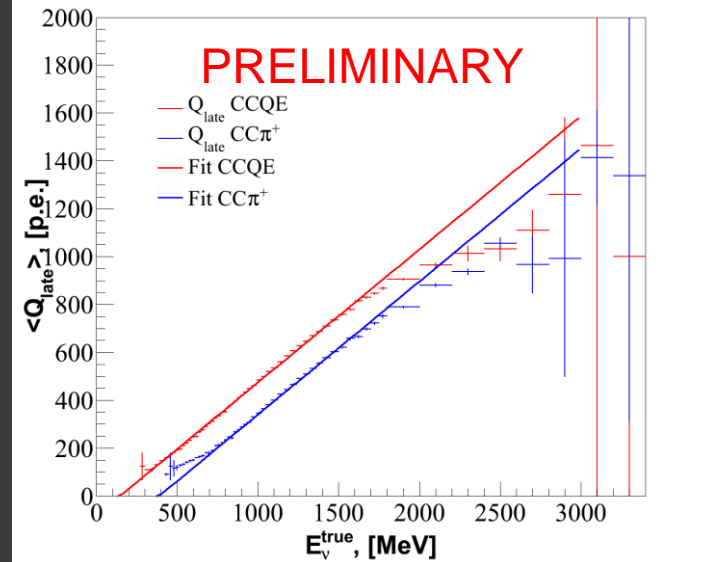
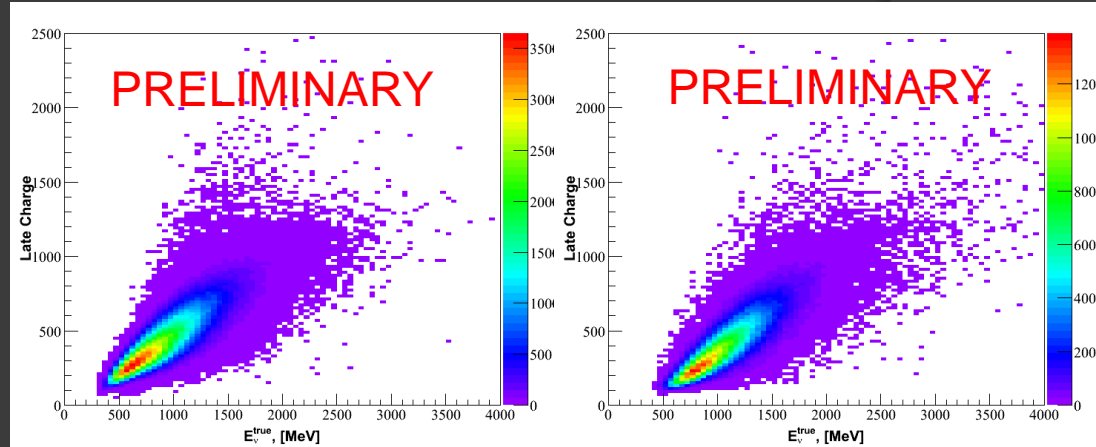
CCQE

CC $\pi^+$

Neutrino energy reconstruction is obtained from the late light charge which is linearly correlated with the true neutrino energy.

The parameters of the reconstruction come from a linear fit to both CCQE and CC $\pi^+$  enhanced samples. the slope parameter is the same in both cases while the Intercept is different.

Energy reconstruction resolution is about 18%.



Plots are from MC.

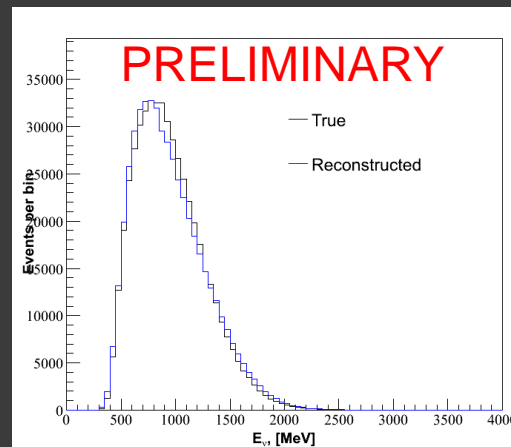
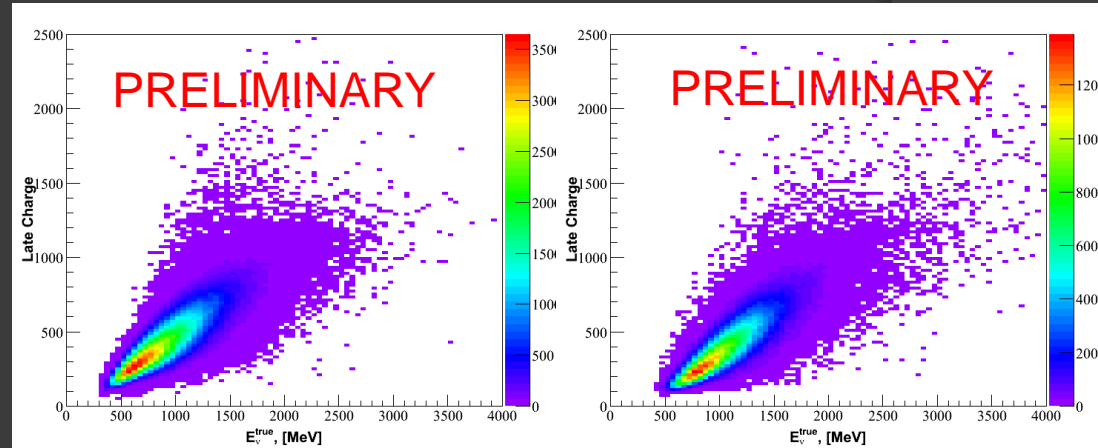


# Neutrino Energy Reconstruction Performance

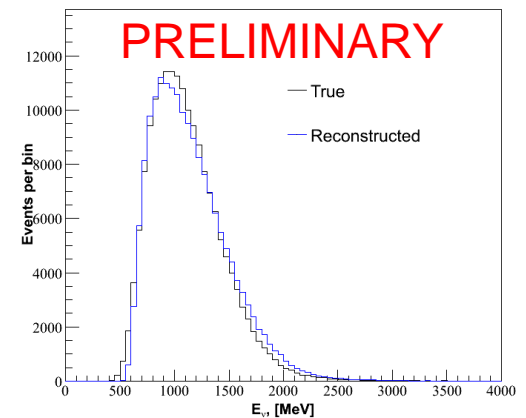
Neutrino energy reconstruction is obtained from the late light charge which is linearly correlated with the true neutrino energy.

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CCQE



$CC\pi^+$

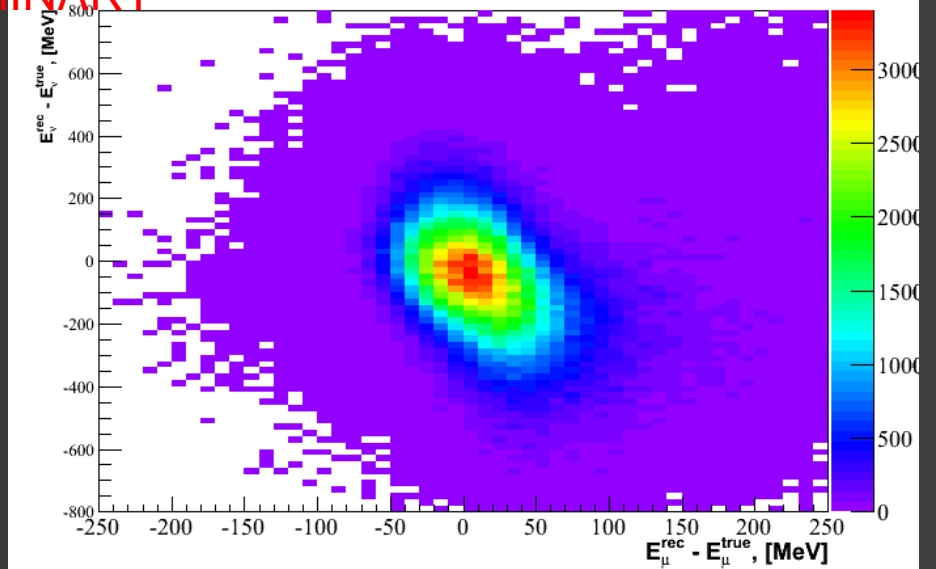
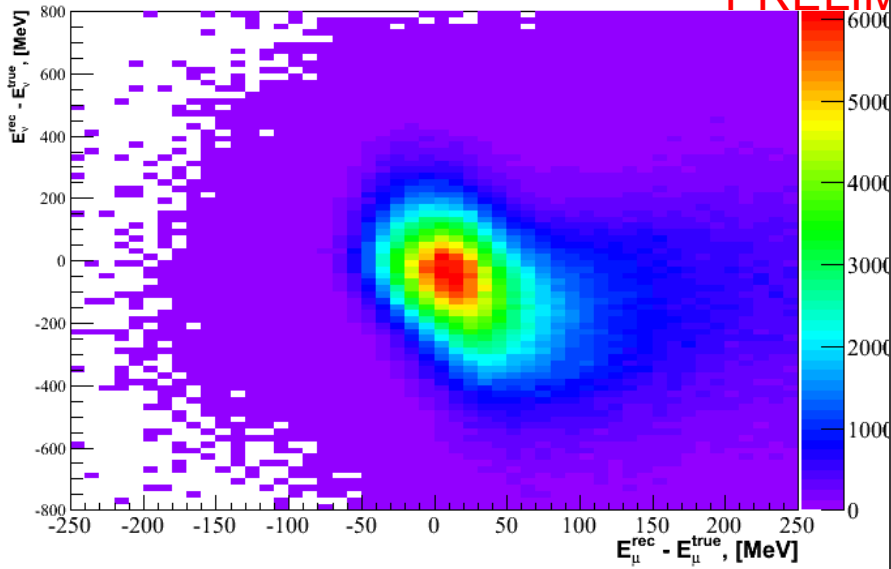
Plots are from MC.

# Reconstruction Correlation

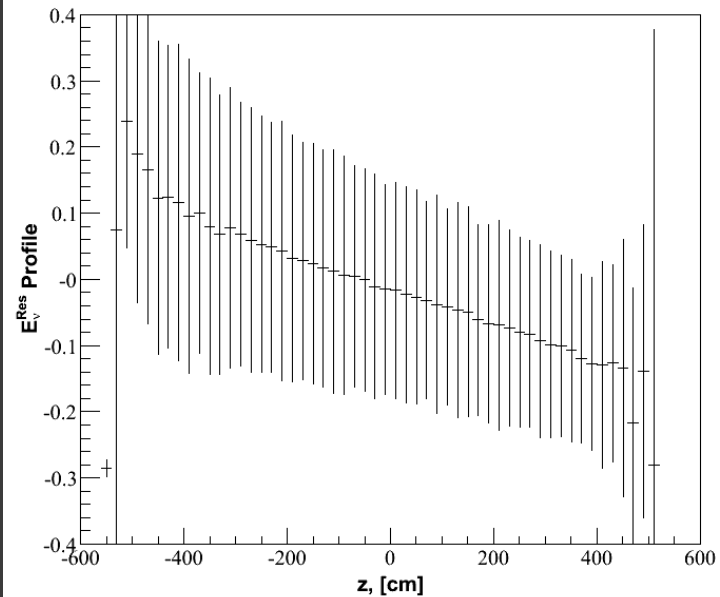
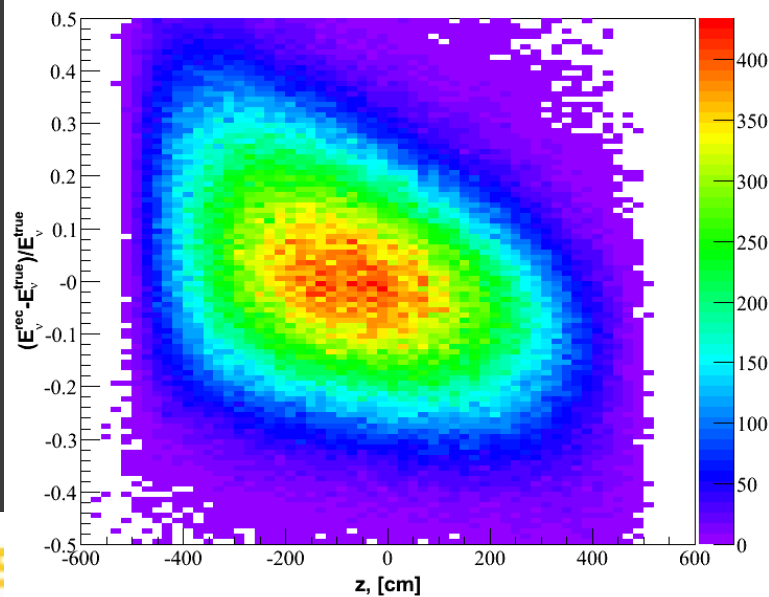
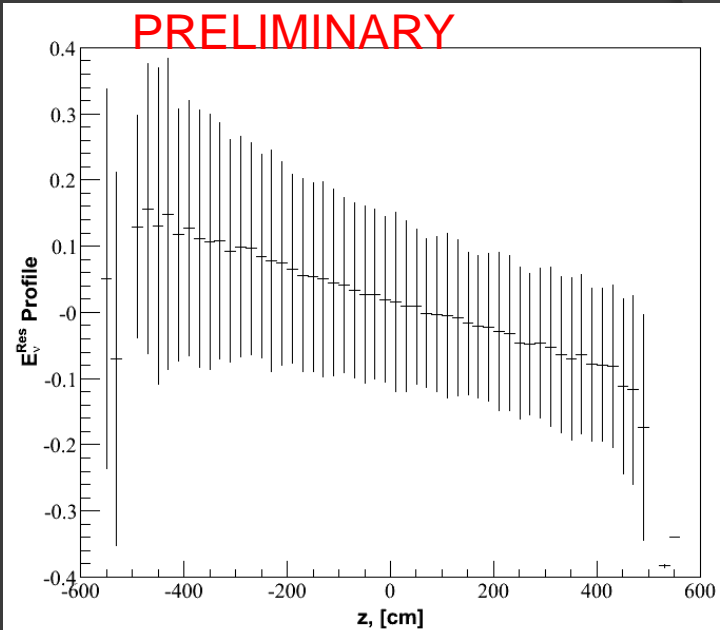
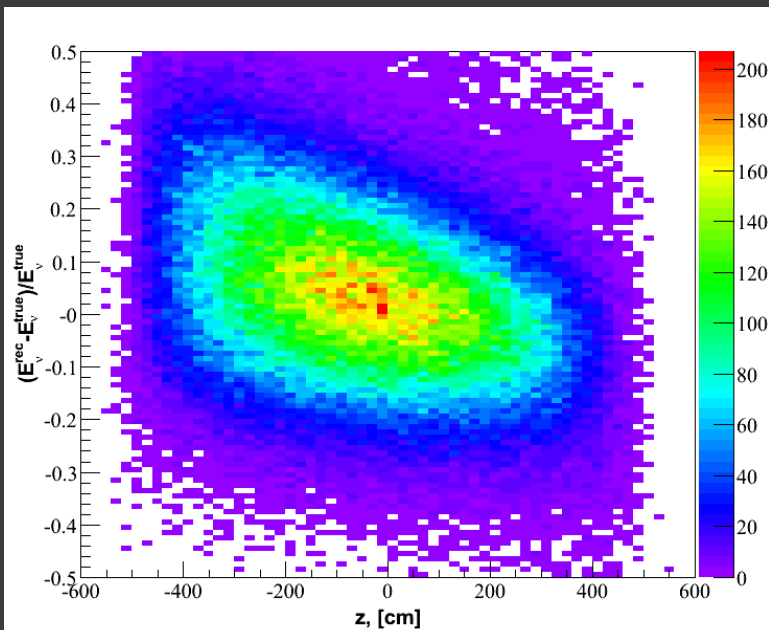
All SE

CCQE

PRELIMINARY



# CC Inclusive Reconstruction – $E_\nu$ Resolution



# Summary

This reconstruction has been successfully applied to SK

The MB reconstruction depends on the “optical” model

Requires very good understanding of the optical properties

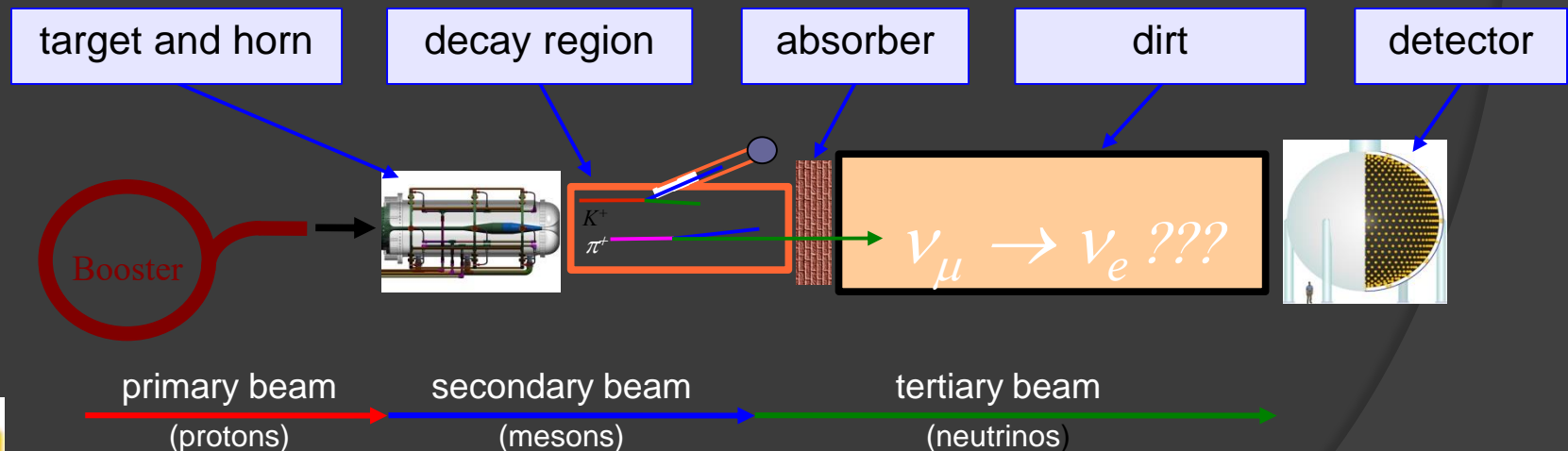
Requires good coverage and good time resolution

Either separate the sci/ser light and/or save the waveforms.

# MiniBooNE Experiment – E898 at Fermilab

Test of LSND within the context of  $\nu_{\mu} \rightarrow \nu_e$  appearance only is an essential first step:

- Keep the same L/E
- Higher energy and longer baseline –  $E=0.5 - 1$  GeV;  $L=500$ m
- Different beam
- Different oscillation signature  $\nu_{\mu} \rightarrow \nu_e$
- Different systematics
- Antineutrino-capable beam



# Neutrino Flux Prediction

- GEANT4 based Monte Carlo simulates the neutrino flux in MiniBooNE beamline,
- high purity  $\nu_\mu$  beam – 99%, small  $\nu_e$  component – intrinsic  $\nu_e$  - background for  $\nu_e$  appearance

$$\nu_\mu \rightarrow \nu_e, \quad \bar{\nu}_e/\nu_\mu = 0.5\%$$

- “Intrinsic”  $\nu_e + \bar{\nu}_e$  sources:

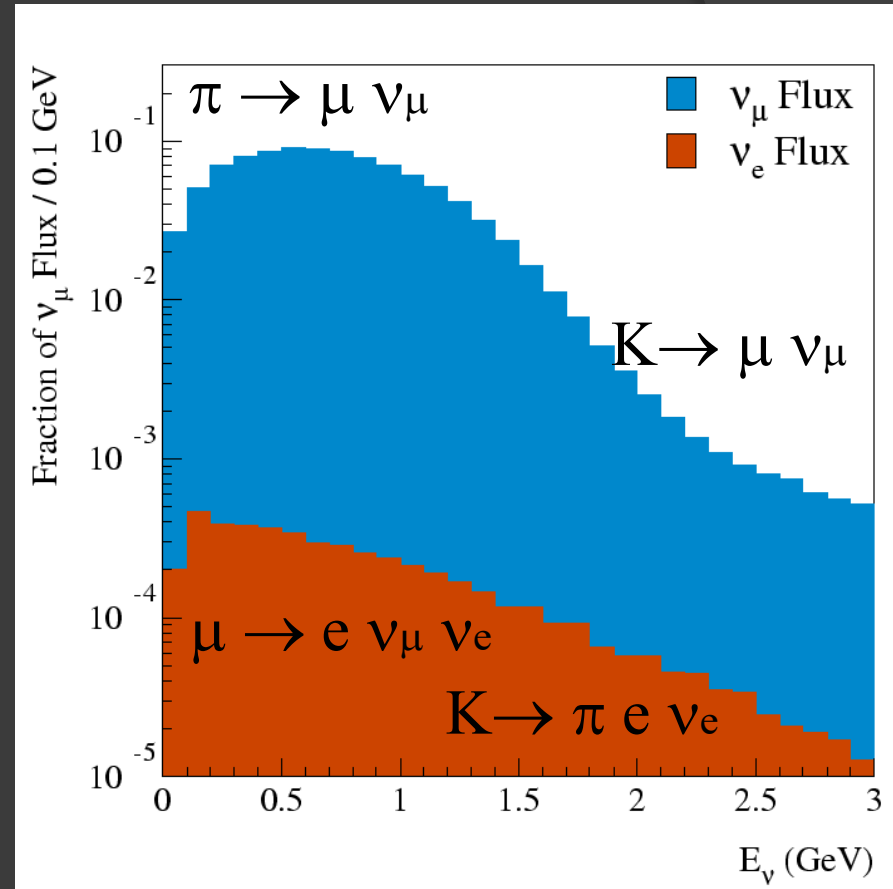
$$\mu^+ \rightarrow e^+ \bar{\nu}_\mu \nu_e \quad (52\%)$$

$$K^+ \rightarrow \pi^0 e^+ \nu_e \quad (29\%)$$

$$K^0 \rightarrow p e \nu_e \quad (14\%)$$

$$\text{Other} \quad (5\%)$$

- Antineutrino content: 6%



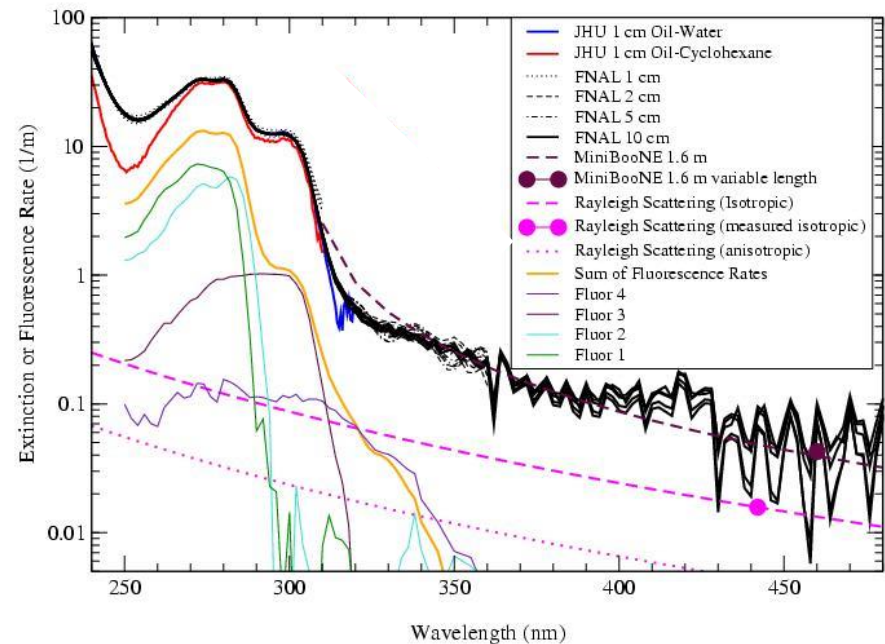
# Detector “Optical” Model

## Primary light sources

- Cherenkov
  - Emitted promptly, in cone known wavelength distribution
- Scintillation
  - Emitted isotropically
  - Several lifetimes, emission modes
  - Studied oil samples using Indiana Cyclotron test beam
  - Particles below Cherenkov threshold still scintillate

*We have developed  
39-parameter  
“Optical Model”  
based on internal calibration  
and external measurement*

Extinction Rate for MiniBooNE Marcol 7 Mineral Oil



- Optical properties of oil, detectors:
- Absorption  
(attenuation length >20m at 400 nm)
  - Rayleigh and Raman scattering
  - Fluorescence
  - Reflections