

Measurement of the proton scalar polarizabilities at MAMI

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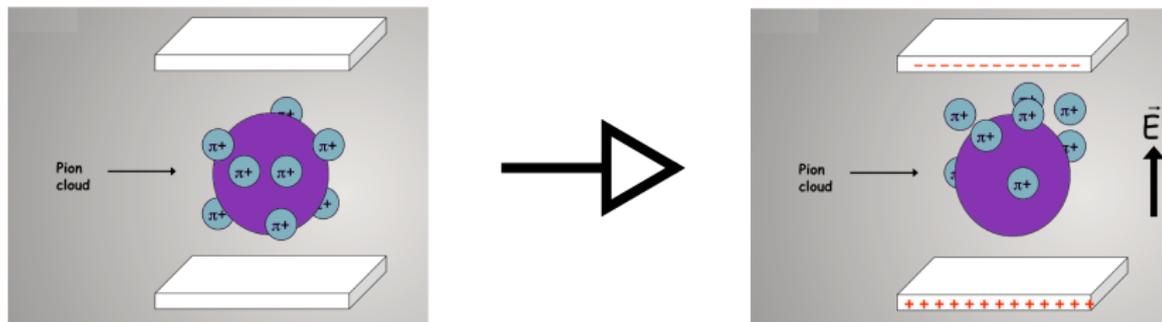


Main goal

Measurement of the scalar polarizabilities α_{E1} and β_{M1} of the proton, using a linearly polarized photon beam

- Fundamental properties of the nucleon
- Closely related to nucleon internal structure
- Important for nuclear physics, atomic physics, astrophysics, spin polarizabilities measurements, etc.

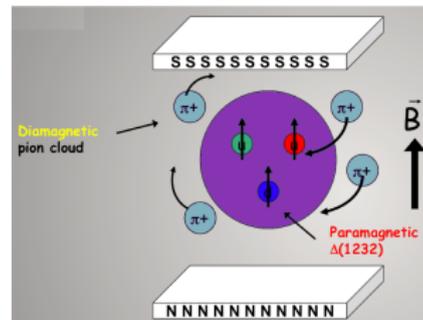
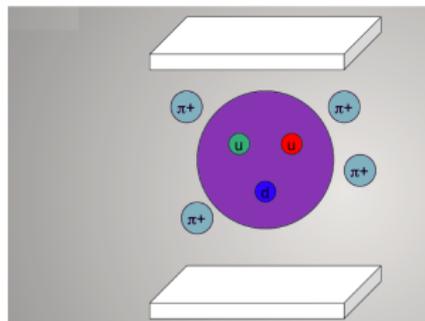
Electric scalar polarizability - α_{E1}



- Electric dipole moment: $\vec{p} = \alpha_{E1} \vec{E}$
- α_{E1} : electric polarizability

- Describes the response of a proton to an applied electric field: "stretchability" of the proton

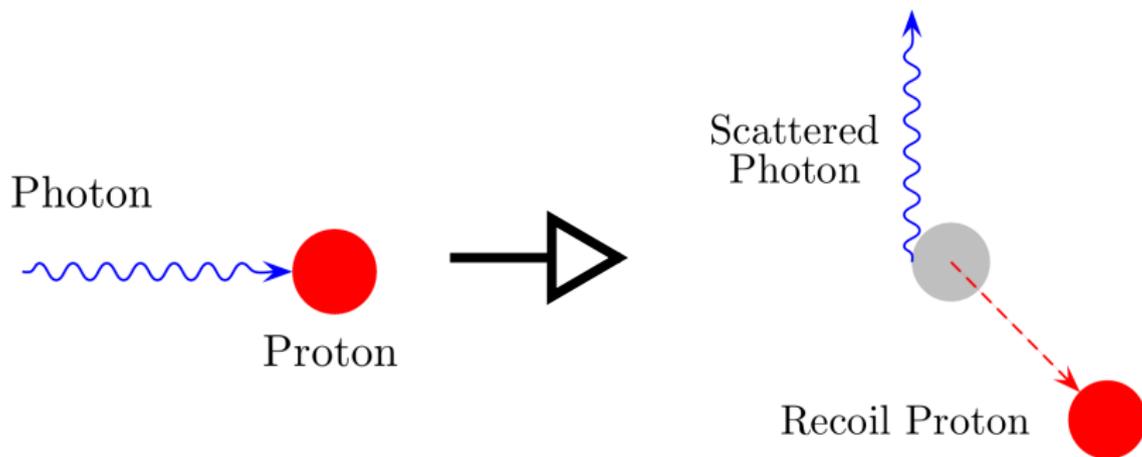
Magnetic scalar polarizability - β_{M1}



- Magnetic dipole moment: $\vec{m} = \beta_{M1}\vec{H}$
- β_{M1} : magnetic polarizability
- Describes the response of a proton to an applied magnetic field: "alignability" of the proton

Compton scattering

- Internal structure can be accessed via Compton scattering



$$\gamma(\mathbf{k}) + \text{P}(\mathbf{p}) \rightarrow \gamma(\mathbf{k}') + \text{P}(\mathbf{p}')$$

Born term

Under the assumption of NO proton internal structure, the effective Hamiltonian can be written in terms of mass, electric charge and anomalous magnetic moment

- Zeroth order: mass and electric charge

$$H_{\text{eff}}^{(0)} = \frac{\vec{\pi}^2}{2m} + e\phi \quad (\text{where } \vec{\pi} = \vec{p} - e\vec{A})$$

- First order: anomalous magnetic moment

$$H_{\text{eff}}^{(1)} = -\frac{e(1+k)}{2m}\vec{\sigma} \cdot \vec{H} - \frac{e(1+2k)}{8m^2}\vec{\sigma} \cdot [\vec{E} \times \vec{\pi} - \vec{\pi} \times \vec{E}]$$

Scalar polarizabilities

Effective Hamiltonian at the second order includes scalar polarizabilities, which are related to the proton internal structure

- Second order: scalar polarizabilities α_{E1} and β_{M1}

$$H_{\text{eff}}^{(2)} = -4\pi \left[\frac{1}{2} \alpha_{E1} \vec{E}^2 + \frac{1}{2} \beta_{M1} \vec{H}^2 \right]$$

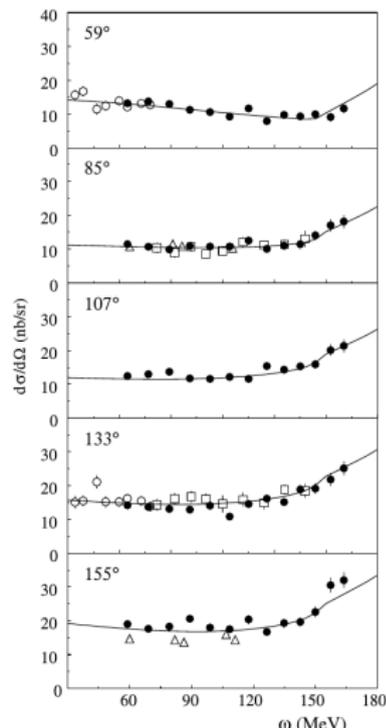
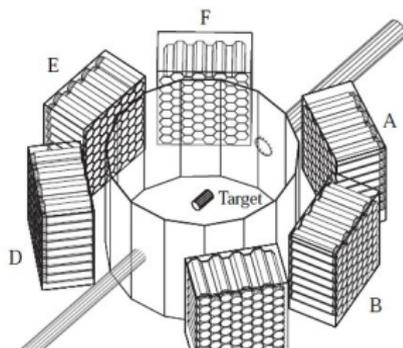
α_{E1} and β_{M1} quantify the response of the proton to a static applied electric or magnetic field

- Third order: spin polarizabilities γ_{E1E1} , γ_{M1M1} , γ_{M1E2} and γ_{E1M2}

$$H_{\text{eff}}^{(3)} = -4\pi \left[\frac{1}{2} \gamma_{E1E1} \vec{\sigma} \cdot (\vec{E} \times \vec{E}) + \frac{1}{2} \gamma_{M1M1} \vec{\sigma} \cdot (\vec{H} \times \vec{H}) \right. \\ \left. - \gamma_{M1E2} E_{ij} \sigma_i H_j + \gamma_{E1M2} H_{ij} \sigma_i E_j \right]$$

Existing data

- Highest statistics published data:
V. Olmos de Leon et al. Eur. Phys. J. A 10, 207-215 (2001)
- 200 hours of Compton scattering
- $E_{\text{beam}} = 180 \text{ MeV}$
- $E_{\gamma} = 55 - 165 \text{ MeV}$, $\theta_{\gamma} = 59^{\circ} - 155^{\circ}$
- 1/3 acceptance of CB system

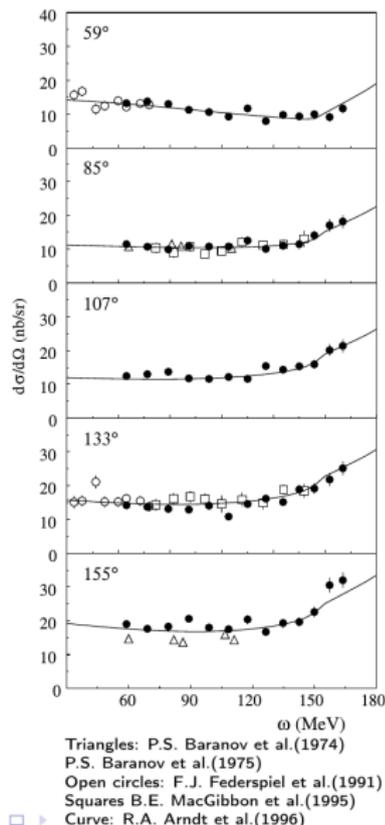


Triangles: P.S. Baranov et al.(1974)
P.S. Baranov et al.(1975)
Open circles: F.J. Federspiel et al.(1991)
Squares B.E. MacGibbon et al.(1995)
Curve: R.A. Arndt et al.(1996)

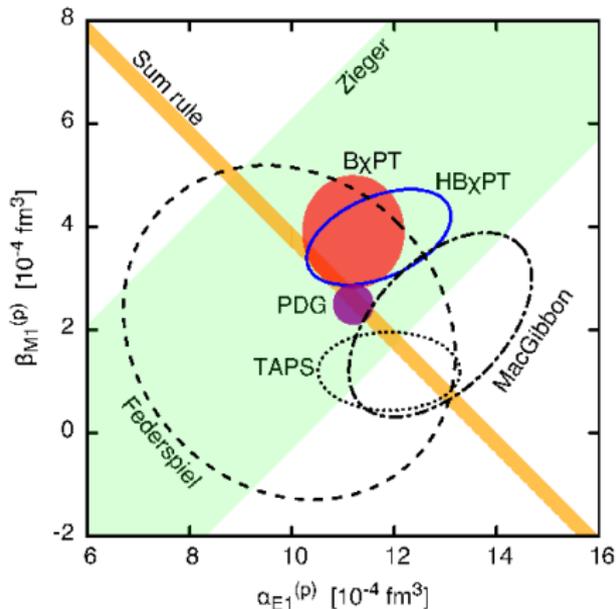
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	with Baldin	without Baldin
α_{E1}	12.1 ± 1.08	11.9 ± 1.39
β_{E1}	1.6 ± 0.89	1.2 ± 0.76



Existing data



PDG (2012) values:

$$\alpha_{E1} = (12.0 \pm 0.6) 10^{-4} \text{ fm}^3$$

$$\beta_{M1} = (1.9 \pm 0.5) 10^{-4} \text{ fm}^3$$

New PDG (2013) values:

$$\alpha_{E1} = (11.2 \pm 0.4) 10^{-4} \text{ fm}^3$$

$$\beta_{M1} = (2.5 \pm 0.4) 10^{-4} \text{ fm}^3$$

Significant change between reviews
without new experimental data

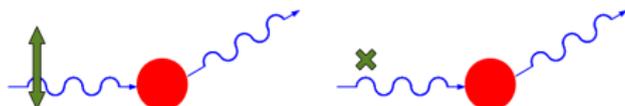
⇒ New high quality data needed!

At low energy (and in certain phase-space regions) Σ_3 is mainly dependent on β_{M1}

Krupina and Pascalutsa, PRL 110, 262001 (2013)

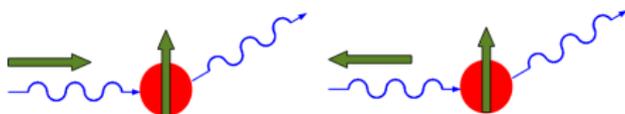
- Linearly polarized beam & unpolarized target

$$\Sigma_3 = \frac{d\sigma_{\perp} - d\sigma_{\parallel}}{d\sigma_{\perp} + d\sigma_{\parallel}}$$



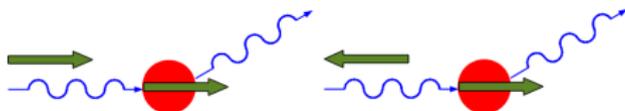
- Circularly polarized beam & transversely polarized target

$$\Sigma_{2x} = \frac{\sigma_{+x}^R - \sigma_{+x}^L}{\sigma_{+x}^R + \sigma_{+x}^L}$$



- Circularly polarized beam & longitudinally polarized target

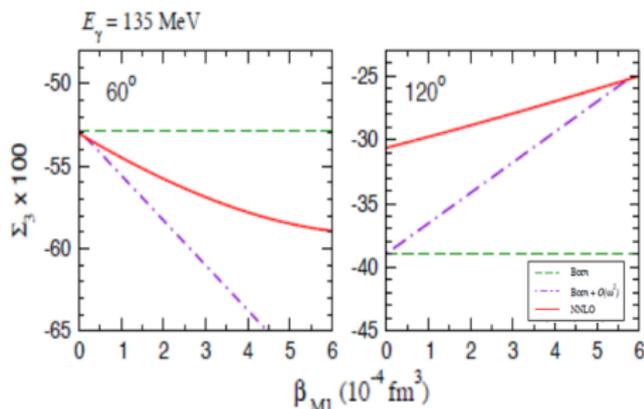
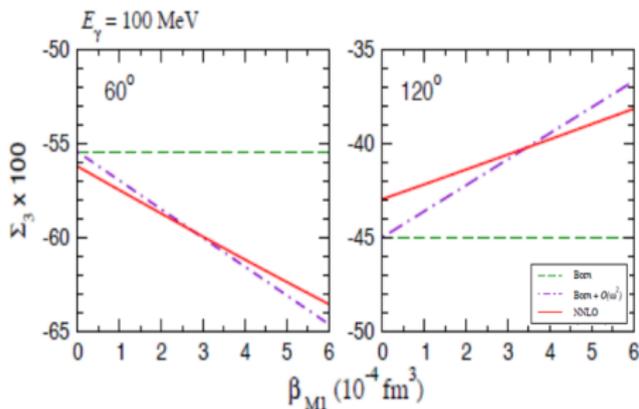
$$\Sigma_{2z} = \frac{\sigma_{+z}^R - \sigma_{+z}^L}{\sigma_{+z}^R + \sigma_{+z}^L}$$



Beam asymmetry Σ_3 and measurement of β_{M1}

At low energy, β_{M1} can be extracted from the beam asymmetry Σ_3 using a linearly polarized photon beam and an unpolarized proton target:

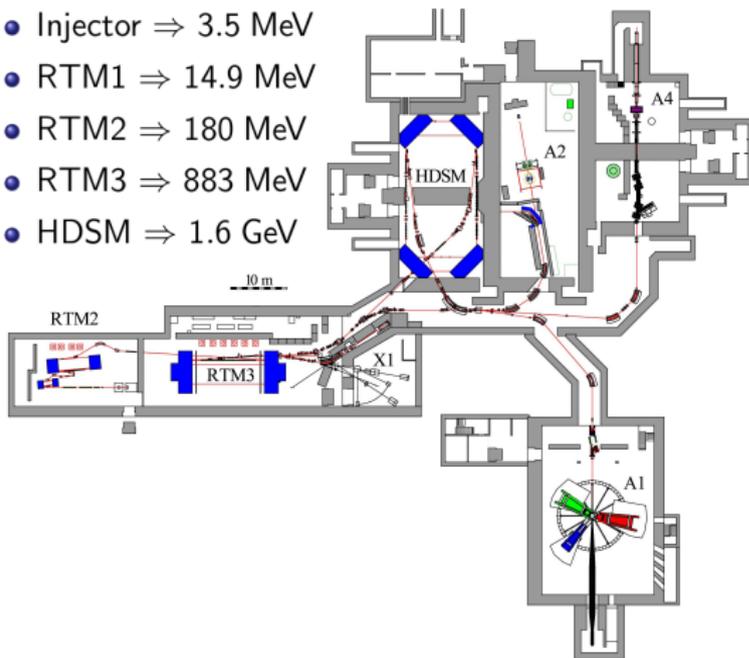
$$\frac{d\sigma}{d\Omega}(\theta, \phi) = \frac{d\sigma}{d\Omega}(\theta) [1 + p_\gamma \Sigma_3 \cos(2\phi)] \quad \text{where} \quad \Sigma_3 = \frac{d\sigma_\perp - d\sigma_\parallel}{d\sigma_\perp + d\sigma_\parallel}$$



Krupina and Pascalutsa, PRL 110, 262001 (2013)

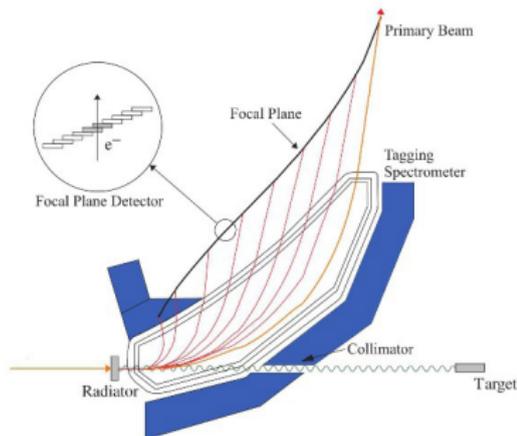
Experimental setup

- Injector \Rightarrow 3.5 MeV
- RTM1 \Rightarrow 14.9 MeV
- RTM2 \Rightarrow 180 MeV
- RTM3 \Rightarrow 883 MeV
- HDSM \Rightarrow 1.6 GeV



High intensity beam of linearly polarized tagged photons:

$$E_{\gamma} = E_0 - E_{e^{-}}$$

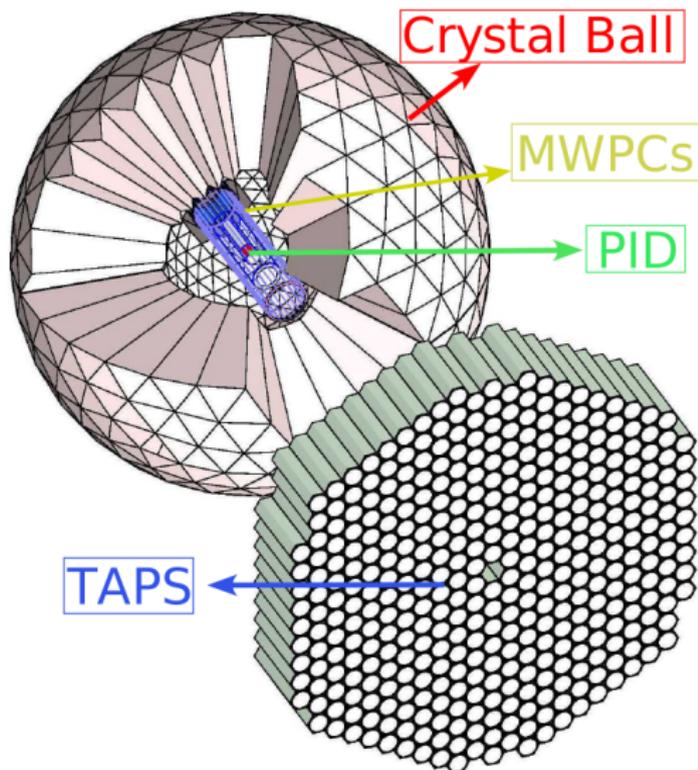


Crystal Ball

- 672 NaI(Tl) crystals
- Particle Identification Detector (**PID**):
24 scintillator paddles
- 2 Multiwire Proportional Chambers (**MWPCs**)

TAPS

- 366 BaF₂ and
72 PbWO₄ crystals
- 384 veto paddles



⇒ The beam asymmetry Σ_3 is an alternative way to extract β_{M1} :

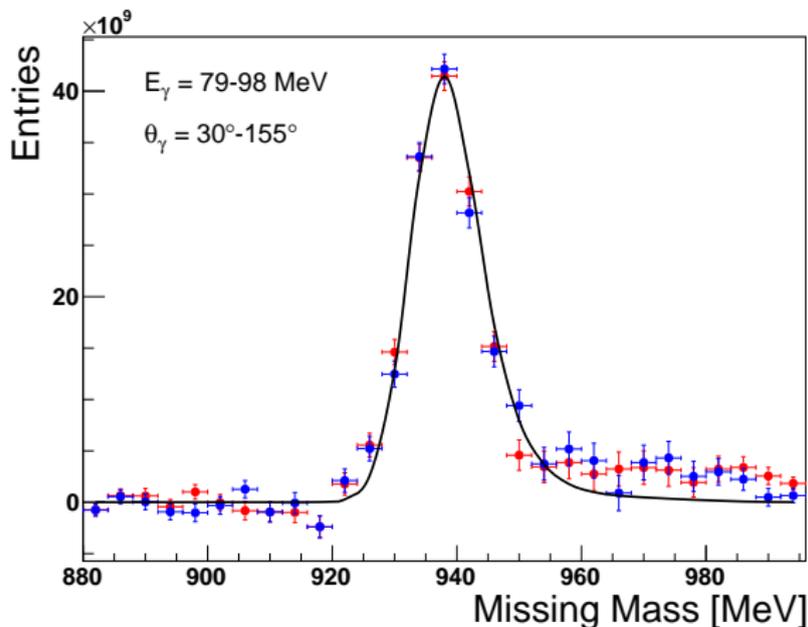
$$\Sigma_3 = \frac{d\sigma_{\perp} - d\sigma_{\parallel}}{d\sigma_{\perp} + d\sigma_{\parallel}}$$

⇒ We selected Compton scattering $\vec{\gamma}p \rightarrow \gamma(p)$ events:

- $E_{\gamma_{\text{beam}}} = 79 - 139 \text{ MeV}$
- $\theta_{\gamma_{\text{out}}} = 30^{\circ} - 155^{\circ}$
- Events with 1 γ in the final state
- Random background subtraction
- Subtraction of empty target contribution
- Missing mass cut

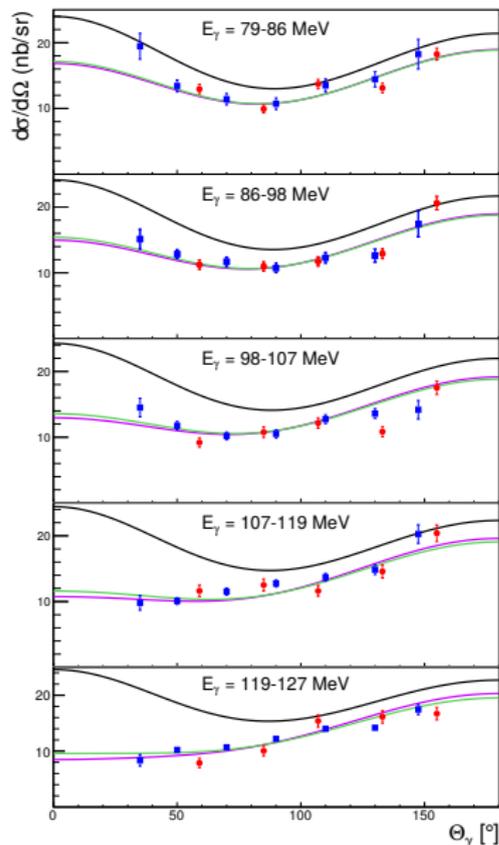
⇒ More than 200,000 Compton scattering events

Example of missing mass distribution



Good agreement between **PARA**, **PERP** and Monte carlo simulation
 \Rightarrow Low background dataset

Unpolarized cross-section



■ New DATA

● V. Olmos de Leon et al., EPJ A 10 (2001)

— Born contribution

— N. Krupina and V. Pascalutsa, PRL 110, 262001 (2013)

— B. Pasquini, D. Drechsel, and M. Vanderhaeghen, Phys. Rev. C 76 (2007)

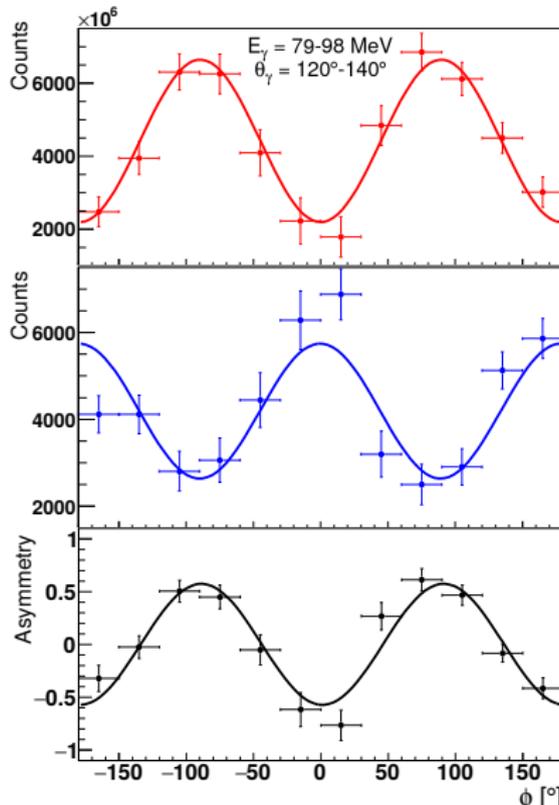
- Agreement with Olmos et al.
- Compatible (or better) statistics in the new data
- Significant deviation from Born contribution
- Extension of the angular coverage in the forward direction

Example of ϕ distribution

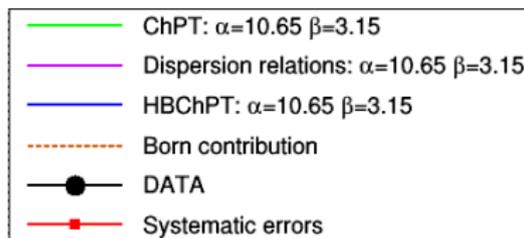
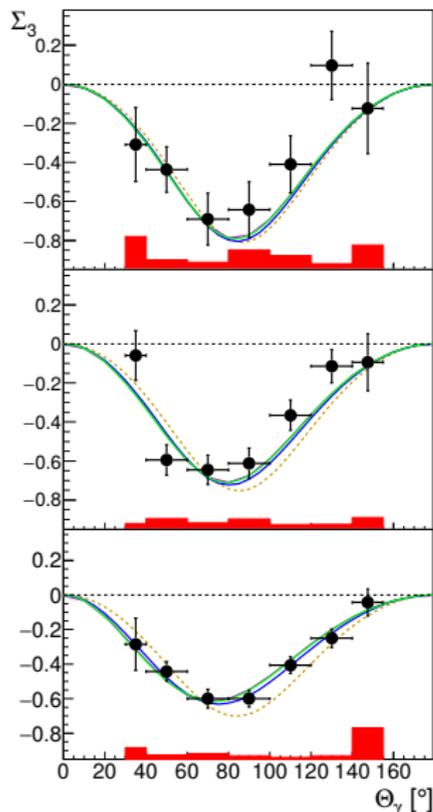
ϕ distribution for **PARA** and **PERP** data

$\cos(2\phi)$ modulation clearly seen

$$\frac{d\sigma}{d\Omega}(\theta, \phi) = \frac{d\sigma}{d\Omega}(\theta) [1 + p_\gamma \Sigma_3 \cos(2\phi)]$$



Σ_3 results



■ N. Krupina and V. Pascalutsa, PRL 110, 262001 (2013)

■ B. Pasquini, D. Drechsel, and M. Vanderhaeghen, Phys. Rev. C 76 (2007)

■ J. McGovern, D. Phillips, H. Griebhammer, EPJA 49, 12 (2013)

Fit on our Σ_3 results using Baldin sum rule constrain gives

BChPT framework:

$$\beta_{M1} = (2.8^{+2.3}_{-2.1}) 10^{-4} \text{ fm}^3$$

HBChPT framework:

$$\beta_{M1} = (3.7^{+2.5}_{-2.3}) 10^{-4} \text{ fm}^3$$

- **Proposed experiment: 600 hours of Compton scattering on proton:**
 - Measurement of the beam asymmetry Σ_3 and unpolarized cross-section
 - 10 cm LH₂ target and a polarized photon beam with $E_{e^-} = 883$ MeV
- **Improvement in statistics:**

	Diamond (both sets)		Copper radiator	
Experiment	Full target	Empty target	Full target	Empty target
Pilot	116 h	110 h	42 h	39 h
Proposed	500 h	70 h	30 h	-

- Tagger upgrade → improvement in rate ~ 5 times

⇒ Decrease of the statistical error ~ 3.5 times

- **Improvement in systematics:**

- Stable linear polarization with the new setup
- Improvement in tagger performance
- Continuous tagging efficiency monitor with pair spectrometer

⇒ Smaller systematic errors

Achieved and expected precision

BChPT fits from P. Martel			Error from ChPT fit (10^{-4} fm^3)					
			With Baldin			Without Baldin		
Experiment	Compton ev.		Σ_3	$\frac{d\sigma}{d\omega}$	$\Sigma_3, \frac{d\sigma}{d\omega}$	Σ_3	$\frac{d\sigma}{d\omega}$	$\Sigma_3, \frac{d\sigma}{d\omega}$
Pilot	≈ 200000	$\Delta\alpha_{E1}$	2.5	1.3	1.1	3.8	1.4	1.3
		$\Delta\beta_{E1}$				2.5	1.7	1.4
Proposed	≈ 4000000	$\Delta\alpha_{E1}$	0.7	0.4	0.3	1.1	0.4	0.4
		$\Delta\beta_{E1}$				0.7	0.5	0.4

⇒ Highest statistic data set: Olmos et al. ($\sim 50\%$ of the existing data):

	with Baldin	without Baldin
α_{E1}	12.1 ± 1.08	11.9 ± 1.39
β_{E1}	1.6 ± 0.89	1.2 ± 0.76

⇒ Word data errors (without double counting):

α_{E1}	0.76	β_{E1}	0.69
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⇒ **Precision improved compared to the existing data based on a single measurement!**

- **Successful pilot experiment!**

- Low background data set in the range 79 – 139 MeV
- First measurement of the beam asymmetry Σ_3 below pion threshold
- Alternative extraction of β_{M1} using ChPT and HBChPT frameworks
- Preliminary measurement of the unpolarized cross-section in agreement with theoretical calculations and existing data

- **Proposed experiment:**

- New measurement of the proton scalar polarizabilities with an unprecedentedly high precision
- Simultaneous measurement of the unpolarized cross-section and beam asymmetry
- Increased of the statistics and reduced of the systematic effects

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THANK YOU FOR YOUR ATTENTION!

BACKUP

⇒ The best way to extract β is the beam asymmetry Σ_3 :

$$\Sigma_3 = \frac{d\sigma_{\perp} - d\sigma_{\parallel}}{P_{\parallel} d\sigma_{\perp} + P_{\perp} d\sigma_{\parallel}}$$

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⇒ More than 200,000 Compton scattering events

$$\Sigma_3 = \Sigma_3^{(B)} - \frac{4M\omega^2 \cos \theta \sin^2 \theta}{\alpha_{em}(1 + \cos^2 \theta)^2} \beta_{M1} + O(\omega^4), \quad (6)$$

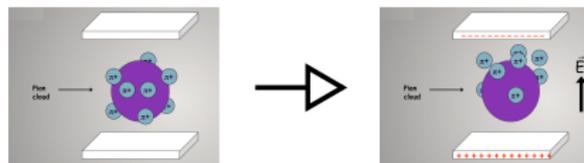
where $\Sigma_3^{(B)}$ is the pure Born contribution, while

$$\omega = \frac{s - M^2 + \frac{1}{2}t}{\sqrt{4M^2 - t}}, \quad \theta = \arccos \left(1 + \frac{t}{2\omega^2} \right) \quad (7)$$

are the photon energy and scattering angle in the Breit (brick-wall) reference frame. In fact, to this order in the LEX the formula is valid for ω and θ being the energy and angle in the lab or center-of-mass frame.

Proton scalar polarizabilities

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- Magnetic dipole moment: $\vec{m} = \beta_{M1} \vec{H}$
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