

Weyl quantization for gauge theories in terms of projective limits of graphs

[joint work w/ Thiemann]

Outline:

1) Motivation (time-dependent Born-Oppenheimer approx., (de)quantization,

↑ QFT on curved spacetimes from QG)

space-adiabatic perturbation theory [Panati, Spohn, Teufel; 2003]

2) The basic construction:

(local) strict deformation quantization of $T^*G \cong G \times \mathfrak{g}$
[Rieffel, Landsman] "Weyl type" ↑ compact Lie group

! \exists (global) version on $G \times \hat{G}$!
[Ruzhansky, Turunen] ↑ equivalence classes of unitary irrep's

3) A projective phase space for Yang-Mills theory

and its quantum counterpart

↑ equivariance of the Weyl quantisation

1) Motivation:

originally: How can we extract QFTs on curved spacetime from (L)QG?

→ mathematical framework?

(beyond $\hbar \rightarrow 0+$, Hamiltonian approach)

→ time dependent Born-Oppenheimer approx.?

(systematics beyond $\mathcal{O}(\hbar)$,

no fibered Hamiltonian)



$$H_\hbar = \int^{\oplus} d\mathfrak{s} H_0(\mathfrak{s}) + f(-i\hbar \nabla_x) \otimes \mathbb{1}$$

$$\mathcal{D}(H_\hbar) \subseteq L^2(\mathbb{R}^d, \mathcal{H}_f)$$

Idea: Use/adapt space-adiabatic perturbation theory

[Panati, Spohn, Teufel; 2003]

Basic ingredient: Pseudo-differential calculus (Weyl-Moyal formalism)

Space-adiabatic perturbation theory: (wish list!)

- **coupled** quantum dynamical system $\mathcal{H}, (\hat{H}, \mathcal{D}(\hat{H}))$
- splitting of the dynamics (controlled by parameter ε)

$$\mathcal{H} = \mathcal{H}_{\text{slow}} \otimes \mathcal{H}_{\text{fast}}$$

- deformation (**de**) quantization (ε -dependent)

$$\hat{\cdot}^{\varepsilon}: \mathcal{S}^{\infty}(\varepsilon; \Gamma, \mathcal{B}(\mathcal{H}_{\text{fast}})) \subset C^{\infty}(\Gamma, \mathcal{B}(\mathcal{H}_{\text{fast}})) \longrightarrow \mathcal{L}(\mathcal{H})$$

↑
symbols

↑
"phase space"

↙ symbol orders

$$\cdot H_{\varepsilon} \sim \sum_{k=0}^{\infty} \varepsilon^k H_k \quad ; \quad H_k \in \mathcal{S}^{\rho^{-k}}(\Gamma, \mathcal{B}(\mathcal{H}_{\text{fast}}))$$

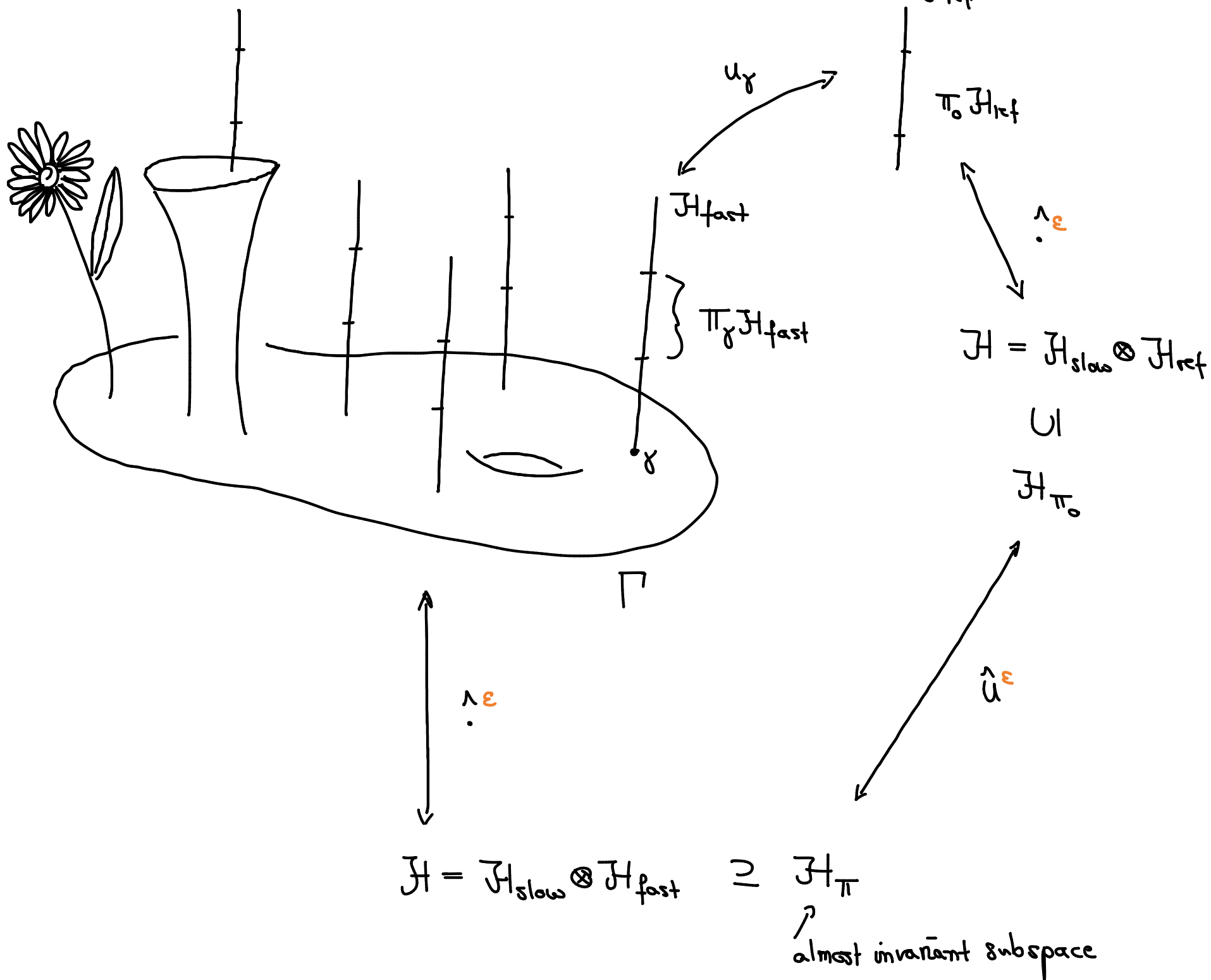
↑ asymptotic expansion (up to $\mathcal{S}^{-\infty}(\Gamma, \mathcal{B}(\mathcal{H}_{\text{fast}}))$)

↑ smoothing operators

$$\hat{H} = \hat{H}_{\varepsilon}^{\varepsilon}$$

- conditions on (point-wise) spectrum $\sigma(H_0) = \{ \sigma_0(H_0(\gamma)) \}_{\gamma \in \Gamma}$

Pictorially:



Upshot: Construct effective dynamics in \mathcal{H}_{π_0}

\uparrow

ϵ -independent space

2) The basic construction

- Γ will be modeled on T^*G \leftarrow compact Lie group G
(structure group of a principal bundle)
- pseudo-differential calculus on T^*G ?

↓

start from a (local) strict deformation quantization
[Rieffel, Landsman]

We have: $T^*G \cong G \times_{\mathfrak{g}}$; $\mathfrak{g} = \text{Lie}(G)$; $n = \dim(G)$
left/right translation

$\exp: \mathfrak{g} \rightarrow G$ gives a diffeomorphism for some

$$\mathfrak{g} \supseteq \underset{0}{U} \longrightarrow \underset{e}{V} \subseteq G$$

relate the Haas measure on G and Lebesgue measure on \mathfrak{g}

$$f \in C_c(V): \int_V f(g) dg = \int_U f(\exp(x)) j(x)^2 dx$$

$$H \in \mathfrak{t} \subseteq \mathfrak{g}: j(H) = \prod_{\alpha \in R^+} \frac{\sin(\alpha(H))}{\alpha(H)}$$

Lie algebra of a max. torus ← positive roots

Construct operators in $L^2(G)$ from convolution kernels:

$$\varepsilon \in (0, 1]: \quad \mathcal{F}_\varepsilon^\varepsilon(h, g) = \int_{\mathfrak{g}^*} \sigma(\theta, g) e^{\frac{i}{\varepsilon} \theta(X_h)} \frac{d\theta}{(2\pi\varepsilon)^n}$$

for $X_h = \exp^{-1}(h)$; $\sigma \in C_{PW, U_\varepsilon}^\infty(\mathfrak{g}^*) \hat{\otimes} C^\infty(G)$; $U_\varepsilon = \varepsilon^{-1}U$

$$\Rightarrow \mathcal{F}_\varepsilon^\varepsilon \in C^\infty(G) \hat{\otimes} C^\infty(G)$$

\uparrow
gives the kernel of a Kohn-Nirenberg Ψ DO

Deform construction to a Weyl-type Ψ DO:

• define a (smooth) square root $\sqrt{\cdot}: V \subseteq \mathfrak{g} \longrightarrow V \subseteq \mathfrak{g}$
 $g \longmapsto \sqrt{g} := \exp\left(\frac{1}{2} X_g\right)$

$$\bullet \quad \mathcal{F}_\varepsilon^{W, \varepsilon}(h, g) = \mathcal{F}_\varepsilon^\varepsilon(h, \sqrt{h^{-1}}g)$$

Weyl quantisation:

$$\forall \psi \in C^\infty(G): \quad (\mathcal{Q}_\varepsilon^W(\sigma)\psi)(g) = \int_G \mathcal{F}_\varepsilon^{W, \varepsilon}(h, g) \psi(h^{-1}g) dh$$

Theorem: [Landsman]

$$C(T^*G) \xrightarrow{C} C(G) \rtimes_L G$$

$$Q_\varepsilon^W: C_{PW, \mu}^\infty(\mathfrak{g}^*) \otimes C^\infty(G) \longrightarrow \mathcal{K}(L^2(G))$$

is a nondegenerate strict deformation quantization on $(0, 1]$ w.r.t. the canonical Poisson structure on T^*G .

$$\sigma, \tau \in C^\infty(T^*G): \{\sigma, \tau\}_{T^*G} = \langle \partial_\theta \sigma, R\tau \rangle - \langle R\sigma, \partial_\theta \tau \rangle + \{\sigma, \tau\}_-$$

(minus) Lie-Poisson structure

Remark: For special choices of σ , we find

"CCR" relations:

$$\sigma_f(\theta, g) = f(g) \quad ; \quad f \in C^\infty(G)$$

$$\sigma_X(\theta, g) = \int_X^{L^*(\cdot)^{-1}}(\theta, g) \quad ; \quad X \in \mathfrak{g}$$

bi-equivariant momentum map of the strongly Hamiltonian left pullback action of G on T^*G

$$Q_\varepsilon^W(\{\sigma_f, \sigma_g\}_{T^*G}) = \frac{i}{\varepsilon} [Q_\varepsilon^W(\sigma_f), Q_\varepsilon^W(\sigma_g)] = 0$$

$$Q_\varepsilon^W(\{\sigma_X, \sigma_f\}_{T^*G}) = \frac{i}{\varepsilon} [Q_\varepsilon^W(\sigma_X), Q_\varepsilon^W(\sigma_f)] = R_X f$$

$$Q_\varepsilon^W(\{\sigma_X, \sigma_Y\}_{T^*G}) = \frac{i}{\varepsilon} [Q_\varepsilon^W(\sigma_X), Q_\varepsilon^W(\sigma_Y)] = i\varepsilon R_{[X, Y]}$$

Define Ψ DO on $C^\infty(G) \subseteq L^2(G)$ via $\mathcal{Q}_\varepsilon^\omega$



Problem: $\exp: \mathfrak{g} \rightarrow G$ is not a diffeomorphism

We need to deal with support properties of $\sigma \in C^\infty(T^*G)$

Solution: Exploit Paley-Wiener-Schwartz characterisation of distributions w/ compact support

Definition: $K \subseteq \mathfrak{g}$ convex, compact

$m \in \mathbb{R}$; $0 \leq \delta \leq \rho \leq 1$

A function $\sigma \in C^\infty(\mathfrak{g}^*, C^\infty(G))$ belongs to the space of Paley-Wiener-Schwartz symbols $S_{PW, \delta, \rho}^{k, m}$ if:

1) $\sigma: \mathfrak{g}^* \rightarrow C^\infty(G)$ is (weakly) holomorphic

ii) $\forall \alpha, \beta \in \mathbb{N}_0^n: \exists C_{\alpha\beta} > 0: \forall \theta \in \mathfrak{g}^*:$

$$\sup_{g \in G} |(R^\alpha \partial_\theta^\beta \sigma)(\theta, g)| \leq C_{\alpha\beta} \langle \theta \rangle^{m - |\beta|\rho + |\alpha|\delta} e^{H_K(\operatorname{Im}(\theta))}$$

supporting function of K

↑
standard regularized distance

We have the expected inclusions:

$$S_{\rho\omega, p, \delta}^{k, m} \subset S_{\rho\omega, p', \delta}^{k, m} \quad ; \quad p \geq p'$$

$$S_{\rho\omega, p, \delta}^{k, m} \subset S_{\rho\omega, p, \delta}^{k', m} \quad ; \quad k \leq k'$$

$$S_{\rho\omega, p, \delta}^{k, m} \subset S_{\rho\omega, p, \delta'}^{k, m} \quad ; \quad \delta \leq \delta'$$

$$S_{\rho\omega, p, \delta}^{k, m} \subset S_{\rho\omega, p, \delta}^{k, m'} \quad ; \quad m \leq m'$$

Set: $S_{\rho\omega, p, \delta}^{k, \infty} := \bigcup_{m \in \mathbb{R}} S_{\rho\omega, p, \delta}^{k, m} \quad ; \quad S_{\rho\omega}^{k, -\infty} := \bigcap_{m \in \mathbb{R}} S_{\rho\omega, p, \delta}^{k, m}$

Corollary: $\sigma \in S_{\rho\omega, p, \delta}^{k, m}$ defines a continuous operator

$\nearrow Q_\varepsilon^w(\sigma) : C^\infty(G) \longrightarrow C^\infty(G)$ for some $0 < \varepsilon_0 \leq 1$
 extends to $\mathcal{D}'(G)$ and all $\varepsilon \in (0, \varepsilon_0]$.
 for $m = -\infty$ (smoothing)

Proposition: The optimal constants $C_{\alpha\beta} > 0$ turn the symbol spaces $S_{\rho\omega, p, \delta}^{k, m}$ into Fréchet spaces.

The strict inductive limit topology turns $S_{\rho\omega, p, \delta}^{k, \infty}$ into an LF-space.

Lemma: Differentiation, multiplication, and the Poisson bracket become continuous maps between suitable symbol spaces.

Theorem: (Asymptotic completeness of Paley-Wiener-Schwartz symbols)

The symbol spaces are asymptotically complete in the following sense:

Given $\{m_k\}_{k=1}^{\infty} \subset \mathbb{R}$ s.t. $\lim_{k \rightarrow \infty} m_k = -\infty$ and $m := \max_k m_k$,

and $\sigma_k \in C^{\infty}(G, S_{\rho, p, \delta}^{k, m_k})$, there exists $\sigma \in C^{\infty}(G, S_{\rho, p, \delta}^{k, m})$

such that:

$$\sigma \sim \sum_{k=1}^{\infty} \sigma_k, \text{ i.e.}$$

↑
resummation

$$\forall M \in \mathbb{R} : \exists k_0 \in \mathbb{N} : \forall k' \geq k_0 : \sigma - \sum_{k=1}^{k'} \sigma_k \in S_{\rho, p, \delta}^{k, M}$$

(uniqueness up to smoothing symbols)

Remark: Non-trivial constructions necessary because of holomorphicity.

↳ no excision functions!

→ kernel cut-off operators (Volterra-Mellin 4DOs, [Kraimer])

Observations: ε -Expansion of the (formal) identity

$$Q_\varepsilon^W(\sigma) Q_\varepsilon^W(\tau) = Q_\varepsilon^W(\sigma *_\varepsilon \tau)$$

becomes an asymptotic series for suitable

$$\sigma \in S_{\mathbb{R}^n, \rho, \delta}^{k, m}, \quad \tau \in S_{\mathbb{R}^n, \rho', \delta'}^{k', m'}$$

But: Is there a resummation of $\sigma *_\varepsilon \tau$ s.t.
the identity is implemented?

- restrictions for compositions of symbols (ρ, δ)

$$\sigma *_\varepsilon \tau \in S_{\mathbb{R}^n, \rho, \delta}^{k+k', m+m'}$$

- fas:
- $\rho > \delta$ (similar to \mathbb{R}^{2n} -case)
 - $\rho > \frac{1}{2}$ (from Lie-Poisson structure)

3) A projective phase space for Yang-Mills theory

$$G \curvearrowright P \xrightarrow{\pi} \Sigma \quad \text{right, semi-analytic, principal } G\text{-bundle}$$

↑
Cauchy surface

Field configurations in terms of **initial data**

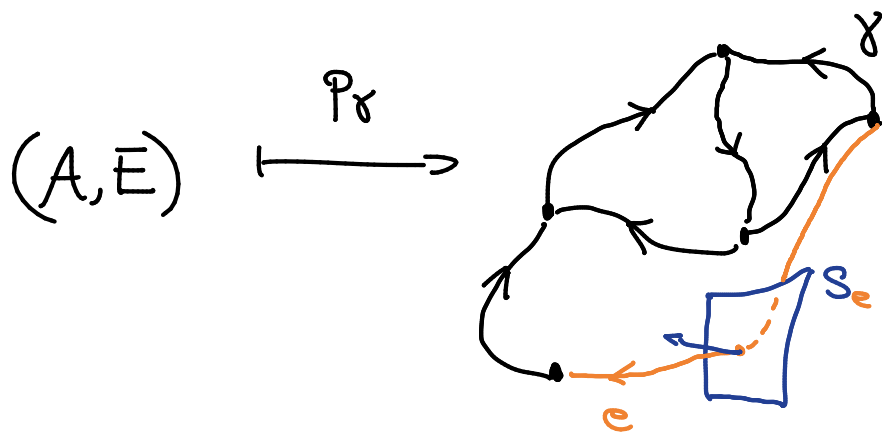
$$(A, E) \in T^*A_P \text{ or } |\Lambda^1|T^*A_P =: \Gamma$$

↑
1-densities

Proceed by construction of finite dimensional projections

$$\Gamma \xrightarrow{P_\gamma} \Gamma_\gamma$$

by localizing (A, E) on finite, compactly supported, semi-analytic oriented, embedded graph $\gamma \in \Gamma_0^{\text{sa}, \uparrow}$



$g_e(A)$ - parallel transport along e

$P_X^e(A, E)$ - flux through S_e weighted by $X \in \mathfrak{g}$

Observation: • Using a suitable regularisation scheme, the symplectic structure on Γ descends to the natural symplectic structure on Γ_γ .

$$(\Gamma, \omega) \xrightarrow{P_\gamma} (\Gamma_\gamma = T^*G^{\times |E(\gamma)|}; \omega_{T^*G}^{\times |E(\gamma)|})$$

is **symplectic**.

• the functionals $(g_e, P_x^e)_{e \in E(\gamma)}$ behave naturally w.r.t. operations on graphs:

$$i) e = e_2 \circ e_1 : g_e(A) = g_{e_2}(A) g_{e_1}(A)$$

(composition)

$$ii) e \mapsto \bar{e} : g_{\bar{e}}(A) = g_e(A)^{-1}$$

$$P_x^{\bar{e}}(A, E) = -P_{\text{Ad}_{g_e(A)}(x)}^e(A, E)$$

(inversion)

Composition is subtler for P_x^e :

$$P_x^e(A, E) = c P_x^{e_2}(A, E) + (1-c) P_{\text{Ad}_{g_{\bar{e}_2}(A)}(x)}^{e_1}(A, E)$$

compatible w.r.t. right action, but only $c=0,1$ valid.

iii) $e \rightarrow \emptyset$: drop the corresponding factors in $\Gamma_\gamma = T^*G^{\times |E(\gamma)|}$
 (removal)

\Rightarrow • $\{\Gamma_\gamma\}_{\gamma \in \Gamma_{\text{sa}, \uparrow}}$ becomes a directed set w.r.t.
 inclusion of oriented, y^* -labeled subgraphs.

• the relative projections $P_{\gamma\gamma'} : \Gamma_{\gamma'} \rightarrow \Gamma_\gamma$; $\gamma \subseteq \gamma'$
 are symplectic surjections. $P_{\gamma\gamma'}^* : C^\infty(\Gamma_\gamma) \rightarrow C^\infty(\Gamma_{\gamma'})$ is continuous.

\uparrow
 this imposes $c=0,1$. We need to choose coherently.
 but breaks reversion symmetry.

• the relative projections $\{P_{\gamma\gamma'}\}_{\gamma \subseteq \gamma'}$ induce injective $*$ -morphisms
 by Weyl quantisation

$$Q_\varepsilon^W(P_{\gamma\gamma'}^*) =: \alpha_{\gamma\gamma'} : L(C^\infty(\Gamma_\gamma)) \longrightarrow L(C^\infty(\Gamma_{\gamma'}))$$

\uparrow
 continuous linear operators on $C^\infty(\Gamma_{\gamma'})$

Viewing elements of $L(C^\infty(\Gamma_\gamma))$ in terms of Schwartz kernels,
 the action on $C^\infty(\Gamma_\gamma)$ comes from the

integrated left regular action ρ_L
 $(Q_\varepsilon^W = \rho_L \circ FW, \varepsilon)$

Some inductive constructions:

We would like to construct an inductive-limit $*$ -algebra to represent the noncommutative analog of Π .

• First try: $\mathcal{O}_\gamma = C(C_\gamma) \rtimes_L C_\gamma \cong \mathcal{K}(L^2(C_\gamma))$
 \parallel
 $\pi_1(\Gamma_\gamma) \cong \mathbb{G}^{|\mathbb{E}(\gamma)|}$

⚡ to small, cannot support the $*$ -morphisms $\{\alpha_{\gamma'\gamma}\}_{\gamma \subseteq \gamma'}$

• Second try: $\mathcal{B}_\gamma := M(\mathcal{O}_\gamma)$
↑
multiplies algebra

✓ works and has some nice extension properties

regarding \mathcal{O}_γ :
i) unique extension of morphisms
ii) embedding of $C(C_\gamma)$ and C_γ

iii) recovery of $S(\mathcal{O}_\gamma)$
↑
state space

as strictly continuous states on \mathcal{B}_γ

$$\mathcal{B} := C^*\text{-}\lim_{\gamma \subseteq \gamma'} \mathcal{B}_\gamma \quad (\text{also } W^*\text{-}\lim \text{ possible})$$

Problems / Questions:

- different extension of \mathcal{O}_γ ? K -theory of $C(C_\gamma) \rtimes_{\text{L}} C_\gamma$?
- control on the state space of the inductive limit?



The only known state comes from gives the inductive-limit Hilbert space:

$$L^2(\overline{A}) := \varinjlim_{\gamma \subseteq \gamma'} L^2(C_\gamma)$$

(spin network Hilbert space)

Final remark:

- gauge transformations act on vertices of γ :

$$\alpha_\gamma^\lambda(\{g_v\}_{v \in V(\gamma)}) (\mathcal{F}) (\{(h_e, g_e)\}_{e \in E(\gamma)})$$

$$= \mathcal{F}(\{\alpha_{g_{e(n)}}^{-1}(h_e), g_{e(n)}^{-1} g_e g_{e(o)}\}_{e \in E(\gamma)})$$

in a compatible way.