

# *Gyromagnetic Faraday Rotation Studies on Polarized $^3\text{He}$*

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Low-energy Probes of New Physics

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MITP, Mainz

# Outline

- Motivation:

New idea to search for Asymmetric Dark Matter?

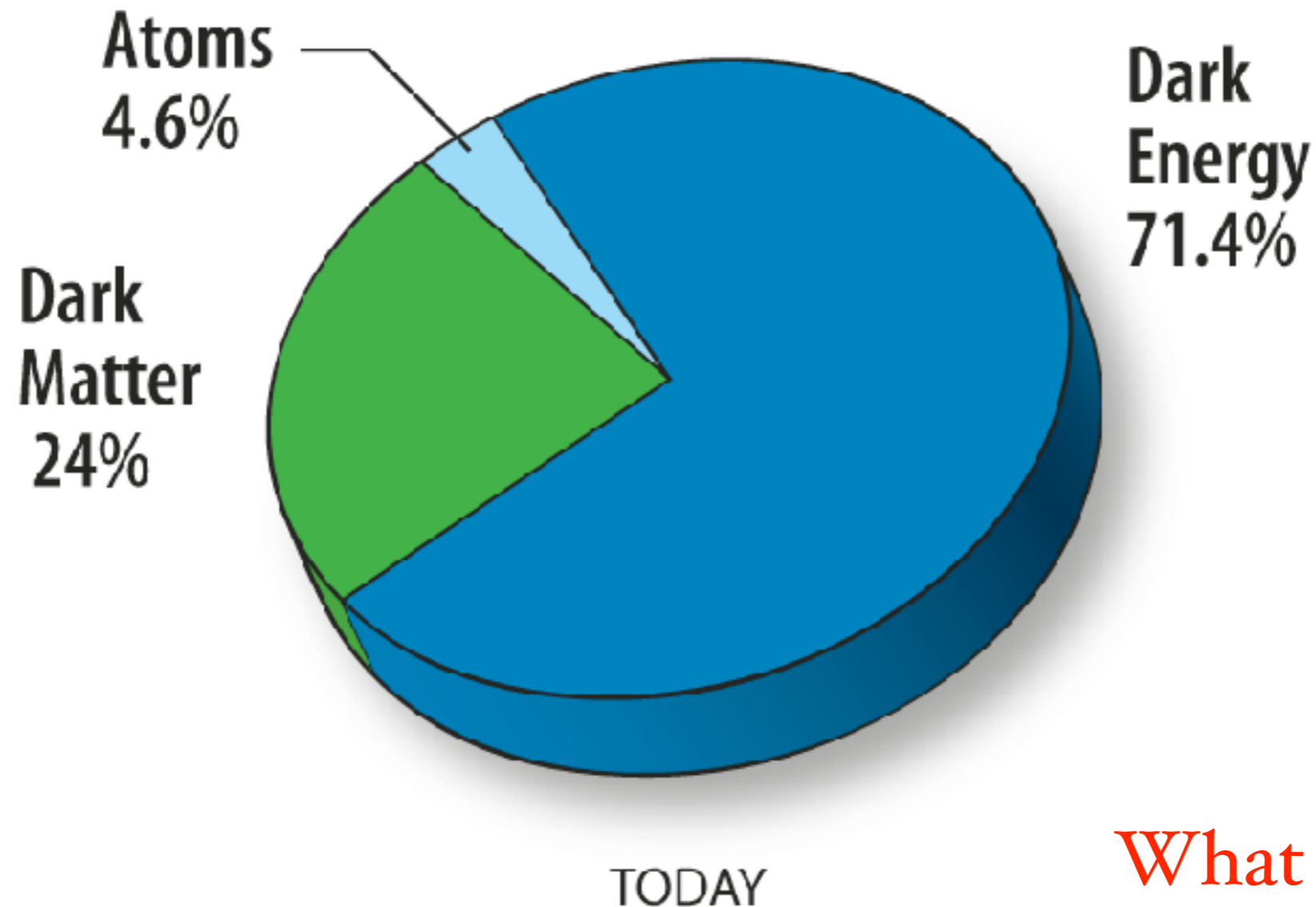
- Test case:

Gyromagnetic Faraday rotation on polarized  $^3\text{He}$

- Summary

# A Dark-Dominated Universe

Most of the cosmic energy budget is in unknown forms!

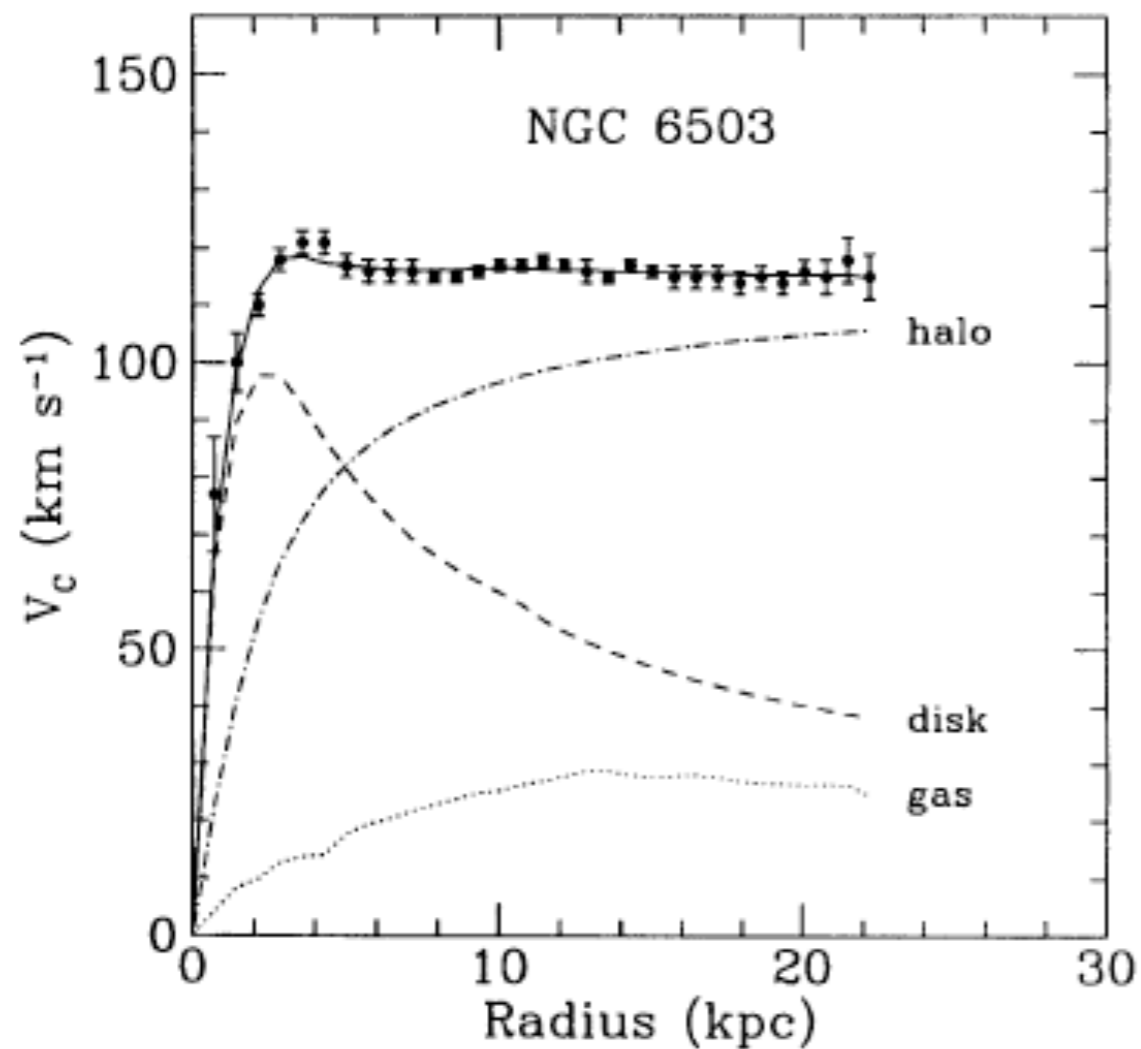


What are they?!

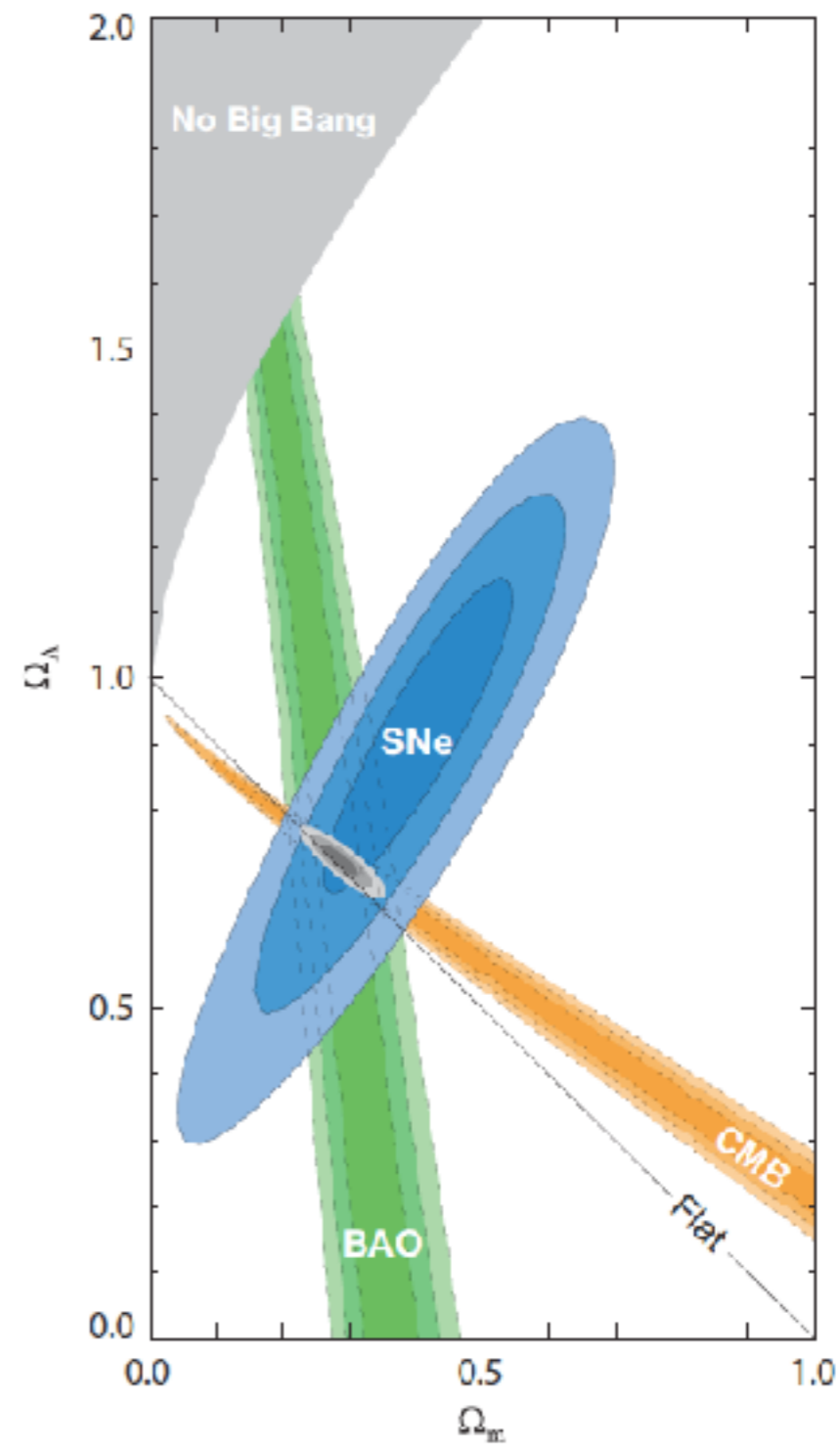
[[map.gsfc.nasa.gov/universe/uni\\_matter.html](http://map.gsfc.nasa.gov/universe/uni_matter.html)]

# Observational Evidence for Dark Matter

**Galactic Rotation Curves:**  
[e.g., from Begeman, Broeils, and Sanders, 1991]



**The observed circular speed does not track the luminous mass.**



**Most of the cosmic energy budget is of an unknown form!**

# Dark Matter (DM) Knowns

## We do know DM must be...

- stable or effectively on Gyr time scales
- not “hot” - i.e., not relativistic at the time it decoupled from matter in the cooling early Universe
- have no substantial strong or electromagnetic charge

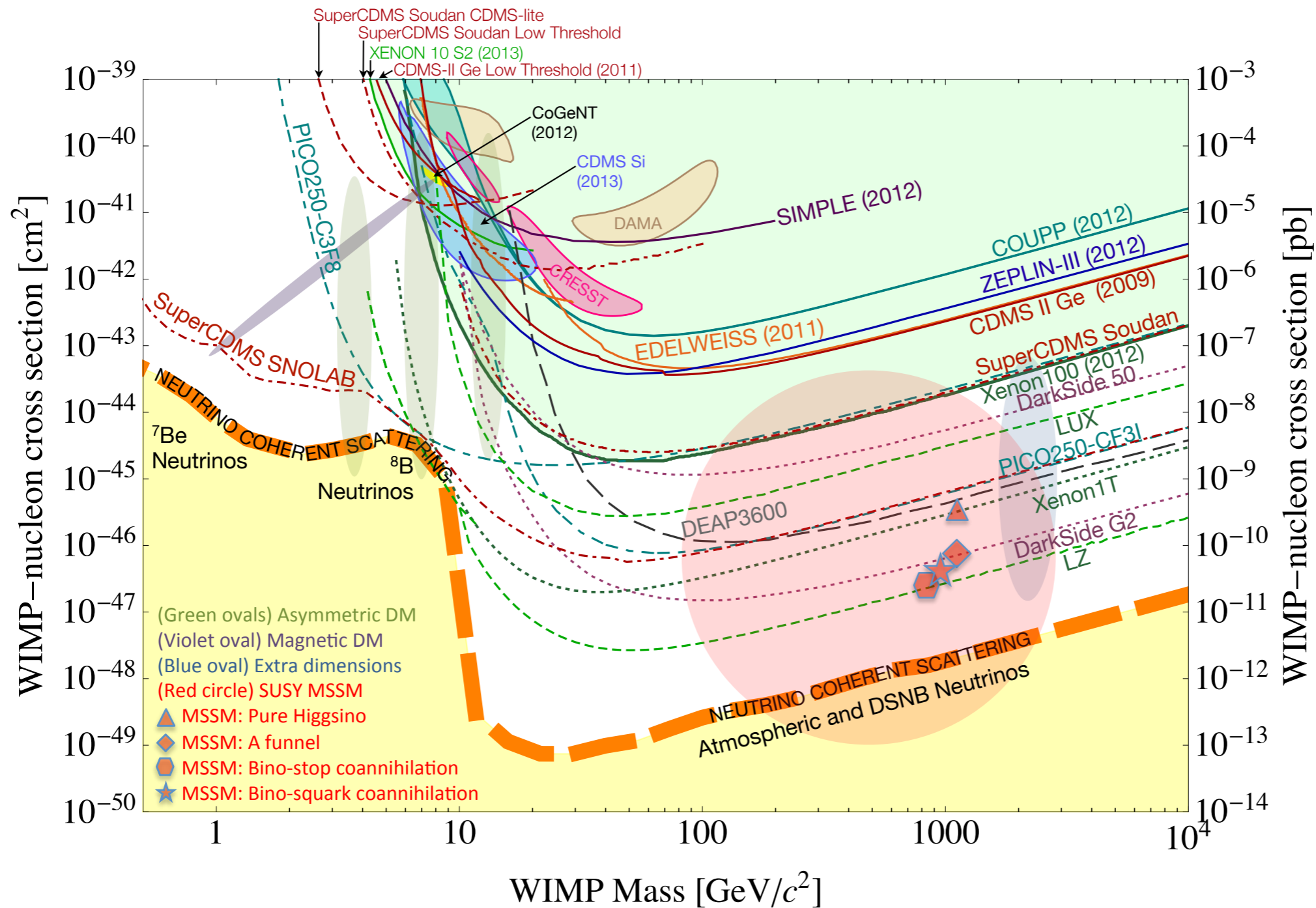
It has long been thought that if DM were produced as a “thermal relic” that it would be a Weakly Interacting Massive Particle or “WIMP”

**Such candidates appear in models with weak-scale supersymmetry (MSSM)**

They can be detected directly in low-background experiments

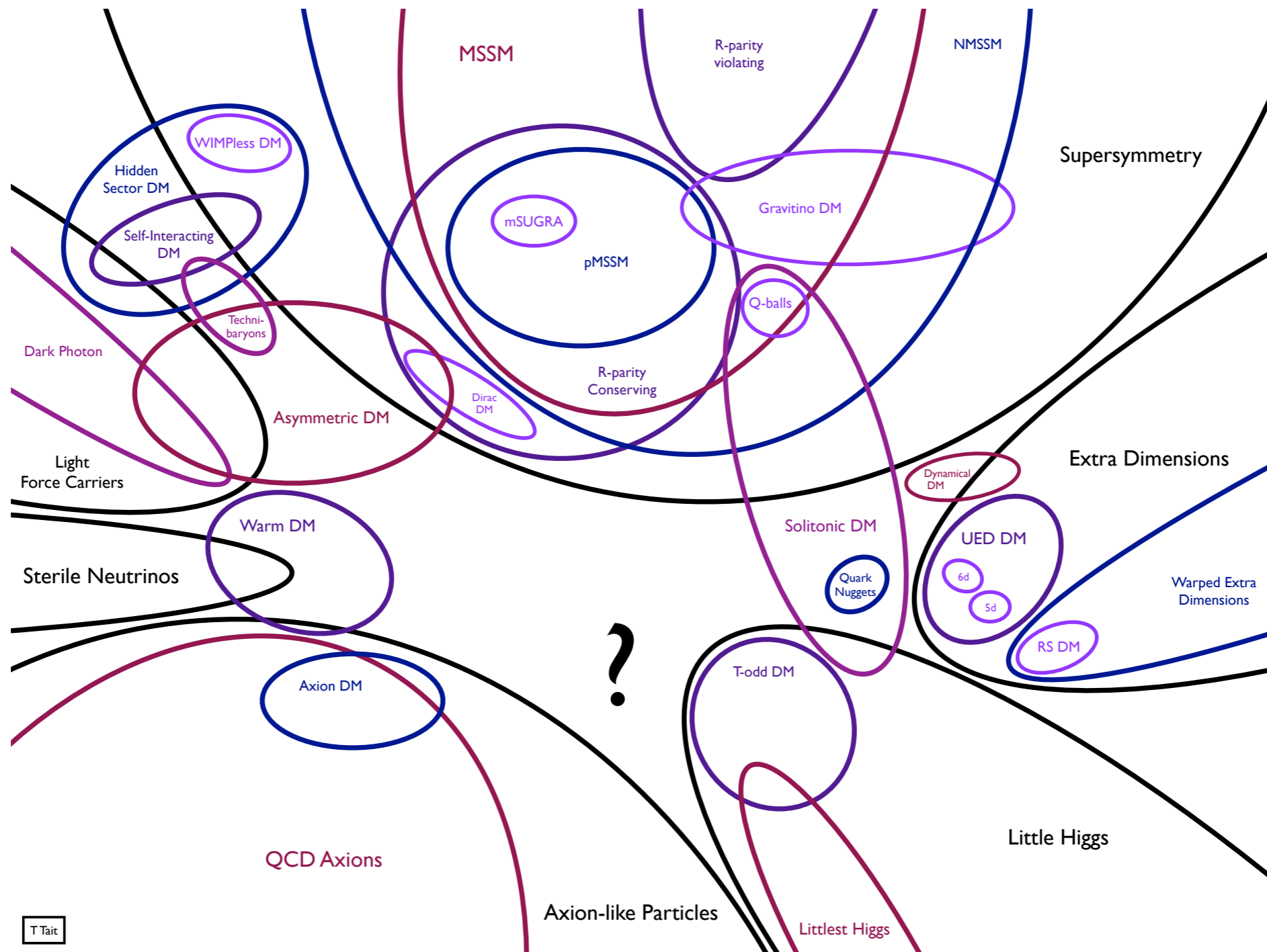
# Direct Detection: Dark Matter “WIMPs”

Limits rely on local DM density and velocity distribution



[arXiv:1401.6085]

# Theory: Possible DM Candidates



[arXiv:1401.6085]

# What if DM is not a WIMP?

Its relic density need not be fixed by thermal freeze-out, and its stability need not be fixed by a discrete symmetry.

**What mechanisms then are operative and how do we discover them?**

- Its stability may be guaranteed by a hidden conserved charge, much as the electron in the SM is stable.

*A complex hidden sector?*

- Its relic density may be related to the cosmic baryon asymmetry.

*The hidden sector could be asymmetric, as baryons are?*

- [axions]....

[Nussinov, PLB 1985; Barr, Chivukula, Farhi, PLB 1990; Harvey and Turner, PRD 1990; Ellis et al., NPB 1992. Rytrov and Sannino, arXiv:0809.0713 [hep-ph]; Kaplan, Luty, Zurek, arXiv:0901.4117 [hep-ph].]



# Asymmetric Dark Matter (ADM)

## Experimental Signatures

- ADM models can give distinctive collider signatures
  - E.g. long-lived metastable states, new charged states at the weak scale, and/or colored states at a TeV. (All not seen as yet.)
- Direct detections signals can arise from interactions that could
  - \* eliminate the symmetric DM component
  - \* transfer the fermion asymmetry

These could be realized through magnetic moment or charge radius couplings.

**Both interactions can give rise to anomalous nuclear recoils....**

[Bagnasco, Dine, and Thomas, PLB 1994; Barger, Keung, Marfatia, arXiv:1007.4345; Banks, Fortin, and Thomas, arXiv:1007.5515]

**A magnetic Faraday effect can also discover dark matter if it possesses a magnetic moment.. and establish asymmetric dark matter.**

S. Gardner, PRL 100, 041303 (2008)

S. Gardner, PRD 79, 055007 (2009)

some reviews: K.M Zurek, Phys. Rep. 537, 91 (2014), D.E. Kaplan et al., PRD 79, 115016 (2009), ...

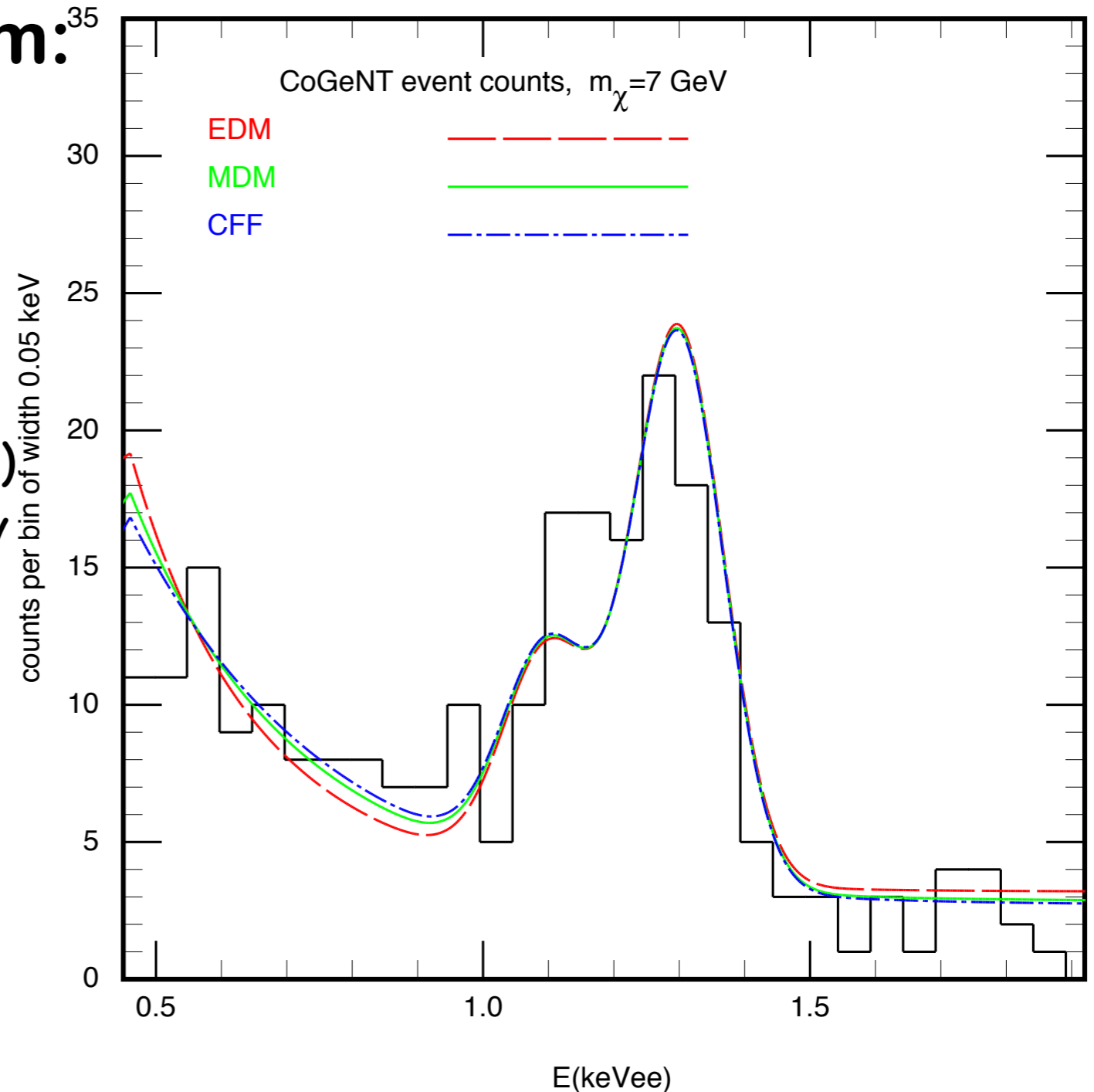
# Experimental Signatures of ADM

## Impact on recoil spectrum: here CoGent

$$m_{\text{DM}} = 7 \text{ GeV}$$

$$\text{EDM: } d_{\chi} = 10^{-20} \text{ e}\cdot\text{cm} \quad (\Lambda_{\text{EDM}} = 1.97 \text{ PeV})$$

$$\text{MDM: } g_{\chi S} = 0.00454 \quad (\Lambda_{\text{MDM}} = 3.09 \text{ TeV})$$



V. Barger et al., arXiv: 1007.4345v2

# Constraints on Magnetic Moments of DM

Properties of DM: Possibility of (anomalous) magnetic moments

Constraints for  $M_{\text{DM}} \leq 1 \text{ MeV}$ :

- precision electroweak measurements  $\rightarrow$  e.w. radiative corrections

anomalous mag. moment ( $\kappa$ ):  $|\kappa| < 4 \times 10^{-6}$  for  $M_{\text{DM}} = m_e$

- $e^+e^- \rightarrow \nu\bar{\nu}\gamma \rightarrow \mu_{\nu\tau} \sim \text{few} \times 10^{-6} \mu_B$

J. Erler and P. Langacker, review in W.M. Yao et al. (Particle Data Group), J. Phys. G 33, 1 (2006)

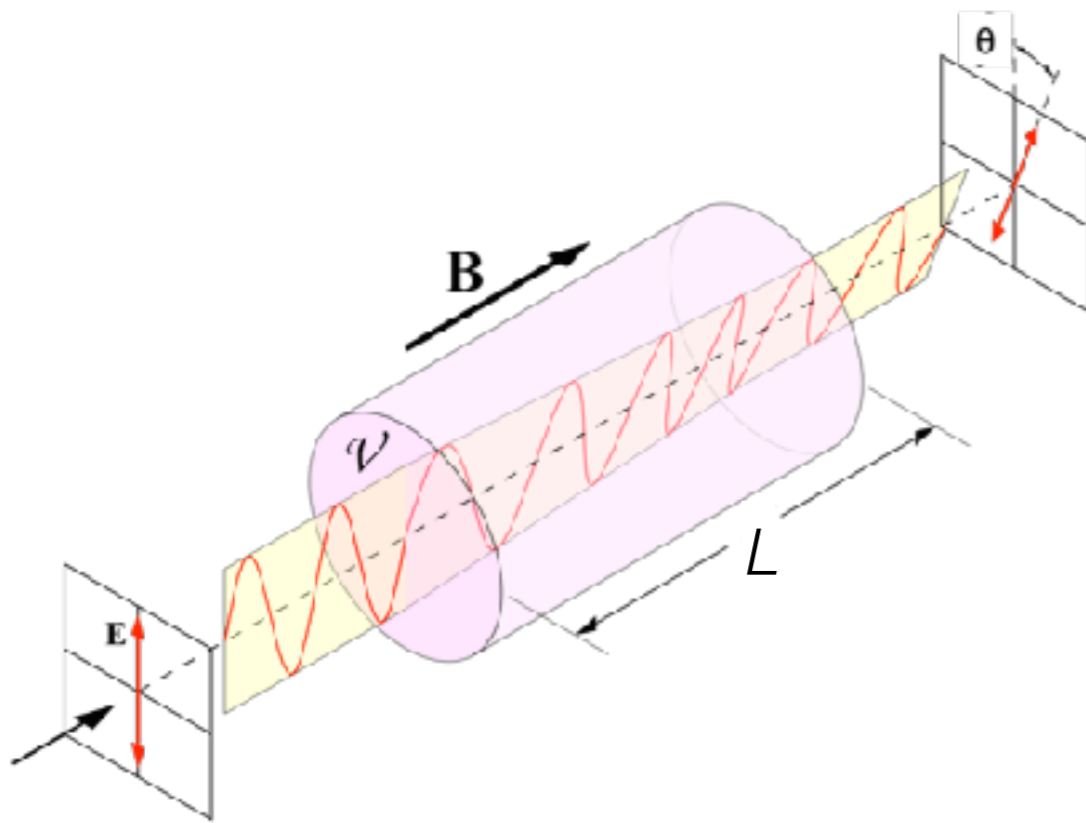
H. Grotch and R.W. Robinett, Z. Phys. C 39, 553 (1988)

**Can Faraday rotation be used to probe  $\mu_{\text{DM}}$ ?**

# Introduction to Faraday Rotation

Faraday Effect: 1846 !!  
Optical activity is induced in matter by magnetic field

Faraday, M., Phil. Mag., 28, 294 (1846) ;  
Phil. Trans. Roy. Soc. London, 1 (1846)



Angle of rotation:

$$\theta = V \cdot B \cdot L = V \int_0^L B dl$$

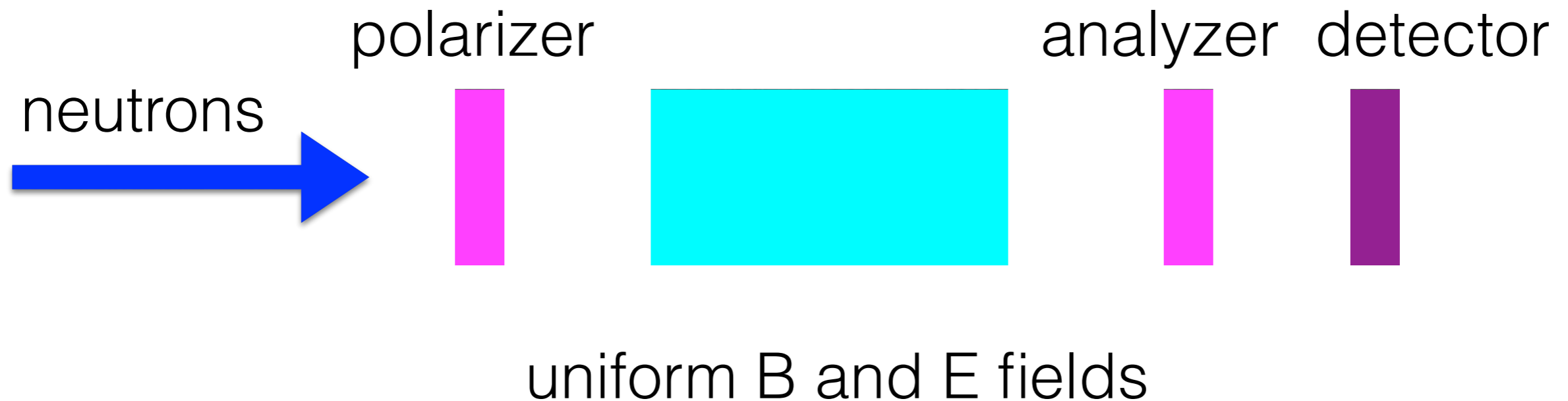
$V$ : Verdet Constant  $\rightarrow$  Contains physics



[https://en.wikipedia.org/wiki/Faraday\\_effect#/media/File:Faraday-effect.svg](https://en.wikipedia.org/wiki/Faraday_effect#/media/File:Faraday-effect.svg)

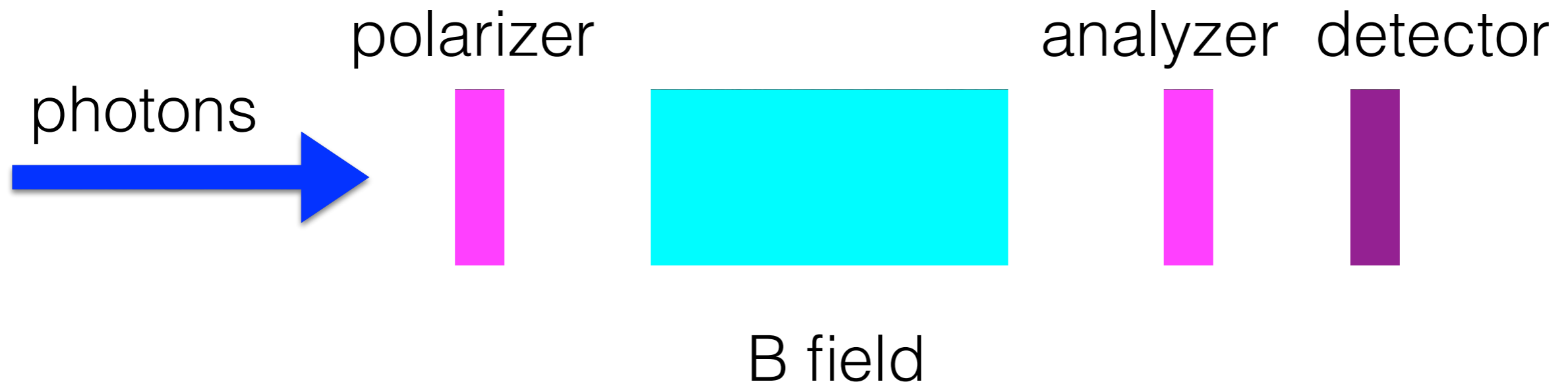
# Measurement Principle

D. Beck's talk: Principle of measuring neutron EDMs



# Measurement Principle

Principle of measuring Faraday rotation



# Similar to PVLAS Experiment

## PVLAS: “Testing the polarization of the vacuum with lasers”

Department of Physics and [National Institute of Nuclear Physics](#) in [Ferrara, Italy](#)

Nonlinear effects due to  
magnetic birefringence  
(of vacuum)

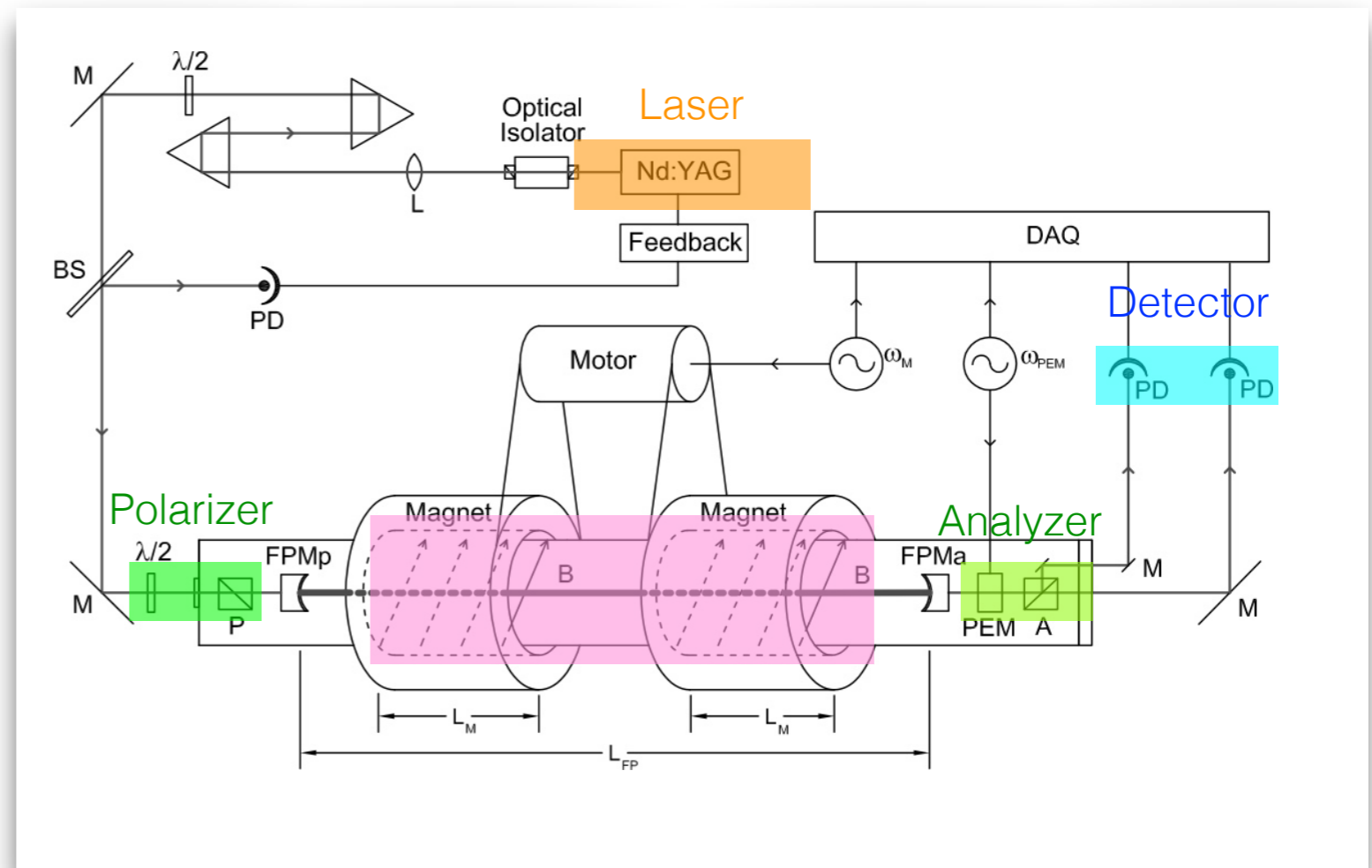
→ sensitivity to axions

$$\Delta n = 3A_e B_{ext}^2$$

“nonlinearity” const.

no effect observed:

$$A_e < 2.9 \cdot 10^{-21} T^{-2}$$



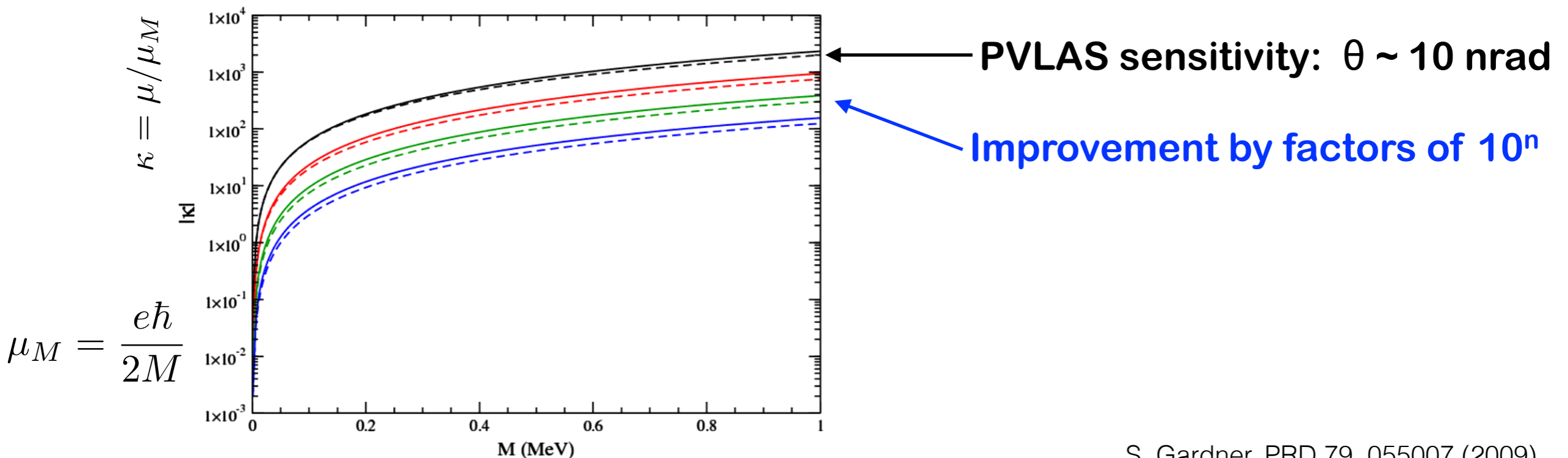
Total sensitivity:  $\Delta\theta \sim 10$  nrad

G. Zavattini et al, Int. J. Mod. Phys. A 27, 1260017 (2012)

# Sensitivity to Magn. Moments of Light DM



Possible anomalous magnetic moment vs.  $m_{DM}$



S. Gardner, PRD 79, 055007 (2009)



# Matter in $\mathcal{E}$ . & $\mathcal{M}$ . Fields

Matter in e.m fields: dielectric displacement and magnetic induction

$$\mathbf{D} = \epsilon \mathbf{E}$$

$$\mathbf{B} = \mu \mathbf{H}$$

$$\epsilon = 1 + 4\pi N \alpha \quad \rightarrow \text{dielectric constant}$$

$$\mu = 1 + 4\pi N \chi \quad \rightarrow \text{permeability}$$

# Magnetic Moments in $\mathcal{E}\&\mathcal{M}$ Fields

Matter in e.m fields: dielectric displacement and magnetic induction

$$\mathbf{D} = \epsilon\mathbf{E} + ig\mathbf{H} + if_1\mathbf{E} \times \mathbf{H}' + k\mathbf{H} \times \mathbf{H}' \quad \mathbf{E}, \mathbf{H} \rightarrow \text{e.m. wave}$$
$$\mathbf{B} = \mu\mathbf{H} - ig\mathbf{E} + if_2\mathbf{H} \times \mathbf{H}' - k\mathbf{E} \times \mathbf{H}' \quad \mathbf{H}' \rightarrow \text{static field}$$

$$\epsilon = 1 + 4\pi N\alpha \quad \rightarrow \text{dielectric constant}$$

$$\mu = 1 + 4\pi N\chi \quad \rightarrow \text{permeability}$$

$$g = \omega 4\pi N\beta \quad \rightarrow \text{natural gyration constant}$$

$$\left. \begin{aligned} f_1 &= \omega 4\pi N\eta \\ f_2 &= \omega 4\pi N\xi \end{aligned} \right\} \rightarrow \text{magneto-optical gyration constants}$$

$$k = 4\pi N\delta \quad \rightarrow \text{contribution to polarizability and birefringence in optically active molecules}$$

# Magnetic Moments in $\mathcal{E}\&\mathcal{M}$ Fields

Precession of  $\mathbf{M}$  in  $\mathbf{B}_0$ : 
$$\frac{d\mathbf{M}}{dt} = \frac{g\mu_M}{\hbar} \mathbf{M} \times \mathbf{B}_0$$

Relativistic eq. (BMT eq. for neutral particles)

$$\gamma \frac{d\mathbf{M}}{dt} = \frac{g\mu_M}{\hbar} (\mathbf{M} \times \mathbf{B}_0 + \mathbf{M}_0 \times (\mathbf{B} - \boldsymbol{\beta} \times \frac{\mathbf{E}}{c}))$$

For LH and RH circularly polarized light:

$$M_{\pm} = \pm \frac{\omega_M(1 + \beta)}{\gamma\omega \pm \omega_B} B_{\pm} \equiv \chi_{\pm} B_{\pm}$$

$$n_{\pm} = \sqrt{1 + \chi_{\pm}}$$

$$\omega \gg \omega_B$$

$$n_- - n_+ = \frac{\omega_M(1 + \beta)}{\gamma c}$$

$$\omega_M = g\mu_M M_0 / \hbar$$

$$\theta = \frac{g\mu_M M_0 L}{2\hbar c}$$

→ gyromagnetic optical rotation  
independent of frequency

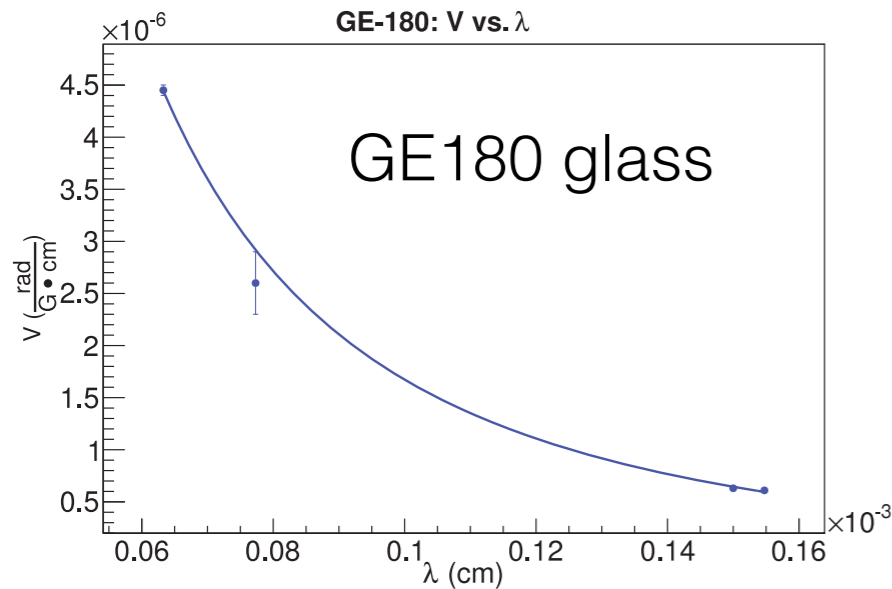
D. Polder, Philos. Mag. 40, 99 (1949)

C.L. Hogan, Rev. Mod. Phys. 25, 253 (1953)

S. Gardner, PRL 100, 041303 (2008)

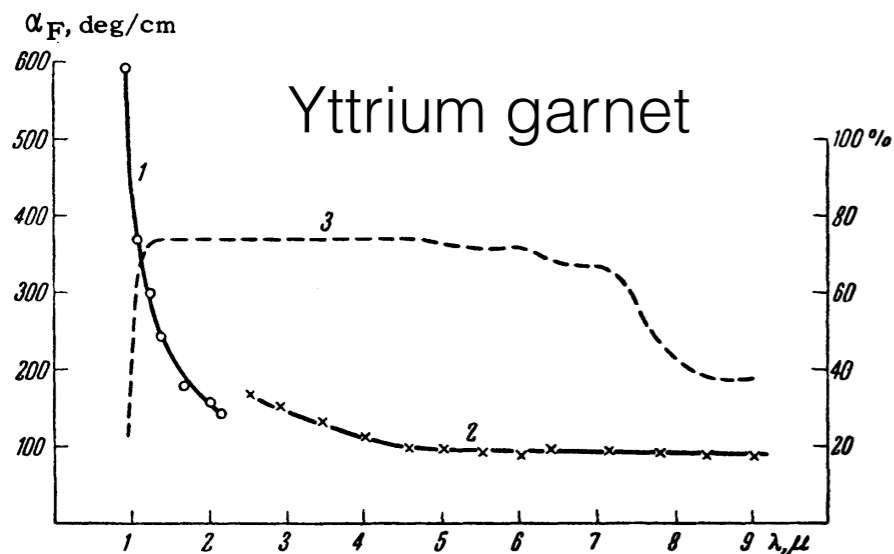
# FR on Different Systems

Here: Far off-resonance FR → different frequency dependence



electric effect:  $\sim \lambda^{-2}$

G. Phelps et al., RSI 86, 073107 (2015)



magnetic effect:  $\lambda^0$

G.S. Krinchik and M.V. Chetkin, JETP 38, 1643 (1960)

Plasma (free electrons)

electric effect:  $\sim \lambda^2$

Faraday rotation: 2<sup>nd</sup> order time-dependent theory

A. D. Buckingham and P.J. Stephens, Ann. Rev. Phys. Chem. 17, 399 (1966)

# Gyromagnetic Faraday Effect

Gyromagnetic FR:

- off-diagonal elements of magnetic susceptibility tensor:  $\chi_M^\pm$
- precession of the  $\vec{M}$  under the influence of the B-field of e.m. wave

$$\theta = \frac{g\mu}{2\hbar c} ML = \frac{g\mu^2}{2\hbar c} NPL$$

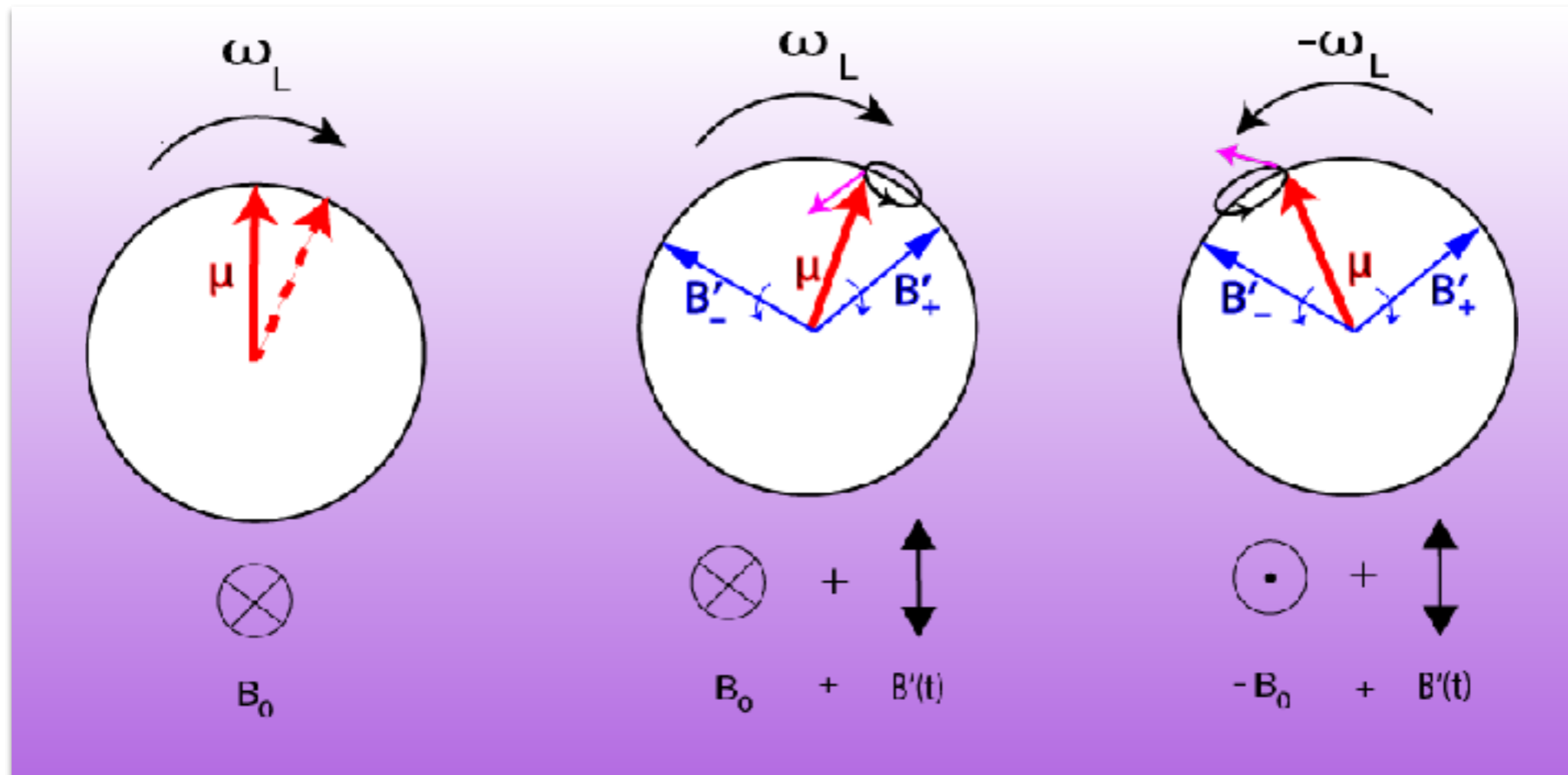
The diagram shows the equation  $\theta = \frac{g\mu}{2\hbar c} ML = \frac{g\mu^2}{2\hbar c} NPL$  inside a yellow rounded rectangle. Arrows point from labels to the variables in the equation: 'magnetic moment' points to  $\mu$  in the first term; 'density' points to  $N$ ; 'polarization' points to  $P$ ; and 'length' points to  $L$ .

G.S. Krinchik & M.V. Chetkin, Sov. Phys. JETP 14, 485 (1961)

S. Gardner, PRD 79, 055007 (2009)

.....

# Test Case: Polarized $^3\text{He}$



- Can produce dense highly polarized targets
- No unpaired electrons  $\rightarrow$  diamagnetic atoms

# Gyromagnetic Faraday Effect on Polarized $^3\text{He}$

Pol.  $^3\text{He}$  as a test case for DM:

Conditions:

- Independent of frequency
- Small rotations:
- $L = 40 \text{ cm}$ ,  $p = < 8 \text{ atm}$ ,  $P < 1$

Expectation: 490 nrad for  $p = 8 \text{ atm}$ ,  $P = 1$

Some order of  
magnitude numbers:

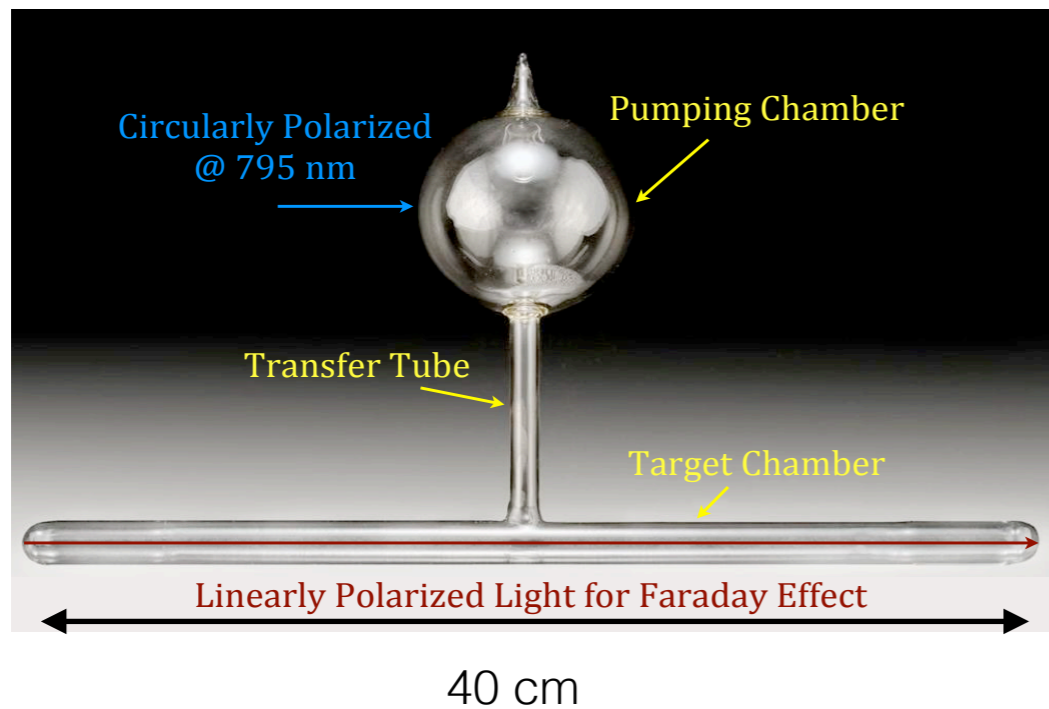
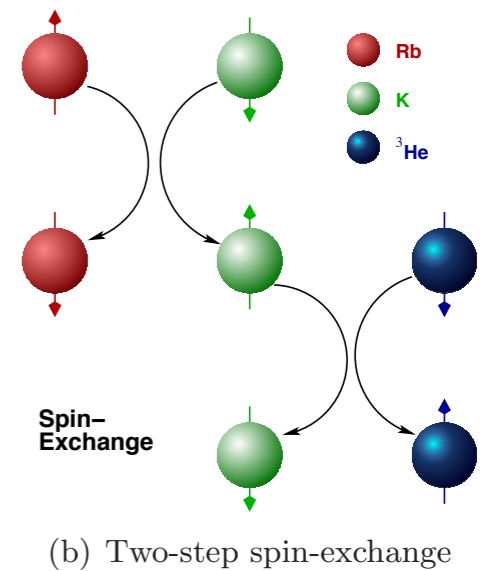
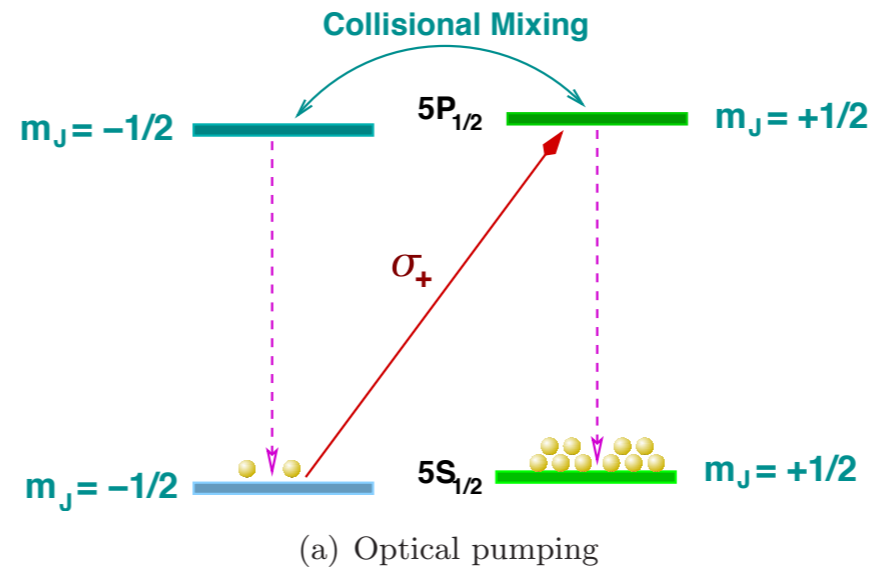
**off-resonance**

Material	Length	$\Theta$ at $B = 4\text{G}$
Flint Glass	1.27 cm	80 $\mu\text{rad}$
Aluminosilicate glass	0.4 cm	6.5 $\mu\text{rad}$
Air	3 m	1.5 $\mu\text{rad}$

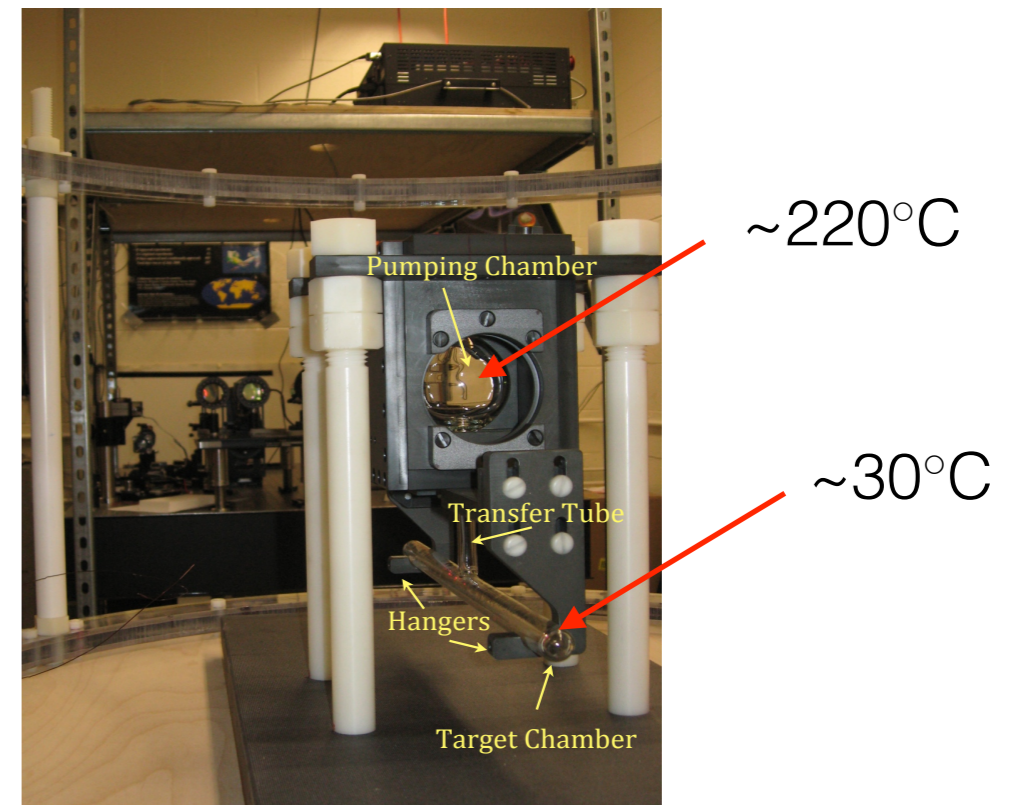
$\lambda = 633 \text{ nm}$

# Producing Densely Polarized $^3\text{He}$ Gas

Polarizing  $^3\text{He}$ :  
spin exchange  
optical pumping  
(SEOP)



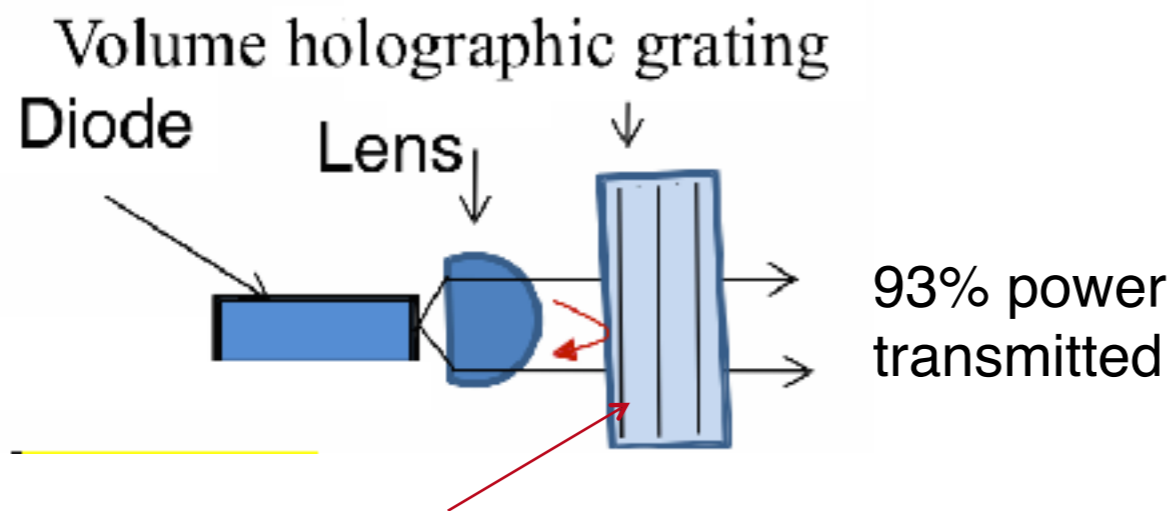
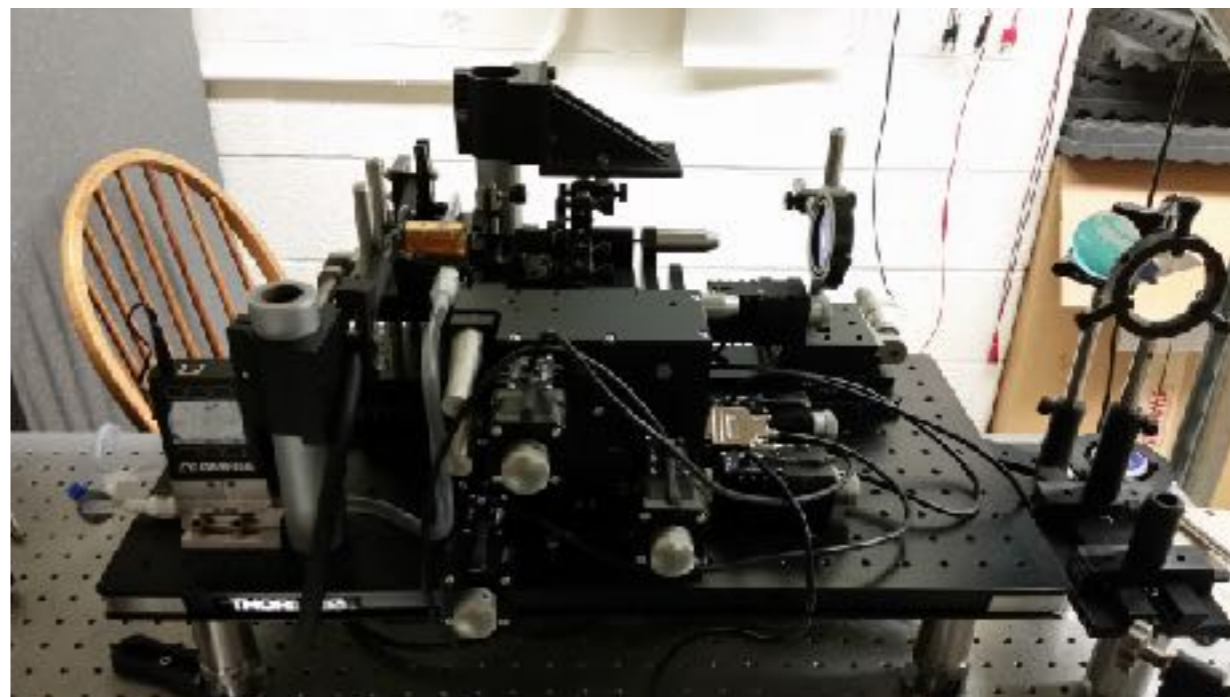
Jefferson Lab geometry



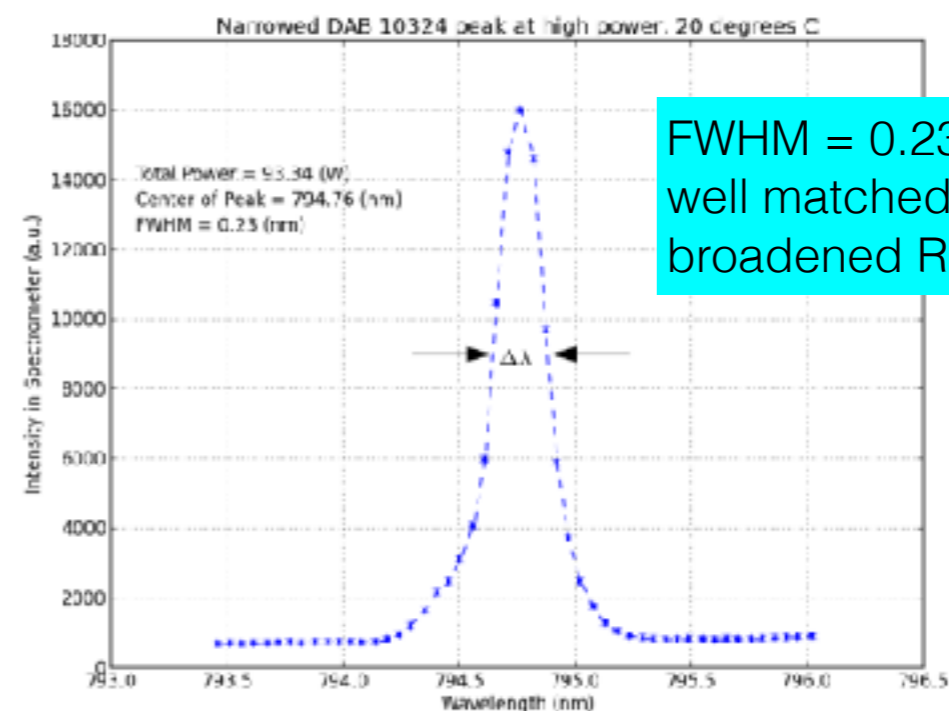


# Polarizing $^3\text{He}$ : SEOP

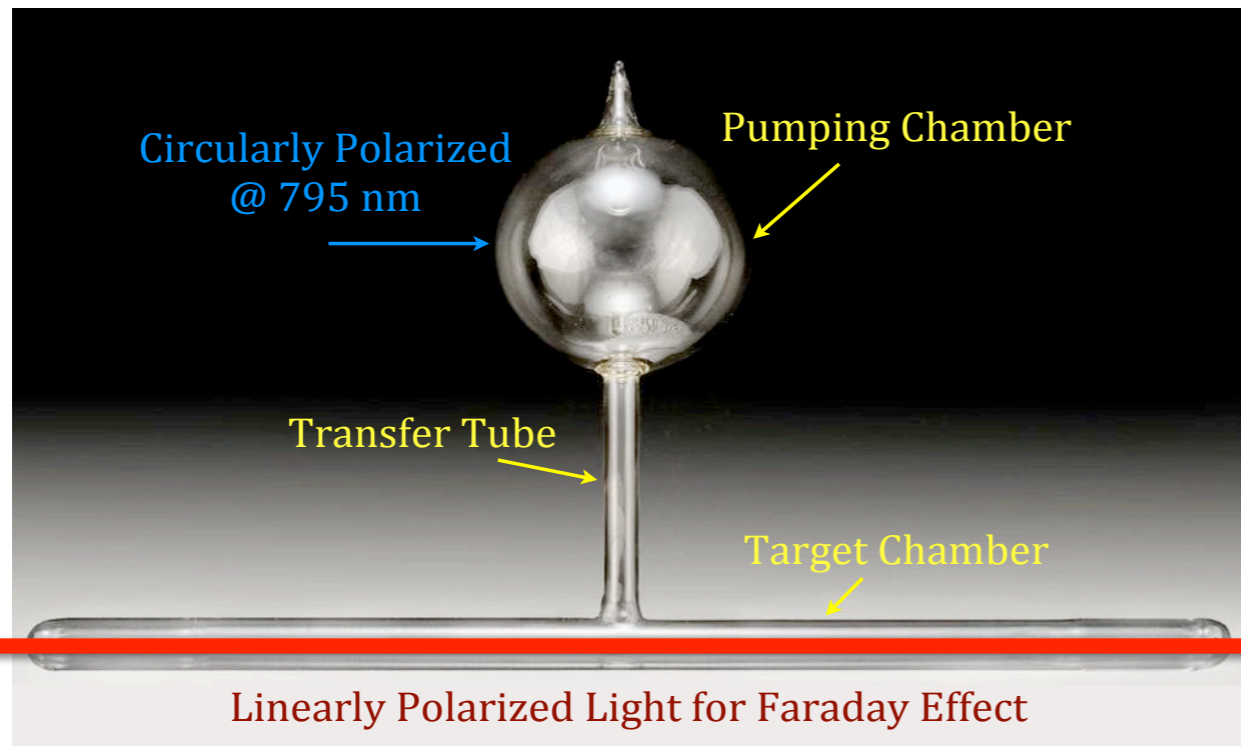
Pumping on D1 line:  
~100 W CW @ 795 nm



Light at 795nm reflected back into diode



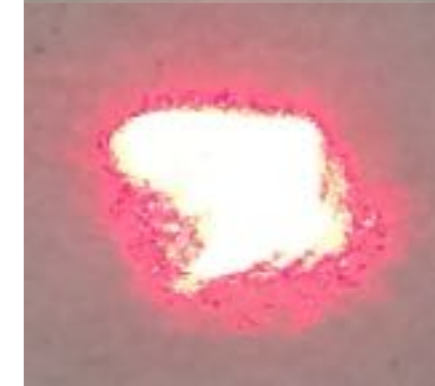
# High Pressure Cells



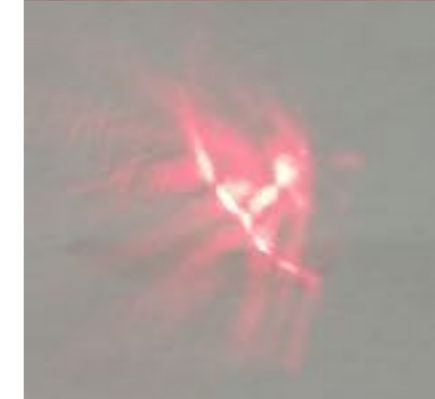
initial probe beam



exiting probe beam (1 pass)

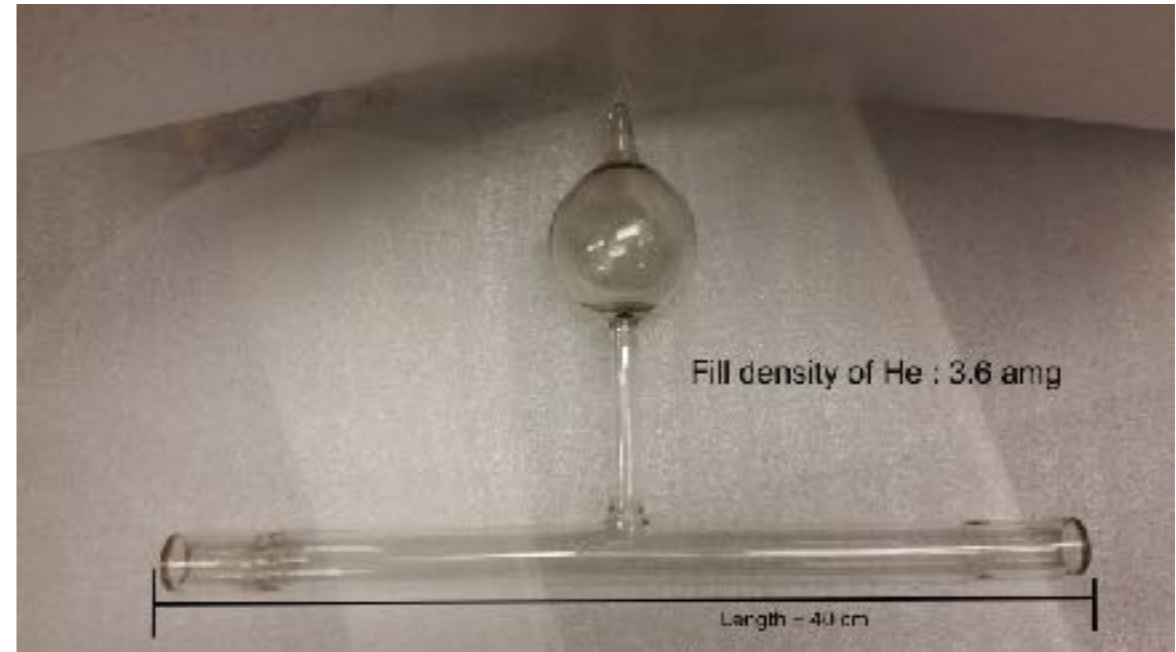


exiting probe beam (3 passes)



# High Pressure Cell for FR Studies

- New SEOP cell: flat, anti-reflective coated end windows designed for multiple passes with laser.
- Target chamber length: 40 cm
- Fill density is 3.6 amg (atm).  
(Lower density due to flat end windows)



rotation angle drops by factor of 2:

→  $\theta \sim 250$  nrad



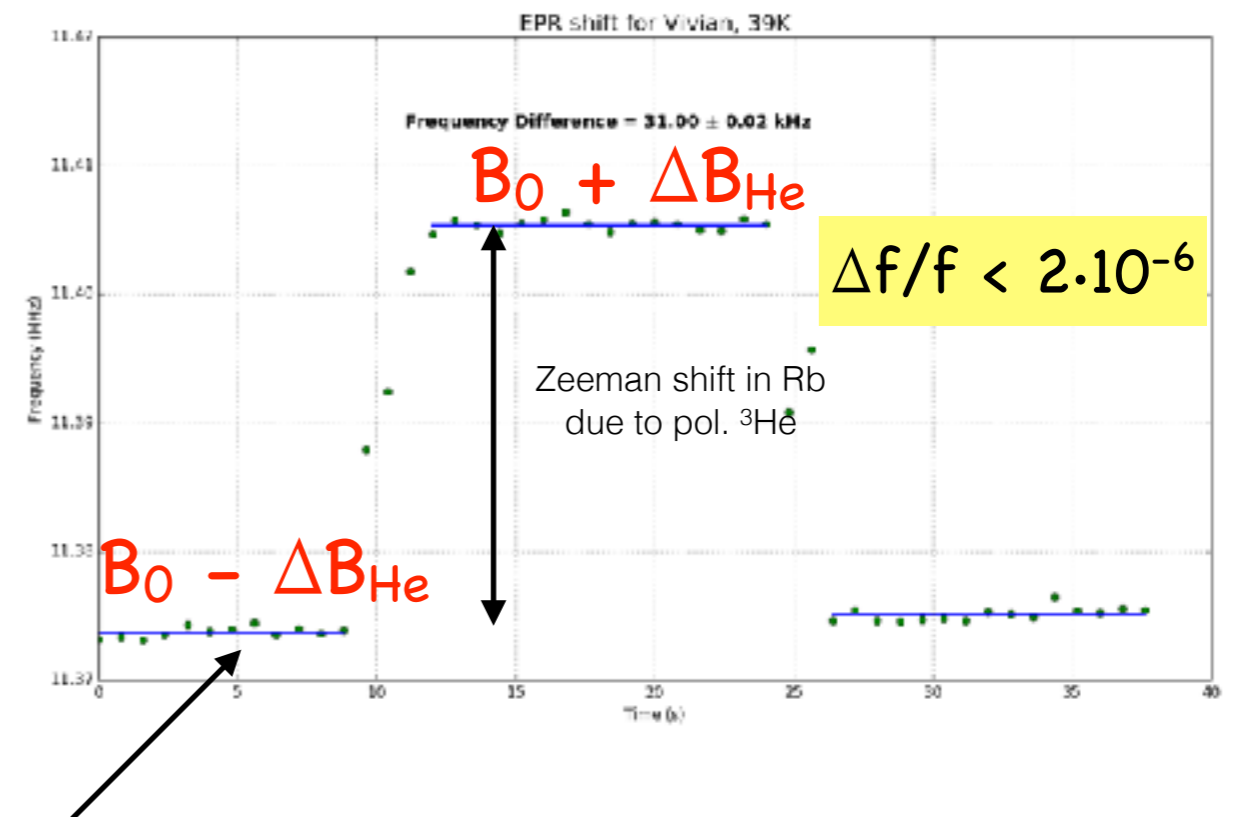
Corning 1720/1723

# Determining the Polarization

## Polarization measurement and monitoring

### Absolute pol. measurement:

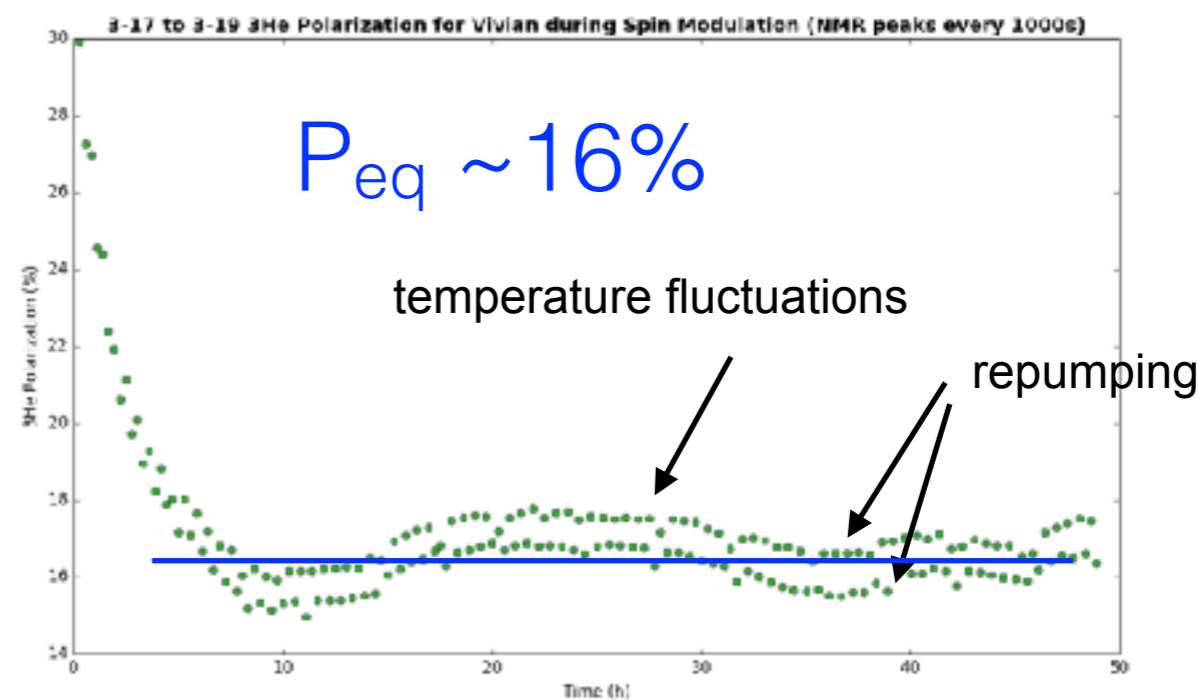
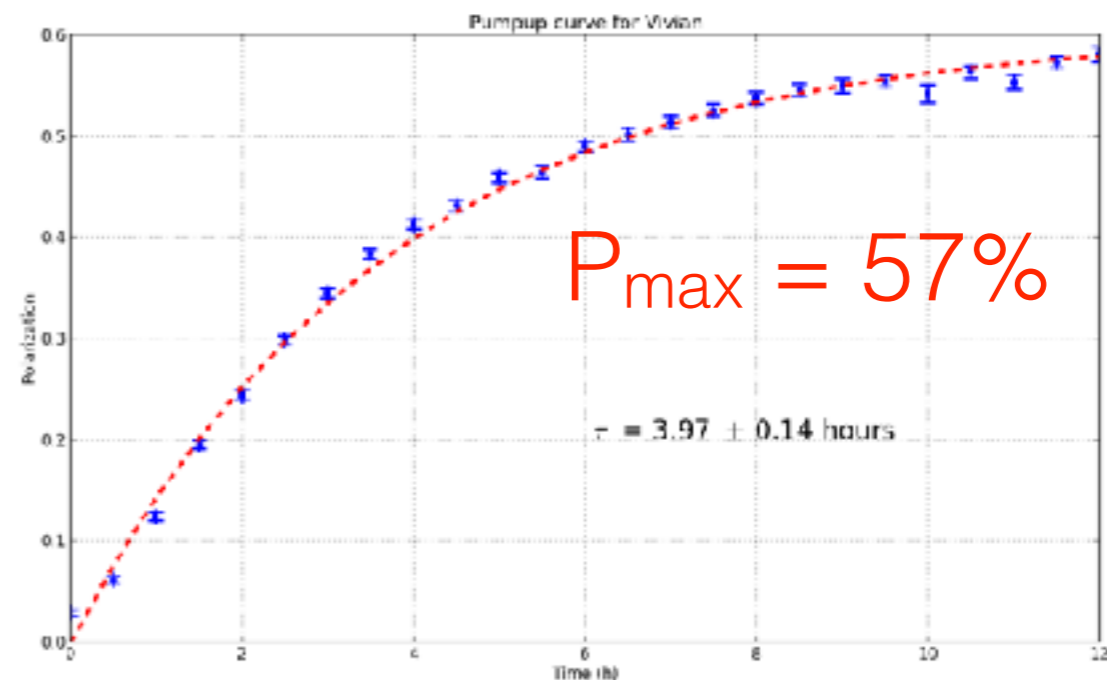
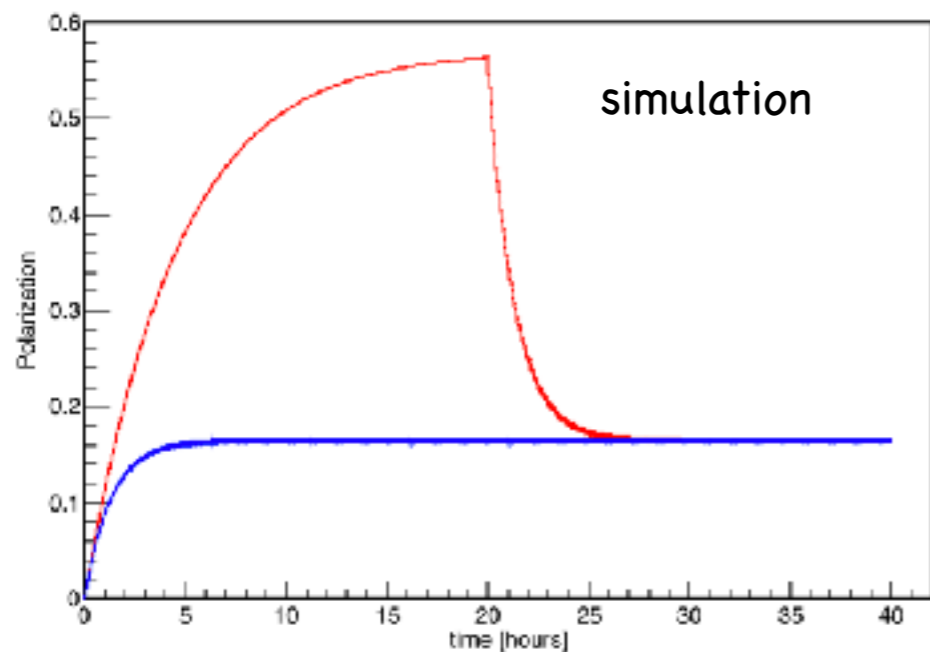
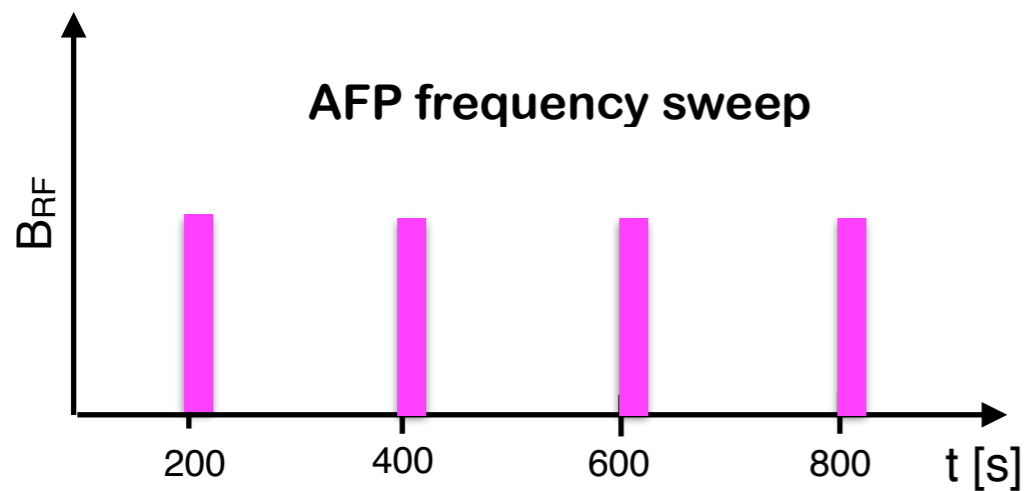
- *Electron Parametric Resonance (EPR) (pumping chamber)*
  - *Zeeman level shift in Rb due to  $^3\text{He}$  “magnetization”*
- *NMR (target chamber)*



active B-field stabilization ( $\Delta B \sim \text{nT}$  level)

# Polarized $^3\text{He}$

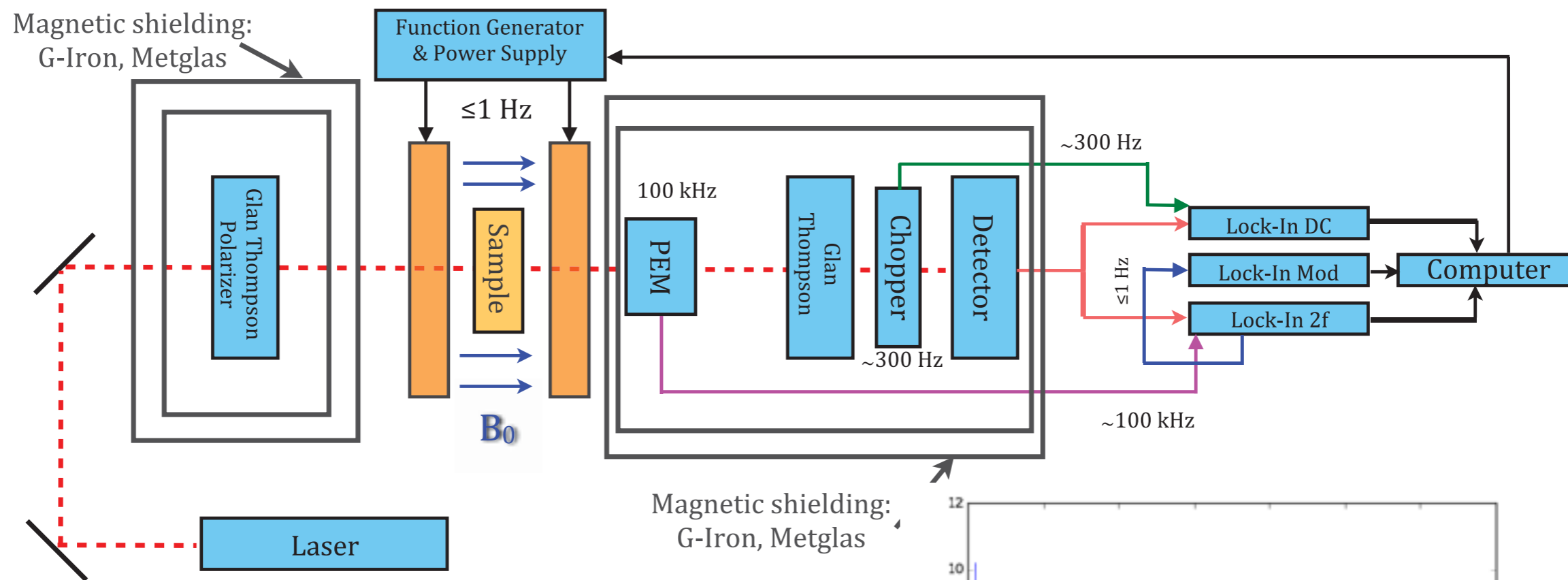
## Target cell polarization studies



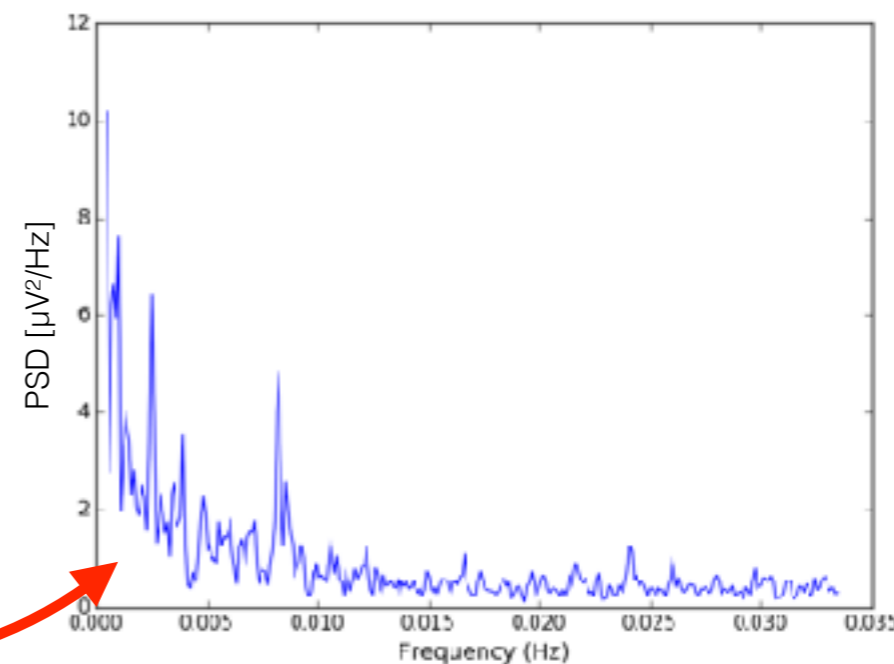
$$P \sim 0.16 \rightarrow \theta \sim 50 \text{ nrad} !!!$$

# Faraday Detection

## Triple modulation technique

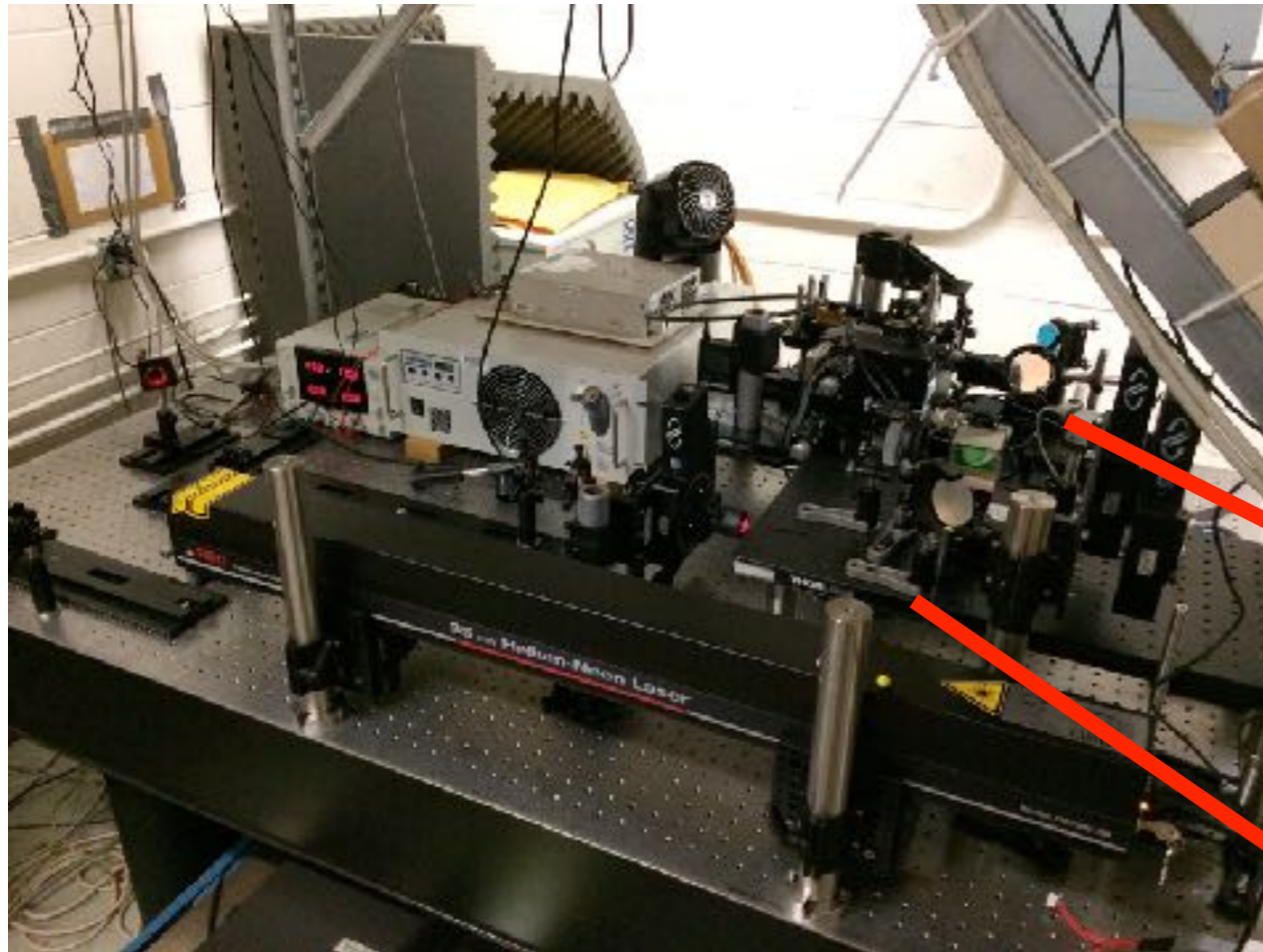


- Three different mod. frequencies:
- polarization mod. of light (100 kHz)
  - intensity mod. of light ( $\sim 100$  Hz)
  - field mod. ( $\sim < 1$  Hz ( $\rightarrow$  few mHz))



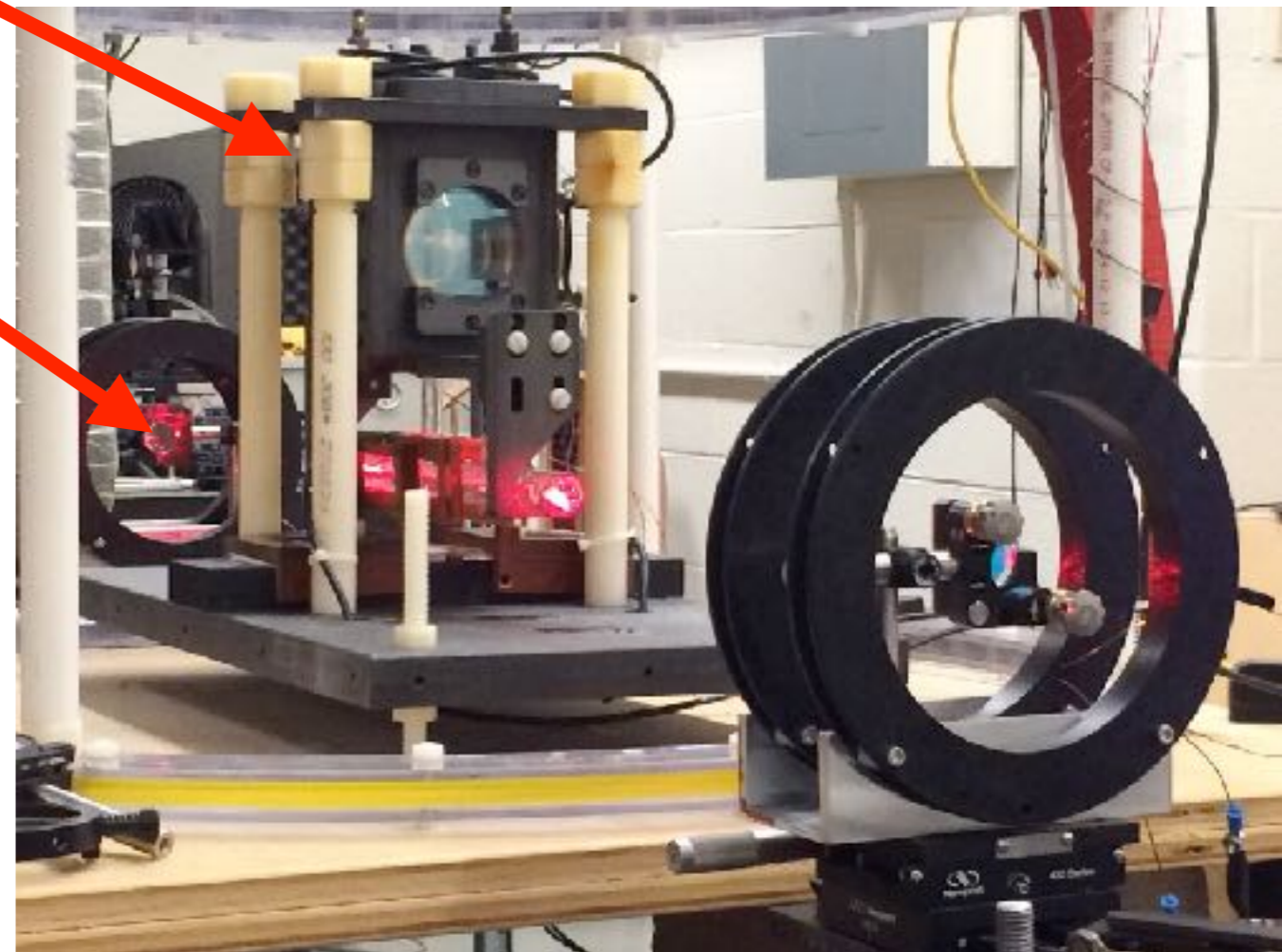
**1/f noise!!**

# Complete Setup



Pump and probe  
laser systems

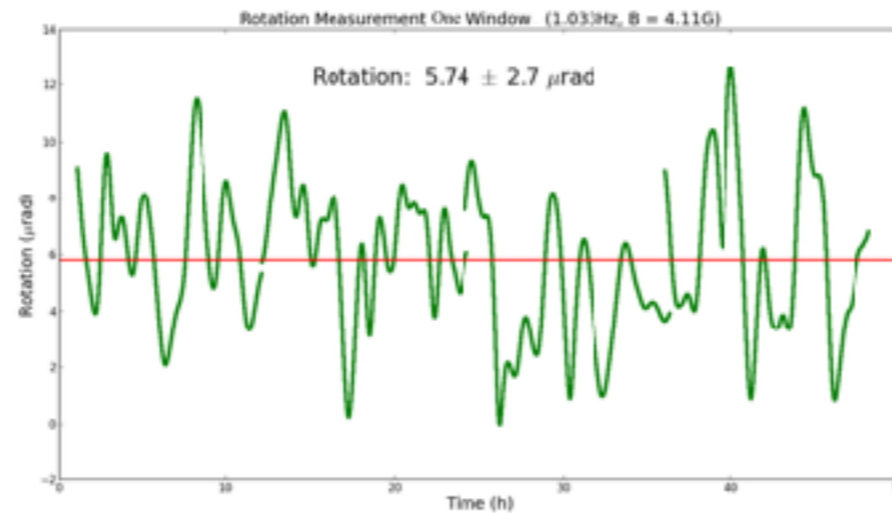
Target region



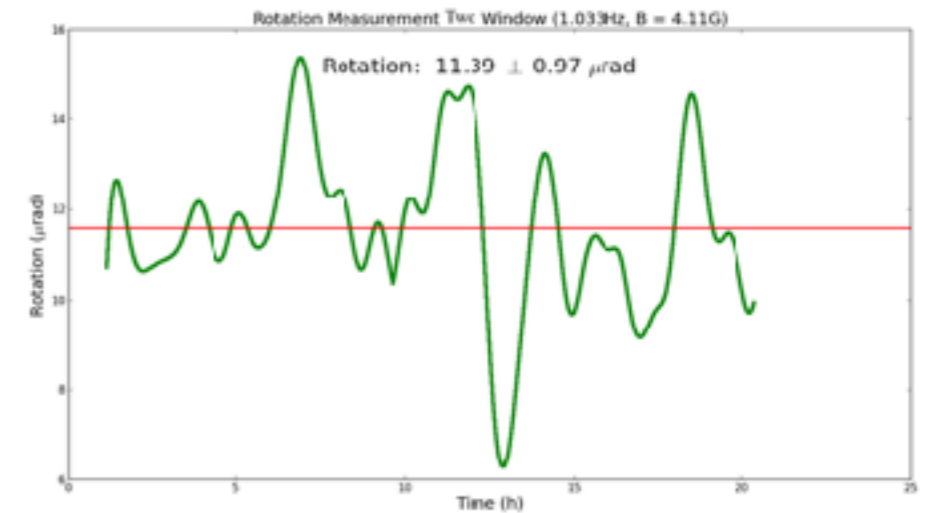
# Some Systematic Checks

- Does the FR from multiple objects in the beam path add up with our measurement method?
- Are there errors associated with windows in the beam path?

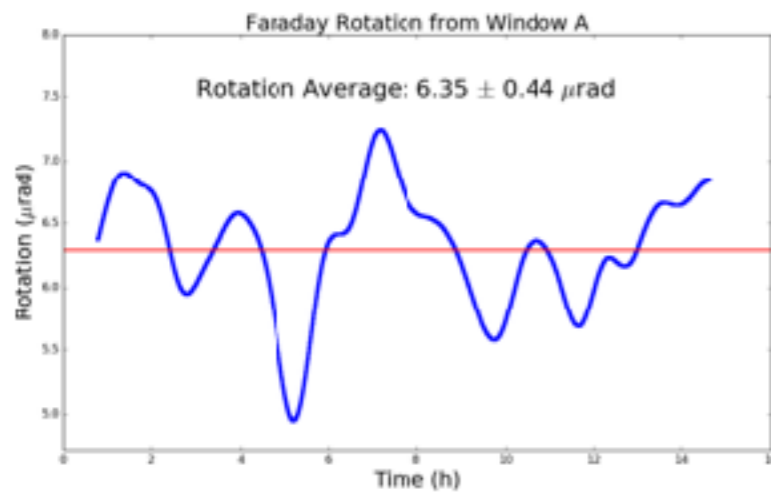
1 window



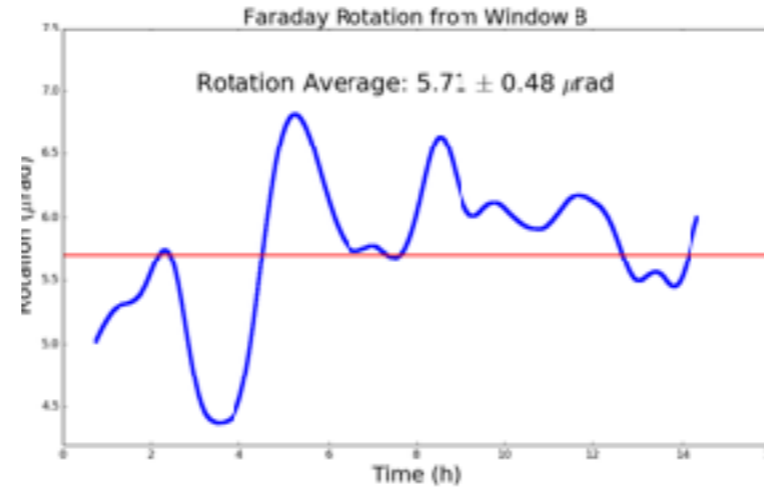
2 windows



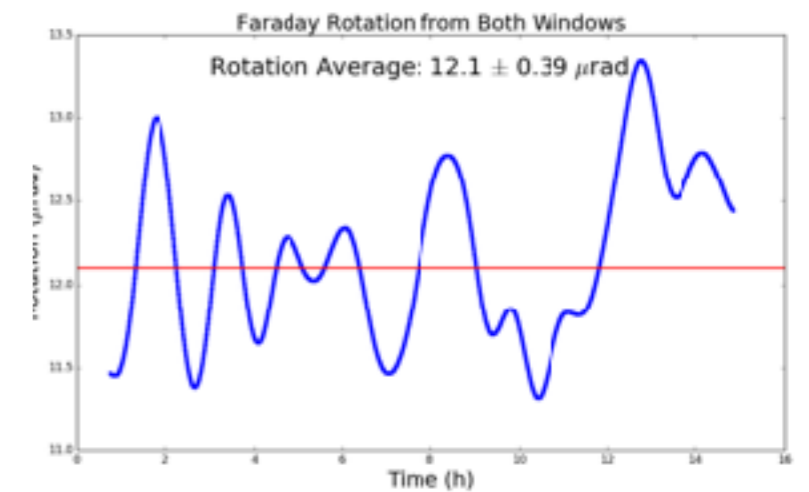
test window A



test window B



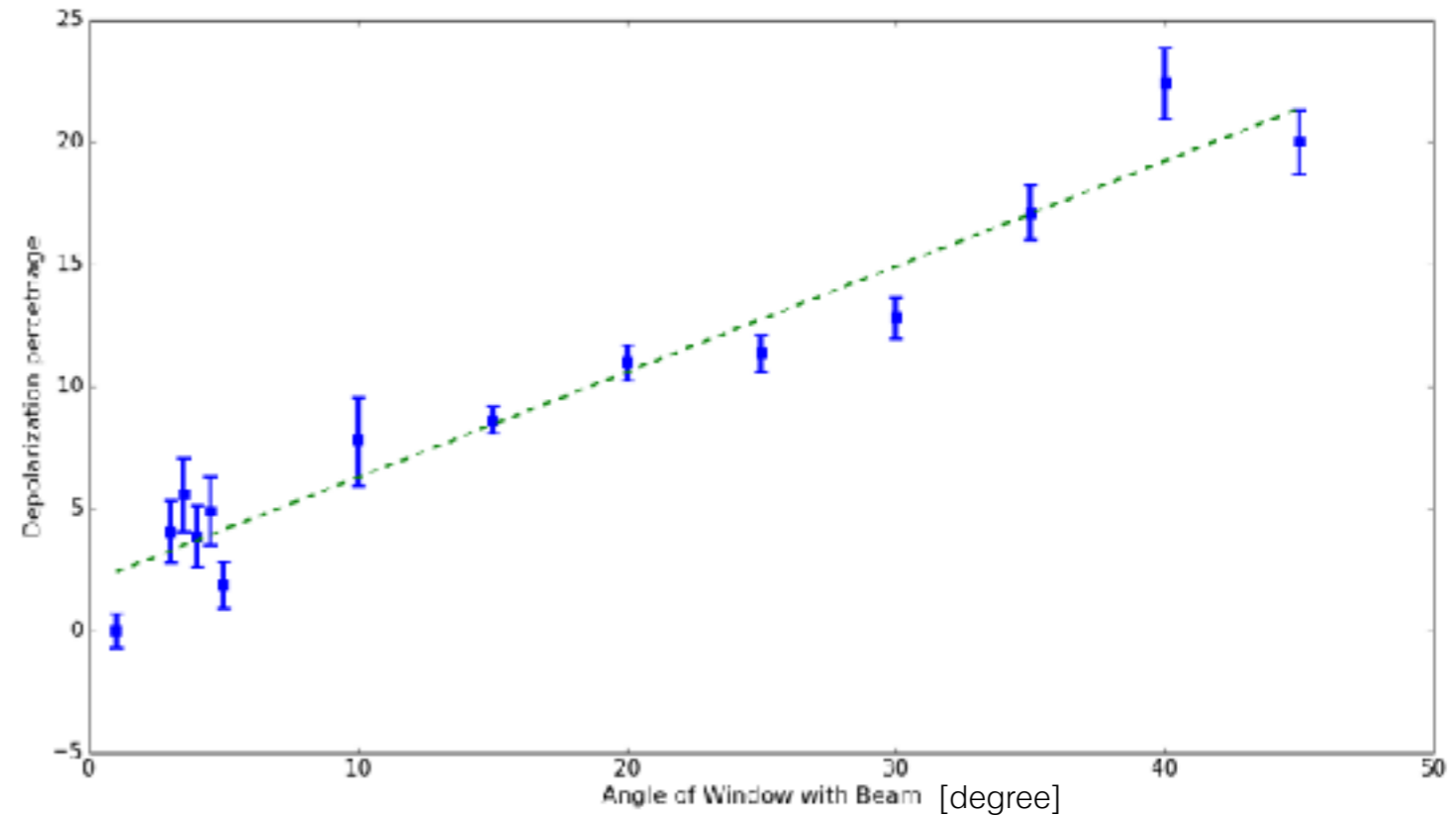
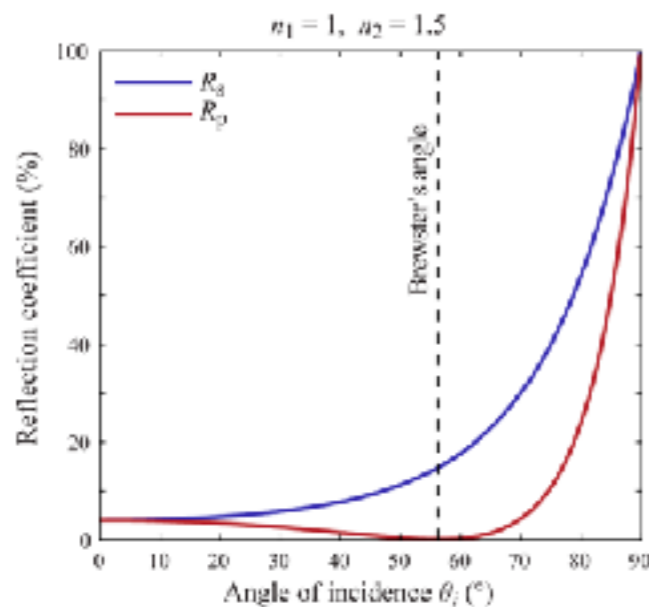
windows A + B





# Light Depolarization

- s- and p-polarizations:
  - depolarization due to reflections
  - effect of magnetic field



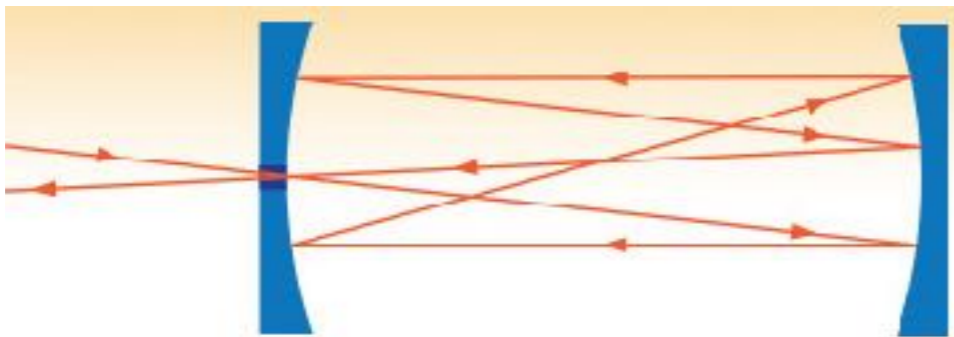
$$\text{Depol \%} = \frac{\text{Measured Value at AOI}}{\text{Measured Value at } 0^\circ \text{ AOI}}$$

Polarization loss: < 3%

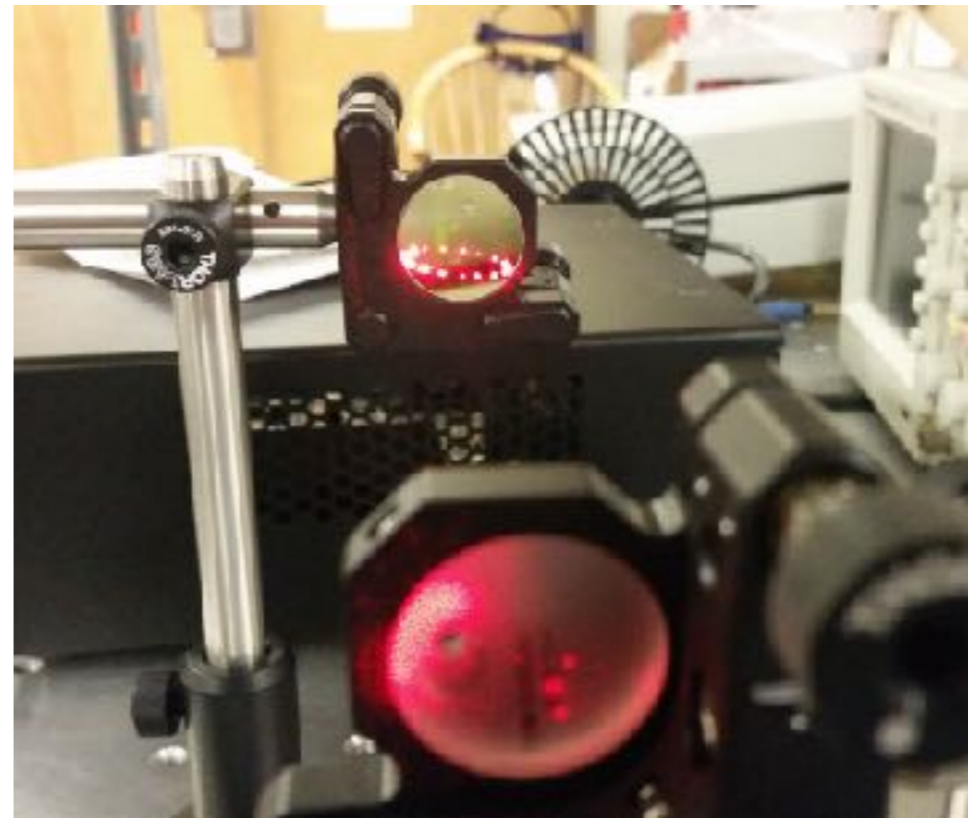
# Increased Sensitivity (Multi-pass Cavity)

Improving SNR:  $\theta_{FR}$  is additive under reflection  $\rightarrow \theta_{FR}$  increases

## Herriott Cavity

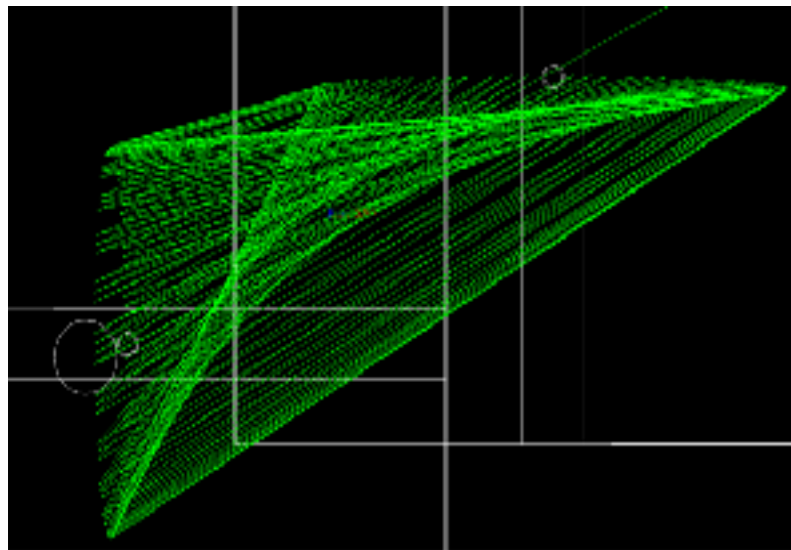
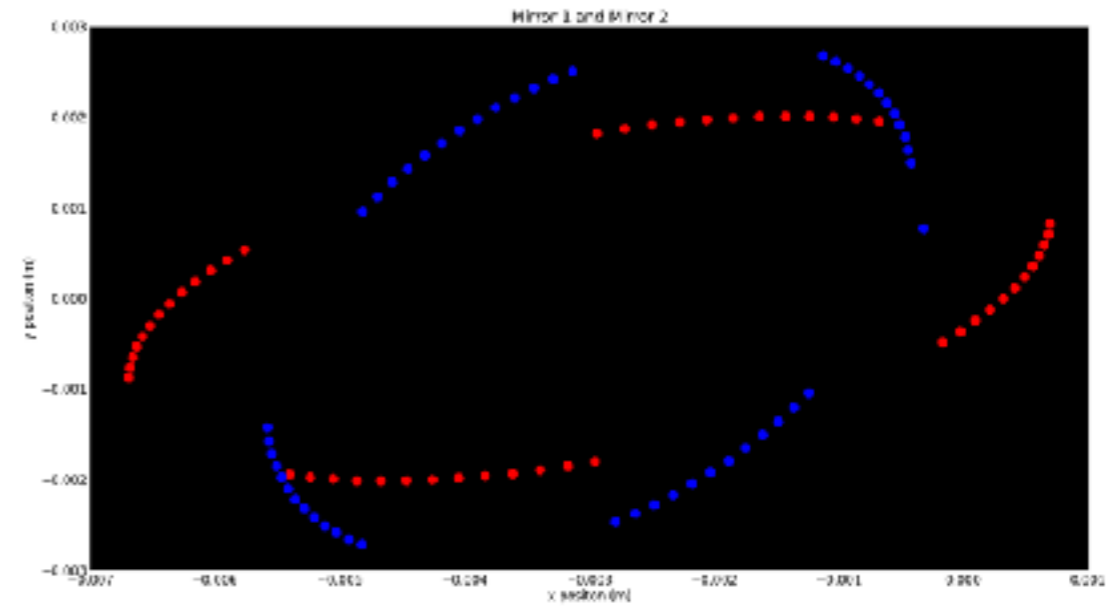
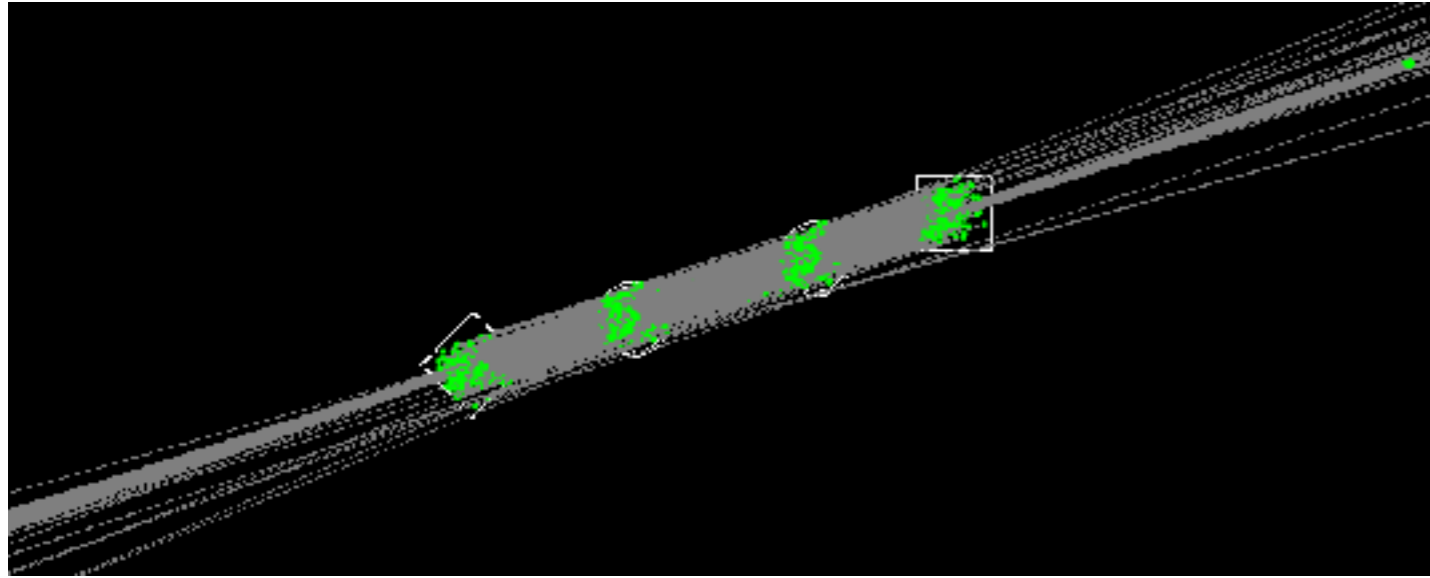


- **Flat mirrors:** not optimal  
need better beam focusing
- **Cylindrical mirrors:** concave ✓
- **Spherical mirrors:** concave ✓



# Optimizing Mirror Geometry

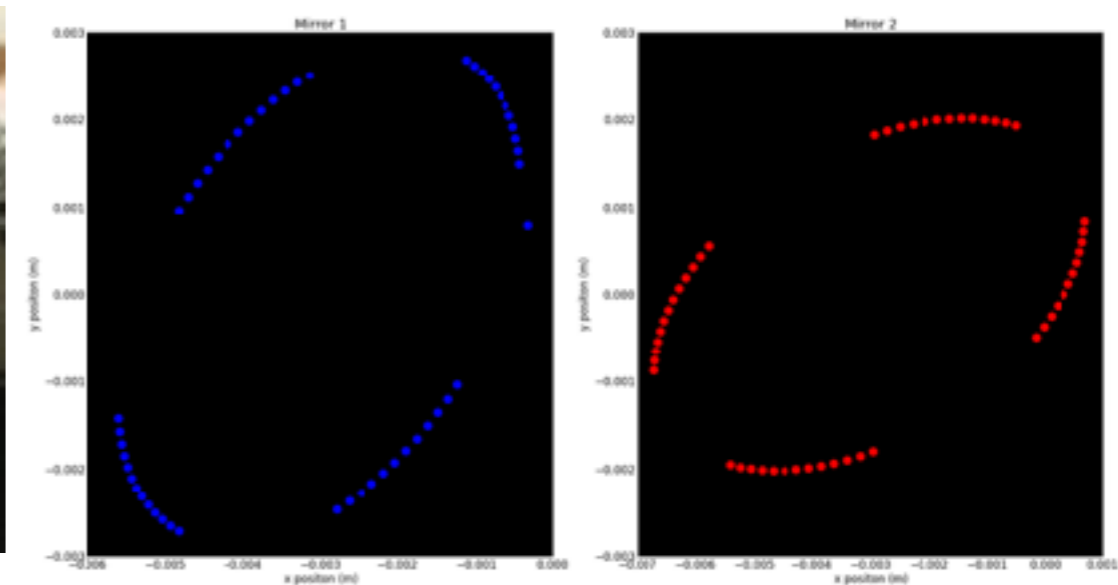
Two spherical mirrors with glass windows



Cylindrical mirrors

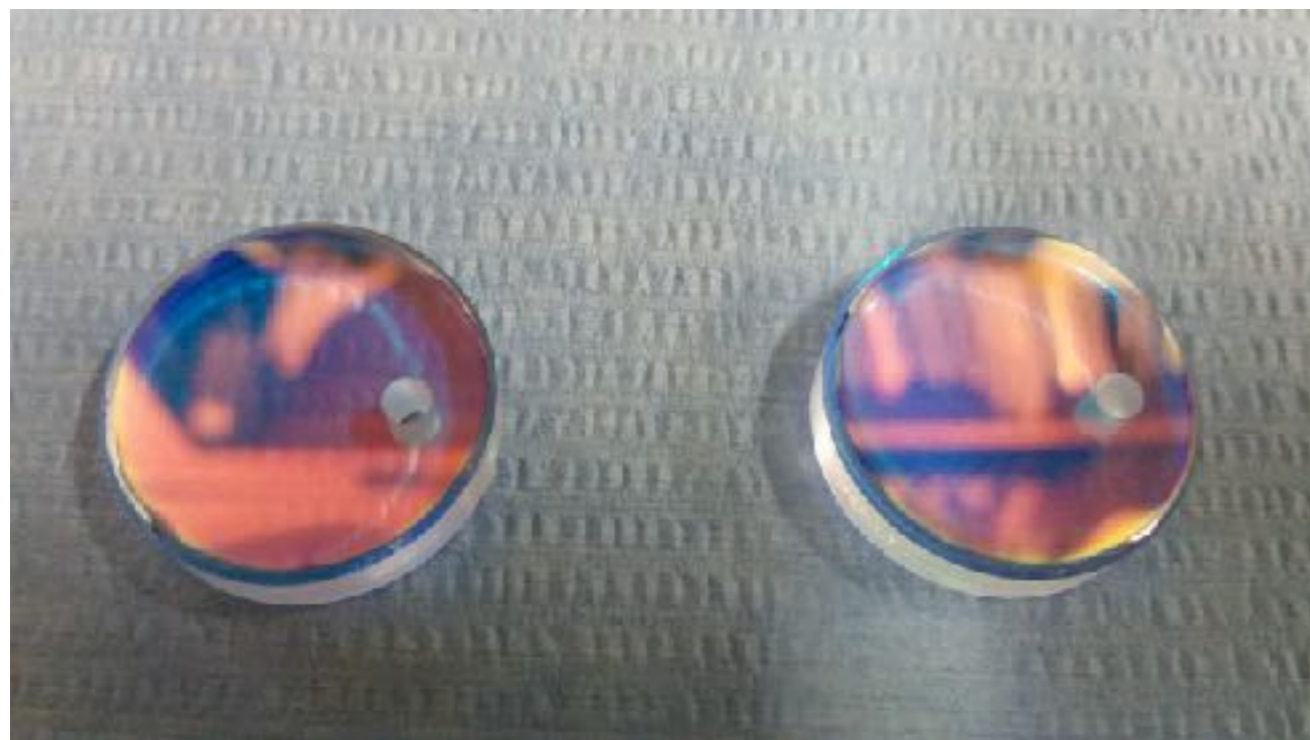


Gold covered spherical mirror

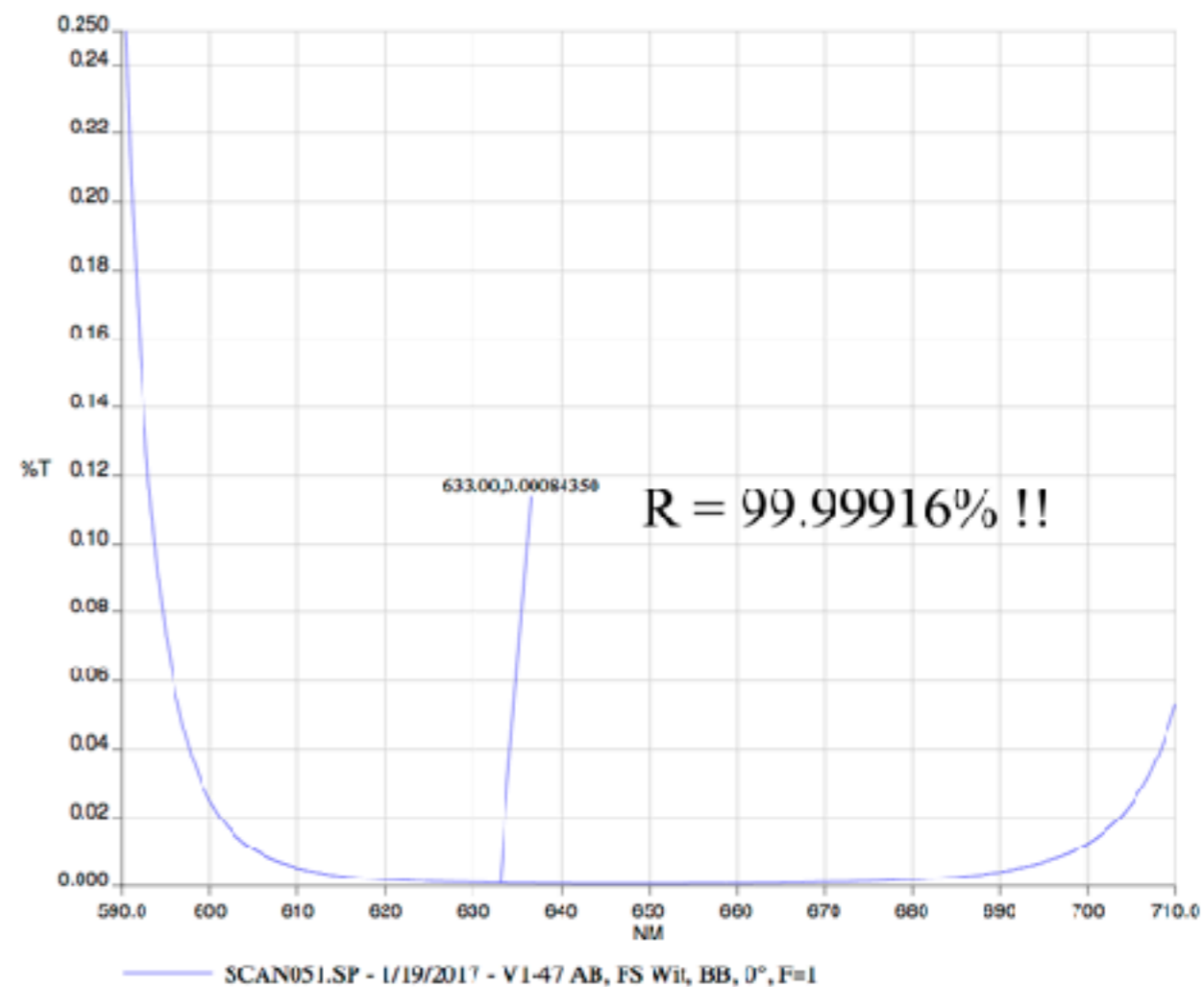


Single photon pattern on mirror surface

# Cavity Mirrors



Date: 1/15/2017 Time: 12:19:12 PM

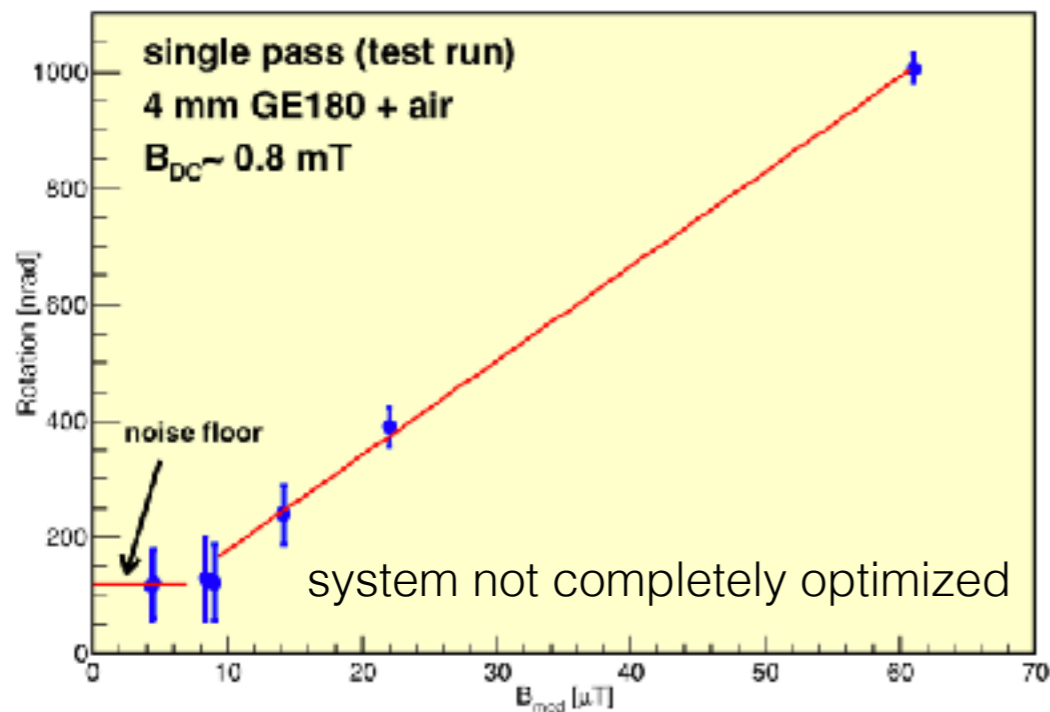


# Experimental Tests

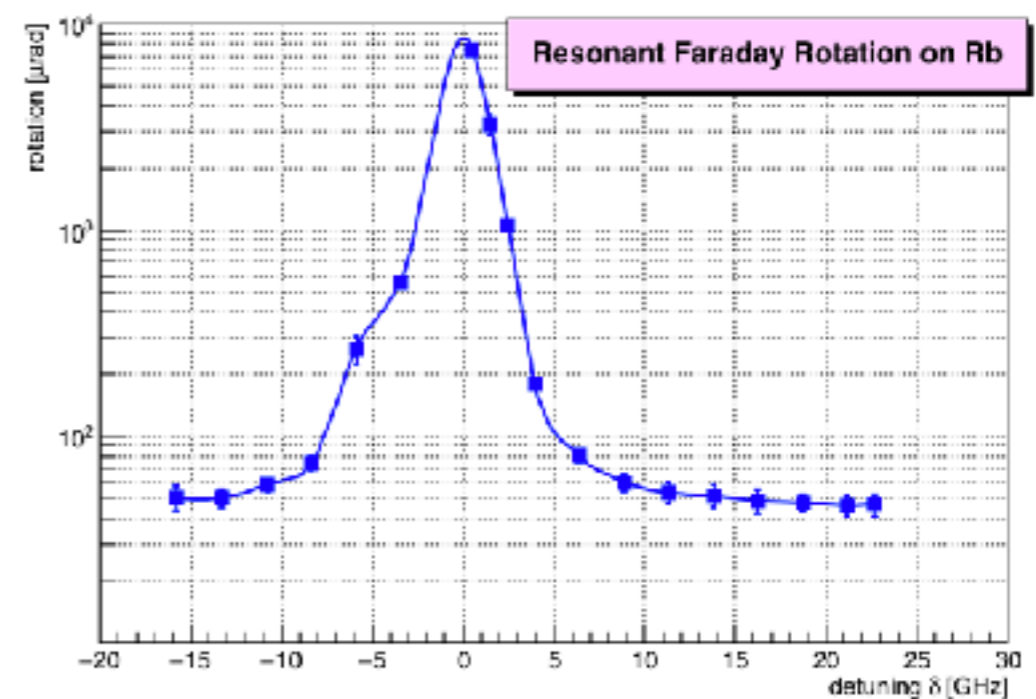
Single pass vs multi pass: 4 mm GE180 + air

# of passes	$\theta_{\text{meas}}$ [ $\mu\text{rad}$ ]	$\theta_{\text{air}}$ [ $\mu\text{rad}$ ]	$\theta_{\text{glass}}$ [ $\mu\text{rad}$ ]
1	8.95	2.15	6.80
15 $\times$ 1	134.25	-	102.00
15	126.05	1.69	100.24

Single pass: B-field scan 4 mm GE180 + air

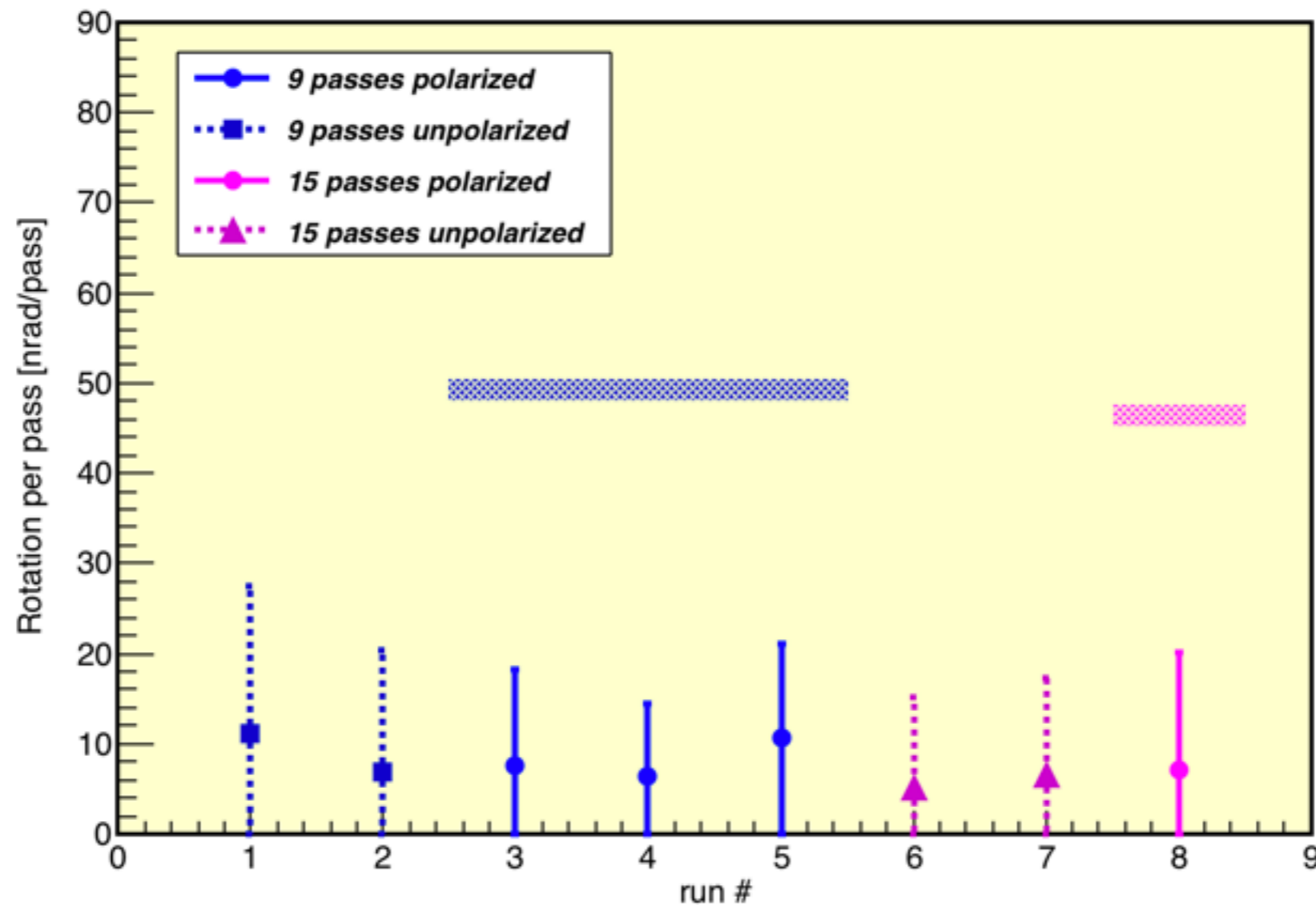


Single pass: On-resonance  
10 cm Rb gas at  $T_{\text{room}}$



# First FR Measurement on Polarized $^3\text{He}$

Preliminary!!



Best measurement so far (w/o  $^3\text{He}$  cell):  $\theta = (15 \pm 9)$  nrad (15 passes)  
 $\rightarrow (1 \pm 2)$  nrad/pass

# Summary

- Gyromagnetic Faraday rotation as a probe of a possible DMAU
- Developed apparatus with a sensitivity of 2 nrad/pass
  - $\Delta n \sim 4 \times 10^{-14}/\text{cm}$  @ 633 nm
- Test case → polarized  $^3\text{He}$ : no effect at level of  $\sim 10$  nrad/40 cm observed yet
- Further systematic studies in progress
- Do results place limits on SME (e.g. limit on  $\vec{\sigma}_{\text{photon}} \cdot \vec{\sigma}_{^3\text{He}}$  correlations ?)