

EFTs & Modified Gravity: the view from below

A liberally conservative point of view

QVG: Testing Gravity in Cosmology
MITP March 2017



CP Burgess

Outline

- EFTs and gravity
- Naturalness issues
 - decoupling and its uses
- Time dependent issues
 - against exceptionalism
- Lessons for tests of gravity
 - some possible surprises

EFTs & Gravity



EFTs & Gravity

- Precision comparison with experiment requires quantification of theoretical error
- $a(\mu\text{on}) = 1159652188.4(4.3) \times 10^{-12}$ (exp)
- $a(\mu\text{on}) = 1159652140(27.1) \times 10^{-12}$ (th)
- QED's renormalizability is important for its calculability, and so underpins theory error

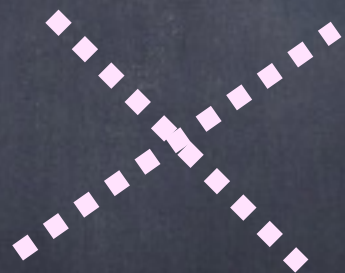
EFTs & Gravity

- GR is also tested with precision
 - $dP/dt = -2.408(10) \times 10^{-12}$ (exp)
 - $dP/dt = -2.40243(5) \times 10^{-12}$ (th)
- Why doesn't nonrenormalizability of GR undermine ability to fix theory error?
 - It would, if we believed nothing could be said at all about quantum corrections in gravity

EFTs & Gravity

- e.g. for graviton scattering on a fixed weakly-curved background:

$$\mathcal{L} = (\partial h)^2 + \frac{1}{M_p} h (\partial h)^2 + \frac{1}{M_p^2} h^2 (\partial h)^2 + \dots$$



$$\mathcal{A}_{\text{classical}} = \frac{Q^2}{M_p^2} + \dots$$

EFTs & Gravity

- e.g. for gravity
weakly-curved

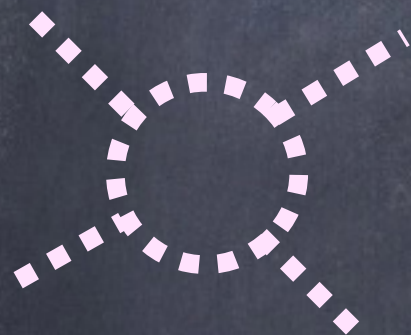
$$\mathcal{L} = (\partial h)^2$$

Need not be expansion about strictly flat space: Q generically denotes size of derivatives, including background curvature

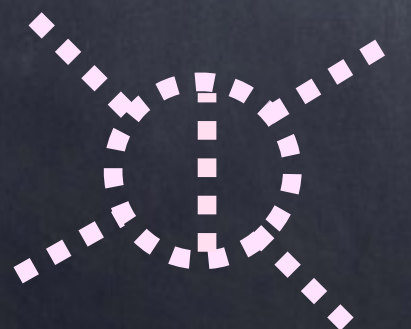
$$A_{\text{classical}} = \frac{Q^2}{M_p^2} + \dots$$

EFTs & Gravity

- Higher order contributions diverge more and more due to dimension of the coupling



$$\mathcal{A}_{1\text{-loop}} = \frac{Q^2}{M_p^4} \int \frac{d^4 p}{(2\pi)^4} \frac{p^6}{(p^2 + Q^2)^4}$$



$$\mathcal{A}_{1\text{-loop}} = \frac{Q^2}{M_p^6} \int \left[\frac{d^4 p}{(2\pi)^4} \right]^2 \frac{p^{10}}{(p^2 + Q^2)^7}$$

EFTs & Gravity

- New divergences cannot be absorbed into G

$$\frac{\mathcal{L}}{\sqrt{-g}} = \Lambda + \frac{M_p^2}{2} R + c_1 R^2 + c_2 R_{\mu\nu} R^{\mu\nu} + \frac{c_3}{m^2} R^3 + \dots$$

- But new divergences **can** be absorbed **if** GR is part of a general derivative expansion involving higher curvatures

EFTs & Gravity

- How to interpret the non-GR terms?

$$\frac{\mathcal{L}}{\sqrt{-g}} = \Lambda + \frac{M_p^2}{2} R + c_1 R^2 + c_2 R_{\mu\nu} R^{\mu\nu} + \frac{c_3}{m^2} R^3 + \dots$$

- As would have arisen after integrating out a collection of particles with masses $m \gg Q$.



EFTs & Gravity

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- As would have arisen after integrating out a collection of particles with masses $m \gg Q$.



- Largest mass (M_p) wins in numerator, but smallest mass (m) wins in denominator

EFTs & Gravity

- As in Wilsonian EFT where effective action (or hamiltonian) is obtained by coarse-graining modes

$$\begin{aligned}\langle \mathcal{O}_1(\ell) \cdots \mathcal{O}_n(\ell) \rangle &= \int \mathcal{D}\ell \mathcal{D}h e^{iS(\ell, h)} \mathcal{O}_1(\ell) \cdots \mathcal{O}_n(\ell) \\ &= \int \mathcal{D}\ell e^{iS_{\text{eff}}(\ell)} \mathcal{O}_1(\ell) \cdots \mathcal{O}_n(\ell)\end{aligned}$$

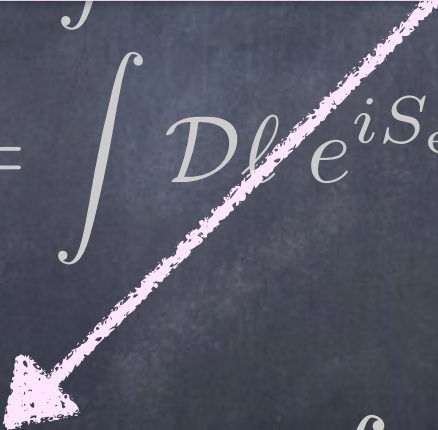
$$e^{iS_{\text{eff}}(\ell)} = \int \mathcal{D}h e^{iS(\ell, h)}$$

EFTs & Gravity

- As in Wilsonian EFT where effective action (or hamiltonian) is obtained by integrating out the high energy modes

S_{eff} local *if* expanded in powers of $1/M$ (due to uncertainty principle)

$$= \int \mathcal{D}\ell \, e^{iS_{\text{eff}}(\ell)} \mathcal{O}_1(\ell) \cdots \mathcal{O}_n(\ell)$$


$$e^{iS_{\text{eff}}(\ell)} = \int \mathcal{D}h \, e^{iS(\ell, h)}$$

EFTs & Gravity

- Predictive despite many terms, provided one recognises one is doing an expansion in Q/m

$$\mathcal{A}_E(Q) \sim \left(\frac{Q^2}{M_p^{E-2}} \right) \left(\frac{Q}{4\pi M_p} \right)^{2L} \prod_{i,d>2} \left(\frac{Q}{M_p} \right)^{2V_{id}} \left(\frac{Q}{m} \right)^{(d-4)V_{id}}$$

- e.g. L -loop amplitude involving E external particles of 'energy' Q , in which V_{id} interactions appear that have i fields and d derivatives

EFTs & Gravity

$$\mathcal{A}_E(Q) \sim \left(\frac{Q^2}{M_p^{E-2}} \right) \left[1 + k \left(\frac{Q}{4\pi M_p} \right)^2 + \dots \right]$$

- Leading contribution:
 - $L=0$ and $V_{id} = 0$ for all $d > 2$
(i.e. Classical GR)

EFTs & Gravity

$$\mathcal{A}_E(Q) \sim$$

Classical GR provides leading low-energy description for almost *any* UV completion!

- Leading contribution:
 - $L=0$ and $V_{ld} = 0$ for all $d > 2$
(i.e. Classical GR)

EFTs & Gravity

$$\mathcal{A}_E(Q) \sim \left(\frac{Q^2}{M_p^{E-2}} \right) \left[1 + k \left(\frac{Q}{4\pi M_p} \right)^2 + \dots \right]$$

- Leading contribution:
 - $L=0$ and $V_{id} = 0$ for all $d > 2$
(i.e. Classical GR)
- Next-to-leading contribution:
 - $L=1$ using only $d=2$ or $L=0$ with $V_{id}=1$ for $d=4$
(i.e. 1-loop GR plus 0-loop with one R^2 interaction)

EFTs & Gravity

$$\mathcal{A}_E(Q) \sim \left(\frac{Q^2}{M_p^{E-2}} \right) \left[1 + k \left(\frac{Q}{4\pi M_p} \right)^2 + \dots \right]$$

- Leading contribution:

These guys renormalise these guys

(i.e. Classical GR)

- Next-to-leading contribution:

- L=1 using only d=2 or L=0 with $V_{\text{id}}=1$ for d=4
(i.e. 1-loop GR plus 0-loop with one R^2 interaction)

EFTs & Gravity

$$\mathcal{A}_E(Q) \sim \left(\frac{Q^2}{4\pi E_{\text{Pl}}^2} \right) \left[1 + k \left(\frac{Q}{4\pi M} \right)^2 + \dots \right]$$

Predictive because only a finite number of unknown coefficients enter at any given order of Q/m

- Leading contribution
 - $L=0$ and $V_{\text{id}}=0$ (i.e. Classical GR)
- Next-to-leading contribution.
 - $L=1$ using only $d=2$ or $L=0$ with $V_{\text{id}}=1$ for $d=4$ (i.e. 1-loop GR plus 0-loop with one R^2 interaction)

EFTs & Gravity

$$\mathcal{A}_E(Q) \sim \left(\frac{Q^2}{M_p^{E-2}} \right) \left[1 + k \left(\frac{Q}{4\pi M_p} \right)^2 + \dots \right]$$

- Leading contribution:

- $L=0$ and $V_{id}=0$
(i.e. Classical GR)

- Next-to-leading

Notice that Q/M_p is loop-counting parameter as well as controlling the derivative expansion

- $L=1$ using only $d=2$ or $L=0$ with $V_{id}=1$ for $d=4$
(i.e. 1-loop GR plus 0-loop with one R^2 interaction)

EFTs & Gravity

- Lessons for testing GR
 - Known to be consistent: GR+light low-spin fields (scalars, vectors); in derivative expansion; possibly higher D ; subject to naturalness constraints.
 - Long-distance implications of many UV theories are captured by limited number of low-dimension interactions

EFTs & Gravity

- Lessons for proposed mods to GR
 - Exotic UV effects?: what is the local effective description at low-energies?
 - Deviations from derivative expansion, e.g. $P(X)$ theories, should check validity of classical approximation (what is m in Q/m ?)
 - Should avoid effects with non-generic & non-negative powers of m (dangerous e.g. for preferred-frame theories)

EFTs & Gravity

- Lessons for proposed mods to GR

- Exotic UV
descriptions

e.g.

$$L = c_1 (\dot{\phi})^2 + c_2 (\nabla \phi)^2$$

- Deviations

theories, should check validity of classical approximation (what is m in Q/m ?)

- Should avoid effects with non-generic & non-negative powers of m (dangerous e.g. for preferred-frame theories)

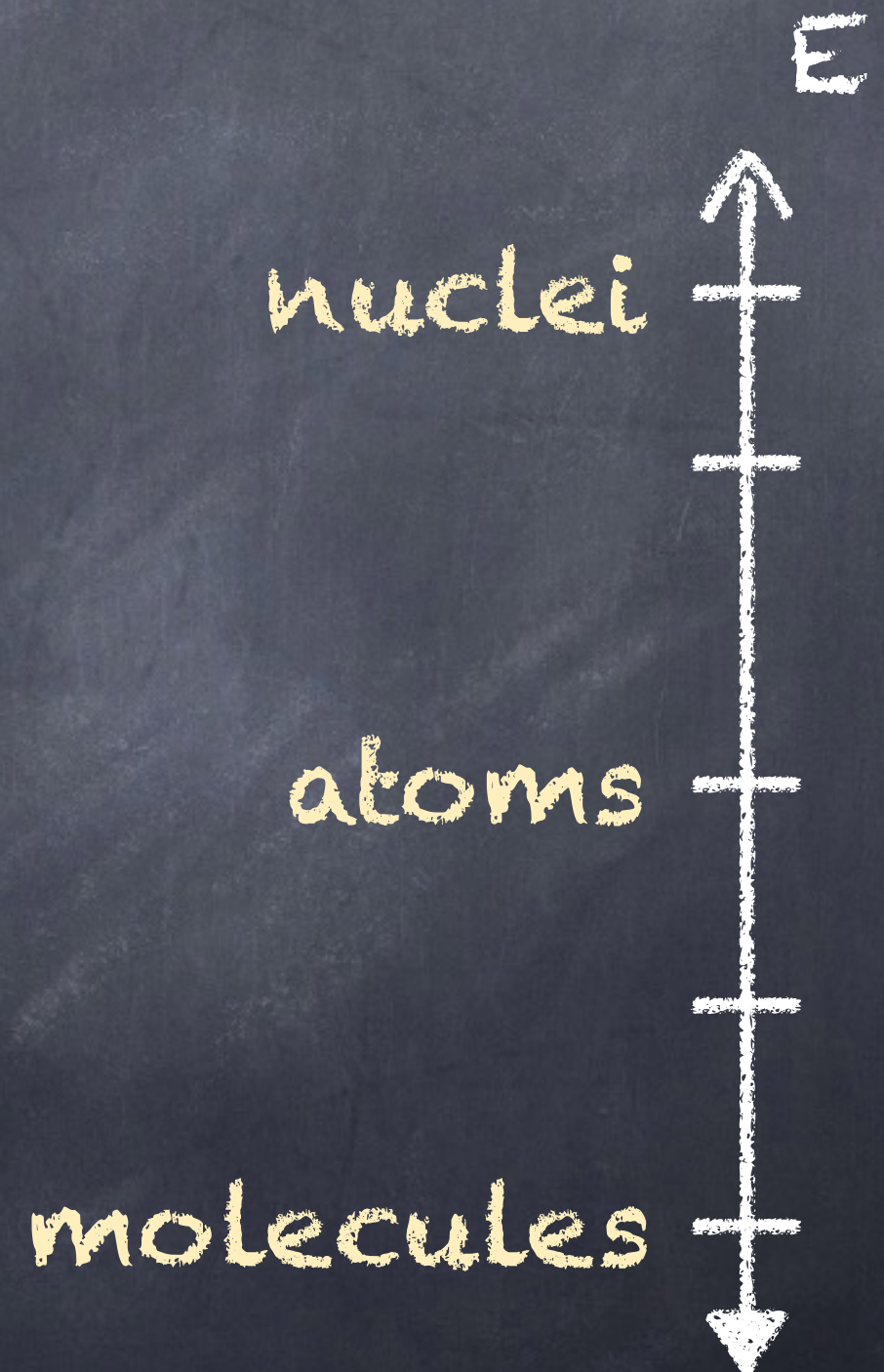
Naturalness



Patron Saint of Naturalness

Naturalness

- Nature comes to us with many scales, and each seems to be understandable on its own terms
- Each is described by an effective theory, obtained by coarse-graining shorter-distance physics



Naturalness

- Nature comes to us with many scales, and each seems to be understandable on its own terms
- Contribution to dimension- D effective interaction $L = c O_D$ after integrating out scale m_i is $c \sim m_i^{4-D}$
- Naturalness: should worry if we find small c when $D < 4$



Naturalness

- Two such interactions in standard theory: one natural one seems not

$$\mathcal{L} = \sqrt{-g} \left[\underbrace{M_p^2 R + \mu^2 H^* H}_{\downarrow} - \lambda (H^* H)^2 - \frac{1}{4} F^2 + \dots \right]$$

$$8\pi G_N = M_p^{-2} \simeq (10^{18} \text{ GeV})^{-2}$$

$$m_H^2 = 2\mu^2 + (\text{loops}) \simeq (125 \text{ GeV})^2$$



Naturalness

- Parameters are specific to a particular effective theory, e.g. for Higgs mass:

$$m_H^2 = 2\mu_1^2 + cM^2 + (\text{loops})$$

$$m_H^2 = 2\mu_0^2 + (\text{loops})$$



Naturalness

- Must cancel to many many decimal places the larger M is

$$m_H^2 = 2\mu_1^2 + cM^2 + (\text{loops})$$


$$m_H^2 = 2\mu_0^2 + (\text{loops})$$

M_P

M

E_W

E

Naturalness

- Technical naturalness:
 - Why is a parameter small in the 'fundamental' theory at very high energies?
 - Why does it remain small when coarse-graining scales down to where it is measured?
- If both answered then 'technically natural'
 - Enhanced symmetry when parameter vanishes provides a simple way to ensure tech. natural
 - Understood hierarchies seem natural in this way

Naturalness

- Useful criterion because suggests kinds of new physics that should not be too distant in energy

- Composite Higgs

(no H field, so no μ , at high E)

Binding energy
EW



Naturalness

- Useful criterion because suggests kinds of new physics that should not be too distant in energy
- Composite Higgs
- Supersymmetric partners
(bose-fermi partners partly cancel)



Naturalness

- Useful criterion because suggests kinds of new physics that should not be too distant in energy
- Composite Higgs
- Supersymmetric partners
- Extra dimensions
(denry quantum gravity enters at M_p)



Naturalness

- Lessons for proposed mods to GR
 - If phenomenology requires small low-dim interactions (eg light scalars in cosmology) should ask why they can be light)

Naturalness

- Lessons for proposed mods to GR
 - If phenomenology requires small low-dim interactions (eg light scalars in cosmology) should ask why they can be light)
 - If symmetry is broken at high energies (eg Lorentz invariance) should ask why it should appear unbroken at low energies

Naturalness

- Lessons for proposed mods to GR

- If phenomenological interaction should appear, why don't $(d\phi/dt)^2$ and $(d\phi/dx)^2$ have coeffs that differ with size $\ln(M/m)$?

- If symmetry is broken at high energies (eg Lorentz invariance) should ask why it should appear unbroken at low energies

CC Problem



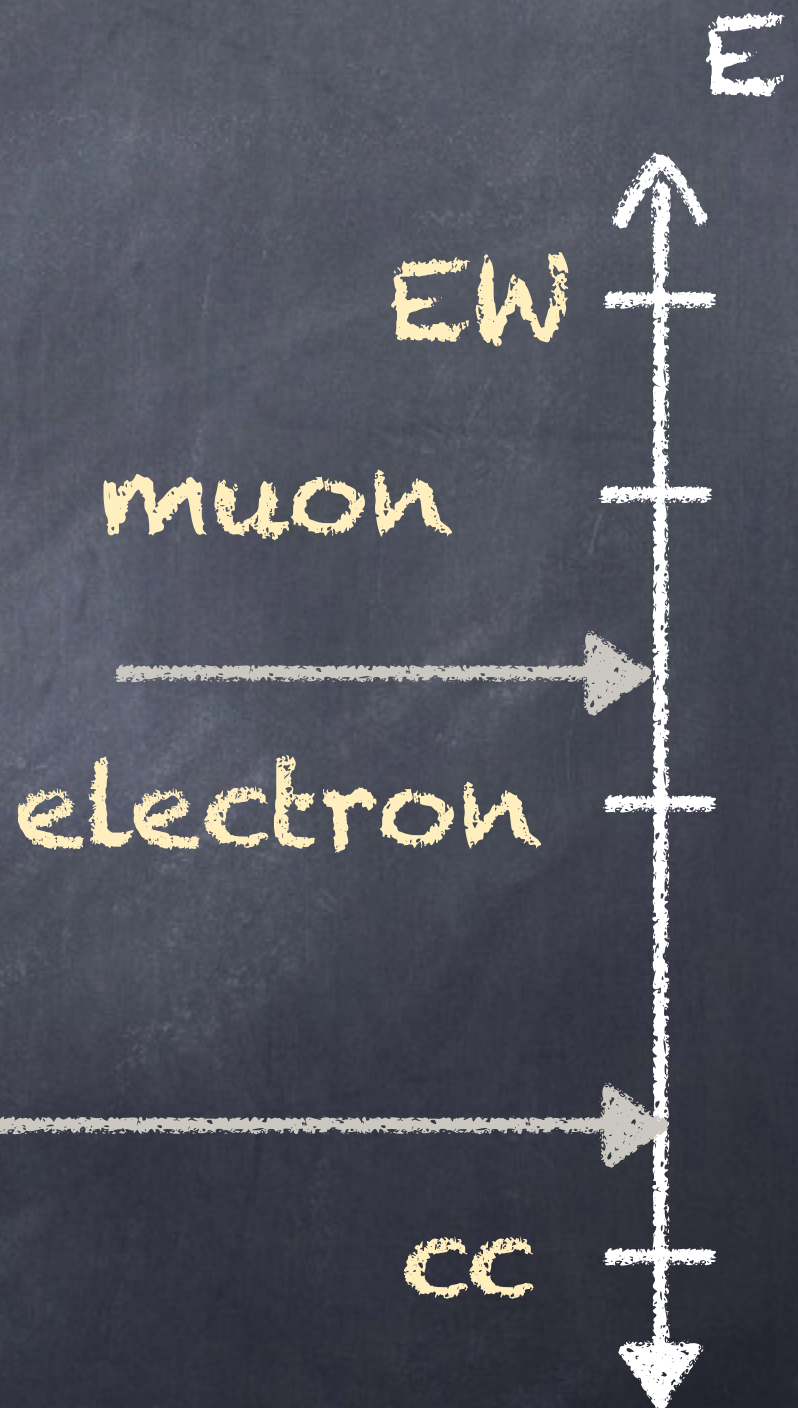
CC Problem

- (old) CC problem: Vacuum energy is also unnatural

$$\rho = \Lambda_1 + cM^4 + (\text{loops})$$



$$\rho = \Lambda_0 + (\text{loops})$$

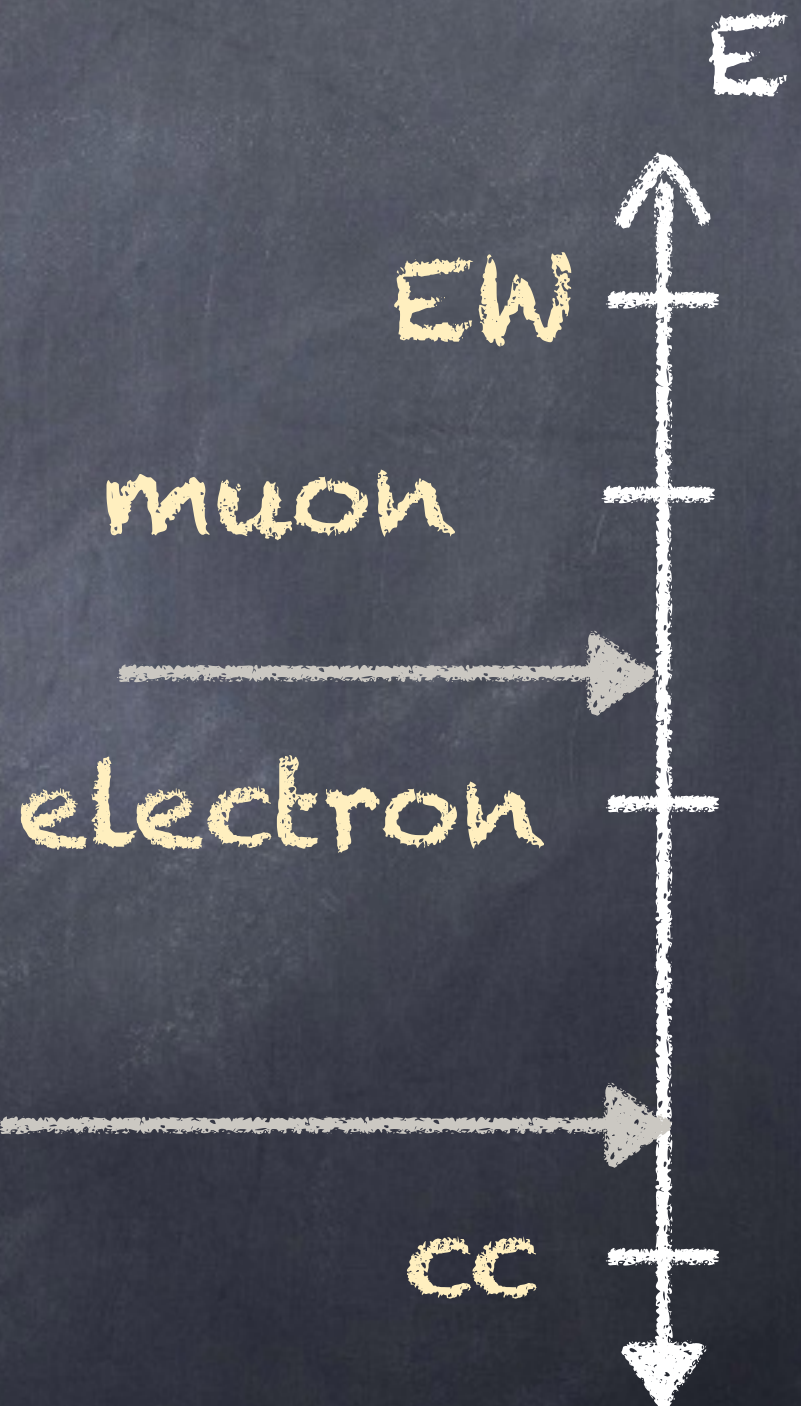


CC Problem

- Now the cancellation occurs at scales we think we understand

$$\rho = \Lambda_1 + cM^4 + (\text{loops})$$

$$\rho = \Lambda_0 + (\text{loops})$$



CC Problem

- Not a problem if we can modify how quantum fluctuations gravitate in vacuum (but NOT also in atoms)
- Any reasonable solution must:
 - go beyond classical approx
 - extend to energies higher than the cc itself
 - do no harm

CC Problem

- No proposals do all three
- Odd situation: no agreed viable proposals yet no no-go result.
- Most common point of view: naturalness arguments may sometimes be wrong or misleading; but when?
 - eg: anthropic proposals
- This is not evidence for failure of EFT itself!

CC Problem

- Some serious contenders exist: e.g. galileons and graviton mass
- Hope to find screening mechanism for cc

$$(\square - m^2)h_{\mu\nu} = \kappa^2 T_{\mu\nu}$$

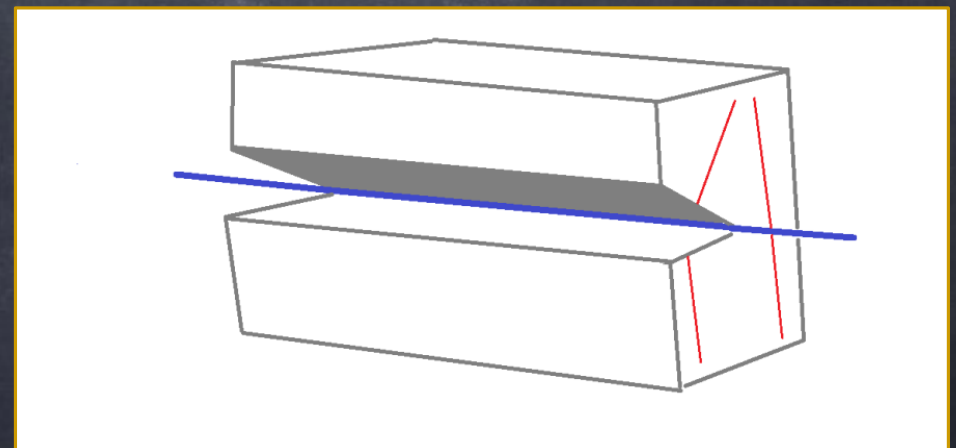
- Inclusion of interactions so far appears to require UV cutoff below cc scale

CC Problem

- My own opinion: Not yet clear conservative scalar-tensor-gauge models cannot work
- Must break link between vacuum energy (which we think is large) and universe's curvature (measured to be small)
- Problem: because vacuum is Lorentz invariant its stress energy $T_{mn} = c g_{mn}$ with Einstein eqs is an obstruction to small curvature

CC Problem

- More opinion: might break this link with extra dimensions of order micron in size (i.e. size of the cc)
- Large 4D Lorentz-invariant tension can curve extra dimensions instead of ours
- no explicit examples work (yet)
- Deviation of inverse square law: smoking gun



Time dependence



Time dependence

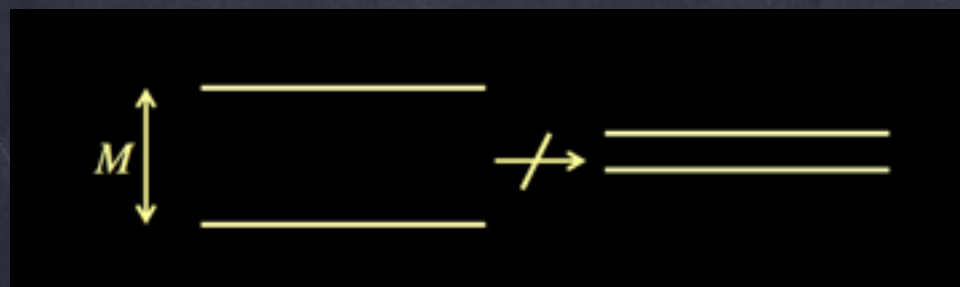
- Can EFTs apply to time-dependent situations where E is not conserved?
- Higher time-derivatives usually imply ghosts; does their absence constrain EFTs?
- What is the most efficient description of fluctuations about t -dependent background?

Time dependence

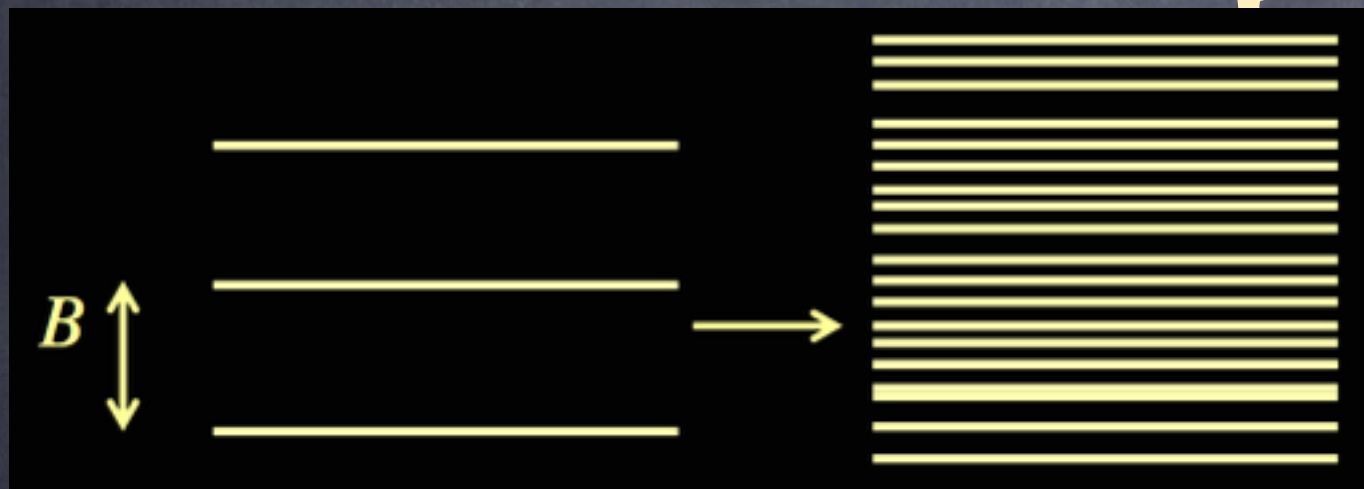
- Can EFTs apply to time-dependent situations?
- E not strictly conserved, but can still apply EFT reasoning if evolution is adiabatic:

$$\frac{\dot{\phi}}{\phi} \ll M$$

- Must also check other conditions (eg low energy) still apply as time evolves



Time dependence



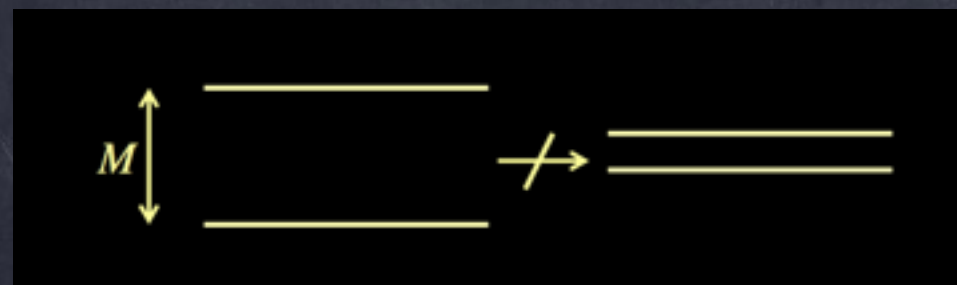
-dependent situations?

ed, but can still apply

EFT reasoning if evolution is adiabatic:

Related example: 'Transplanckian' issues; are not unique to gravity

- Must also check other conditions (eg low energy) still apply as time evolves



Time dependence

- Must EFTs be constrained not to have higher time derivatives? (Implicit to Horndesky-type models)
- Dangerous ghosts generically absent at fixed order in $1/M$

$$L = \dot{q}^2 + \ddot{q}^2/M^2$$

$$\ddot{q} + \dddot{q}/M^2 = 0$$

$$q(t) = A + Bt + Ce^{Mt} + De^{-Mt}$$

Time dependence

- Related (but not identical) issue: what EFT best describes fluctuations about time-dependent backgrounds: e.g. EFT for inflationary fluctuations
- Exploit breaking of time-translation invariance by background to identify leading low-energy contributions to CMB
- Reasoning similar to EFT for goldstone bosons in QCD and in condensed matter

Time

$$t \rightarrow t + \xi^0(t, \vec{x}) \quad \delta\phi \rightarrow \delta\phi + \dot{\phi} \xi^0$$

$$L = M_p^2 \dot{H} (\partial_\mu \pi) (\partial^\mu \pi)$$

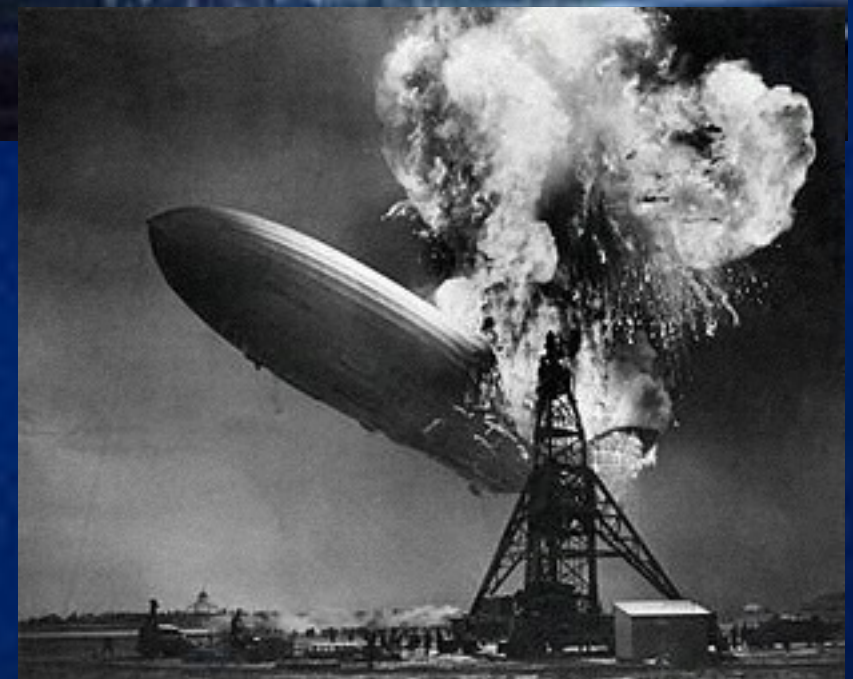
$$+ 2M_s^4 \left[\dot{\pi}^2 + \dot{\pi}^3 - \frac{\dot{\pi} (\partial_i \pi \partial^i \pi)}{a^2} \right] + \dots$$

- Related (but describes fluctuations about background)

Most constraining for single-field models, where few terms possible (Cheung et al)

- Reasoning similar to EFT for goldstone bosons in QCD and in condensed matter

Might there be Surprises?



EFT Surprises?

- No evidence for gravitational exceptionalism
- But gravitational situations explore aspects of EFTs in different regimes than in particle physics and so can contain surprises, some to do with t -dependence:
 - Adiabatic requirements for t -dependent EFTs
 - Instabilities can be features not bugs
 - Fluid-like systems, such as arise in LSS

EFT Surprises?

- Gravitational environments closer to effective description of particle in a medium than to traditional low-energy Wilsonian EFT
 - Are open systems when horizons are present, since degrees of freedom are excluded not based on conservation laws (so can entangle)
 - Generic difficulties computing late-time behaviour due to 'secular' effects and breakdown of perturbative tools
 - EFT exterior to black hole possibly nonlocal over horizon scales? (usual arguments against neednt apply)

Summary

- EFTs: Love them or Hate them, but use them!
- Embedding gravity into broader context allows assessment of theoretical error and contains useful clues
- Tools developed elsewhere in physics can be useful in gravitational applications
- Gravitational problems provide mind-broadening examples for EFT applications