Ab Initio Valence-Space Hamiltonians for Exotic Nuclei



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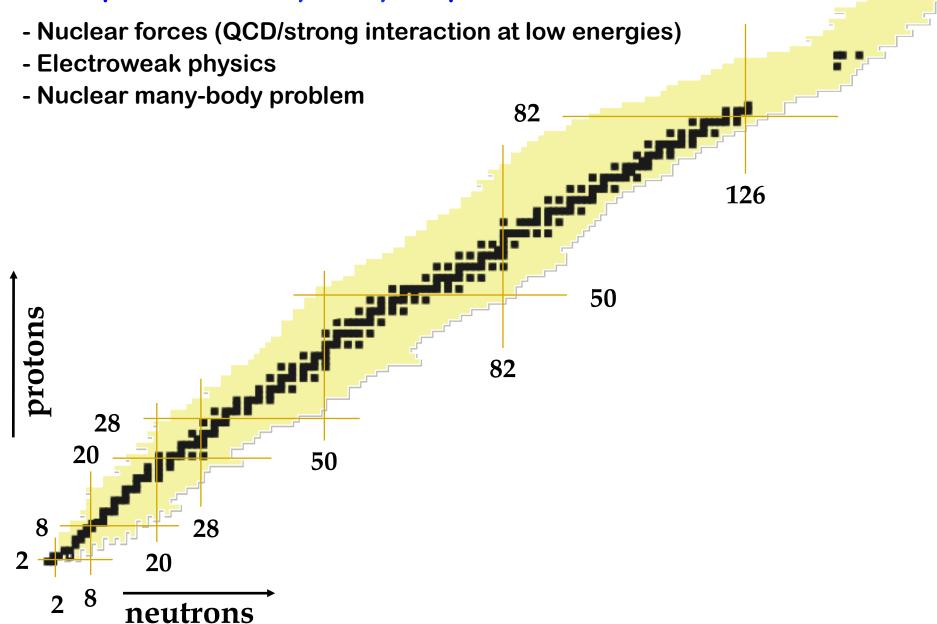


J. Menendez

Frontiers and Impact of Nuclear Science

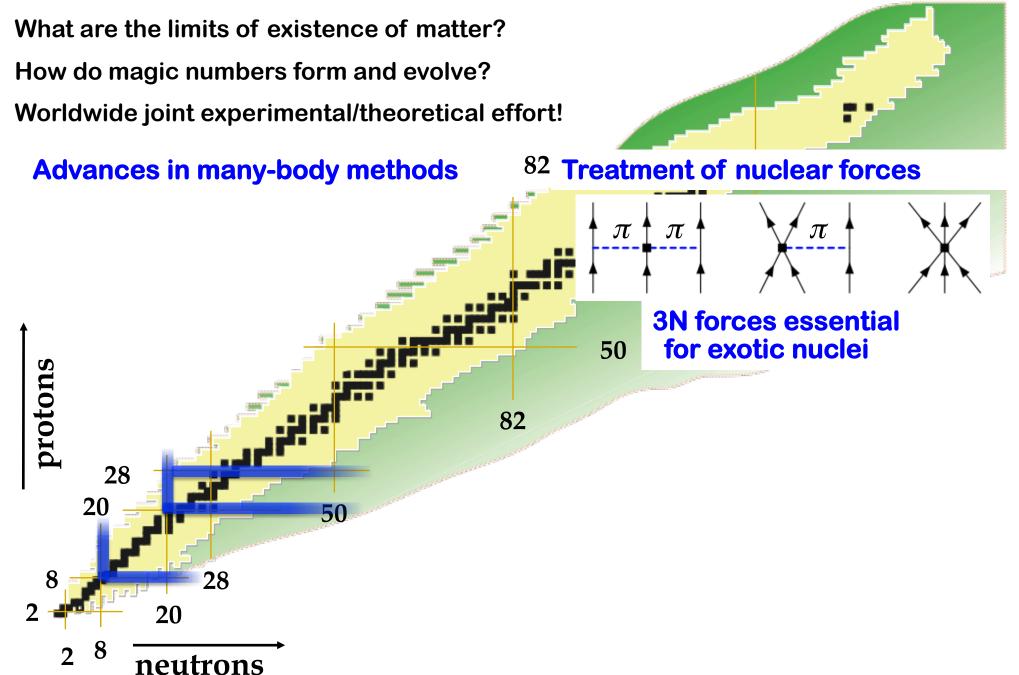
Aim of ab initio nuclear theory:

Develop unified first-principles picture of structure and reactions



Medium- and Heavy-Mass Exotic Nuclei

What are the properties of proton/neutron-rich matter?



The Nuclear Many-Body Problem

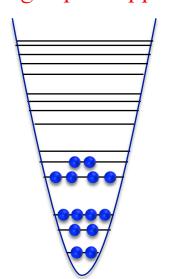
Nucleus strongly interacting many-body system -A-body problem impossible

$$H\psi_n = E_n \psi_n$$

Quasi-exact solutions in light nuclei (GFMC, (IT)NCSM, ...)

Large space: controlled approximations to full Schrödinger Equation

Large-space approach



Limited range:

Closed shell ±1

Even-even

Limited properties:

Ground states only

Some excited state

Coupled Cluster

In-Medium SRG

Green's Function

The Nuclear Many-Body Problem

Nucleus strongly interacting many-body system -A-body problem impossible

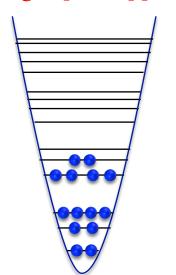
$$H\psi_n = E_n \psi_n$$

Quasi-exact solutions in light nuclei (GFMC, (IT)NCSM, ...)

Large space: controlled approximations to full Schrödinger Equation

Valence space: diagonalize exactly with reduced number of degrees of freedom

Large-space approach



Limited range:

Closed shell ±1

Even-even

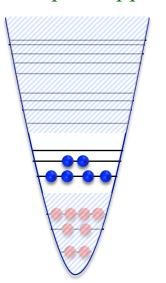
Limited properties:

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Coupled Cluster
In-Medium SRG
Green's Function

Valence-space approach



All nuclei near closed-shell cores

All properties:

Ground states

Excited states

EW transitions

Coupled Cluster
In-Medium SRG
Perturbation Theory

In-Medium Similarity Renormalization Group

Continuous unitary trans (basis change) decouples "off-diagonal" physics

$$H(s) = U(s)HU^{\dagger}(s) \equiv H^{\mathrm{d}}(s) + H^{\mathrm{od}}(s) \to H^{\mathrm{d}}(\infty)$$

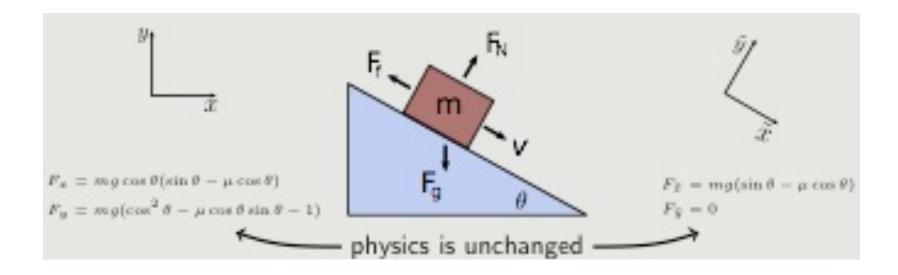
Interaction in new basis is simple

In-Medium SRG

Continuous unitary trans (basis change) decouples "off-diagonal" physics

$$H(s) = U(s)HU^{\dagger}(s) \equiv H^{\mathrm{d}}(s) + H^{\mathrm{od}}(s) \to H^{\mathrm{d}}(\infty)$$

Interaction in new basis is simple



Can always write $U = e^{\eta}$, for some generator η

For incline plane:
$$\eta = \begin{pmatrix} 0 & \theta \\ -\theta & 0 \end{pmatrix}$$

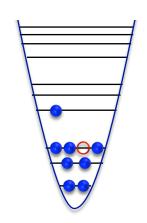
Life Is Difficult: Particle/Hole Excitations

Consider basis states as excitations from uncorrelated reference state

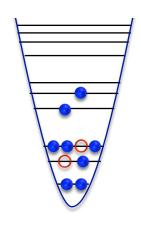
Ref. Slater Determinant

Unoccupied (Particles) Occupied (Holes) $|\Phi_0\rangle = \prod a_i^{\dagger} |0\rangle$

1p-1h excitation 2p-2h excitation



$$\Phi_i^a \rangle = a_a^{\dagger} a_i |\Phi_0 \rangle$$

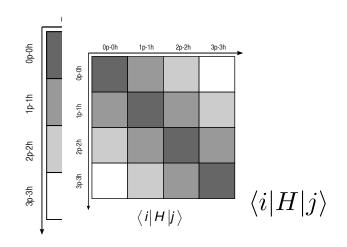


$$|\Phi_i^a\rangle = a_a^{\dagger} a_i |\Phi_0\rangle \qquad |\Phi_{ij}^{ab}\rangle = a_a^{\dagger} a_i a_b^{\dagger} a_j |\Phi_0\rangle$$

Hamiltonian schematically in terms of ph excitations

Ground-state coupled to excitations is difficult

$$H^{\text{od}} = \langle p|H|h\rangle + \langle pp|H|hh\rangle + \cdots + \text{h.c.}$$

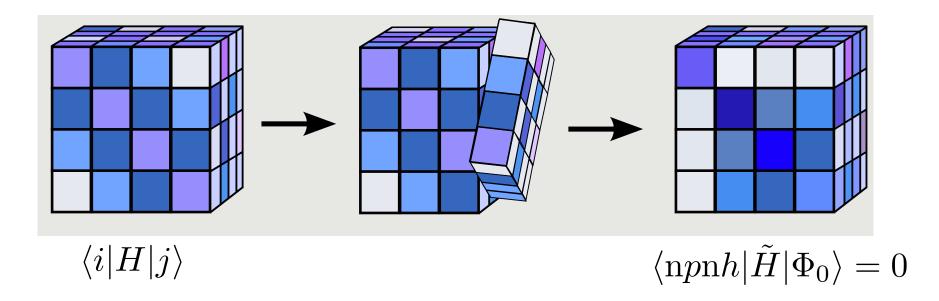


In-Medium SRG for Nuclei

For nuclear Hamiltonian, take

$$U = e^{\eta}$$
 with $\eta = \frac{H_{\text{od}}}{\Delta} + \text{h.c.}$

Perform multiple rotations $U_N = e^{\eta_N} \cdots e^{\eta_2} e^{\eta_1}$ until $\eta_N = 0$



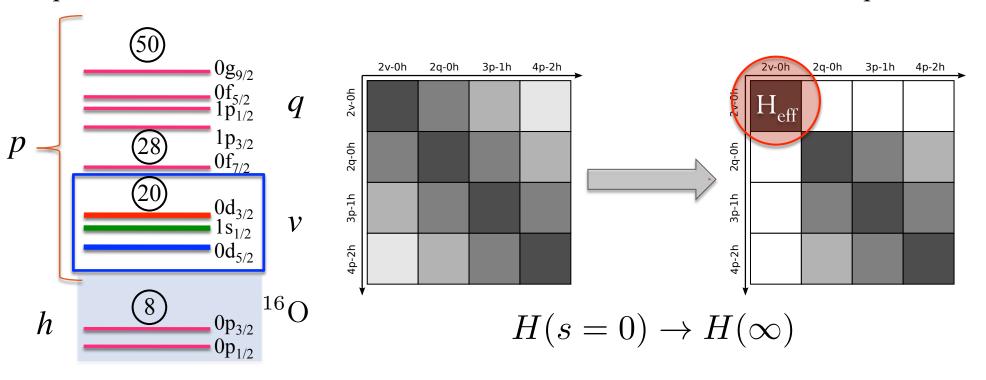
Fully correlated ground state: one matrix element $\langle \Phi_0 | \hat{H} | \Phi_0 \rangle$

Also flow equation approach
$$\frac{dH(s)}{ds} = [\eta(s), H(s)]$$

IM-SRG for Valence-Space Hamiltonians

Tsukiyama, **Bogner**, Schwenk, PRC (2012)

Separate p states into valence states (v) and those above valence space (q)

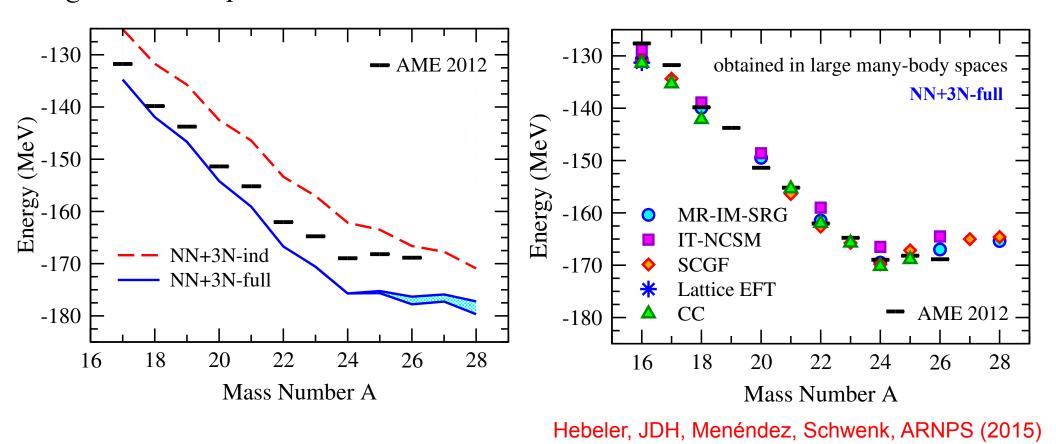


Redefine H^{od} to decouple valence space from excitations outside v

$$H^{\mathrm{od}} = \langle p|H|h\rangle + \langle pp|H|hh\rangle + \langle v|H|q\rangle + \langle pq|H|vv\rangle + \langle pp|H|hv\rangle + \mathrm{h.c.}$$
 Core Energy Single-particle energies Two-body valence particle interaction matrix elements

Ground-State Energies in Oxygen Isotopes

Large/valence-space methods with same SRG-evolved NN+3N-full forces



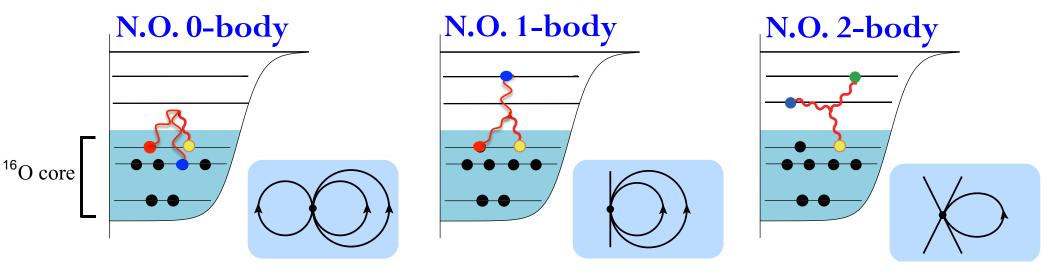
Agreement between all methods with same input forces

Clear improvement with NN+3N-full

Still significant discrepancy between valence/large-space results

How Do We Handle 3N Forces?

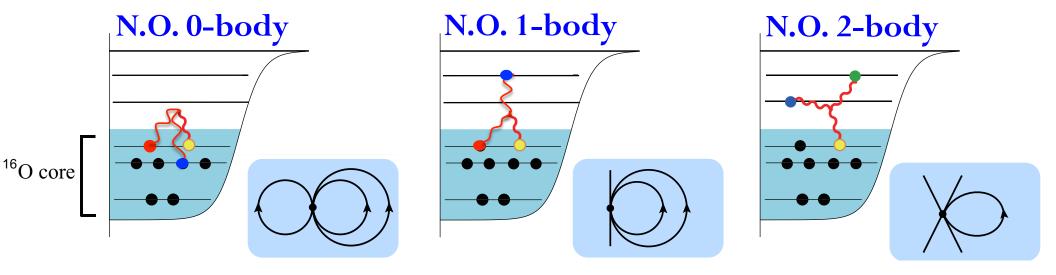
Normal-ordered 3N: contribution from core with valence particles



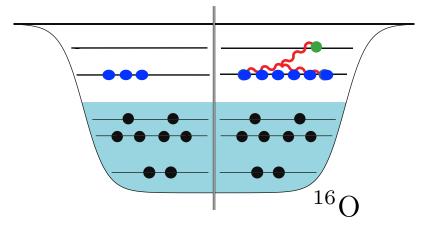
Neglect 3N forces between valence particles – significant as $N_v \sim N_c$

Targeted Normal Ordering

Normal-ordered 3N: contribution from core with valence particles

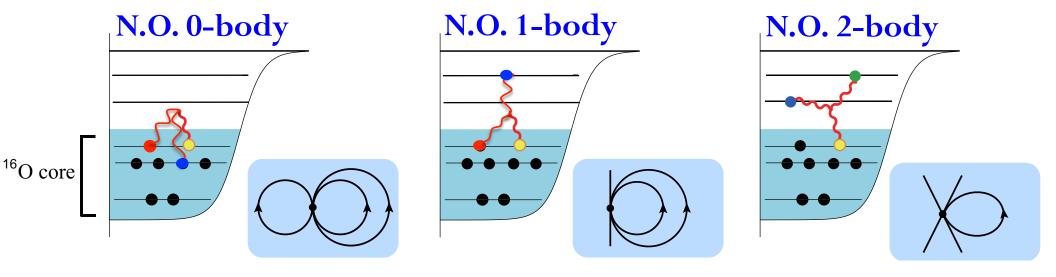


Neglect 3N forces between valence particles – significant as $N_v \sim N_c$ Capture these effects with new Targeted N.O.

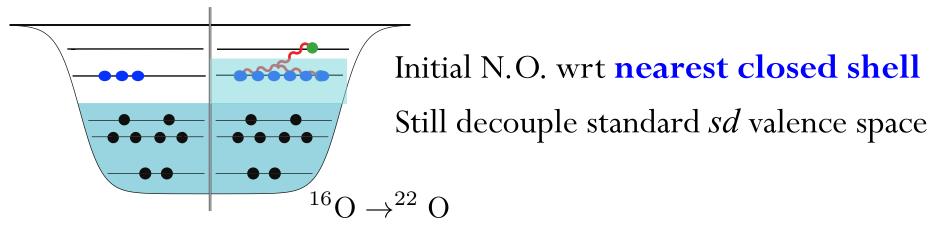


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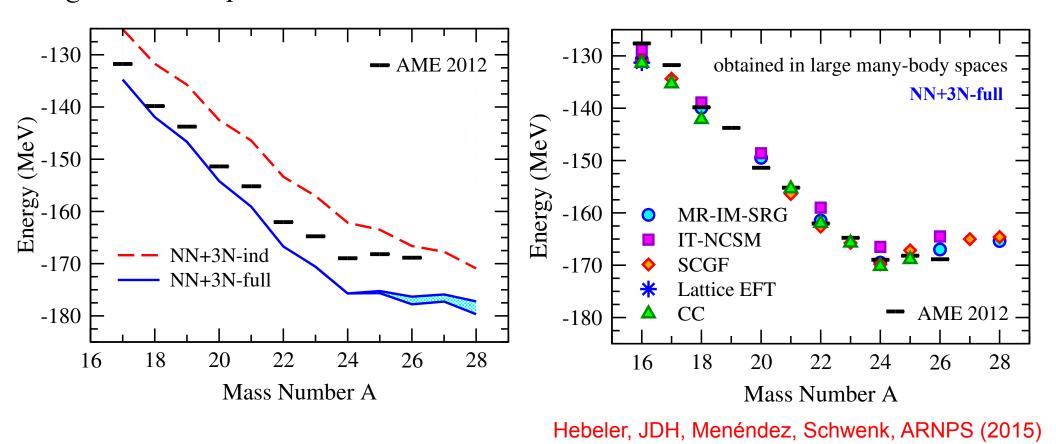


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Ground-State Energies in Oxygen Isotopes

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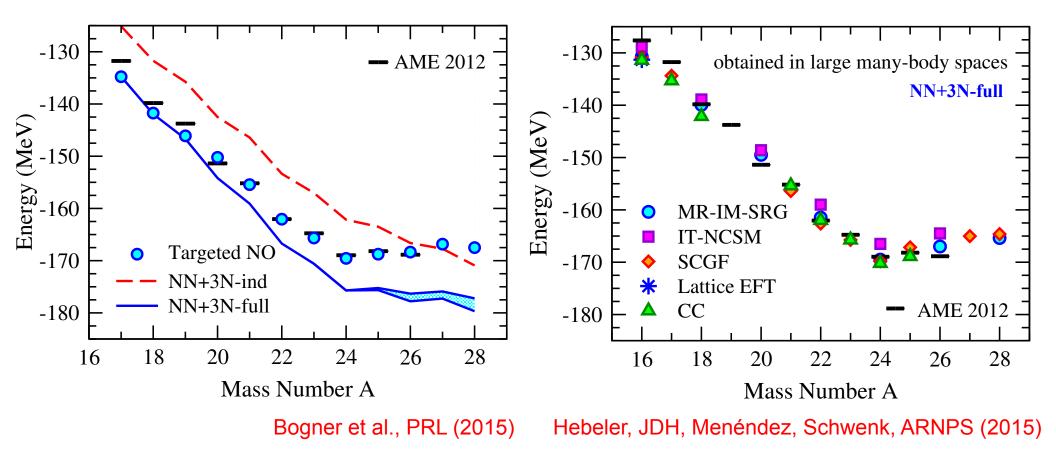
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Targeted N.O. in Oxygen Isotopes

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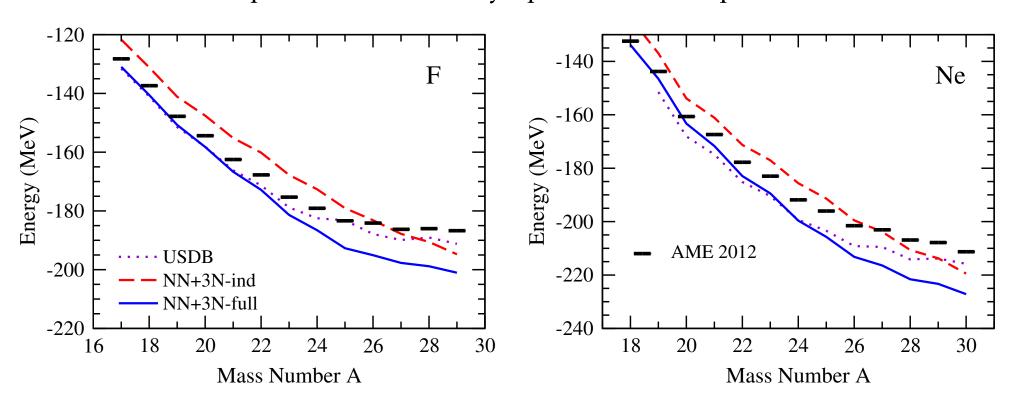


Improved method to capture neglected 3N forces in valence space

"Targeted" IMSRG results agree well with data and large-scale methods

Beyond Semi-Magic: Ground States of F/Ne

IM-SRG valence-space results for fully open F/Ne isotopes

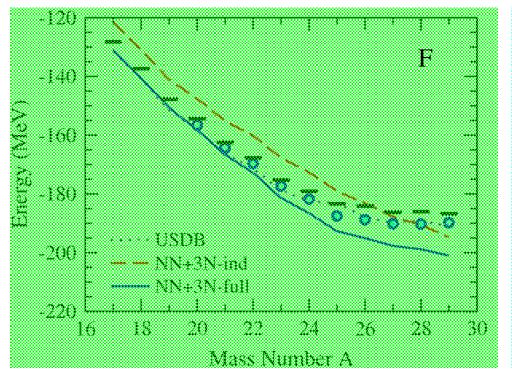


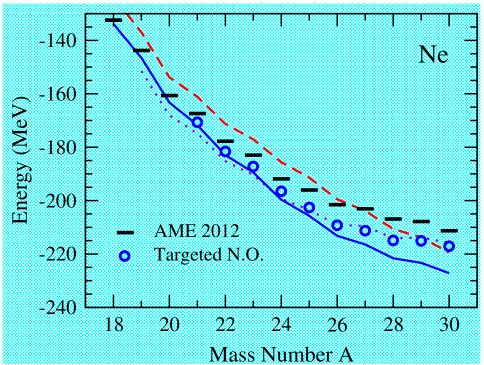
Stroberg et al., arXiv:1511.03802

NN+3N-full improves agreement with experiment; overbound past N=14

Beyond Semi-Magic: Ground-States of F/Ne

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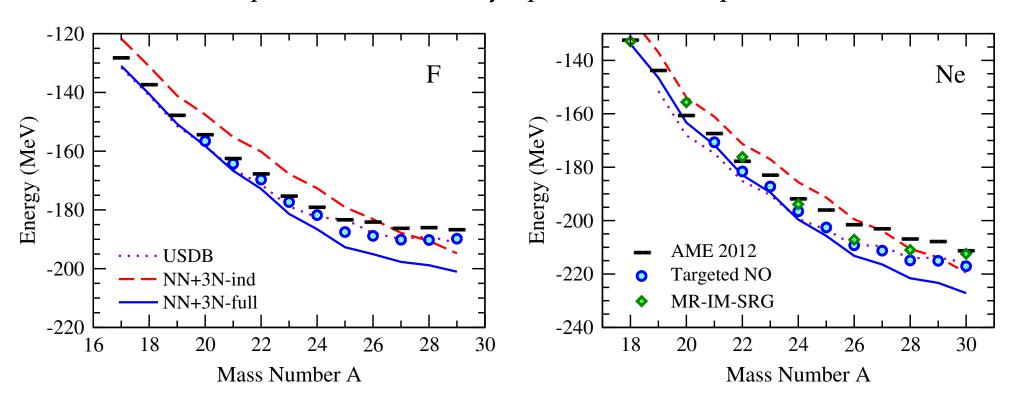


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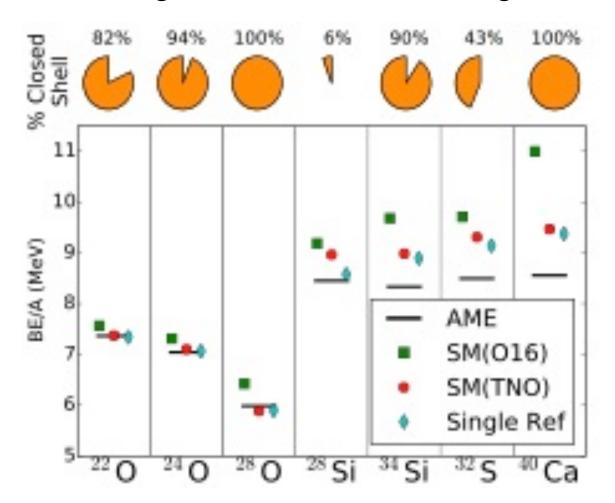


Stroberg et al., arXiv:1511.03802

NN+3N-full improves agreement with experiment; overbound past N=14 Targeted N.O. results further improved — similar to phenomenology Good agreement with large-space MR-IM-SRG!

Ground States from Oxygen to Calcium

3N force effects significant as N_v becomes large



Stroberg et al., in prep

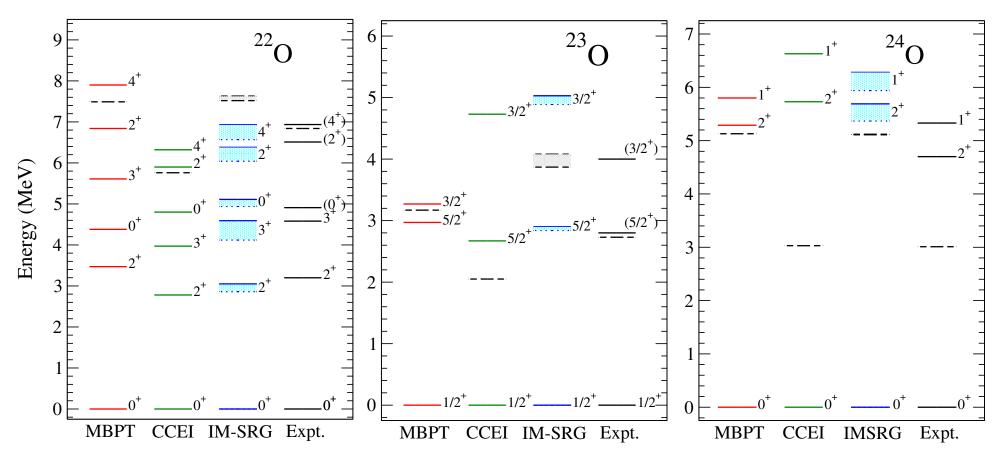
Targeted N.O. valence-space results agrees with large-space in all cases!

²⁸Si not good closed shell (single ref. incorrect)

Discrepancy with experiment from initial nuclear interactions

Comparison with MBPT/CCEI Oxygen Spectra

Oxygen spectra: Effective interactions from Coupled-Cluster theory



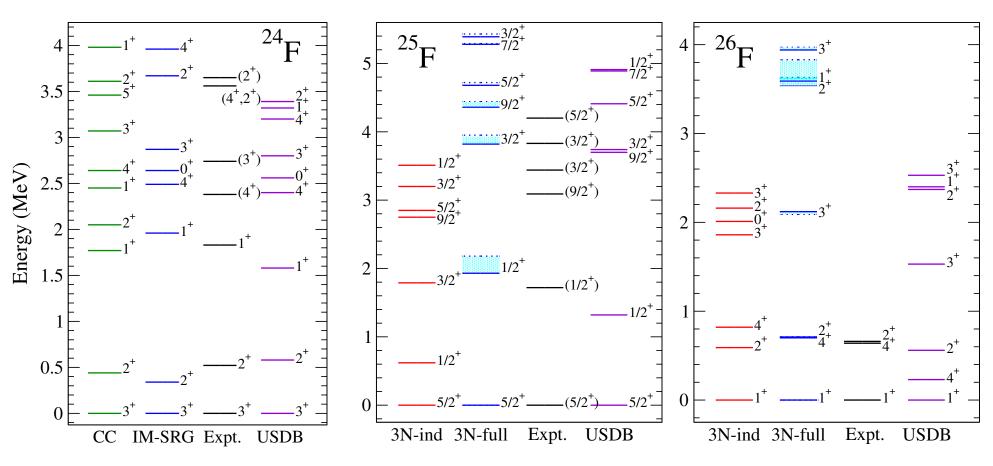
Hebeler, JDH, Menéndez, Schwenk, ARNPS (2015)

MBPT in extended valence space

IM-SRG/CCEI spectra agree within ~300 keV

Doubly Open Shell: Neutron-Rich F Spectra

Fluorine spectroscopy: **MBPT** and **IM-SRG** (*sd* shell) from NN+3N forces

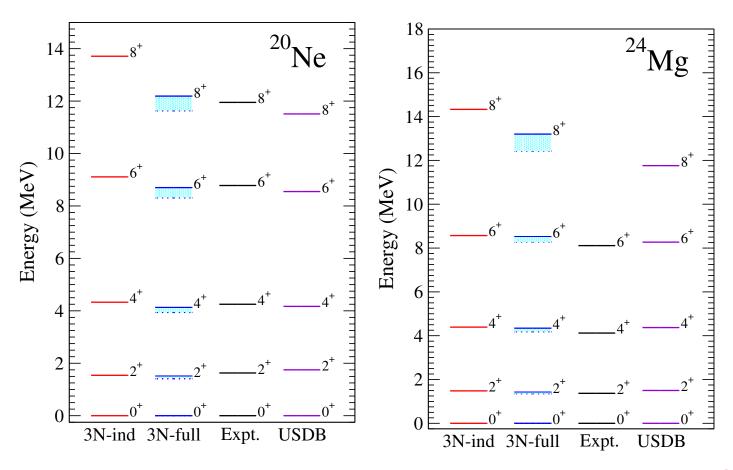


Stroberg et al., arXiv:1511.03802

IM-SRG: competitive with phenomenology, good agreement with data

Deformed Systems: ²⁰Ne and ²⁴Mg

Ground-state rotational band for well-known deformed nuclei

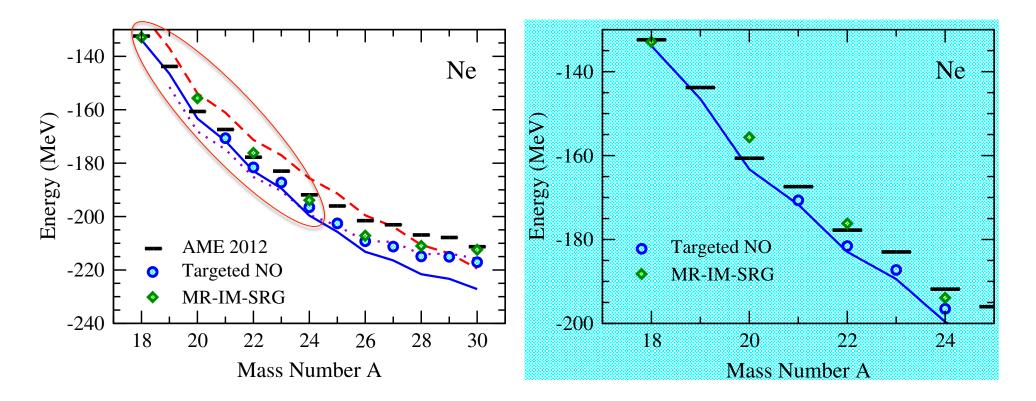


Stroberg et al., arXiv:1511.03802

IM-SRG: **competitive with phenomenology**, good agreement with data Further observables (quadrupole moments, E2 transitions) needed

Deformation with Large-Space MR-IM-SRG?

Ground states in light neon isotopes – clear discrepancies in ^{20,22}Ne

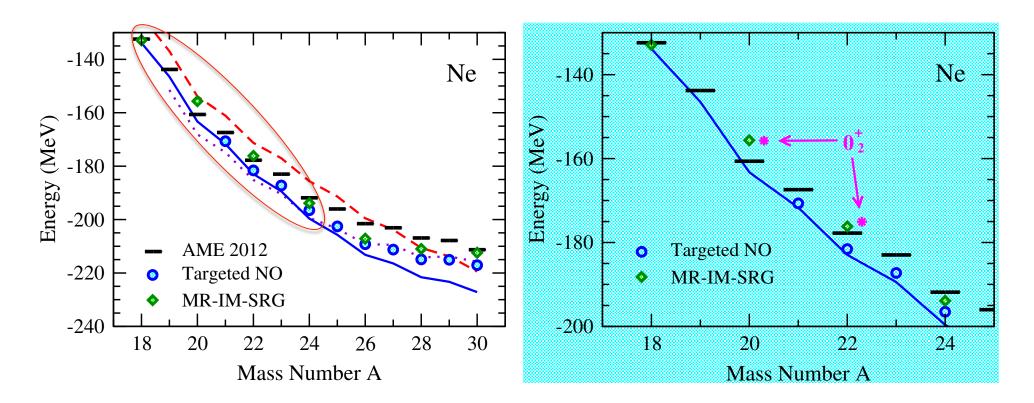


MR-IM-SRG built on spherical reference state

Not expected to produce deformed ground states – not a problem for SM

Deformed Systems: ²⁰Ne and ²⁴Mg

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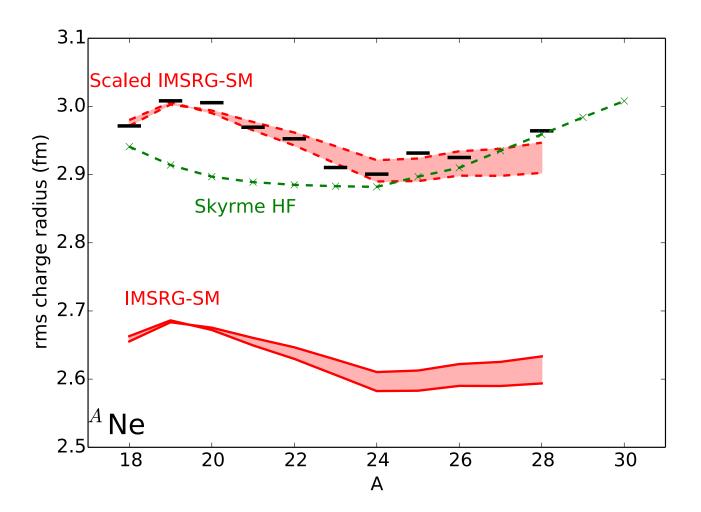
MR-IM-SRG built on spherical reference state

Not expected to produce deformed ground states — not a problem for SM First (likely spherical) excited 0⁺ SM state agrees remarkably with MR-IM-SRG Indicates SM captures physics of deformed ground state

RMS Charge Radii in sd Shell Model

Previous SM radii calculations rely on empirical input or as relative to core

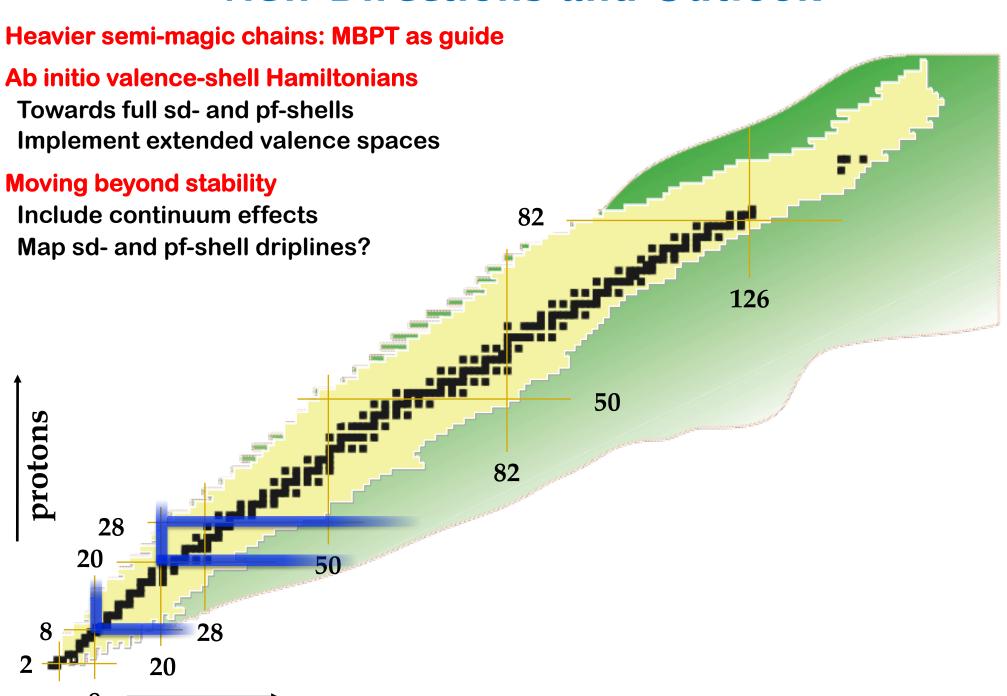
Radii for stable sd-shell nuclei calculated in shell model NN+3N



Stroberg et al., in prep

Next: general tensor operators M1, E2, GT, double-beta decay

New Directions and Outlook



neutrons

New Directions and Outlook

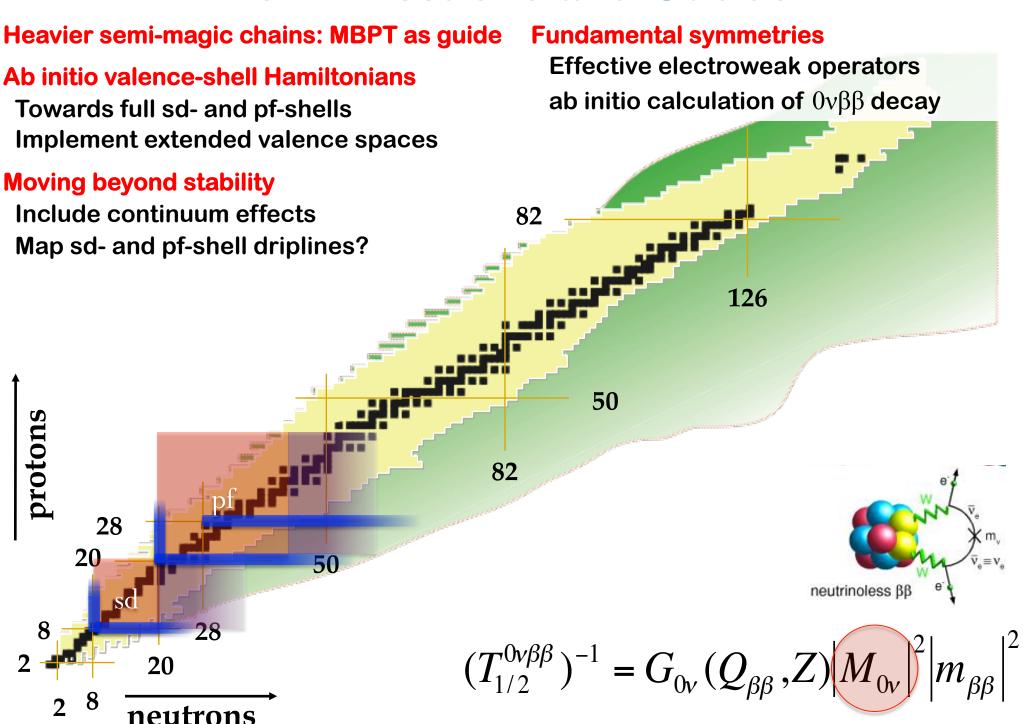
Heavier semi-magic chains: MBPT as guide **Ab initio valence-shell Hamiltonians** Towards full sd- and pf-shells Implement extended valence spaces **Moving beyond stability** Include continuum effects 82 Map sd- and pf-shell driplines? 126 50 protons 82 28 20 50 8

28

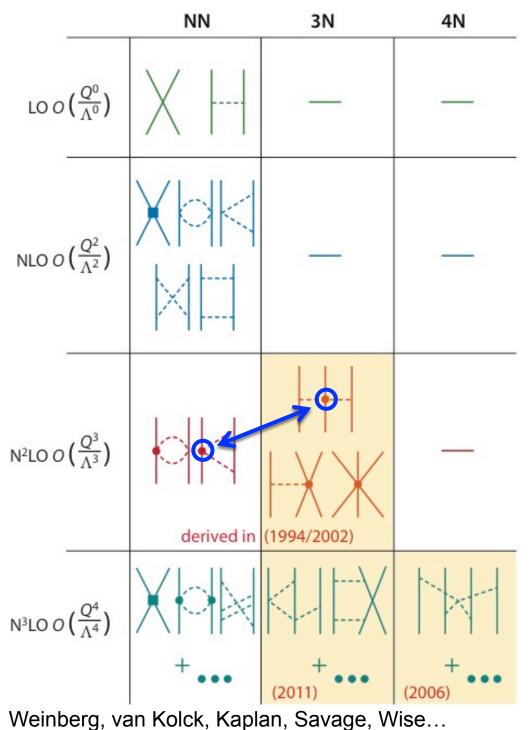
neutrons

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New Directions and Outlook



Chiral Effective Field Theory: Nuclear Forces



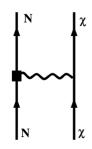
Nucleons interact via pion exchanges and contact interactions

Consistent treatment of NN, 3N,...

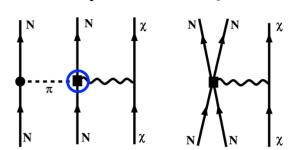
NN couplings fit to scattering data 3N couplings fit to 3/4-body systems

Consistent EW/WIMP interactions

one-body currents at Q^0 and Q^2



+ two-body currents at Q³



Normal-Ordered Hamiltonian

Now rewrite exactly the initial Hamiltonian in normal-ordered form

$$H_{\text{N.O.}} = E_0 + \sum_{ij} f_{ij} \left\{ a_i^{\dagger} a_j \right\} + \frac{1}{4} \sum_{jkl} \Gamma_{ijkl} \left\{ a_i^{\dagger} a_j^{\dagger} a_l a_k \right\} + \frac{1}{36} \sum_{ijklmn} W_{ijklmn} \left\{ a_i^{\dagger} a_j^{\dagger} a_k^{\dagger} a_l a_m a_n \right\}$$

Normal-ordered Hamiltonian w.r.t. reference state

Loop = sum over occupied states

Include dominant 1-,2-,3-body physics in NO

IM-SRG: Flow Equation Formulation

Define U(s) implicitly from particular choice of generator:

$$\eta(s) \equiv (dU(s)/ds) U^{\dagger}(s)$$

chosen for desired decoupling behavior - e.g.,

$$\eta_{\scriptscriptstyle I}(s) = \left[H^{\mathrm{d}}(s), H^{\mathrm{od}}(s)
ight]$$
 Wegner (1994)

Solve **flow equation** for Hamiltonian (coupled DEs for 0,1,2-body parts)

$$\frac{\mathrm{d}H(s)}{\mathrm{d}s} = [\eta(s), H(s)] \qquad H(s) = E_0(s) + f(s) + \Gamma(s) + \cdots$$

Hamiltonian and generator truncated at 2-body level: IM-SRG(2)

0-body flow drives uncorrelated ref. state to fully correlated ground state

$$E_0(\infty) \to \operatorname{Core} \operatorname{Energy}$$

Ab initio method for energies of **closed-shell systems**

New Approach: Magnus Expansion

Morris, Parzuchowski, Bogner, PRC (2015)

Magnus expansion: explicitly construct unitary transformation

$$U(s) = \exp \Omega(s)$$

With flow equation:

$$\frac{\mathrm{d}\Omega(s)}{\mathrm{d}s} = \eta(s) + \frac{1}{2} \left[\Omega(s), \eta(s)\right] + \frac{1}{12} \left[\Omega(s), \left[\Omega(s), \eta(s)\right]\right] + \dots$$

Leads to commutator expression for evolved Hamiltonian

$$H(s) = e^{\Omega(s)}He^{-\Omega(s)} = H + \frac{1}{2}\left[\Omega(s), H\right] + \frac{1}{12}\left[\Omega(s), \left[\Omega(s), H\right]\right] + \cdots$$

Nested commutator series – in practice truncate numerically

All calculations truncated at normal-ordered two-body level

Effective Operators

Keep unitary transformation from evolution of Hamiltonian

Can generalize to arbitrary operators

$$H(s) = e^{\Omega(s)}He^{-\Omega(s)} = H + \frac{1}{2}\left[\Omega(s), H\right] + \frac{1}{12}\left[\Omega(s), \left[\Omega(s), H\right]\right] + \cdots$$



$$\mathcal{O}^{\Lambda}(s) = e^{\Omega(s)} \mathcal{O}^{\Lambda} e^{-\Omega(s)} = \mathcal{O}^{\Lambda} + \frac{1}{2} \left[\Omega(s), \mathcal{O}^{\Lambda} \right] + \frac{1}{12} \left[\Omega(s), \left[\Omega(s), \mathcal{O}^{\Lambda} \right] \right] + \cdots$$

Must work out normal-ordered operators in *J*-coupled basis

First apply to scalar operators

EO Transitions and Radii

Seldom calculated in nuclear shell model

In single HO shell:

$$|\langle f|\rho_{E0}|i\rangle|^2 \propto \delta_{ij} \text{ where } \rho_{E0} = \frac{1}{e^2R} \sum_i e_i r_i^2$$

Must resort to phenomenological gymnastics

IM-SRG: straightforward to calculate effective valence-space operator:

$$\rho_{E0}(s) = e^{\Omega(s)} \rho_{E0} e^{-\Omega(s)} = \rho_{E0} + \frac{1}{2} [\Omega(s), \rho_{E0}] + \cdots$$

Commutators induce important higher-order and two-body parts

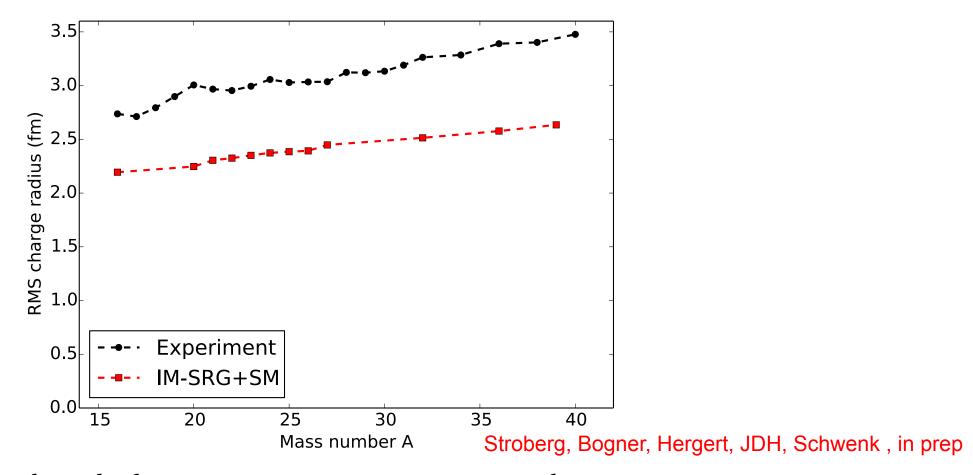
$$\Theta$$
 + $\frac{1}{\Omega}$ Θ + $\frac{1}{\Omega}$ + ...

Quantify importance of induced higher-body contributions!

RMS Charge Radii in sd Shell Model

Previous SM radii calculations rely on empirical input or as relative to core

Absolute radii for entire sd shell calculated in shell model NN+3N



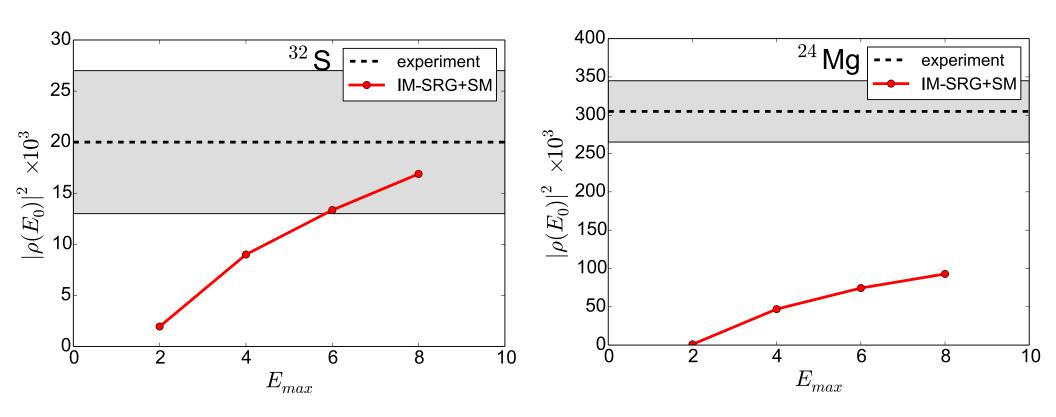
Benchmarked against NCSM in various SM codes

~10% too small – deficiencies expected to come from initial Hamiltonian

Two-body part important 15-20%

EO Transitions in sd Shell Model

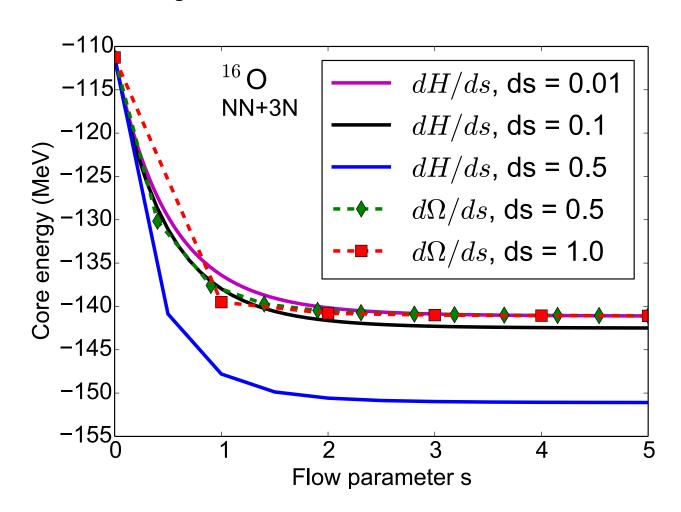
Preliminary results in *sd* shell:



Promising but need additional benchmarks

Magnus vs Flow-Equation

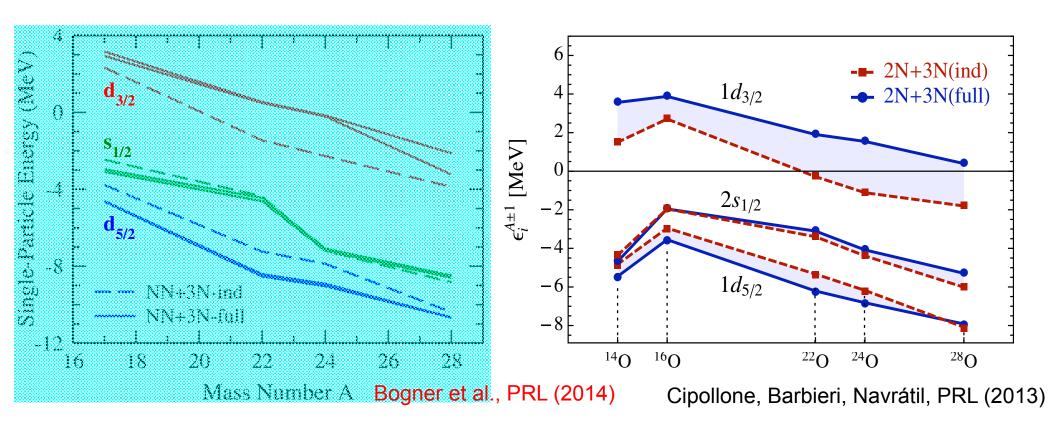
Variation of step size



Evident error accumulation in flow-equation for small step sizes Magnus: rapid convergence, independent of step size

Oxygen Dripline Mechanism

Self-consistent Green's Function with same SRG-evolved NN+3N forces



Robust mechanism driving dripline behavior

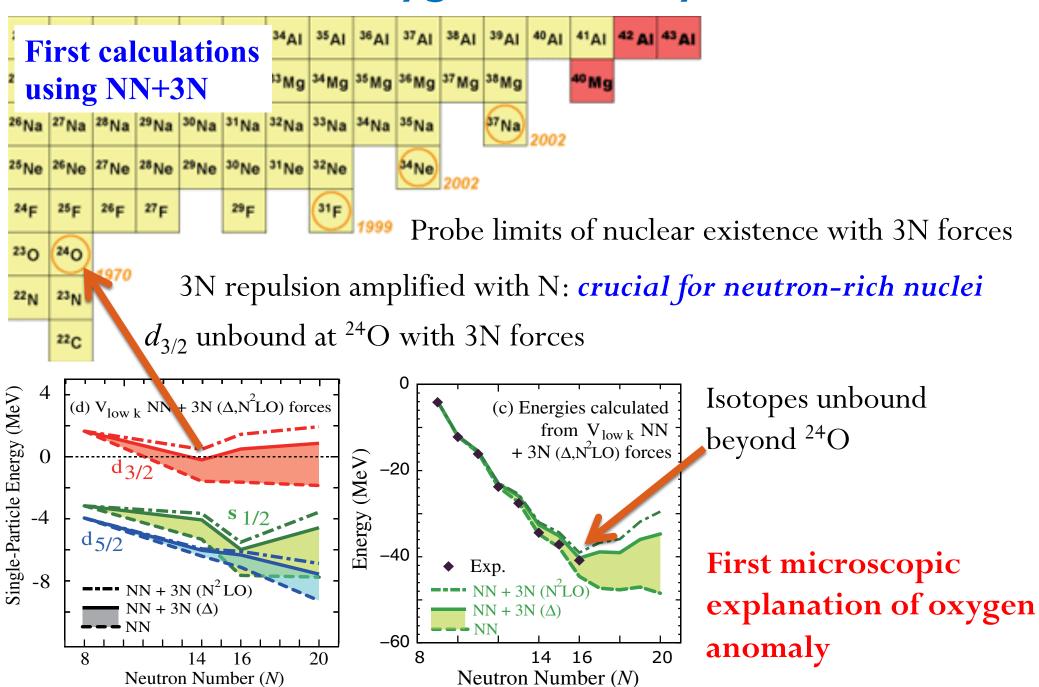
3N repulsion raises $d_{3/2}$, lessens decrease across shell

Similar to first MBPT NN+3N calculations in oxygen

Perturbative Approach

- 1) Effective Hamiltonian: sum excitations outside valence space to MBPT(3)
- 2) Self-consistent single-particle energies
- 3) Harmonic-oscillator basis of 13-15 major shells: converged
- 4) NN and 3N forces from chiral EFT

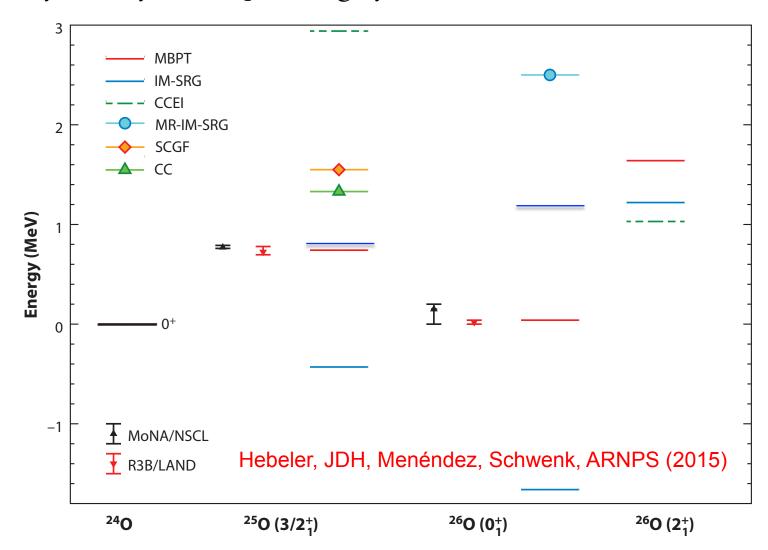
Oxygen Anomaly



Otsuka, Suzuki, JDH, Schwenk, Akaishi, PRL (2010)

Beyond the Oxygen Dripline

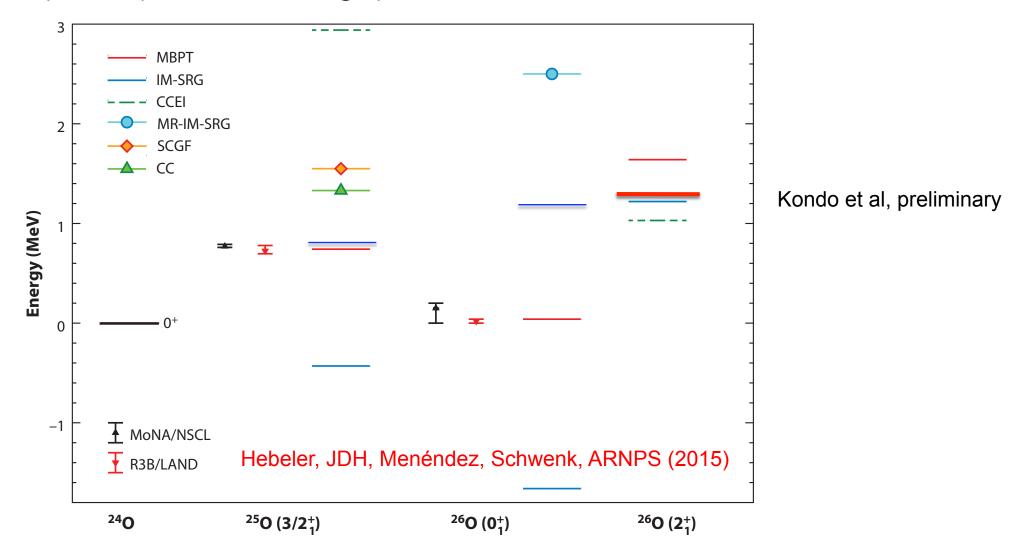
Physics beyond dripline highly sensitive to 3N forces and continuum effects



Prediction of low-lying 2⁺ in ²⁶O (recently measured at RIKEN)

Beyond the Oxygen Dripline

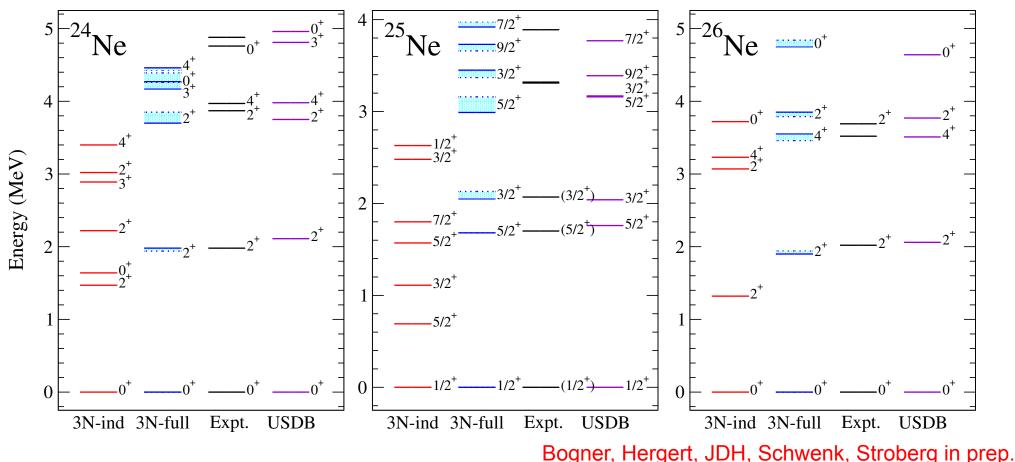
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Doubly Open Shell: Neutron-Rich Ne Spectra

Neon spectra: extended-space MBPT and IM-SRG (sd shell)



Bogner, Hergert, 3DH, Schwenk, Stroberg in prep

NN+3N-ind: clear deficiencies

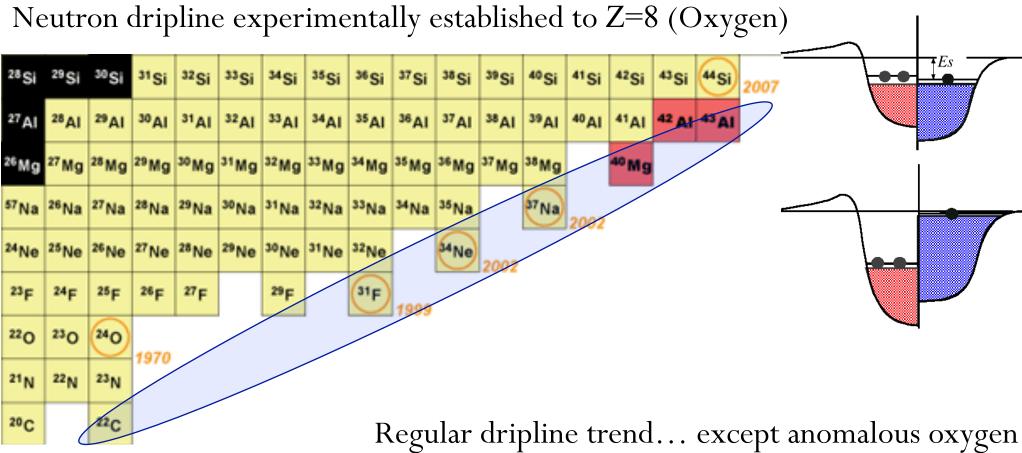
NN+3N-full: **competitive with phenomenology**, good agreement with data

Limits of Nuclear Existence: Oxygen Anomaly

Where is the nuclear dripline?

Limits defined as last isotope with positive neutron separation energy

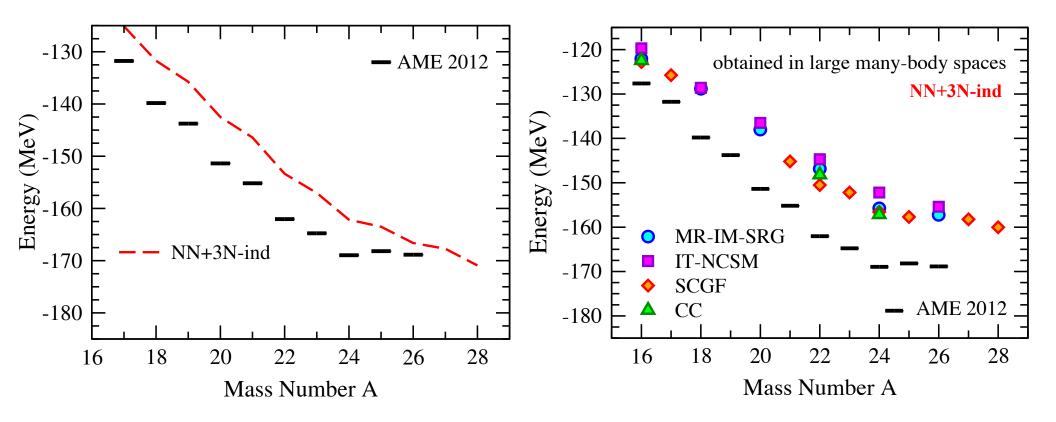
- Nucleons "drip" out of nucleus



Adding one proton binds 6 additional neutrons

Ground-State Energies in Oxygen Isotopes

Large/valence-space methods with same SRG-evolved NN+3N-ind forces



Agreement between all methods with same input forces

No reproduction of oxygen dripline in any case