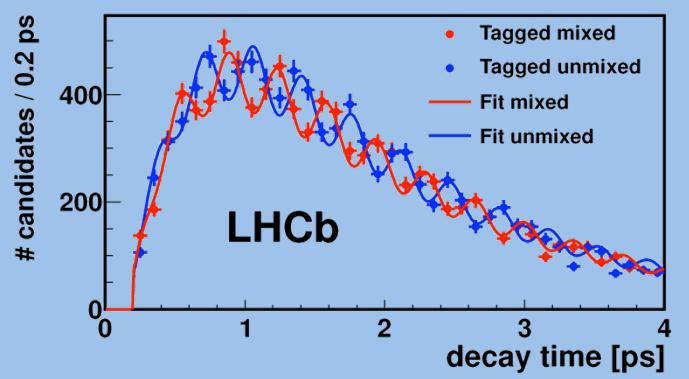


B-Flavour Physics in LHCb



Marcel Merk
On behalf of the LHCb collaboration
54th International Winter Meeting on Nuclear Physics
Bormio, Jan 25-29, 2016

Outline

1. Flavour Physics

2. The LHCb Experiment

3. CP Violation Measurements

- a) CP Violation in B-mixing: a_{sl}^d and a_{sl}^s
- b) Direct CP Violation: CKM angle γ
- c) Time dependent CP violation: CKM angles ϕ_d and ϕ_s

4. B Decay Rates

- a) Very Rare Decays: $B \rightarrow \mu\mu$
- b) Rare Decays: $b \rightarrow s$ quark transition
- c) Lepton Flavour Universality tests

5. Outlook

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Flavour Physics: perhaps a bit difficult...?

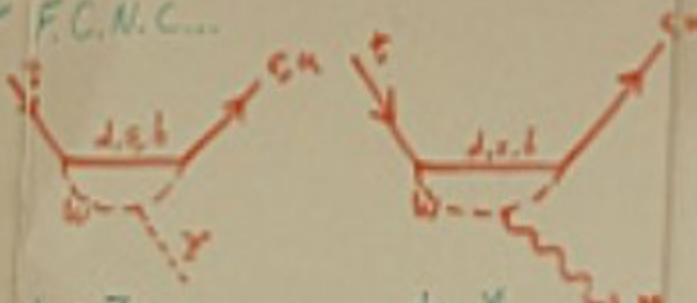
$$t \rightarrow W b$$
$$BR(t \rightarrow W b) = \frac{\Gamma(t \rightarrow W b)}{\Gamma(t \rightarrow W q)}$$



$$\hat{=} \frac{|V_{tb}|^2}{|V_{cb}|^2 + |V_{cb}|^2 + |V_{ts}|^2}$$

$$\approx \frac{(0.9945)^2}{(0.0079)^2 + (0.04)^2 + (0.7745)^2} \\ \approx 99.82\%$$

but F.C.N.C...



$t \rightarrow Z_c$
 $t \rightarrow Z_h$

$t \rightarrow Y_c$
 $t \rightarrow Y_h$

wavy line

$C_{tb} C_{t\bar{b}}$...

$- S_{tb} C_{t\bar{b}} - C_{tb} S_{t\bar{b}} S_{t\bar{b}} e^{i\gamma} \dots$

$\dots - S_{tb} C_{t\bar{b}} - C_{tb} S_{t\bar{b}} S_{t\bar{b}} e^{i\gamma} \dots$

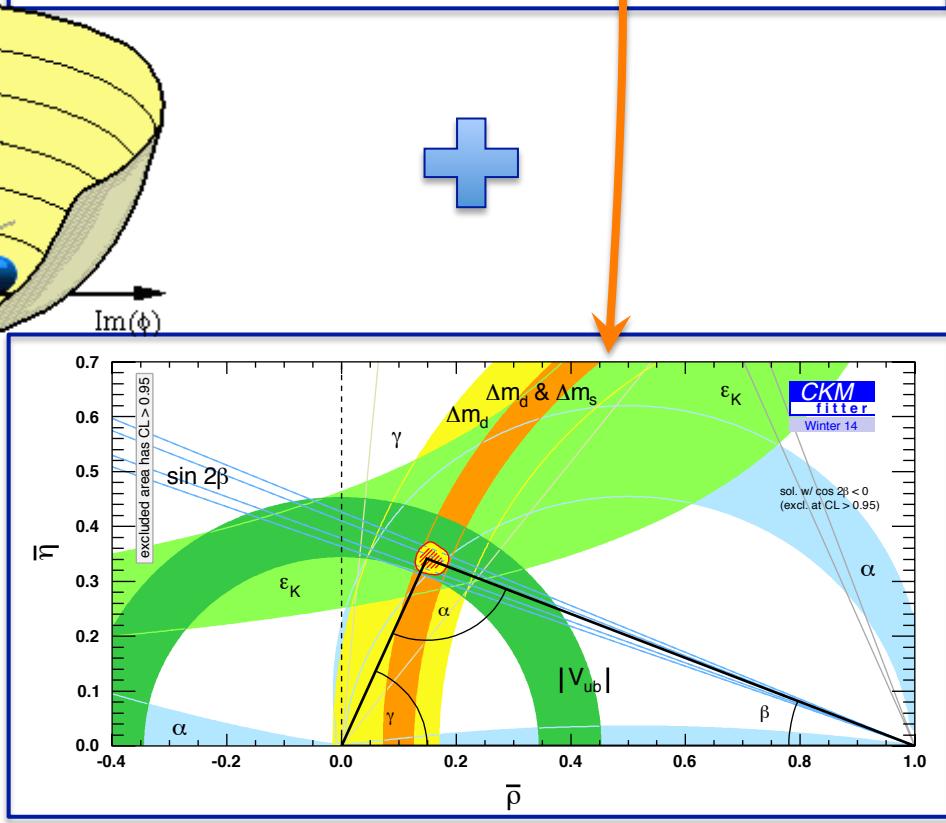
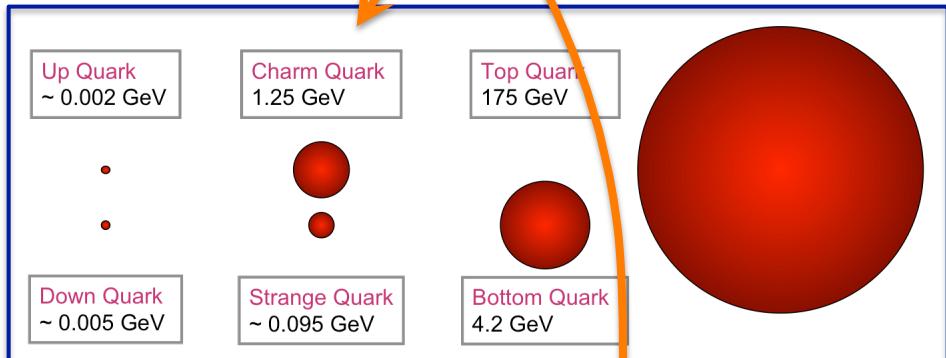
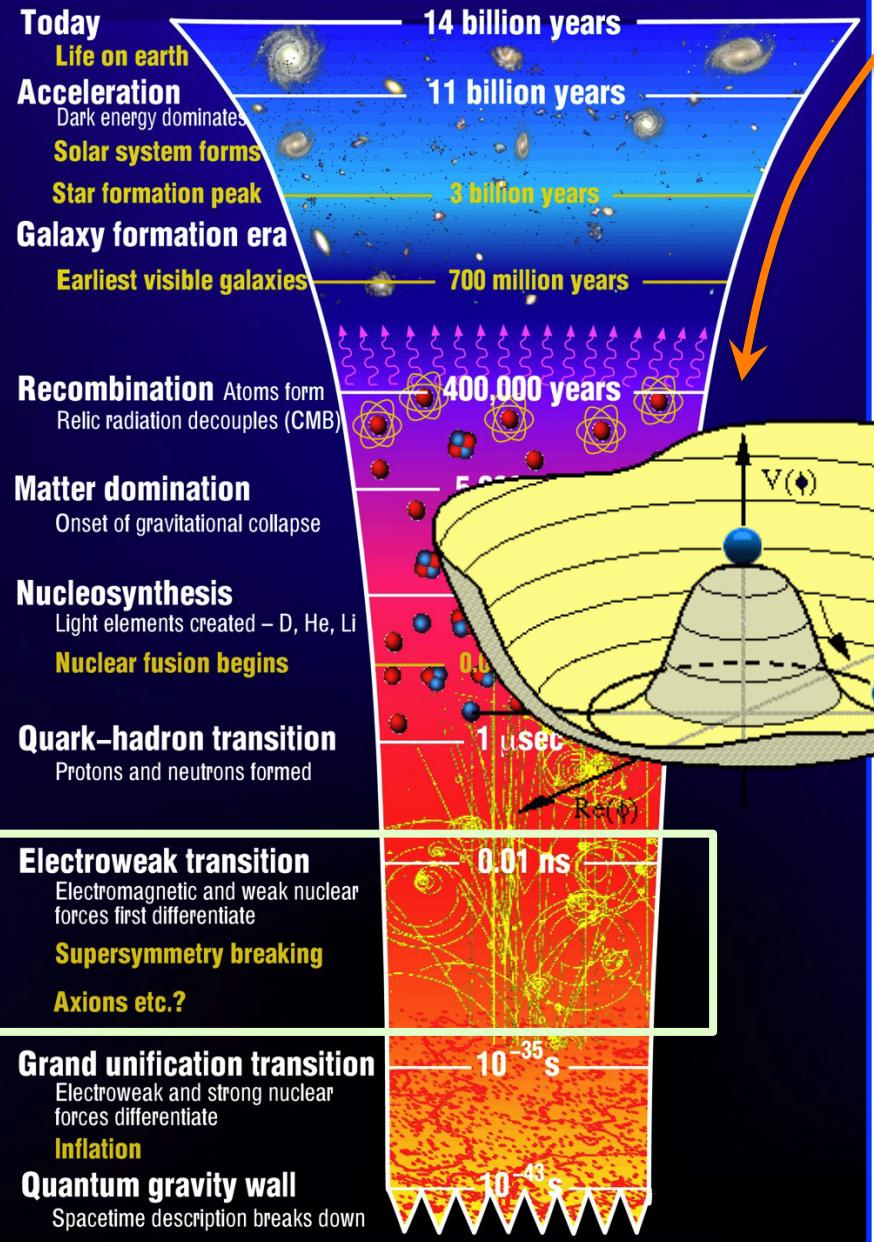
$\dots - S_{tb} C_{t\bar{b}} - C_{tb} S_{t\bar{b}} S_{t\bar{b}} e^{i\gamma} \dots$

$\dots - S_{tb} C_{t\bar{b}} - C_{tb} S_{t\bar{b}} S_{t\bar{b}} e^{i\gamma} \dots$

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$\dots - S_{tb} C_{t\bar{b}} - C_{tb} S_{t\bar{b}} S_{t\bar{b}} e^{i\gamma} \dots$

$$\mathcal{L}_{EW} = i\bar{\psi}_L \gamma_\mu D^\mu \psi_L + V(\phi) + Y_{ij} (\bar{\psi}_L^i \phi) \psi_R^j$$



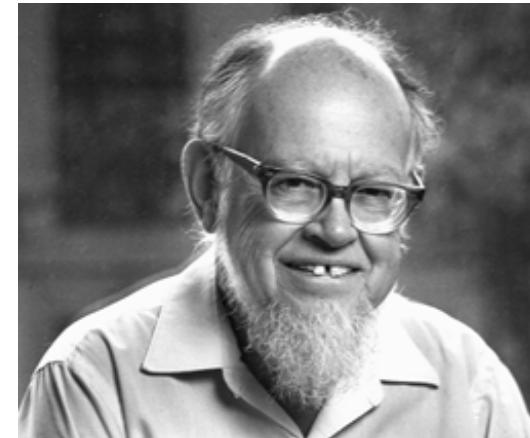
Mass vs Weak Eigenstates

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

- The Higgs interaction $Y_{ij} (\bar{\psi}_L \phi) \psi_R$ and the W interaction $g \bar{\psi}_L (\gamma^\mu W_\mu) \psi_L$ do not agree on the quark eigenstates.
 - Relation is the CKM matrix V_{CKM} ,
 - Leads to three mixing angles and one free imaginary CP phase in the CKM.
- Quark CP violation emerged together with mass.

The CKM Matrix V_{CKM}

$$\begin{matrix}
 & d & s & b \\
 u & V_{ud} & V_{us} & V_{ub} \\
 c & V_{cd} & V_{cs} & V_{cb} \\
 t & V_{td} & V_{ts} & V_{tb}
 \end{matrix}$$



Wolfenstein parametrization: V_{CKM}

$$\begin{pmatrix}
 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\
 -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\
 A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1
 \end{pmatrix}$$

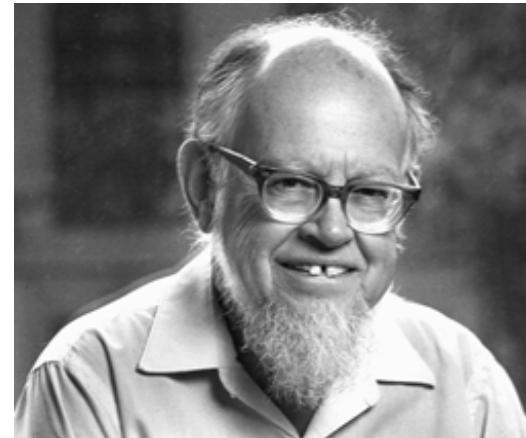
From unitarity ($V_{\text{CKM}} V_{\text{CKM}}^\dagger = 1$)
CKM has four free parameters:

3 real: λ (≈ 0.22), A (≈ 1), ρ
1 imaginary: $i\eta$

Particle \rightarrow Antiparticle: $V_{ij} \rightarrow V_{ij}^*$
 \Rightarrow 1 CP Violating phase

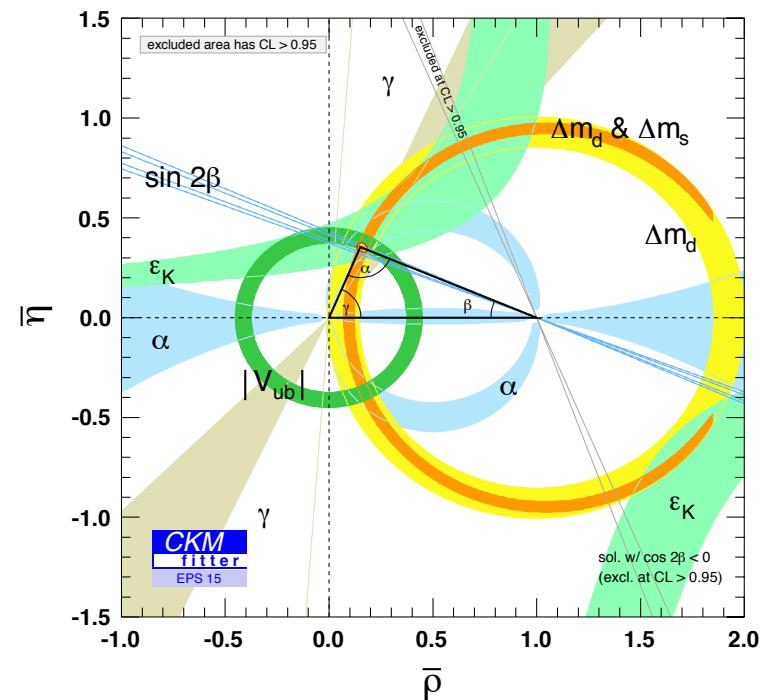
The CKM Matrix V_{CKM}

$$\begin{array}{c}
 \text{d} & \text{s} & \text{b} \\
 \text{u} & \boxed{V_{ud}} & V_{us} & V_{ub} \\
 \text{c} & V_{cd} & \boxed{V_{cs}} & V_{cb} \\
 \text{t} & V_{td} & V_{ts} & \boxed{V_{tb}}
 \end{array}$$



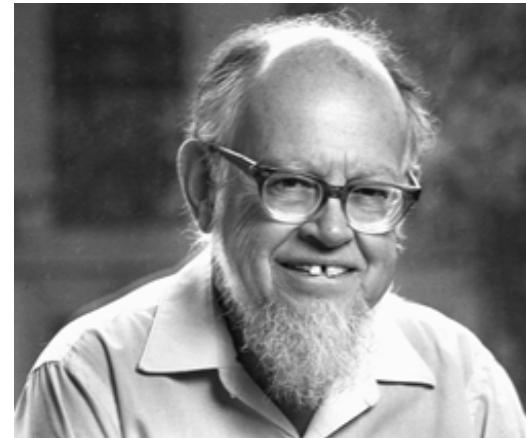
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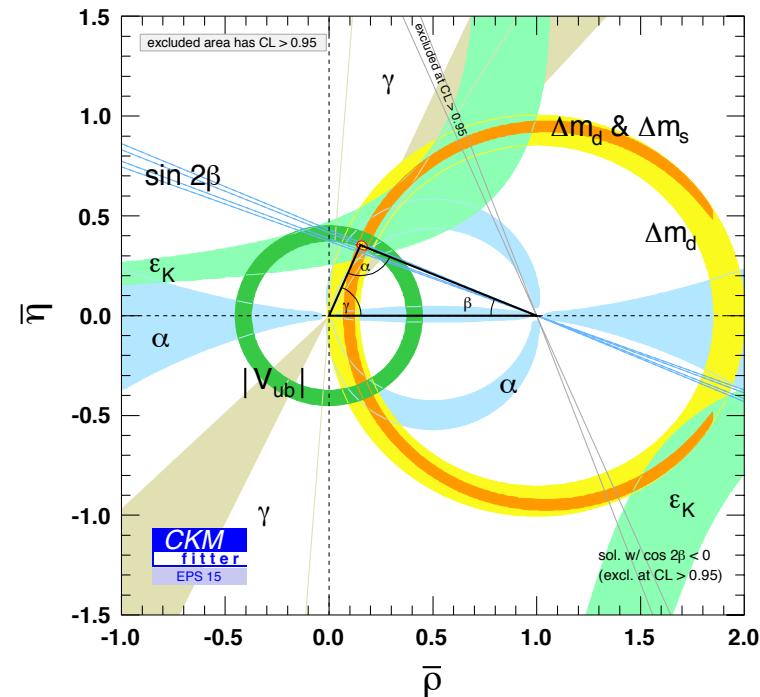
The CKM Matrix V_{CKM}

$$\begin{array}{c}
 \text{d} & \text{s} & \text{b} \\
 \text{u} & \left(\begin{array}{ccc}
 V_{ud} & V_{us} & V_{ub} \\
 V_{cd} & V_{cs} & V_{cb} \\
 V_{td} & V_{ts} & V_{tb}
 \end{array} \right) \\
 \text{c} \\
 \text{t}
 \end{array}$$



Wolfenstein parametrization: V_{CKM}

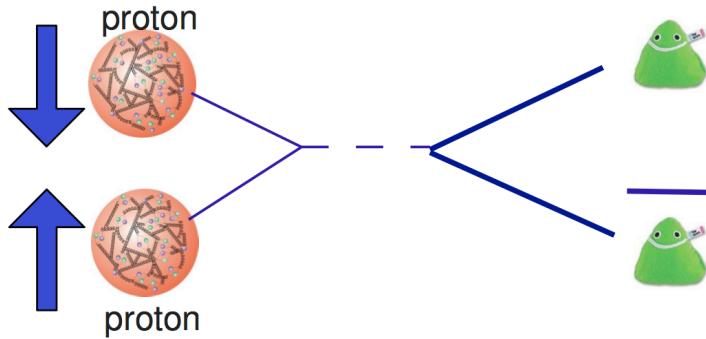
$$\left(\begin{array}{ccc}
 |V_{ud}| & |V_{us}| & |V_{ub}| e^{-i\gamma} \\
 -|V_{cd}| & |V_{cs}| & |V_{cb}| \\
 |V_{td}| e^{-i\beta} & -|V_{ts}| e^{i\beta_s} & |V_{tb}|
 \end{array} \right)$$



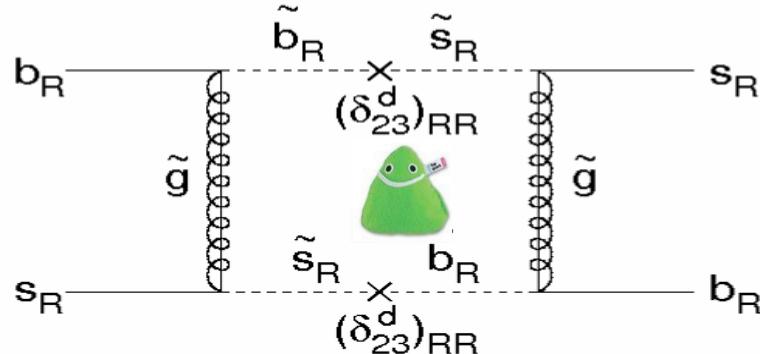
Can flavour physics explain the matter - antimatter asymmetry? → Requires New Physics!

Direct vs Indirect search

The **absolute** energy frontier:



The **virtual** energy frontier:



- Direct observation:
 - Produce particles on-shell and detect decay products
 - More intuitive(?), “really” produced
 - Limited by collision energy

- Indirect observation:
 - Less intuitive(?), quantum level
 - Limited by precision, not by collision energy
 - CP observables sensitive to imaginary couplings

- Indirect observations in the past:
 - Kaon decay $K^0 \rightarrow \mu\mu$ hints at c-quark via GIM in 1970 (J/ ψ produced in 1974)
 - 3rd quark family predicted 1972 to explain CP violation (b produced in 1977, t in 1994)
 - Neutral current discovered in neutrino experiment in 1973 (Z-boson produced in 1983)
 - $B\bar{B}$ mixing (1987) hints at large top mass, LEP predicts top mass (1990) (top in 1994)

Indirect evidence for top mass and charm quark

$K^0 \rightarrow \mu\mu$ pointed to the **charm** quark

GIM, Phys. Rev. D2, 1285, 1970

Weak Interactions with Lepton-Hadron Symmetry*

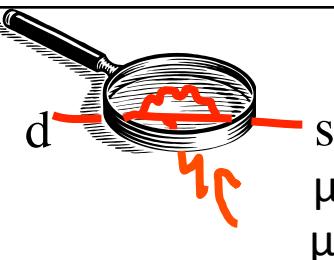
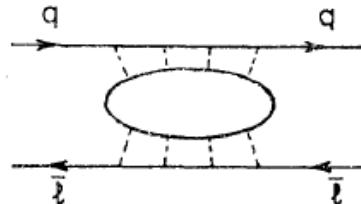
S. L. GLASHOW, J. ILIOPoulos, AND L. MAIANI†
Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02139
(Received 5 March 1970)

We propose a model of weak interactions in which the currents are constructed out of four basic quark fields and interact with a charged massive vector boson. We show, to all orders in perturbation theory, that the leading divergences do not violate any strong-interaction symmetry and the next to the leading divergences respect all observed weak-interaction selection rules. The model features a remarkable symmetry between leptons and quarks. The extension of our model to a complete Yang-Mills theory is discussed.

splitting, beginning at order $G(G\Lambda^2)$, as well as contributions to such unobserved decay modes as $K_2 \rightarrow \mu^+ + \mu^-$, $K^+ \rightarrow \pi^+ + l + \bar{l}$, etc., involving neutral lepton ...

We wish to propose a simple model in which the divergences are properly ordered. Our model is founded in a quark model, but one involving four, not three, fundamental fermions; the weak interactions are mediators ...

new quantum number C for charm.



B^0 mixing pointed to the **top** quark mass

ARGUS Coll, Phys. Lett. B192:245, 1987

DESY 87-029

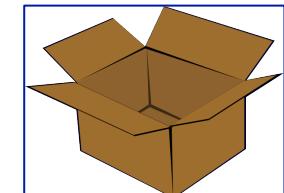
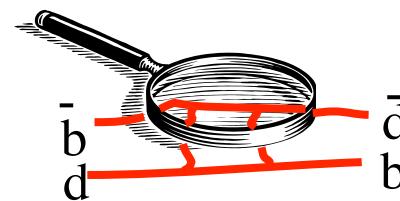
April 1987

OBSERVATION OF $B^0 - \bar{B}^0$ MIXING

The ARGUS Collaboration

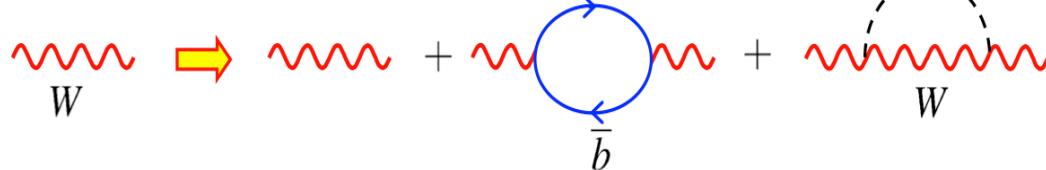
In summary, the combined evidence of the investigation of B^0 meson pairs, lepton pairs and B^0 meson-lepton events on the $\Upsilon(4S)$ leads to the conclusion that $B^0 - \bar{B}^0$ mixing has been observed and is substantial.

Parameters	Comments
$r > 0.09$ 90% CL	This experiment
$x > 0.44$	This experiment
$B^{\frac{1}{2}} f_B \approx f_\pi < 160$ MeV	B meson (\approx pion) decay constant
$m_b < 5\text{GeV}/c^2$	b -quark mass
$\tau_b < 1.4 \cdot 10^{-12}$ s	B meson lifetime
$ V_{tb} < 0.018$	Kobayashi-Maskawa matrix element
$ m_{cb} < 0.86$	QCD correction factor [17]
$m_t > 50\text{GeV}/c^2$	t quark mass



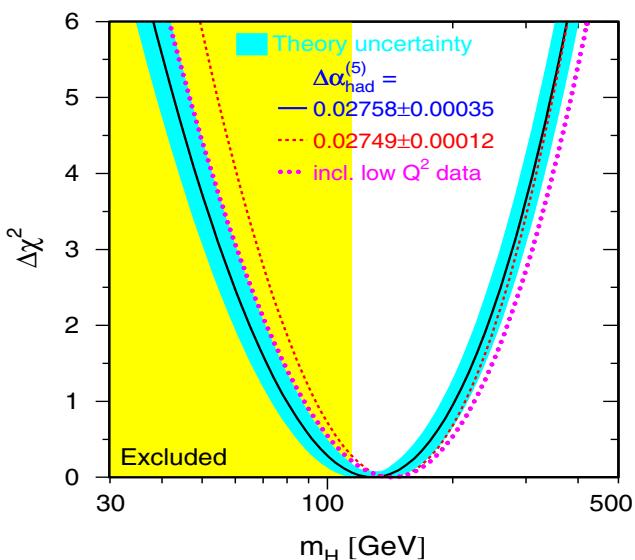
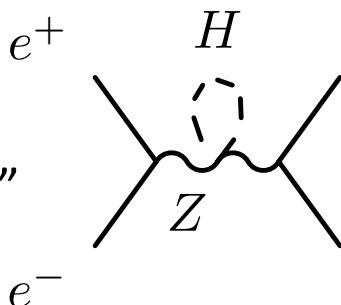
Higgs mass and Electroweak precision measurements

W propagator (Z similar):



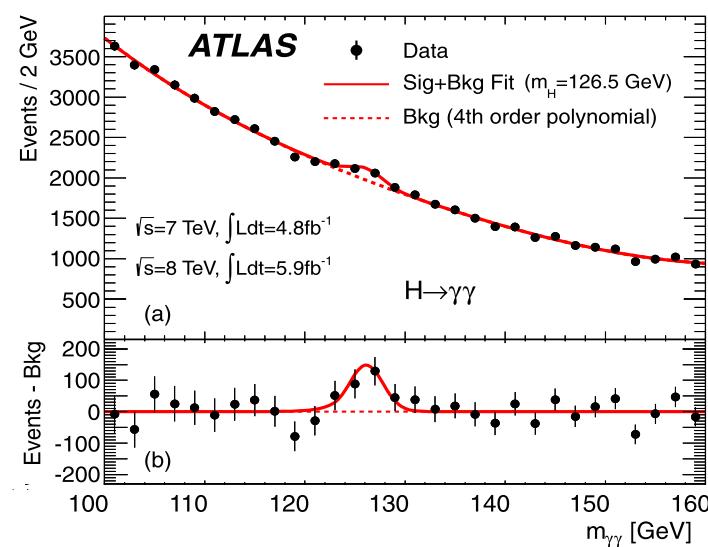
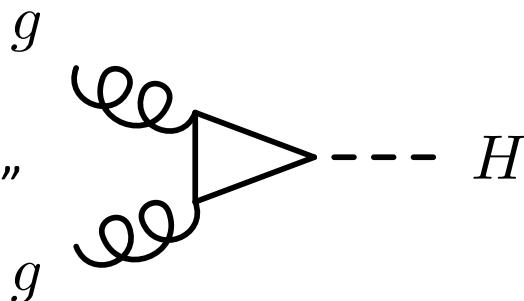
$$M_W = M_W^0 + a m_t^2 + b \ln\left(\frac{M_H}{M_W}\right)$$

- LEP
“Indirect”



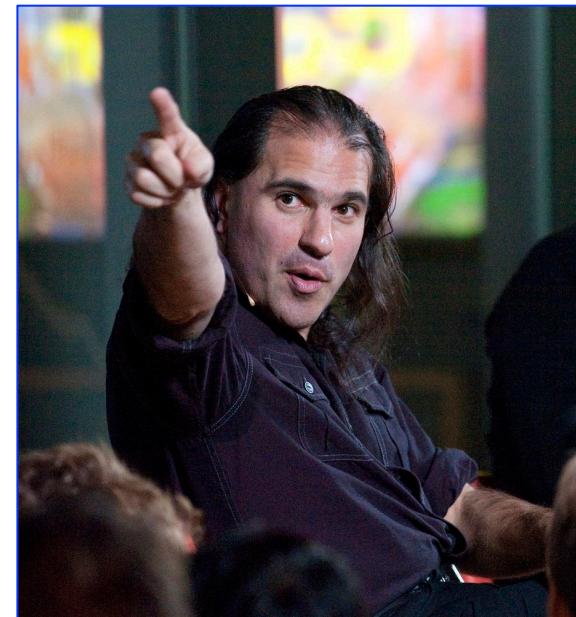
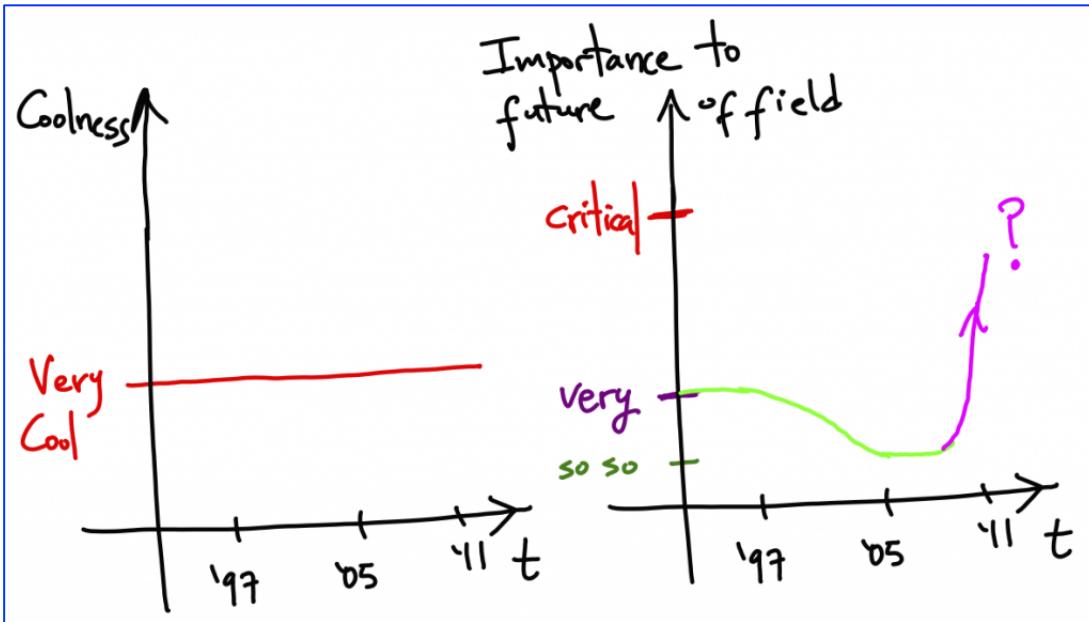
LEP: Phys.Rep.427 (2006) 257

- LHC
“Direct”



Atlas coll. Phys.Lett.B716 (2012) 1

Why work on Flavor Physics?

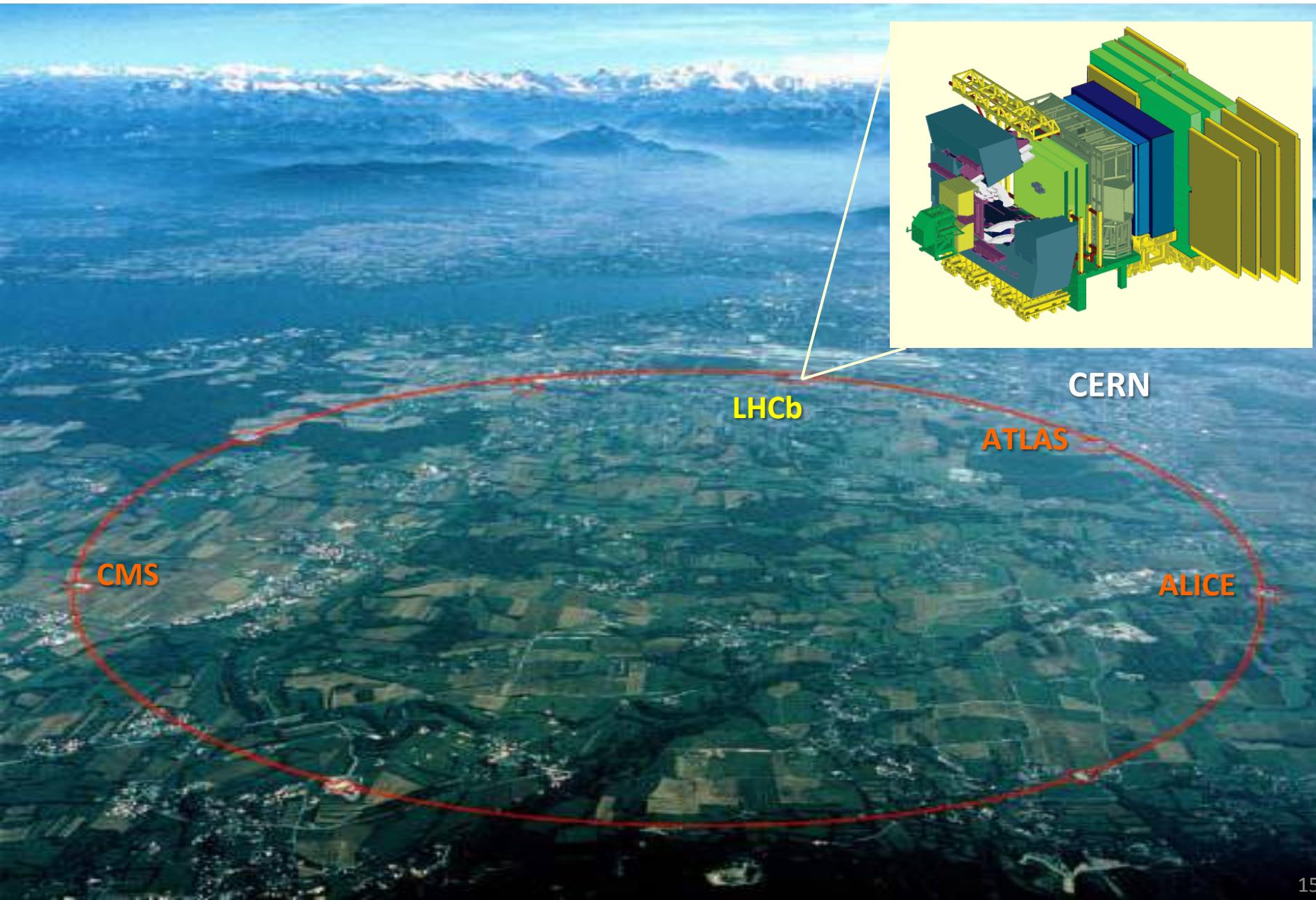


(Nima Arkani-Hamed likes it too)

Outline

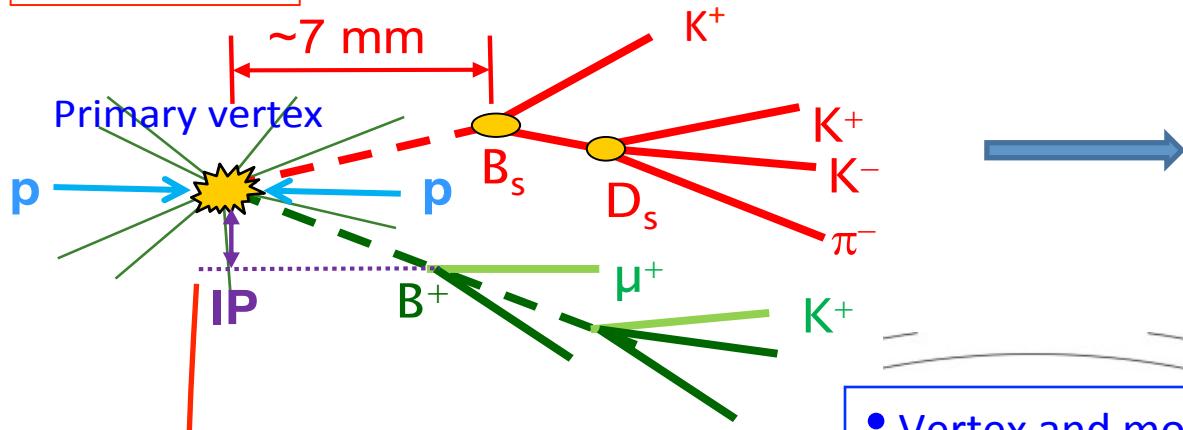
- 1. Flavour Physics
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- 5. Outlook

LHCb @ LHC



B Physics @ LHCb

$B_s \rightarrow D_s^- K^+$

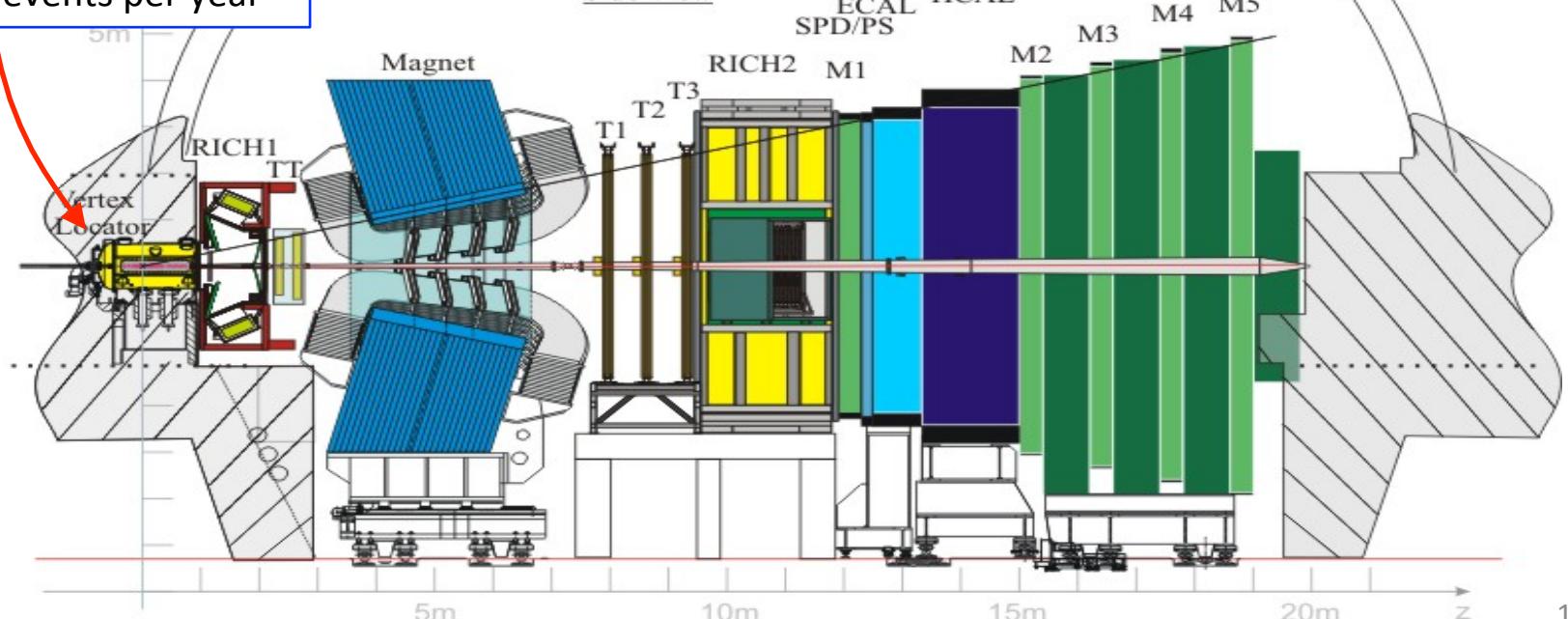


- Background Suppression
- Flavour tagging
- Decay time measurement

Design:

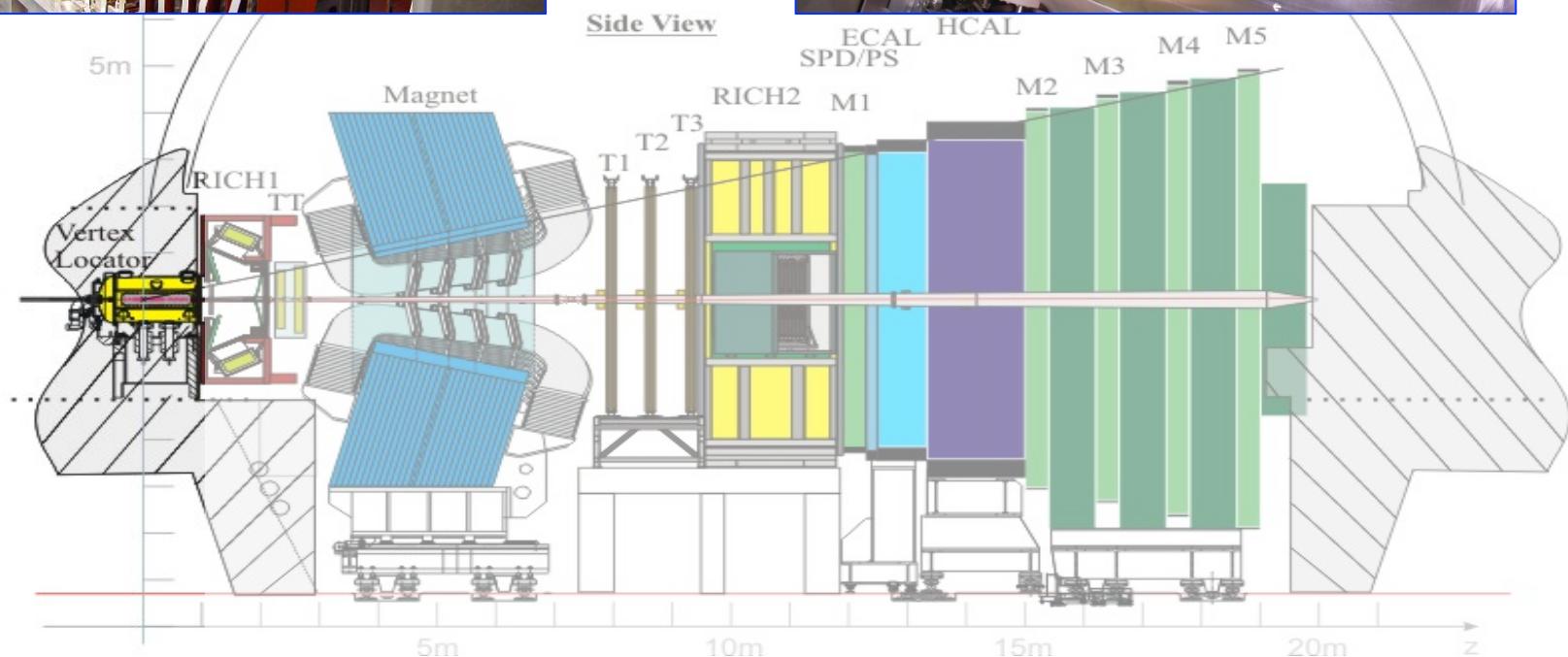
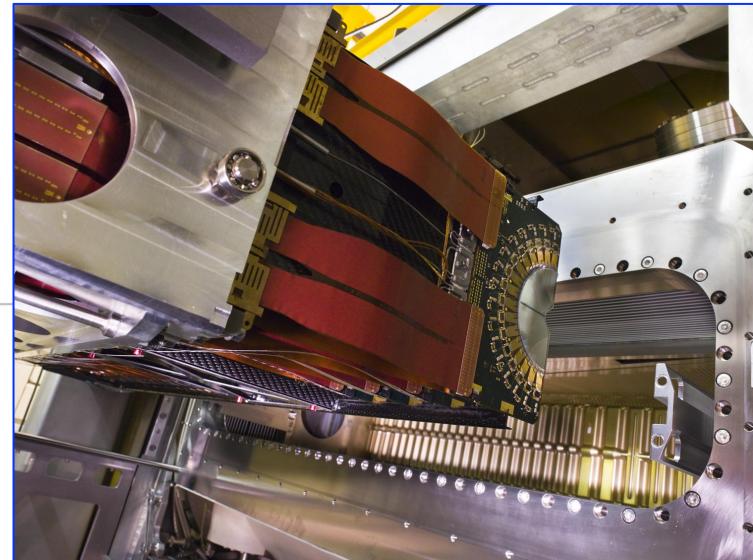
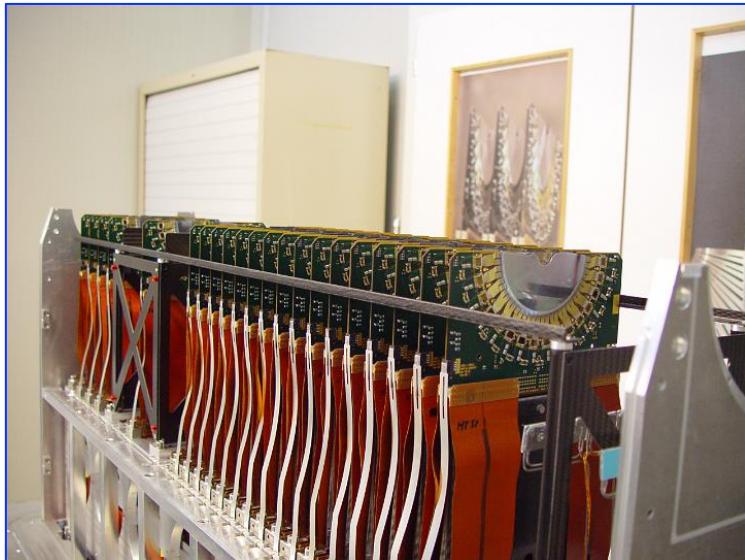
- LHCb event rate: $\sim 30 \text{ MHz}$
- 1 in 160 is a b-b event
- $10^{12} \text{ b-b events per year}$

- Vertex and momenta reconstruction
- Particle identification (π, K, μ, e, γ)
- Trigger



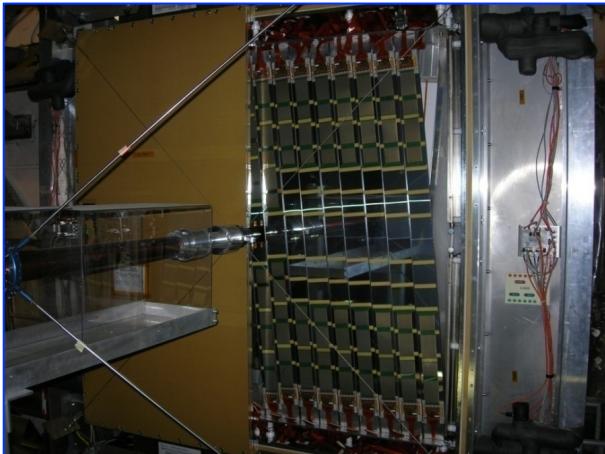
Vertex Topology

Vertex Locator

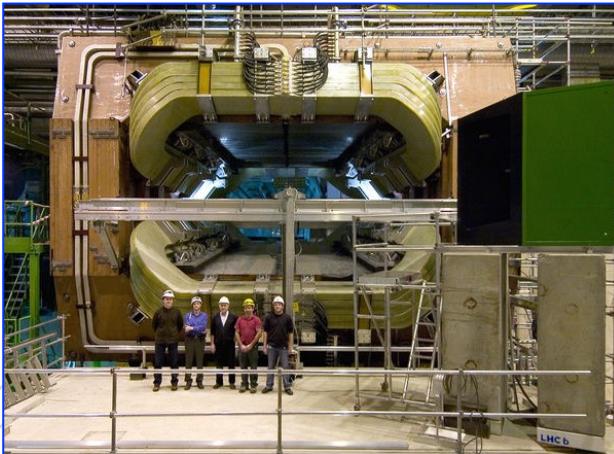


Momentum and Mass Reconstruction

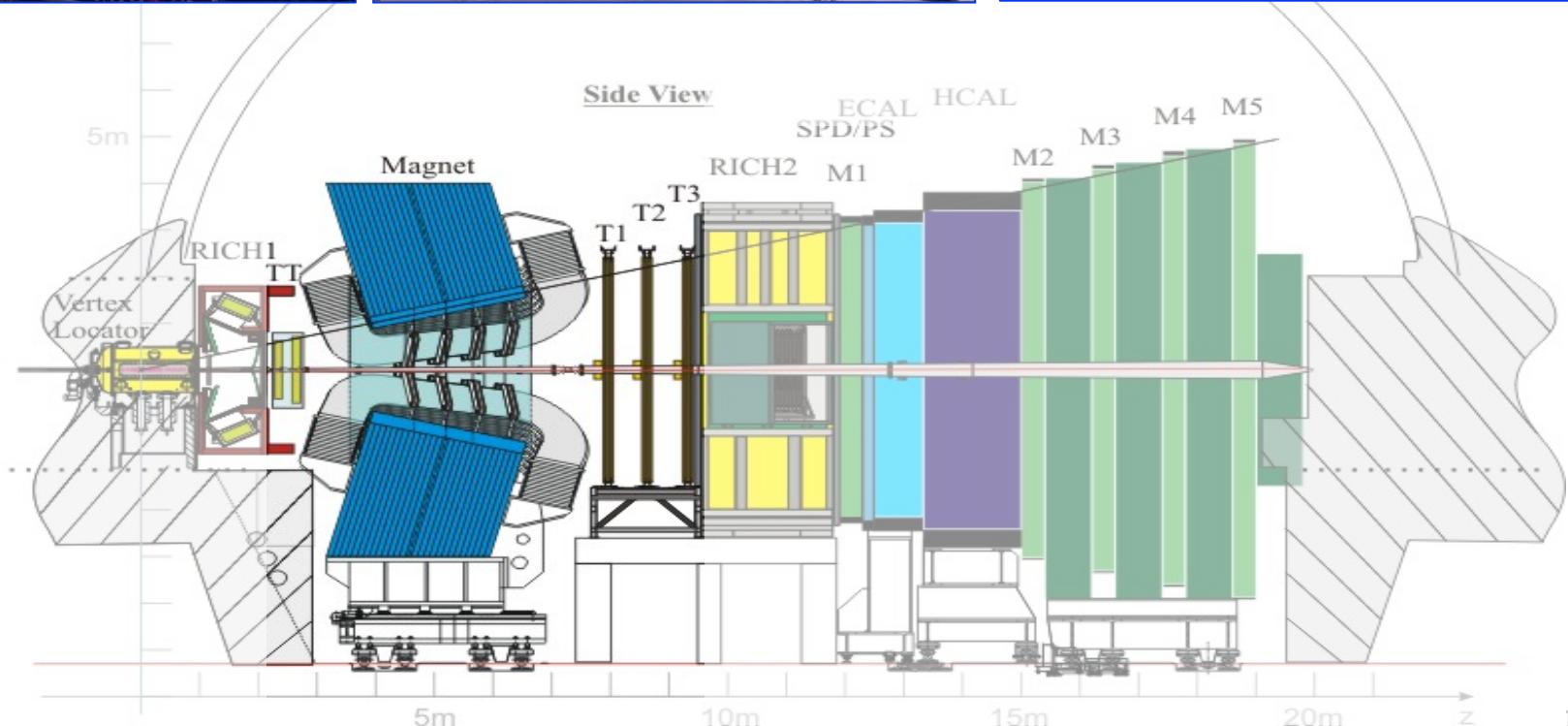
Trigger Tracker



Magnet



Outer Tracker

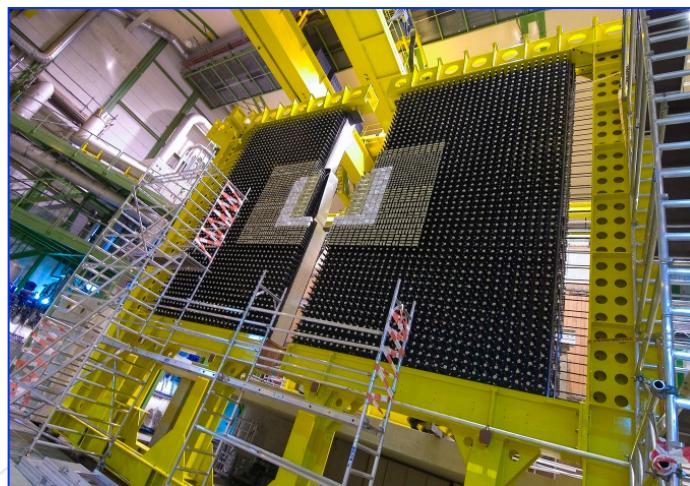


Particle Identification: π , K , μ , γ , e

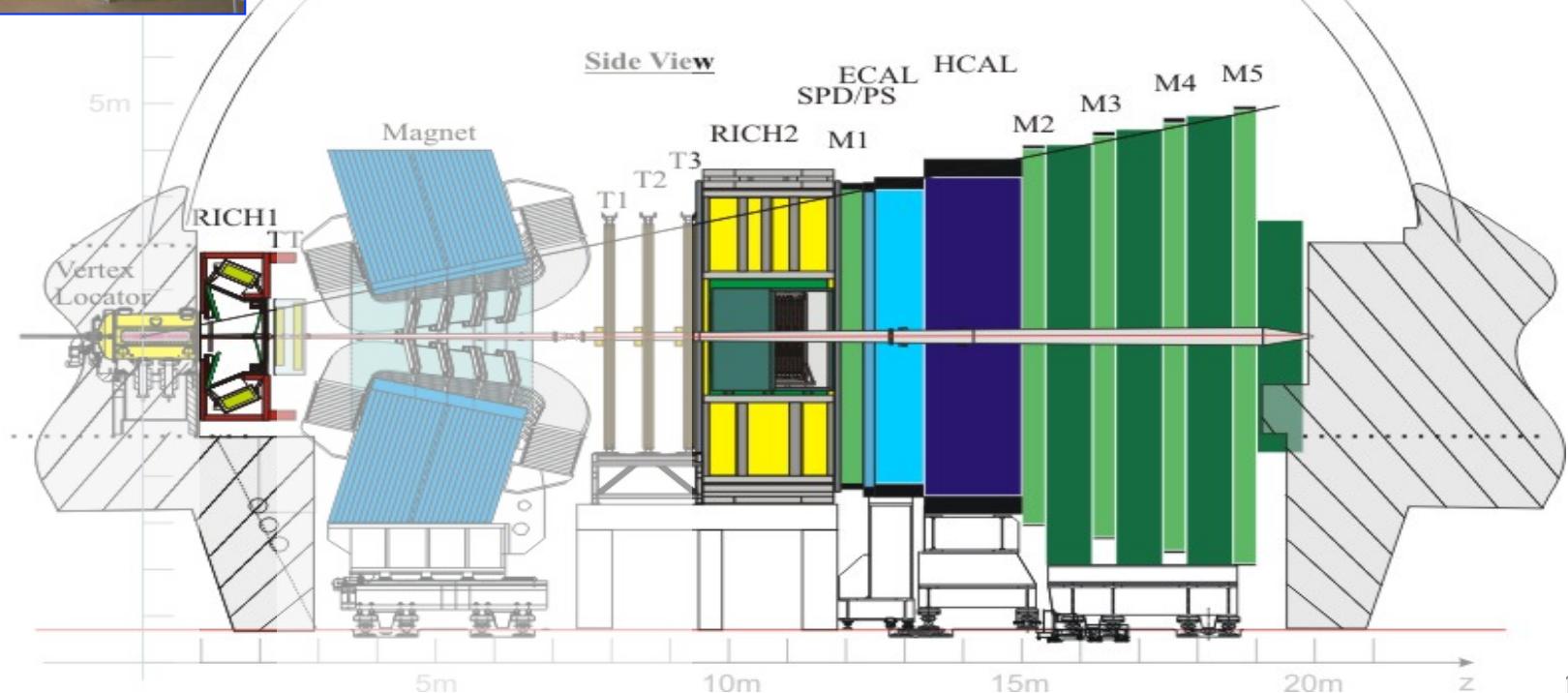
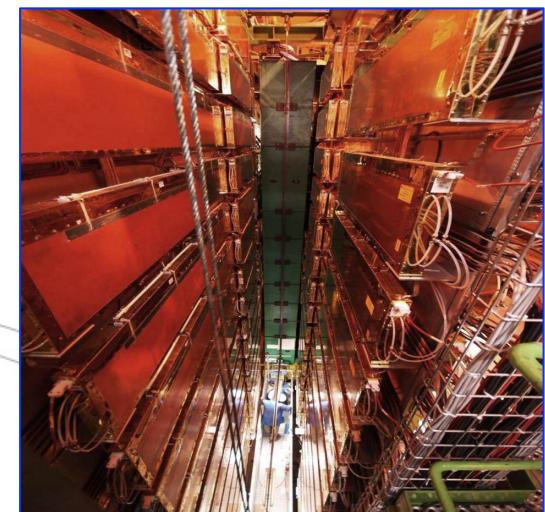
RICH



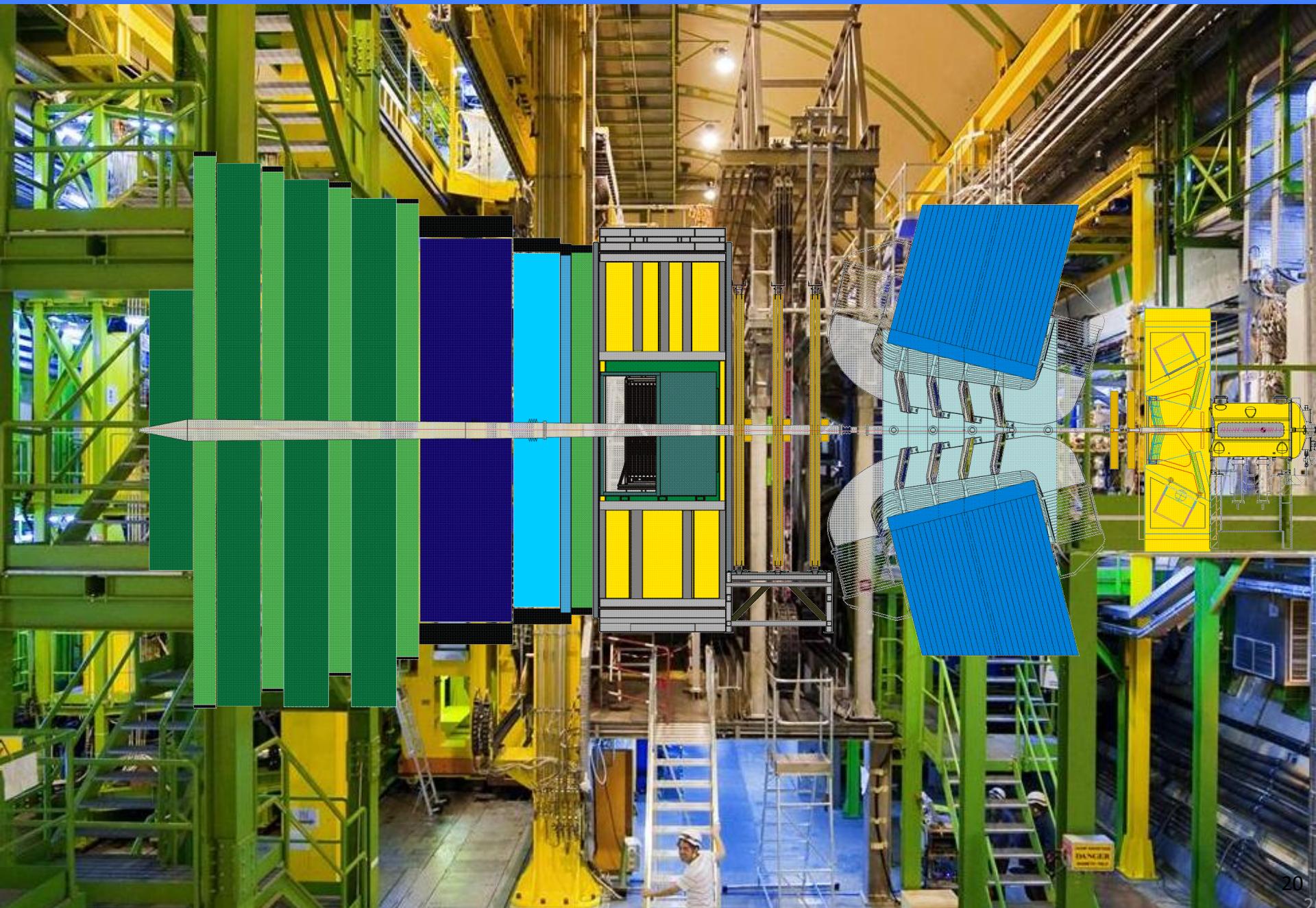
Calorimeter



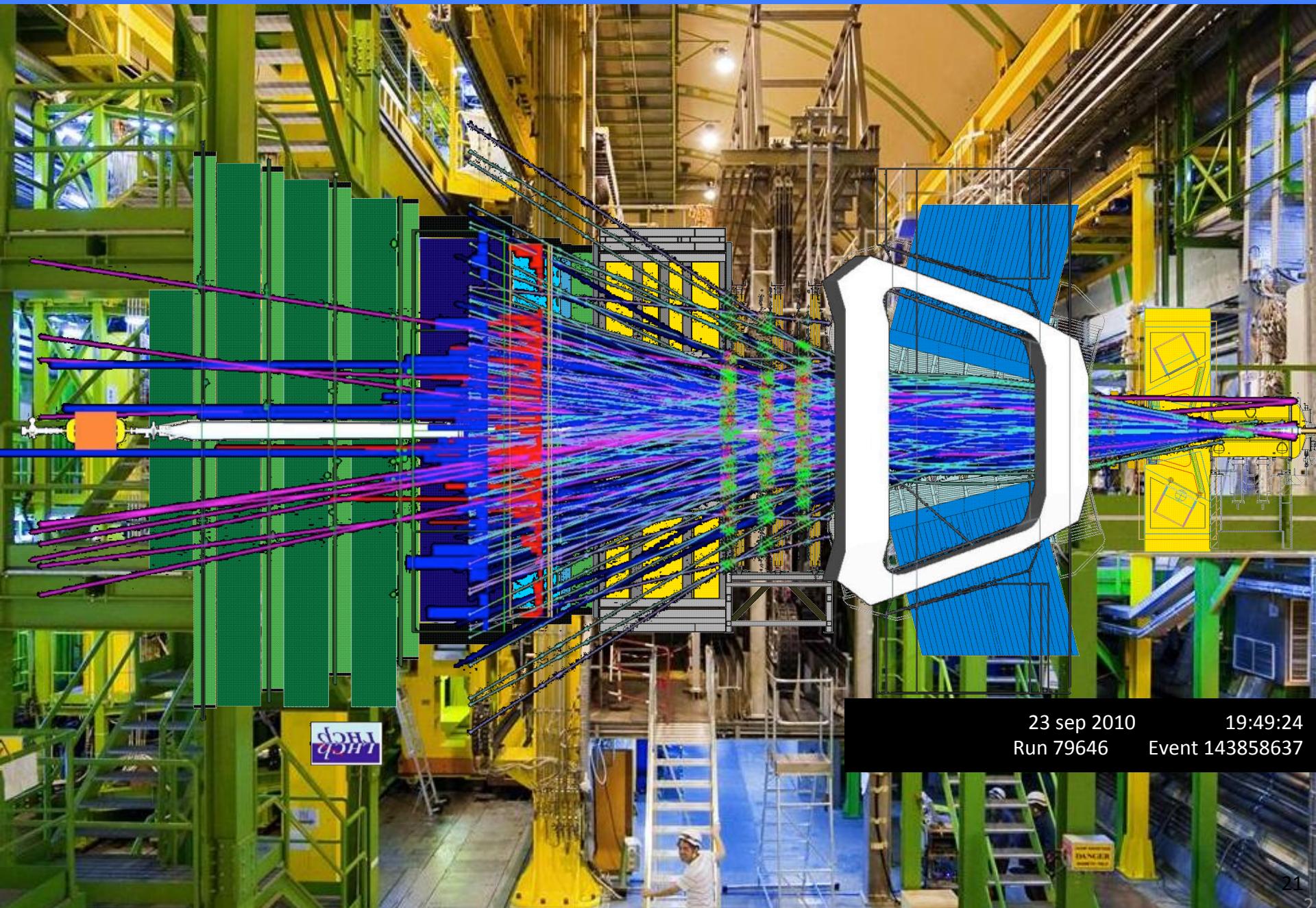
Muon



The LHCb Detector



The LHCb Detector



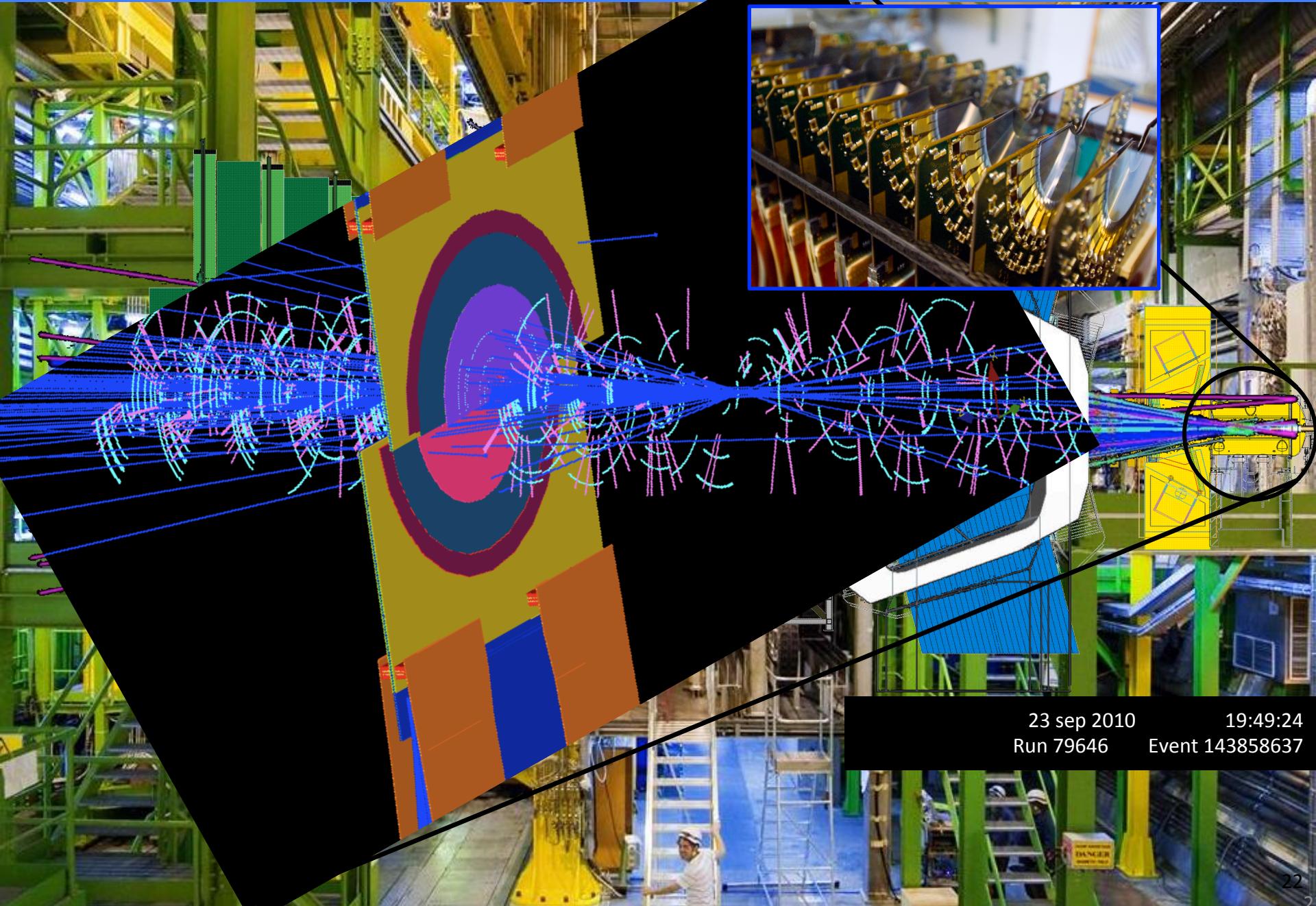
23 sep 2010

Run 79646

19:49:24

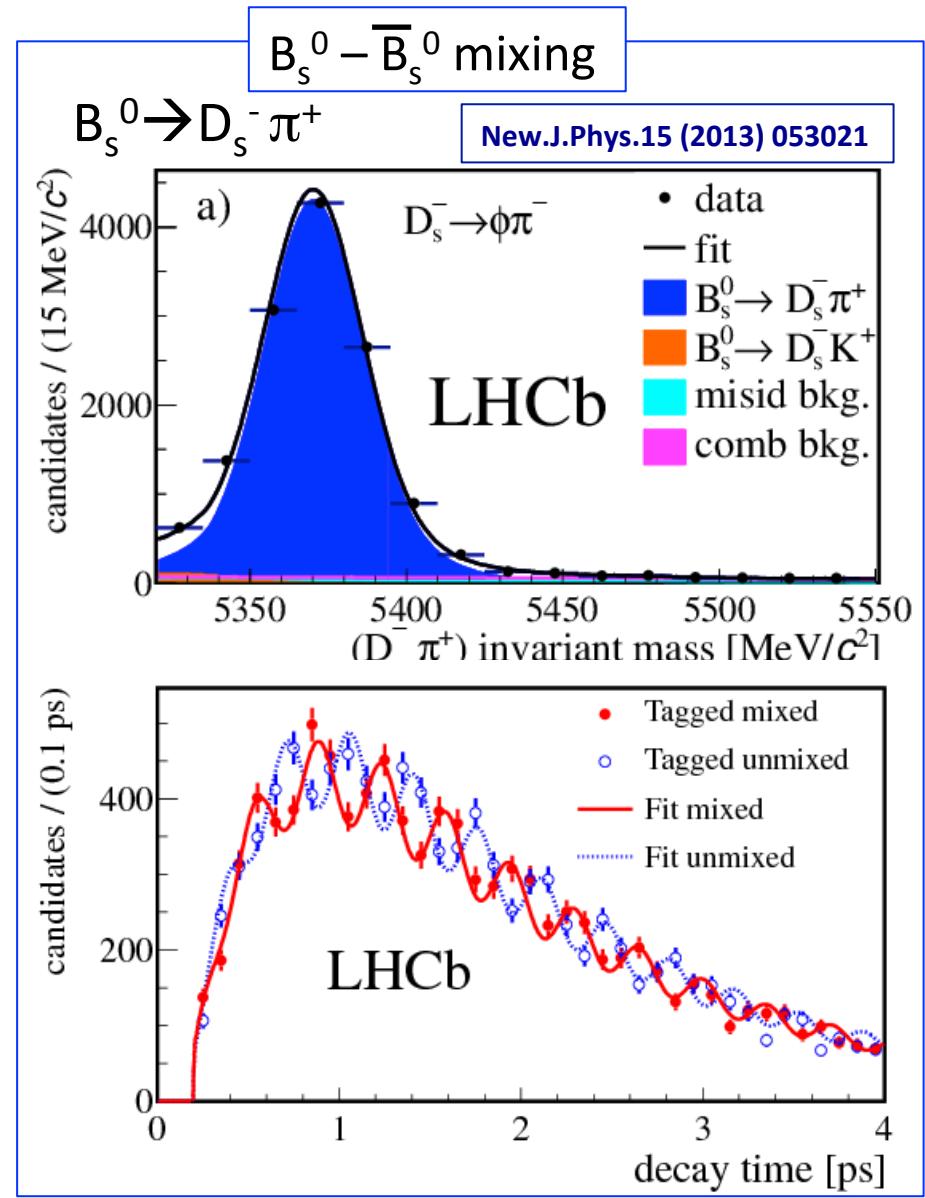
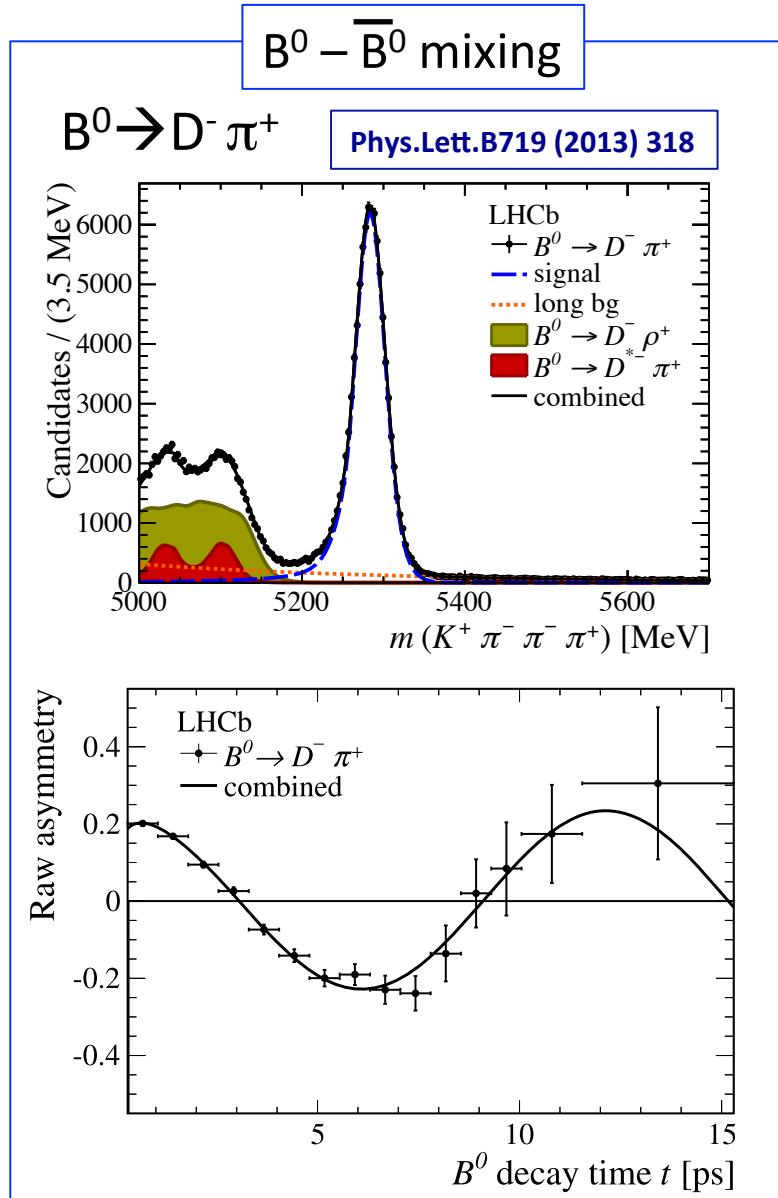
Event 143858637

The LHCb Detector



LHCb Performance: $B_{d/s}$ Mixing

Trigger, background suppression, decay time reconstruction, flavour tagging: ***it works!***



Outline

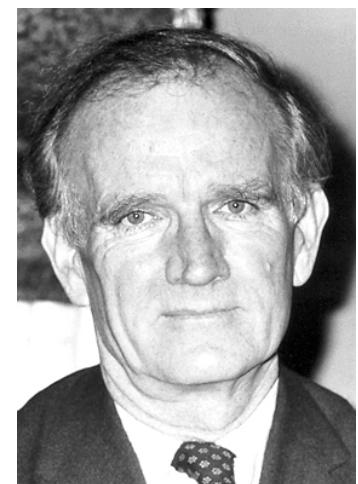
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A slide of History on CP Violation

- Charge Parity Violation in particle physics:
 - a) **1964** (CCFT): Discovered in neutral Kaon beam "indirect CP violation"
 - Also called: **CPV in mixing**
 - $\text{Prob}(\text{K}^0 \rightarrow \bar{\text{K}}^0) \neq \text{Prob}(\bar{\text{K}}^0 \rightarrow \text{K}^0)$
 - $|\varepsilon| = (2.228 \pm 0.011) \times 10^{-3}$ (PDG 2014)
 - b) **1999** (NA48 & KTeV): Seen in Kaon decays "direct CP violation"
 - Also called: **CPV in decay**
 - Decay rates $\Gamma(\text{K}^0 \rightarrow \pi^+ \pi^-) \neq \Gamma(\bar{\text{K}}^0 \rightarrow \pi^+ \pi^-)$
 - $\text{Re}(\varepsilon'/\varepsilon) = (1.65 \pm 0.26) \times 10^{-3}$ (PDG 2014)
 - c) **2001** (Belle & Babar): Observed in B mesons decays "CP violation in interference"
 - Also called: **mixing induced CPV**
 - $\sin 2\beta = 0.682 \pm 0.019$ (PDG 2014)



James Watson Cronin



Val Logsdon Fitch

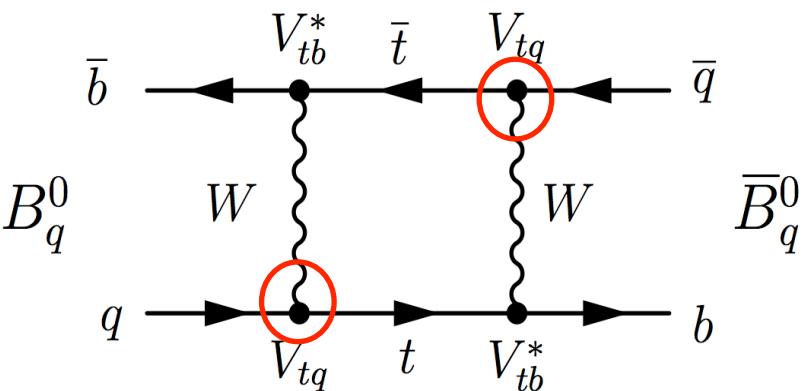
→ Nobel prize 1980

Observing CP Violation in B decays

- CP Violation occurs in *interference of two quantum amplitudes* that have a different CP odd (as well as CP even) phase.
 - Use the phenomenon of B - \bar{B} mixing:
 - Interference of direct decay and decay via mixing leads to CP violation.

Standard Model:

- B^0 - \bar{B}^0 phase $\Phi_d = \arg(V_{td}) = 2\beta$
- B_s - \bar{B}_s phase $\Phi_s = \arg(V_{ts}) = -2\beta_s$

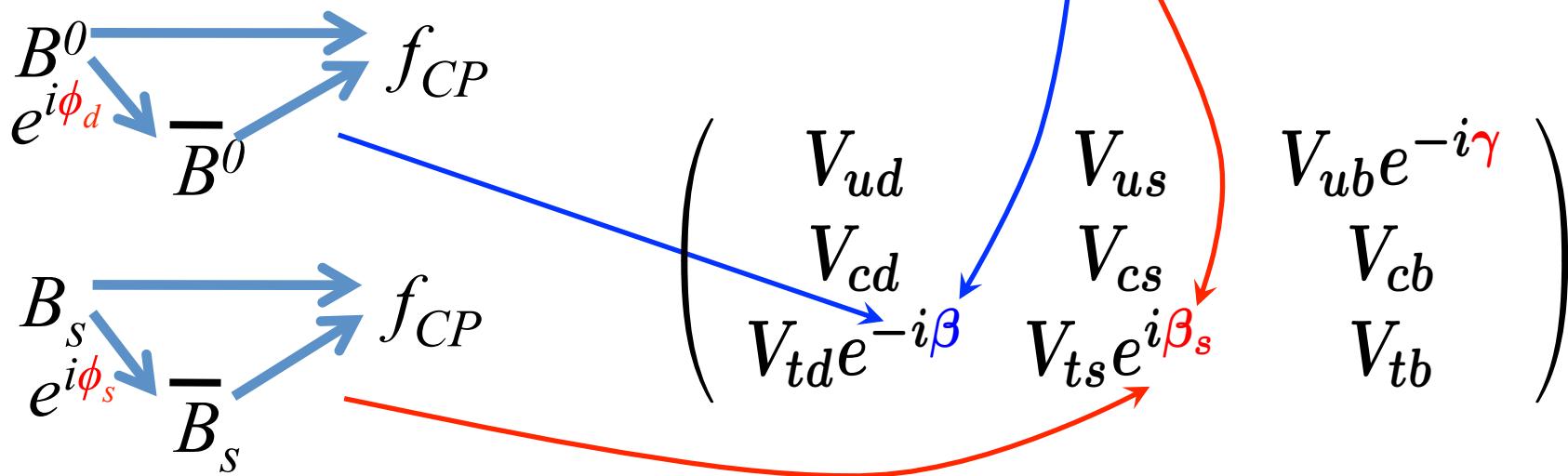
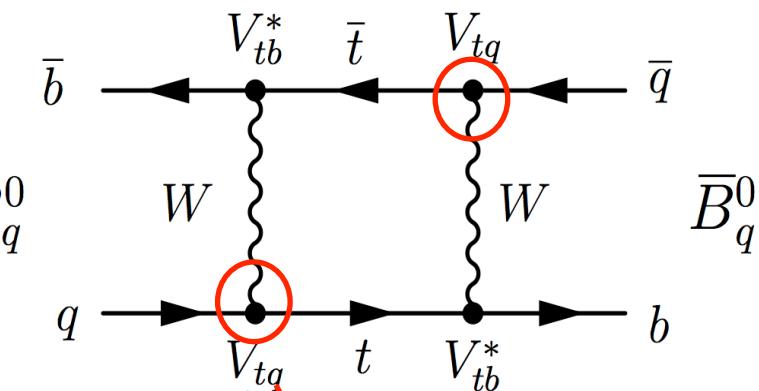


Observing CP Violation in B decays

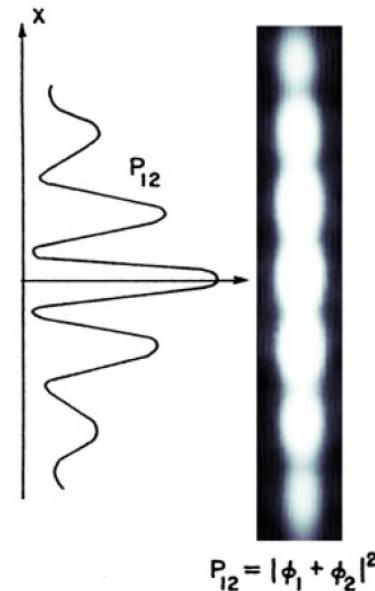
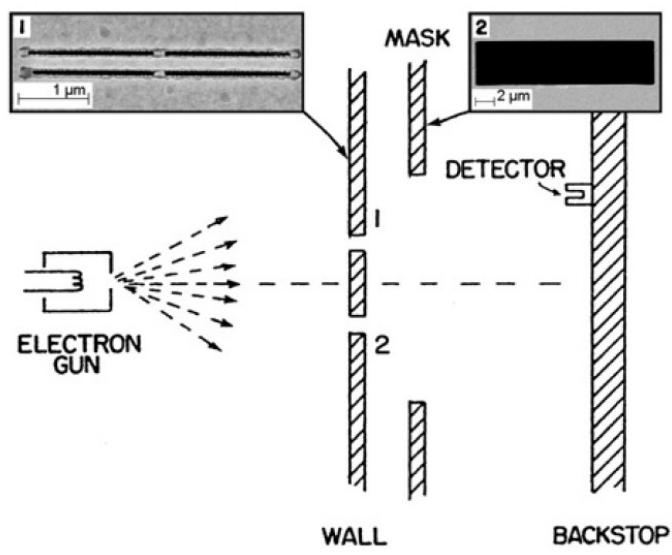
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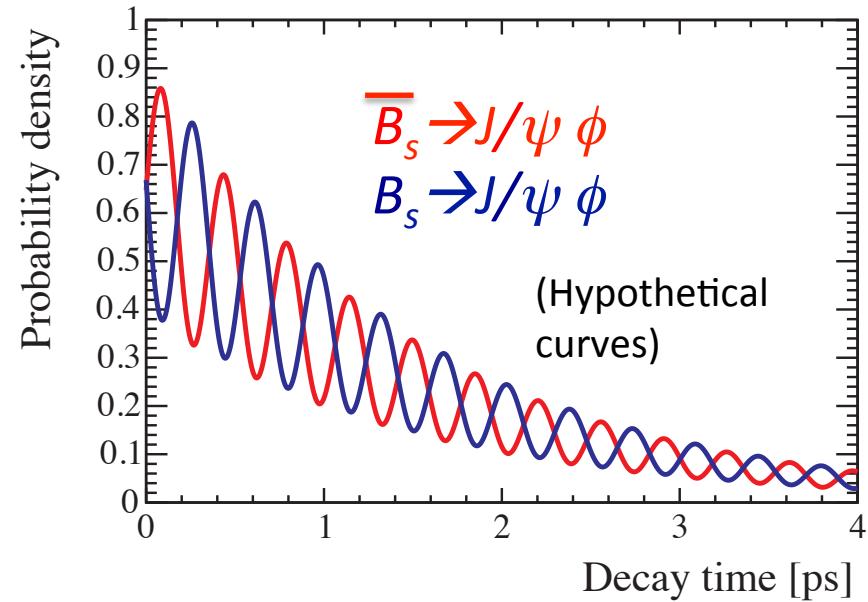
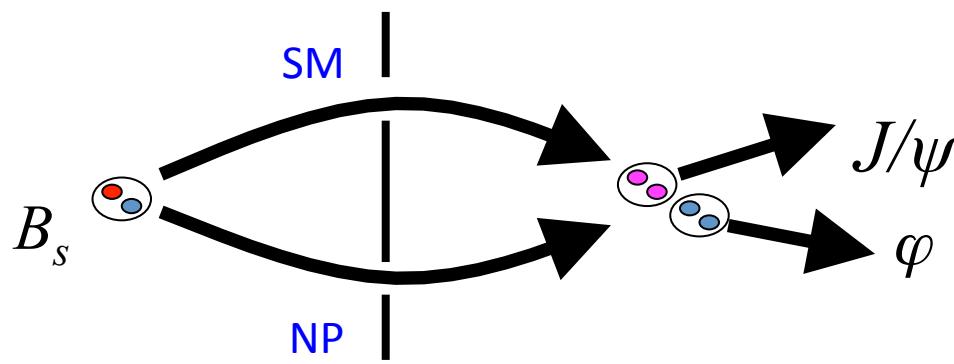
- B^0 - \bar{B}^0 phase $\phi_d = \arg(V_{td}) = 2\beta$
- B_s - \bar{B}_s phase $\phi_s = \arg(V_{ts}) = -2\beta_s$



Feynman: “Remember the double slit experiment”



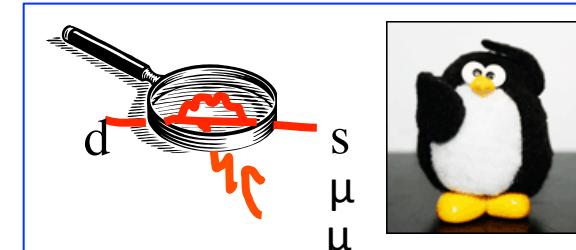
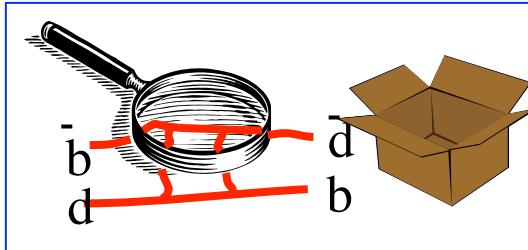
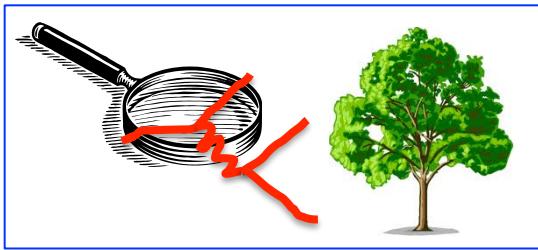
LHCb is a completely analogous interference experiment using B-mesons.



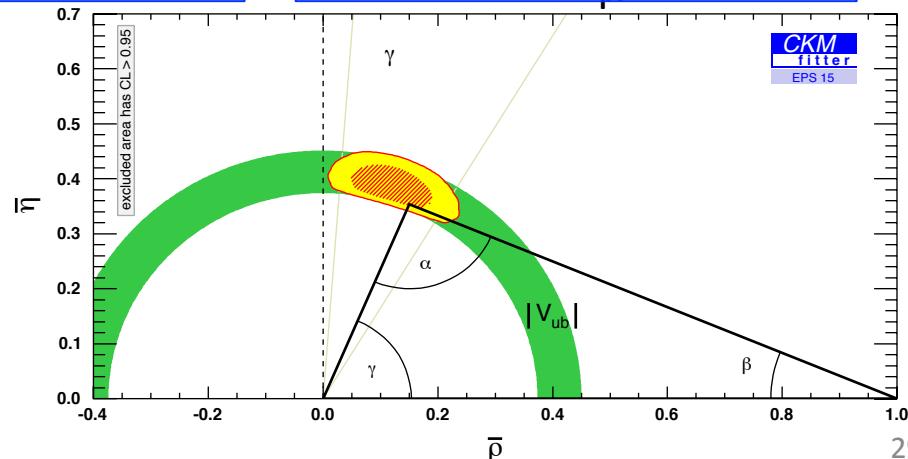
Look for interference of SM with NP

Probing new CP-violation with B mesons

- Look for *interference* of CKM with BSM amplitudes:
 - SM Prediction: No CP violation a la kaons (“*indirect*” or “*mixing*”)
 - $K \leftrightarrow \bar{K} = \text{epsilon} \neq 0$ but $B \leftrightarrow \bar{B} = 0$
 - Precise measurement of angle γ “*with trees*” vs γ “*with loops*” (“*direct*”)
 - Many decay modes with tree and loop diagrams are available
 - SM Prediction: B_s box diagram phase $\phi_s = -2\beta_s \approx 0$ (“*mixing induced*”)
 - New Physics “in the box”?



$$\left(\begin{array}{ccc} V_{ud} & V_{us} & V_{ub} e^{-i\gamma} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} e^{-i\beta} & V_{ts} e^{i\beta_s} & V_{tb} \end{array} \right)$$



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5. Outlook

CP Violation in Mixing

- Interference of the *dispersive* and *absorptive* amplitudes

– Similar to measurement ϵ in kaon physics

- Hamiltonian:

$$i \frac{d}{dt} \begin{pmatrix} B_s^0 \\ \bar{B}_s^0 \end{pmatrix} = \begin{pmatrix} M_{11} - \frac{i}{2}\Gamma_{11} & M_{12} - \frac{i}{2}\Gamma_{12} \\ M_{12}^* - \frac{i}{2}\Gamma_{12}^* & M_{22} - \frac{i}{2}\Gamma_{22} \end{pmatrix} \begin{pmatrix} B_s^0 \\ \bar{B}_s^0 \end{pmatrix}$$

- Eigenstates:

$$M_L, M_H \quad ; \quad \Delta M = M_H - M_L$$

$$\Gamma_L, \Gamma_H \quad ; \quad \Delta\Gamma = \Gamma_H - \Gamma_L$$

$$\phi_{12} = \arg(-M_{12}/\Gamma_{12})$$

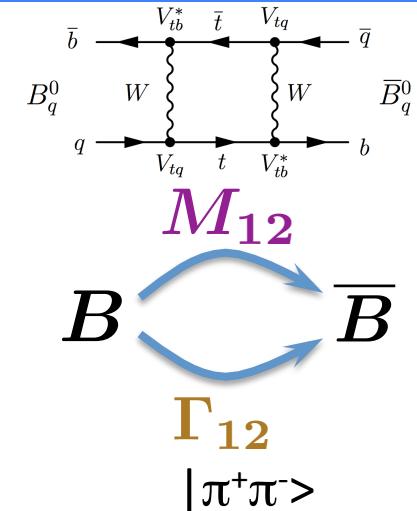
- Observable:

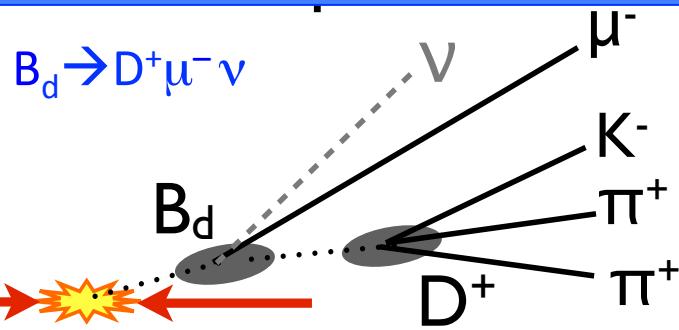
$$a_{sl} = \frac{\Gamma(\bar{B}(t) \rightarrow f) - \Gamma(B(t) \rightarrow \bar{f})}{\Gamma(\bar{B}(t) \rightarrow f) + \Gamma(B(t) \rightarrow \bar{f})} = \frac{\Delta\Gamma}{\Delta M} \tan \phi_{12}$$

- In the Standard Model $\phi_{12} \approx 0.2^\circ$ for B_s

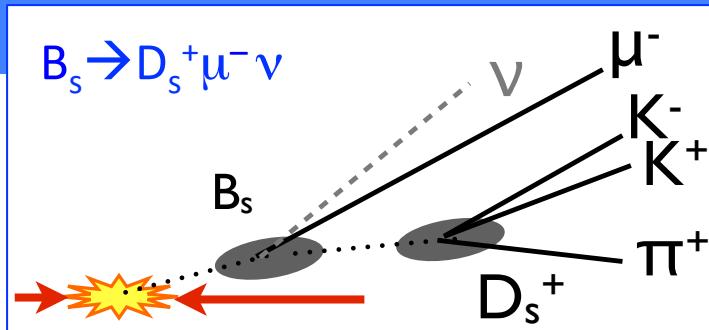
Lenz, Nierste JHEP.06 (2007) 072

- $B_{d,s}$: measure with flavor specific decays: $B_{d,s} \rightarrow D_{(s)}^+ \mu^- \nu$ and C.C.





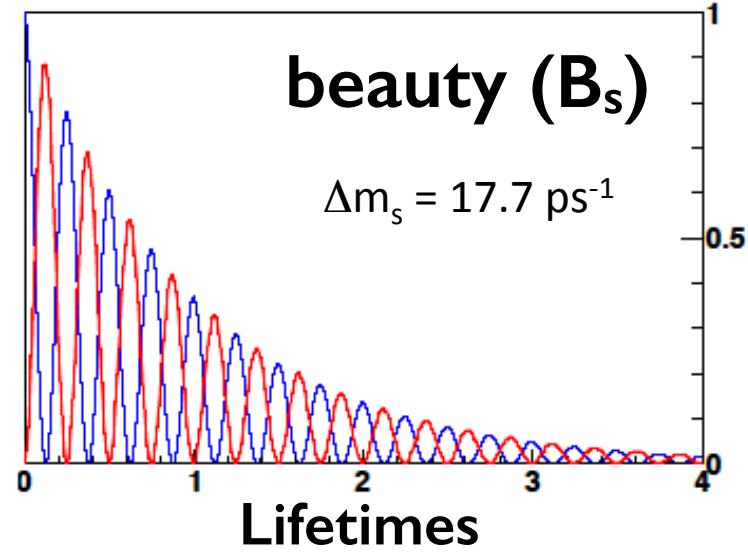
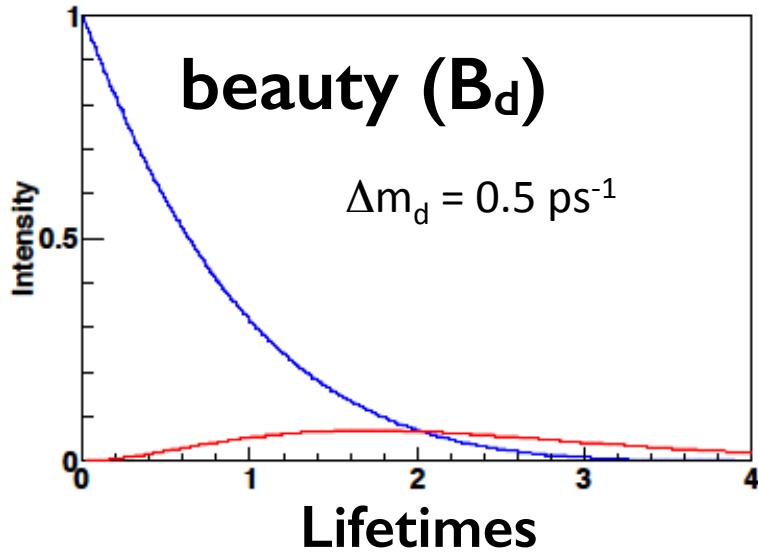
$B_d \text{ vs } B_s$



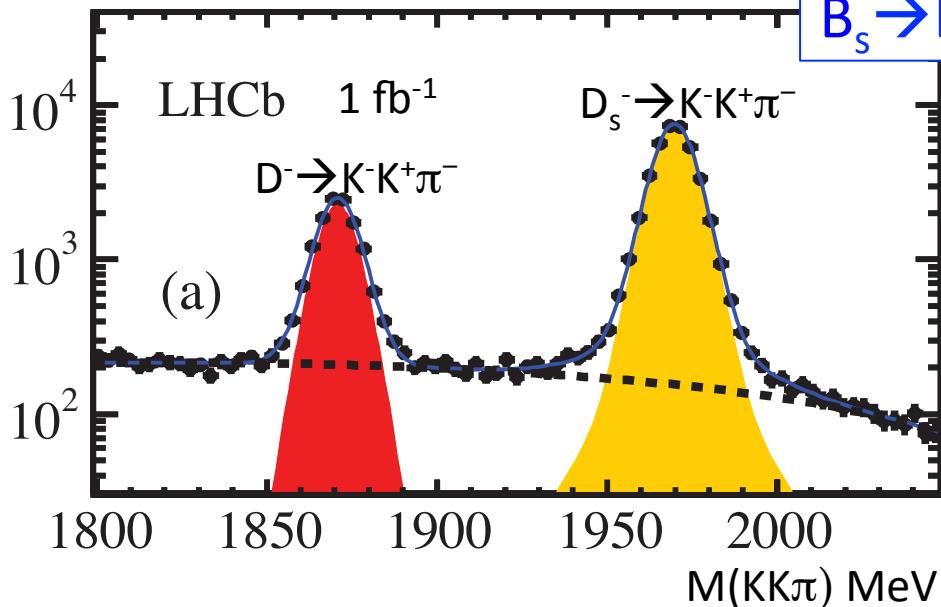
$$\mathcal{A}_{B\bar{B}} = \frac{N(D_{(s)}^-\mu^+, t) - N(D_{(s)}^+\mu^-, t)}{N(D_{(s)}^-\mu^+, t) + N(D_{(s)}^+\mu^-, t)} = \frac{a_{sl}}{2} + \left(a_p + \frac{a_{sl}}{2} \right) \frac{\cos \Delta M t}{\cosh \frac{\Delta \Gamma}{2} t}$$

a_{sl}^d (B_d) : must measure time-dependent

a_{sl}^s (B_s) : washout of fast oscillation factor and production asymmetry



Semileptonic charge asymmetry with B_s : a_{sl}^s



$B_s \rightarrow D_s^+ \mu^- \nu$

$$\mathcal{A}_{\text{meas}} = \frac{\Gamma [D_s^- \mu^+] - \Gamma [D_s^+ \mu^-]}{\Gamma [D_s^- \mu^+] + \Gamma [D_s^+ \mu^-]} = \frac{a_{sl}^s}{2}$$

Correct for detection asymmetry:

$$\frac{\epsilon(D_s^- \mu^+)}{\epsilon(D_s^+ \mu^-)}$$

LHCb:

PLB 728C (2014) 607

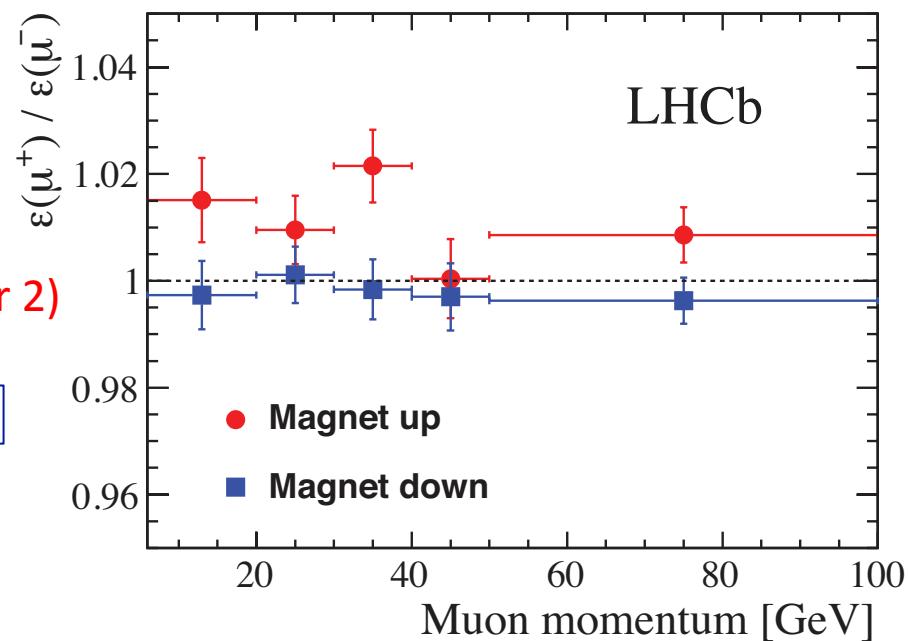
$$a_{sl}^s = (-0.06 \pm 0.50 \pm 0.36) \%$$

(Update coming soon; error reduces by factor 2)

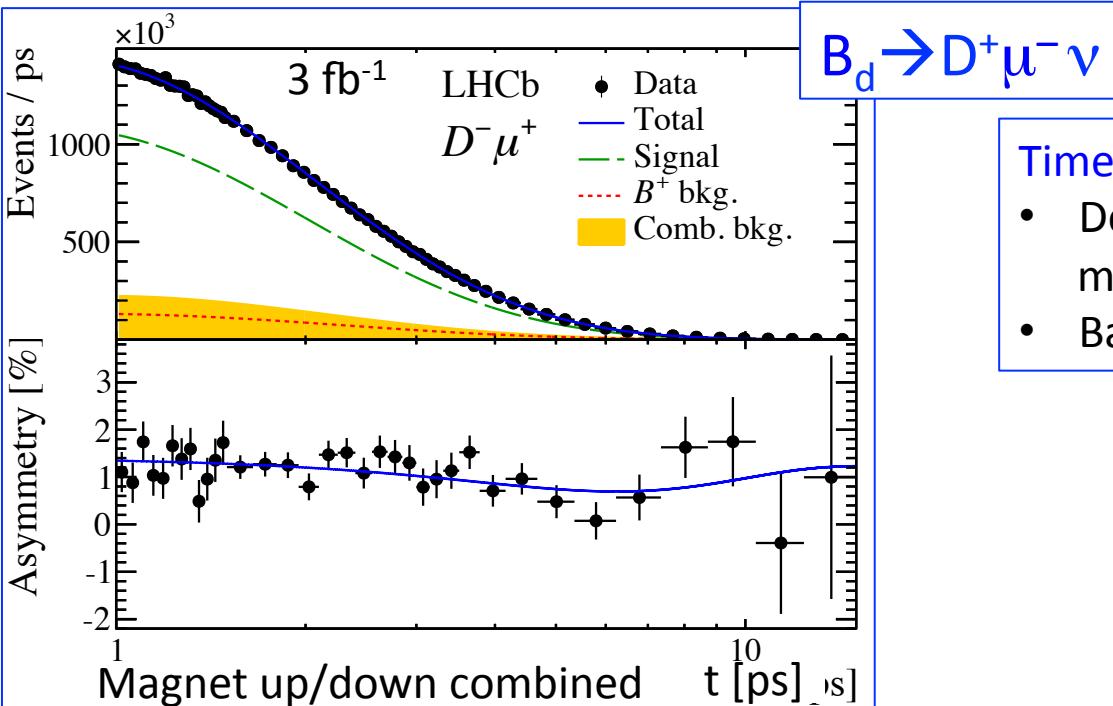
Standard Model:

Lenz, Nierste JHEP.06 (2007) 072

$$a_{sl}^s = (1.9 \pm 0.3) \times 10^{-5}$$



Semileptonic charge asymmetry with B_d : a_{sl}^d



Time dependent analysis:

- Decay time reconstruction with missing neutrino
- Background modelling

LHCb:

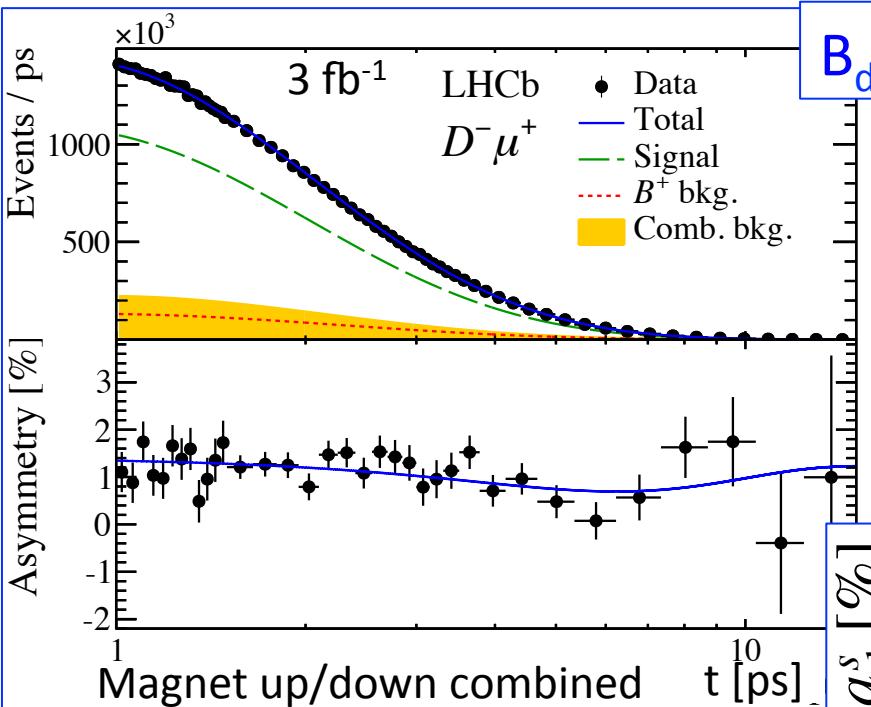
PRL 114 (2015) 041601

$$a_{sl}^d = (-0.02 \pm 0.19 \pm 0.30) \%$$

Standard Model: Lenz, Nierste JHEP.06 (2007) 072

$$a_{sl}^d = (-4.1 \pm 0.6) \times 10^{-4}$$

Semileptonic charge asymmetry with B_d : a_{sl}^d

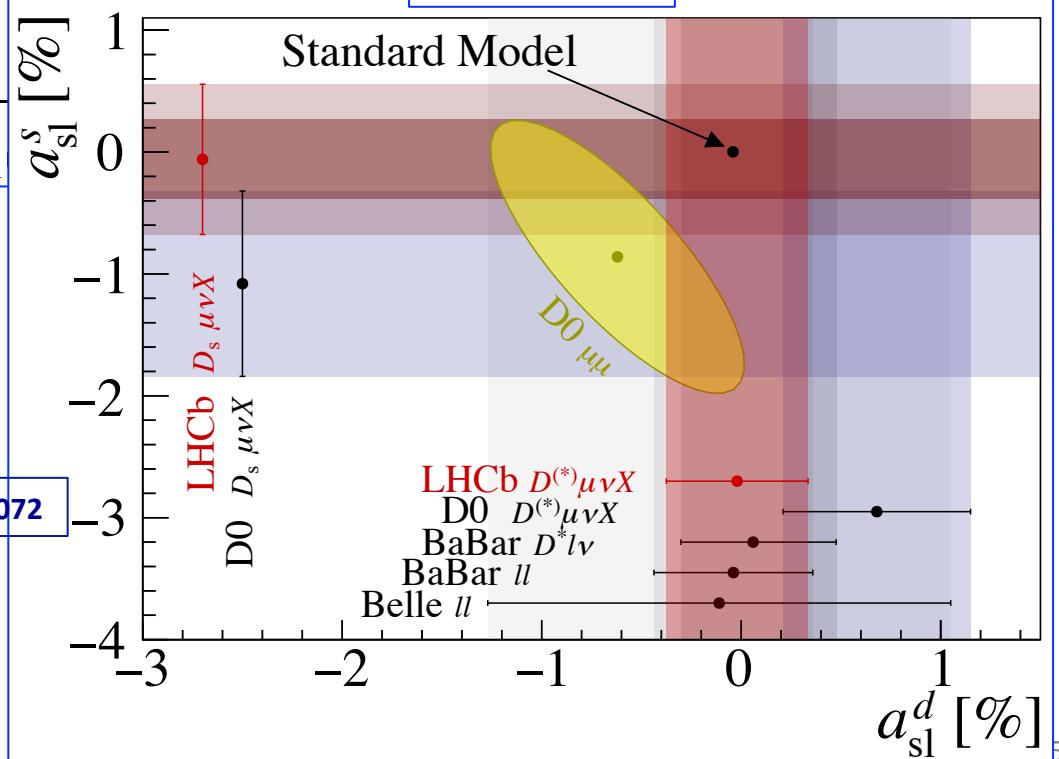


$B_d \rightarrow D^+ \mu^- \nu$

Time dependent analysis:

- Decay time reconstruction with missing neutrino
- Background modelling

a_{sl}^s vs a_{sl}^d



LHCb:

PRL 114 (2015) 041601

$$a_{sl}^d = (-0.02 \pm 0.19 \pm 0.30) \%$$

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Lenz, Nierste JHEP.06 (2007) 072

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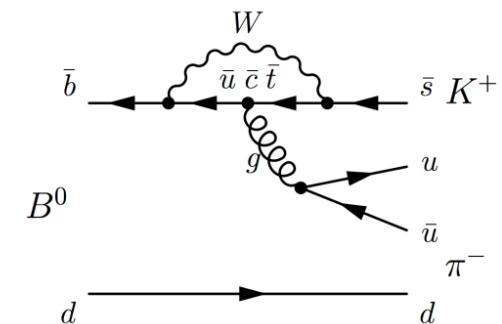
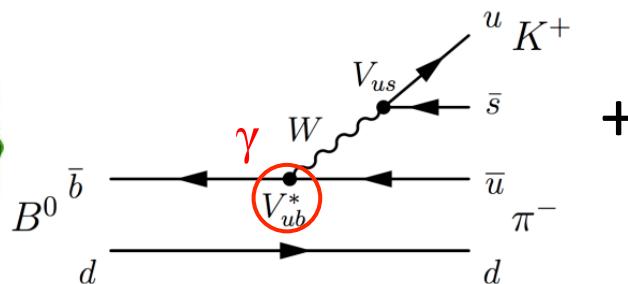
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Direct CP Violation with $B_{(s)} \rightarrow K\pi$ (“charmless”)



- Interference of trees and penguins
 - New physics can contribute to penguin loop
- Sensitive to V_{ub} phase: CKM angle γ
- Measure the *untagged* CP asymmetry for B^0 and B_s decays:

$$\mathcal{A}_{CP} = \frac{(N_{\bar{B} \rightarrow \bar{f}} - N_{B \rightarrow f})}{(N_{\bar{B} \rightarrow \bar{f}} + N_{B \rightarrow f})}$$

$$B \rightarrow f = \begin{cases} B^0 \rightarrow K^+ \pi^- \\ B_s \rightarrow \pi^+ K^- \end{cases}$$

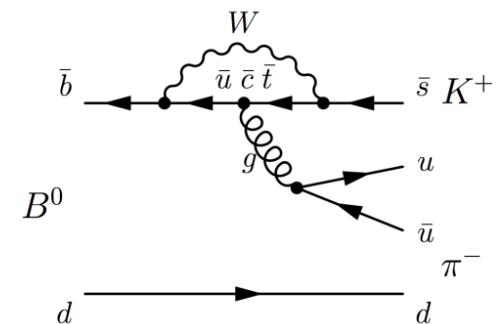
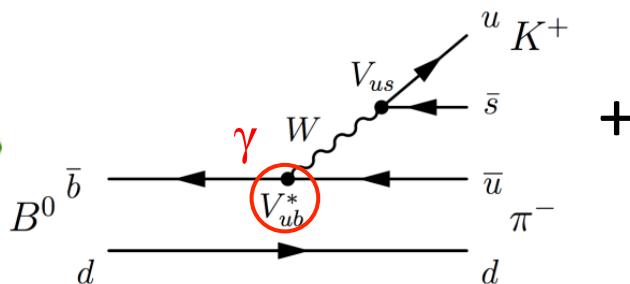
$$\mathcal{A}_{\text{raw}} = \mathcal{A}_{CP} + \mathcal{A}_{\text{det}} + \kappa \cdot \mathcal{A}_{\text{prod}}$$

(The correction factor is $O(1\%)$, measured from data)

Instrumental
asymmetry

Mixing dilution times
production asymmetry

Direct CP Violation with $B_{(s)} \rightarrow K\pi$ (“charmless”)



$B^0 \rightarrow K\pi$

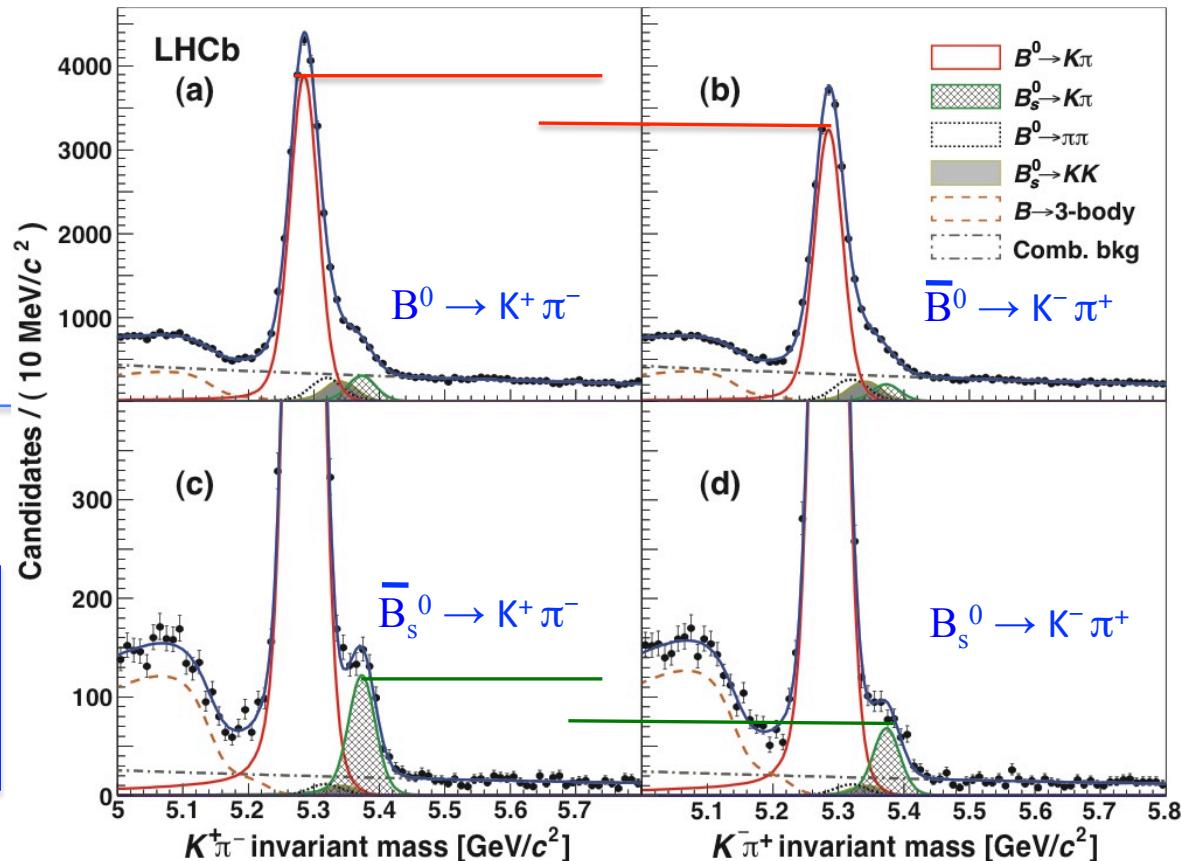
PRL 110.221601

$A_{CP} = -0.080 \pm 0.007 \pm 0.003$
Most precise measurement
of CP violation in a hadronic
Machine.

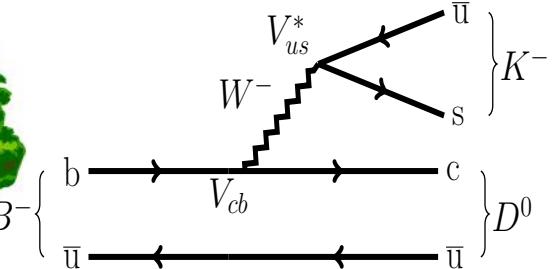
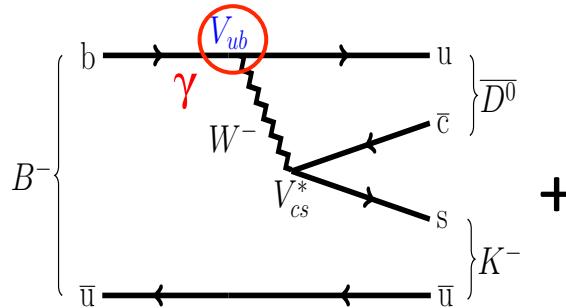
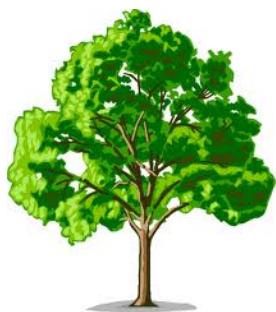
$B_s \rightarrow \pi K$

PRL 110.221601

$A_{CP} = 0.27 \pm 0.04 \pm 0.01$
First observation of a
CP asymmetry in B_s decays



Direct CP Violation with $B \rightarrow D\bar{K}$ (“charmed”)



- Interference of trees and trees
 - Interfere decays $b \rightarrow c$ with $b \rightarrow u$ to final states common to D^0 and \bar{D}^0
- Sensitive to V_{ub} phase: CKM angle γ :
- Many methods

$$\frac{A(B^- \rightarrow \bar{D}^0 K^-)}{A(B^- \rightarrow D^0 \bar{K}^-)} = r_B e^{i\delta_B} e^{-i\gamma}$$

GLW method:

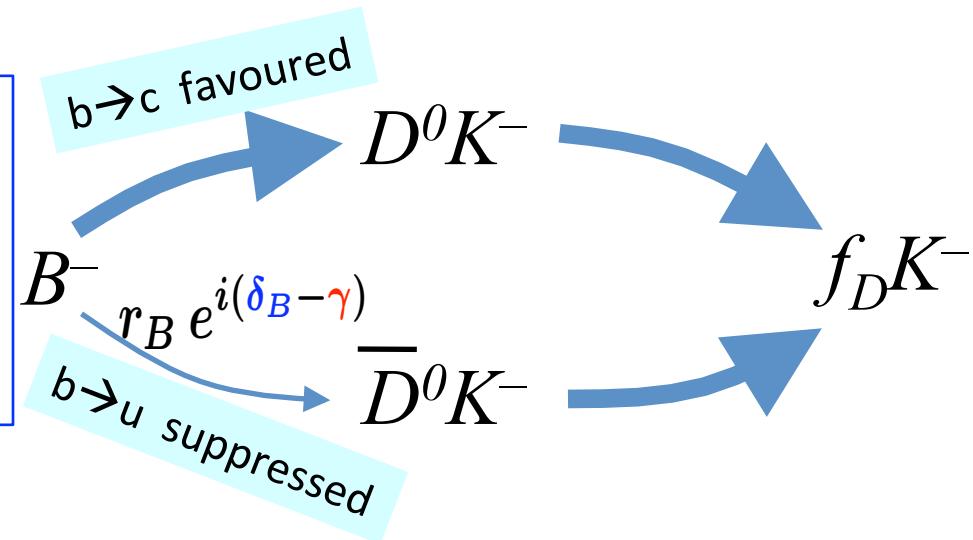
Gronau, London, Wyler:
Phys. Lett. B 265, 172 (1991)

f_D is a CP eigenstate common to

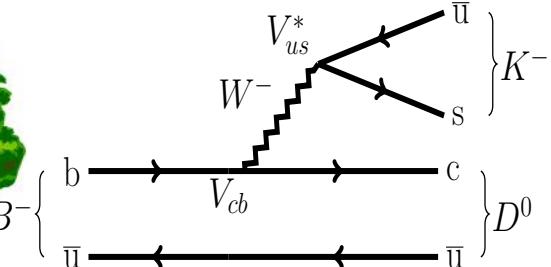
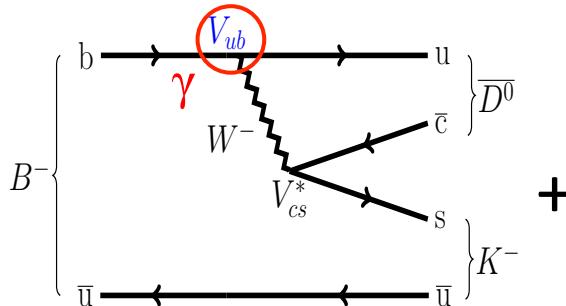
D^0 and \bar{D}^0 : $f_D = K^+ K^-$, $\pi^+ \pi^-$, ...

Interfere: $B \rightarrow D^0 K$, $B \rightarrow \bar{D}^0 K$

• Large event rate; small interference



Direct CP Violation with $B \rightarrow D\bar{K}$ (“charmed”)



- Interference of trees and trees
 - Interfere decays $b \rightarrow c$ with $b \rightarrow u$ to final states common to D^0 and \bar{D}^0
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$$\frac{A(B^- \rightarrow \bar{D}^0 K^-)}{A(B^- \rightarrow D^0 \bar{K}^-)} = r_B e^{i\delta_B} e^{-i\gamma}$$

ADS method:

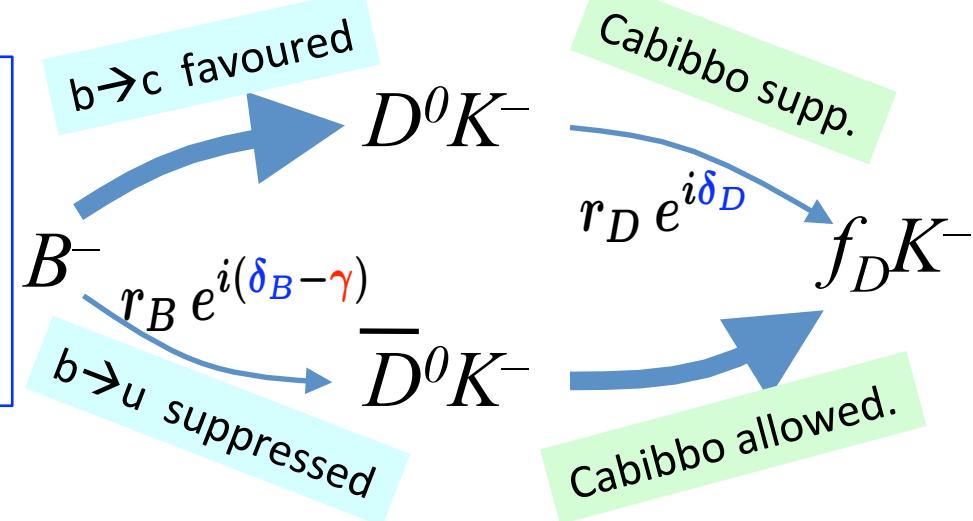
Atwood, Dunietz, Soni:
 Phys Rev Lett 78, 3257 (1997)
 Phys Rev D 63, 036005 (2001)

Use common flavour state $f_D = (K^+ \pi^-)$

Note: decay $D^0 \rightarrow K^+ \pi^-$ is double

Cabibbo suppressed

• Lower event rate; large interference

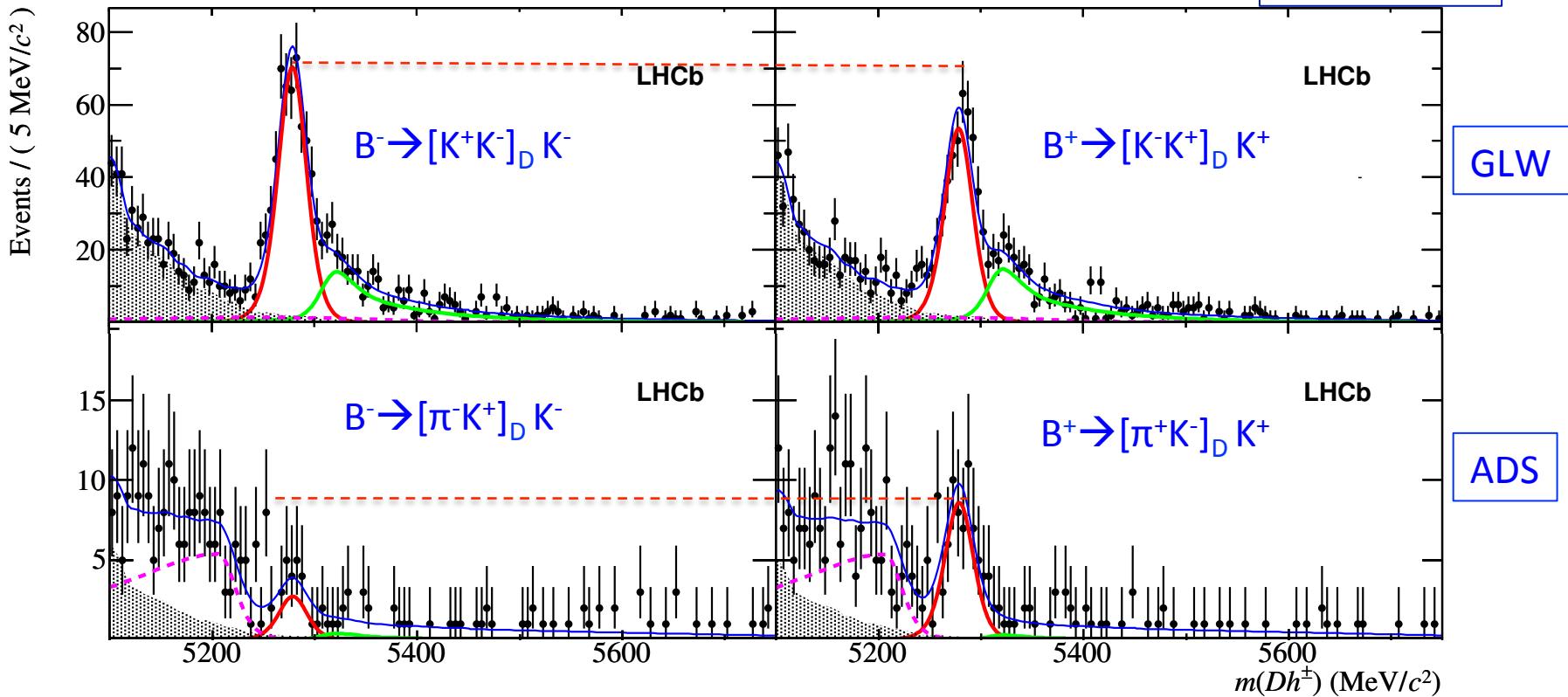


Cabibbo allowed.

Direct CP Violation with $B \rightarrow D K$ (“charmed”)

Observe ADS mode $B^- \rightarrow [\pi^- K^+]_D K^-$ with 10σ significance

PLB 712 (212) 203
arXiv:1203.3662
LHCb: 1.0 fb^{-1}



GLW:

$$A_{CP+} = 0.145 + 0.032 + 0.010$$

$$R_{CP+} = 1.007 + 0.038 + 0.012$$

ADS:

$$A_{ADS} = -0.52 + 0.15 + 0.02$$

$$R_{ADS} = 0.0152 + 0.0020 + 0.0004$$

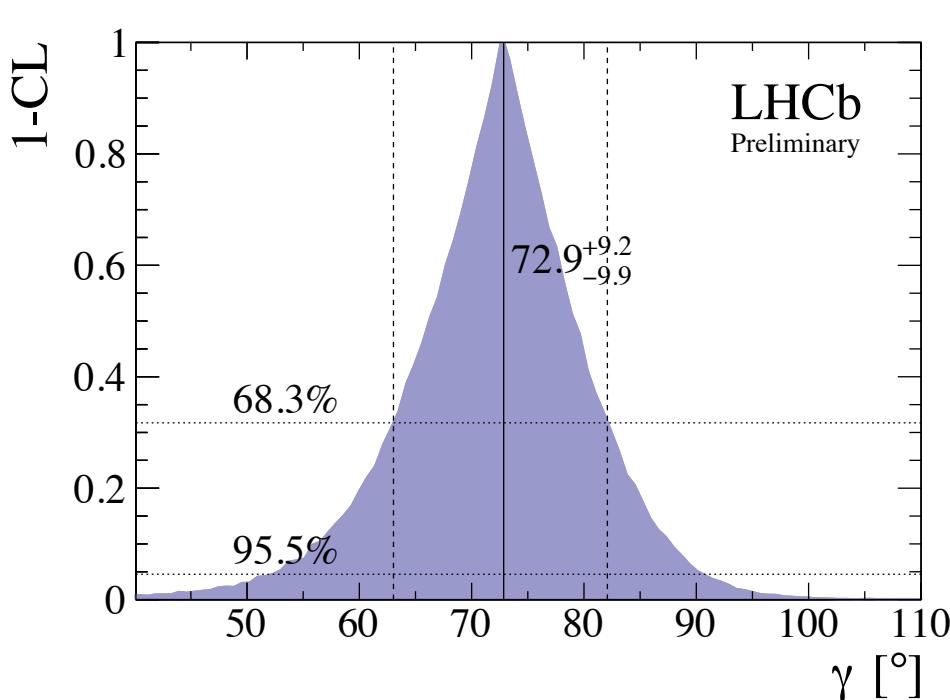
→ CP violation is observed with 5.8σ

γ from $B \rightarrow D K$ and $B \rightarrow D \pi$

More methods/modes available: GLW, ADS, GGSZ, GLS...:

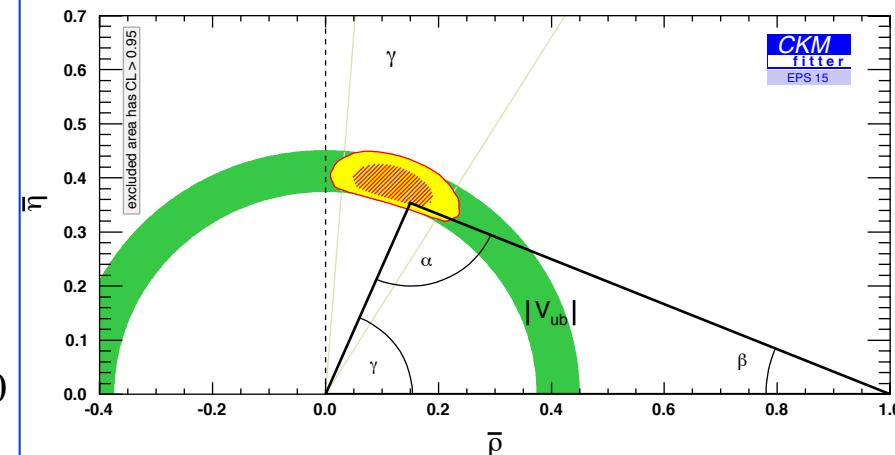
- $B^+ \rightarrow D h^+$ with $D \rightarrow hh$; GLW/ADS method (2011 data): Phys.Lett.B712(212)203
- $B^+ \rightarrow D h^+$ with $D \rightarrow K\pi\pi\pi$; ADS method (2011 data) : Phys.Lett.B723(2013)44
- $B^+ \rightarrow D K^+$ with $D \rightarrow K_s^0 hh$; GGSZ method (2011+2012 data) : arXiv:1408.2748
- $B^+ \rightarrow D K^+$ with $D \rightarrow K_s^0 K\pi$; GLS method (2011+2012 data) : Phys.Lett B 733 (2014) 36
- $B^0 \rightarrow D^0 K^{*0}$ with $D \rightarrow hh$; GLW/ADS (2011+2012 data) : arXiv:1407.8136

PLB 726 (2013) 151, LHCb-CONF-2014-004



Most precise γ determination:

$$\gamma_{B \rightarrow Dh} = (72.9^{+9.2}_{-9.9})^\circ$$



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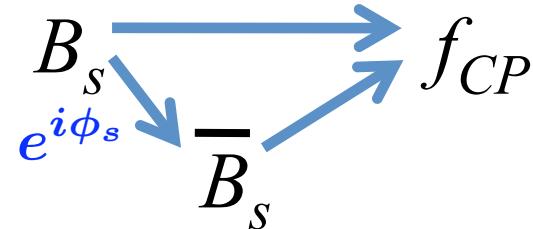
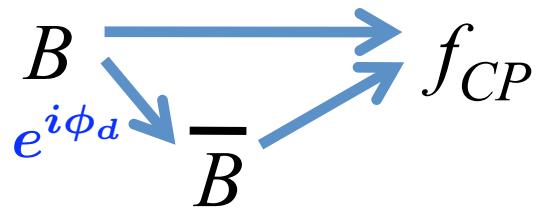
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Time Dependent CP Violation: mixing phase

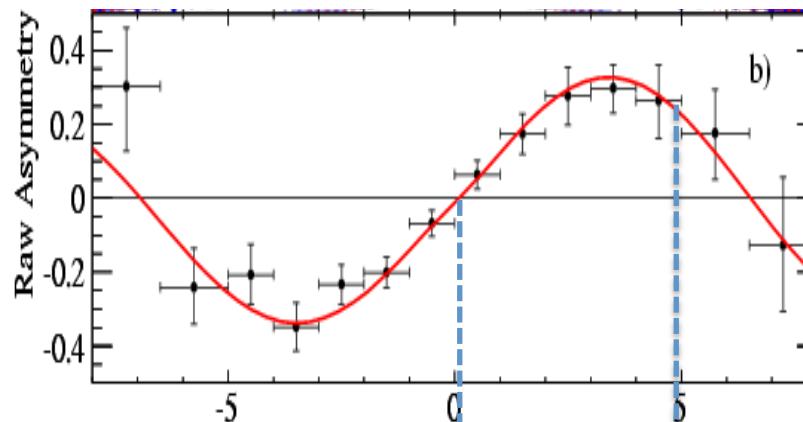


$$\mathcal{A}_{CP}(t) = \frac{\Gamma_{\bar{B} \rightarrow f}(t) - \Gamma_{B \rightarrow f}(t)}{\Gamma_{\bar{B} \rightarrow f}(t) + \Gamma_{B \rightarrow f}(t)} = \frac{\mathcal{A}^{dir} \cos(\Delta M t) + \mathcal{A}^{mix} \sin(\Delta M t)}{\cosh\left(\frac{\Delta \Gamma}{2} t\right) - \mathcal{A}^{\Delta \Gamma} \sinh\left(\frac{\Delta \Gamma}{2} t\right)}$$

$$(\mathcal{A}^{dir})^2 + (\mathcal{A}^{mix})^2 + (\mathcal{A}^{\Delta \Gamma})^2 = 1$$

- Extraction of \mathcal{A}^{dir} and \mathcal{A}^{mix} require flavour tagging: knowing the flavour of the B -meson at production
- Measure ϕ_d (SM: $\mathcal{A}^{mix} = \sin 2\beta$) and ϕ_s (SM: $\mathcal{A}^{mix} = -\sin 2\beta_s$)

ϕ_d : B-factory golden mode: $B^0 \rightarrow J/\psi K_s$



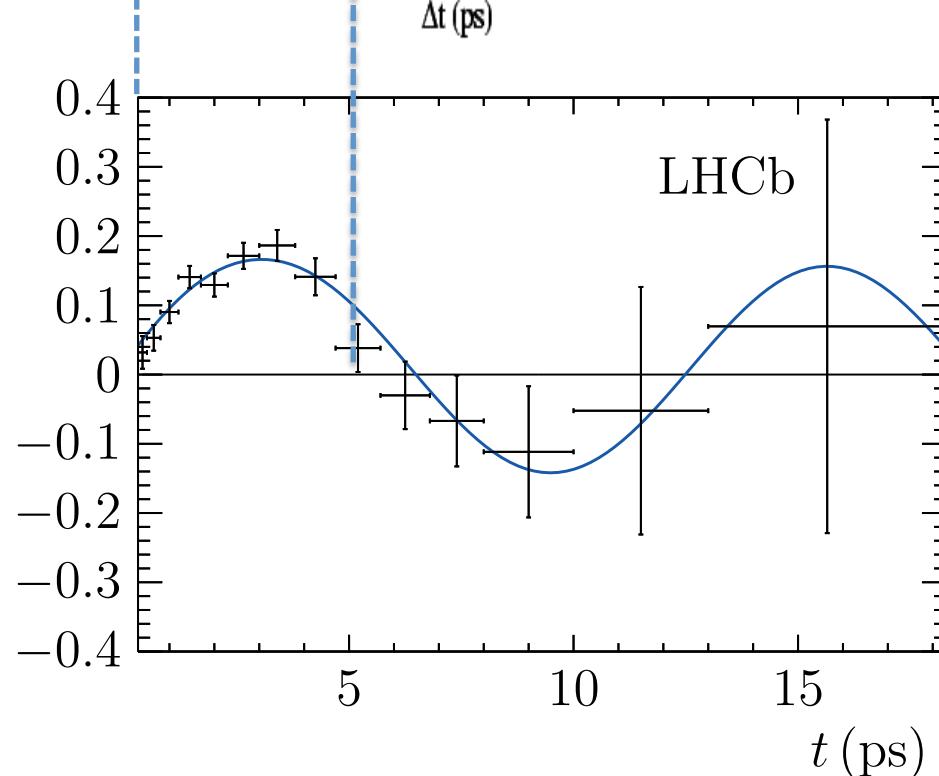
B-factories:

- Coherent production, clean flavour tag

Babar: Phys.Rev.D79 (2009) 072009
Belle: PRL 108 (2012) 171802

- Babar:
 $\sin 2\beta = 0.662 \pm 0.039 \pm 0.012$
- Belle:
 $\sin 2\beta = 0.670 \pm 0.029 \pm 0.013$
- LHCb: PRL. 115 (2015) 031601
(Not so golden mode for LHCb:
 K_s decays often outside Velo)
 $\sin 2\beta = 0.731 \pm 0.035 \pm 0.020$

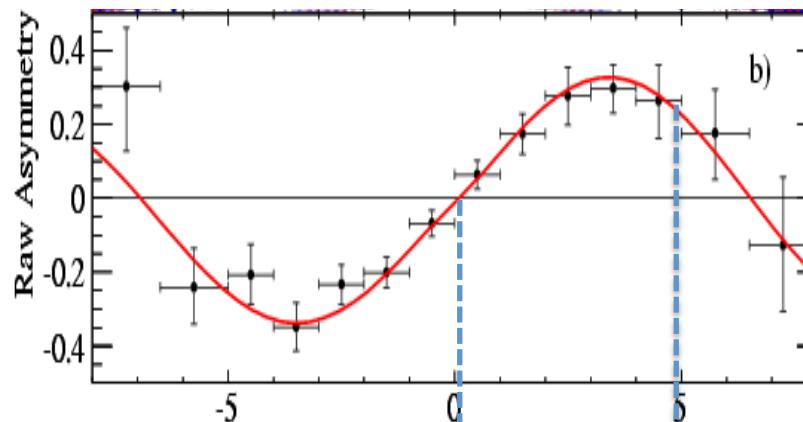
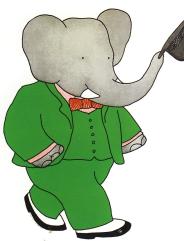
Signal yield asymmetry



(CKM Mixing Phase: $\phi_d = 2\beta$)

ϕ_d : B-factory golden mode: $B^0 \rightarrow J/\psi K_s$

BABAR



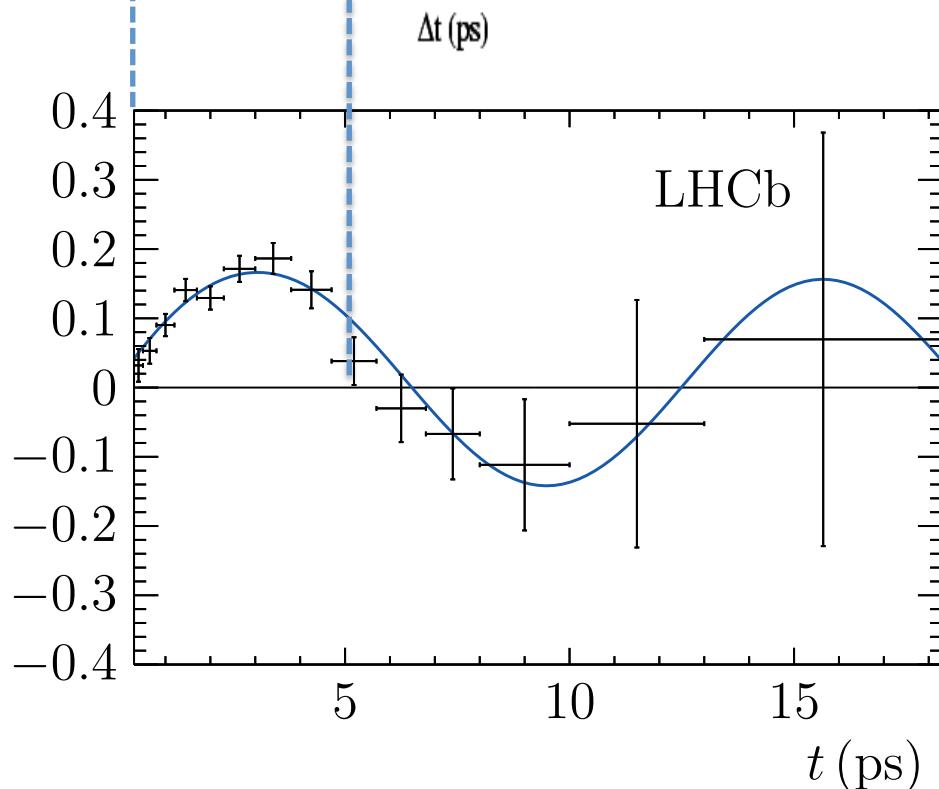
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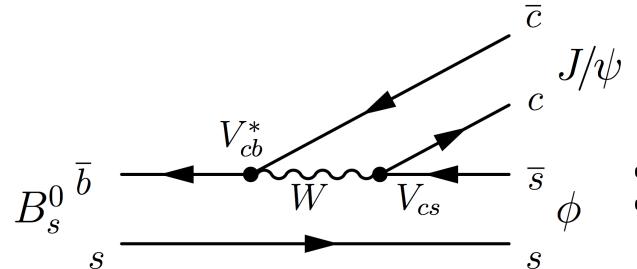
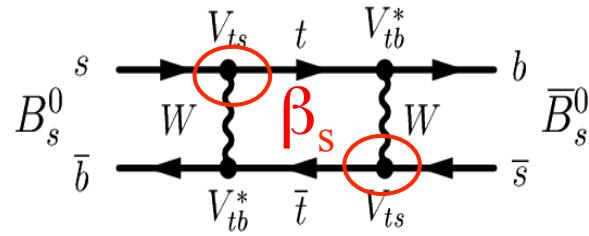
(CKM Mixing Phase: $\phi_d = 2\beta$)



ϕ_s : LHCb golden mode: $B_s \rightarrow J/\psi \phi$

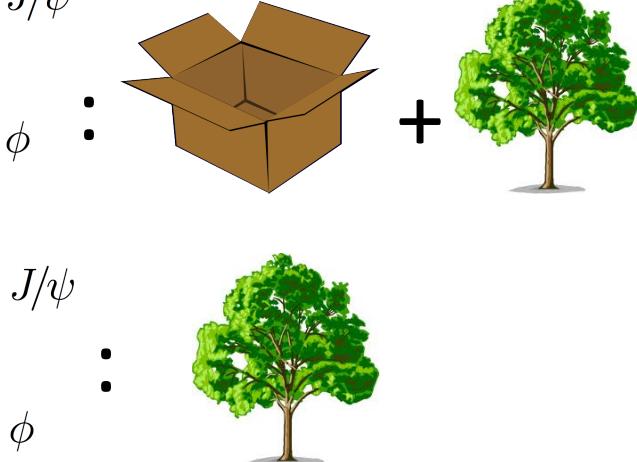
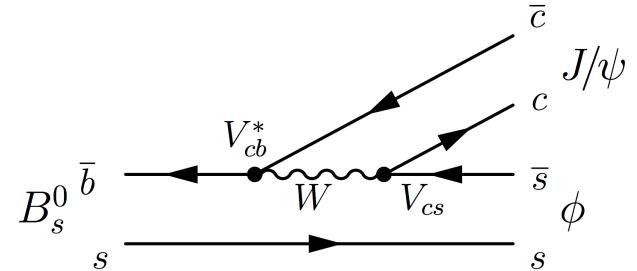
Two decay amplitudes to the same CP eigenstate final state.

1) via mixing $B_s \rightarrow \bar{B}_s \rightarrow J/\psi \phi$:



(CKM Mixing Phase:
 $\phi_s = -2\beta_s \approx 0$)

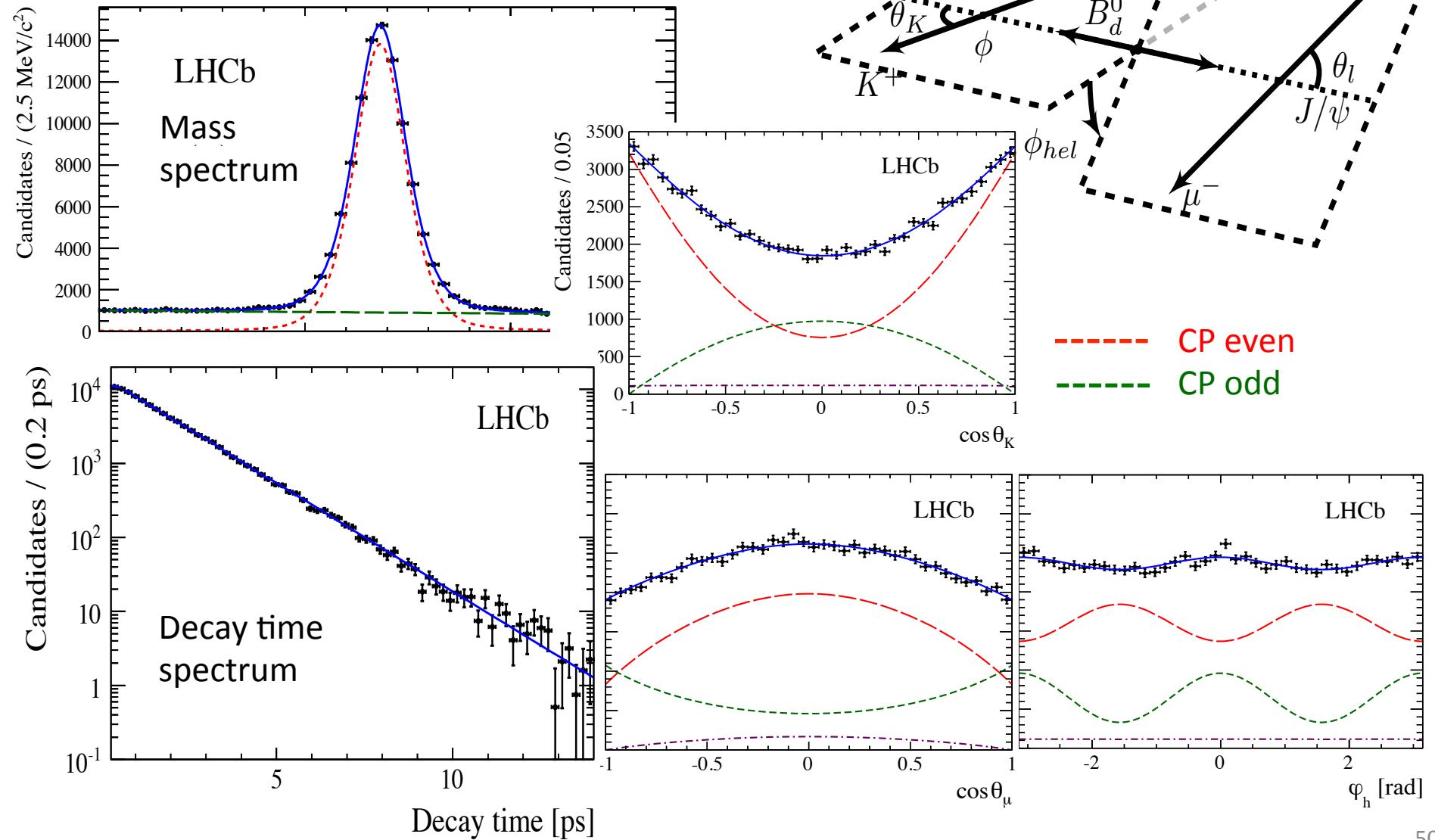
2) direct: $B_s \rightarrow J/\psi \phi$:



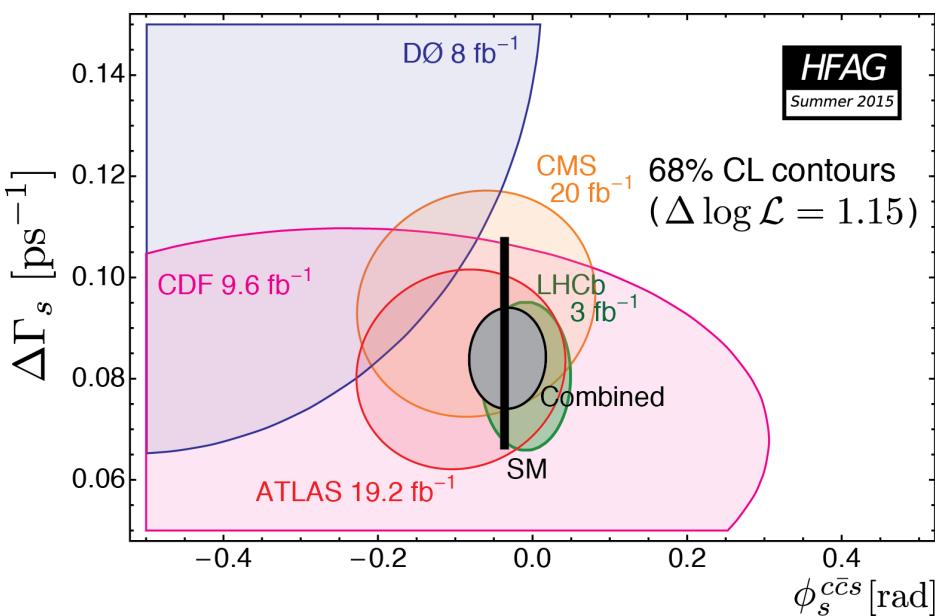
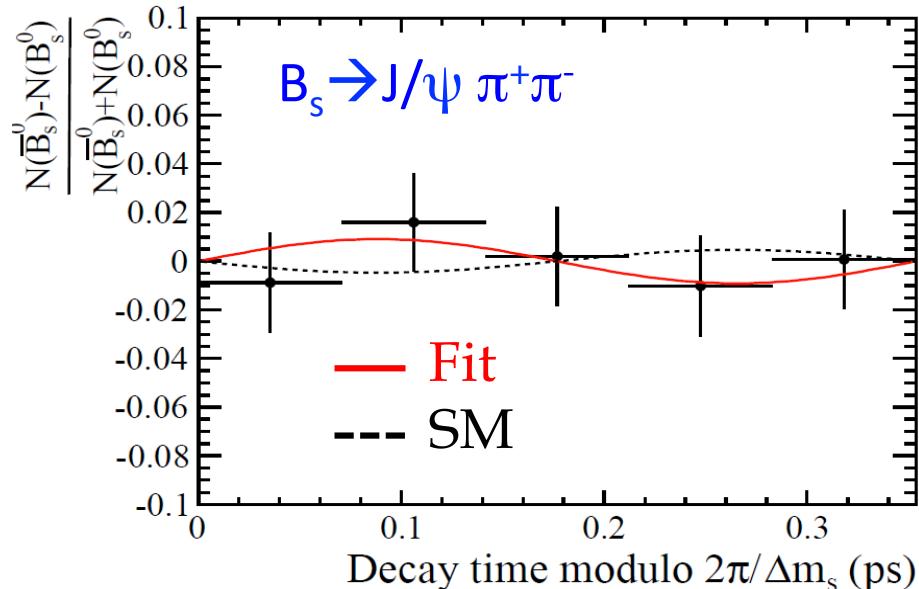
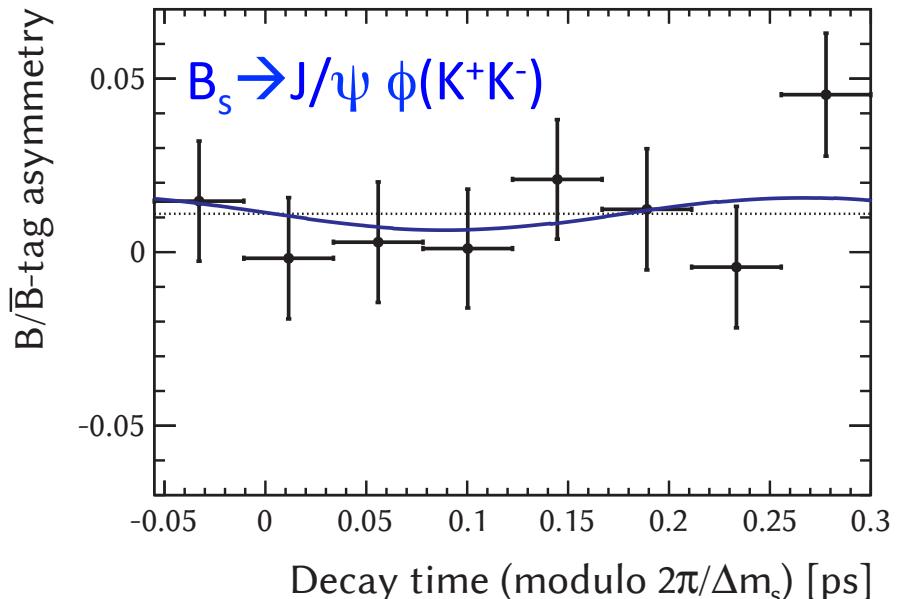
- Golden mode for a time dependent measurement of “ ϕ_s via mixing diagram”
 - Hadronic uncertainties small
- Sensitive to new physics that enters via the mixing loop
- B_s is a *pseudoscalar* ($s=0$) while J/ψ and ϕ are *vector* particles ($s=1$)
 - Final state is superposition of CP even ($L=0$ and $L=2$) CP odd ($L=1$)
 - Requires angular analysis to disentangle
- Alternative analysis in pure CP odd eigenstate $B_s \rightarrow J/\psi \pi^+ \pi^-$

ϕ_s : LHCb golden mode: $B_s \rightarrow J/\psi \phi$

- Requires combined decay time and angular analysis to statistically disentangle the CP even and CP odd amplitudes...



ϕ_s : LHCb golden mode: $B_s \rightarrow J/\psi \phi$



LHCb: **PRL 114 (2015) 041801**
 $\phi_s = -0.010 \pm 0.039 \text{ [rad]}$

Standard Model: **CKMFitter**
 $\phi_s = -2\beta_s = -0.0376 \pm 0.0008 \text{ [rad]}$

Consistent.

2011 + 2012
data: 3 fb^{-1}

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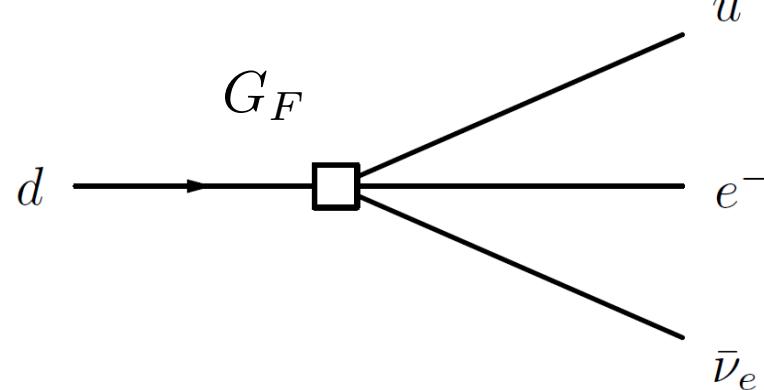


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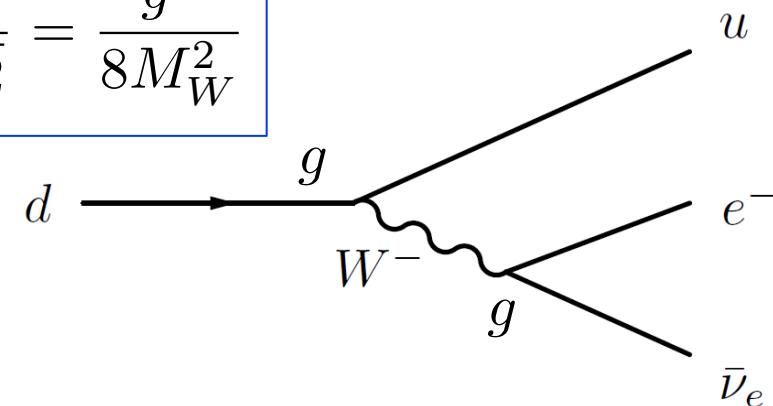
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Effective Couplings

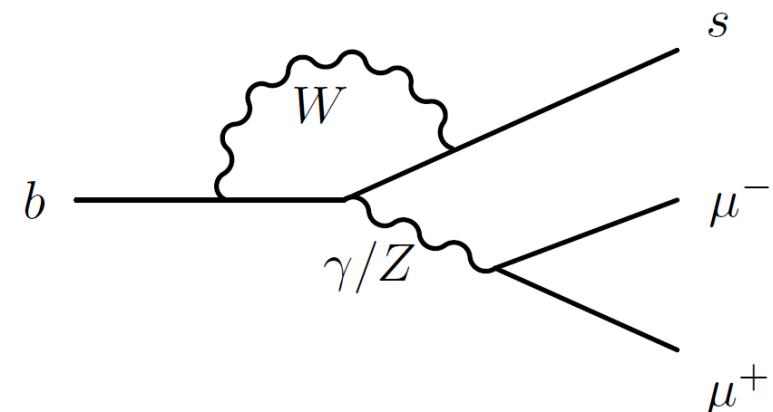
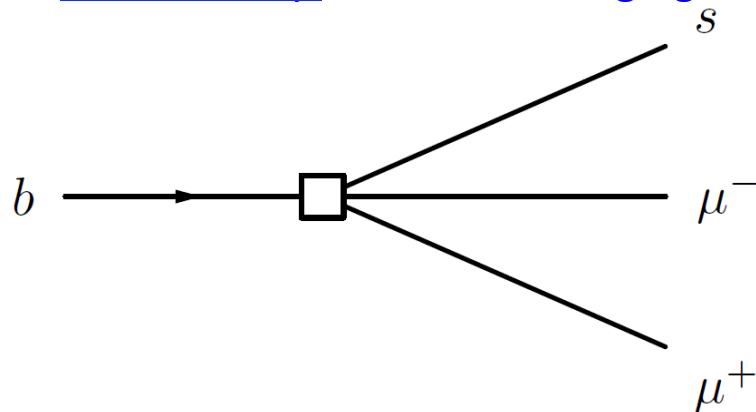
- Beta decay: “charged current”:



$$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8M_W^2}$$



- Rare B decay: “Flavour changing neutral current”:



- Effective Hamiltonian

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{CKM} \sum_i \mathcal{C}_i \mathcal{O}_i$$

Effective local Operators \mathcal{O}_i with Wilson coefficients \mathcal{C}_i predicted by the Standard Model.

Outline

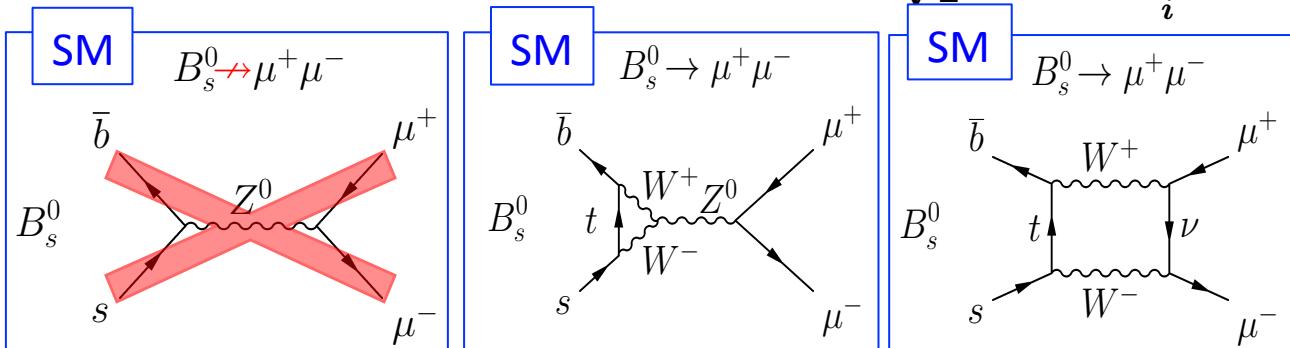
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Very Rare Decays

$$B_s^0 \rightarrow \mu^+ \mu^-$$

$$B_d^0 \rightarrow \mu^+ \mu^-$$

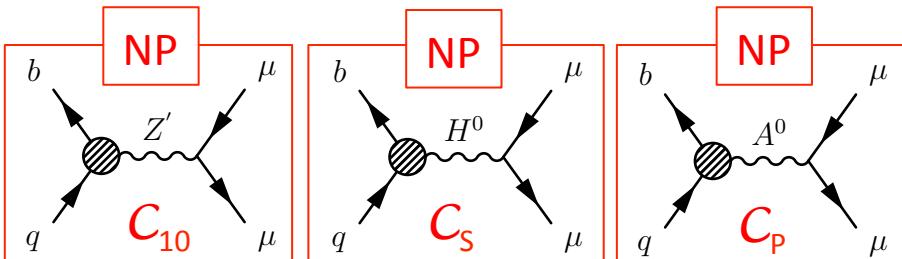
SM: CKM and helicity suppressed: very small B.R.
 → Axial vector coupling C_{10}



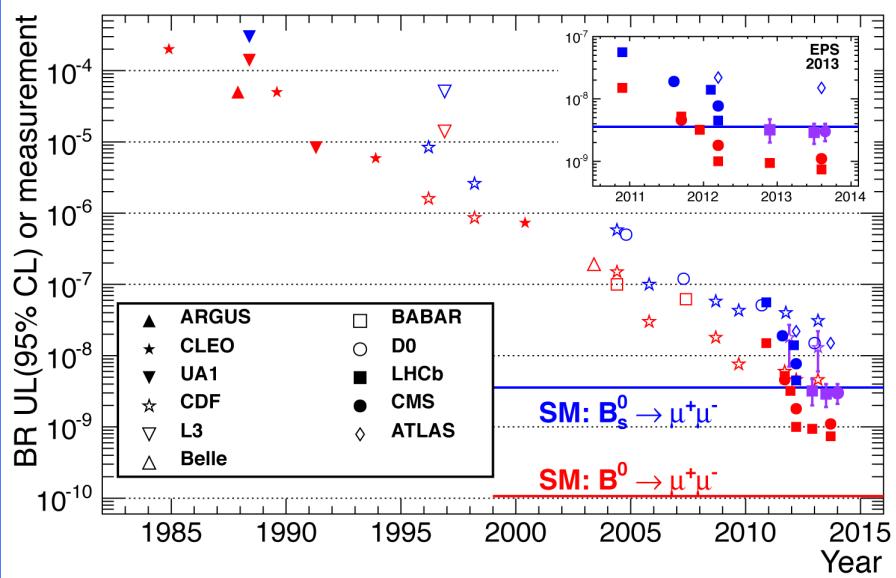
$$\text{BR}(B_s^0 \rightarrow \mu^+ \mu^-) = (3.66 \pm 0.23) \times 10^{-9}$$

$$\text{BR}(B_d^0 \rightarrow \mu^+ \mu^-) = (1.06 \pm 0.09) \times 10^{-10}$$

NP: Sensitive to new particles via additional (C_{10} , C_S , C_P) couplings.
 → e.g.: Z' , (pseudo-)scalars, ...



A 30-year quest for very rare decays....



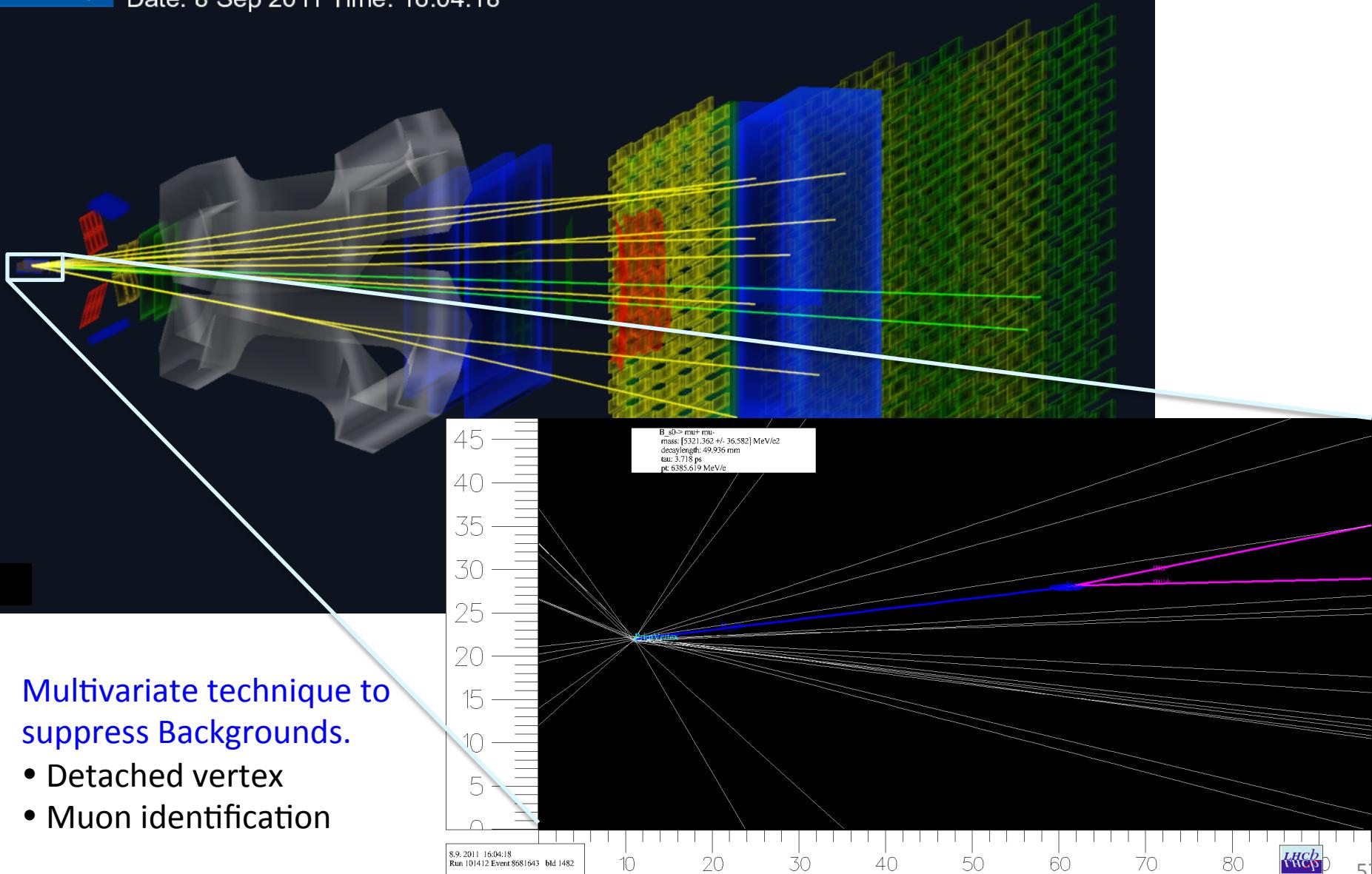
$$BR \propto |V_{tb} V_{tq}|^2 \left[\left(1 - \frac{4m_\mu^2}{M_B^2}\right) |C_S - C'_S|^2 + |(C_P - C'_P) + \frac{2m_\mu}{M_B^2} (C_{10} - C'_{10})|^2 \right]$$

Very Rare Decays



LHCb experiment

Run: 101412 Event: 8681643
Date: 8 Sep 2011 Time: 16:04:18



Multivariate technique to suppress Backgrounds.

- Detached vertex
- Muon identification

Very Rare Decays

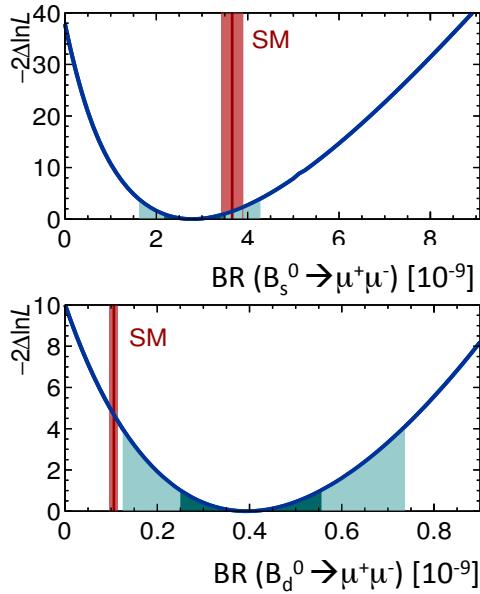
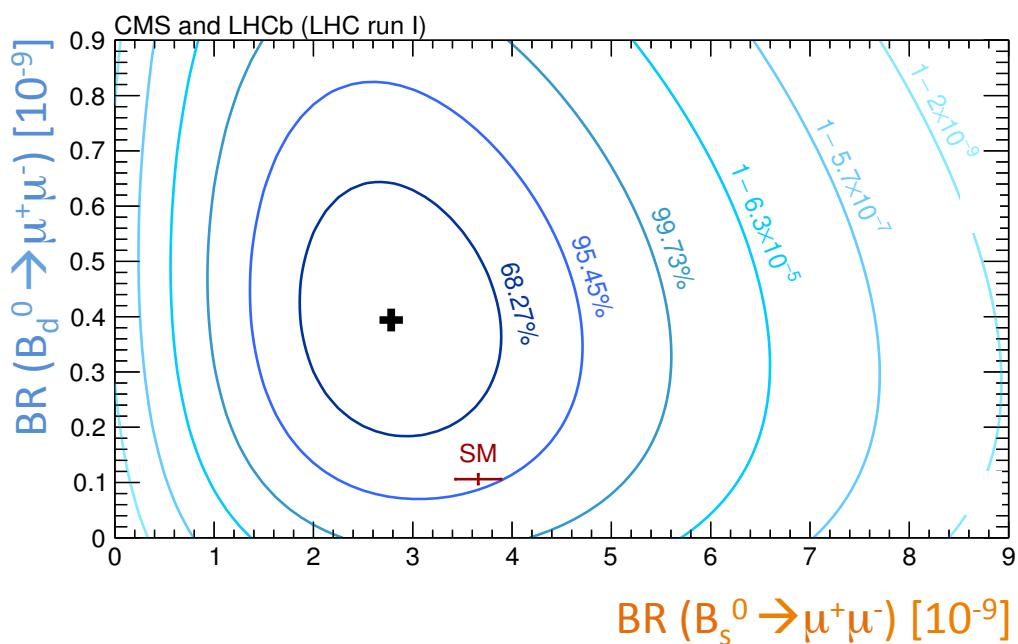
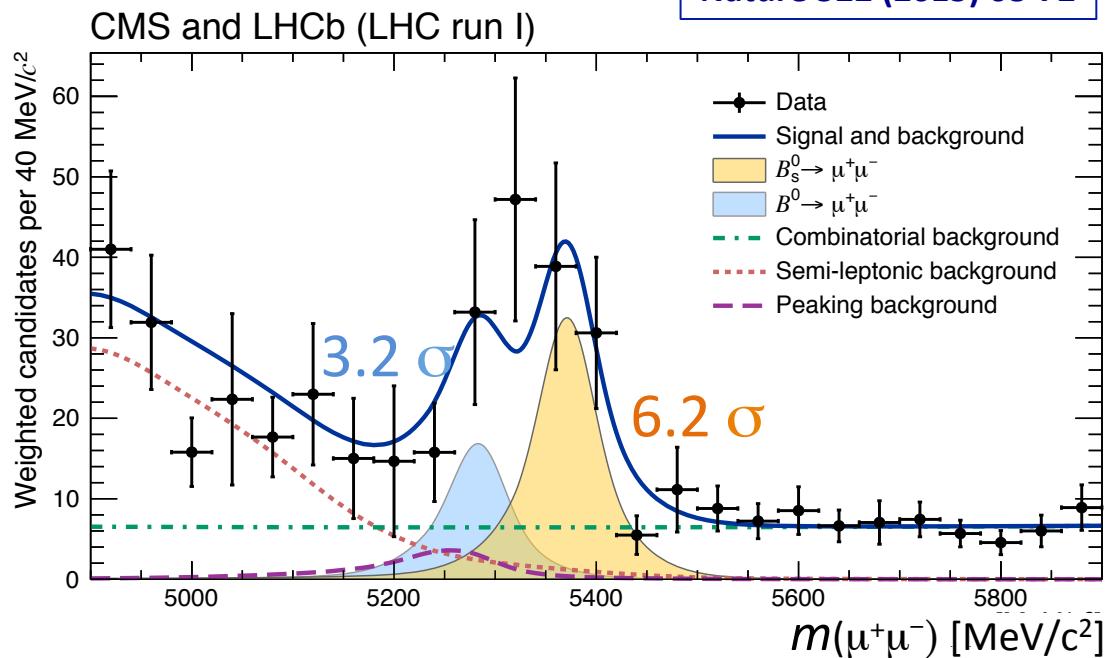
Nature 522 (2015) 68-72

$$B^0_{d/s} \rightarrow \mu^+ \mu^-$$

- LHCb + CMS combined signal:

$$\begin{aligned} BR(B_s^0 \rightarrow \mu^+ \mu^-) &= (2.8 \pm 0.7) \times 10^{-9} \\ BR(B_d^0 \rightarrow \mu^+ \mu^-) &= (3.9 \pm 1.5) \times 10^{-10} \end{aligned}$$

Consistent with SM at $\sim 2\sigma$
Hot topic for LHC Run-2!



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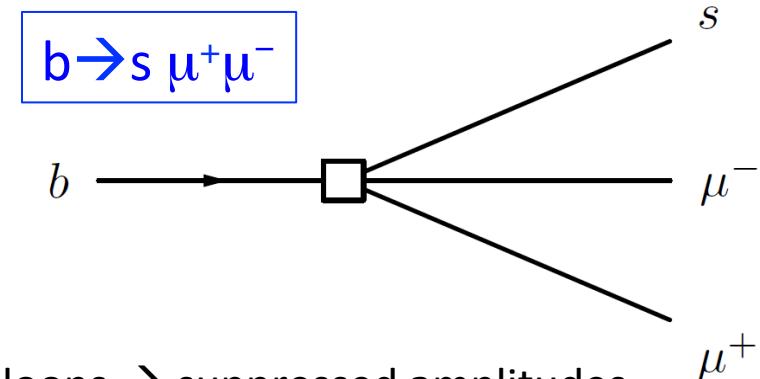
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$b \rightarrow s$ quark transition

- Effective 4-fermion coupling:

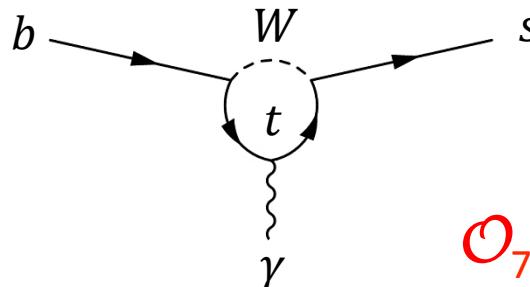
$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i C_i \mathcal{O}_i$$



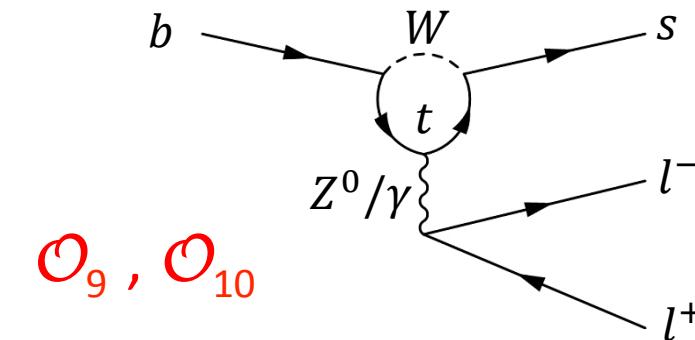
- Standard Model:

- No flavour changing neutral currents: loops \rightarrow suppressed amplitudes

Photon penguin:

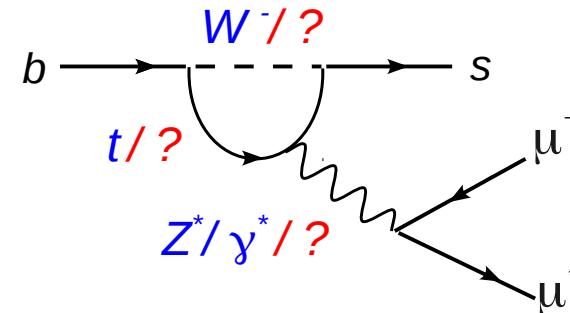


Vector, Axial vector:



- New Physics:

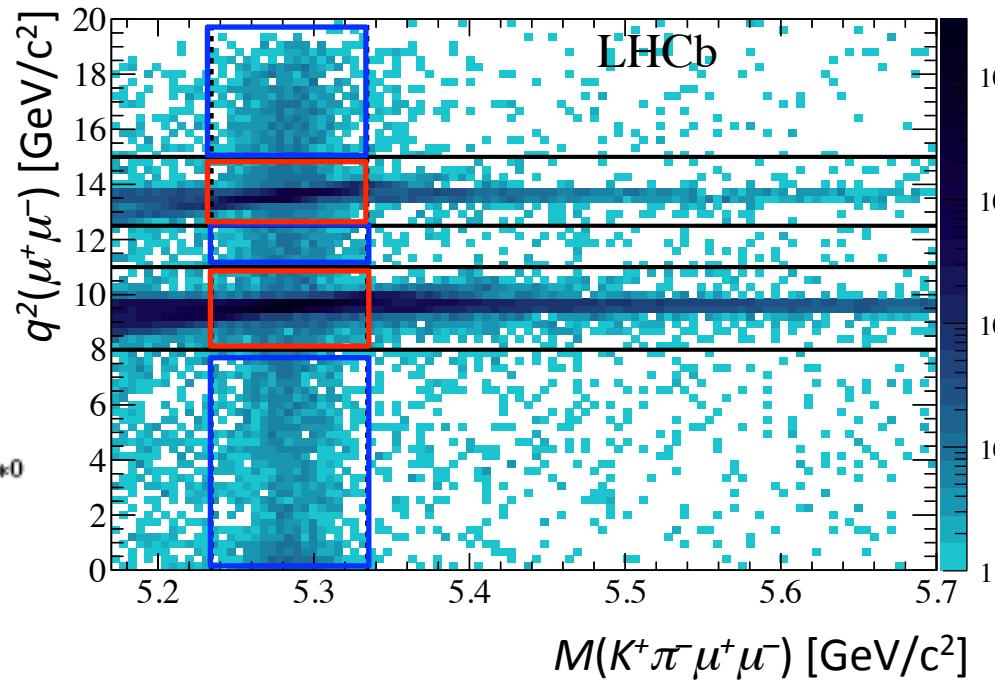
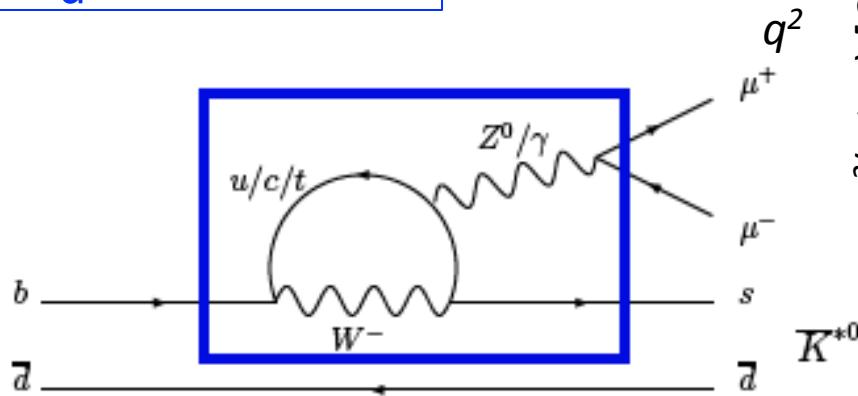
- Sensitivity for NP in Wilson coefficients C_7, C_9, C_{10}



Rare Decays: $b \rightarrow s \mu^+ \mu^-$

arXiv:1512.044442

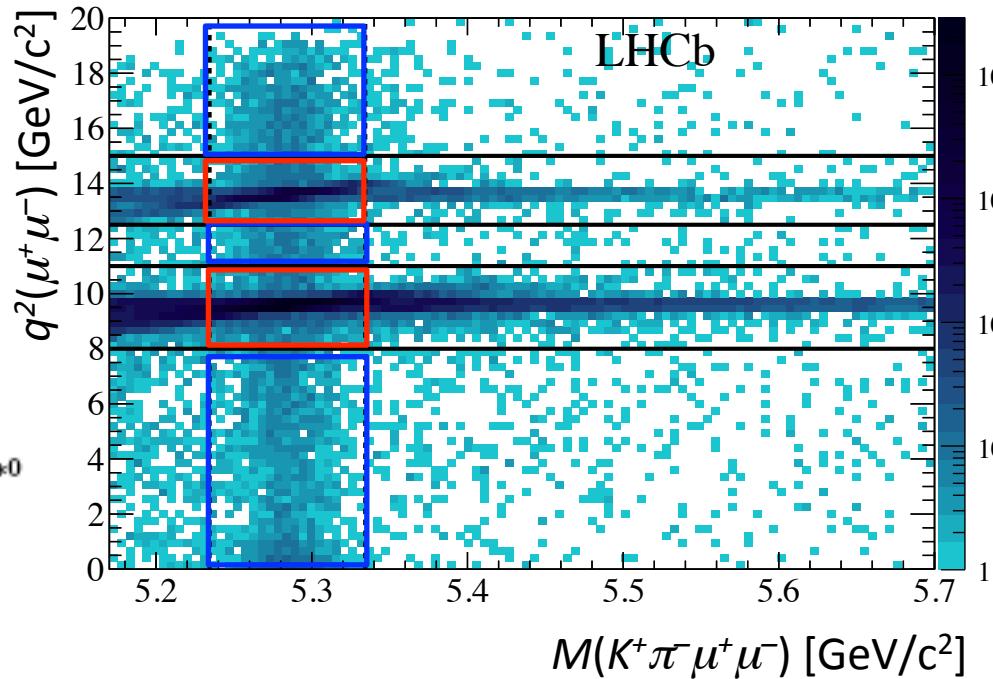
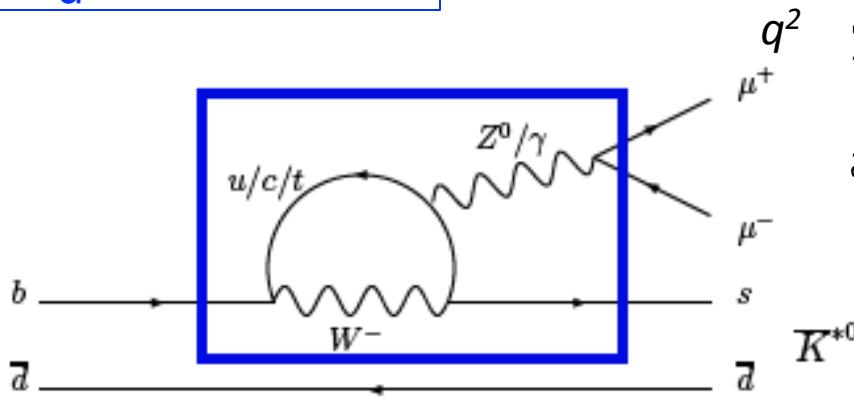
$$B_d^0 \rightarrow K^* \mu^+ \mu^-$$



Rare Decays: $b \rightarrow s \mu^+ \mu^-$

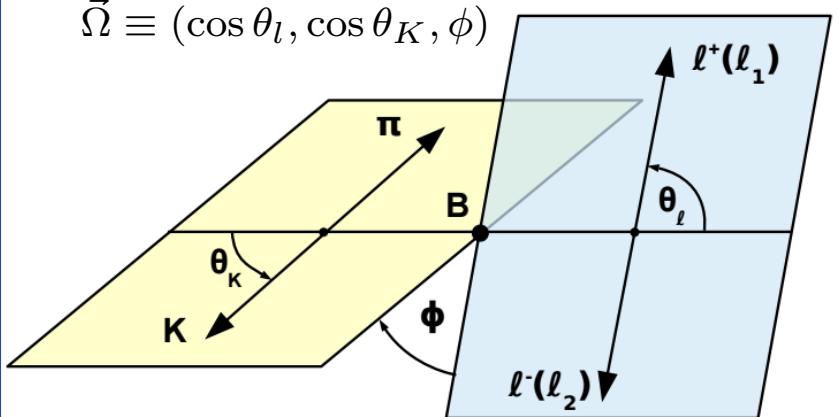
arXiv:1512.044442

$$B_d^0 \rightarrow K^* \mu^+ \mu^-$$



- Study the angular distribution of the final state particles:

$$\vec{\Omega} \equiv (\cos \theta_l, \cos \theta_K, \phi)$$



$$\frac{d^4 \Gamma[\bar{B}^0 \rightarrow K^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_j \mathbf{I}_j(q^2) f_j(\vec{\Omega})$$

$\mathbf{I}_j \rightarrow \bar{\mathbf{I}}_j$ for B^0

- Observables:

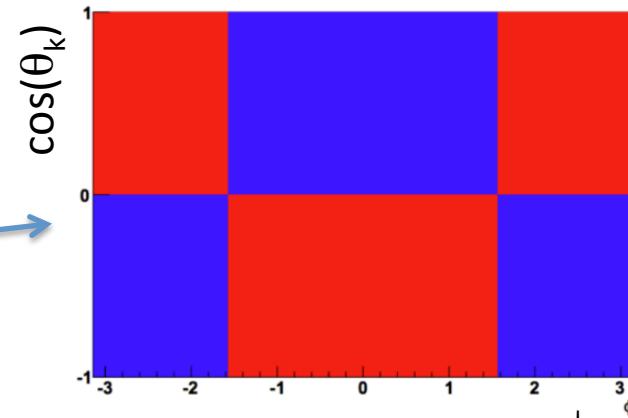
$$S_j = (I_j + \bar{I}_j) / \left(\frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right)$$

$$A_j = (I_j - \bar{I}_j) / \left(\frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right)$$

$B_d^0 \rightarrow K^* \mu^+ \mu^-$

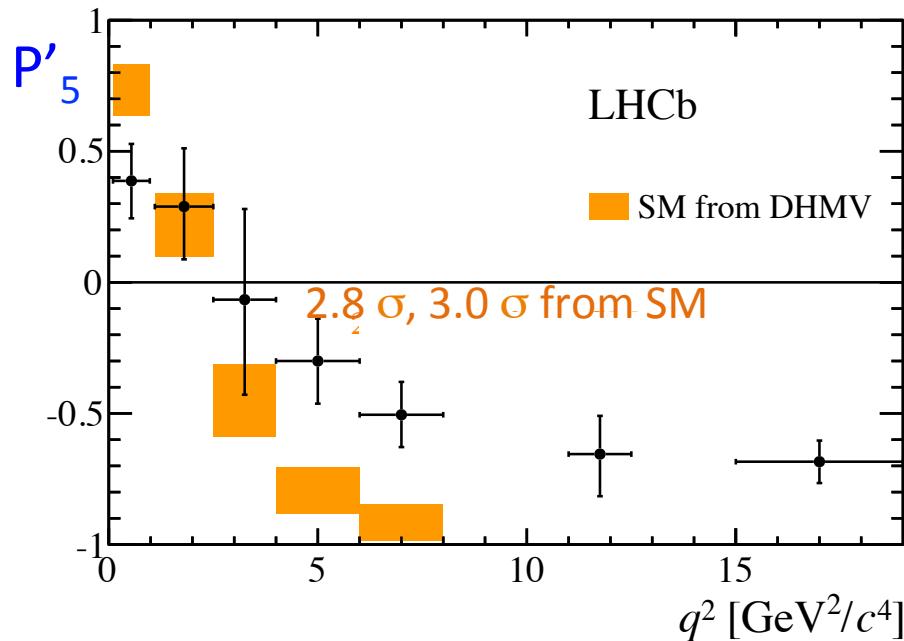
Observables S_5, P'_5 :

- S_5 variable: count blue minus red
- “Robust” P'_5 variable: $P'_5 = \frac{S_5}{\sqrt{F_L (1 - F_L)}}$
(F_L = longitudinal polarisation K^*)

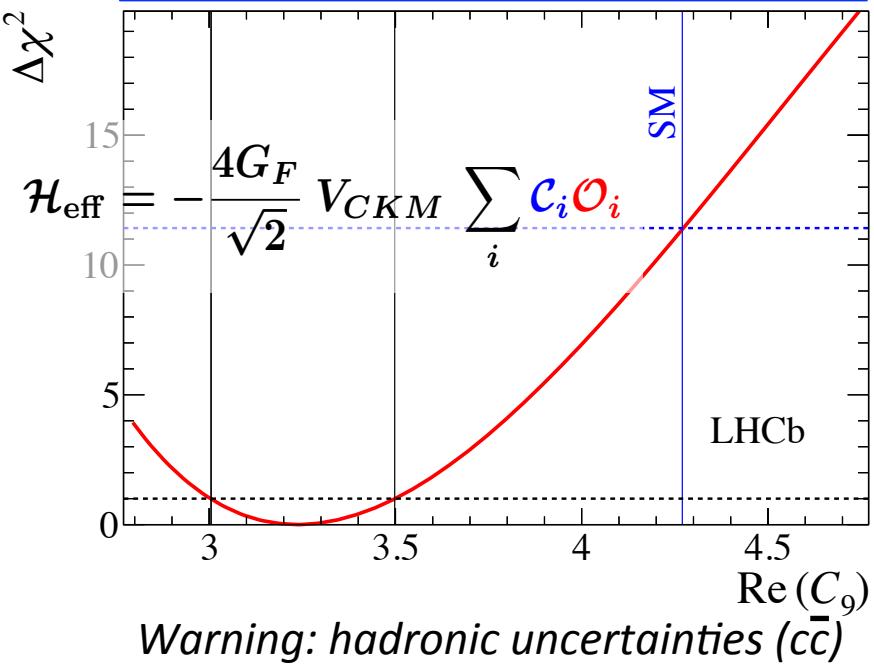


See talk Matthias Neubert: ϕ

P'_5 vs q^2 shows deviations from SM:



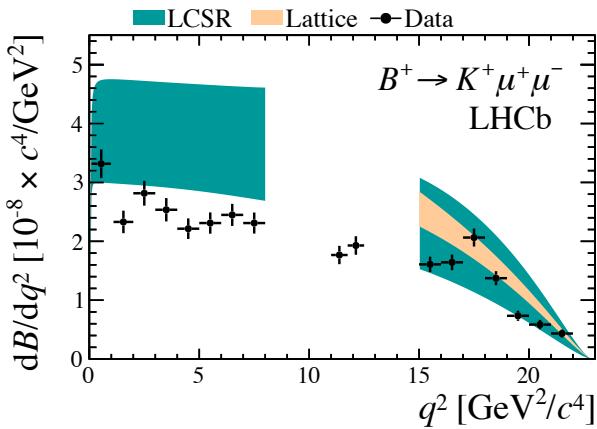
Fit the value of Wilson coefficient:
 $C_9 = C_9^{\text{SM}} + C_9^{\text{NP}}$ points at $C_9^{\text{NP}} \sim -1$



Rare Decays: $b \rightarrow s \mu^+ \mu^-$

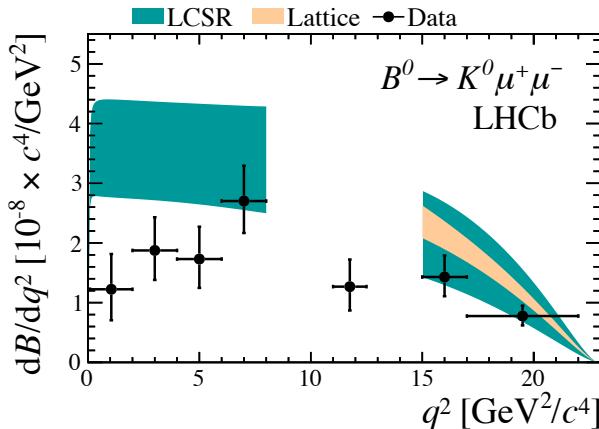
$$B^+ \rightarrow K^+ \mu^+ \mu^-$$

JHEP 06 (2014) 133



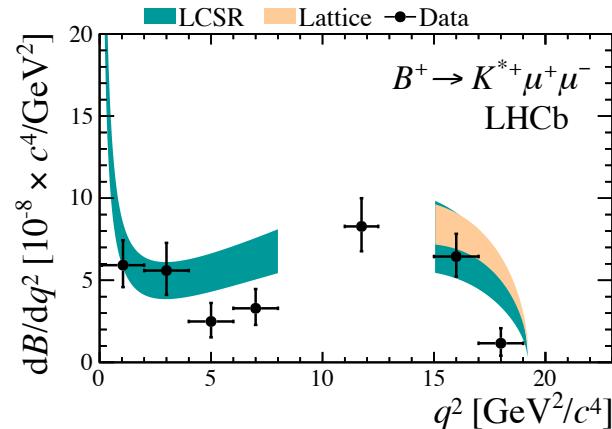
$$B^0 \rightarrow K^0 \mu^+ \mu^-$$

JHEP 06 (2014) 133



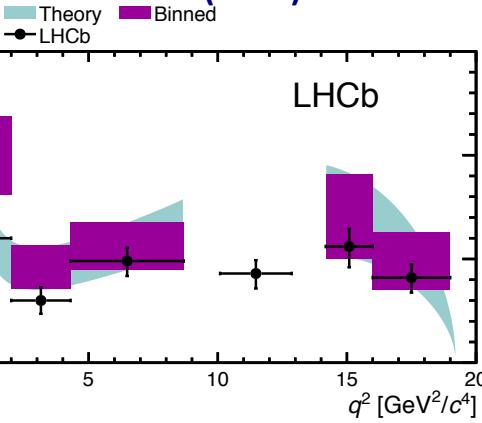
$$B^+ \rightarrow K^{*+} \mu^+ \mu^-$$

JHEP 06 (2014) 133



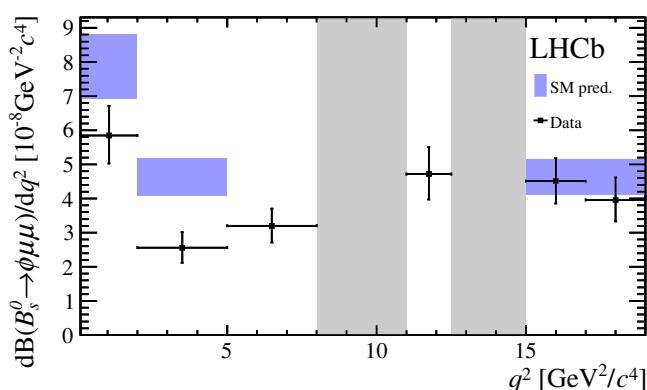
$$B^0 \rightarrow K^{*0} \mu^+ \mu^-$$

JHEP 08 (2013) 131



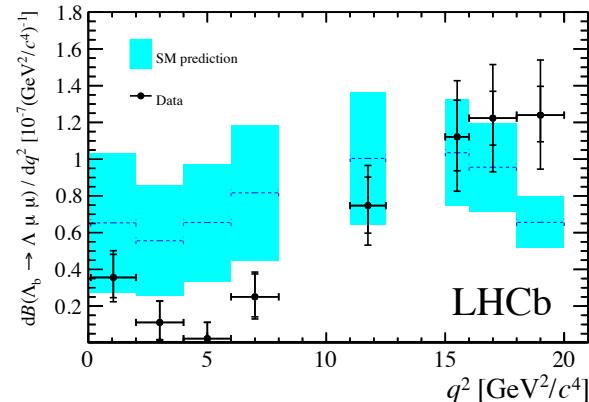
$$B_s \rightarrow \phi \mu^+ \mu^-$$

JHEP 09 (2015) 179



$$\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$$

JHEP 06 (2015) 115



- Branching fractions related to $b \rightarrow s \mu^+ \mu^-$ transition consistently lower than predicted.

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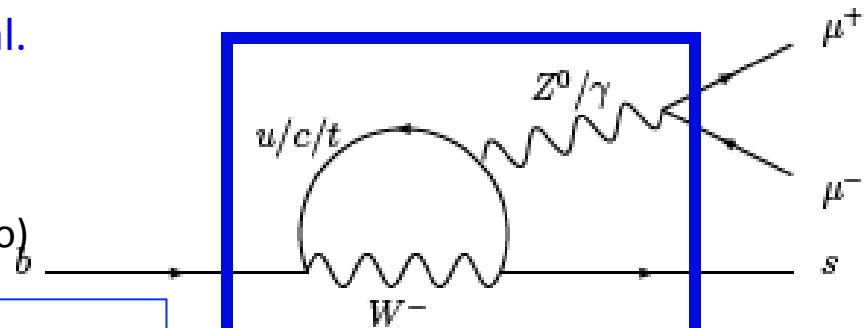
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Lepton Universality: $B^+ \rightarrow K^+ \mu^+ \mu^-$ / $B^+ \rightarrow K^+ e^+ e^-$

- Standard Model: $b \rightarrow s l^+ l^-$ is flavour universal.
- Expect equal branching fractions for $b \rightarrow s \mu^+ \mu^-$ and $b \rightarrow s e^+ e^-$
 - (Hadronic uncertainties largely cancel in the ratio)

$$R_K \equiv \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)} \stackrel{\text{SM}}{=} 1 \pm \mathcal{O}(10^{-2})$$

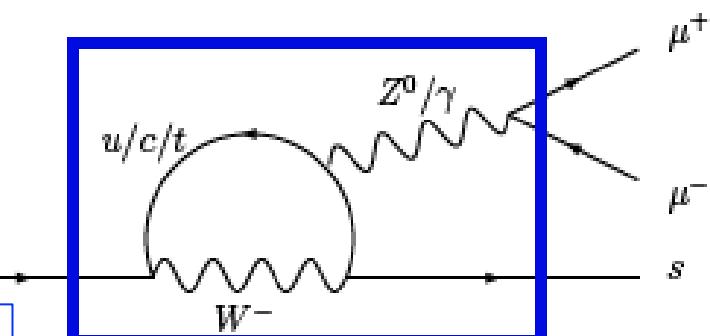


(R_{K^*} coming soon)

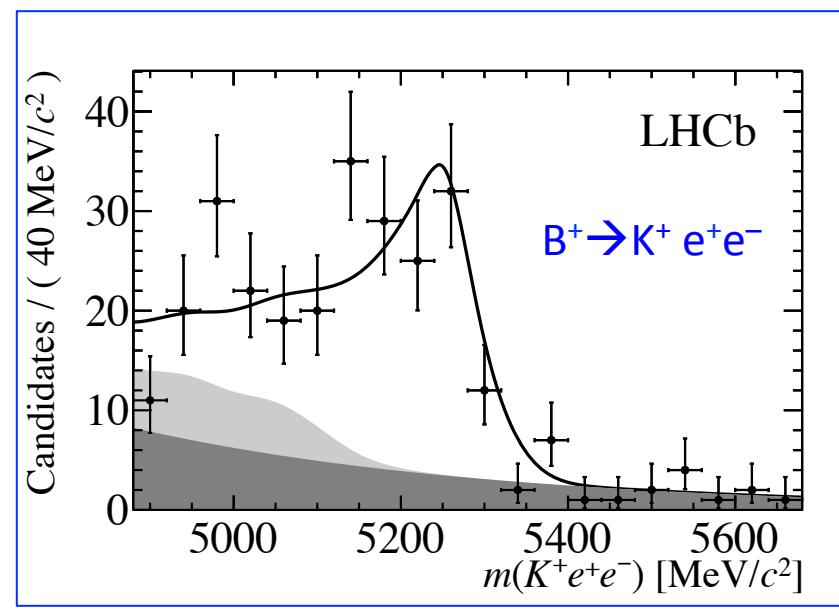
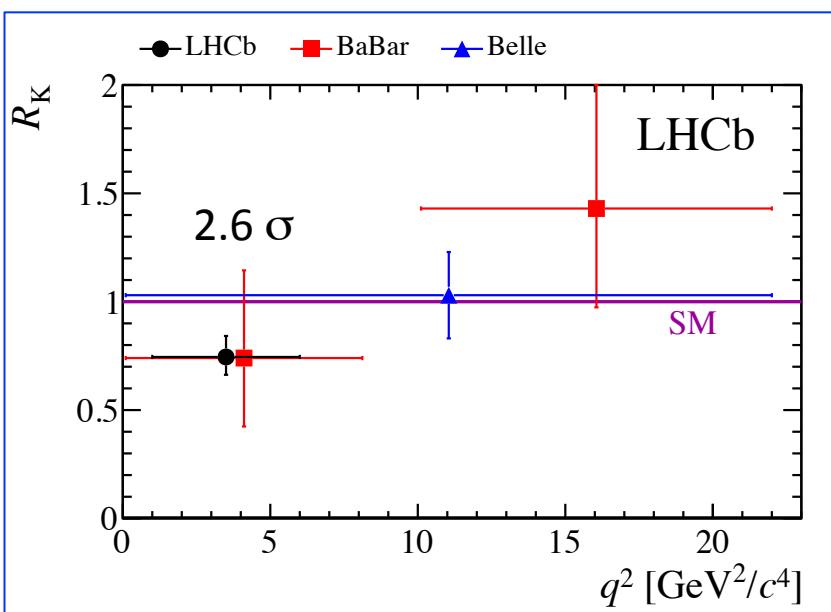
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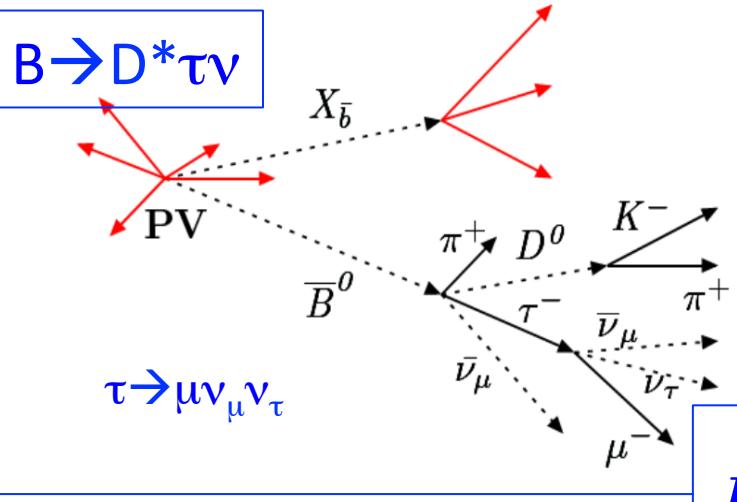
(R_{K^*} coming soon)



$$R_K = 0.745^{+0.090}_{-0.074} \pm 0.036$$

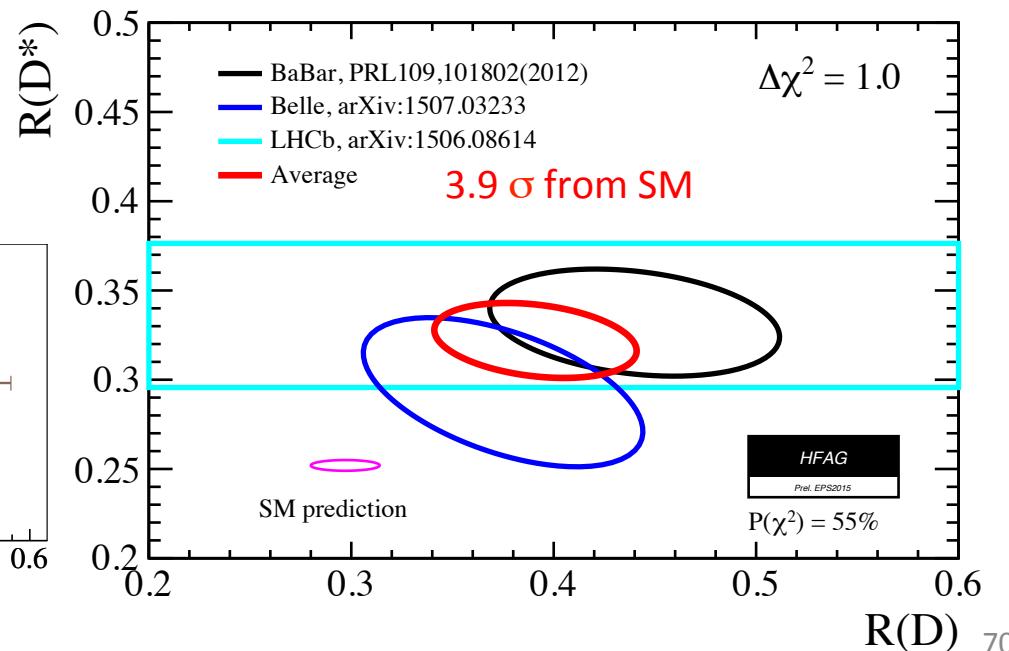
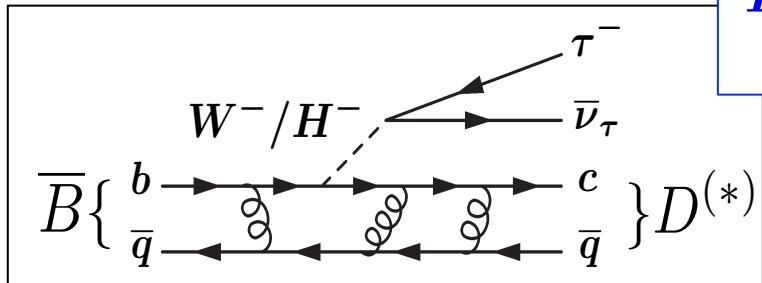
$R_K \neq 1$ would imply NP with non-universal lepton couplings.

Lepton Universality: $B \rightarrow D^* \tau \nu$ / $B \rightarrow D^* \mu \nu$

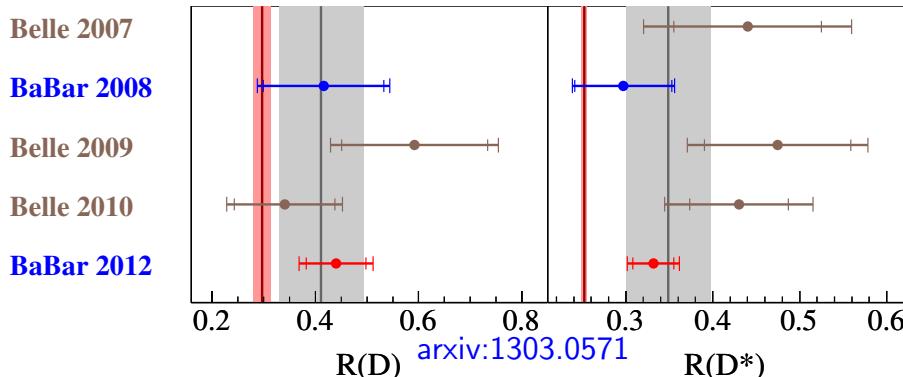


$$R_{D^*} \equiv \frac{\mathcal{B}(B \rightarrow D^* \tau \nu)}{\mathcal{B}(B \rightarrow D^* \mu \nu)}$$

(# 1 due to phase space!)



Before LHCb:



Outline

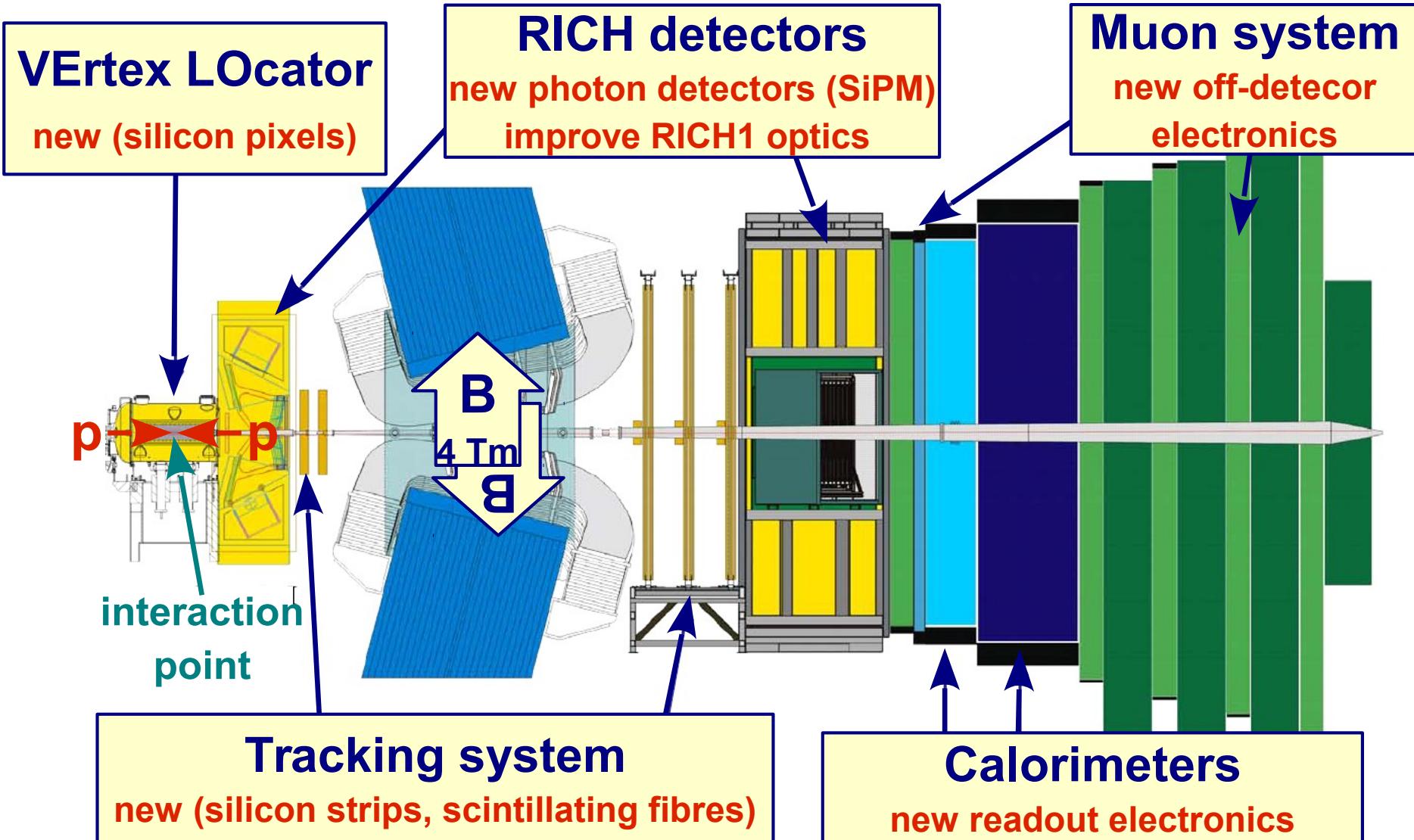
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The Future – LHCb Upgrade



The Future – LHCb Upgrade



www.jolyon.co.uk

**Triggers
today**



**Triggers
in the future**

Expected statistical uncertainties before and after upgrade

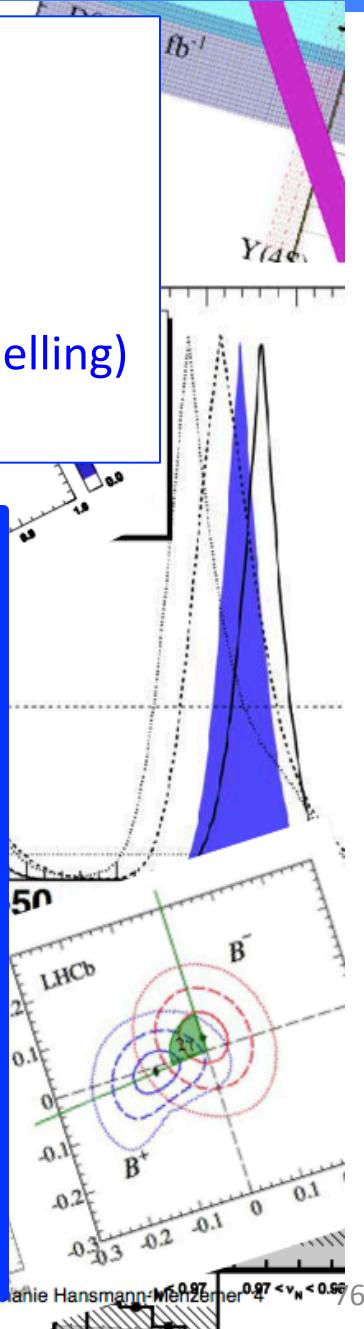
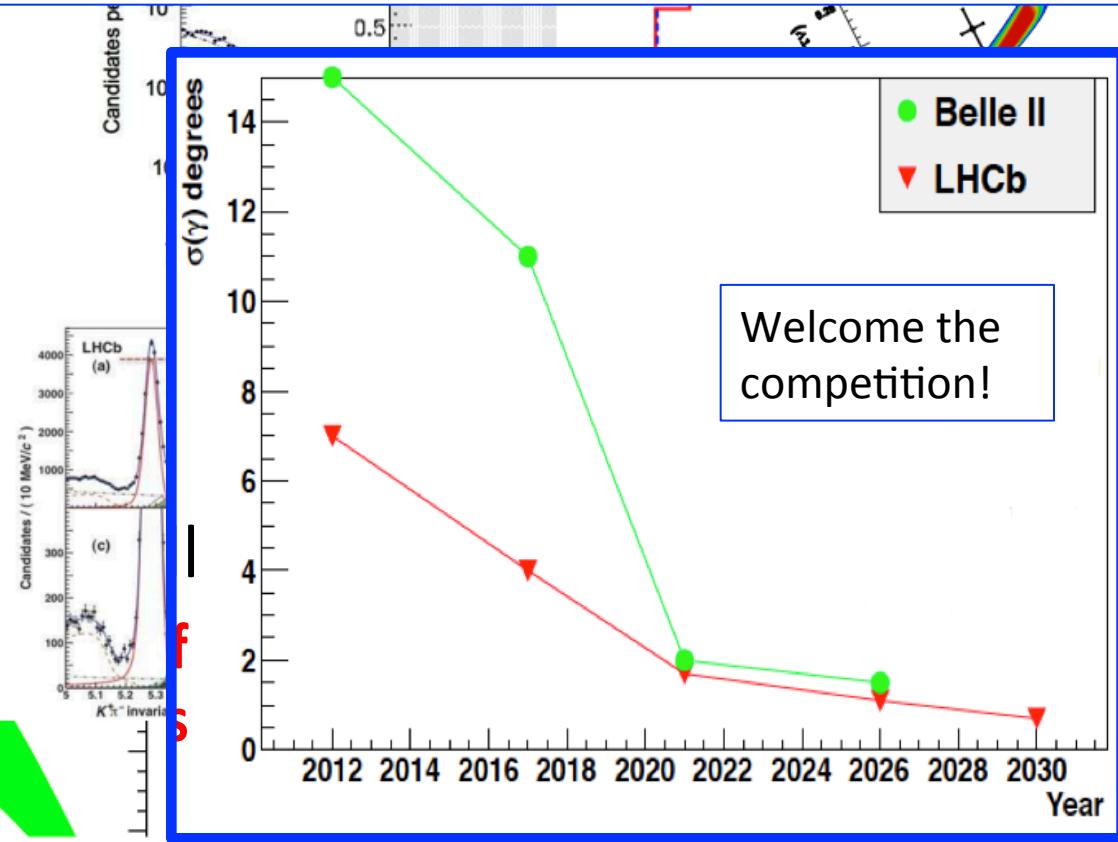
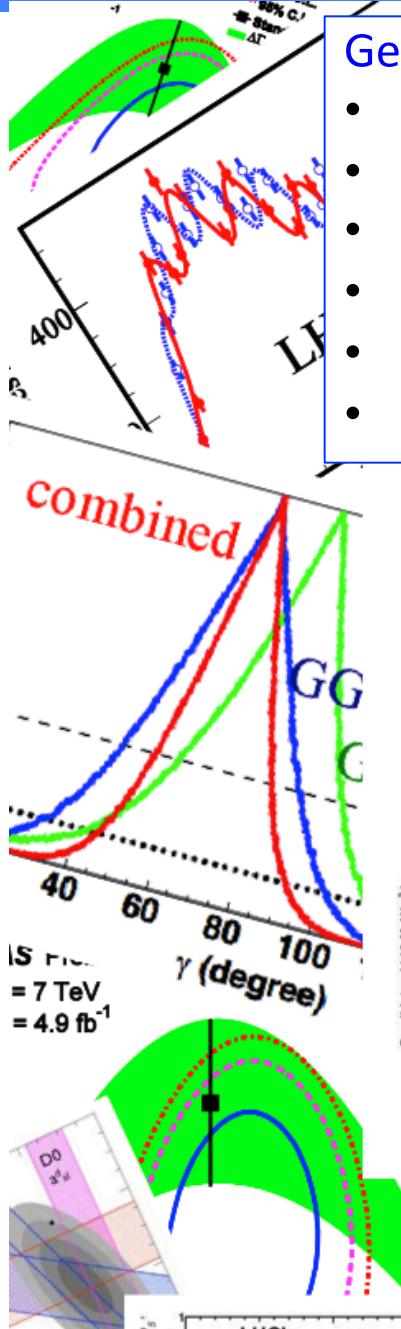
Type	Observable	LHC Run 1	before upgrade	after upgrade	Theory uncertainty
			LHCb 2018	Upgrade 50/fb	
B_s^0 mixing	$\phi_s(B_s^0 \rightarrow J/\psi \phi)$ (rad)	0.049	0.025	0.009	~ 0.003
	$\phi_s(B_s^0 \rightarrow J/\psi f_0(980))$ (rad)	0.068	0.035	0.012	~ 0.01
	$A_{\text{sl}}(B_s^0) (10^{-3})$	2.8	1.4	0.5	0.03
Gluonic penguin	$\phi_s^{\text{eff}}(B_s^0 \rightarrow \phi\phi)$ (rad)	0.15	0.10	0.018	0.02
	$\phi_s^{\text{eff}}(B_s^0 \rightarrow K^{*0} \bar{K}^{*0})$ (rad)	0.19	0.13	0.023	< 0.02
	$2\beta^{\text{eff}}(B^0 \rightarrow \phi K_S^0)$ (rad)	0.30	0.20	0.036	0.02
Right-handed currents	$\phi_s^{\text{eff}}(B_s^0 \rightarrow \phi\gamma)$ (rad)	0.20	0.13	0.025	< 0.01
	$\tau^{\text{eff}}(B_s^0 \rightarrow \phi\gamma)/\tau_{B_s^0}$	5%	3.2%	0.6%	0.2 %
Electroweak penguin	$S_3(B^0 \rightarrow K^{*0} \mu^+ \mu^-; 1 < q^2 < 6 \text{ GeV}^2/c^4)$	0.04	0.020	0.007	0.02
	$q_0^2 A_{\text{FB}}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)$	10%	5%	1.9%	$\sim 7\%$
	$A_I(K \mu^+ \mu^-; 1 < q^2 < 6 \text{ GeV}^2/c^4)$	0.09	0.05	0.017	~ 0.02
	$\mathcal{B}(B^+ \rightarrow \pi^+ \mu^+ \mu^-)/\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)$	14%	7%	2.4%	$\sim 10\%$
Higgs penguin	$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) (10^{-9})$	1.0	0.5	0.19	0.3
	$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)/\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$	220%	110%	40%	$\sim 5\%$
Unitarity triangle angles	$\gamma(B \rightarrow D^{(*)} K^{(*)})$	7°	4°	0.9°	negligible
	$\gamma(B_s^0 \rightarrow D_s^\mp K^\pm)$	17°	11°	2.0°	negligible
	$\beta(B^0 \rightarrow J/\psi K_S^0)$	1.7°	0.8°	0.31°	negligible
Charm CP violation	$A_\Gamma(D^0 \rightarrow K^+ K^-) (10^{-4})$	3.4	2.2	0.4	—
	$\Delta A_{CP} (10^{-3})$	0.8	0.5	0.1	—

Comparable precision for experiment & theory after upgrade.

LHCb is more than b-physics...

General Purpose detector in the forward region:

- Charm physics
- Electroweak Physics: W, Z
- Exotica & Majorana particles searches
- Spectroscopy (see S. Neubert)
- QCD & Heavy Ion physics (Pb-Pb, p-Pb) (see M. Schmelling)
- Focus on b-physics remains



Thank you for your attention!

