

The study of nuclear structure far from stability

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Introduction

Stable nuclei are qualitatively described by “simple” models

- (semi-empirical) **liquid-drop** model
- (basic) **shell** model

New techniques enable **ab-initio** methods (A -body models)

What happens far from stability ?

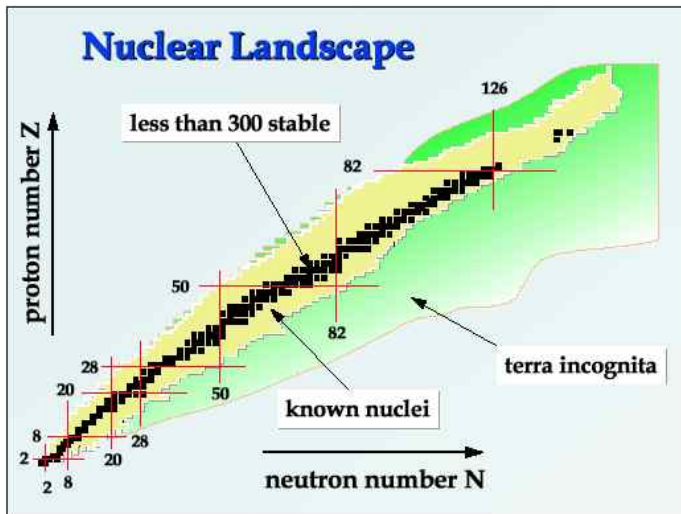
Experimentally, Radioactive-Ion Beams (RIB) available since 80s

⇒ study of structure **far from stability**

⇒ discovery of **exotic** structures

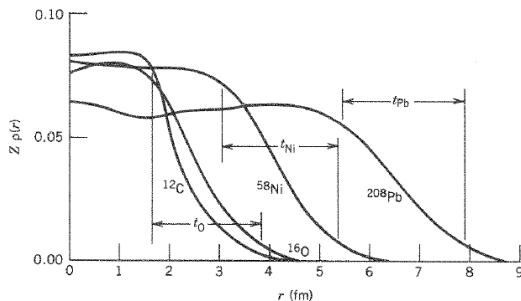
- **super-heavy** elements
- **halo** nuclei
- shell inversions

Nuclear Landscape



- 1 Basic features in nuclear structure
 - Liquid-drop model
 - Shell model
- 2 Ab-initio nuclear models
- 3 Superheavy nuclei
- 4 Radioactive-Ion Beams
- 5 Oddities far from stability : halo nuclei
- 6 Experimental techniques
 - Knockout reactions
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Charge distributions in (stable) nuclei



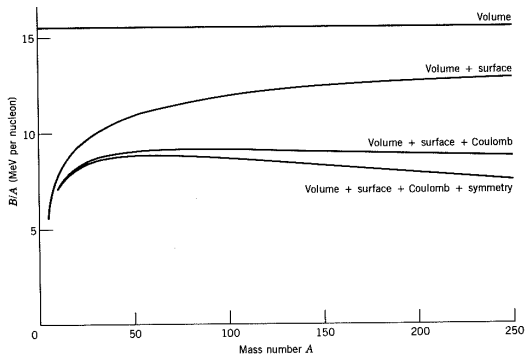
- constant density ρ_0 out to the surface (saturation)
- same skin thickness t

(Stable) nuclei look like **liquid drops** of radius $R \propto A^{1/3}$

Liquid-drop model

Bethe-Weizsäcker semi-empirical mass formula

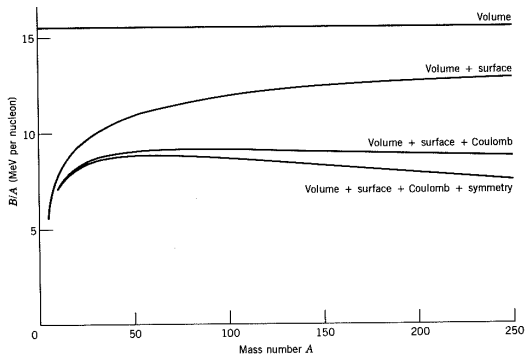
$$B(Z, N) = a_v A$$



Liquid-drop model

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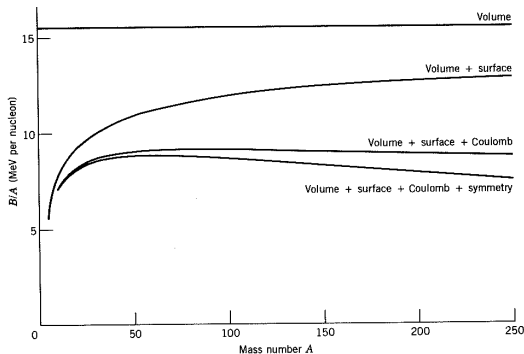
$$B(Z, N) = a_v A - a_s A^{2/3}$$



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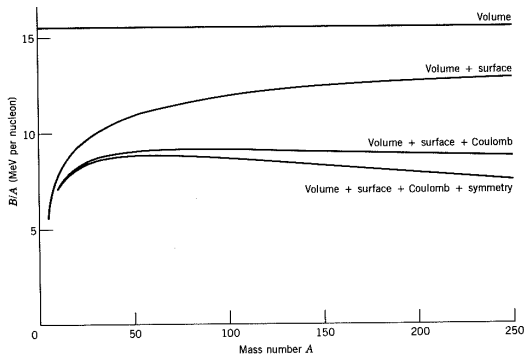
$$B(Z, N) = a_v A - a_s A^{2/3} - a_c \frac{Z(Z-1)}{A^{1/3}}$$



Liquid-drop model

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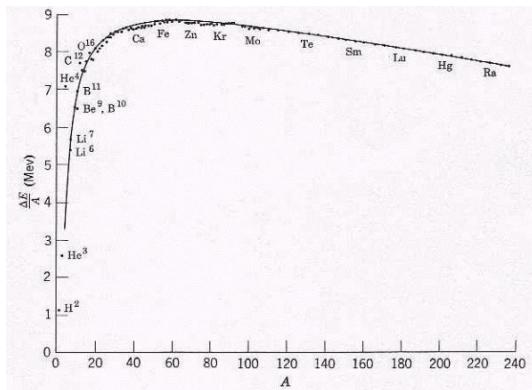
$$B(Z, N) = a_v A - a_s A^{2/3} - a_c \frac{Z(Z-1)}{A^{1/3}} - a_{\text{sym}} \frac{(A-2Z)^2}{A}$$



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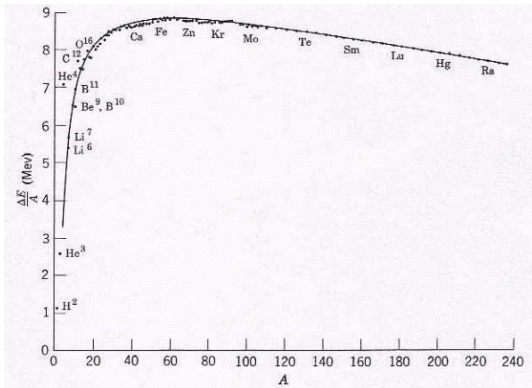
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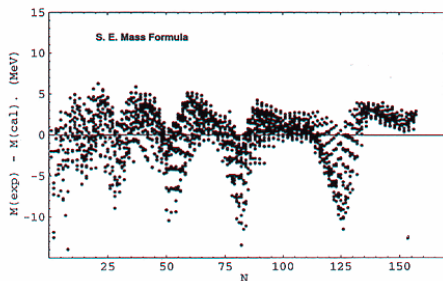
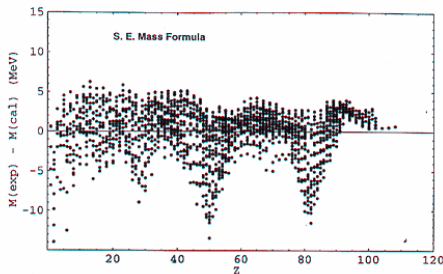
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Exoenergetic reactions :

- **fission** of heavy nuclei
(nuclear power plants,
atomic bomb)
- **fusion** of light nuclei
(stars, thermonuclear
weapons)



Variation from the semi-empirical mass formula



More bound systems at Z or $N = 2, 8, 28, 50, 82, 126$

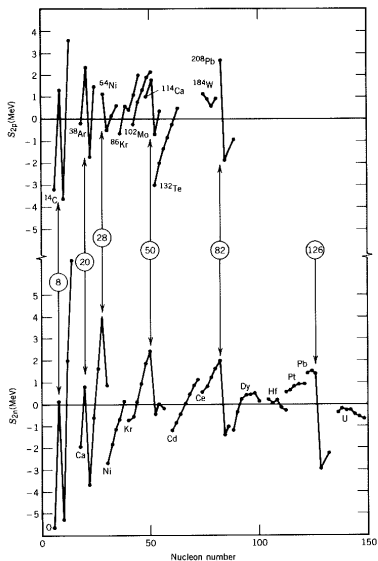
magic numbers

⇒ *shell structure* in nuclei as in atoms?

Two-nucleon separation energy

Same magic numbers in S_{2p} and S_{2n}

⇒ more bound at shell closure
cf. ionisation energies of atoms

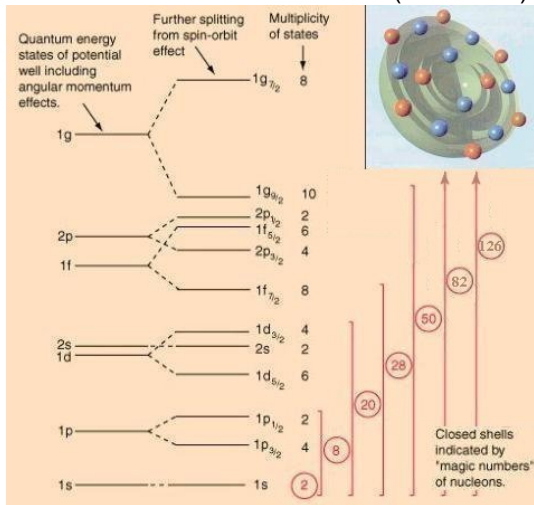


Shell model

Developed in 1949 by M. Goeppert Mayer, H. Jensen and E. Wigner (NP 1963)

As electrons in atoms, nucleons in nuclei feel a **mean field** and arrange into **shells**

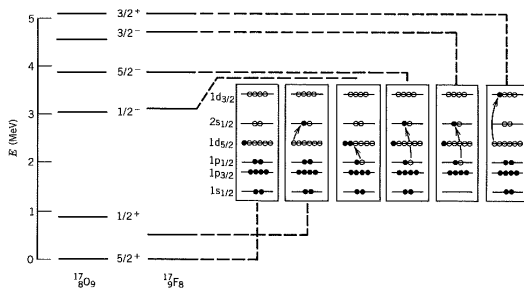
Spin-orbit coupling is crucial to get right ordering of shells



Example

Shell model explains the higher stability at some Z and N

It predicts the spin and parity of ground state of most nuclei and some of their excited levels, e.g. ^{17}O and ^{17}F



Confirmed within coupled-cluster calculation

[G. Hagen *et al.* PRL 104, 182501 (2010)]

Nowadays

Can we go beyond these models ?

Can we build **ab-initio** models ?

i.e. based on first principles

- nucleons as building blocks
- realistic N - N interaction

What happens away from stability ?

- Is there an island of stability for heavy nuclei ?
- Is nuclear density similar for radioactive nuclei ?
- Are magic numbers conserved ?

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A -body Hamiltonian

Nuclear-structure calculations : A **nucleons** (Z protons+ N neutrons)

Relative motion described by the A -body **Hamiltonian**

$$H = \sum_{i=1}^A T_i + \sum_{j>i=1}^A V_{ij}$$

\Rightarrow solve the A -body Schrödinger equation

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle$$

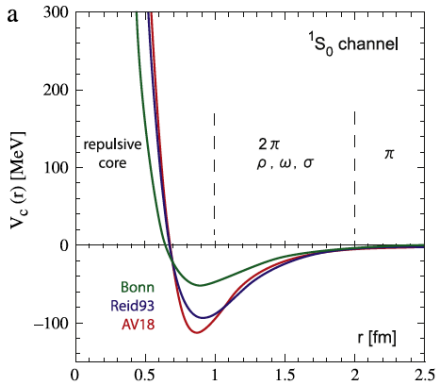
$\{E_n\}$ is the nucleus spectrum

Realistic N - N interactions

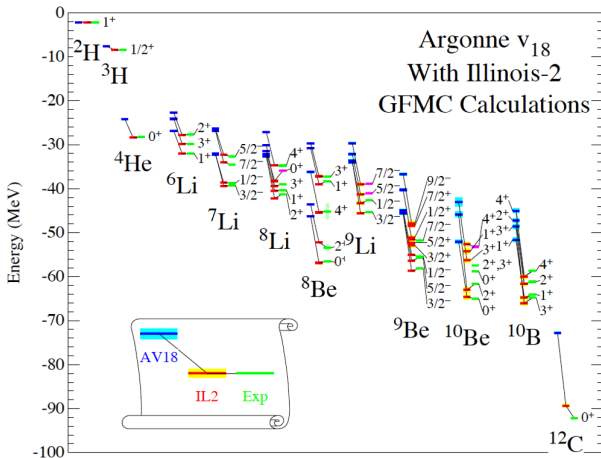
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d binding energy,
 N - N phaseshifts

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 Ex. : Argonne V18, CD-Bonn, ...



Light nuclei calculations



[R. Wiringa, Argonne]

Three-body force

Need three-body forces to get it right. . .

$$H = \sum_{i=1}^A T_i + \sum_{j>i=1}^A V_{ij} + \sum_{k>j>i=1}^A V_{ijk} + \dots$$

But there is no such thing as three-body force. . .

Three-body force

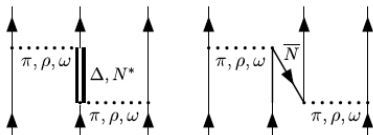
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But there is no such thing as three-body force. . .

They simulate the **non-elementary** character of nucleons

⇒ include virtual Δ resonances, \bar{N} . . .



Phenomenological 3-body interaction fitted on $A > 2$ levels : IL2
Alternatively, derived from EFT

Effective Field Theory

EFT is an **effective** quantum field theory based on QCD symmetries with resolution scale Λ that selects appropriate degrees of freedom : nuclear physics is not built on quarks and gluons, but on **nucleons** and **mesons**

EFT provides the nuclear force with a systematic expansion in Q/Λ

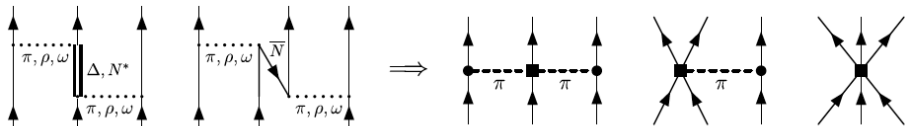
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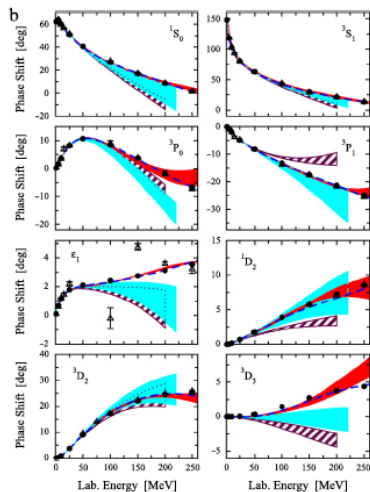
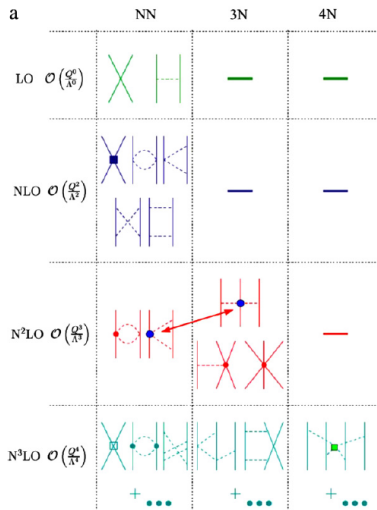
EFT provides the nuclear force with a systematic expansion in Q/Λ

- gives an estimate of theoretical uncertainty
- naturally includes many-body forces



[see G. Hagen's talk on Wednesday morning
and J. Holt on Tuesday afternoon]

Expansion of the EFT force



Solving the Schrödinger equation

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle$$

Ψ usually developed on a **basis** $\{|\Phi_{[v]}\rangle\}$:

$$|\Psi_n\rangle = \sum_{[v]} \langle \Phi_{[v]} | \Psi_n \rangle |\Phi_{[v]}\rangle$$

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Solving the Schrödinger equation reduces to matrix diagonalisation

$$\begin{aligned} \langle \Phi_{[\mu]} | H | \Psi_n \rangle &= \sum_{[\nu]} \langle \Phi_{[\mu]} | H | \Phi_{[\nu]} \rangle \langle \Phi_{[\nu]} | \Psi_n \rangle \\ &= E_n \langle \Phi_{[\mu]} | \Psi_n \rangle \end{aligned}$$

No-core shell model

Slater determinants of 1-body mean-field wave functions ϕ_{ν_i}

$$\langle \xi_1 \xi_2 \dots \xi_A | \Phi_{[\nu]} \rangle = \mathcal{A} \phi_{\nu_1}(\xi_1) \phi_{\nu_2}(\xi_2) \dots \phi_{\nu_A}(\xi_A)$$

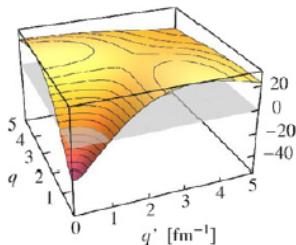
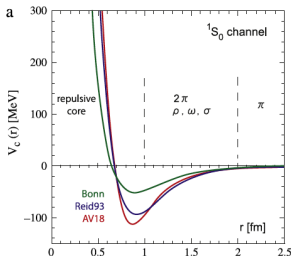
But short-range correlations couple low and high momenta

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But short-range correlations couple low and high momenta



\Rightarrow requires large basis $|\Phi_{[v]}\rangle$

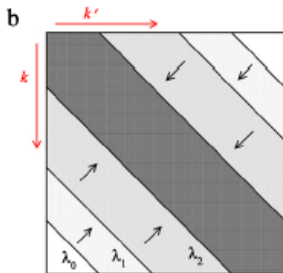
Similarity Renormalisation Group

Idea : apply a unitary transformation

$$\begin{aligned} |\widetilde{\Phi}_{[v]}\rangle &= U|\Phi_{[v]}\rangle \\ \Leftrightarrow H_{\text{eff}} &= U^\dagger H U \end{aligned}$$

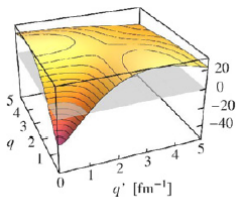
- keeps the same spectrum (unitary)
- keeps the same on-shell properties (phaseshifts)
- removes the short-range correlations

This has a costs : induces “unphysical” three-body forces

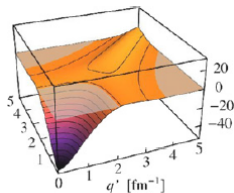


SRG : example on ^4He

SRG lowers correlations



(a) V_{AV18} .

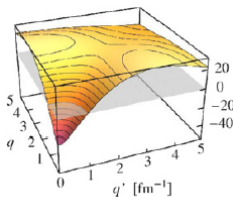


(b) V_{SRG} .

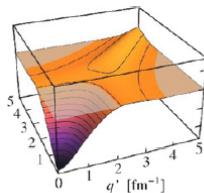
[see G. Hagen's talk on Wednesday morning
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SRG : example on ${}^4\text{He}$

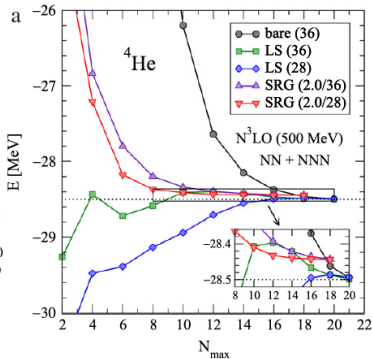
SRG lowers correlations \Rightarrow fastens convergence



(a) V_{AV18} .



(b) V_{SRG} .



[see G. Hagen's talk on Wednesday morning
and J. Holt on Tuesday afternoon]

What happens far from stability ?

Liquid-drop and shell models are fair models of stable nuclei

What happens away from stability ?

Are there **superheavy** nuclei ? Recent discovery of new elements
($Z = 113, 115, 117$, and 118)

In 80s Radioactive-Ion Beams were developed

Enable study of nuclear **structure** [see A. Obertelli's talk on Friday]

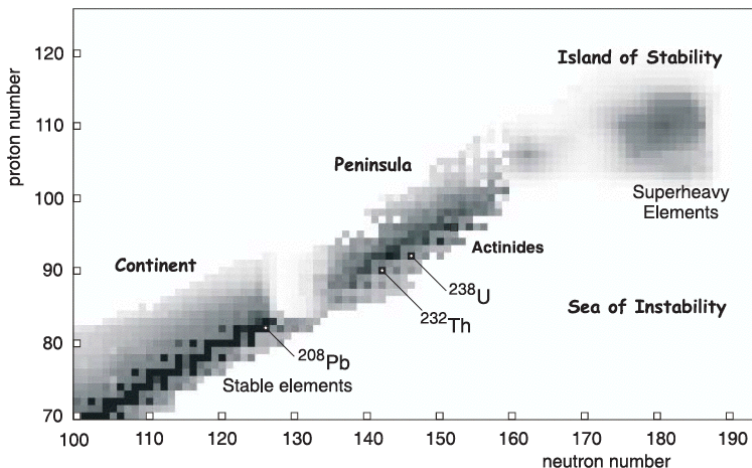
- are radioactive nuclei compact ?
- are shells conserved far from stability ?

Reactions involving radioactive nuclei useful in **astrophysics**

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Superheavy nuclei

Does the stability end with U? Or is there an **island of stability**?
Is $Z \sim 114 - 126$ a new magic number?



[Oganessian Rad. Act. 99, 429 (2011)]

Synthesis of Super Heavy Elements

Search elements heavier than U has started in the 40's

Pu produced by $U+d$ and $U+n$ (identified by Seaborg in 1941)

High flux of n on U or Pu \rightarrow ${}_{100}\text{Fm}$

$Z = 101-106$ obtained by bombarding actinides by light elements

Heavier elements ($Z = 107-113$) by **cold fusion**

${}^{208}\text{Pb}$ or ${}^{209}\text{Bi}$ + massive projectile ($A \geq 50$) \rightarrow high Z + n

Nowadays, **hot fusion** : ${}^{48}\text{Ca}$ on actinide target \rightarrow SHE + 4–5 n

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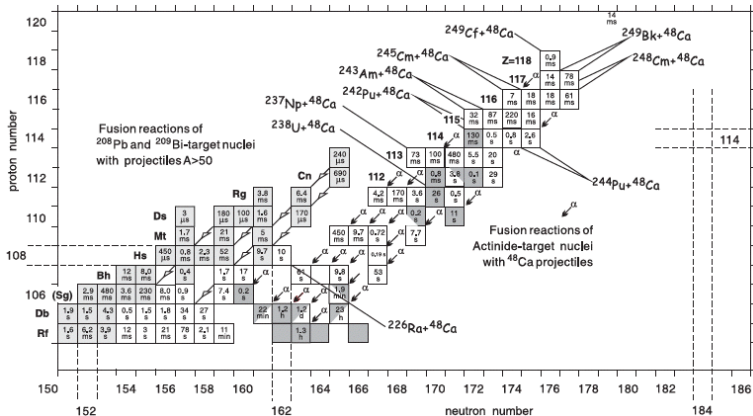
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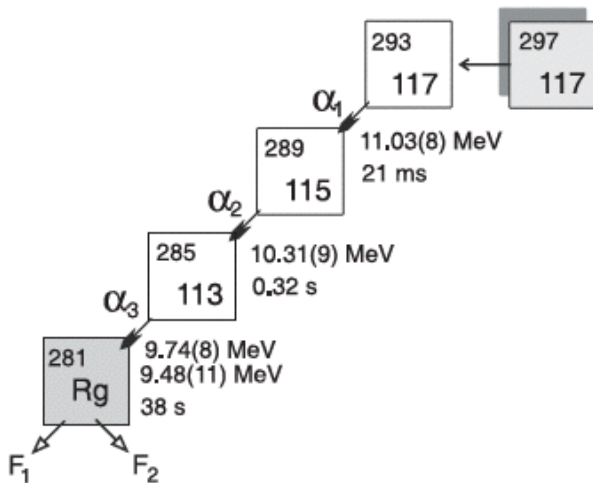
Synthesis of Super Heavy Elements

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Z = 113, 115, 117, and 118 recently recognised by IUPAC



[Oganessian Rad. Act. 99, 429 (2011)]

Identification by α decay chain



[Oganessian Rad. Act. 99, 429 (2011)]

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- ISOL : Fire a proton at a heavy nucleus



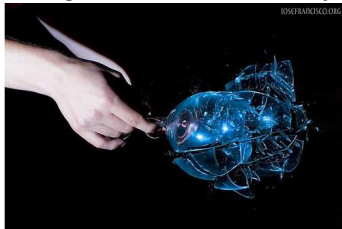
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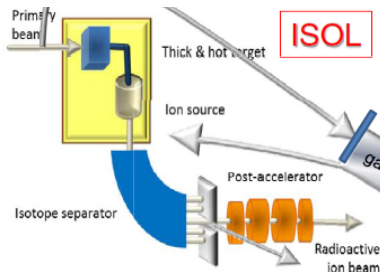
- In-flight : Smash a heavy nucleus on a target



Where ?

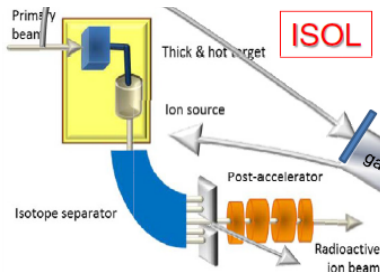


ISOL : Isotope Separation On Line



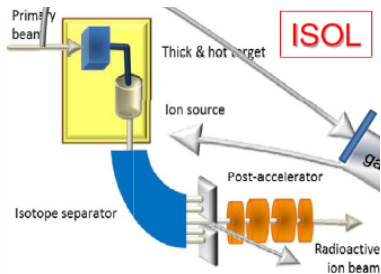
high-energy/intensity **primary beam** of light nuclei (e.g. protons)
 on **thick target** of heavy elements (Ta or UC_x)
 \Rightarrow spallation/fragmentation produces **exotic fragments**

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Diffuse in the target and effuse to an **ion source**

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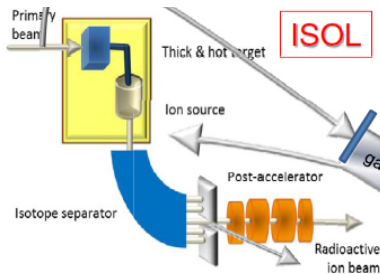
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Then **selected** using dipole magnet (A/Q)

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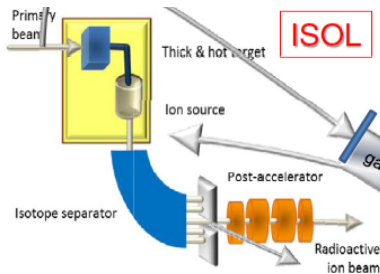
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or post-accelerated for reactions (e.g. astrophysical energy)

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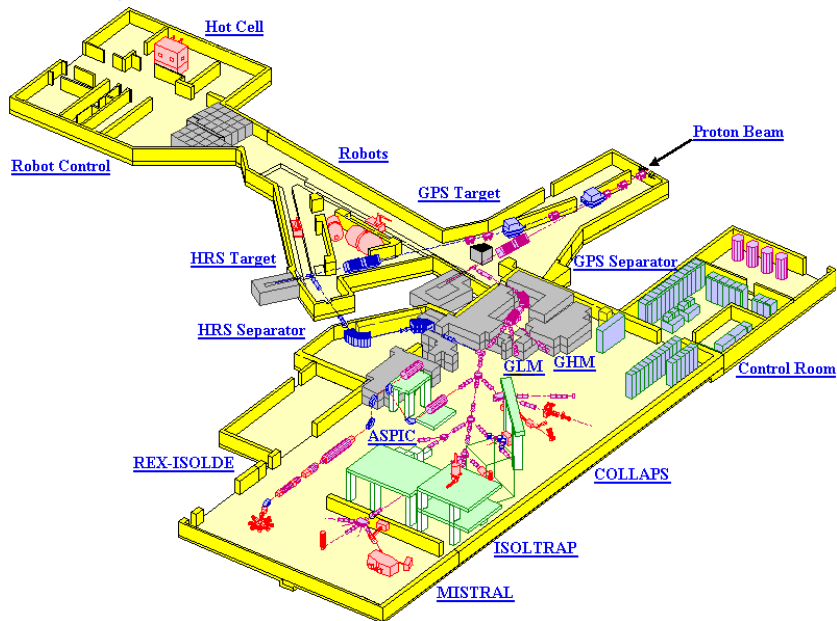
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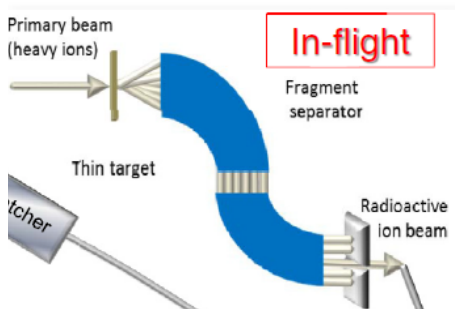
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Examples : ISOLDE (CERN), TRIUMF, SPIRAL (GANIL)

ISOLDE @ CERN



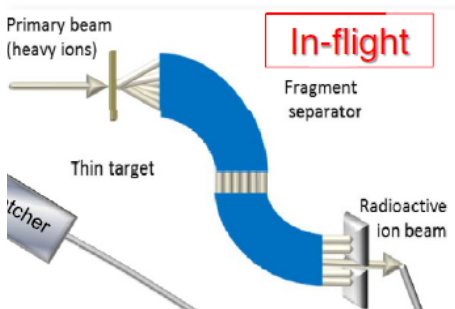
In-flight projectile fragmentation



high-energy **primary beam** of heavy ions (e.g. ^{18}O , ^{48}Ca , U...)
on **thin** target of light elements (Be or C)

⇒ fragmentation/fission produces many **exotic fragments** at $\approx v_{\text{beam}}$

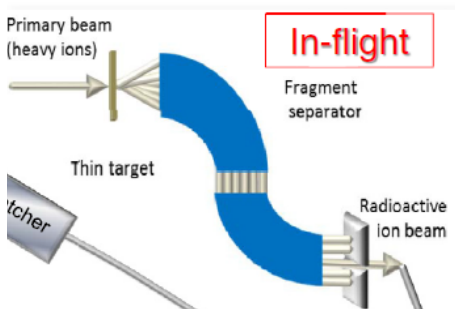
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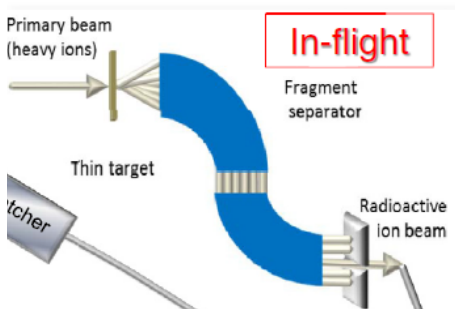
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Used for **high-energy** reactions (KO, breakup...)

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Used for **high-energy** reactions (KO, breakup...)

Examples : NSCL (MSU), GSI, RIKEN, GANIL

Existing NSCL



3a K500
Cyclotron



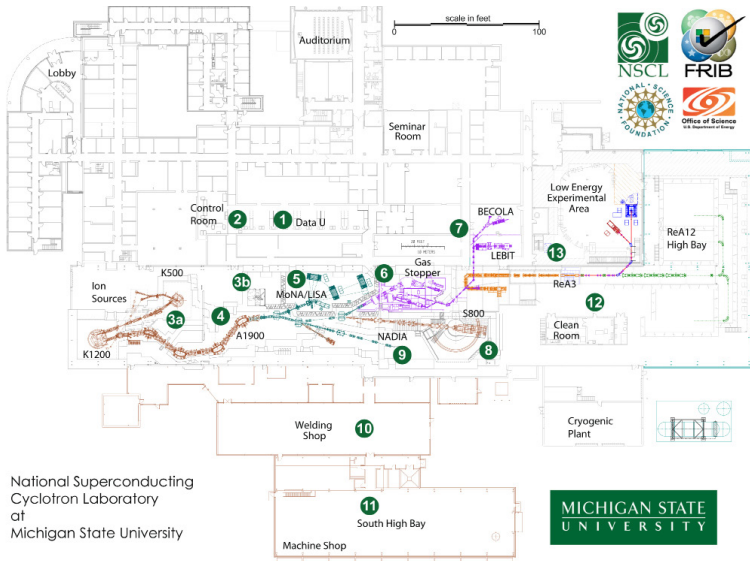
4 A1900
Fragment Separator



5 Modular
Neutron Array



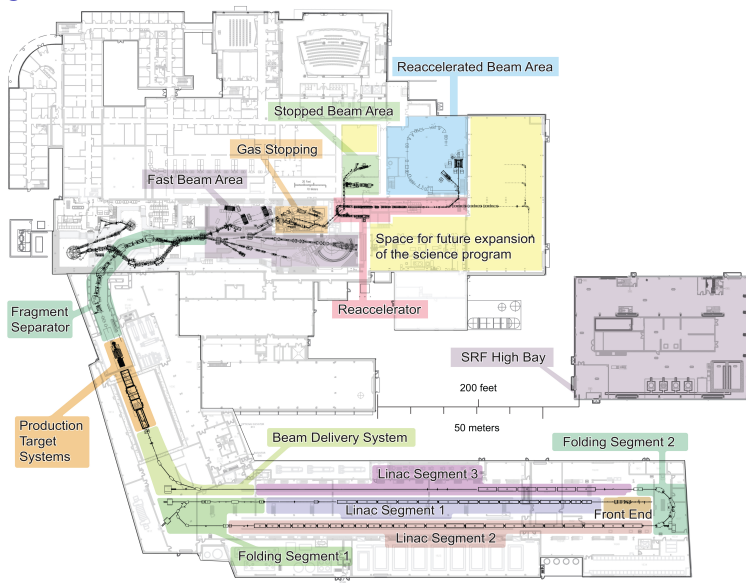
8 S800
Spectrograph



National Superconducting
Cyclotron Laboratory
at
Michigan State University



Future : FRIB



Properties

ISOL

- Low beam energy
may require post-acceleration
- **Low** beam intensity
- Not all elements produced
 - ▶ Slow
 - ▶ Chemically limited
- **Good** beam quality :
can use chemistry to select
fragments

Properties

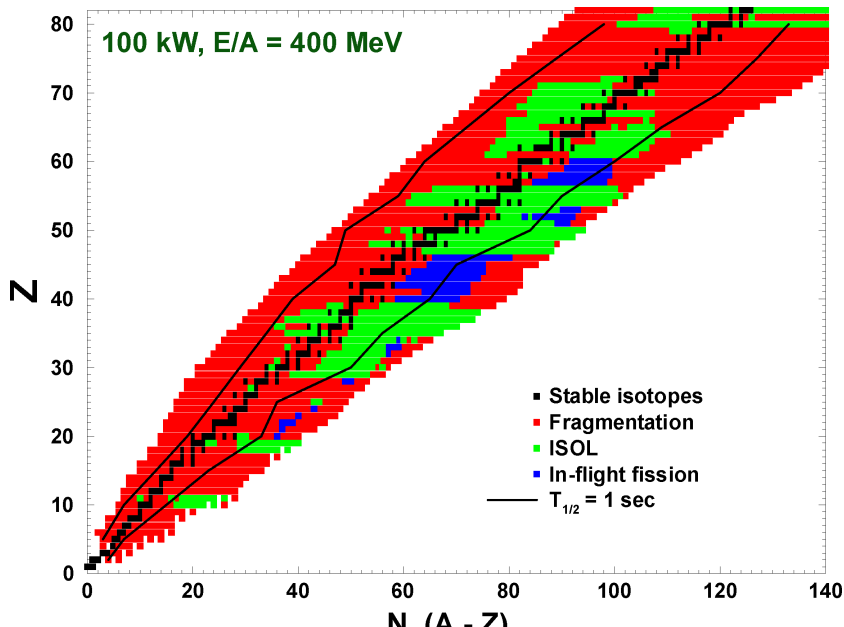
ISOL

- Low beam energy
may require post-acceleration
- **Low** beam intensity
- Not all elements produced
 - ▶ Slow
 - ▶ Chemically limited
- **Good** beam quality :
can use chemistry to select
fragments

In-flight

- High beam energy
 $v_{\text{fragments}} \approx v_{\text{beam}}$
- **High** beam intensity
- **Efficient** production
 - ▶ Fast
 - ▶ Chemically independent
- Many **fragments** in beam
⇒ need ion ID

Choose according what you want to measure



- 1 Basic features in nuclear structure
 - Liquid-drop model
 - Shell model
- 2 Ab-initio nuclear models
- 3 Superheavy nuclei
- 4 Radioactive-Ion Beams
- 5 **Oddities far from stability : halo nuclei**
- 6 Experimental techniques
 - Knockout reactions
 - In-beam γ spectroscopy
 - Transfer reactions
- 7 Summary

Halo structure

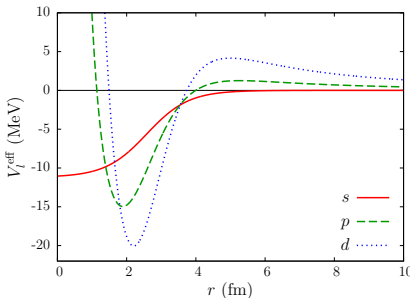
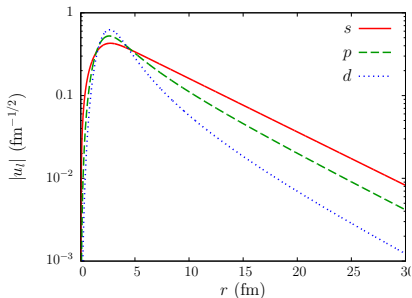
Seen as **core** + one or two **neutrons** at large distance

[P. G. Hansen and B. Jonson, Europhys. Lett. 4, 409 (1987)]

Peculiar structure of nuclei due to small S_n or S_{2n}

⇒ neutrons tunnel far from the **core** to form a **halo**

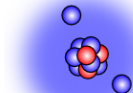
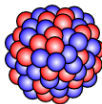
Halo only appears for low centrifugal barrier (low ℓ)



see F. Colomer's poster on Monday

Halo nuclei

- Light, **neutron-rich** nuclei
- small S_n or S_{2n}
- low- ℓ orbital

 ^{11}Li  ^{208}Pb 

One-neutron halo

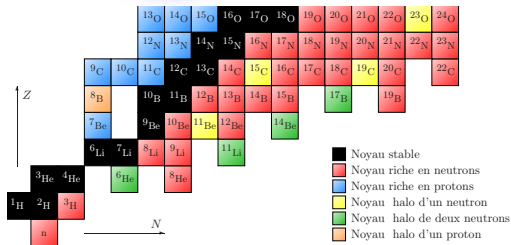
$$^{11}\text{Be} \equiv ^{10}\text{Be} + n$$

$$^{15}\text{C} \equiv ^{14}\text{C} + n$$

Two-neutron halo

$$^6\text{He} \equiv ^4\text{He} + n + n$$

$$^{11}\text{Li} \equiv ^9\text{Li} + n + n$$

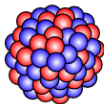


Proton haloes are possible but less probable : ^8B , ^{17}F

Halo nuclei

- Light, **neutron-rich** nuclei
- small S_n or S_{2n}
- low- ℓ orbital

 ^{11}Li

 ^{208}Pb


One-neutron halo

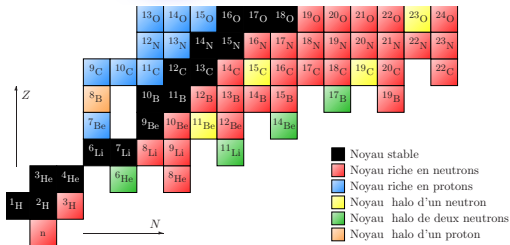
$$^{11}\text{Be} \equiv ^{10}\text{Be} + n$$

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Two-neutron halo

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Proton haloes are possible but less probable : ^8B , ^{17}F

Two-neutron halo nuclei are **Borromean**...

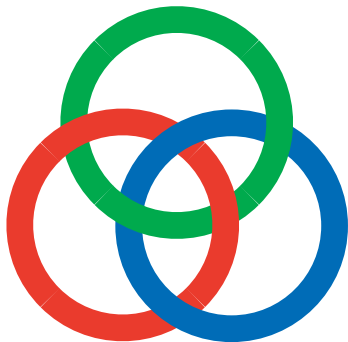
$c+n+n$ is bound but not two-body subsystems

e.g. ^6He bound but not ^5He or 2n

Borromean nuclei

Named after the Borromean rings. . .

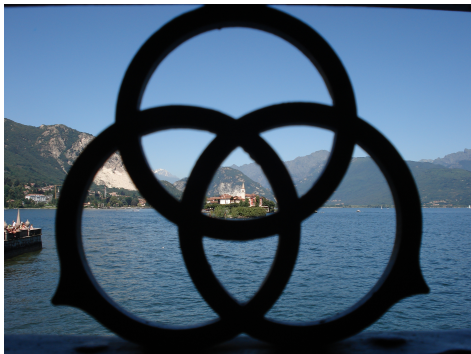
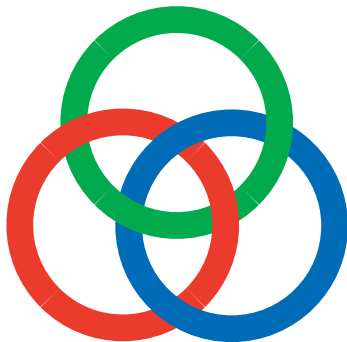
[M. V. Zhukov *et al.* Phys. Rep. 231, 151 (1993)]



Borromean nuclei

Named after the Borromean rings. . .

[M. V. Zhukov *et al.* Phys. Rep. 231, 151 (1993)]



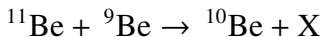
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Knockout reactions

Knockout reaction : direct removal of one nucleon

For **exotic** nuclei, measured in **inverse kinematics**

on light target (C or Be) at intermediate energy (70 A MeV)



Inclusive measurement (KO nucleon not measured)

Used to probe the shell-model wave function of nuclei

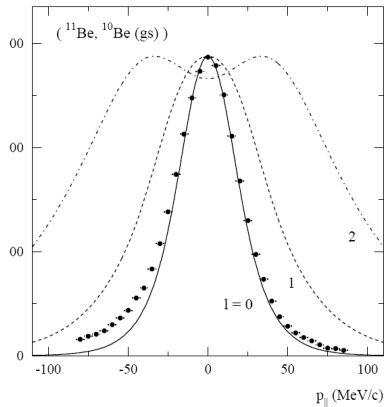
The **valence nucleon** occupies a shell outside a **core** :

$$^A Y(J^\pi) = {}^{A-1} X(J_c^\pi) \otimes \psi_{lm} + \dots$$

Analysis of reaction within **eikonal** model (sudden approximation)

- **Parallel-momentum** distribution of the core gives nucleon shell
- Total σ_{-N} provides **Spectroscopic Factor** (SF)

$^9\text{Be}(^{11}\text{Be}, ^{10}\text{Be} + \gamma)\text{X} @ 60\text{A MeV}$



[Aumann *et al.* PRL 84, 35 (2000)]

- Shell model predicts

$$^{11}\text{Be} \equiv ^{10}\text{Be}(0^+) \otimes \nu_{0p1/2}$$

- Experiment shows **shell inversion**

$$^{11}\text{Be} \equiv ^{10}\text{Be}(0^+) \otimes \nu_{1s1/2}$$

- **Narrow** momentum distribution
 \Leftrightarrow **extended** spacial distribution
 \Rightarrow confirms the neutron **halo**

Study of nuclear spectroscopy

Reaction models rely on **single-particle** model :

$$[T_r + V_{cf}(r) - \epsilon]\phi_{nlm}(\mathbf{r}) = 0 \quad \text{with } \|\phi_{nlm}\| = 1$$

In reality, there is admixture of configurations :

$${}^A Y(J^\pi) = {}^{A-1} X(J_c^\pi) \otimes \psi_{lm} + \dots$$

where ψ_{lm} is the **overlap wave** function

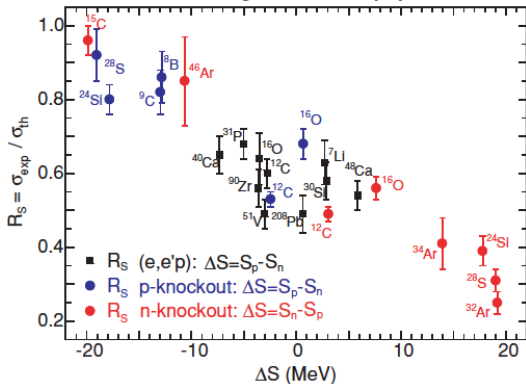
Spectroscopic Factor : $\mathcal{S}_l = \|\psi_{lm}\|^2$

Single-particle approximation $\equiv \psi_{lm} = \sqrt{\mathcal{S}_l} \phi_{nlm}$

\Rightarrow usual idea : $\mathcal{S}_l = \sigma_{\text{bu}}^{\text{exp}} / \sigma_{\text{bu}}^{\text{th}}$

Reduction of the SF

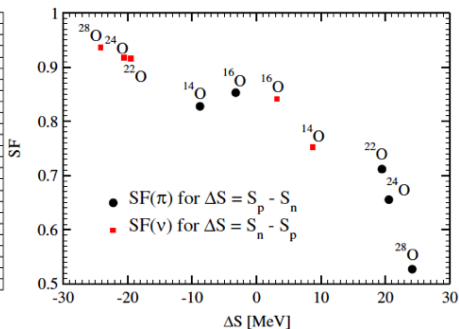
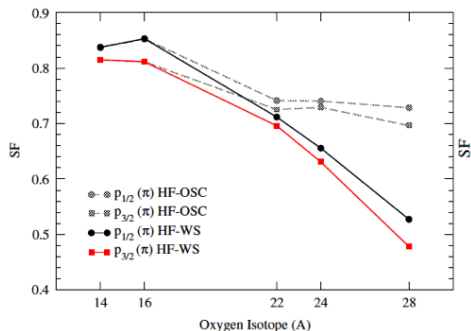
Predicted shell-model SF too high for deeply-bound nucleons



[Gade *et al.* PRC 77, 044306 (2008)]

- Something missing in usual shell-model calculations ?
- Problem in the reaction modelling ?

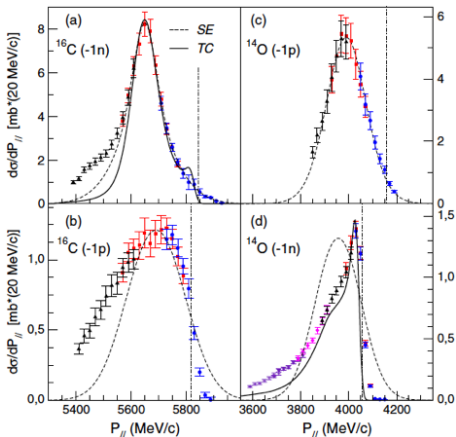
Quenching of SF



[Jensen *et al.* PRL 107, 032501 (2011)] ; see also G. Hagen's talk on Wednesday
 Inclusion of **continuum** in shell-model basis **reduces** SF
 \Rightarrow solves partly the problem

KO reaction dynamics

- For loosely-bound nucleons
 $S_n(^{16}\text{C}) = 4.25 \text{ MeV}$
 $S_p(^{14}\text{O}) = 4.63 \text{ MeV}$
 eikonal model works fine
- For deeply-bound nucleons
 $S_p(^{16}\text{C}) = 22.6 \text{ MeV}$
 $S_n(^{14}\text{O}) = 23.2 \text{ MeV}$
 energy conservation must be taken into account



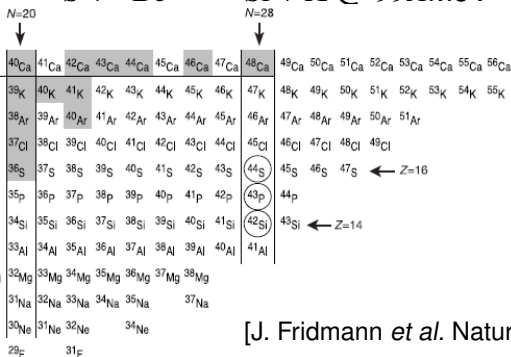
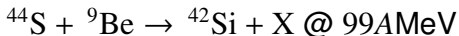
[Flavigny et al. PRL 108, 252501 (2012)]

see A. Obertelli's talk on Friday

⇒ solves another part of the problem ?

New magic numbers ?

New magic $Z = 14$ in ^{42}Si ?



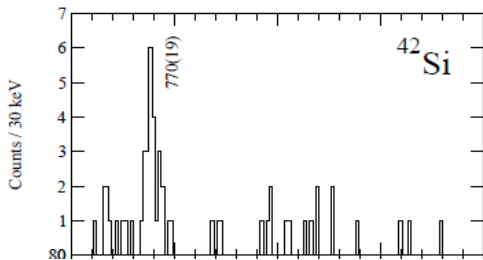
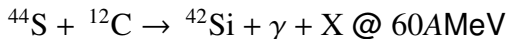
[J. Fridmann *et al.* Nature 435, 922 (2005)]

Analysis within shell-model calculation suggests $Z = 14$ to be magic far from stability...

However...

Measurement of γ in coincidence with proton stripping

see A. Obertelli's and L. Atar's talks on Friday

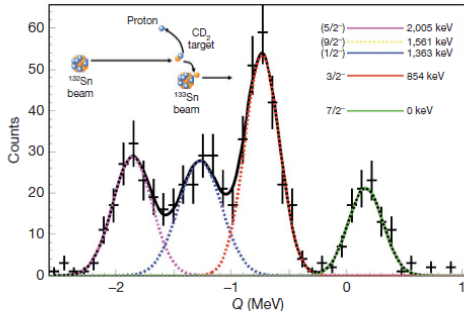
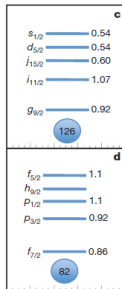
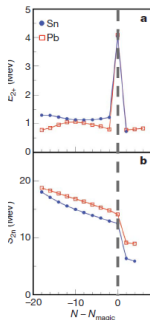
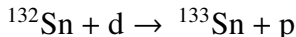


[B. Bastin *et al.* PRL 99, 022503 (2007)]

Indicates a low-lying state in $^{42}\text{Si} \Rightarrow Z = 14$ not magic
and collapse of $N = 28$ far from stability

Confirmed magic numbers

Transfer reaction can be used to study shell structure



[K. Jones *et al.* Nature 465, 454 (2010)]

Confirms the single-particle structure of ^{133}Sn
and hence that ^{132}Sn is magic (i.e. $N = 82$ at $Z = 50$)

Summary

Liquid-drop and shell model describe qualitatively stable nuclei
Nowadays **ab-initio** nuclear-structure models from first principles

RIB enable study nuclear structure far from stability

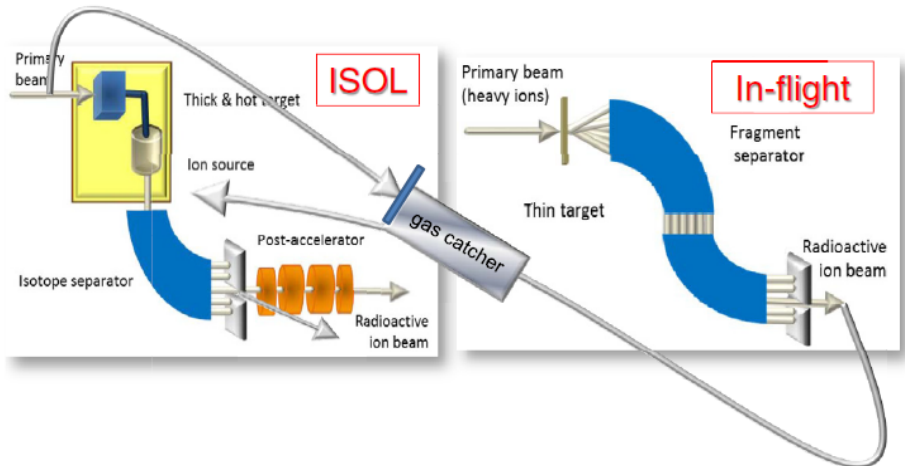
Low intensities require new experimental techniques :

KO reactions, in-beam γ spectroscopy,...

- discovery of **halo** nuclei
diffuse halo around a compact core
- **shell inversions** or **shell collapse**

RIB can be used to study reactions of **astrophysical** interest

Combined with a gas stopper



- can use thin target in ISOL
- can study low-energy reaction with in-flight fragmentation

C($^{31}\text{Ne}, ^{30}\text{Ne} + \gamma$)X @ 230A MeV

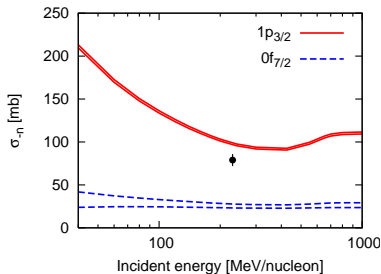
Shell model predicts ^{31}Ne ground state to be $7/2^-$ ($^{30}\text{Ne} \otimes \nu_{0f7/2}$)

One-neutron knock-out measured at RIKEN

$$^{31}\text{Ne} + \text{C} \rightarrow ^{30}\text{Ne} + \text{X} @ 230\text{A MeV}$$

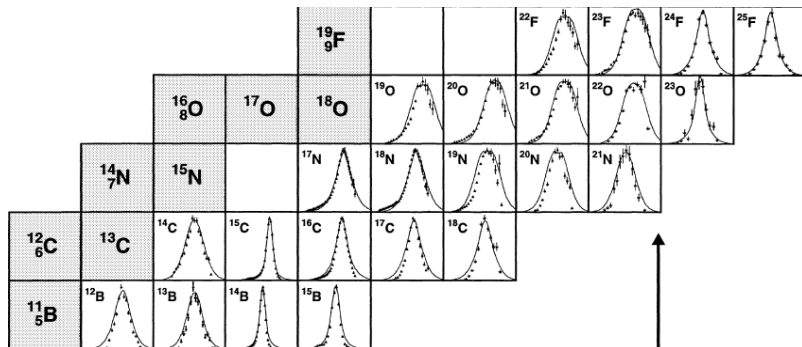
[T. Nakamura *et al.* PRL 103, 262501 (2009)]

Theoretical analysis suggests $3/2^-$ ($^{30}\text{Ne} \otimes \nu_{p3/2}$)



[W. Horiuchi *et al.* PRC 81, 024606 (2010)]

Systematic study



[Sauvan *et al.* PLB 491, 1 (2000)]