

Part 2: Jet Algorithms

Last time:

Clusters of
partons

$\sim$
 $\sim$
 $\uparrow$
up to Λ_{QCD}
effects

Clusters of
hadrons

Jets are proxies for underlying quarks / gluons
Essential for new physics searches
Fundamentally ambiguous, but opportunity to optimize.

Today: Jet Algorithms as one way to study QCD processes
(Not only way! E.g. event shapes, probabilistic methods.)
Offers unique event interpretation.

Generic Jet Algorithm

$$\{p_1, p_2, p_3 \dots p_n\}_{\text{hadrons}} \Rightarrow \{p_1, p_2 \dots p_N\}_{\text{jets}}$$

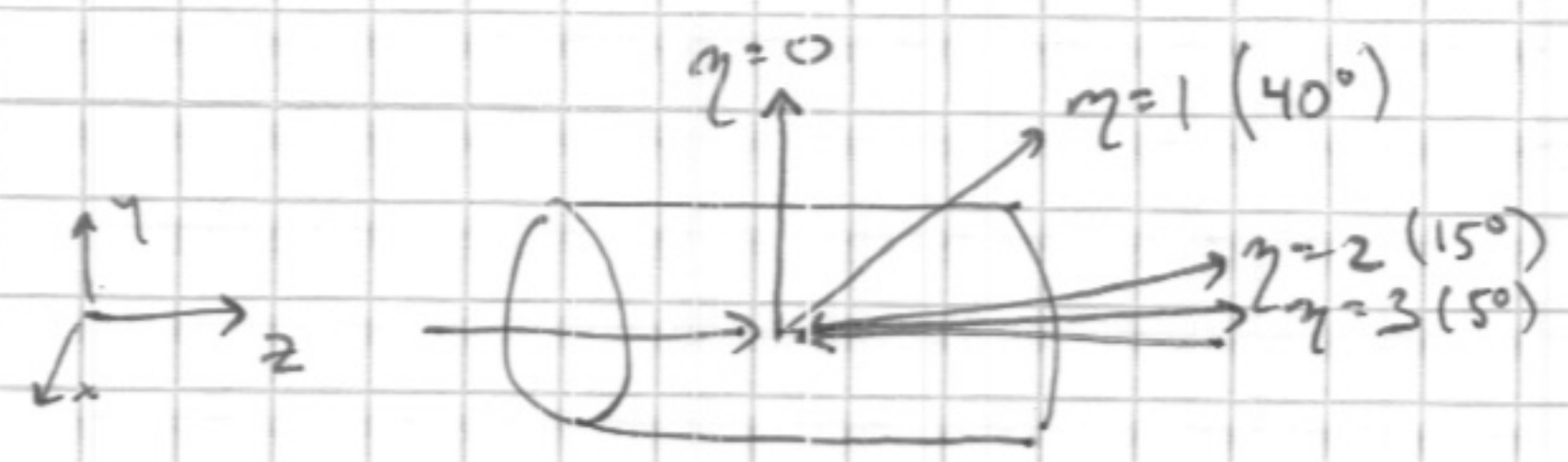
Usually: $p_{\text{jet}} = \sum_{i \in \text{jet}} p_i$ (E-scheme recombination)

Note: $p_{\text{jet}}^2 = \left(\sum_i p_i \right)^2 \geq \sum_i m_i^2$

Jets are full fourvectors with non-zero mass.

Mass has to be included for correct resonance reconstruction. (e.g. $Z' \rightarrow \text{dijets}$)

Coordinate system for hadron colliders.



Typically:
longitudinally
boost invariance

$$p_T = \sqrt{p_x^2 + p_y^2}$$

$$\phi = \arctan \frac{p_y}{p_x}$$

$$\eta = \operatorname{arctanh} \frac{p_z}{|\vec{p}|}$$

$$y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z}$$

↑ true rapidity

pseudorapidity
(only boost invariant
difference for massless particles)
 $= -\ln \tan \frac{\theta}{2}$

(longitudinal)

Boost-invariant distances

$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$$

Jets from cone algorithms?



Punch holes of radius R.

Immediate challenges:
Which holes?
Split/Merge ambiguity

Cone algorithms are intuitive: flow of energy in
a given direction within a given angle

Subtleties encountered in Tevatron era
(e.g. IRC unsafe seeds)

Should probably be revisited (more later)

What is ΔR ? (Beyond being just boost invariant.)

(16)

Standard opening angle: $d\Omega^2 = d\Theta^2 + \sin^2\Theta d\phi^2$

Exercise: Using $\sin\Theta = \frac{1}{\cosh\eta}$ for massless particles

$$\text{Show: } d\Omega^2 = \frac{1}{\cosh^2\eta} \underbrace{(d\eta^2 + d\phi^2)}_{\approx dR^2}$$

So ΔR is just like opening angle, rescaled by factors of $\cosh\eta$. ($\cosh(\eta=0) = 1$)

For massless particles: $E = p_T \cosh\eta$

$$\Rightarrow E \Delta\Omega \approx p_T \Delta R \quad (\text{nice thing to remember})$$

$\Rightarrow$ Jets are smaller in forward ($|\eta|$) direction.

Q: Why $p_T, \Delta R$ instead of $E, \Delta\Omega$?

(17)

longitudinally
boost invariant

used extensively
in e^+e^- colliders

Ok answer: PDFs mean we don't know longitudinal momentum fraction of colliding quark/gluons.

Better answer: Rapidity Plateau: in "minimum bias" collisions, number of particles roughly constant as a function of η .

Q: Why fixed R ?

Actually, there is interest in variable R , dynamical R , multiple R ...

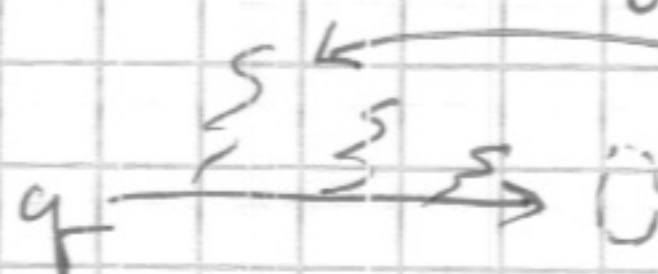
Right now, $R=0.4$ anti- k_T jets very standard
 $\hookrightarrow$ ($\sim 22^\circ$ half angle, $\sim 45^\circ$ full angle)

vs. 0.1×0.1 HCAL granularity

Fixed $R \Rightarrow$ roughly uniform jet contamination as a function of η .

Q: Why not very small R ?

Perturbative $\log R$ resummation needed
 Sensitivity to hadronization, though

 "out of cone" corrections.

The way I think of jet algorithms: Optimization Problem

What do you want to optimize? (many choices...)

→ Hard partons should be clustered into jets

→ Jet axis $\approx$ jet momentum direction

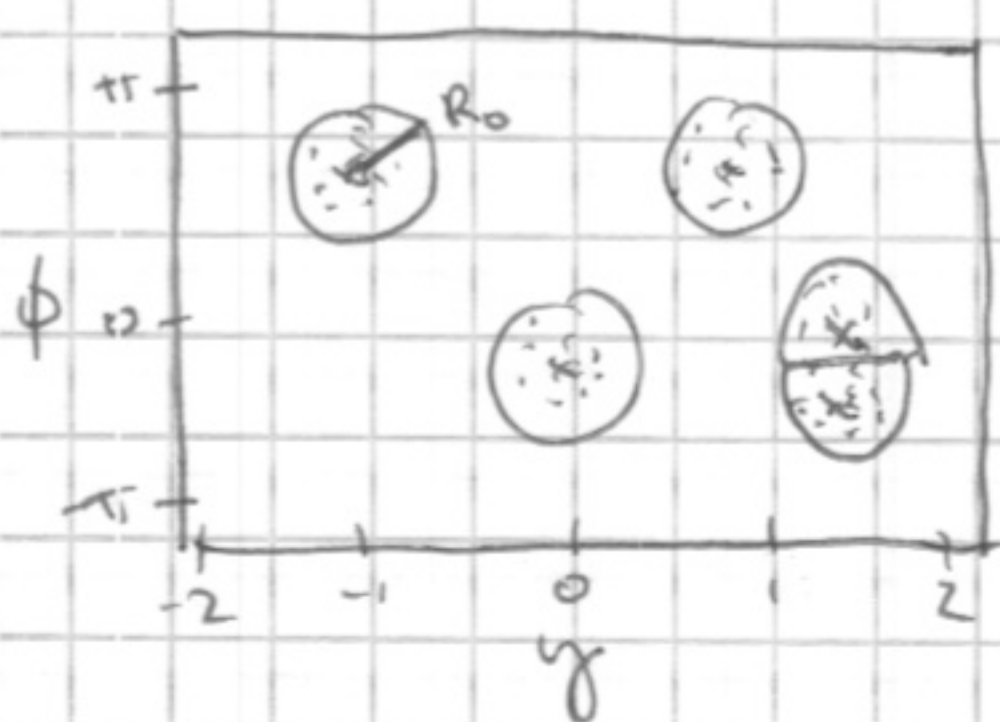
→ Roughly equal area (for contamination)

All of these can / should be revisited.

E.g. I want exactly N jets with axes: $\hat{n}_1, \hat{n}_2, \dots, \hat{n}_N$

Optimize this: $\mathcal{Z}_N = \sum_{i \in \text{event}} \min \left\{ p_{T,i}, p_{T,i} \left(\frac{\Delta R_{i1}}{R_0} \right)^2, \dots, p_{T,i} \left(\frac{\Delta R_{iN}}{R_0} \right)^2 \right\}$

$\uparrow$ $\uparrow$ $\uparrow$ $\uparrow$
 N -jetness penalty function not in a jet in first jet in N -th jet



Minimize $\mathcal{Z}_N$

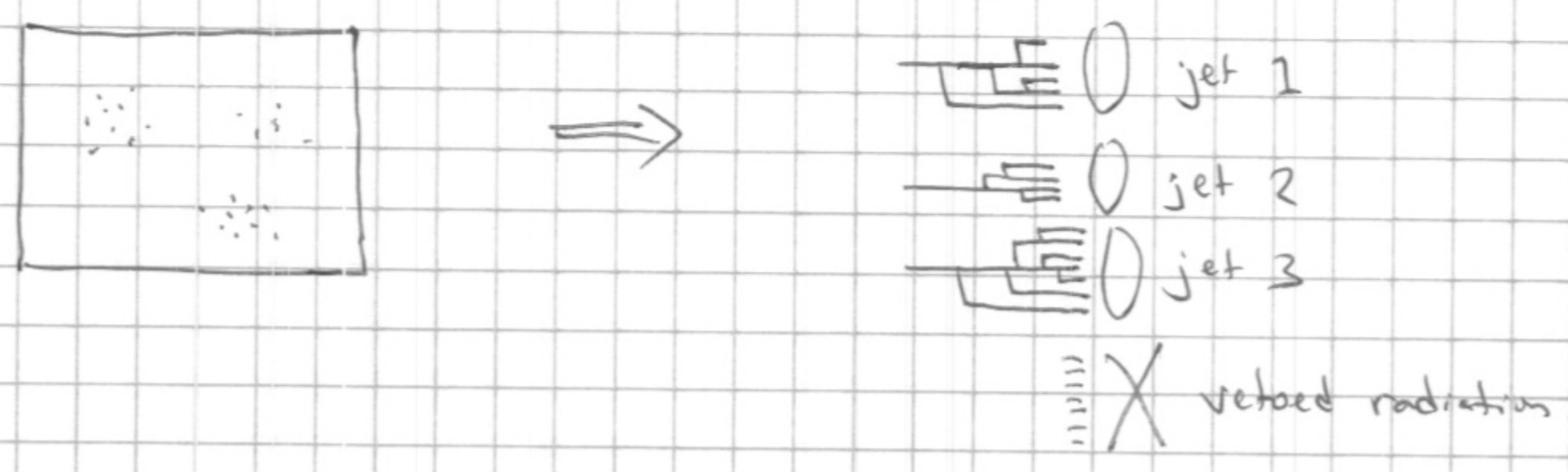
$\Rightarrow$ axes that align with hardest (by p_T) clusters

Basis for X-Cone jet algorithm.

Challenge: Computational expensive to fully optimize
Have to choose N in advance.

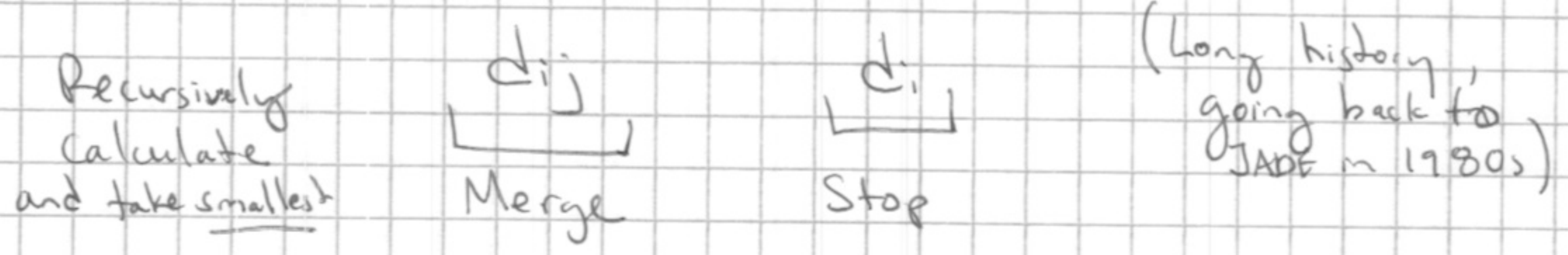
(Standard problems for cluster optimization in computer science.)

Workhorse at the LHC: Recursive Jet Algorithm, (a.k.a hierarchical clustering trees)



Instead of global optimization, "local" optimization
↳ in tree-space.

Should I combine two protojets together or
should I declare a protojet finished?



- Nice properties:
- Works the same on partons/hadrons/calor cells
 - Computationally very efficient ($n \log n$)
 - Defined perturbatively to all α_s orders
 - (Mostly) pathology-free.

Note: $p_{ij} = p_i + p_j$ in E-scheme

Let's make anti- k_t ! (Almost every LHC jet is anti- k_t) 20

Goal 1: Prioritize hardest jets

$$d_i = \frac{1}{p_{Ti}^2}$$

First remove jets of highest p_T

Goal 2: Jets should soak up all radiation within radius R_0 .

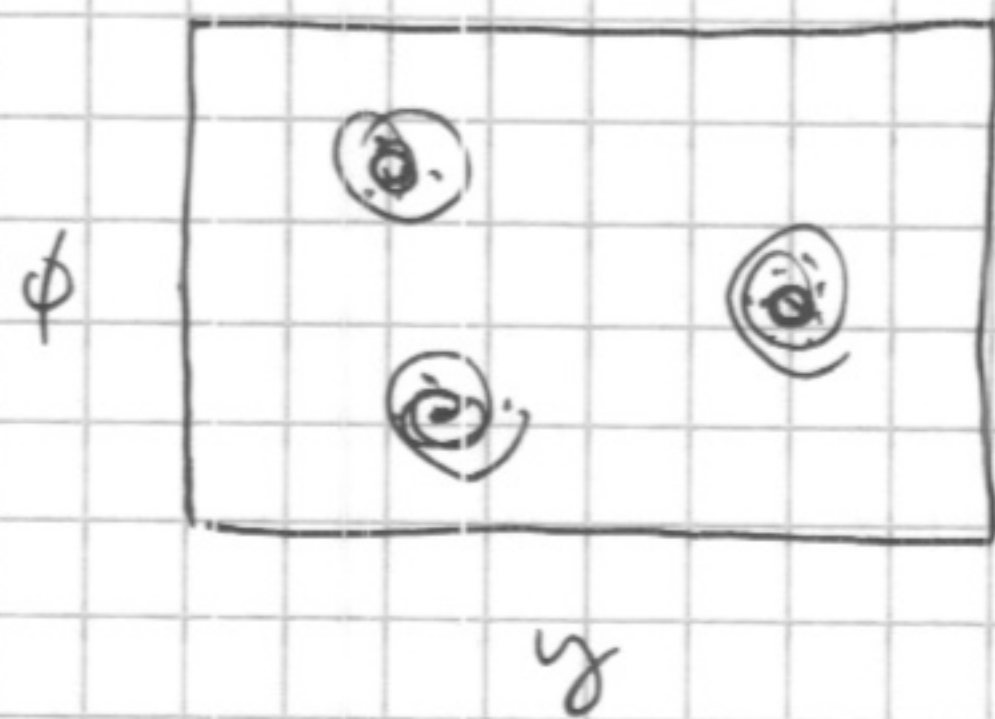
Want $d_{ij} < d_i$ if $R_{ij} < R_0$

$$d_{ij} = \min \{ d_i, d_j \} \left(\frac{R_{ij}}{R_0} \right)^2$$

$$d_{ij} = \frac{1}{\max \{ p_{Ti}^2, p_{Tj}^2 \}} \left(\frac{R_{ij}}{R_0} \right)^2$$

Done!
That's anti- k_T

Anti- k_t : Builds jets from hard to soft



Roughly: take hardest particle and merge with nearest neighbor until R_0

For me: A revelation when I read this paper in 2008
(Why did it take so long?)

Different goals?

k_t

Cluster according
to singularity
structure of
QCD

$$d_{ij} = \min \{ p_{Ti}^2, p_{Tj}^2 \} \left(\frac{R_{ij}}{R_0} \right)^2 \quad d_i = p_{Ti}^2$$

C/A

Cluster by
angle alone

$$d_{ij} = \frac{R_{ij}^2}{R_0^2}$$

$$d_i = 1$$

Both of these yield "amoeba jets"
with irregular shapes

Still, very important for
jet substructure.

Choice of jet algorithm depends on physics of interest!

E.g.

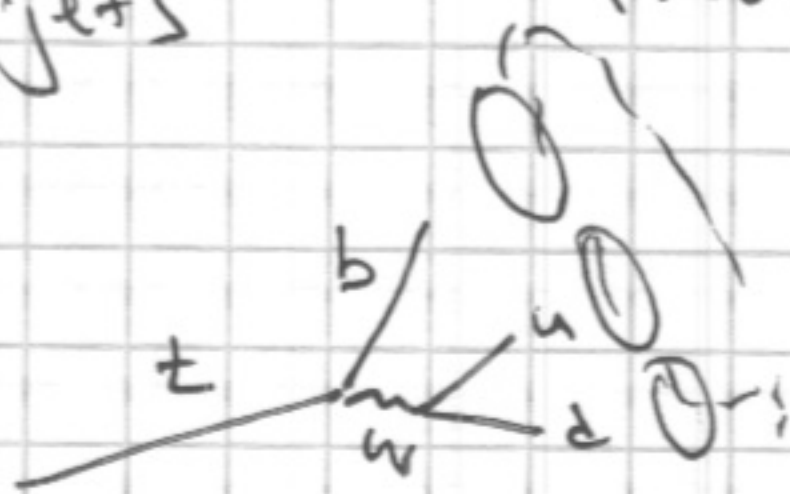
Anti- k_T

$R=1.0$
jets

3 k_T

$\Rightarrow$ reclustered $\Rightarrow$ subjects
in do

for



boosted top
identification.

E.g.

Anti- k_t
 $R=0.4$

$\Rightarrow$

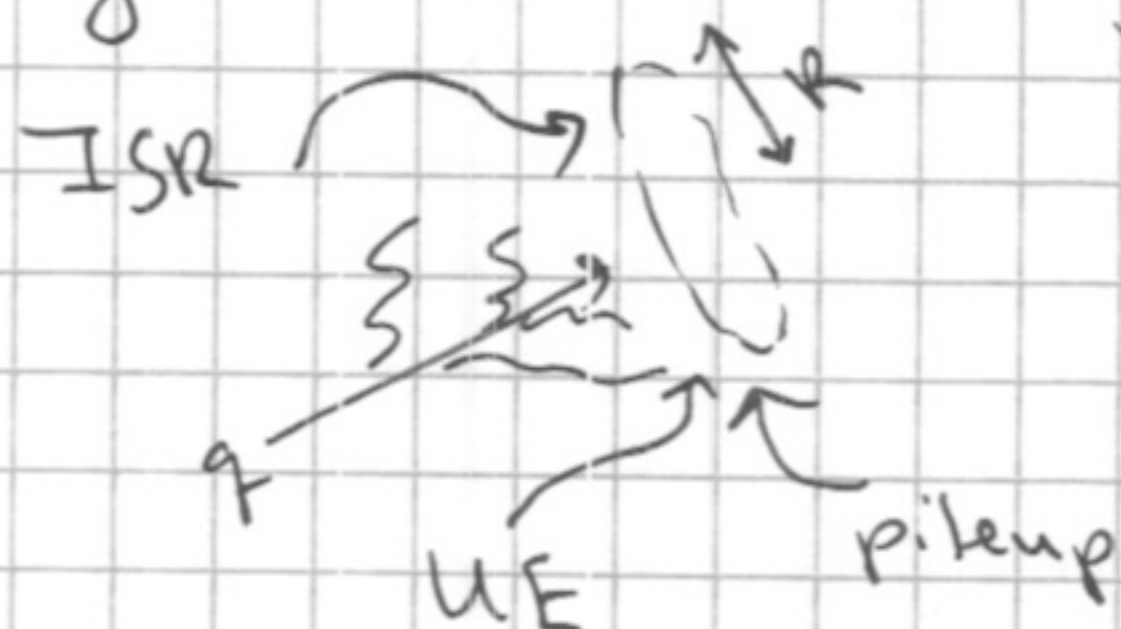
Good balance between in-and-out-
at cone corrections.
Well defined jet area (usually)

Periodically useful to revisit standard jet lore

- Longitudinal boost invariance?
- Manifest energy/momentum conservation? ($P_{12} = p_1 + p_2$)
- Fixed jet radius?
- Seedless jet algorithms?
- Unique assignment of hadrons to clusters?
- Unique jet reconstruction per event?

Your job is to question your elders!

E.g. Well-defined jet area?



Want to subtract contamination from pileup

$$P_{jet}^{corr} = P_{jet} - \rho A_{jet}$$

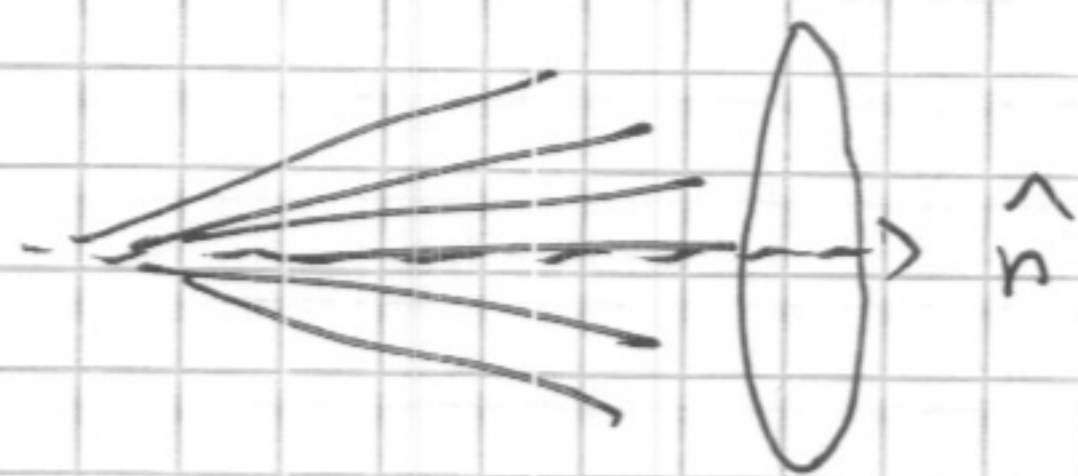
ρ ← jet area
 contamination per unit η, ϕ .
 (rough 1 GeV per unit per PV)

But why not build an algorithm that refuses to cluster soft particles to begin with?

"Zero Area Jet Algorithms?"

→ under active investigation

E.g. Jet axis parallel to jet momentum?



$$\hat{n}_{jet} = \frac{\vec{p}_{jet}}{|\vec{p}_{jet}|}$$

In terms of optimization, very natural:

$$\chi_1 = \sum_i p_{Ti} \left(\underbrace{(\phi_i - \phi_A)^2 + (y_i - y_A)^2}_{R_{iA}^2} \right)$$

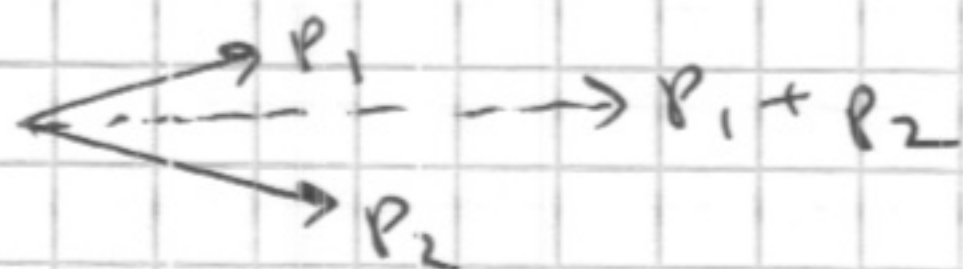
$$\frac{\partial \chi_1}{\partial \phi_A} \stackrel{?}{=} 0 = \sum_i 2 p_{Ti} (\phi_i - \phi_A)$$

$$\hookrightarrow \text{solved for } \phi_A = \frac{\sum_i p_{Ti} \phi_i}{\sum_i p_{Ti}} \quad \text{mostly the same.}$$

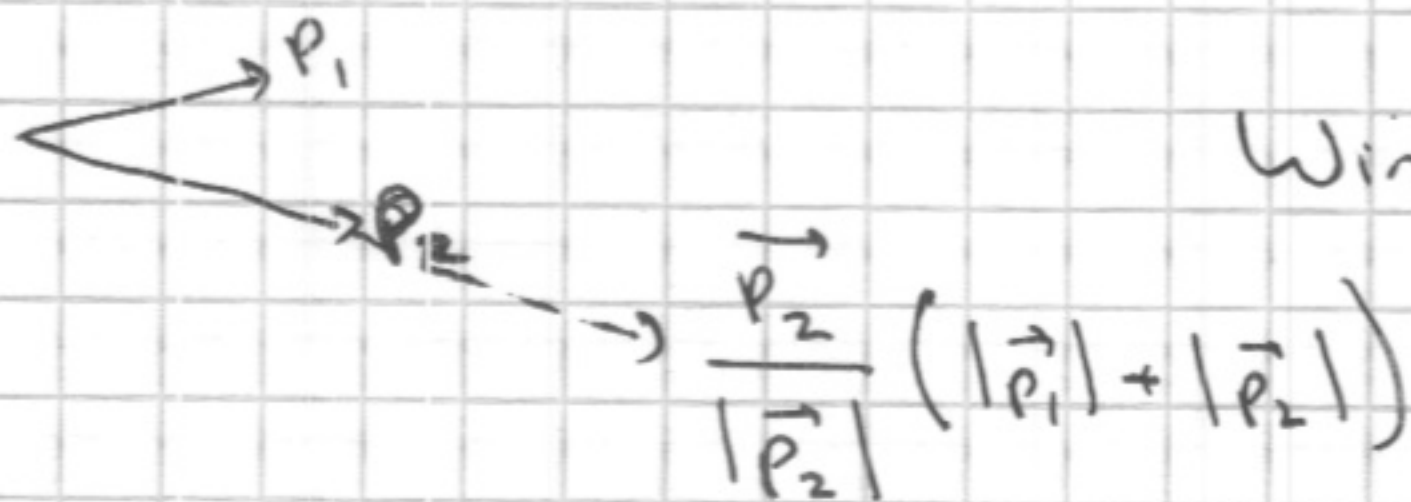
Same as finding "mean"

Why not "median"? $\chi_1 = \sum_i p_{Ti} R_{iA}^1$

In language of clustering:



E-scheme axis ("mean")



Winner-take-all axis ("median")

Like median being robust to outliers,
WTA axis robust to contamination.

One lore that you challenge at your peril!

Infrared / Collinear Safety (more next lecture)

(By the way, I've broken IRC safety with some success. Still, beware!)

Why IRC Safety?

Key success of QCD:

Fixed-order (and resummed) calculations for jet cross sections.

Multi-loop, multi-leg, multi-log. (A whole separate lecture series.)

Famously: Loop diagrams have UV divergences (fix with renormalization)

and IR divergences (cancel against real emission diagrams)

$$\frac{d\sigma}{de} = \sum_N \int d\mathcal{I}_N \frac{d\sigma}{d\mathcal{I}_N} \delta(e - \hat{e}(\mathcal{I}_N))$$

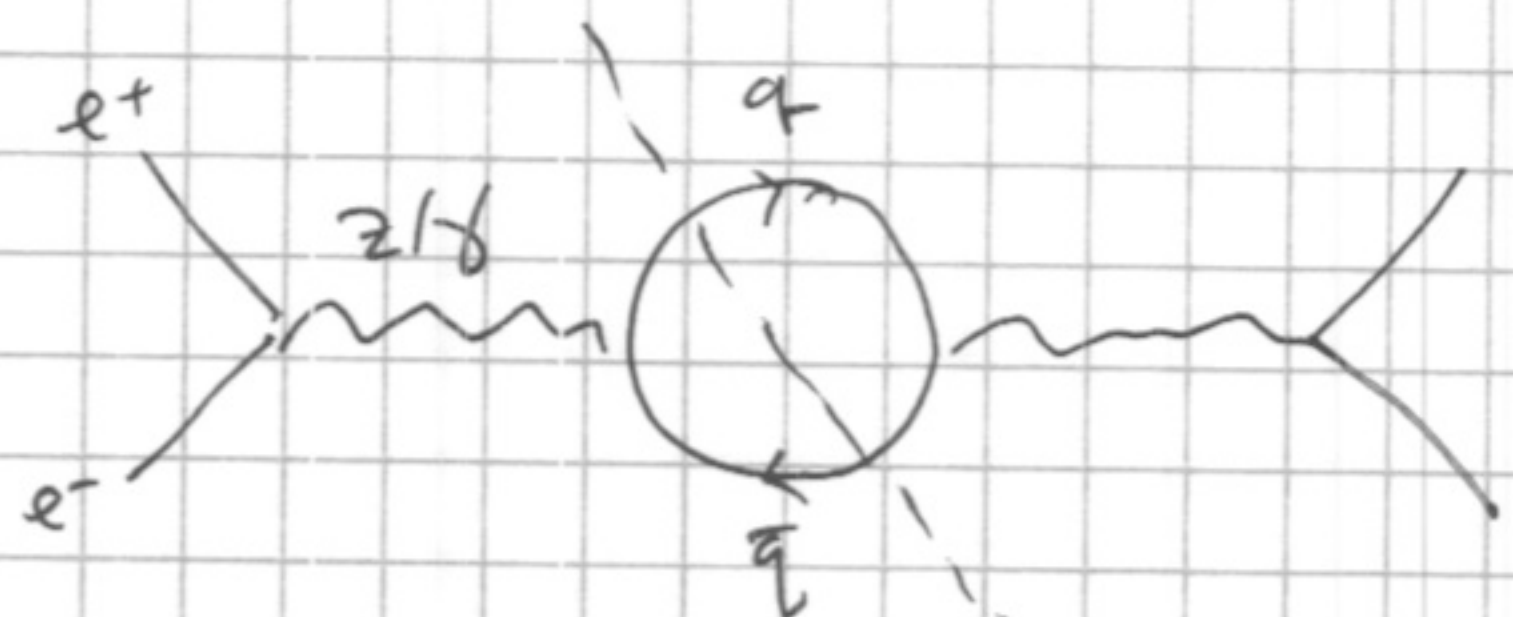
$\uparrow$ Finite if e is IRC safe observable
 $\uparrow$ N-body phase space
 $\uparrow$ IR divergences at every order
 $\uparrow$ measurement function.

IRC Safety required for sensible perturbative calculations.
 ↳ well, not quite (e.g. Sudakov safety)

For total cross sections, IRC safety easy to understand by optical theorem.

$$\sigma_{2 \rightarrow \text{anything}} \iff \text{Im } M_{\text{forward}}$$

e.g. $e^+e^- \rightarrow \text{hadrons} \iff e^+e^- \rightarrow e^+e^-$



$$\mathcal{O}(\alpha_s^0)$$

Imaginary part from cutting rules.



$$\mathcal{O}(\alpha_s^1)$$

$e^+e^- \rightarrow q\bar{q}g$ (tree)
 $e^+e^- \rightarrow q\bar{q}$ (one-loop)
 Singularities have to cancel!

Somewhat counter intuitively:

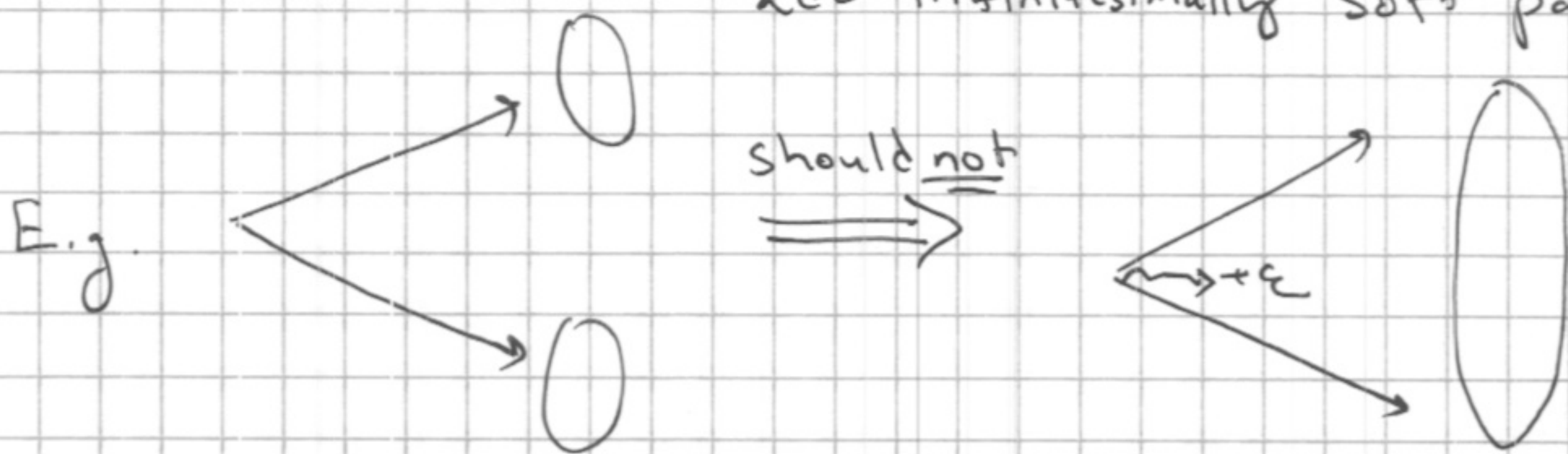
$2 \rightarrow 2$ tree-level cross section is good approximation to total cross section
 (not exclusive 2 jet xsec)

$$2 \rightarrow 2_{\text{loop}} + 2 \rightarrow 3_{\text{tree}} \text{ largely cancel.}$$

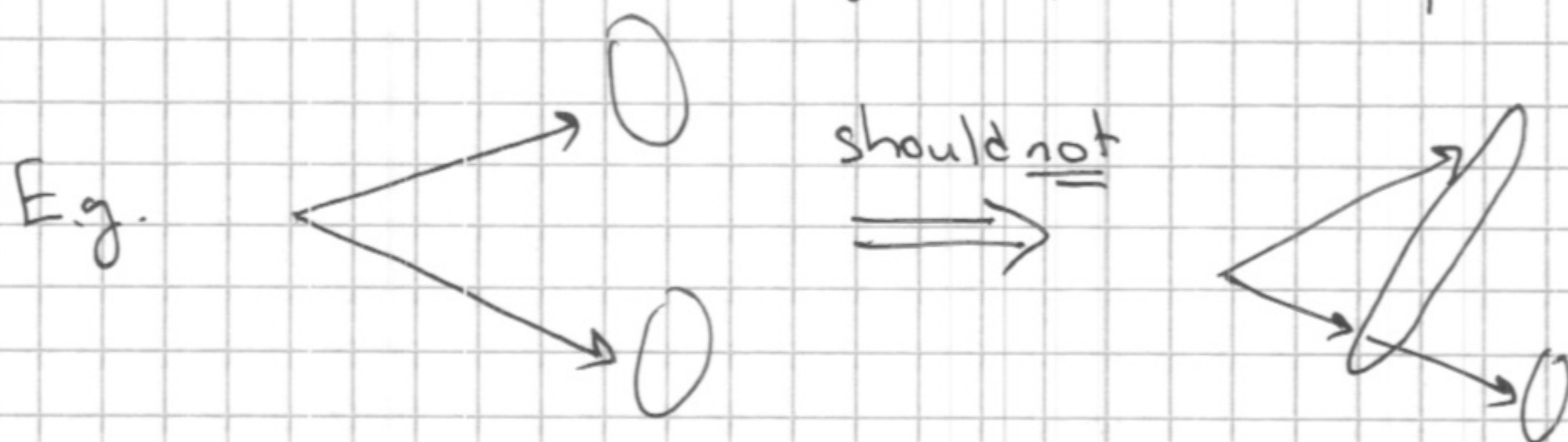
For jet cross sections, IRC safety is a bit non-trivial, and people have gotten in trouble over the years.

Heuristic for IRC Safety (almost always works)

Soft safety: Answer should not change if you add infinitesimally soft particle



Collinear safety: Answer should not change if you split a particle exactly in two



Among jet aficionados, IRC Safety is the first thing we check about any new idea.

Not safe? Not calculable in fixed-order perturbative QCD!

(Perhaps other strategies available, though)

Is anti- k_t safe?

$$d_{ij} = \frac{1}{\max\{p_{Ti}^2, p_{Tj}^2\}} \frac{R_{ij}^2}{R_0^2} \quad d_i = \frac{1}{p_{Ti}^2}$$

Soft emission? $d_i^{\text{soft}} \rightarrow \infty$ (no effect)

d_{ij}^{soft} could be important!
but infinitesimally soft particle
does not change jet p_T

Depends on question!

Anti- k_t jet p_T ? Soft safe

Anti- k_t jet multiplicity of constituents? Soft unsafe.

Collinear split? $d_i^{\text{collinear}}$ changes!

But $d_{ij}^{\text{collinear}} = 0$, so first thing you
do is put split back together.

So yes, anti- k_t is IRC safe.

NNLO jet cross sections being actively developed.

Next time: What exactly are soft/collinear
singularities, and can we use
them to understand quark/gluon tagging.