Gravitational Origin of Dark Matter

mostly based on arXiv 1604.08564

Angnis Schmidt-May



MINI-WORKSHOP ON "BEYOND THE WIMP PARADIGM" June 8, 2016 Mainz Institute for Theoretical Physics

In collaboration with ...

1604.08564

Eugeny Babichev Luca Marzola Martti Raidal Federico Urban Hardi Veermäe Mikael von Strauss



Yashar Akrami Oliver Baldacchino Laura Bernard Cédric Deffayet Jonas Enander Fawad Hassan Mikica Kocic Frank Könnig Edvard Mörtsell Rachel Rosen Adam Solomon

Activation

Contents

- ☆ Massless & Massive Spin-2 Fields
- ☆ The Ghost-Free Theory
- Deviations from General Relativity
- ☆ Spin-2 Dark Matter
- ☆ Conclusions



Philosophy

Approach I

- 1. invent a model to explain observations
 - many possibilities
- 2. check if model is consistent, fits into larger framework, has motivations besides cosmology, etc.

Philosophy

Approach I

- 1. invent a model to explain observations
 - many possibilities
- 2. check if model is consistent, fits into larger framework, has motivations besides cosmology, etc.

Approach 2

- construct a consistent model guided by a fundamental question
 few possibilities
- 2. check if it can explain (part of) an observational phenomenon

Philosophy

Approach I

invent a model to explain observations
 many possibilities

2. check if model is consistent, fits into larger framework, has motivations besides cosmology, etc.

Approach 2

construct a consistent model guided by a fundamental question
 few possibilities

2. check if it can explain (part of) an observational phenomenon

Consistent Field Theories

Standard Model of Particle Physics & General Relativity

Spin 0: Higgs boson ϕ

- Spin 1/2: leptons, quarks ψ^a
- Spin 1: gluons, photon, W- & Z-boson A_{μ}

Spin 2: graviton $g_{\mu\nu}$

Consistent Field Theories

Standard Model of Particle Physics & General Relativity

Spin 0: Higgs boson ϕ

- Spin 1/2: leptons, quarks ψ^a
- Spin 1: gluons, photon, W- & Z-boson A_{μ}

 $g_{\mu
u}$

Spin 2: graviton

MASSLESS !

massive &

massless



How do we describe massive spin-2 fields ?

Massless + Massive Spin-2 Fields

General Relativity (GR)

Einstein-Hilbert action:

$$S_{\rm EH}[g] = M_{\rm P}^2 \int \mathrm{d}^4 x \sqrt{g} \left(R(g) - 2\Lambda \right)$$

Einstein's equations: $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + \Lambda g_{\mu\nu} = 0$

Around maximally symmetric solutions, $\bar{R}_{\mu\nu} = \Lambda \bar{g}_{\mu\nu}$, the propagating degrees of freedom are those of a <u>massless spin-2 field</u>.

Also nonlinearly GR contains two dynamical degrees of freedom.



General Relativity

nonlinear theory of massless spin-2

Nonlinear Mass Term

... should not contain derivatives nor loose indices.

But if we try to contract the indices of the metric, we get: $g^{\mu\nu}g_{\mu\nu} = 4$ This is not a mass term.

Simplest way out: Introduce second "metric" to contract indices:

$$g^{\mu\nu}f_{\mu\nu} = \text{Tr}(g^{-1}f) \qquad f^{\mu\nu}g_{\mu\nu} = \text{Tr}(f^{-1}g)$$

Nonlinear Mass Term

... should not contain derivatives nor loose indices.

But if we try to contract the indices of the metric, we get: $g^{\mu\nu}g_{\mu\nu} = 4$ This is not a mass term.

Simplest way out: Introduce second "metric" to contract indices:

$$g^{\mu\nu}f_{\mu\nu} = \text{Tr}(g^{-1}f) \qquad f^{\mu\nu}g_{\mu\nu} = \text{Tr}(f^{-1}g)$$



4/15

Massive gravity action is of the form

$$S_{\rm MG}[g] = S_{\rm EH}[g] - \int {\rm d}^4 x \, V(g, f)$$

kinetic term mass term

mass term



Bimetric Theory

Isham, Salam & Strathdee (1971)

Nonlinear bimetric action:

$$S_{\rm b}[g,f] = m_g^2 \int \mathrm{d}^4 x \sqrt{g} \left(R(g) - 2\Lambda \right) + m_f^2 \int \mathrm{d}^4 x \sqrt{f} \left(R(f) - 2\tilde{\Lambda} \right) - \int \mathrm{d}^4 x \, V(g,f)$$

- ☆ both metrics are dynamical and treated on equal footing
- should describe massive & massless spin-2 field (5+2 d.o.f.)



This looks nice ...

... but unfortunately the general theory has ghosts!





Ghosts

Ghost = field with negative kinetic energy

$$\mathcal{L} = (\partial_t \phi)^2 \cdots$$
 healthy $\mathcal{L} = -(\partial_t \phi)^2 \cdots$ ghost

consequences: classical instability, negative probabilities at quantum level
 must be avoided!



Ghosts

Ghost = field with negative kinetic energy

- $\mathcal{L} = (\partial_t \phi)^2 \cdots$ healthy $\mathcal{L} = -(\partial_t \phi)^2 \cdots$ ghost
- consequences: classical instability, negative probabilities at quantum level
 must be avoided!

Modifications of General Relativity tend to be haunted by ghosts. Modifying gravity is EXTREMELY difficult!



Ghosts

Ghost = field with negative kinetic energy

- $\mathcal{L} = (\partial_t \phi)^2 \cdots$ healthy $\mathcal{L} = -(\partial_t \phi)^2 \cdots$ ghost
- consequences: classical instability, negative probabilities at quantum level
 must be avoided!

Modifications of General Relativity tend to be haunted by ghosts. Modifying gravity is EXTREMELY difficult!

Boulware & Deser (1972): Nonlinear massive gravity always has a ghost!

6/15



Massive gravity stinks. If you want to modify gravity, try something else...

Quote from lecture notes by Kurt Hinterbichler, 2010 (now turned into a very nice review!)

The Ghost-Free Theory

Development

Creminelli, Nicolis, Papucci, Trincherini (2005): attempt to construct ghost-free candidate theory; fails only because of unfortunate sign mistake

de Rham, Gabadadze, Tolley (2010): construction of candidate theory for massive gravity in flat reference frame; ghost-free in "decoupling limit"

Hassan, Rosen, ASM, von Strauss (2011/12): proof of absence of ghost in fully nonlinear theory

Hassan & Rosen (2011): generalisation to ghost-free bimetric theory

Hassan & Rosen (2011)

$$S_{b}[g,f] = m_{g}^{2} \int d^{4}x \sqrt{g} R(g)$$

+ $m_{f}^{2} \int d^{4}x \sqrt{f} R(f) - \int d^{4}x V(g,f)$

$$\left(V(g,f) = m^4 \sqrt{g} \sum_{n=0}^4 \beta_n e_n \left(\sqrt{g^{-1}f} \right) = m^4 \sqrt{f} \sum_{n=0}^4 \beta_{4-n} e_n \left(\sqrt{f^{-1}g} \right) \right)$$

$$e_1(S) = \operatorname{Tr}[S] \qquad e_2(S) = \frac{1}{2} \left((\operatorname{Tr}[S])^2 - \operatorname{Tr}[S^2] \right)$$
$$e_3(S) = \frac{1}{6} \left((\operatorname{Tr}[S])^3 - 3 \operatorname{Tr}[S^2] \operatorname{Tr}[S] + 2 \operatorname{Tr}[S^3] \right)$$

E

- free bimetric theory

Hassan & Rosen (2011)

$$S_{b}[g,f] = m_{g}^{2} \int d^{4}x \sqrt{g} R(g)$$

+ $m_{f}^{2} \int d^{4}x \sqrt{f} R(f) - \int d^{4}x V(g,f)$

$$V(g,f) = m^4 \sqrt{g} \sum_{n=0}^4 \beta_n e_n \left(\sqrt{g^{-1}f}\right) = m^4 \sqrt{f} \sum_{n=0}^4 \beta_{4-n} e_n \left(\sqrt{f^{-1}g}\right)$$

more details: ASM, Mikael von Strauss; 1512.00021

- free bimetric theory



What is the physical content of ghost-free bimetric theory ?

Hassan, ASM, von Strauss (2012)

Mass spectrum

Proportional solutions: $\bar{f}_{\mu\nu} = c^2 \bar{g}_{\mu\nu}$ with c = const.

Maximally symmetric backgrounds with $~R_{\mu
u}(ar{g})=\Lambda_gar{g}_{\mu
u}$

Perturbations:
$$g_{\mu\nu} = \bar{g}_{\mu\nu} + \delta g_{\mu\nu}$$
 $f_{\mu\nu} = c^2 \bar{g}_{\mu\nu} + \delta f_{\mu\nu}$

Can be diagonalised into mass eigenstates ($lpha\equiv m_f/m_g$) :

$$\delta G_{\mu
u} \propto \delta g_{\mu
u} + \alpha^2 \delta f_{\mu
u}$$
 massless (2 d.o.f.)
 $\delta M_{\mu
u} \propto \delta f_{\mu
u} - c^2 \delta g_{\mu
u}$ massive (5 d.o.f.)
with Fierz-Pauli mass $m_{\rm FP} = m_{
m FP}(\alpha, \beta_n, c)$



Ghost-free bimetric theory

nonlinear theory of massless & massive spin-2

Deviations from General Relativity



What is the physical metric ?

How does matter couple to the tensor fields ?

Matter coupling

Yamashita, de Felice, Tanaka; de Rham, Heisenberg, Ribeiro (2015)

Only one metric can couple to matter!

$$\begin{split} S_{gf} &= m_g^2 \int \mathrm{d}^4 x \sqrt{g} \; R(g) &+ m_f^2 \int \mathrm{d}^4 x \sqrt{f} \; R(f) \\ &- m^4 \int \mathrm{d}^4 x \sqrt{g} \; \sum_{n=0}^4 \beta_n e_n \left(\sqrt{g^{-1} f} \right) \\ &+ \int \mathrm{d}^4 x \sqrt{g} \; \mathcal{L}_{\mathrm{matter}}(g, \phi) \end{split}$$



only coupling that does not re-introduce the ghost

 $g_{\mu\nu}$ is gravitational metric

Matter coupling

Yamashita, de Felice, Tanaka; de Rham, Heisenberg, Ribeiro (2015)

Only one metric can couple to matter!

$$S_{gf} = m_g^2 \int d^4x \sqrt{g} R(g) + m_f^2 \int d^4x \sqrt{f} R(f) - m^4 \int d^4x \sqrt{g} \sum_{n=0}^4 \beta_n e_n \left(\sqrt{g^{-1}f}\right) + \int d^4x \sqrt{g} \mathcal{L}_{matter}(g,\phi)$$



only coupling that does not re-introduce the ghost

• $g_{\mu\nu}$ is gravitational metric

The gravitational metric is not massless !!

Physical Interpretation

Baccetti, Martin-Moruno, Visser (2012); Hassan, ASM, von Strauss (2012/14); Akrami, Hassan, Koennig, ASM, Solomon (2015)

Bimetric theory = General Relativity (GR) + corrections

Recall:
$$\delta g_{\mu
u}\propto\delta G_{\mu
u}-lpha^2\delta M_{\mu
u}$$

Assume that $\alpha = m_f/m_g$ is small (i.e. weak gravity!)



the gravitational metric is almost massless



the massive spin-2 field interacts only weakly with matter

Physical Interpretation

Baccetti, Martin-Moruno, Visser (2012); Hassan, ASM, von Strauss (2012/14); Akrami, Hassan, Koennig, ASM, Solomon (2015)

Bimetric theory = General Relativity (GR) + corrections

Recall:
$$\delta g_{\mu
u}\propto\delta G_{\mu
u}-lpha^2\delta M_{\mu
u}$$

Assume that $\alpha = m_f/m_g$ is small (i.e. weak gravity!)



the gravitational metric is almost massless



the massive spin-2 field interacts only weakly with matter

 $\alpha \rightarrow 0$ is the General Relativity limit of bimetric theory



Ghost-free bimetric theory

General Relativity + additional (heavy?) tensor field



Dark Matter

Babichev, Marzola, Raidal, ASM, Urban, Veermäe, von Strauss (2016); Aoki, Mukohyama (2016)

Redefining the nonlinear fields gives: $G \sim \bar{g} + \delta G$ $M \sim \delta M$

$$\begin{split} \mathcal{L} \sim \sqrt{G} \begin{bmatrix} m_{\mathrm{Pl}}^2 R(G) + K(G, \nabla \nabla M) + V(G, M) + \mathcal{L}_{\mathrm{m}}(G, \alpha M, \phi) \end{bmatrix} \\ & \text{Gravity} & \text{Dark Matter} & \text{Baryonic Matter} \end{split}$$

In the General Relativity (GR) limit of bimetric theory, $\alpha \rightarrow 0$: massive spin-2 field decouples from matter, interacts only with gravity.

A large spin-2 mass further suppresses deviations from GR.

Consistency checks

×~

Our spin-2 dark matter is part of gravity (!) and...

- 🔊 ... gravitates just like baryonic matter
 - ... does not decay into gravitons and its decay rate into Standard Model fields is sufficiently small: $\Gamma(\delta M \to XX) \sim \frac{\alpha^2 m_{\rm FP}^3}{m_{\odot}^2}$
 - automatically stable
- ... has interactions with baryonic matter which are naturally suppressed by the Planck scale
- \gg ... can be produced thermally for a mass of 1 10⁸ TeV

Detection

- ☆ not observable in current indirect and direct detection experiments
- massive spin-2 field may gravitate differently in curved backgrounds
 - non-standard behaviour of dark matter around massive objects ?
- dark matter self-interactions: could be observable in cluster collisions and in power spectrum



correlations with gravitational waves Aoki, Mu

Lessons learned

🔊 is one of the few known consistent modifications of General Relativity

vert can be interpreted as gravity in the presence of an extra spin-2 field

contains an interesting dark matter candidate whose coupling to baryonic matter is suppressed by the Planck scale



"Screening" does not work. But maybe extra symmetries?

Could have spin-2!

25% – Dark Matter

Viable cosmology with self-accelerating solutions

Akrami, Hassan, Könnig, ASM, Solomon (2015); Könnig, Patil, Amendola (2014); Akrami, Koivisto, Mota, Sandstad (2013); Volkov; von Strauss, ASM, Enander, Mörtsell, Hassan; Comelli, Crisostomi, Nesti, Pilo (2011)

70% Dark Energy

> 5% normal matter