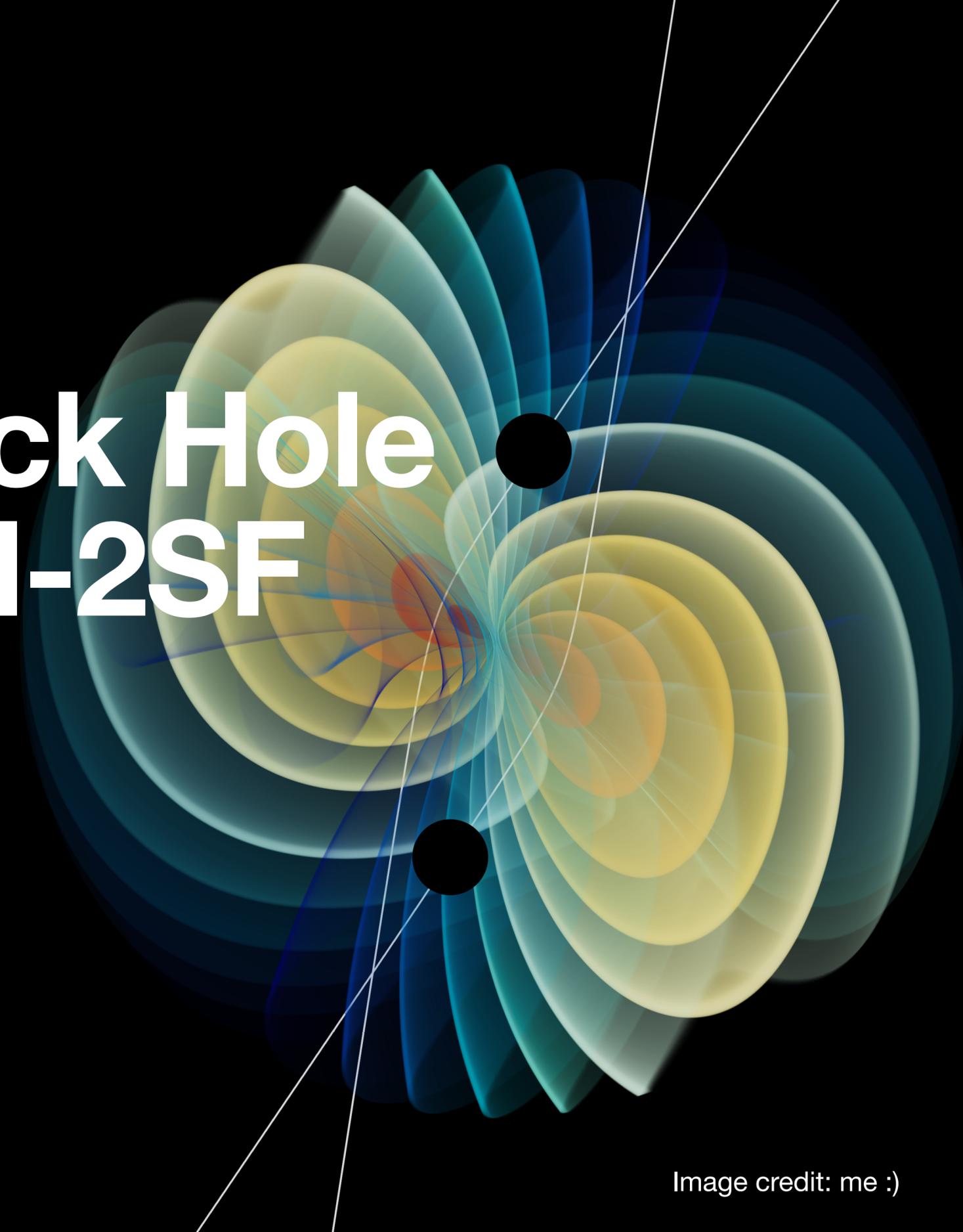


Conservative Black Hole Scattering @ 5PM-2SF



Based on [2601.xxxxx]

in collaboration with:

Gustav Jakobsen, Gustav Mogull, Christoph Nega, Jan
Plefka, Benjamin Sauer, Johann Usovitsch

Mathias Driesse, Lost in Translation 19.01.2026

Image credit: me :)

By the end of this talk, you will know:

Motivation

WQFT formalism for gravity

Comparisons of 5PM-1SF to SF and NR

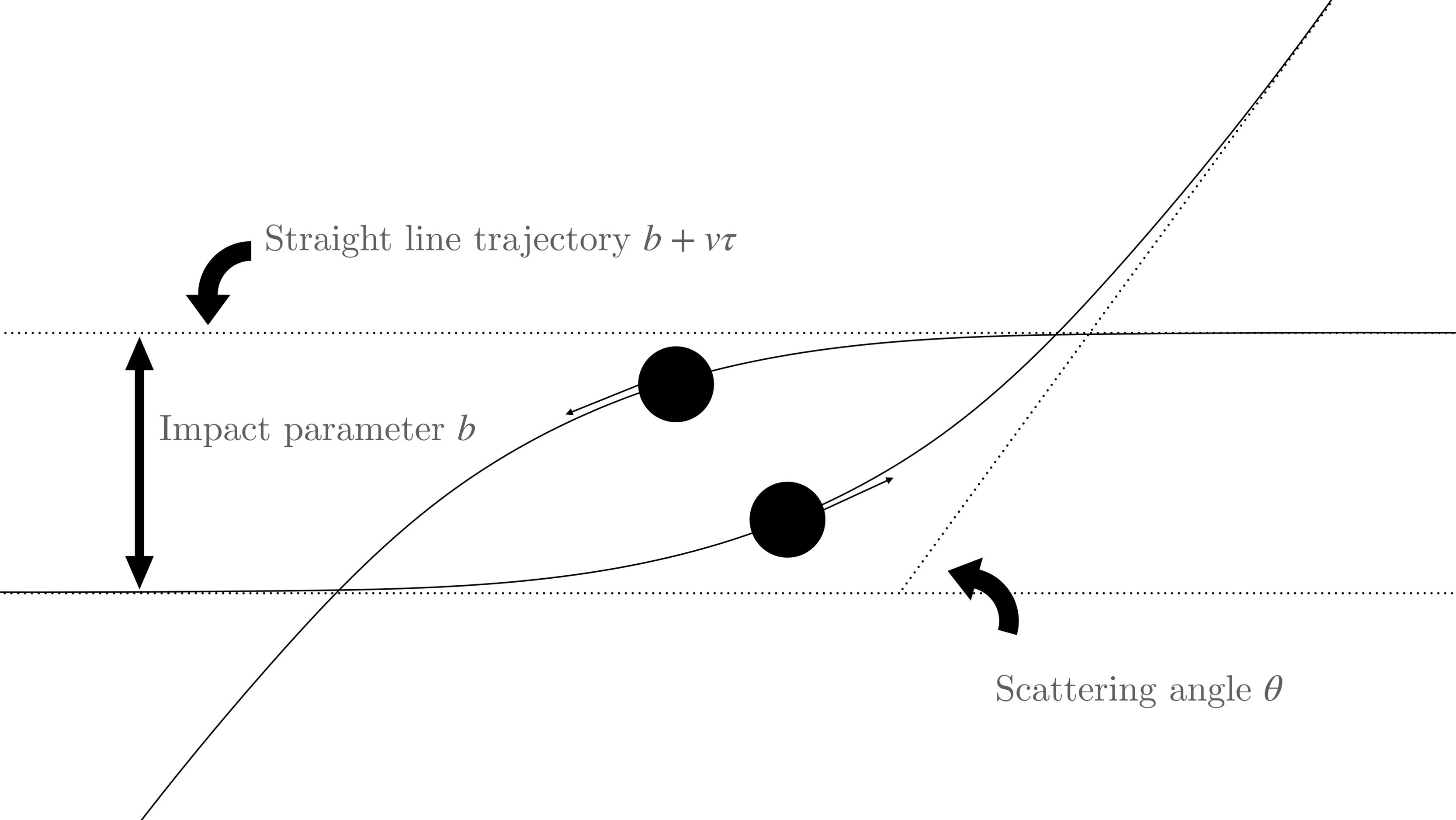
Preliminary 2SF results

The Big Picture



Building Analytical Models

- Effective-One-Body [Buonanno, Damour 1999]
- Self-force/PN/PM hybridization [Honet, Pound, Compère 2025]
- Phenomenological models [many people, ~2007]
- That's all I know of

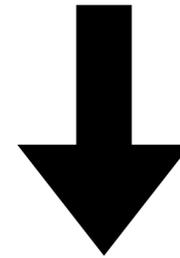


[Mogull, Plefka, Steinhoff 2021]

Worldline Quantum Field Theory

“Tree-level one-point functions solve the classical equations of motion”

[Boulware, Brown 1967]



$$S = \int d\tau \text{ point-particle} + \int d^D x \text{ gravity}$$

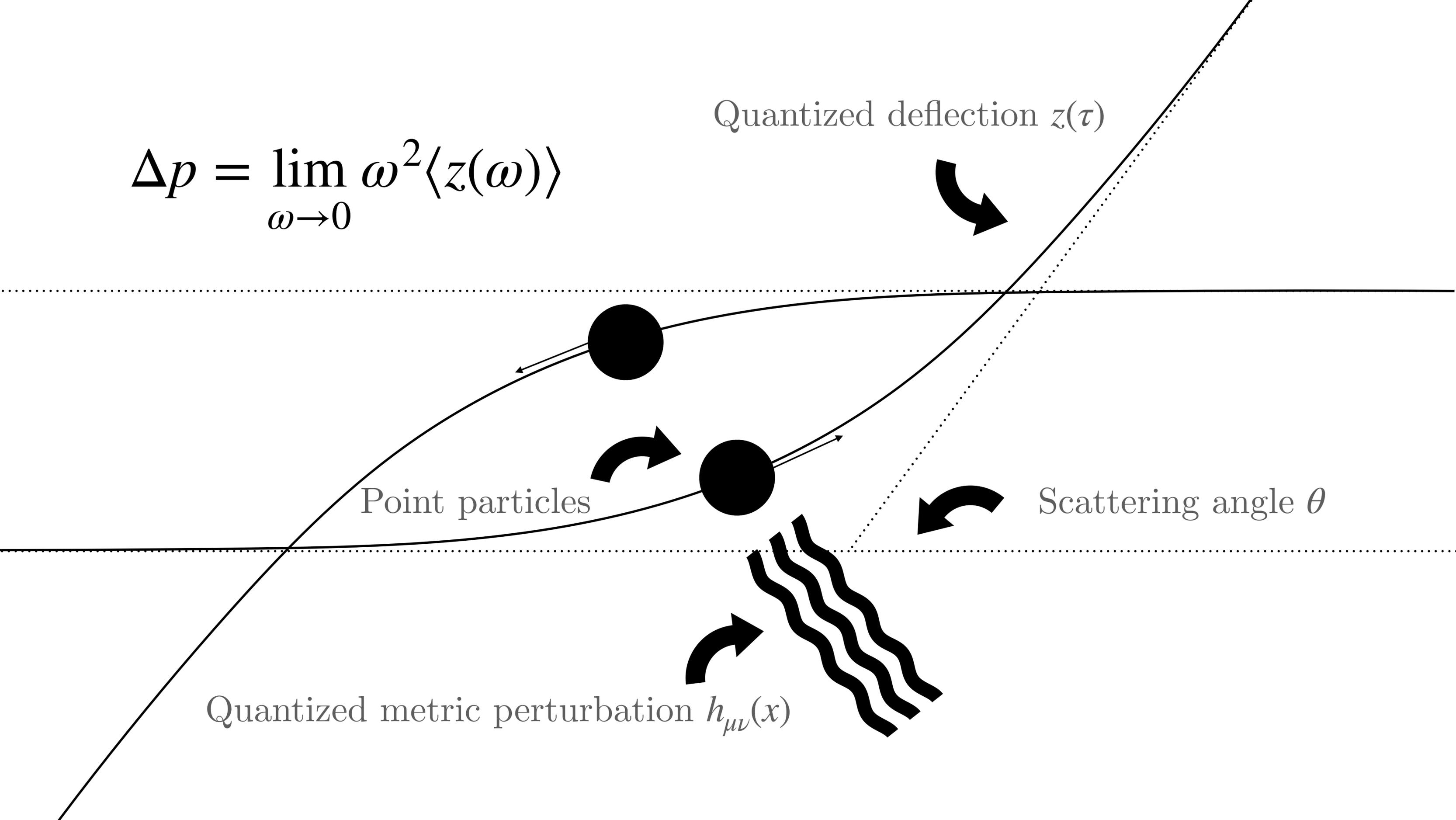
$$\Delta p = \lim_{\omega \rightarrow 0} \omega^2 \langle z(\omega) \rangle$$

Quantized deflection $z(\tau)$

Point particles

Scattering angle θ

Quantized metric perturbation $h_{\mu\nu}(x)$



The post-Minkowskian expansion

Basic idea: Post-Minkowskian expansion in $\frac{Gm}{b}$

$$\Delta p = G^1 \Delta p^{(1)} + G^2 \Delta p^{(2)} + \dots$$

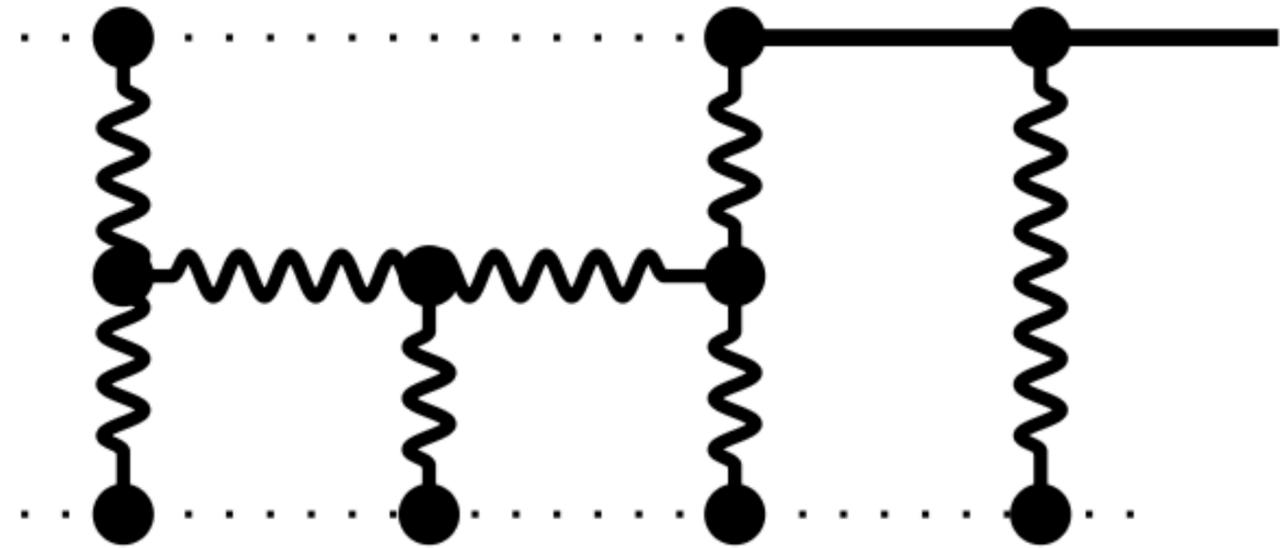
ONLY function of $\frac{1}{\sqrt{1-v^2}} = \gamma = \frac{\left(x + \frac{1}{x}\right)}{2}$

Features of the WQFT formalism

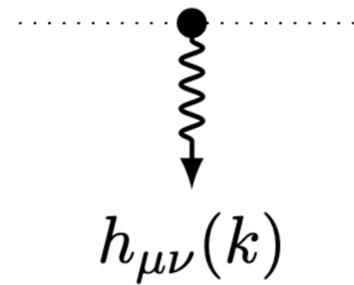
- Tree-level integrand \rightarrow **recursion**
- Schwinger-Keldysh in-in formalism \rightarrow **causality, retarded propagators**
- Energy is conserved at worldline vertices \rightarrow **loop integration**

Example graph:

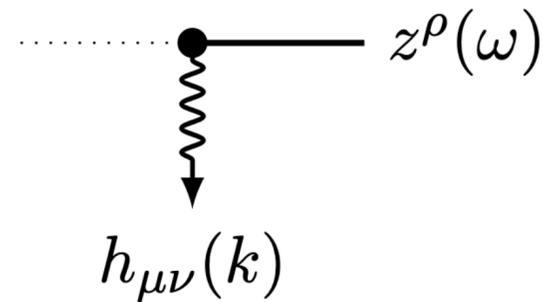
$$\begin{aligned} \text{---} \xrightarrow{\ell^\mu} \text{---} &= \frac{1}{\ell \cdot v_j + i0^+} \\ \text{---} \xrightarrow{\ell^\mu} \text{---} &= \frac{1}{(\ell^0 + i0^+)^2 - \ell^2} \\ \text{---} \xrightarrow{\ell^\mu} \text{---} &= \delta(\ell \cdot v_j) \end{aligned}$$



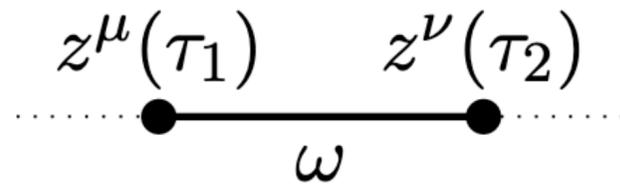
The self-force expansion from WQFT



$$\sim m$$



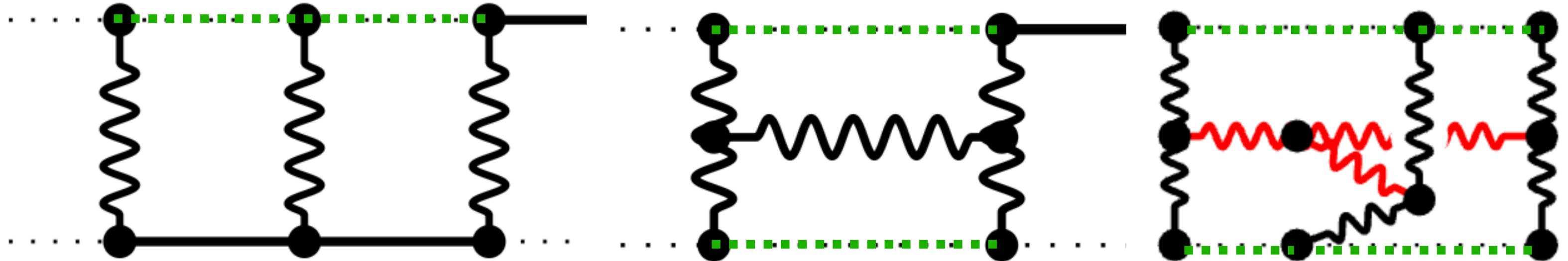
$$\sim m$$



$$\sim \frac{1}{m}$$

Polynomial mass
dependence!

The self-force expansion



0SF

1SF

2SF

The self-force expansion of the impulse

	0SF	1SF	2SF
1PM	✓	—	—
2PM	✓	—	—
3PM	✓	✓	—
4PM	✓	✓	—
5PM	✓	✓	✓ / ✗

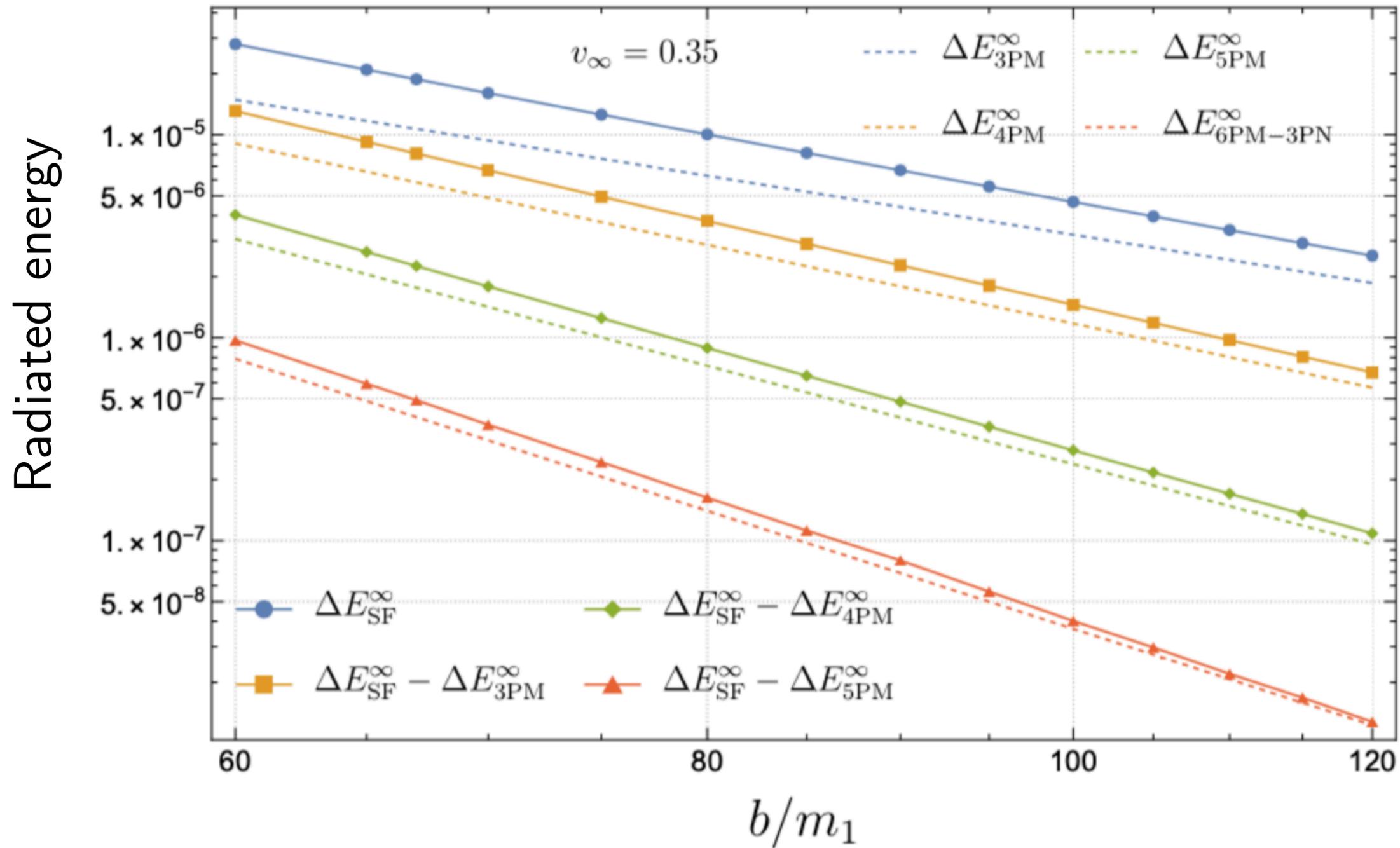
[Westpfahl, 1985]

[Bern, Parra-Martinez, Roiban, Ruf, Shen, Solon, Zeng]
 [Dlapa, Källin, Liu, Porto]
 [Vanhove, Damgaard, Plante]
 [Jakobsen, Mogull, Plefka, Sauer]

[Bern, Cheung, Roiban, Shen, Solon, Zeng]
 [Kälin, Liu, Porto][Di Vecchia, Heissenberg, Russo, Veneziano]
 [Bjerrum-Bohr, Vanhove, Damgaard]
 [Brandhuber, Chen, Travaglini, Wen]
 [Jakobsen, Mogull, Plefka, Sauer]

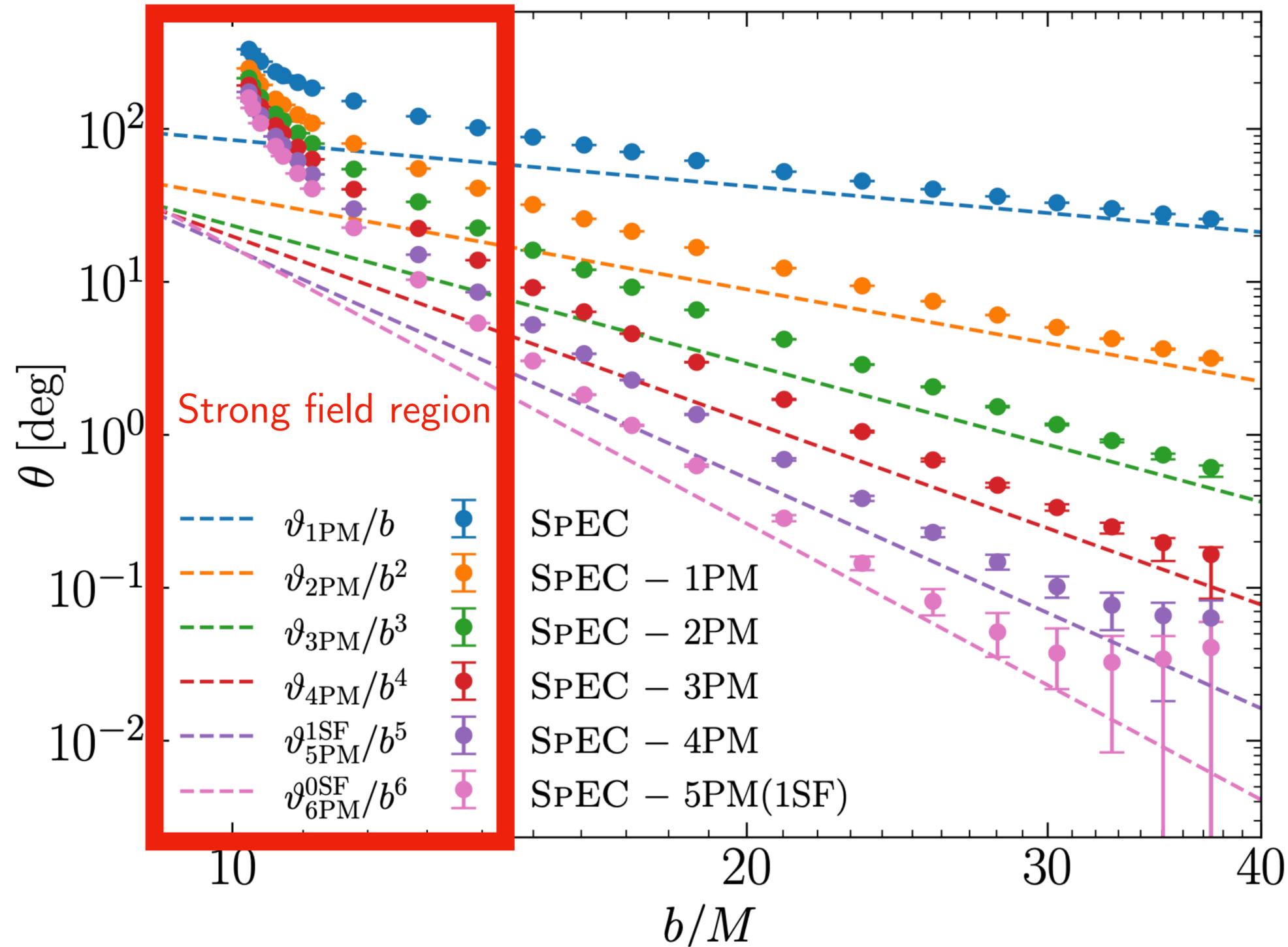
[**MD**, Jakobsen, Mogull, Plefka, Sauer, Usovitsch]
 [**MD**, Jakobsen, Klemm, Mogull, Nega, Plefka, Sauer, Usovitsch]
 [Bern, Hermann, Roiban, Ruf, Smirnov, Smith, Zeng]

Matching PM to Self-Force



[Warburton 2025]

Matching PM to Numerical Relativity



[Long, Pfeiffer,
Kidder, Scheel 2025]

Method of differential equations

Idea: relate derivatives of integrals to themselves

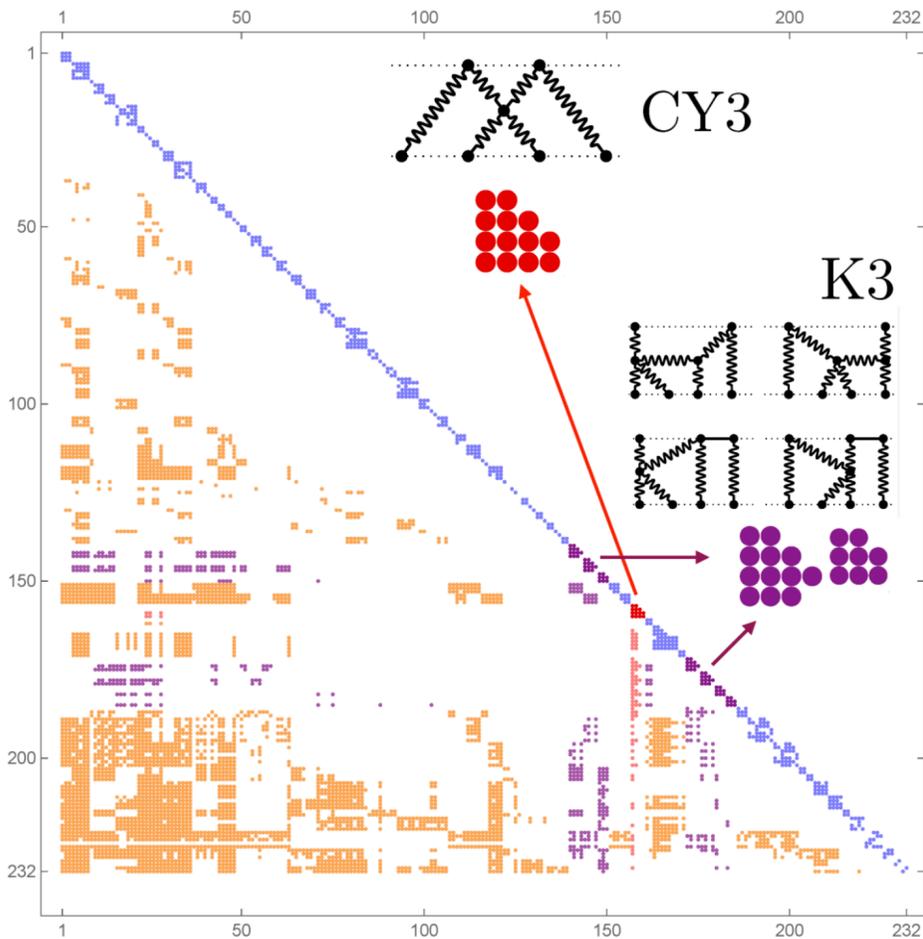
$$\frac{d}{dx} \vec{I} = M(x, \epsilon) \vec{I} \quad [\text{Gehrmann, Remiddi 1999}]$$

Problem: non-trivial x, ϵ dependence

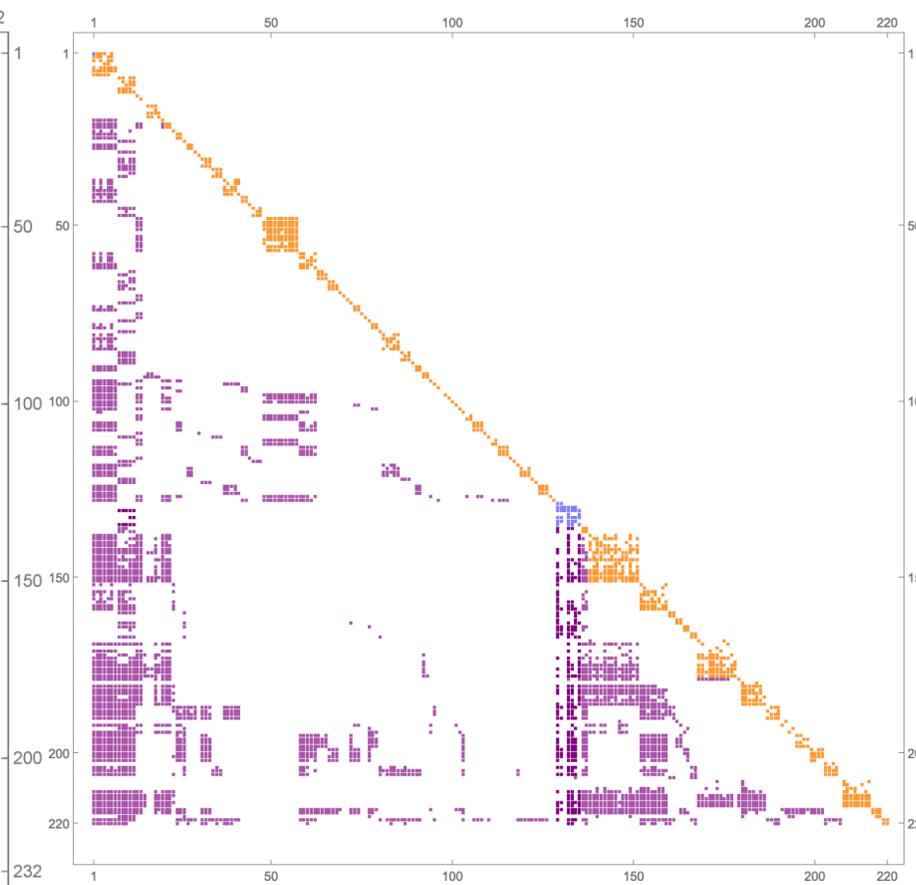
Solution: find basis transformation s.t.

$$\frac{d}{dx} \vec{I} = \epsilon M'(x) \vec{I} \quad [\text{Henn 2013}]$$

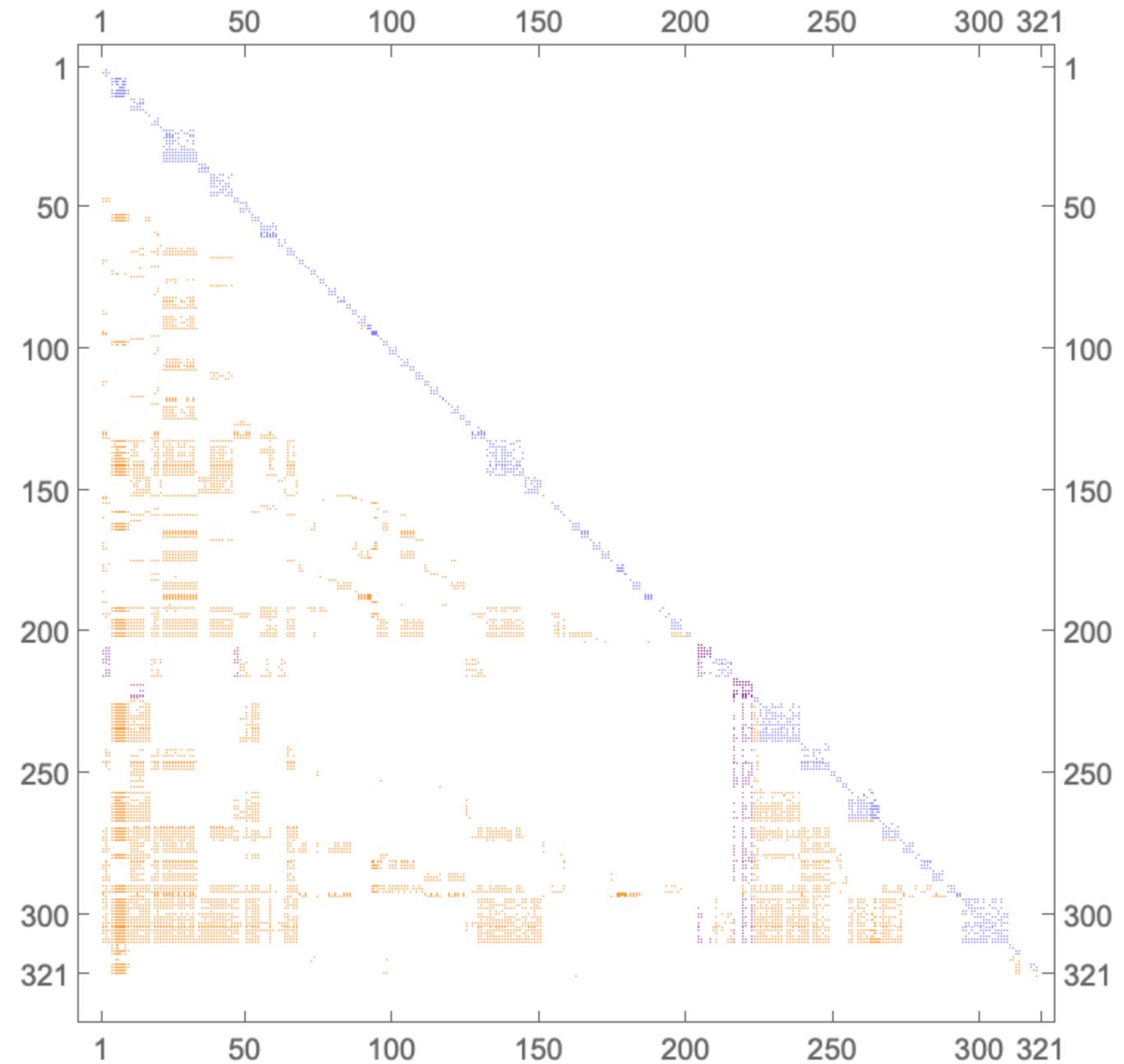
Differential equations



Odd, 1SF



Even, non-
planar, 2SF

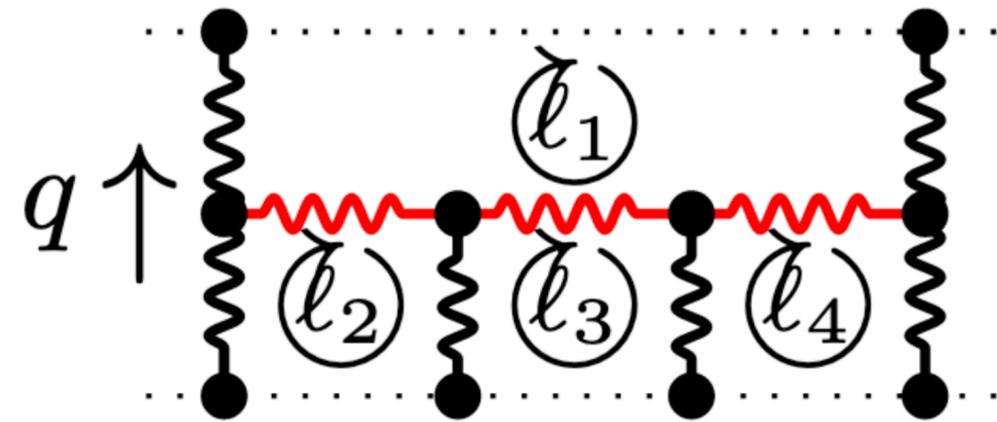


Even, planar, 2SF

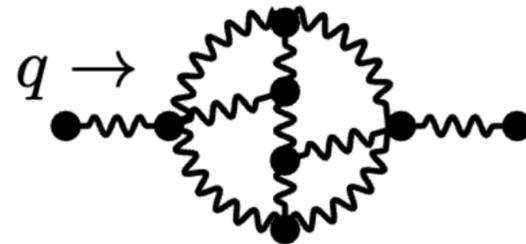
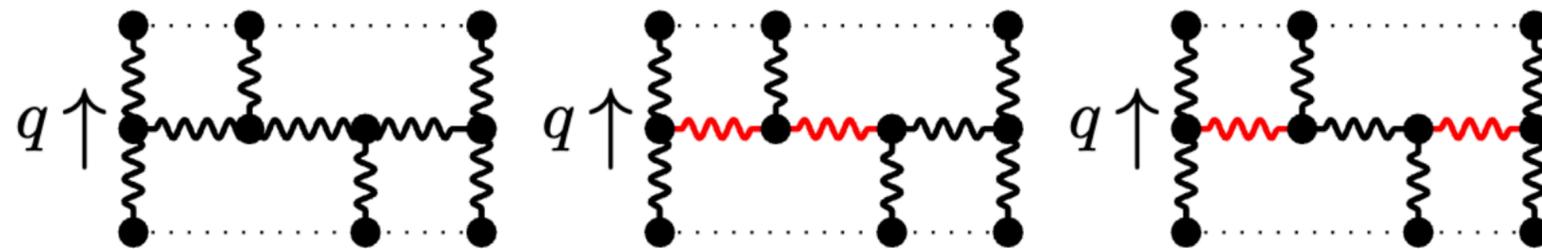
Boundary conditions: Expansion by regions

$$\ell_P \sim (l_0, \vec{\ell}) \sim (1-x, 1)$$

$$\ell_R \sim (l_0, \vec{\ell}) \sim (1-x, 1-x)$$

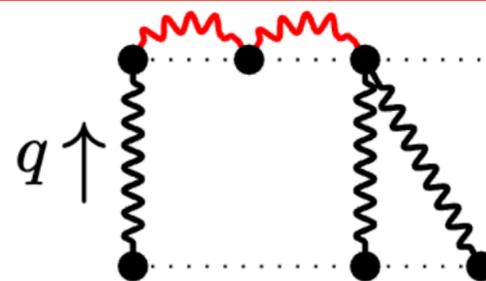


Expansion by regions

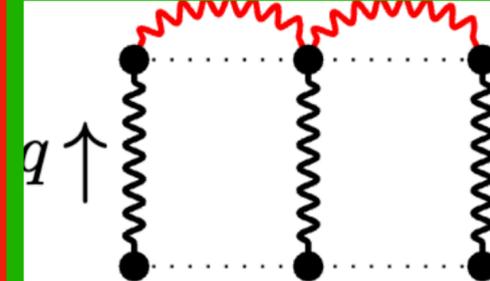


(a) Potential region

$$\propto \frac{1}{\epsilon} + \frac{1}{\sqrt{\gamma - 3}}$$



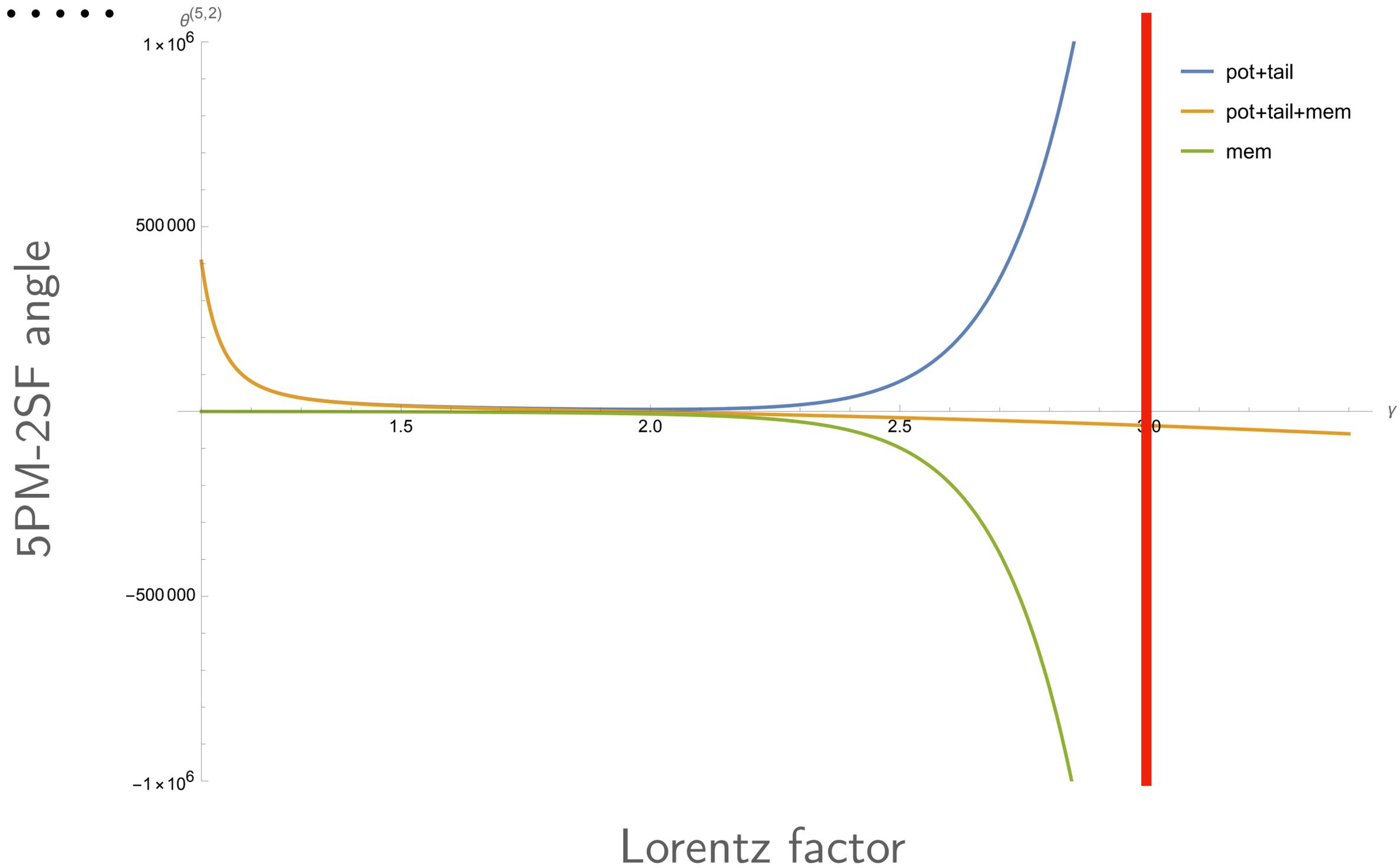
(b) Tail region



(c) Memory region

$$\propto -\frac{1}{\epsilon} - \frac{1}{\sqrt{\gamma - 3}}$$

Cancelations



Functions

- Iterated integrals over the following:

$$\left\{ \frac{1}{x}, \frac{1}{1-x^2}, \frac{x}{1-x^2}, \frac{1}{1+x^2}, \frac{x}{1+x^2}, \frac{1-x^2}{1+x^2+x^4}, \frac{x+2x^3}{1+x^2+x^4}, \frac{(1+x^2)\varpi_{\mathbb{K}3}(x)}{x}, \frac{1}{xr(x)\varpi_{\mathbb{K}3}(x)} \right\}$$

Where $r(x)$ is $\sqrt{-1+34x^2-x^4}$

- Calabi-Yau 3-folds (even more complicated) drop out

Contributions

- Agreement with SF/NR
- First analytic 2SF results
- Cancellation of Calabi-Yaus
- Double cancellation of poles

What now?

- 2SF dissipative — the full answer coming soon™
- Comparisons!
- Analytic continuation to bound states!
- Building waveform models!

Conclusion

“Lost in Translation is a great conference that brings together researchers from different fields. I am very hopeful for the future of analytical gravitational-wave physics in the LISA/Einstein Telescope era.”

- Albert Einstein

