

V_{us} DETERMINATION(S) USING INCLUSIVE
HADRONIC τ DECAY DATA

with

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OUTLINE

- V_{us} from inclusive $FB \tau$ FESRs
 - * The “conventional” implementation and the $> 3\sigma$ low $|V_{us}|$ puzzle
 - * OPE issues: continuum and lattice investigations
 - * A new implementation strategy (and $|V_{us}|$ results)
- A new approach to V_{us} using lattice and us $V+A \tau$ data

CONTEXT: VARIOUS V_{us} DETERMINATIONS

- From 3-family unitarity: Hardy-Towner 2014 $0^+ \rightarrow 0^+$ nuclear β decay analysis $|V_{ud}| = 0.97417(21)(?) \Rightarrow$

$$|V_{us}| = 0.2258(9)(?)$$

- From $K_{\ell 3}$: Current lattice $n_f = 2 + 1$, $n_f = 2 + 1 + 1$ $f_+(0)$ inputs yield

$$|V_{us}| = 0.2235(4)_{exp(9)latt} \quad (RBC/UKQCD \ 2015)$$

$$|V_{us}| = 0.2231(4)_{exp(7)latt} \quad (FNAL/MILC \ 2014)$$

$$|V_{us}| = 0.2230(4)_{exp(11)latt} \quad (ETM \ 2016)$$

- From $\Gamma[K_{\mu 2}]/\Gamma[\pi_{\mu 2}]$: Using lattice f_K/f_π , HT2014 $|V_{ud}|$ input

* FLAG 2015 $n_f = 2 + 1 + 1$, $n_f = 2 + 1$ assessments
(both $f_K/f_\pi = 1.194(5)$) \Rightarrow

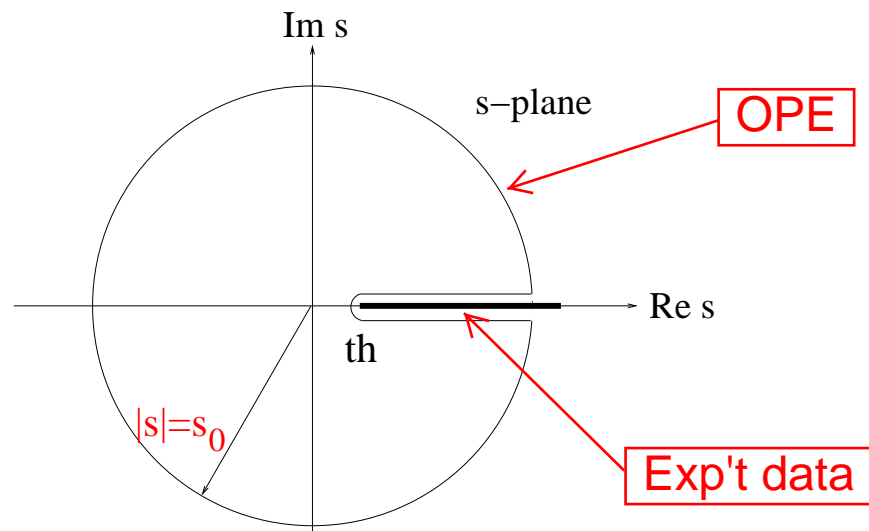
$$|V_{us}| = 0.2250(4)_{exp}(9)_{latt}$$

- C.f. “conventional” FB τ kinematic-weight FESR determination, inclusive ud , us BF input [Gamiz et al.]
 - * With HFAG input: $|V_{us}| = 0.2176(19)_{exp}(10??)_{th}$
(Passemar, CKM14 τ $|V_{us}|$ summary)
 - * 3.4σ low c.f. 3-family unitarity expectations [interesting **if real**, but theory systematics?]
 - * In this talk:
 - Continuum, lattice input on OPE input/treatment issues, lattice to quantify OPE errors
 - Alternate implementation of FB FESR analysis with continuum, lattice lessons in mind

REMINDER RE V_{us} FROM FB FESRs

- Basic theoretical tool: FESRs (*Cauchy's Thm*)
- For any s_0 , analytic $w(s)$, kinematic-singularity-free Π :

$$\int_{s_{th}}^{s_0} ds w(s) \rho(s) = \frac{-1}{2\pi i} \oint_{|s|=s_0} ds w(s) \Pi(s)$$



- **Inclusive FB τ $|V_{us}|$ determination:** From FESRs for V,A current 2-pt function polarizations $\Pi_{ud-us;V+A}^{(J)}(Q^2)$

* **Experimental spectral input:** Scaled spectral functions $|V_{ij}|^2 \rho_{ij;V/A}^{(J)}(s)$ from $dR_{ij;V/A}/ds$, with $R_{ij;V/A} \equiv \frac{\Gamma[\tau \rightarrow \nu_\tau \text{ hadrons}_{ij;V/A}(\gamma)]}{\Gamma[\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e(\gamma)]}$ (+ small us $J = 0$ subtraction) **[SM “kinematic weight” w_τ]**

* **$R_{ij;V/A}^w(s_0)$:** Re-weighted $R_{ij;V/A}$ analogue, integrated to variable upper endpoint s_0 in spectrum

$$R_{ij;V/A}^w(s_0) \sim \int_{th}^{s_0} ds \frac{dR_{ij;V/A}}{ds} \frac{w(s/s_0)}{w_\tau(s/m_\tau^2)}$$

* FESR for $\delta R^w \equiv \frac{R_{ud;V+A}^w}{|V_{ud}|^2} - \frac{R_{us;V+A}^w}{|V_{us}|^2}$ yields

$$|V_{us}| = \sqrt{\frac{R_{us;V+A}^w(s_0)}{\frac{R_{ud;V+A}^w(s_0)}{|V_{ud}|^2} - [\delta R^w(s_0)]^{OPE}}}$$

valid for arbitrary analytic $w(y)$, $s_0 \leq m_\tau^2$

* “Self-consistency tests”:

- $|V_{us}|$ independent of s_0 , w (tests control, understanding of OPE, experimental systematics)
- E.g., integrated $D = 2k + 2$ OPE $\sim 1/s_0^k$: errors in higher D treatment $\leftrightarrow s_0$ -instability

- Below: FESRs with variable s_0 , w , $\Delta\Pi_\tau \equiv \Pi_{ud-us;V+A}^{(0+1)}$

- The conventional implementation:

- * $w = w_\tau$ only, $s_0 = m_\tau^2$ only

- * w_τ degree 3 \Rightarrow OPE contributions up to $D = 8$; $D = 2$, $D = 4$ known, $D = 6$ OPE estimated with “VSA” (very small), $D = 8$ neglected

- * Resulting $[\delta R^{w_\tau}(m_\tau^2)]^{OPE}$ ($D \geq 2$) estimate a few to several % of individual $R_{ud,us;V+A}$

- \Rightarrow modest OPE accuracy enough for precision $|V_{us}|$

- * Advantage of $w = w_\tau$, $s_0 = m_\tau^2$ restriction: $R_{ud,us;V+A}$ from inclusive ud , us BFs only [BUT also precludes $w(s)$ -, s_0 -independence tests \Rightarrow price to pay in unquantified potential systematic errors]

- Testing the conventional implementation

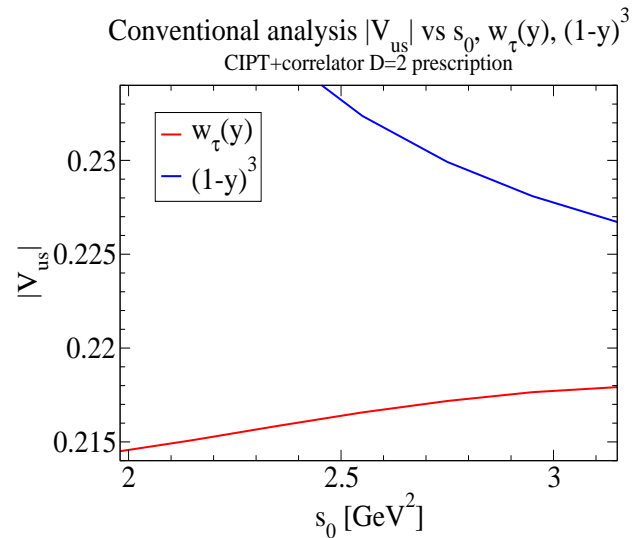
- * Variable $s_0 \leq m_\tau^2$, $|V_{us}|$ s_0 -stability check

- * A (targeted) $|V_{us}|$ weight-independence check:

- $\hat{w}(y)$ c.f. $w_\tau(y)$ $\Delta\Pi_\tau$ $|V_{us}|$ results, $y = (s/s_0)$,
 $w_\tau(y) = 1 - 3y^2 + 2y^3$, $\hat{w}(y) = 1 - 3y + 3y^2 - y^3$

- ($D = 6, 8$ for \hat{w} : $-1, -1/2$ \times those for w_τ)

- **IF** $D > 4$ assumptions OK for w_τ FESR \Rightarrow also for \hat{w} FESR $\Rightarrow |V_{us}|$ agreement. **If NOT, opposite-sign s_0 -instabilities for w_τ, \hat{w} FESRs, decreasing with s_0 for both**



Which is it? You decide ...

- Candidate self-consistency problem sources:
 - * **Experiment:** The less-well-known us distribution
 - * **Theory:** Conventional $D > 4$ assumptions (w_τ , \hat{w} comparison); slow $D = 2$ OPE convergence

- The $D = 2, 4$ series and slow $D = 2$ convergence

$$\left[\Delta \Pi_\tau(Q^2) \right]_{D=2}^{OPE} = \frac{3}{2\pi^2} \frac{\bar{m}_s^2}{Q^2} \left[1 + 2.333\bar{a} + 19.933\bar{a}^2 + 208.746\bar{a}^3 + \dots \right]$$

$$\left[\Delta \Pi_\tau(Q^2) \right]_{D=4}^{OPE} = \frac{[\langle m_\ell \bar{\ell} \ell \rangle - \langle m_s \bar{s} s \rangle]}{Q^4} \left(2 - 2\bar{a} - \frac{26}{3}\bar{a}^2 \right)$$

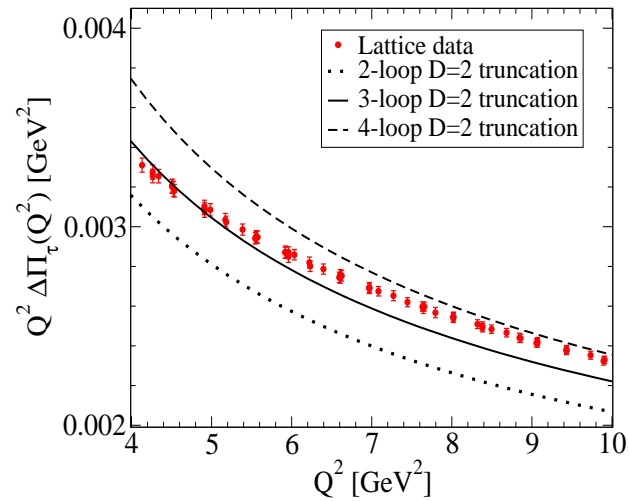
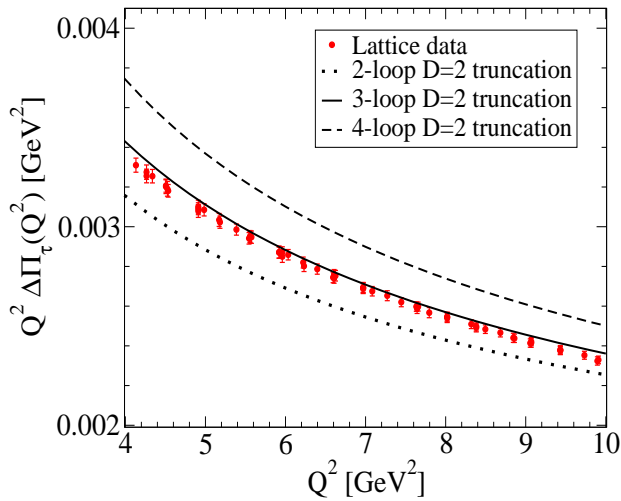
with running \overline{MS} quantities $\bar{a} = \frac{\alpha_s(Q^2)}{\pi}$, $\bar{m}_s = m_s(Q^2)$

- * $\bar{a}(m_\tau^2) > 0.10 \Rightarrow$ slowly converging $D = 2$ series at ALL scales accessible in τ decay
- * Slow $D = 2$ convergence a potential issue for conventional $D = 2$ truncation error estimates

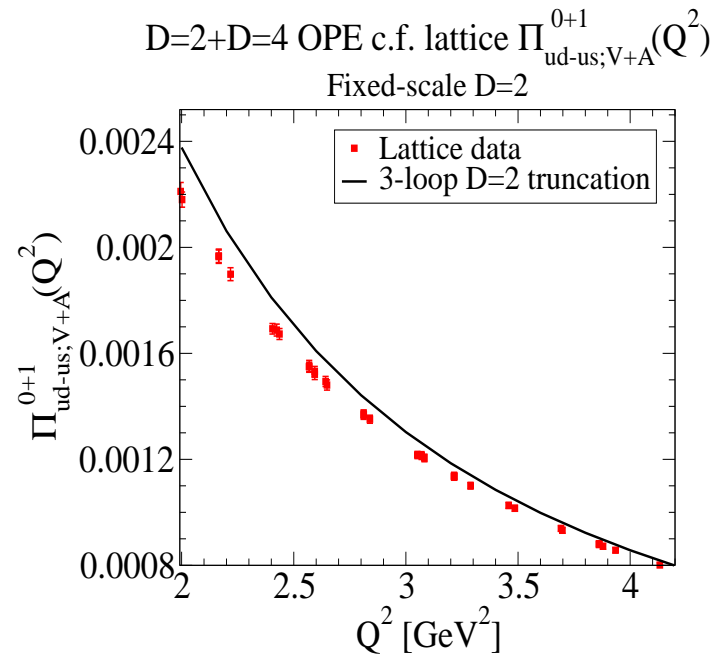
- More on the $D > 4$ assumptions
 - * No natural rapid convergence with D for general $w(s)$
 - * VSA ($D = 6$) known to be VERY crude, channel-dependent from ud V, A studies
 - * Strong cancellation in individual ud, us V+A sums in VSA
 - * Even stronger cancellation in resulting $ud - us$ VSA difference
 - * w_τ, \hat{w} comparison STRONGLY suggests problem(s) with $D > 4$ assumptions

LATTICE RE $D = 2$ SERIES, $D > 4$ OPE ISSUES

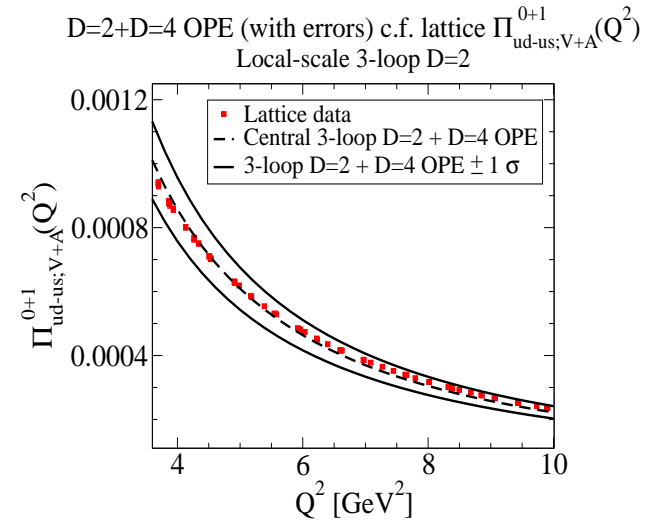
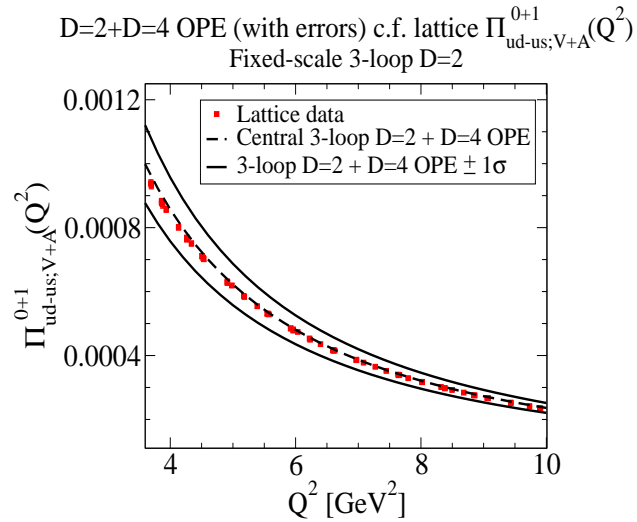
- OPE c.f. RBC/UKQCD lattice $\Delta\Pi_\tau(Q^2)$
 - * Here: lightest m_π , fine ($1/a = 2.38 \text{ GeV}$) $32^3 \times 64$ $2 + 1$ ensemble ($m_\pi \sim 300 \text{ MeV}$, $m_\pi L \sim 4.1$)
 - * Tight cylinder cut for continuum correlator behavior (width from stabilization in ud V channel α_s study)
 - * Large Q^2 : $D = 2 + 4$ OPE exploration; lower Q^2 : re possible non-negligible $D > 4$ at τ decay scales
 - * Fixed- or local-scale $D = 2$? [c.f. “FOPT” vs “CIPT” issue for FESRs]



- Higher Q^2 : best (excellent) lattice vs $D = 2 + 4$ OPE match for 3-loop-truncated, fixed-scale $D = 2$
- Fixed scale suggests FOPT for FESR $D = 2$



- Onset of $D > 4$ contributions below $\sim 4 \text{ GeV}^2$



- Standard $D = 2$, $D = 4$ error estimates conservative, despite very slow convergence of the $D = 2$ series

A NEW IMPLEMENTATION STRATEGY BASED ON LATTICE/CONTINUUM LESSONS

- No assumptions re $D > 4$; include contributions and fit effective condensates C_D as part of analysis
- 3-loop-truncated FOPT for $D = 2$ contribution; standard error estimates for $D = 2 + 4$
- Use range of s_0 and fit both $|V_{us}|$, C_D
- Fits simplest for $w(y) = w_N(y) = 1 - \frac{y}{N-1} + \frac{y^N}{N-1}$ (C_{2N+2} , $|V_{us}|$ as only fit parameters)
- Self-consistency check: $|V_{us}|$ from different w_N

OPE, SPECTRAL INPUT

- PDG, FLAG, HPQCD input for $D = 2, 4$ OPE
- ud V+A spectral data from ALEPH 2013
- us V+A spectral data from sum over exclusive modes [$> 90\%$ of B_{us}^{TOT} from $K_{\ell 2}$, Belle, BaBar $K\pi$, $K\pi\pi$, $3K$ results; residual: 1999 ALEPH]
- Here, for brevity, $K\pi$ normalization including preliminary BaBar $B[\tau \rightarrow K^- \pi^0 \nu_\tau]$ update (Adametz thesis)

RESULTS OF THE NEW ANALYSES

- Fitted $|V_{us}|$ (as expected) lies between (s_0 -unstable) w_τ and \hat{w} conventional implementation results
- Fitted C_D show FB cancellation in comparison to fitted ud $V+A$ analogues (qualitative self-consistency test)
- Excellent stability of both $|V_{us}|$, C_{2N+2} , wrt variation of fit window size; excellent agreement of central $|V_{us}|$ from different w_N

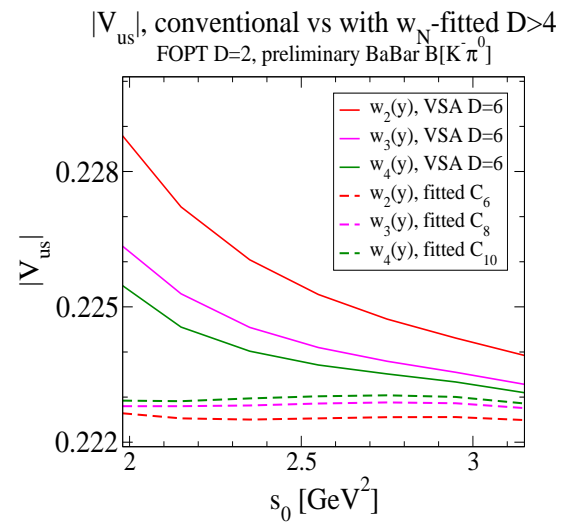
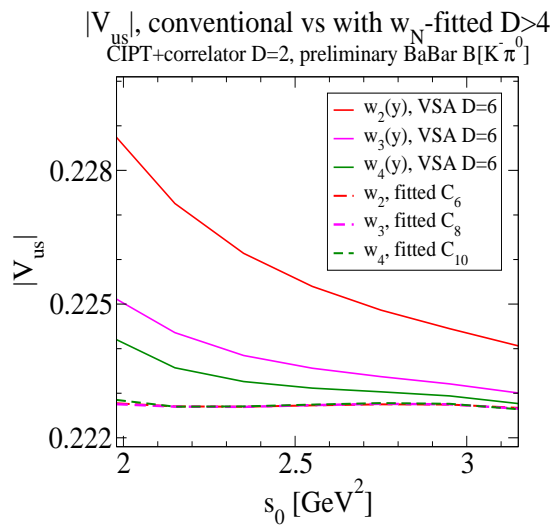
- Results using preliminary $K\pi$ BF update

Weight	$ V_{us} $	$ V_{us} $
	CIPT+corr $D = 2$	FOPT $D = 2$
w_2	0.22271(228)	0.22252(228)
w_3	0.22271(228)	0.22282(228)
w_4	0.22271(229)	0.22296(229)

Error budget, 3-loop-truncated FOPT $D = 2$ fits

Source	$\delta V_{us} $ (w_2 FESR)	$\delta V_{us} $ (w_3 FESR)	$\delta V_{us} $ (w_4 FESR)	
$\delta\alpha_s$	0.00001	0.00004	0.00004	
$\delta m_s(2 \text{ GeV})$	0.00017	0.00019	0.00019	
$\delta\langle m_s \bar{s}s \rangle$	0.00035	0.00035	0.00035	← Theory
$\delta(\text{long corr})$	0.00009	0.00009	0.00009	
ud exp	0.00027	0.00028	0.00028	
us exp	0.00226	0.00227	0.00227	← Exp't

Fit quality: s_0 -stability for conventional implementation, except with fitted C_{2N+2}



SUMMARY, CONVENTIONAL FB FESR ANALYSIS

- Continuum, lattice \Rightarrow conventional $D > 4$ assumptions untenable, must fit $D > 4$ effective condensates
- Fitting $D > 4$ resolves s_0 -, $w(y)$ -dependence problems
- Fitted $D > 4$, when included, raise $|V_{us}|$ by ~ 0.0020
- $D > 4$ fit needs variable $s_0 \Rightarrow$ **BF-only approach untenable** ($\sim .0020 - 0.0030$ additional systematic)
- Lattice \Rightarrow 3-loop-truncated FOPT for $D = 2$ series

- Lattice $\Rightarrow D = 2 + 4$ OPE error estimate conservative
- $|V_{us}|$ errors then strongly dominated by us experimental uncertainties, subject to experimental improvement
- New preliminary $K^- \pi^0$ BF normalization, 3-weight combined fit result $|V_{us}| = 0.2228(23)_{exp}(5)_{th}$ in agreement with other determinations (especially $K_{\ell 3}$)
(c.f. $0.2200(23)_{exp}(5)_{th}$ with 2014 HFAG $B[K^- \pi^0]$)
- Theory error (conservative) ~ 0.0005 c.f. 0.0009 for $K_{\ell 3}$, $\Gamma[K_{\mu 2}]/\Gamma[\pi_{\mu 2}]$ (lattice crucial here)
- \Rightarrow **FB FESR method competitive with sufficient improvement to us experimental errors**

FUTURE PROSPECTS

- Current us spectral integral errors δB normalization uncertainty dominated \Rightarrow improvement through BF error reduction (*B-factory $K^-\pi^0\pi^0$ distribution useful too*)
- An upcoming alternate to the FB FESR approach
 - * KM, Hudspith, Izubuchi, Lewis, Ohki + RBC/UKQCD
 - * Generalized dispersion rel'ns for products of us $V+A$ polarization, weights with poles at Euclidean Q^2
 - * Theory: Lattice us 2-point function data (no OPE)
 - * Experiment: $dR_{us;V+A}/ds$ (no $J = 0$ subtraction)

More on the lattice-inclusive us τ approach

- Experimental $dR_{us;V+A}/ds$ distribution yields $|V_{us}|^2 \tilde{\rho}(s)$,

$$\tilde{\rho}_{us;V+A}(s) \equiv \left(1 + 2\frac{s}{m_\tau^2}\right) \rho_{us;V+A}^{(J=1)}(s) + \rho_{us;V+A}^{(J=0)}(s)$$

(no continuum us $J = 0$ subtraction)

- Associated (kinematic-singularity-free) polarization

$$\tilde{\Pi}_{us;V+A}(Q^2) \equiv \left(1 - 2\frac{Q^2}{m_\tau^2}\right) \Pi_{us;V+A}^{(J=1)}(Q^2) + \Pi_{us;V+A}^{(J=0)}(Q^2)$$

- For weights $w_N(s) \equiv \frac{1}{\prod_{k=1}^N (s+Q_k^2)}$, $N \geq 3$, obtain convergent, unsubtracted 'dispersion relation'

$$\int_{th}^{\infty} ds w_N(s) \tilde{\rho}_{us;V+A}(s) = \sum_{k=1}^N \frac{\tilde{\Pi}_{us;V+A}(Q_k^2)}{\prod_{j \neq k} (Q_j^2 - Q_k^2)}$$

- * Lattice data for $\tilde{\Pi}_{us;V+A}(Q_k^2)$ on RHS
- * LHS from experimental $dR_{us;V+A}/ds$, up to $|V_{us}|^2$
- * $w_N(s)$: rapid fall-off in s if all $Q_k^2 < 1 \text{ GeV}^2$ ($\Rightarrow K, K\pi$ dominate spectral integral, near-endpoint multi-particle, $s > m_\tau^2$ contributions strongly suppressed)
- * Increasing $\{Q_k^2\}$: decreasing lattice $\tilde{\Pi}_{us;V+A}(Q_k^2)$ error, increasing experimental error

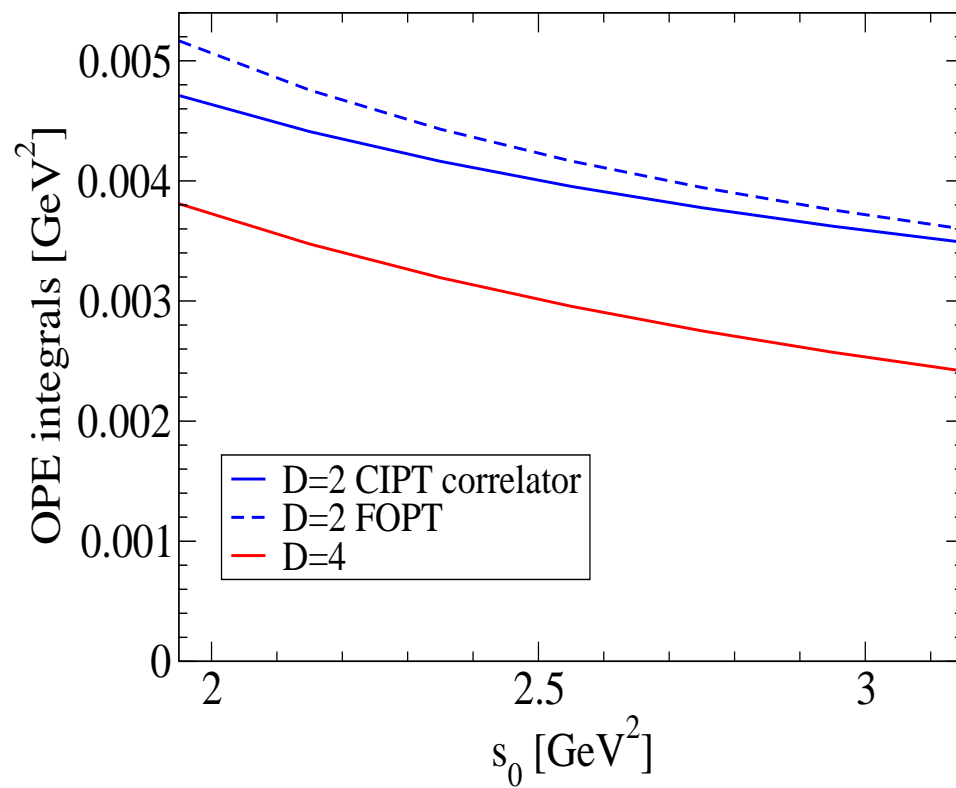
- An illustration of expected precision: $N = 4$, $\{Q_k^2\} = \{0.25, 0.30, 0.35, 0.40 \text{ GeV}^2\}$
 - * Spectral integral error $\sim 0.5\%$
 - * Statistical error, sum of residues for near-physical-point ensembles: $\sim 1/3\%$ for $48^3 \times 96$, $1/a = 1.73 \text{ GeV}$, $\sim 1/2\%$ for $64^3 \times 128$, $1/a = 2.36 \text{ GeV}$
 - * Quadrature sum $\Rightarrow \delta|V_{us}| \sim 0.0007, \sim 0.0008$ for near-physical-point $48^3, 64^3$ ensembles, respectively
- Optimization/systematics under current study

(THE END)

BACKUP SLIDES

- Are $D = 6, 8$ OPE contributions likely to be small?
 - * $D = 4 \ll D = 2$ for $w_\tau(y) = 1 - 3y^2 + 2y^3$, $y = s/s_0$, “accidental” [$O(\alpha_s^2)$ suppression due to absence of term linear in y in $w_\tau(y)$]
 - * Comparison of $D = 2, 4$ OPE contributions for $w(y) = (1 - y)^2$ (a case without this suppression) to see natural relative sizes

D=2, 4 OPE integrals
for weight $w(y)=(1-y)^2 (y=s/s_0)$



- Spectral functions from hadronic τ decay distributions

With $R_{ij;V/A} \equiv \frac{\Gamma[\tau \rightarrow \nu_\tau \text{ hadrons}_{ij;V/A}(\gamma)]}{\Gamma[\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e(\gamma)]}$

$$\frac{dR_{V/A;ij}}{ds} = \frac{12\pi^2 |V_{ij}|^2 S_{EW}}{m_\tau^2} \left[w_\tau(y_\tau) \rho_{ij;V/A}^{(0+1)}(s) + w_L(y_\tau) \rho_{ij;V/A}^{(0)}(s) \right]$$

$$w_\tau(y) = (1-y)^2(1+2y), w_L(y) = -2y(1-y)^2$$

$$y_\tau = s/m_\tau^2$$

More on the $D > 6$ VSA guesstimate

- VSA VERY crude: sizable (e.g. $\sim 4 - 5$), channel-dependent violation in ud V, A channels
- Crudeness \Rightarrow double very close cancellation (in individual ud , us V+A sums, and in FB $ud - us$ difference) dangerous to rely on
- Integrated $D > 4 \ll$ than integrated $D = 2$ for w_τ (no $O(y)$ term) but comparable for other $w(y) \Rightarrow$ smallness of doubly-cancelled $D = 6$ VSA not “natural”

MORE ON THE us DATA

- K pole via $f_K|V_{us}|$ from $K_{\ell 2}$
- Rather precise unit-normalized $K^-\pi^0$, $\bar{K}^0\pi^-$, $K^-\pi^+\pi^-$, $\bar{K}^0\pi^-\pi^0$, $3K$ distributions from Belle, BaBar (main uncertainties from BFs)
- K , B-factory modes over 90% of B_{us}^{TOT}
- Residual us exclusive mode contributions from 1999 ALEPH data, covariances

Alternative $K\pi$ BF normalizations

- Existing HFAG $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau] = 0.0126$
- Existing $B[K^-\pi^0\nu_\tau] = 0.00431(15)$ value \rightarrow preliminary BaBar (Adametz thesis) result $0.00500(15)$ yields $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau] = 0.0133$
- Central $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau]$ from $K_{\ell 3}$, dispersion rel'n expectations [ACLP13] also 0.0133
- 0.07% difference “small” but represents $\sim 2.4\%$ of B_{us}^{TOT} , hence $\sim 1.2\%$ increase in $|V_{us}|$

Results for $|V_{us}|$ for current $K\pi$ BFs:

Weight	$ V_{us} $	
	CIPT+corr $D = 2$	FOPT $D = 2$
w_2	0.21985(230)	0.21966(230)
w_3	0.21985(231)	0.21966(231)
w_4	0.21985(231)	0.22009(231)

Error budget, existing $K\pi$ BFs

Source	$\delta V_{us} $ (w_2 FESR)	$\delta V_{us} $ (w_3 FESR)	$\delta V_{us} $ (w_4 FESR)
$\delta\alpha_s$	0.00001	0.00003	0.00005
$\delta m_s(2 \text{ GeV})$	0.00017	0.00018	0.00020
$\delta\langle m_s \bar{s}s \rangle$	0.00034	0.00034	0.00034
$\delta(\text{long corr})$	0.00009	0.00009	0.00009
ud exp	0.00027	0.00027	0.00027
us exp	0.00229	0.00229	0.00230

Stability of $|V_{us}|$ with fitted C_{2N+2} input, existing $K\pi$ BF normalization

