



Advancing (Non)perturbative Many-Body Methods For Strongly Correlated Nuclear Matter

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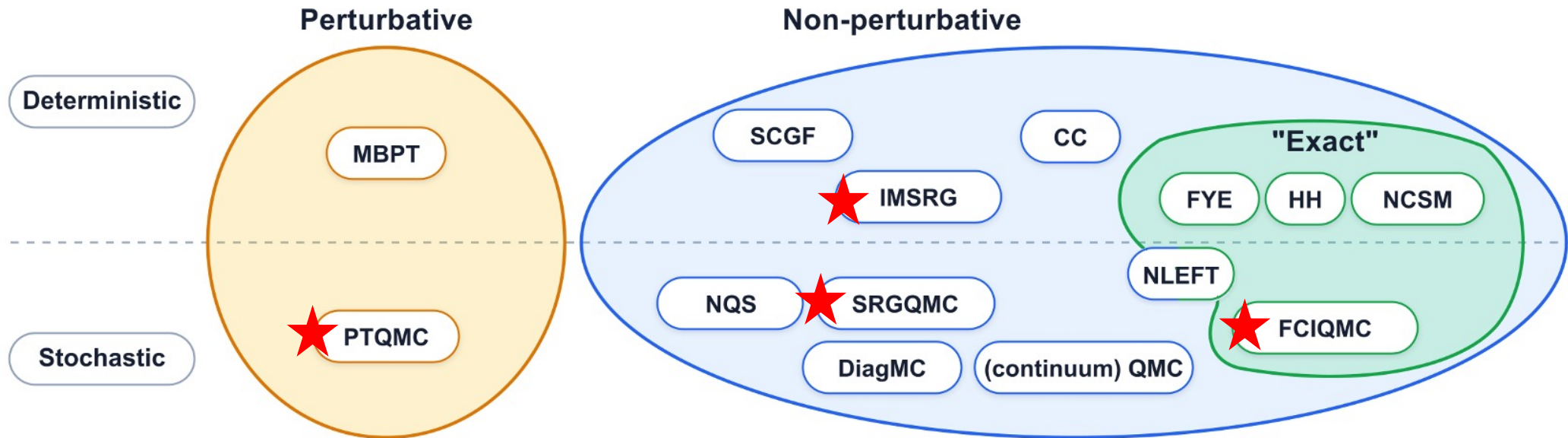
Many-Body Theory: *Nuclear Physics Meets Quantum Chemistry*

Deterministic and Stochastic Routes to Many-Body Correlations

$$\begin{aligned}\hat{H} = & E_0 + \sum_{pq} f_{pq} a_p^\dagger a_q \\ & + \frac{1}{4} \sum_{pqrs} \bar{\Gamma}_{pq,rs} a_p^\dagger a_q^\dagger a_s a_r \\ & + \frac{1}{36} \sum_{pqrst} \bar{W}_{pqr,stu} a_p^\dagger a_q^\dagger a_r^\dagger a_u a_t a_s.\end{aligned}$$

→ Normal-ordered with respect to a reference state, typically **Hartree–Fock**.

The challenge: capturing higher-order correlations beyond mean-field.

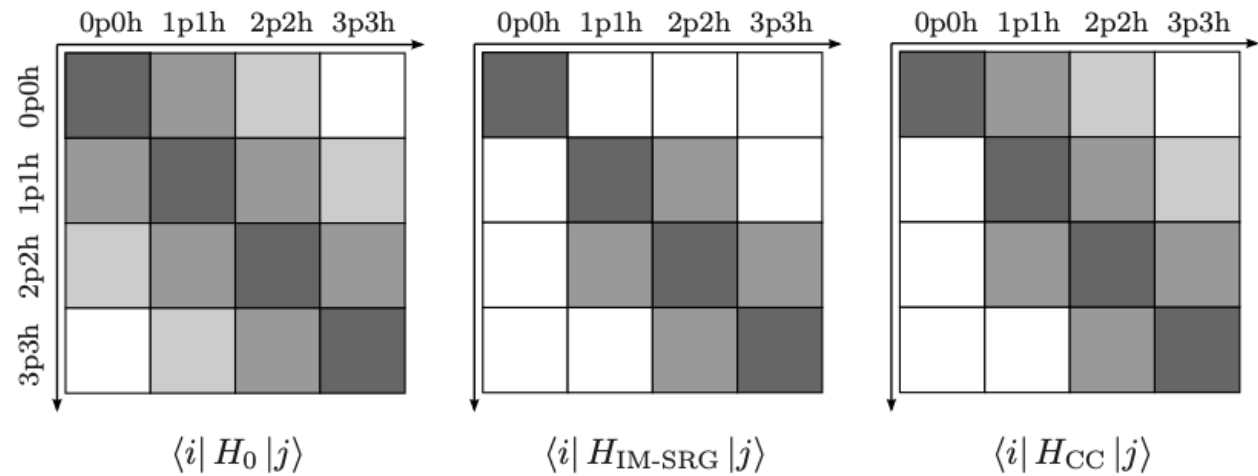


Complementary Routes: IMSRG & FCIQMC

IMSRG: Truncated Continuous Unitary Transformations (Scalable and Productive)

$$H(s) = U(s)H(0)U^\dagger(s)$$

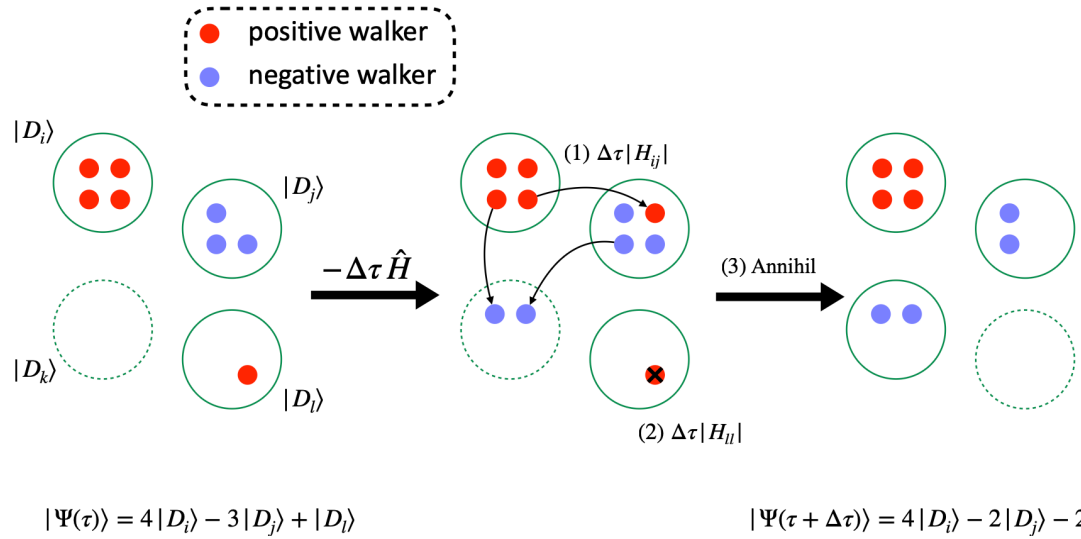
$$\frac{dH(s)}{ds} = [\eta(s), H(s)] \rightarrow \frac{dO(s)}{ds} = [\eta(s), O(s)]$$



Off-diagonal couplings suppressed.

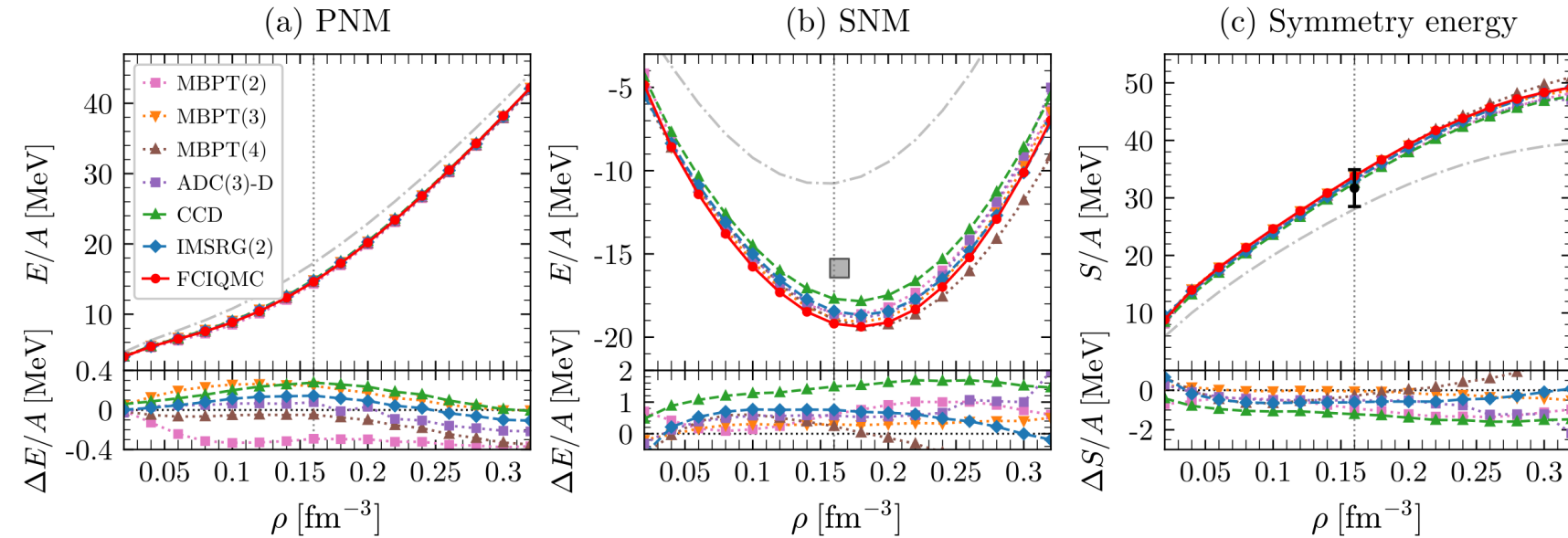
FCIQMC: Stochastic Imaginary-Time Projection (Accurate as a Benchmark)

$$N_i(\tau + \Delta\tau) - N_i(\tau) = -\Delta\tau(H_{ii} - S)N_i(\tau) - \Delta\tau \sum_{j \neq i} H_{ij}N_j(\tau)$$



Signed walkers stochastically project out the exact ground state.

For the Equation of State: More Correlations ... ?



X. Zhen, R. Hu, *et al.*

Phys. Lett. B, **862**, 139350 (2025).

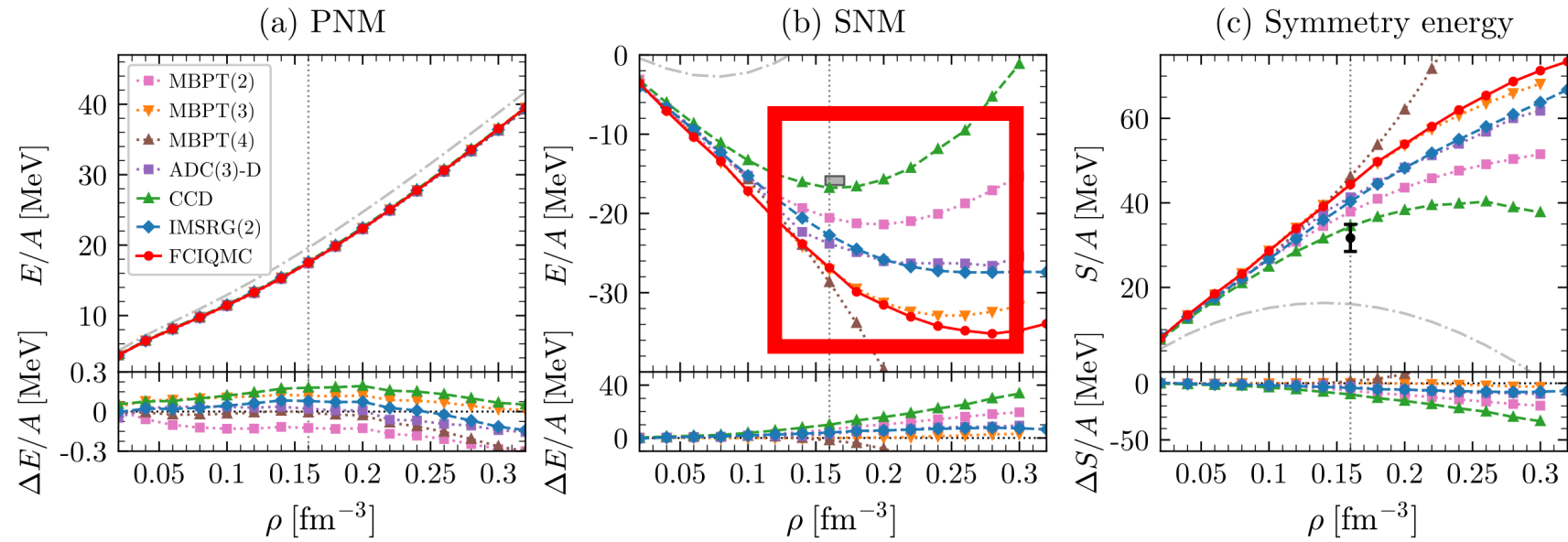
R. Hu, S. Jin, X. Zhen, *et al.*

Phys. Rev. Lett., **137**, 062503 (2026).

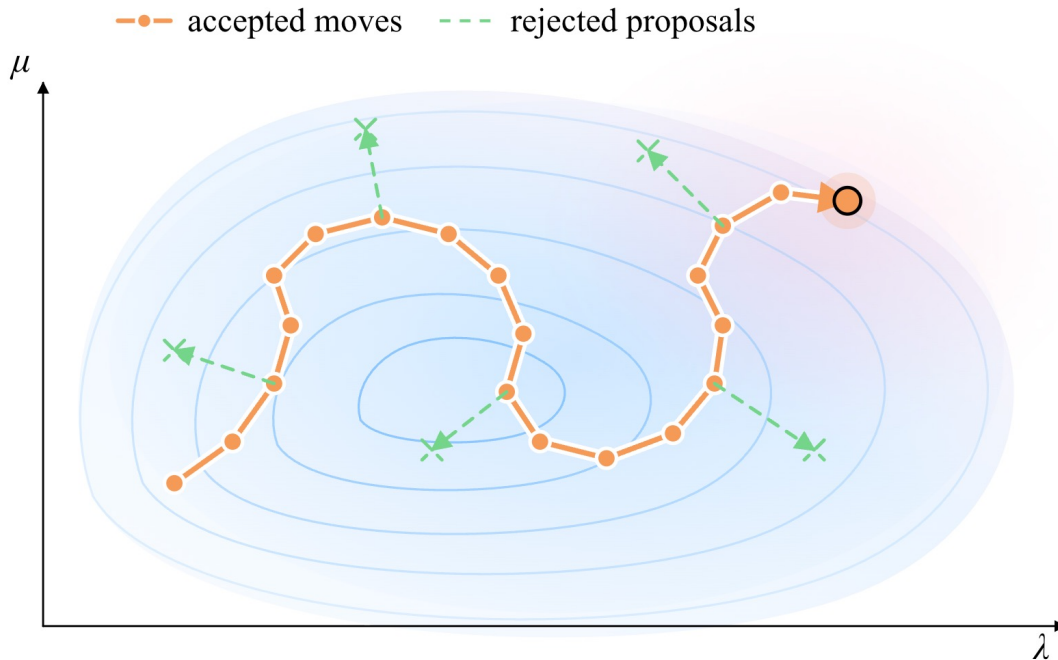
Similar Results for
Softer Forces...

But for Harder forces ...

Systematically
Underestimated!



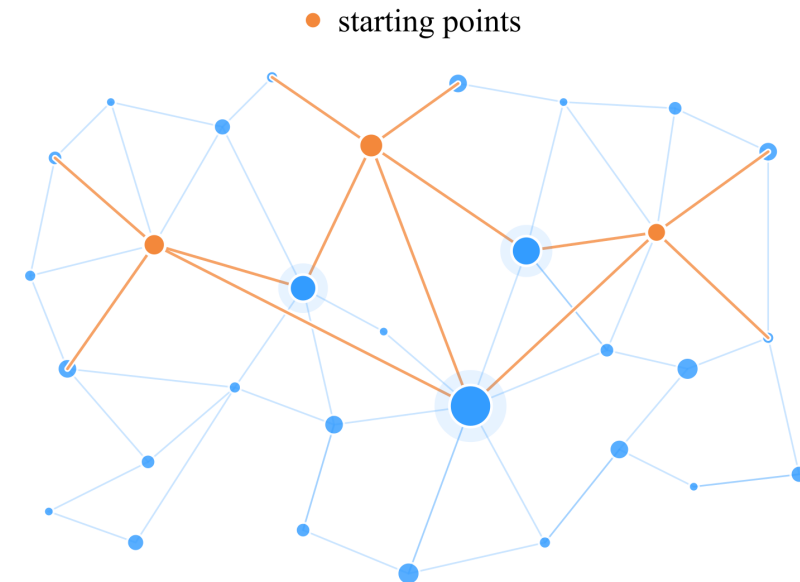
Continuous Sampling: Single point in high-dimensional space



Examples (~MC):

- Variational
- Green's Function
- Diffusion
- Auxiliary Field

Discrete Sampling: Distribution of discrete objects



Examples (~MC):

- Diagrammatic
- Full Configuration Interaction
- Perturbation Theory
- Similarity Renormalization Group

PT : Why do we need Monte Carlo ?

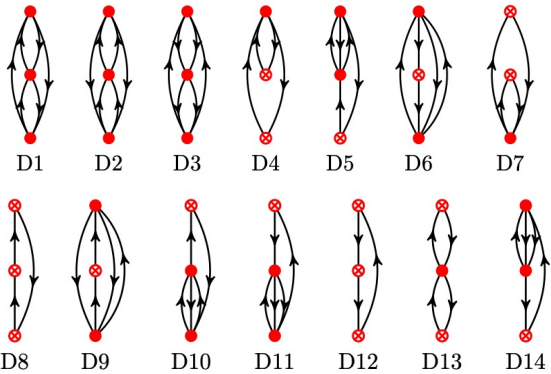
$$E^{(0)}|\Phi\rangle + \sum_{p=1}^{\infty} \lambda^p \left(H_1 |\Psi^{(p-1)}\rangle + H_0 |\Psi^{(p)}\rangle \right) = E^{(0)}|\Phi\rangle + \sum_{p=1}^{\infty} \lambda^p \left(\sum_{j=0}^p E^{(j)} |\Psi^{(p-j)}\rangle \right).$$

collect terms at the same order...

Configuration Space (computationally prohibitive) :

$$C_{m0}^{(p)} = \frac{1}{E^{(0)} - E_m^{(0)}} \left[\sum_q \langle \Phi_m | H_1 | \Phi_q \rangle C_{q0}^{(p-1)} - \sum_{j=1}^p E^{(j)} C_{m0}^{(p-j)} \right]$$

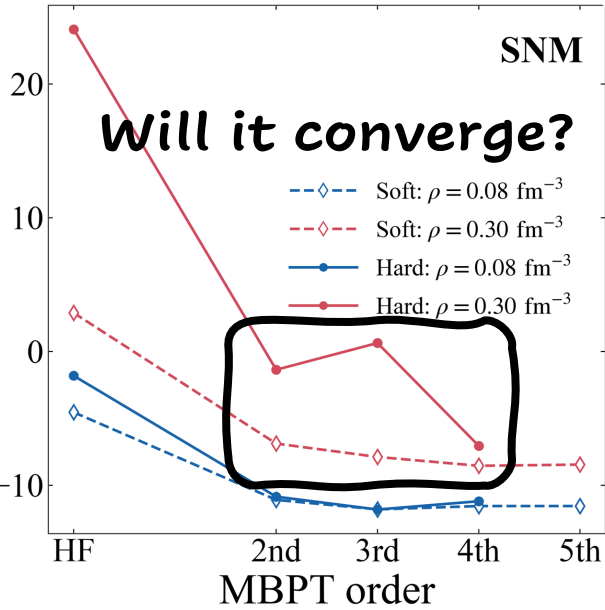
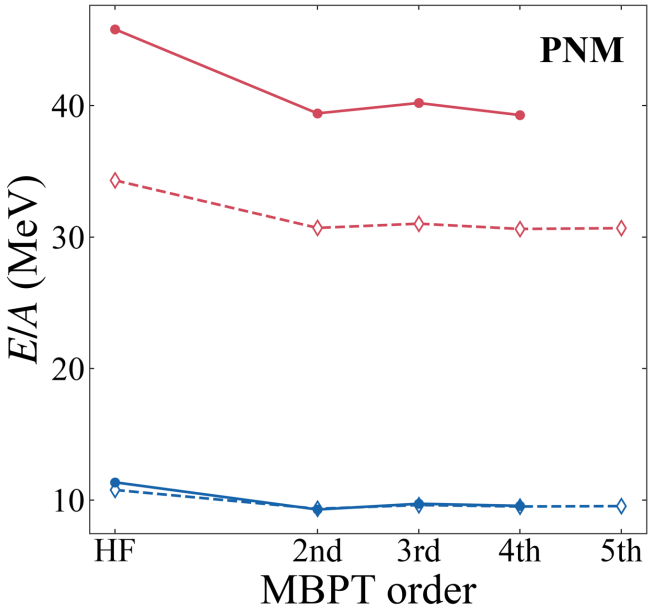
Operator Space (currently up to fifth order) :



思想自由 兼容并包

level	MBPT(<i>n</i>)	<i>N</i> _{diag} ^a	<i>N</i> _{diag} ^(eff)	<i>N</i> _{dim}	<i>N</i> _{grp}
two-body	HF	1	1	6	1
	2	1	1	9	1
	3	3	3	18	2
	4	39	24	24	4
	5	840	375	30	5
	6 ^b	27300	11269	36	6
three-body	2	1	1	18	1
	3	19	14 ^c	21; 24; 27	14

$$H(\lambda) \equiv H_0 + \lambda H_1, \\ E(\lambda) \equiv E^{(0)} + \lambda E^{(1)} + \lambda^2 E^{(2)} + \dots, \\ |\Psi(\lambda)\rangle \equiv |\Psi^{(0)}\rangle + \lambda |\Psi^{(1)}\rangle + \lambda^2 |\Psi^{(2)}\rangle + \dots$$

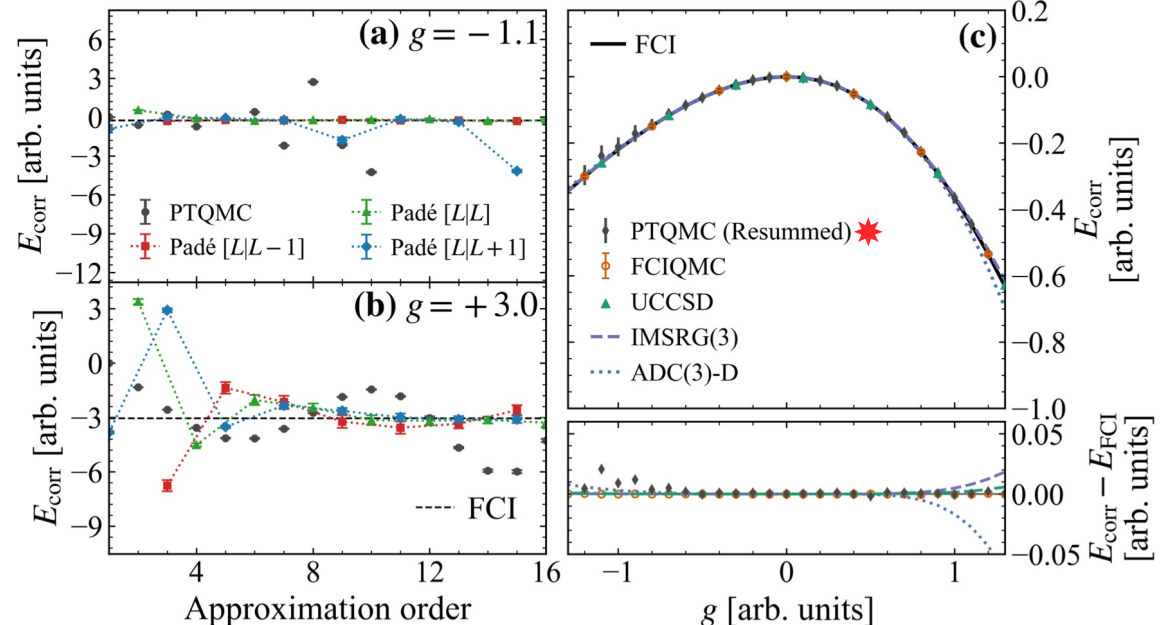
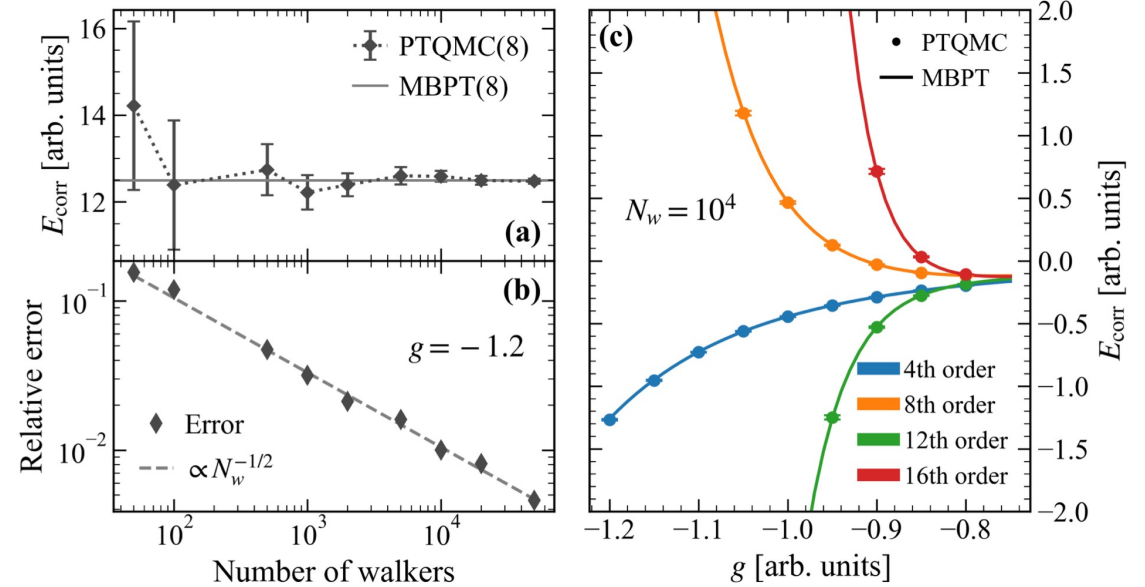
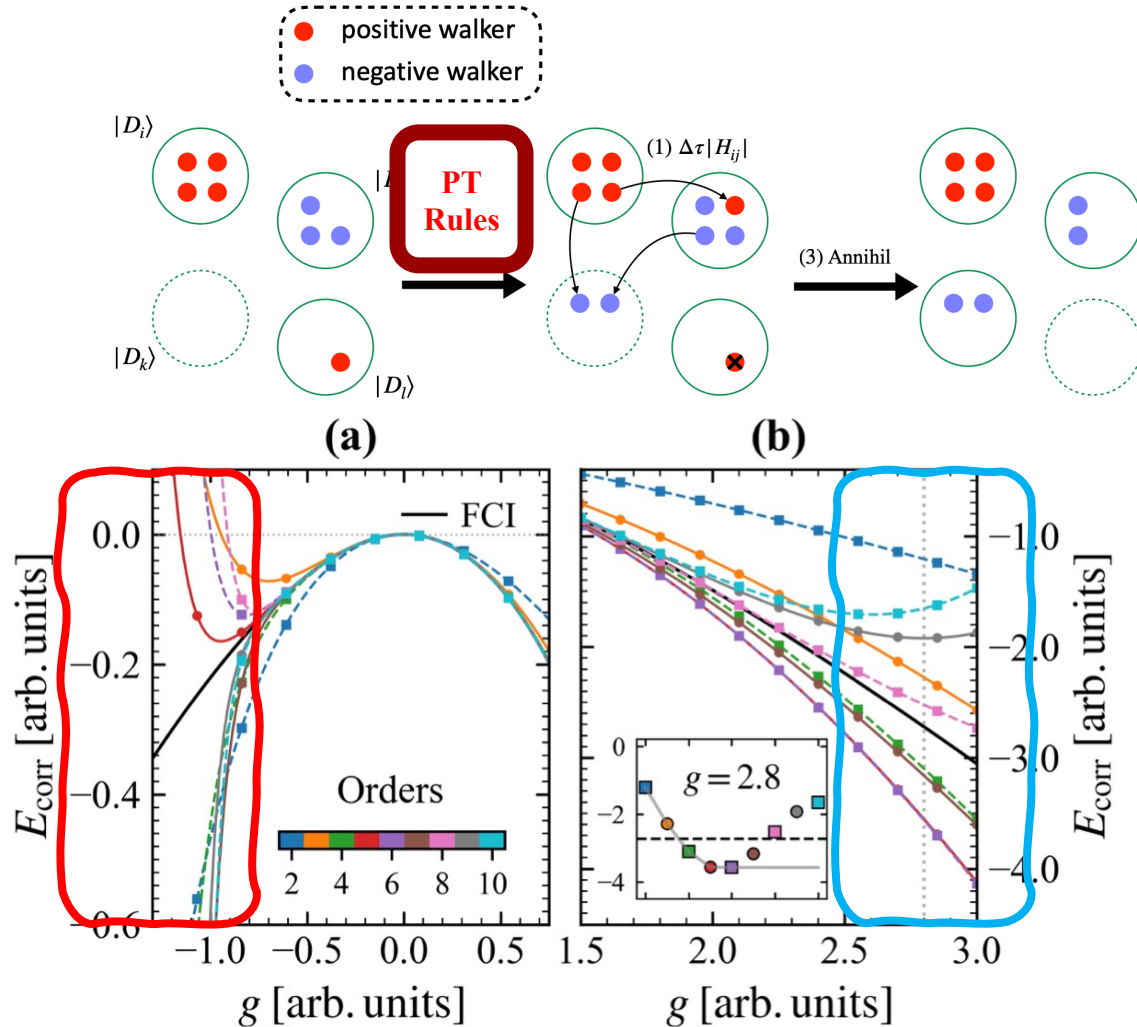


Can we go to even higher orders together with many-nucleon forces?

A. Tichai, R. Roth, *et al.* *Front. Phys.*, 8, 164 (2020).
C. Drischler, K. S. McElvain, *et al.* *arXiv:2603.24532* (2026).
Z. Li, A. Tichai, *et al.* *arXiv:2607.00119* (2026). < 6 >

Perturbation Theory Quantum Monte Carlo (PTQMC)

Recursive spawning among configurations

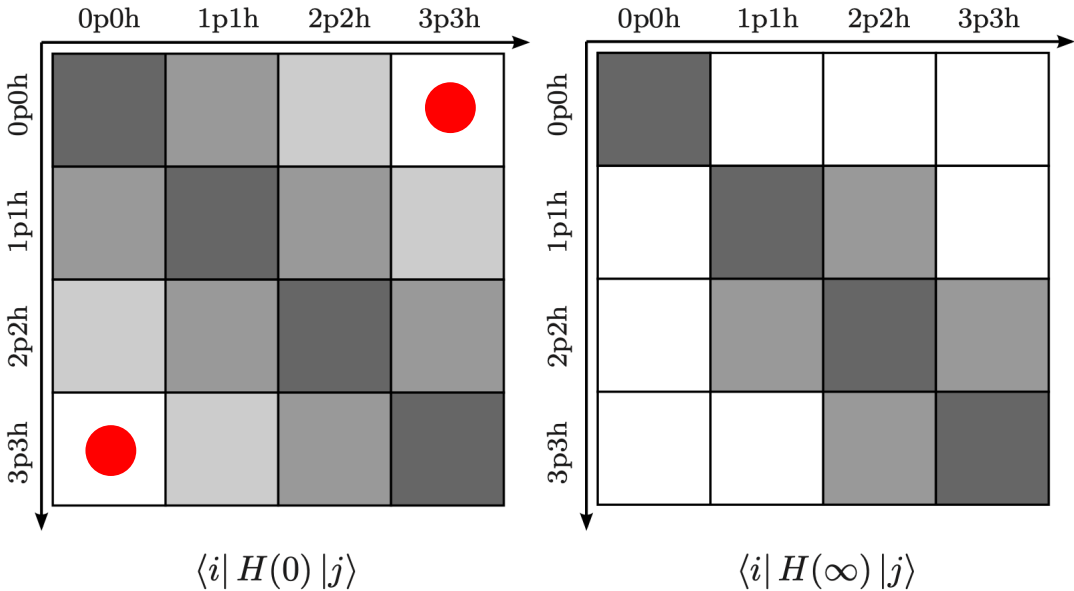


IMSRG : Can we go to higher orders?

H. Hergert, *et al. Phys. Rep.*, **621**, 165–222 (2016).

M. Heinz, *et al. Phys. Rev. C*, **103**, 044318 (2021).

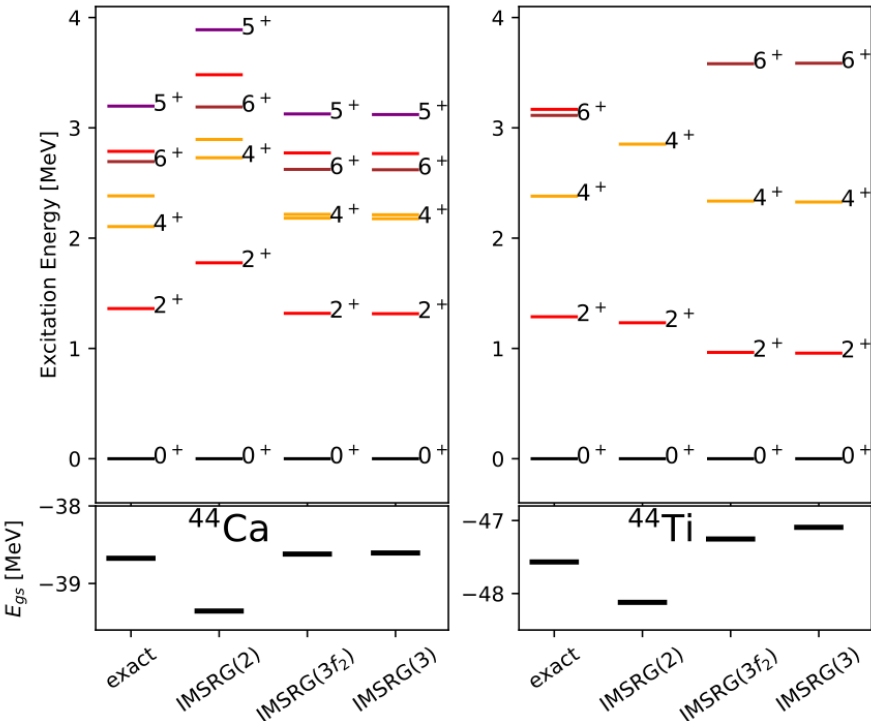
B. C. He, S. R. Stroberg, *Phys. Rev. C*, **110**, 044317 (2024).



Three-Body Forces:
Important in nuclear physics but extremely expensive.

IMSRG(3f2) $O(N^6)$
[2, [2,2] $\rightarrow$ 3] $\rightarrow$ 1, 2

Factorization for higher-order two-body correlations ! But...



Commutator	Cost	Included in ...				
		IMSRG(3)-MP4	IMSRG(3)- N^7	IMSRG(3)- N^8	IMSRG(3)- g^5	IMSRG(3)
[2, 2] $\rightarrow$ 3	$O(N^7)$	✓	✓	✓	✓	✓
[2, 3] $\rightarrow$ 2	$O(N^7)$	✓	✓	✓	✓	✓
[1, 3] $\rightarrow$ 3	$O(N^7)$	✓	✓	✓	✓	✓
[3, 3] $\rightarrow$ 0	$O(N^6)$	✓	✓	✓	✓	✓
[2, 3] $\rightarrow$ 1	$O(N^6)$		✓	✓	✓	✓
[1, 3] $\rightarrow$ 2	$O(N^6)$		✓	✓	✓	✓
[3, 3] $\rightarrow$ 1	$O(N^7)$		✓	✓		✓
[2, 3] $\rightarrow$ 3	$O(N^8)$			✓	✓	✓
[3, 3] $\rightarrow$ 2	$O(N^8)$			✓	✓	✓
[3, 3] $\rightarrow$ 3	$O(N^9)$					✓

Polynomial Scaling



Similarity Renormalization Group Monte Carlo (SRGQMC)

- Stochastic evolution of flow equations

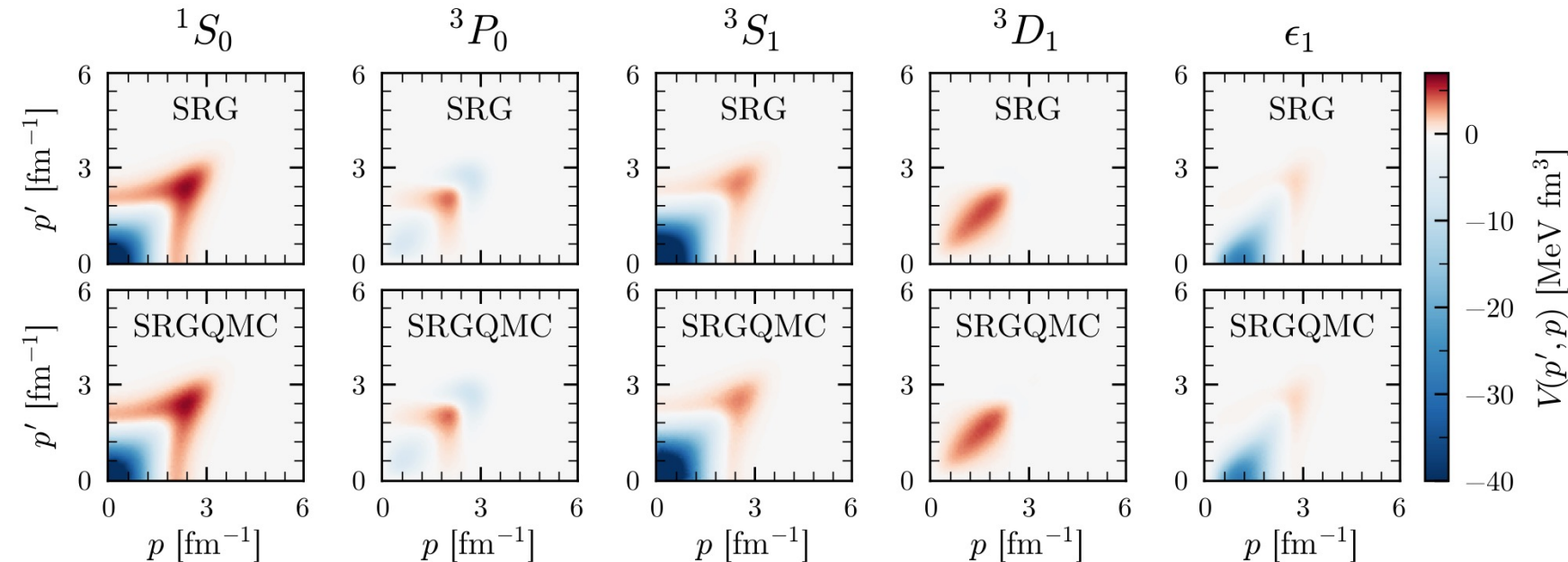
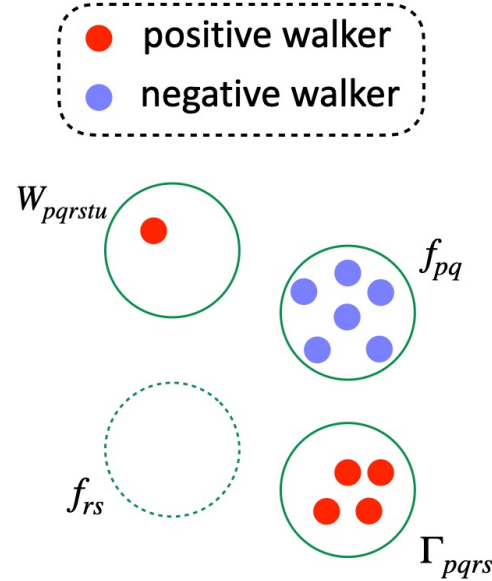
$$\frac{dC}{d\tau} = L(H)C$$

master equation ($\alpha = 1$)

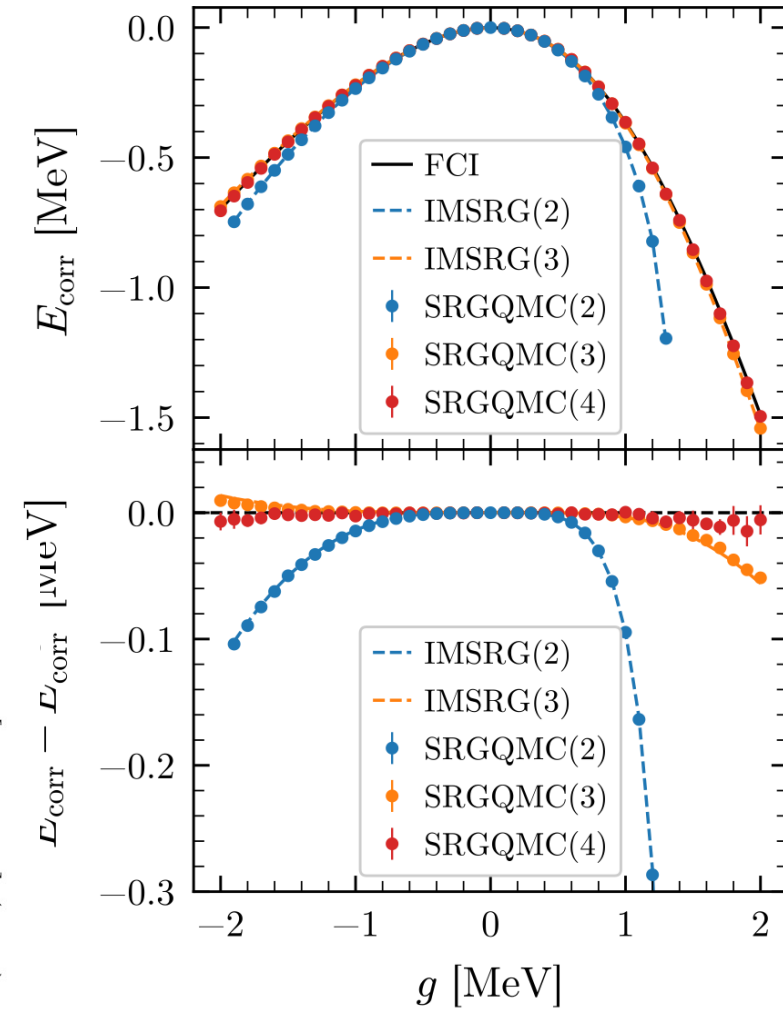
Source-free, homogeneous first-order ODEs !

flow equation ($\alpha = q + 1$)

$$\frac{dH_s}{ds} = \eta(H_s)H_s - H_s\eta(H_s), \quad \eta(\lambda H_s) = \lambda^q \eta(H_s)$$



R. Hu*, X. Zhen*, et al. arXiv:2607.08830 (2026).



More on the way !

We are closing the loop !

Requirements

Physics Targets

Controlled equation of state:

Predict symmetric and neutron matter with cross-benchmarked calculations.

Separated uncertainties:

Quantify errors from nuclear interactions and many-body truncations.

Many-body forces:

Investigate the impact of three- and higher-body forces, especially at high densities.

Many-body Techniques

FCIQMC — Exact benchmarks:
Establish reliable anchors for strongly correlated systems.

PTQMC — High-order perturbative:
Access very high orders and diagnose convergence statistically.

SRGQMC — High-order nonperturbative:
Recover higher-rank correlations through stochastic flow evolution.

Next — ?

Brutal high-order, hybrid methods, multi-reference versions ...

Validations

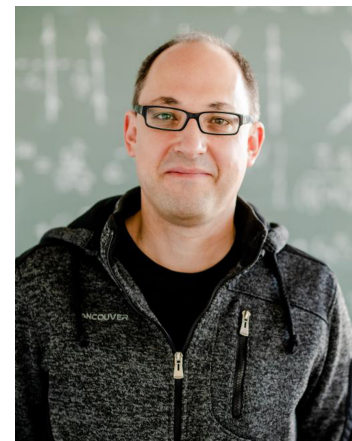
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