



北京航空航天大学
BEIHANG UNIVERSITY

Reviving primordial black hole formation in slow first-order phase transitions

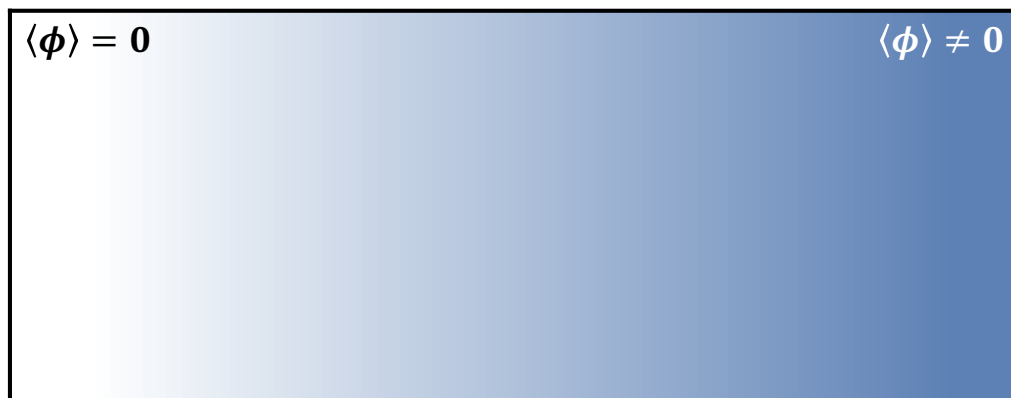
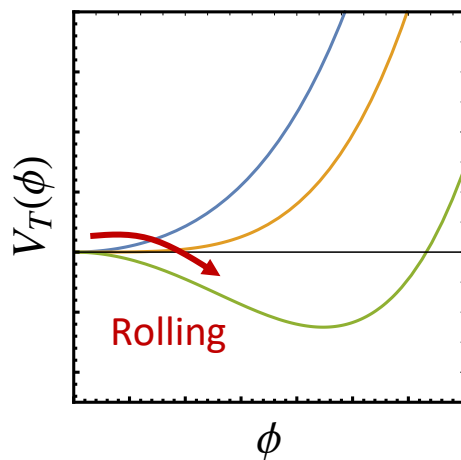
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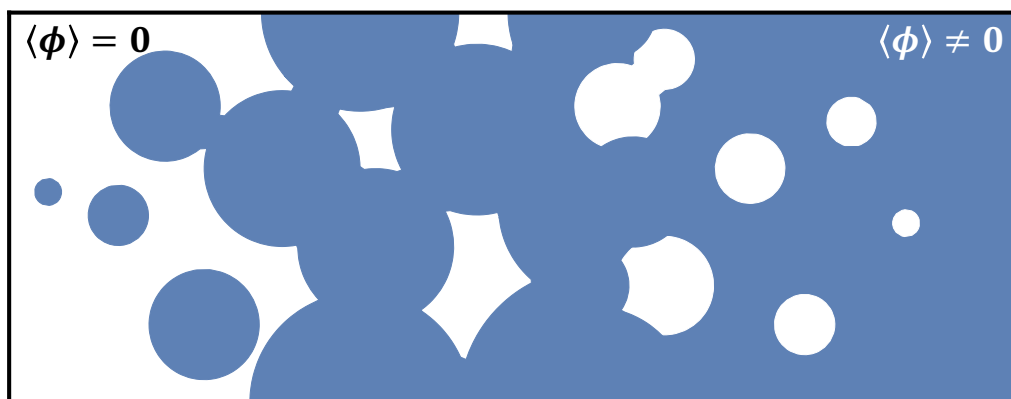
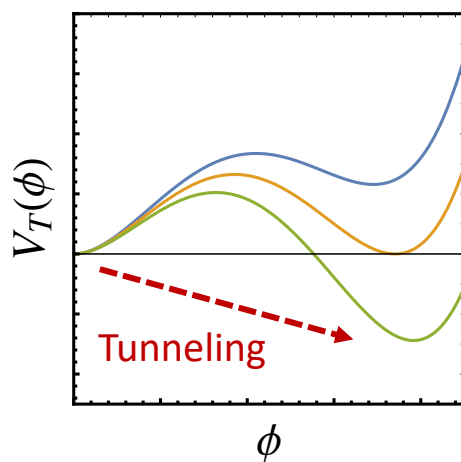
2026.7.20, The Phenomenology of Primordial Black Holes, Mainz, Germany

With Wen-Yuan Ai, [arXiv: 2605.11332](https://arxiv.org/abs/2605.11332)

Patterns of symmetry breaking



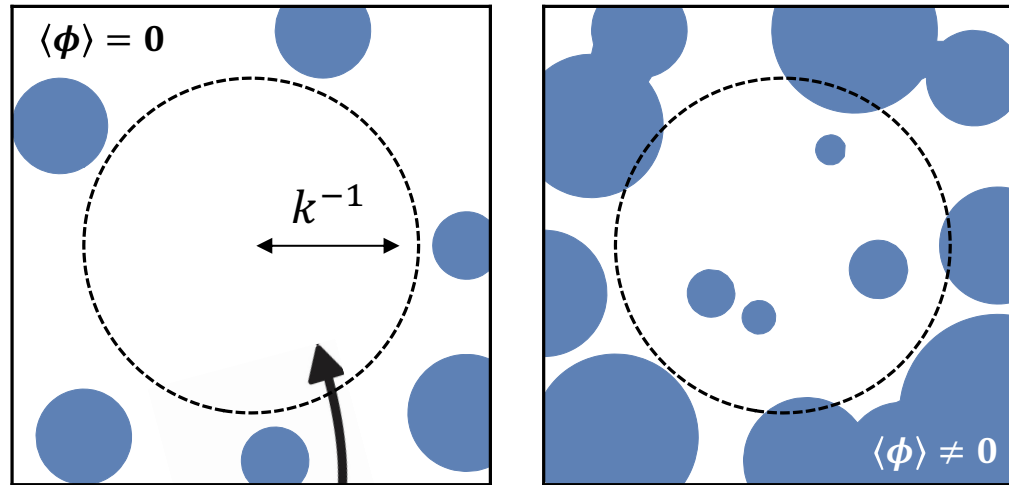
Second-order phase transition / Crossover



First-order phase transition (FOPT)

Generation of overdensities

Randomness of bubble nucleation

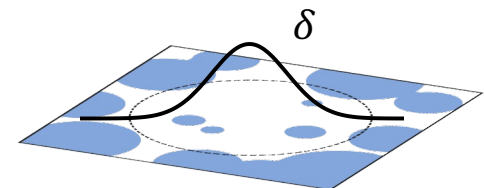


Delayed-decay region
More vacuum energy, $\rho_V \propto a^0$

Normal region
More radiation energy, $\rho_r \propto a^{-4}$

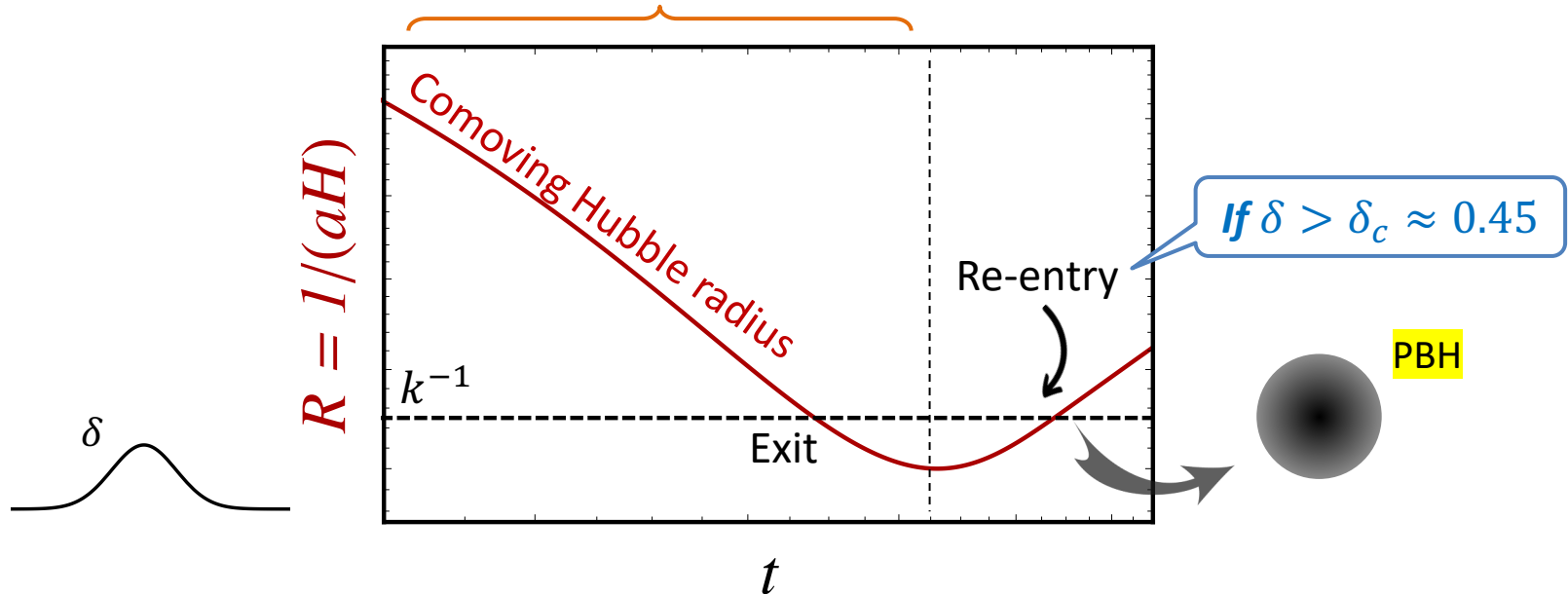
Energy density contrast

$$\delta = \frac{\rho - \bar{\rho}}{\bar{\rho}} > 0$$



Black hole formation

Supercooling: temporary vacuum domination before FOPT



Tracing back to [Kodama et al, Prog. Theor. Phys. 68 \(1982\) 1979](#)

Recent **big progress** in [Liu et al, Phys.Rev.D 105 \(2022\) 2, L021303](#)

A popular mechanism since 2022

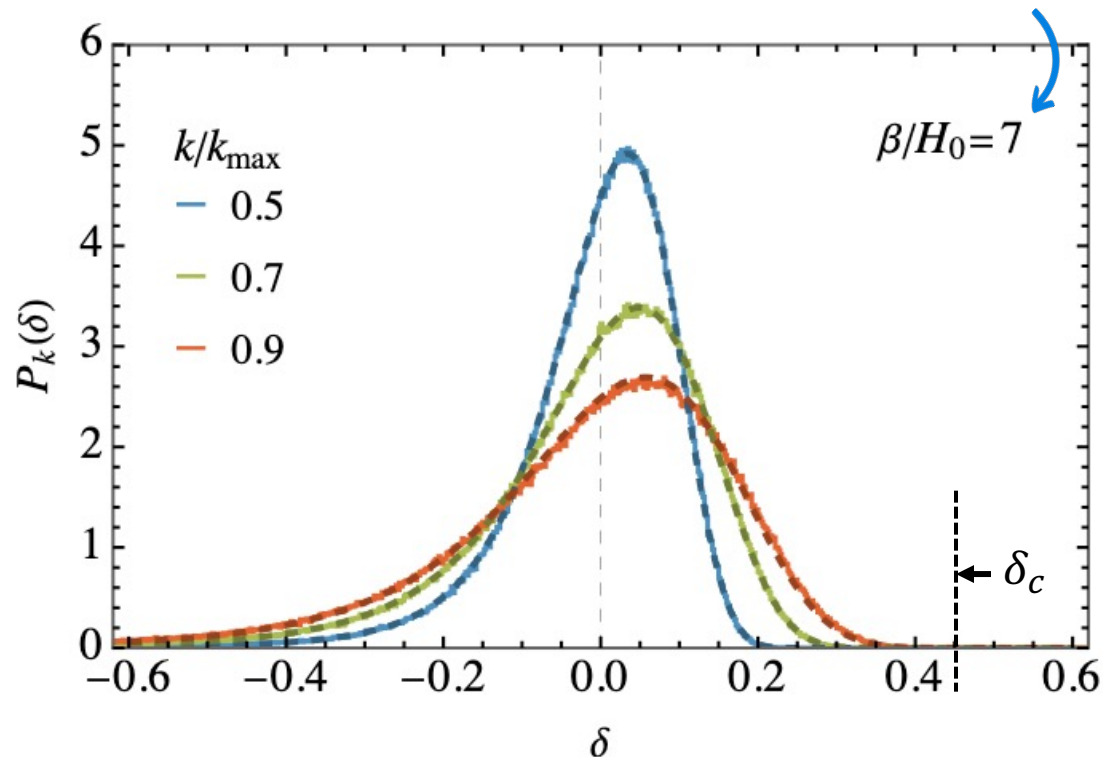
[Hashino et al, PLB 838 \(2023\) 137688](#); [Lewicki et al, JHEP 09 \(2023\) 092](#); [He et al, SCPMA 67 \(2024\) 4, 240411](#); [Cai et al, SCPMA 67 \(2024\) 9, 290411](#); [Flores et al, PRD 110 \(2024\) 1, 015005](#); [Kanemura, Tanaka, KPX, JHEP 06 \(2024\) 036](#); [Kawana et al, PRD 108 \(2023\) 103531](#); [Gouttenoire et al, PRD 110 \(2024\) 4, 4](#); [Banerjee et al, JCAP 07 \(2025\) 007](#); [Huang et al, PRD 113 \(2026\) 5, 055013](#); etc

Calculation framework

Bubble nucleation history simulation [Lewicki et al, Phys.Rev.Lett. 133 \(2024\) 22, 221003](#)

- Probability distribution of δ (non-Gaussian)

(Hubble time)/(Transition duration)

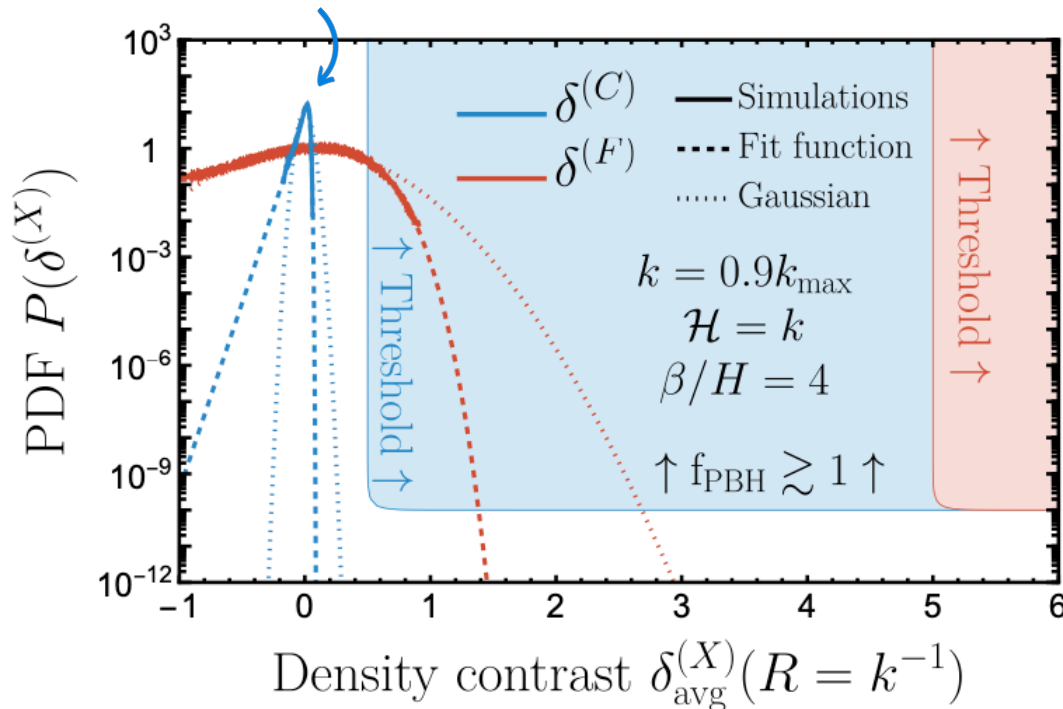


- PBH abundance obtained by integrating $P_k(\delta)$ over $\delta > \delta_c$
- Slow transitions** with $\beta/H_n \lesssim 10$ lead to abundant PBH formation

Debates

δ depends on coordinate choice [Franciolini et al, Phys.Rev.Lett. 136 \(2026\) 17, 171404](#)

- Bubble simulation $\rightarrow$ *flat-gauge* value $\delta^{(F)}$
- PBH threshold $\rightarrow$ *comoving-gauge* value $\delta^{(C)} \gtrsim 0.45$
- $\delta^{(C)}$ **too small** to form PBHs!



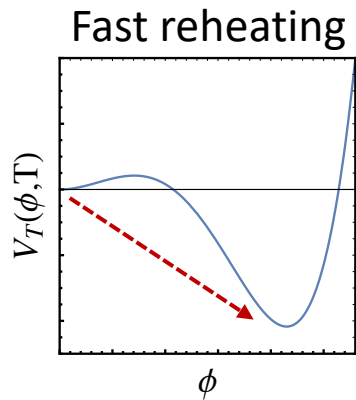
Similar result obtained independently in [Wang et al, Phys.Rev.D 113 \(2026\) 12, 123026](#)

Conclusion: this mechanism is ruled out ???

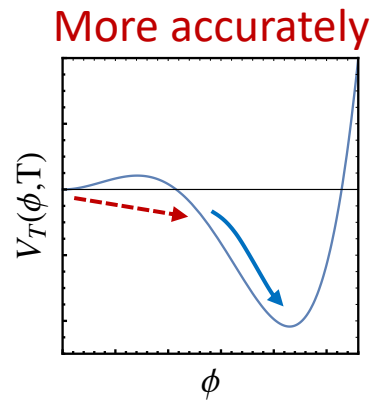
Our insight

A *new* gauge-invariant approach to calculate $\delta^{(C)}$ [this work]

- Confirming $\delta^{(C)} \ll \delta_c \approx 0.45$
- But this does **not** preclude the formation of PBHs



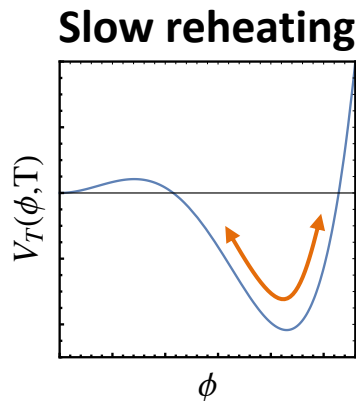
$$\Gamma_\phi / H_* > 1$$



Key variable:

$$\frac{\Gamma_\phi}{H_*} \quad \begin{array}{l} \Gamma_\phi \text{ Decay width} \\ H_* \text{ Hubble at transition} \end{array}$$

$$\Gamma_\phi / H_* < 1$$



- ϕ coherently **oscillates**, decaying to SM particles, reheating the Universe
- **Early matter era** with $\rho_\phi \propto a^{-3}$, pressure ≈ 0
- Small $\delta^{(C)}$ grows and collapses, **forming PBHs** ✓

A unified treatment

	Condition	Energy transfer
Fast reheating	$\Gamma_\phi/H_* > 1$	Vacuum $\rightarrow$ Radiation
Slow reheating	$\Gamma_\phi/H_* < 1$	Vacuum $\rightarrow \phi$ oscillation $\rightarrow$ Radiation

Generation of δ

Two-fluid system

1. Vacuum, $P_V = -\rho_V$
 2. Species f with $P_f = w_f \rho_f$
- $\left\{ \begin{array}{l} \text{Fast reheating} \rightarrow f = r \ (w_r = 1/3) \\ \text{Slow reheating} \rightarrow f = \phi \ (w_\phi = 0) \end{array} \right.$

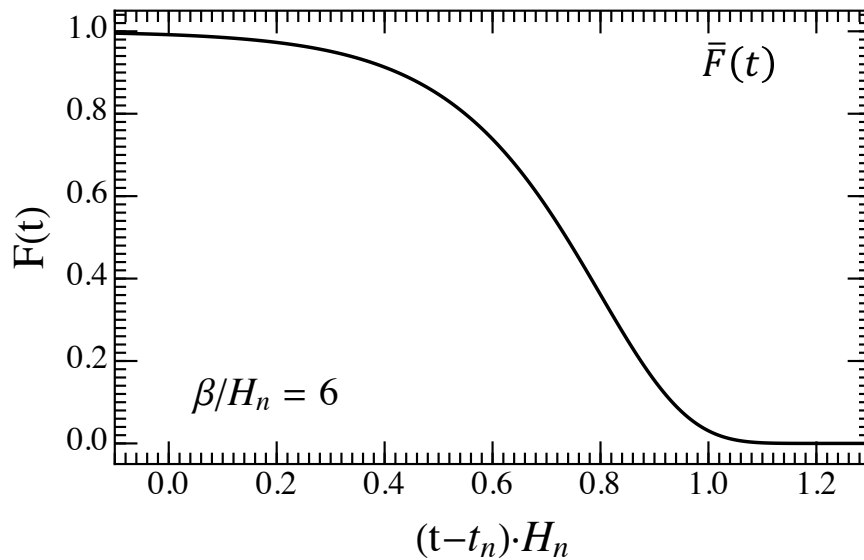
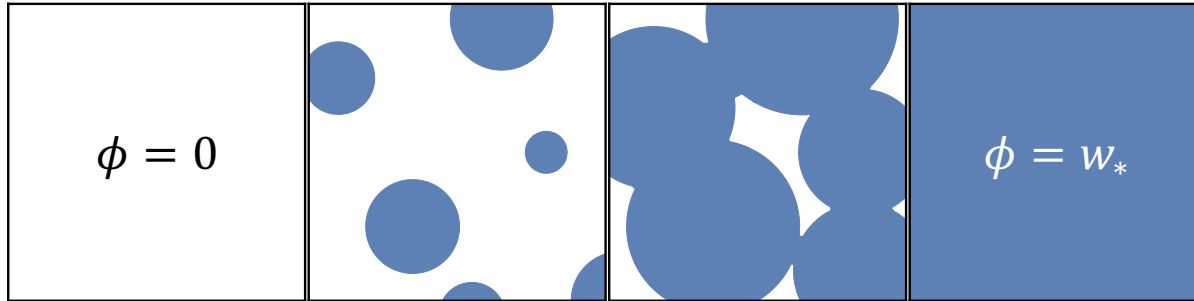
Relativistic perturbation theory

$$\rho(t, \mathbf{k}) = \bar{\rho}(t) + \delta\rho(t, \mathbf{k}); \quad \delta = \frac{\delta\rho}{\bar{\rho}} = \frac{\delta\rho_V + \delta\rho_f}{\bar{\rho}_V + \bar{\rho}_f}$$

- FOPT yields $\delta\rho_V \rightarrow$ sourcing $\delta \rightarrow$ PBH formation

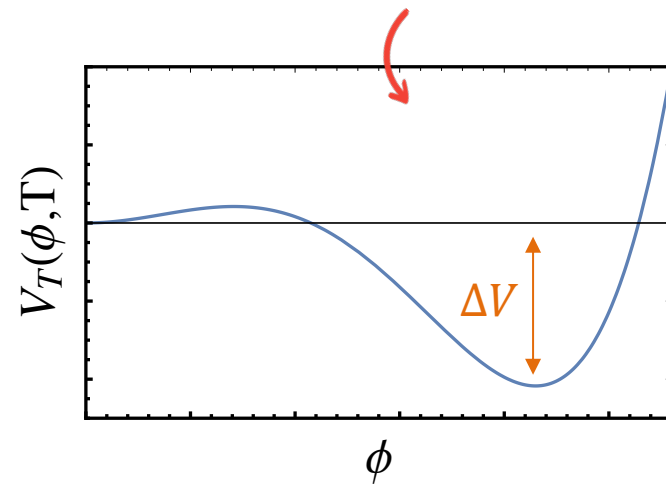
Vacuum energy

False vacuum volume fraction: the whole Universe



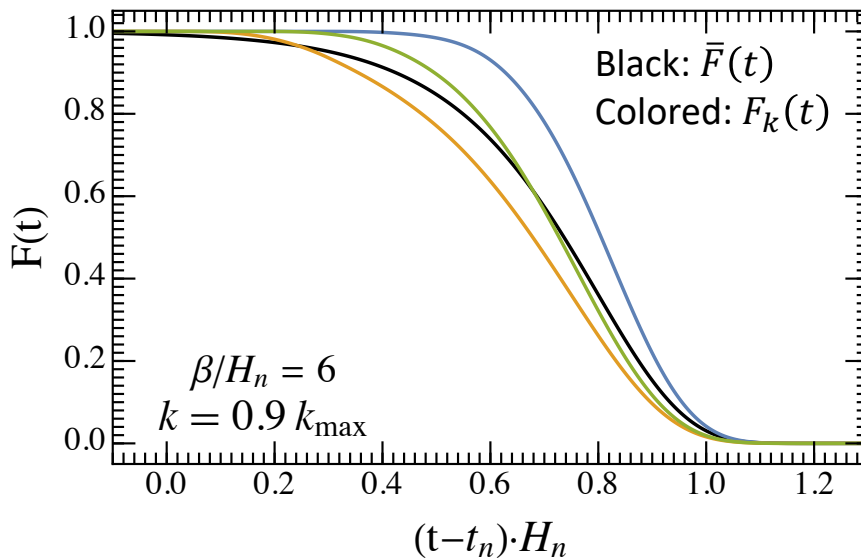
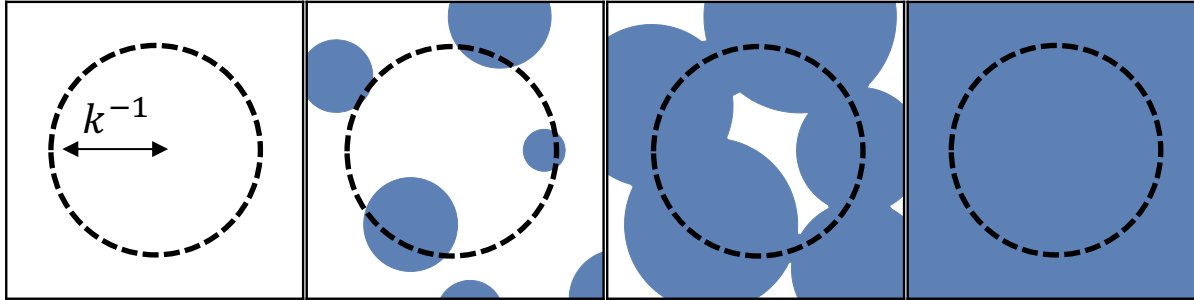
Background vacuum energy

$$\bar{\rho}_V(t) = \Delta V \cdot \bar{F}(t)$$



Vacuum energy perturbation

False vacuum volume fraction: a finite sphere



Comoving volume $4\pi k^{-3}/3$
 $\rho_V(t) = \Delta V \cdot F_k(t)$

Obtained from **bubble nucleation simulation**

Fluctuation $\delta\rho_V = \rho_V - \bar{\rho}_V = \Delta V \cdot (F_k(t) - \bar{F}(t))$



Perturbation theory (1)

Textbook: **gauge-invariant** equations

$$\dot{\Delta} - 3H \frac{\bar{P}}{\bar{\rho}} \dot{\Delta} = \left(1 + \frac{\bar{P}}{\bar{\rho}}\right) \frac{k^2 \mathcal{V}}{a}; \quad \dot{\mathcal{V}} + H\mathcal{V} = -\frac{1}{a(\bar{\rho} + \bar{P})} \left(\delta P_{\text{nad}} + \frac{\dot{\bar{P}}}{\dot{\bar{\rho}}} \bar{\rho} \Delta \right) - \frac{\Phi}{a}$$

Density contrast

$$\Delta = \delta + \frac{a\dot{\rho}(v + B)}{\bar{\rho}}$$

Velocity potential

$$\mathcal{V} = v + a\dot{E}$$

with $\mathbf{v} = \nabla v$

Non-adiabatic pressure
perturbation

$$\delta P_{\text{nad}} = \delta P - \frac{\dot{\bar{P}}}{\dot{\bar{\rho}}} \delta \rho$$

Perturbation theory (2)

Textbook: **gauge-invariant** equations

$$\dot{\Delta} - 3H \frac{\bar{P}}{\bar{\rho}} \dot{\Delta} = \left(1 + \frac{\bar{P}}{\bar{\rho}}\right) \frac{k^2 \mathcal{V}}{a}; \quad \dot{\mathcal{V}} + H\mathcal{V} = - \frac{1}{a(\bar{\rho} + \bar{P})} \left(\delta P_{\text{nad}} + \frac{\dot{\bar{P}}}{\dot{\bar{\rho}}} \bar{\rho} \Delta \right) - \frac{\Phi}{a}$$

Our work: two-fluid system

- Evaluating δP_{nad} in **flat gauge**

$$\dot{Q} + (4 + 3w_f)HQ = \frac{1 + w_f}{a} \delta \rho_V - \left[\frac{w_f}{a} \bar{\rho} + \frac{aH}{2k^2} (\dot{\bar{\rho}} - 3(1 + w_f)\dot{\bar{\rho}}_V) \right] \Delta$$


 Momentum perturbation
 $Q = (\bar{\rho} + \bar{P})\mathcal{V}$


 From bubble simulations

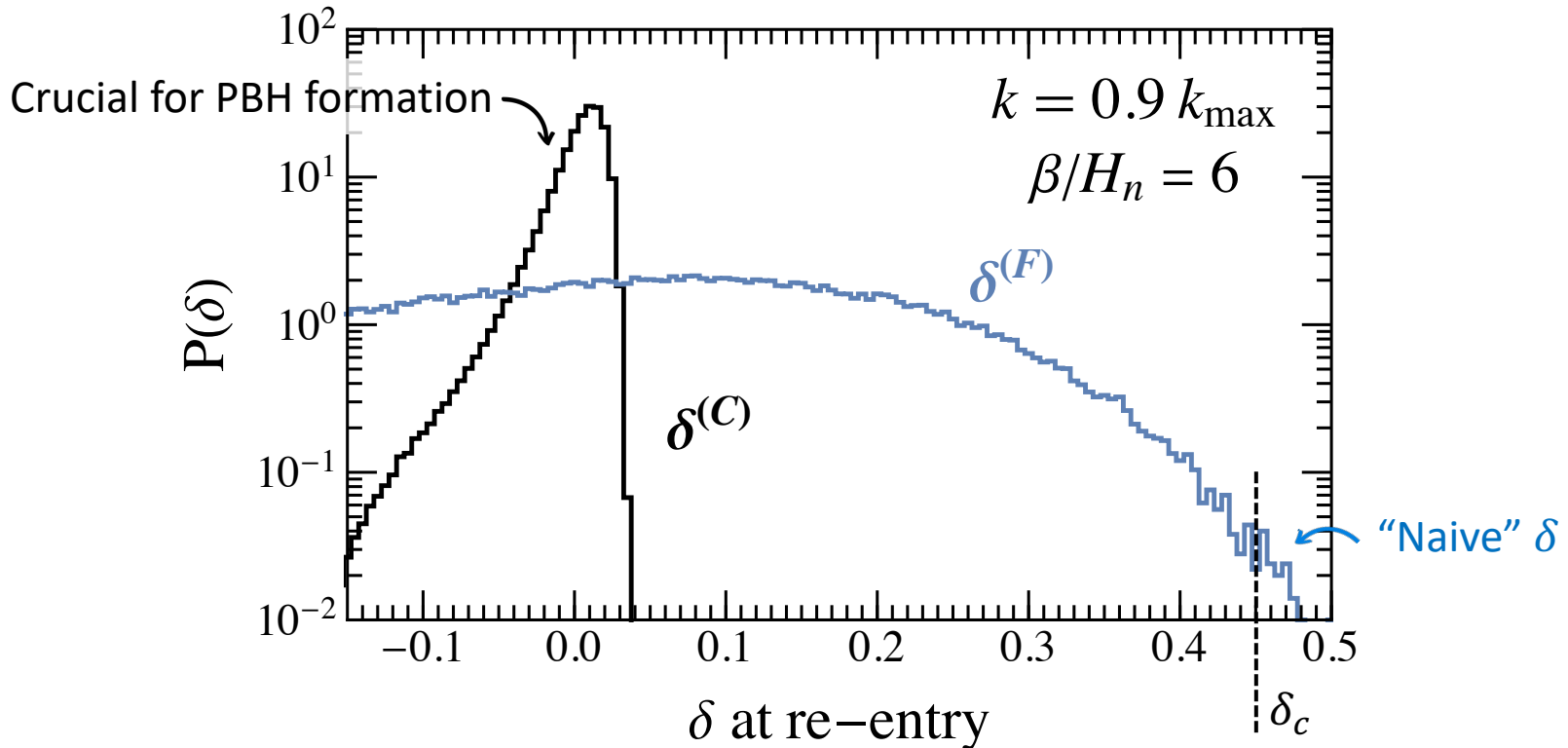
$\Delta \equiv \delta^{(C)}$ [**comoving gauge**] resolved, evolved to horizon re-entry

- Single realization $\rightarrow$ one $\delta^{(C)}$
- Multiple realizations $\rightarrow$ probability distribution of $\delta^{(C)}$

Distribution of δ (1)

δ **very sensitive** to gauge choice

$$\delta^{(C)} \ll \delta^{(F)}$$



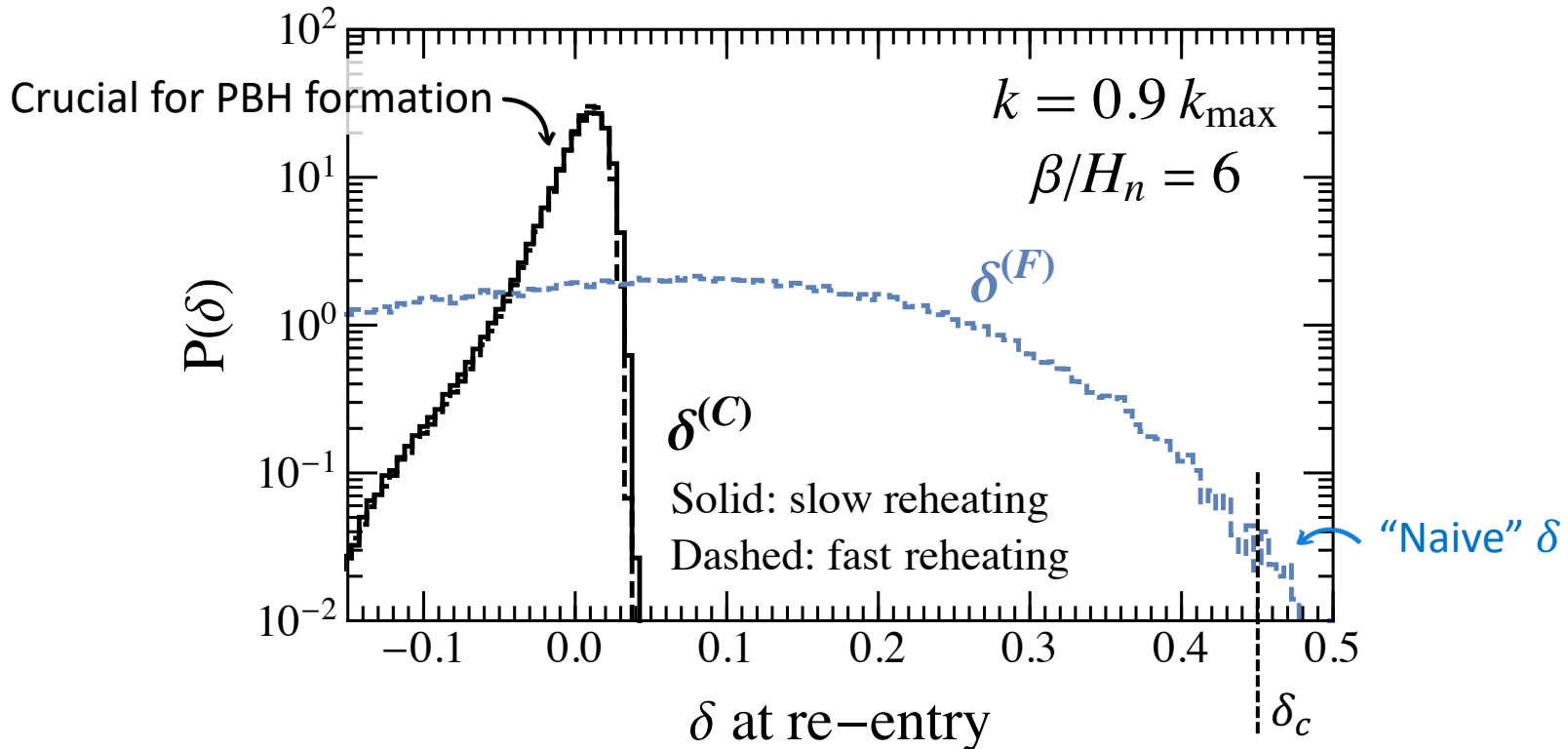
Fast reheating:

- Consistent with [Phys.Rev.Lett. 136 \(2026\) 17, 171404](#); [Phys.Rev.D 113 \(2026\) 12, 123026](#)

Distribution of δ (2)

δ **not sensitive** to the efficiency of post-FOPT reheating

$$\delta^{(C)}(f = r) \approx \delta^{(C)}(f = \phi)$$



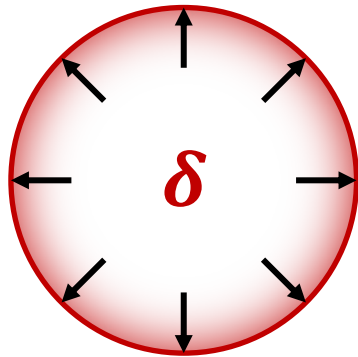
Slow reheating:

- **New** in our paper, crucial for PBH formation

Hereafter, δ solely denotes $\delta^{(C)}$

Fate of the density contrast

Fast reheating: pressure $P_r = \rho_r/3$ resists collapse



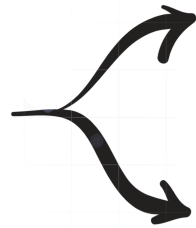
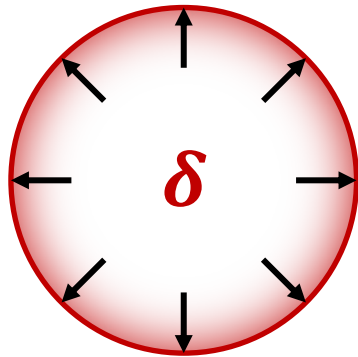
$\delta > \delta_c$: instantaneous collapse
Black hole formation



$\delta < \delta_c$: no collapse
Overdensity dissipates

Fate of the density contrast

Fast reheating: pressure $P_r = \rho_r/3$ resists collapse

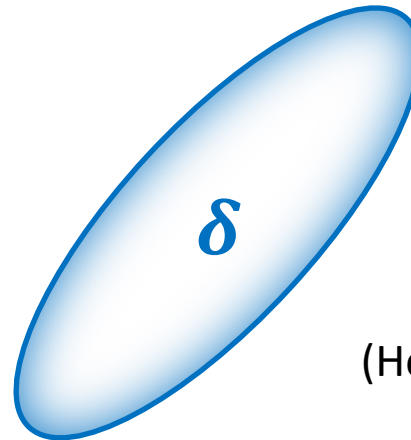
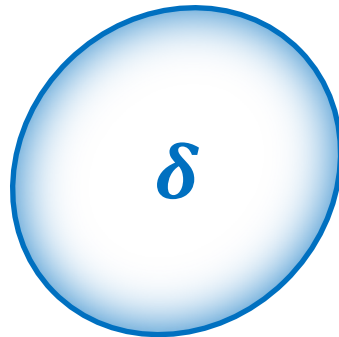


$\delta > \delta_c$: instantaneous collapse
Black hole formation



$\delta < \delta_c$: no collapse
Overdensity dissipates

Slow reheating: $P_\phi \approx 0$, perturbation δ increases with time



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Non-sphericity &
angular momentum grows

(Hoop conjecture)

“Pancake collapse” at late time, **black hole** formation

PBH formation in the matter era

Perturbation shape describe by the **deformation tensor** $\mathcal{D}_{3 \times 3}$

Our crucial assumption

$$P(\mathcal{D}) = P(\delta) \times P_{\tilde{\mathcal{D}}}(\tilde{\mathcal{D}}_{ij})$$

From our simulation (Non-Gaussian) $\leftarrow \delta = -\text{tr } \mathcal{D}$ Trace part $\leftarrow \tilde{\mathcal{D}}_{ij} = \mathcal{D}_{ij} + \frac{\delta_{ij}}{3} \delta \rightarrow$ Assumed to be Gaussian Traceless part

(Calling for future detailed simulations!)

We improve Refs. [Harada et al, Astrophys. J. 833 \(2016\) no. 1, 61, PRD 96 \(2017\) no. 8, 083517](#)

PBH formation criterion: $\Gamma_{\phi}/H_* \lesssim 10^{-3}$ (long matter era)

Probability

$$\beta_k \approx \int_{\delta_{\text{th}}}^{\infty} d\delta \int_{-\infty}^{\infty} dy \int_{2|y|}^{\infty} dz \left[P(\delta) \frac{9\sqrt{2}}{\sqrt{\pi}\sigma_3^5} (z^3 - 4zy^2) e^{-(2y^2+3z^2/2)/\sigma_3^2} \right]$$

From hoop conjecture, typically $\delta_{\text{th}} \ll 0.45$

PBH profile

Key variables $\left\{ \begin{array}{l} \text{New physics scale } \Lambda_{\text{NP}} \approx \Delta V^{1/4} \\ \beta/H_n = (\text{Hubble time})/(\text{Transition duration}) \\ \Gamma_\phi/H_* = (\text{Hubble time})/(\text{Early matter era duration}) \end{array} \right.$

PBH mass & dark matter fraction

No accretion [Ballesteros et al, JCAP 06 \(2020\) 014](#)

$$M_{\min} \sim 10^{20} \text{ g} \times \left(\frac{10^6 \text{ GeV}}{\Lambda_{\text{NP}}} \right)^2$$

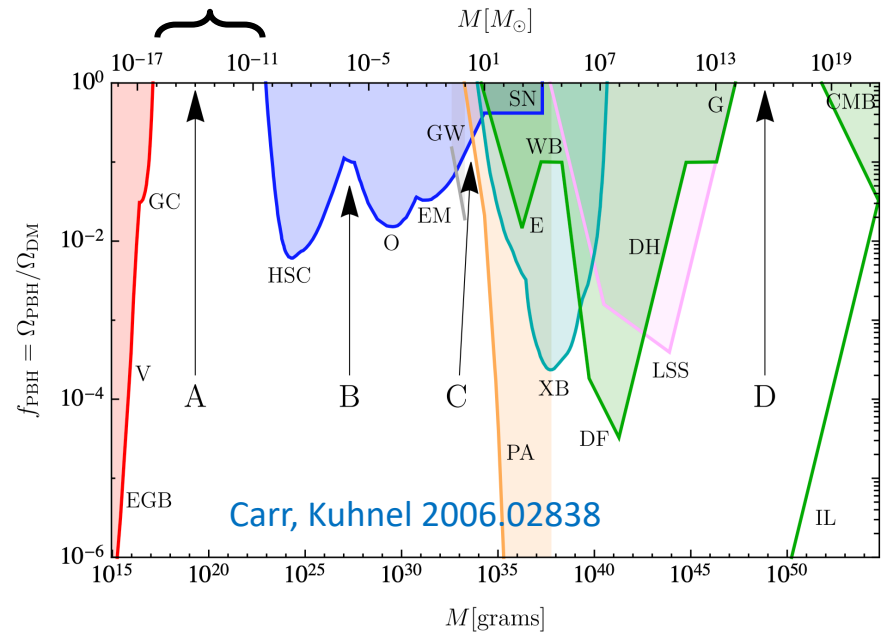
$$f_{\min} \sim 10^{0.4 \cdot (12 - \beta/H_n)} \left(\frac{\Lambda_{\text{NP}}}{10^6 \text{ GeV}} \right) \left(\frac{\Gamma_\phi}{H_*} \right)^{1/2}$$

Maximal accretion [de Jong et al, JCAP 03 \(2022\) 03, 029](#)

$$M_{\max} \sim 10^{20} \text{ g} \times \left(\frac{10^6 \text{ GeV}}{\Lambda_{\text{NP}}} \right)^2 \left(\frac{H_*}{\Gamma_\phi} \right)$$

$$f_{\max} \sim 10^{0.4 \cdot (12 - \beta/H_n)} \left(\frac{\Lambda_{\text{NP}}}{10^6 \text{ GeV}} \right) \left(\frac{H_*}{\Gamma_\phi} \right)^{1/2}$$

Asteroid-mass ($10^{20} \text{ g} - 10^{22} \text{ g}$),
allowing for 100% **dark matter**



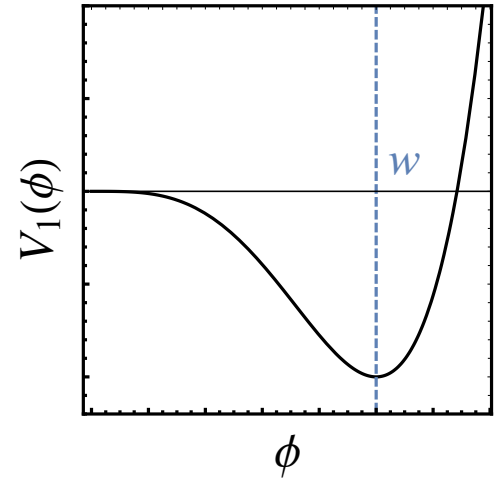
Model-building guide

Supercooled and slow FOPT

- Classically conformal principle

$$V_1(\phi) = \frac{B}{4} \phi^4 \left(\log \frac{\phi}{w} - \frac{1}{4} \right)$$

- Motivated by hierarchy problem [Bardeen, FERMILAB-CONF-95-391-T](#); [Meissner et al, PLB 660 \(2008\) 260-266](#)
- Leads to supercooling [Iso et al, PRL 119, 141301 \(2017\)](#)



Slow (post-FOPT) reheating

- Suppressed $\phi \rightarrow \text{SM SM}$ decay

Dark $U(1)_X$ model

[Babu et al, Phys. Rev. D 57 \(1998\) 6788–6792](#)

- Scalar Φ , boson X
- Kinetic mixing

$$\mathcal{L} \sim -\frac{\epsilon}{2} X_{\mu\nu} B^{\mu\nu}$$

- $\phi \rightarrow X^* X^* \rightarrow f \bar{f} f' \bar{f}', \Gamma_\phi \propto g_X^{10} \epsilon^4$

Minimal $U(1)_{B-L}$ model

[Iso et al, Phys.Lett.B 676 \(2009\) 81-87](#)

- Scalar Φ , gauge boson Z'
- Higgs mixing

$$\theta \approx \frac{v_{ew}}{w} \frac{m_h^2}{m_h^2 - m_\phi^2} \sim g_{B-L}^2 \left(\frac{v_{ew}}{m_\phi} \right)^3$$

- $\Gamma_\phi \propto \theta^2$

Conclusion

This work Ai and KPX, arXiv: 2605.11332

- Gauge-invariant derivation of $\delta^{(C)}$ from FOPT without super-horizon approximations
- **PBH formation revived** via slow reheating

Outlook

- Full simulations of $\mathcal{D}_{3 \times 3}$ & PBH spin distribution
- Gravitational waves from FOPT (with matter-domination imprint) & non-spherical collapse
- Particle phenomenology, e.g., collider, neutrino oscillation, etc

*“Problems worthy of attack
prove their worth by fighting back.”*
— Piet Hein

Thank you!