

A Brief Introduction to Nuclear Physics

Pierre Capel



JOHANNES GUTENBERG
UNIVERSITÄT MAINZ



18 January 2026

Introduction

Stable nuclei are qualitatively described by “simple” models

- (semi-empirical) **liquid-drop** model
- (basic) **shell** model

New techniques enable ***ab initio*** methods (A -body models)

What happens far from stability ?

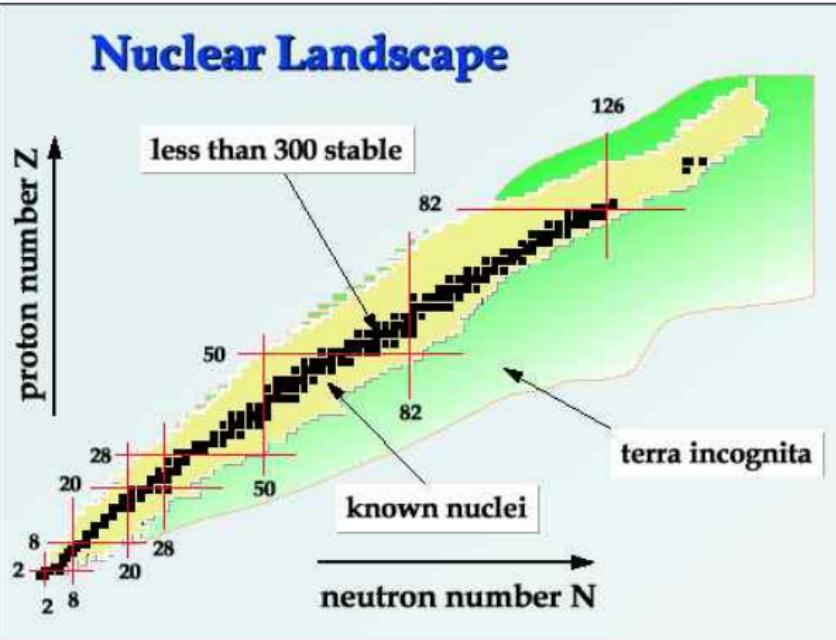
Experimentally, Radioactive-Ion Beams (RIB) available since 80s

⇒ study of structure **far from stability**

⇒ discovery of **exotic** structures

- **halo** nuclei
- shell inversions
- ...

Nuclear Landscape



- 256 stable nuclei
@ $Z \approx N$ up to ^{40}Ca
@ $N > Z$ for $A > 40$
- stable nuclei
 - ▶ compact
 - ▶ *magic numbers*
- RIB allow to study radioactive nuclei
- Terra incognita between driplines
 n -dripline unknown beyond ^{34}Ne

1 Basic features in nuclear structure

- Liquid-drop model
- Shell model

2 Ab initio nuclear models

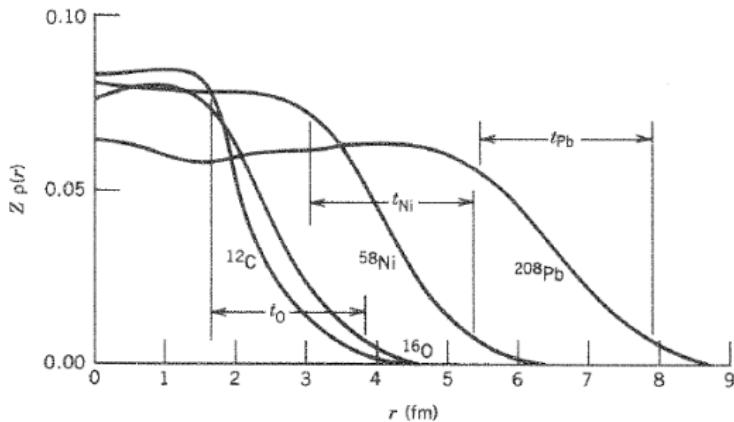
3 Radioactive-Ion Beams

4 Oddities far from stability

- Halo nuclei

5 Summary

Charge distributions in (stable) nuclei

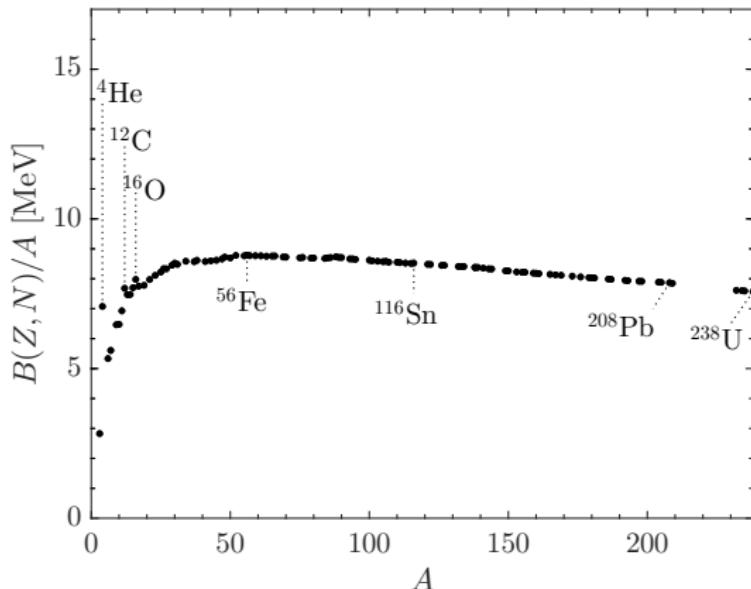


- constant density ρ_0 out to the surface (saturation)
- same skin thickness t

(Stable) nuclei look like **liquid drops** of radius $R \propto A^{1/3}$

Liquid-drop model

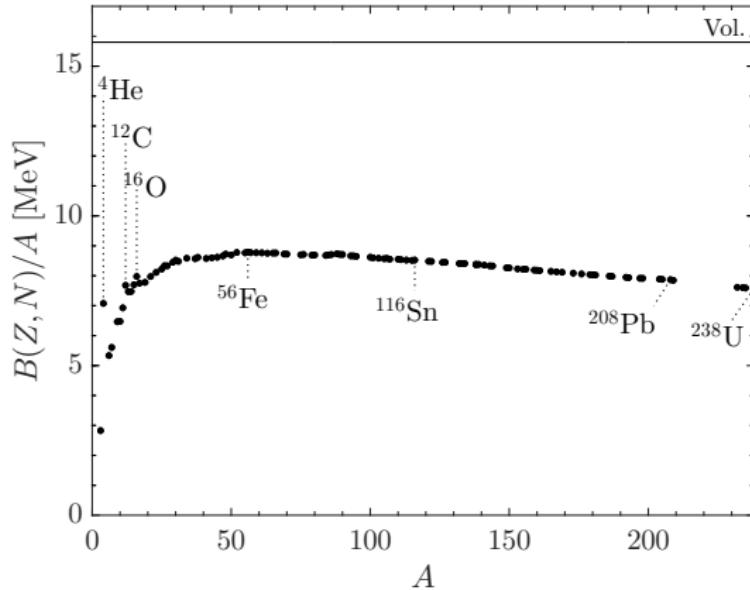
Binding energy per nucleon $B(Z, N)/A$ has smooth behaviour



Liquid-drop model

Binding energy per nucleon $B(Z, N)/A$ has smooth behaviour
Bethe-Weizsäcker semi-empirical mass formula

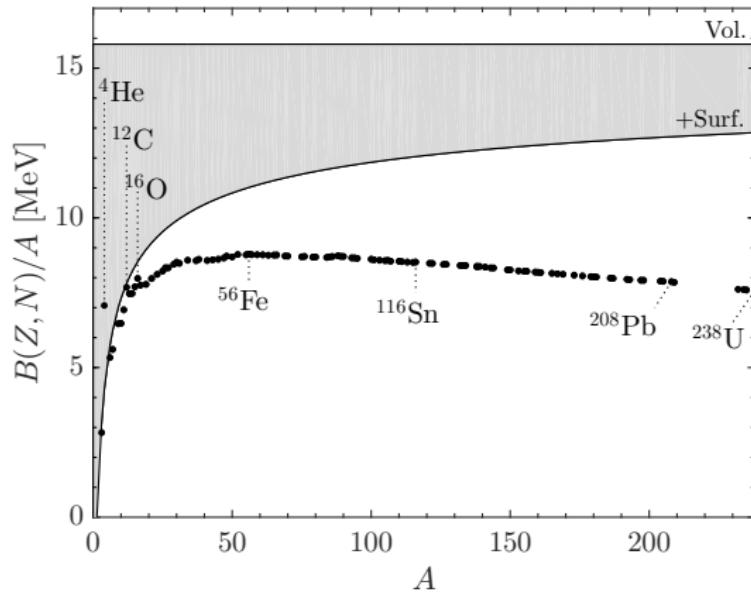
$$B(Z, N) = a_v A$$



Liquid-drop model

Binding energy per nucleon $B(Z, N)/A$ has smooth behaviour
Bethe-Weizsäcker semi-empirical mass formula

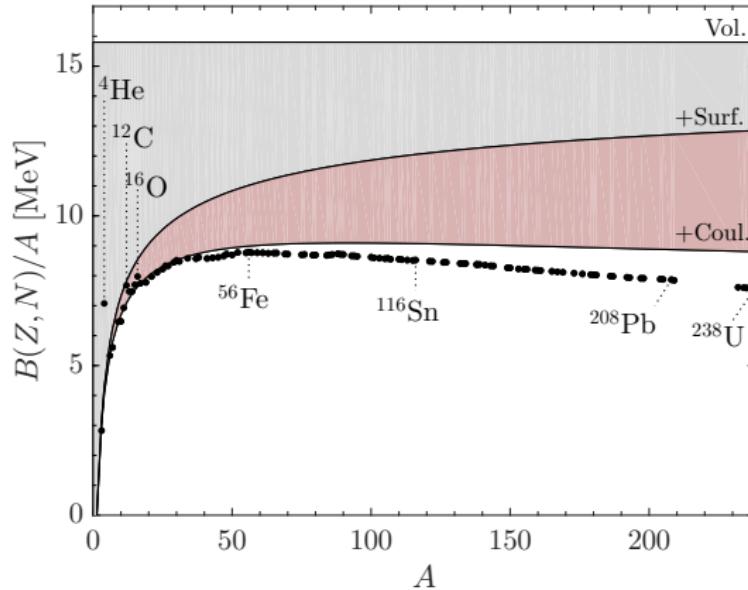
$$B(Z, N) = a_v A - a_s A^{2/3}$$



Liquid-drop model

Binding energy per nucleon $B(Z, N)/A$ has smooth behaviour
Bethe-Weizsäcker semi-empirical mass formula

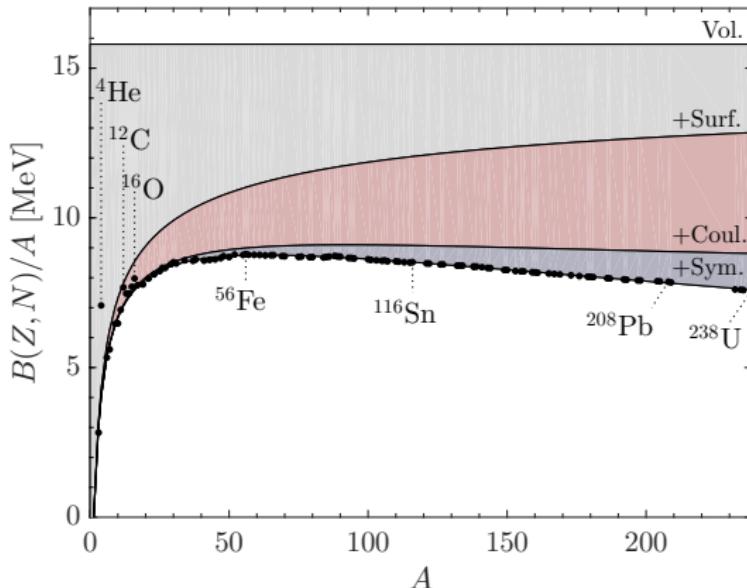
$$B(Z, N) = a_v A - a_s A^{2/3} - a_C \frac{Z(Z-1)}{A^{1/3}}$$



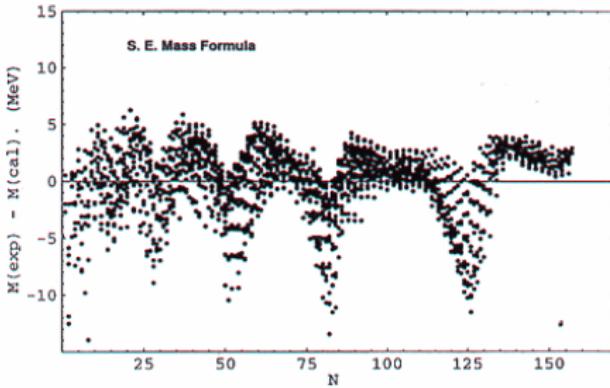
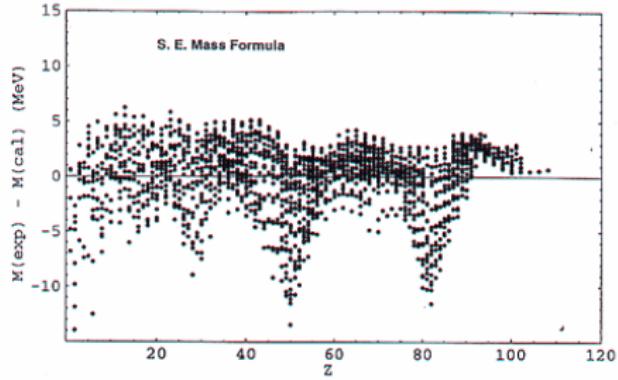
Liquid-drop model

Binding energy per nucleon $B(Z, N)/A$ has smooth behaviour
 Bethe-Weizsäcker semi-empirical mass formula

$$B(Z, N) = a_v A - a_s A^{2/3} - a_C \frac{Z(Z-1)}{A^{1/3}} - a_{\text{Sym}} \frac{(A-2Z)^2}{A}$$



Variation from the semi-empirical mass formula



More bound systems at Z or $N = 2, 8, 20, 28, 50, 82, 126$
magic numbers

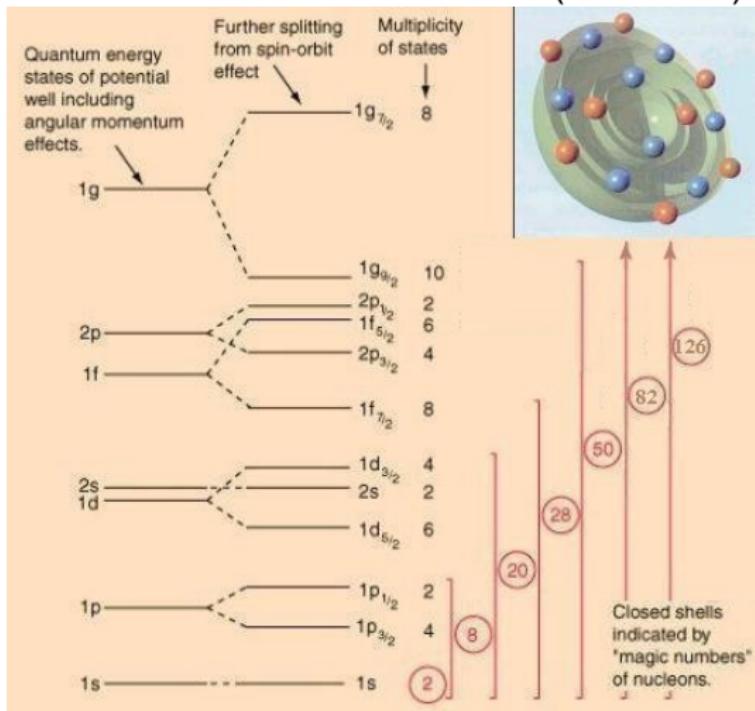
⇒ shell structure in nuclei as in atoms ?

Shell model

Developed in 1949 by M. Goeppert Mayer and H. Jensen

As electrons in atoms,
nucleons in nuclei
feel a **mean field**
and arrange into **shells**

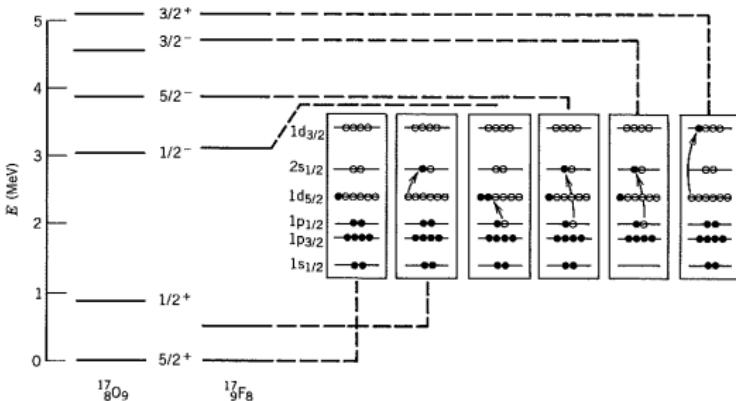
Spin-orbit coupling is crucial
to get right ordering of shells



Example

Shell model explains the higher stability at some Z and N

It predicts the spin and parity of ground state of most nuclei and some of their excited levels, e.g. ^{17}O and ^{17}F



Nowadays

Can we go beyond these models ?

Can we build *ab initio* models ?

i.e. based on first principles

- nucleons as building blocks
- realistic N - N interaction

1 Basic features in nuclear structure

- Liquid-drop model
- Shell model

2 Ab initio nuclear models

3 Radioactive-Ion Beams

4 Oddities far from stability

- Halo nuclei

5 Summary

A-body Hamiltonian

Nuclear-structure calculations : *A* nucleons (*Z* protons+*N* neutrons)

Relative motion described by the *A*-body Hamiltonian

$$H = \sum_{i=1}^A T_i + \sum_{j>i=1}^A V_{ij}$$

⇒ solve the *A*-body Schrödinger equation

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle$$

$\{E_n\}$ is the nucleus spectrum

Realistic N - N interactions

V_{ij} not (yet) deduced from QCD

⇒ **phenomenological** potentials

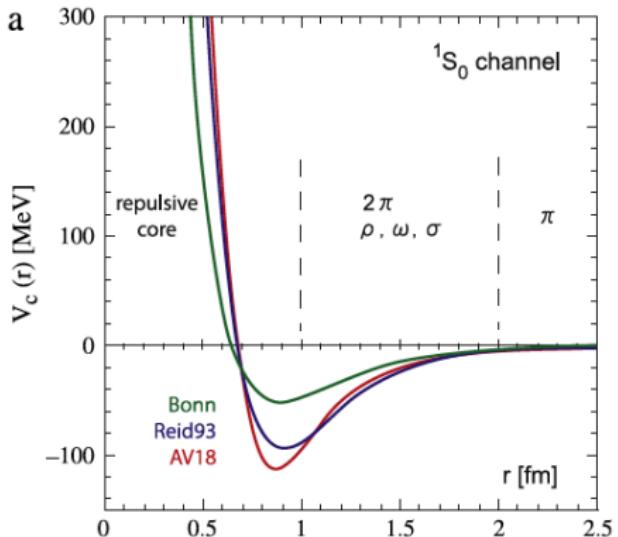
fitted on N - N observables :

binding energy,

N - N phaseshifts

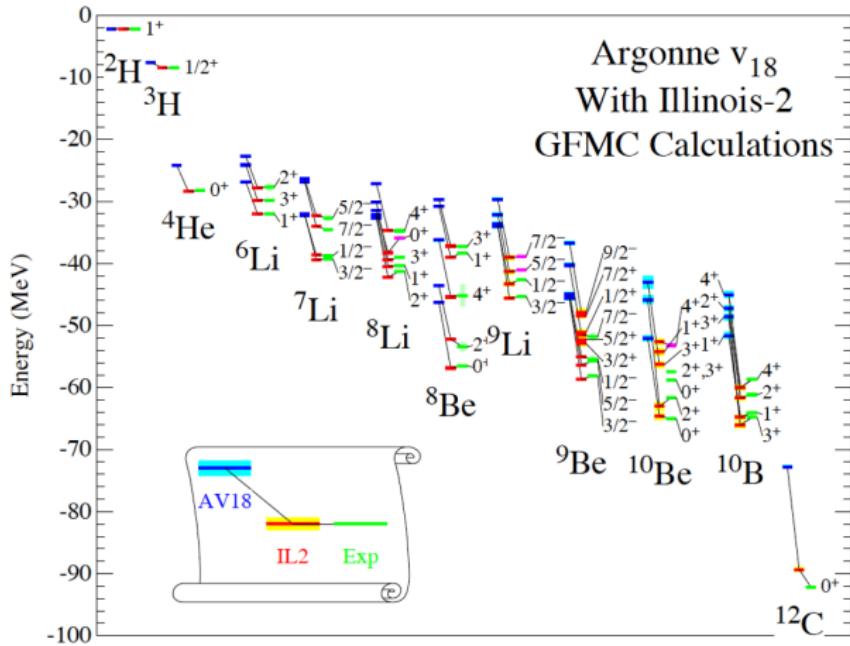
Realistic N - N interactions

V_{ij} not (yet) deduced from QCD
 \Rightarrow **phenomenological** potentials
 fitted on N - N observables :
 d binding energy,
 N - N phaseshifts
 Ex. : Argonne V18, CD-Bonn,...



[see A. Schmidt's talk on Friday]

Light nuclei calculations



[R. Wiringa, Argonne]

Three-body force

Need three-body forces to get it right...

$$H = \sum_{i=1}^A T_i + \sum_{j>i=1}^A V_{ij} + \sum_{k>j>i=1}^A V_{ijk} + \dots$$

But there is no such thing as three-body force...

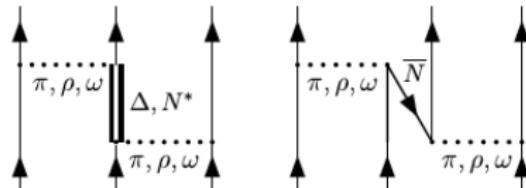
Three-body force

Need three-body forces to get it right...

$$H = \sum_{i=1}^A T_i + \sum_{j>i=1}^A V_{ij} + \sum_{k>j>i=1}^A V_{ijk} + \dots$$

But there is no such thing as three-body force...

They simulate the **non-elementary** character of nucleons
 \Rightarrow include virtual Δ resonances, \bar{N} ...



Phenomenological 3-body interaction fitted on $A > 2$ levels : IL2
 Alternatively, derived from EFT

Effective Field Theory

EFT is an **effective** quantum field theory based on QCD symmetries with resolution scale Λ that selects appropriate degrees of freedom : nuclear physics is not built on quarks and gluons, but on **nucleons** and **mesons**

EFT provides the nuclear force with a systematic expansion in Q/Λ

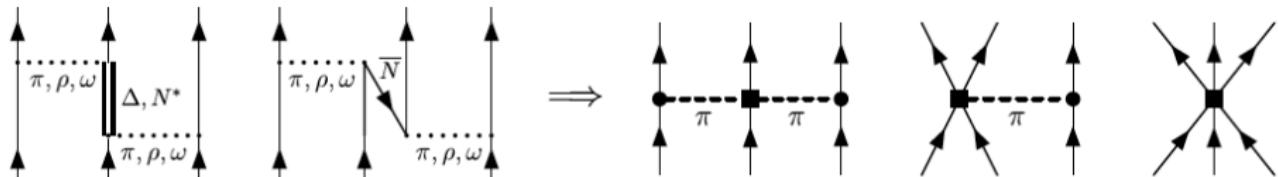
- gives an estimate of theoretical uncertainty

Effective Field Theory

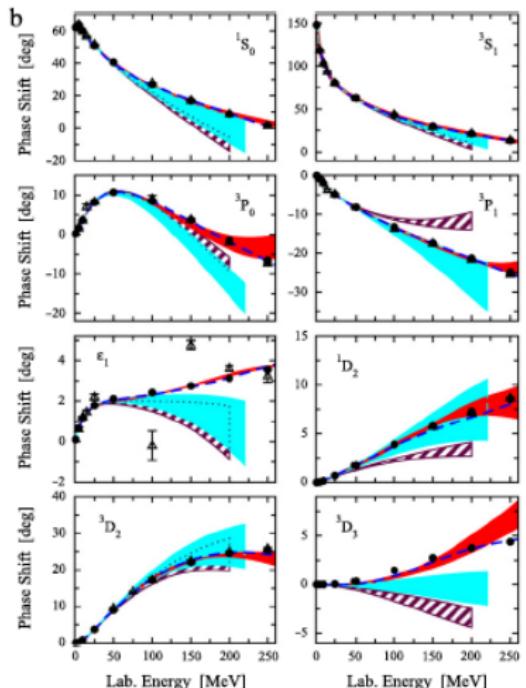
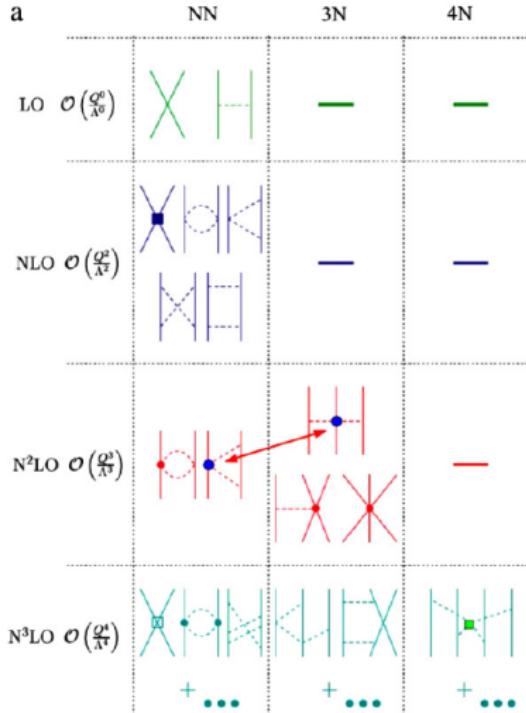
EFT is an **effective** quantum field theory based on QCD symmetries with resolution scale Λ that selects appropriate degrees of freedom : nuclear physics is not built on quarks and gluons, but on **nucleons** and **mesons**

EFT provides the nuclear force with a systematic expansion in Q/Λ

- gives an estimate of theoretical uncertainty
- naturally includes many-body forces



Expansion of the EFT force



Solving the Schrödinger equation

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle$$

Ψ usually developed on a **basis** $\{|\Phi_{[v]}\rangle\}$:

$$|\Psi_n\rangle = \sum_{[v]} \langle \Phi_{[v]} | \Psi_n \rangle |\Phi_{[v]}\rangle$$

Solving the Schrödinger equation

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle$$

Ψ usually developed on a **basis** $\{|\Phi_{[\nu]}\rangle\}$:

$$|\Psi_n\rangle = \sum_{[\nu]} \langle \Phi_{[\nu]} | \Psi_n \rangle |\Phi_{[\nu]}\rangle$$

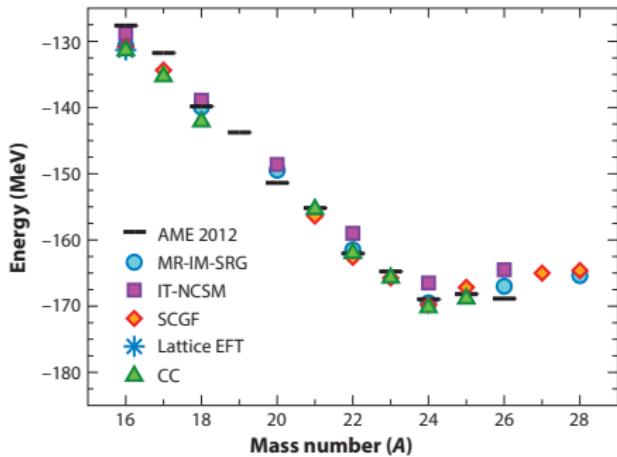
Solving the Schrödinger equation reduces to matrix diagonalisation

$$\begin{aligned} \langle \Phi_{[\mu]} | H | \Psi_n \rangle &= \sum_{[\nu]} \langle \Phi_{[\mu]} | H | \Phi_{[\nu]} \rangle \langle \Phi_{[\nu]} | \Psi_n \rangle \\ &= E_n \langle \Phi_{[\mu]} | \Psi_n \rangle \end{aligned}$$

⇒ need to build an efficient set of basis states $\{|\Phi_{[\nu]}\rangle\}$

Clear short review paper : [Bacca EPJ Plus **131**, 107 (2016)]

Example : oxygen isotopes



[Hebeler *et al.* Annu. Rev. Nucl. Part. Sci. **65**, 457 (2015)]

Different *ab initio* models predict similar result
All require $3N$ forces to reproduce the dripline at ^{24}O

What happens far from stability ?

Liquid-drop and shell models are fair models of stable nuclei

What happens away from stability ?

In 80s Radioactive-Ion Beams were developed

[see A. Spyrou on Monday & T. Nilsson on Thursday]

Enable study of nuclear **structure**

- are radioactive nuclei compact ?
- are shells conserved far from stability ?

Reactions involving radioactive nuclei useful in **astrophysics**

[see 2nd part

& A. Spyrou on Monday & E. Lopez Saavedra on Wednesday]

1 Basic features in nuclear structure

- Liquid-drop model
- Shell model

2 Ab initio nuclear models

3 Radioactive-Ion Beams

4 Oddities far from stability

- Halo nuclei

5 Summary

How ?

Idea : break a heavy nuclei into pieces to produce exotic isotopes

How ?

Idea : break a heavy nuclei into pieces to produce exotic isotopes

- ISOL : Fire a proton at a heavy nucleus



How ?

Idea : break a heavy nuclei into pieces to produce exotic isotopes

- ISOL : Fire a proton at a heavy nucleus



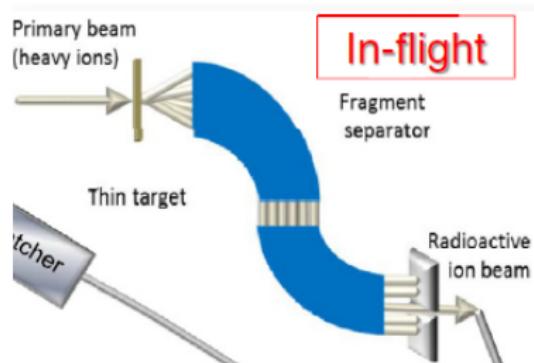
- In-flight : Smash a heavy nucleus on a target



Where ?

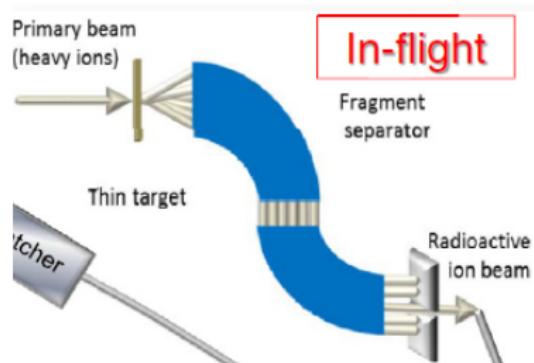


In-flight projectile fragmentation



high-energy **primary beam** of heavy ions (e.g. ^{18}O , ^{48}Ca , U...) on **thin** target of light element (Be or C)
⇒ fragmentation/fission produces many **exotic fragments** at $\approx v_{\text{beam}}$
Sorted in **fragment separator**

In-flight projectile fragmentation



high-energy **primary beam** of heavy ions (e.g. ^{18}O , ^{48}Ca , U...) on **thin** target of light element (Be or C)
⇒ fragmentation/fission produces many **exotic fragments** at $\approx v_{\text{beam}}$
Sorted in **fragment separator**

Used for **high-energy** reactions (KO, breakup...)

Examples : FRIB (MSU), RIKEN, GSI/FAIR, GANIL

[see A. Spyrou on Monday & T. Nilsson on Thursday]

FRIB @ MSU

FACILITY FOR RARE ISOTOPE BEAMS (FRIB)



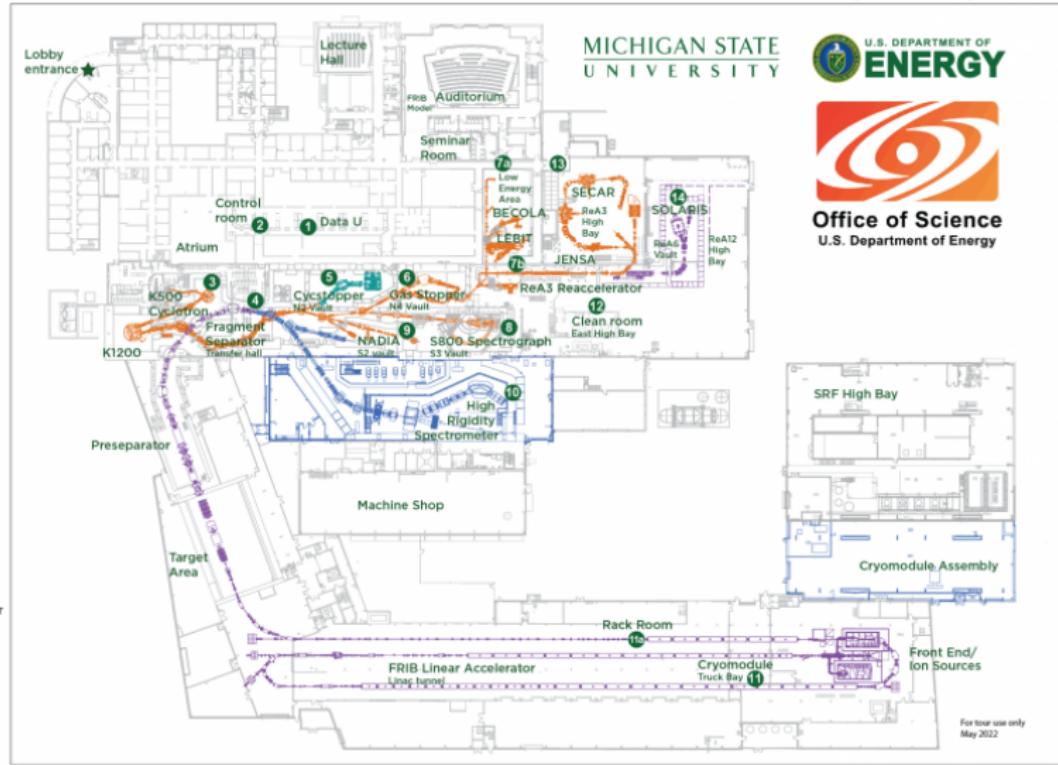
4 Fragment Separator



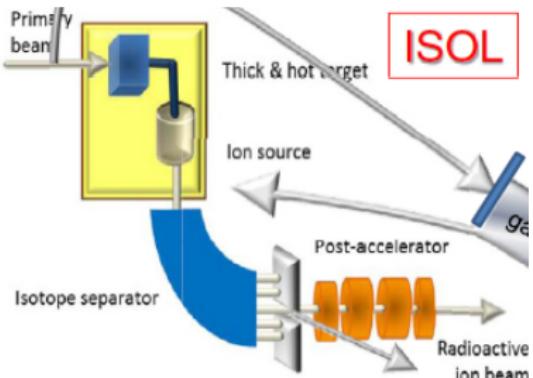
5 Cycstopper



8 S800 Spectrograph & GRETA Gamma-ray detector

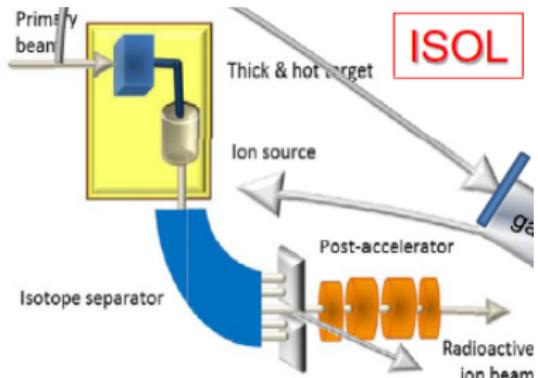


ISOL : Isotope Separation On Line



high-energy/intensity **primary beam**
of light nuclei (e.g. protons)
on **thick target** of heavy elements (UC_x)
⇒ spallation/fragmentation produces
exotic fragments
Diffuse in the target and
effuse to an **ion source**
Selected using dipole magnet (A/Q)

ISOL : Isotope Separation On Line

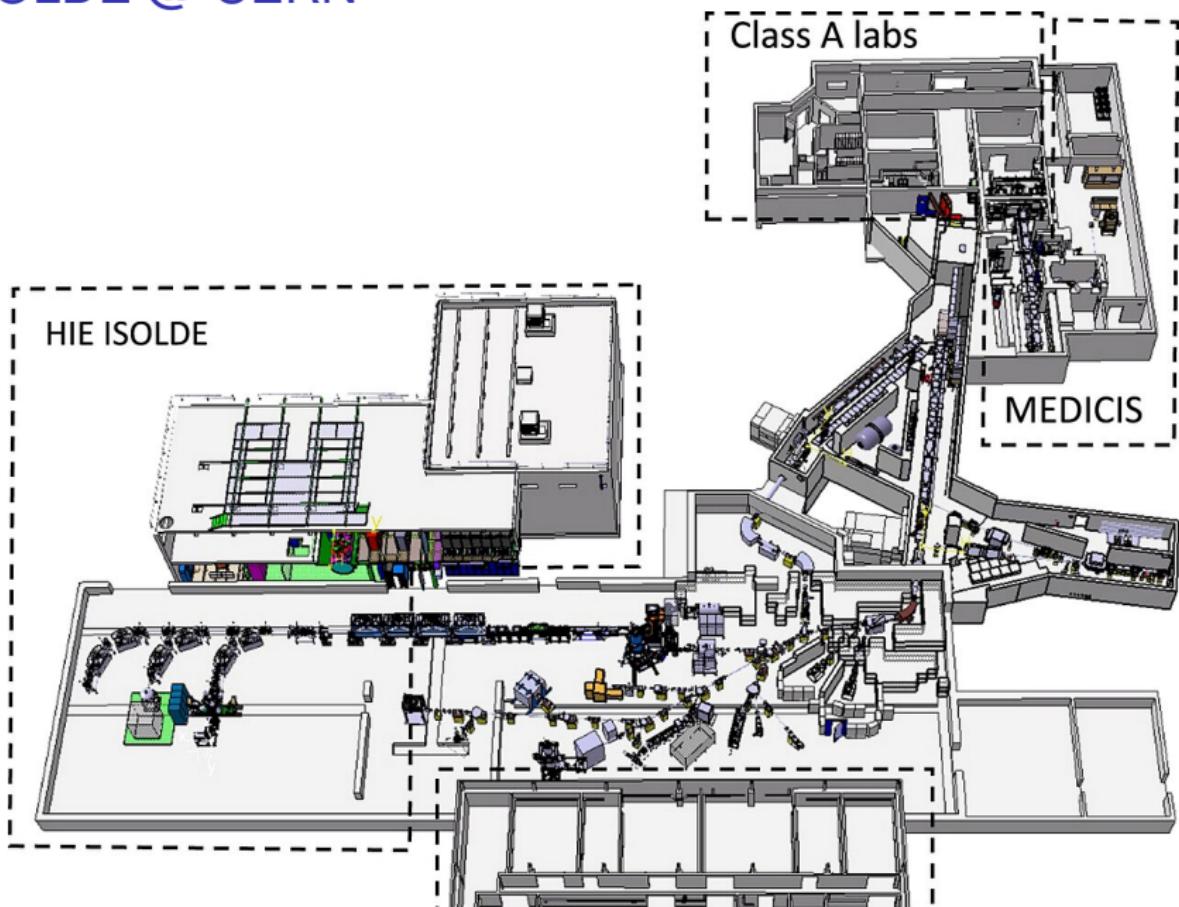


high-energy/intensity **primary beam**
 of light nuclei (e.g. protons)
 on **thick target** of heavy elements (UC_x)
 ⇒ spallation/fragmentation produces
exotic fragments
Diffuse in the target and
 effuse to an **ion source**
Selected using dipole magnet (A/Q)

Either used directly (mass measurement, radioactive decay...) or post-accelerated for reactions (e.g. astrophysical energy)

Examples : ISOLDE (CERN), TRIUMF, SPIRAL (GANIL)

ISOLDE @ CERN



1 Basic features in nuclear structure

- Liquid-drop model
- Shell model

2 Ab initio nuclear models

3 Radioactive-Ion Beams

4 Oddities far from stability

- Halo nuclei

5 Summary

Halo nuclei

Exotic structure discovered by I. Tanihata

[PLB 160, 380 (1985)]

Very large matter radius ($R \gg A^{1/3}$)

Seen as **core** + one or two **neutrons** at large distance

- Light, **neutron-rich** nuclei
- small S_n or S_{2n}

One-neutron halo

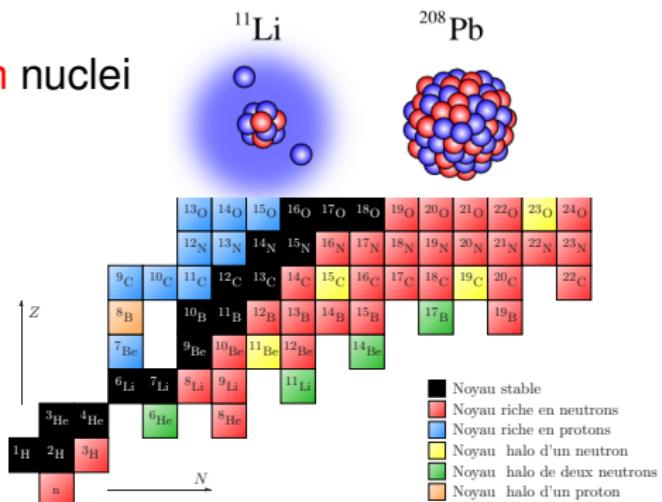
$^{11}\text{Be} \equiv ^{10}\text{Be} + \text{n}$

$^{15}\text{C} \equiv ^{14}\text{C} + \text{n}$

Two-neutron halo

$^6\text{He} \equiv ^4\text{He} + \text{n} + \text{n}$

$^{11}\text{Li} \equiv ^9\text{Li} + \text{n} + \text{n}$



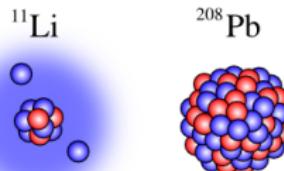
[R. Garreau's talk on Tuesday & Q. Bozé's talk on Thursday]

Halo nuclei

Exotic structure discovered by I. Tanihata [PLB 160, 380 (1985)]

Very large matter radius ($R \gg A^{1/3}$)

Seen as **core** + one or two **neutrons** at large distance



- Light, **neutron-rich** nuclei
- small S_n or S_{2n}

One-neutron halo

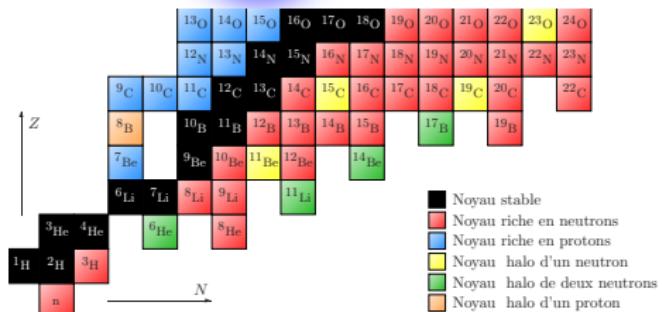
$^{11}\text{Be} \equiv ^{10}\text{Be} + \text{n}$

$^{15}\text{C} \equiv ^{14}\text{C} + \text{n}$

Two-neutron halo

$^6\text{He} \equiv ^4\text{He} + \text{n} + \text{n}$

$^{11}\text{Li} \equiv ^9\text{Li} + \text{n} + \text{n}$



[R. Garreau's talk on Tuesday & Q. Bozé's talk on Thursday]

Two-neutron halo nuclei are **Borromean**...

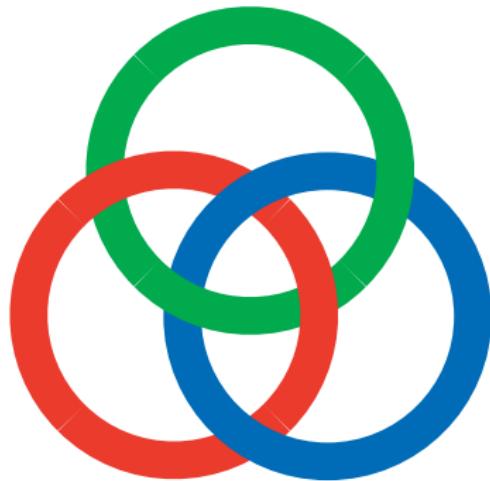
$c+n+n$ is bound but not two-body subsystems

e.g. ^6He bound but not ^5He nor ^2n

Borromean nuclei

Named after the Borromean rings...

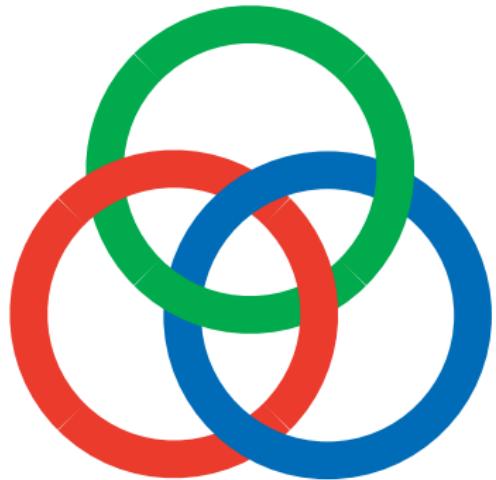
[M. V. Zhukov *et al.* Phys. Rep. 231, 151 (1993)]



Borromean nuclei

Named after the Borromean rings...

[M. V. Zhukov *et al.* Phys. Rep. 231, 151 (1993)]



Summary

Liquid-drop and shell model describe qualitatively stable nuclei
Nowadays *ab initio* nuclear-structure models from first principles

RIBs enable us to study nuclear structure far from stability
New exotic structure discovered :

- **halo** nuclei
diffuse halo around a compact core
- **shell inversions** or **shell collapse**
- nuclei beyond the **dripline** (resonant ground state)

RIB can be used to study reactions of **astrophysical** interest...

Properties ISOL

- Low beam energy
may require post-acceleration
- **Low** beam intensity
- Not all elements produced
 - ▶ Slow
 - ▶ Chemically limited
- **Good** beam quality :
can use chemistry and atomic
physics to select fragments

Properties

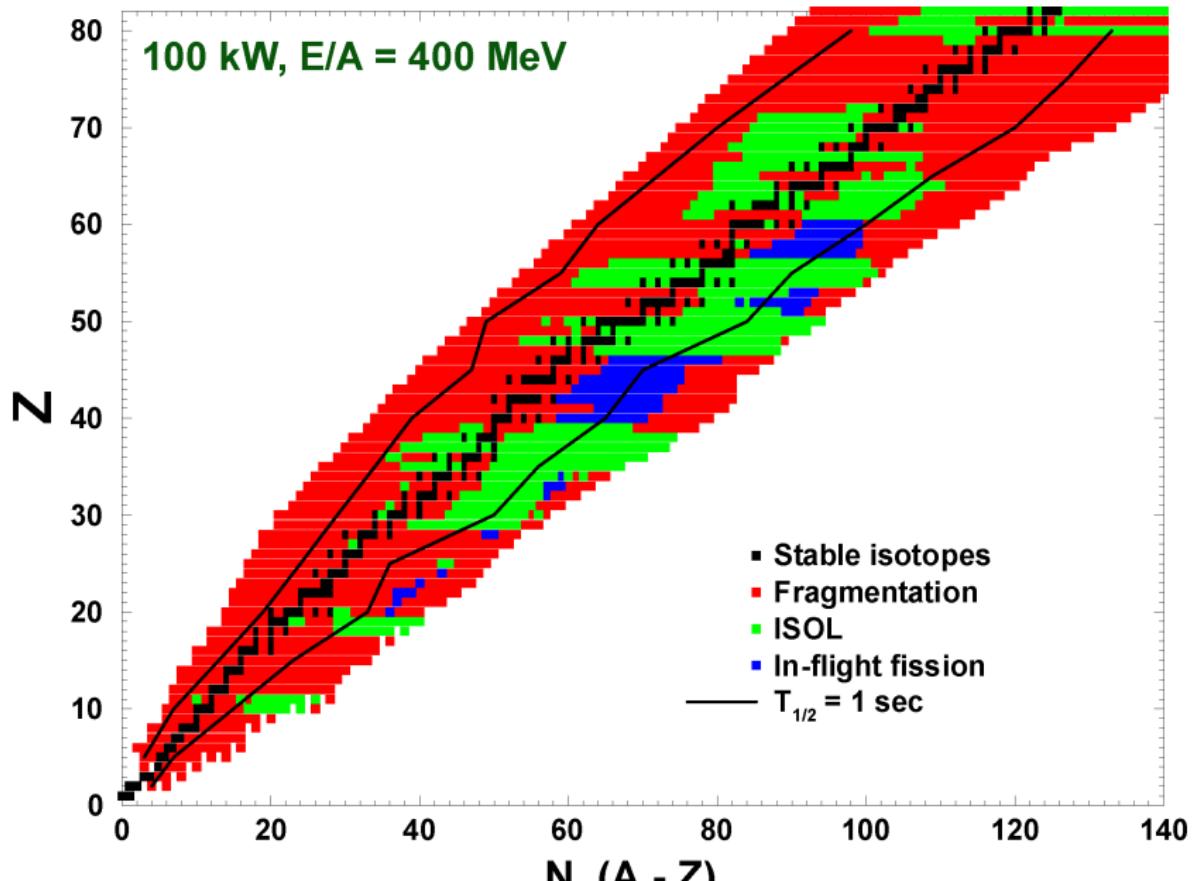
ISOL

- Low beam energy
may require post-acceleration
- **Low** beam intensity
- Not all elements produced
 - Slow
 - Chemically limited
- **Good** beam quality :
can use chemistry and atomic
physics to select fragments

In-flight

- High beam energy
 $v_{\text{fragments}} \approx v_{\text{beam}}$
- **High** beam intensity
- **Efficient** production
 - Fast
 - Chemically independent
- Many **fragments** in beam
⇒ need ion ID

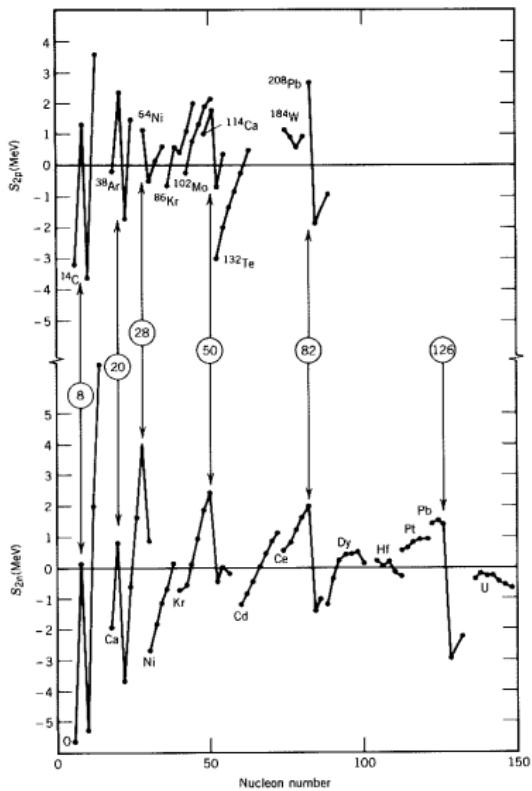
Choose according what you want to measure



Two-nucleon separation energy

Same magic numbers in S_{2p} and S_{2n}

⇒ more bound at shell closure
cf. ionisation energies of atoms



No-Core Shell Model

One should be able to account for the fermion nature of nucleons
 \Rightarrow wave function must be antisymmetric
 \Rightarrow basis states built as **Slater determinants**
 of 1-body mean-field wave functions ϕ_{ν_i}

$$\begin{aligned} \langle \xi_1 \xi_2 \dots \xi_A | \Phi_{[\nu]} \rangle &= \mathcal{A} \phi_{\nu_1}(\xi_1) \phi_{\nu_2}(\xi_2) \dots \phi_{\nu_A}(\xi_A) \\ &= \frac{1}{\sqrt{A!}} \begin{vmatrix} \phi_{\nu_1}(\xi_1) & \phi_{\nu_1}(\xi_2) & \cdots & \phi_{\nu_1}(\xi_A) \\ \phi_{\nu_2}(\xi_1) & \phi_{\nu_2}(\xi_2) & \cdots & \phi_{\nu_2}(\xi_A) \\ \vdots & \vdots & \ddots & \vdots \\ \phi_{\nu_A}(\xi_1) & \phi_{\nu_A}(\xi_2) & \cdots & \phi_{\nu_A}(\xi_A) \end{vmatrix} \end{aligned}$$

The shell model uses harmonic-oscillator wave functions for ϕ_{ν_i}
 The basis size increases with $A \Rightarrow$ limited to light nuclei