Zeta Functions via Periods of Calabi-Yau Hypersurfaces in Non-Fano Toric Varieties

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Main Question:

How can one obtain $\zeta_p(X_{\varphi}, T)$ from $\vartheta_1(\mathbb{P}_{\Delta}, X)$?

"Zeta from theta"

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Mirror symmetry

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Toric Varieties

Toric varieties are algebraic varieties that contain an dense algebraic torus [Dan78, CLS11].

- Two constructions:
 - Newton polytope $\Delta \rightsquigarrow \mathbb{P}_{\Delta}$ (vertices \iff monomials)
 - Fan or spanning polytope $\Sigma \rightsquigarrow \mathbb{P}_{\Sigma}$ (vectices \iff divisors)
- When $\Sigma^{(1)}$ are normal to facets of Δ , we have $\mathbb{P}_{\Delta} \cong \mathbb{P}_{\Sigma}$.



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Calabi-Yau hypersurfaces in toric varieties

- In string compactifications, spacetime is of the form $\mathbb{R}^{3,1} \times X$ where X is a compact Calabi-Yau manifold [CHSW85].
- By the adjunction formula, the anticanonical divisor $-K_{\mathbb{P}_{\Delta}} = \sum_{\rho=1}^{r} D_{\rho}$ has trivial canonical class and hence $-K_{\mathbb{P}_{\Delta}} = X$ is a Calabi-Yau manifold.

Example 1.2 (Anticanonical hypersurfaces in projective space)

$$-K_{\mathbb{P}^n} = (n+1)H \implies \begin{cases} -K_{\mathbb{P}^2} \text{ is the cubic elliptic curve} \\ -K_{\mathbb{P}^3} \text{ is the quartic K3} \\ -K_{\mathbb{P}^4} \text{ is the quintic CY3} \\ \vdots \end{cases}$$

where H is the hyperplane class.

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Gauged linear sigma models

Witten introduced gauged linear sigma models (GLSMs) – $U(1)^s$ gauge theories whose charge matrix is given by the *Mori* vectors of a toric variety [Wit93].

Definition 1.3 (Mori vectors of a toric variety)

Let $\Sigma^{(1)} = \{v_1, \ldots, v_r\}$ be the 1 dimensional cones of a toric variety. Let $\overline{v}_i = (v_i, 1)$ and $v_0 = \vec{0}$. Then the *Mori vectors* $\ell_i^{(a)}$ are defined by

$$\sum_{i=0}^{7} \ell_i^{(a)} \overline{v}_i = 0$$

Example 1.4 (\mathbb{P}^2)

$$\Sigma_{\mathbb{P}^2}^{(1)} = \{(1,0), (0,1), (-1,-1)\} \implies \ell^{(1)} = (-3;1,1,1)$$

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Fano, Semi-Fano, Calabi-Yau, non-Fano

 $c_1(\mathbb{P}_\Delta)$ is the Poincarè dual to $-K_{\mathbb{P}_\Delta}.$ Write

$$-K_{\mathbb{P}_{\Delta}} = \sum_{a=1}^{h^{1,1}(\mathbb{P}_{\Delta})} (-\ell_0^{(a)}) D_a$$

in a basis of ample divisors.

Definition 1.5 (Fano, Semi-Fano, Calabi-Yau, non-Fano) • If $(-\ell_0^{(a)}) > 0$ then \mathbb{P}_Δ is Fano. • If $(-\ell_0^{(a)}) \ge 0$ then \mathbb{P}_Δ is semi-Fano. • If $\ell_0^{(a)} = 0$ then \mathbb{P}_Δ is Calabi-Yau. • If $\ell_0^{(a)} \in \mathbb{Z}$ then \mathbb{P}_Δ is "of general type" or non-Fano.

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Hirzebruch scrolls

• The Hirzebruch scrolls $\mathcal{F}_m^{(n+1)} = \mathbb{P}(\mathcal{O}_{\mathbb{P}^n}(m) \oplus \mathcal{O}_{\mathbb{P}^1})$ are \mathbb{P}^n fibrations over \mathbb{P}^1 .



•
$$-K_{\mathcal{F}_m^{(n+1)}} = (n+1)S + (2-m)F$$

• Non-Fano for $m>2$



Mori vectors for $\mathcal{F}_m^{(2)}$:

$$\ell^{(1)} = (-2; 1, 1, 0, 0)$$

$$\ell^{(2)} = (-(2-m); -m, 0, 1, 1)$$

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Kreuzer-Skarke database of reflexive polytopes

The Kreuzer-Skarke database lists all reflexive polytopes in 3 and 4 dimensions [KS].

Definition 1.6

- A polytope Δ is reflexive if
 - All facets are supported by an affine hyperplane of the form $\{m\in \mathbb{R}^n\,|\,\langle m,n\rangle=-1\}$
 - $\bullet~\Delta$ has the orgin as its unique interior point.

Theorem 1.7

 Δ is reflexive if and only if \mathbb{P}_{Δ} is Fano.

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Numbers of reflexive polytopes

• 16 in two dimensions

- 4,319 in three dimensions
- 473,800,776 in four dimensions
- ?? in five+ dimensions

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Sigma models

GLSMs have phases, among which are NLSMs and LG models.

Definition 1.8 (Non-linear Sigma Model)

Given an embedding $\phi: \mathsf{WS} \to X$ into a Calabi-Yau,

$$\begin{split} S_{\text{NLSM}} &= i \int \left(\frac{1}{2} (g_{j\bar{k}} + iB_{j\bar{k}}) \partial_{\tau} \phi^{j} \partial_{\bar{\tau}} \bar{\phi}^{\bar{k}} + \frac{1}{2} (g_{j\bar{k}} - iB_{j\bar{k}}) \partial_{\bar{\tau}} \phi^{j} \partial_{\tau} \bar{\phi}^{\bar{k}} \\ &+ \frac{i}{2} g_{j\bar{j}} \overline{\psi}_{-}^{\bar{j}} \nabla_{\tau} \psi_{-}^{j} + \frac{i}{2} g_{j\bar{j}} \overline{\psi}_{+}^{j} \nabla_{\bar{\tau}} \psi_{+}^{j} \\ &+ \frac{1}{4} R_{j\bar{j}k\bar{k}} \psi_{+}^{j} \overline{\psi}_{+}^{\bar{j}} \psi_{-}^{k} \overline{\psi}_{-}^{\bar{k}} \right) d\tau \wedge d\bar{\tau}. \end{split}$$



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LG models

Landau-Ginzburg models can be thought of as sigma models that are deformed by a potential W which is a function of superfields.

Definition 1.9 (Landau-Ginzburg Model)

Given $\phi : WS \to X$ and a superpotential $W : (\mathbb{C}^*)^n \to \mathbb{C}$,

$$\begin{split} S_{\mathsf{L}\mathsf{G}} &= \int \left(\partial_{\tau} \phi^{j} \partial_{\bar{\tau}} \overline{\phi}^{\bar{j}} + \partial_{\bar{\tau}} \phi^{j} \partial_{\tau} \overline{\phi}^{\bar{j}} \\ &+ i \overline{\psi}_{-}^{\bar{j}} \nabla_{\tau} \psi_{-}^{j} + i \overline{\psi}_{+}^{\bar{j}} \nabla_{\bar{\tau}} \psi_{+}^{j} \\ &- \frac{1}{4} \partial_{j} W \partial_{\bar{j}} \overline{W} - \frac{1}{2} \partial_{j} \partial_{k} W \psi_{+}^{j} \psi_{-}^{k} - \frac{1}{2} \partial_{\bar{j}} \partial_{\bar{k}} \overline{W} \overline{\psi}_{-}^{\bar{j}} \overline{\psi}_{+}^{\bar{k}} \right) d\tau \wedge d\bar{\tau} \end{split}$$

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BRST cohomology

A sigma model has supersymmetry generated by

$$\begin{aligned} Q_{+} &= \psi^{j} p_{j} = i \overline{\partial}^{\dagger} & Q_{-} &= \overline{\psi}^{j} p_{j} = -i \partial \\ \overline{Q}_{+} &= \overline{\psi}^{\overline{j}} p_{\overline{j}} = -i \overline{\partial} & \overline{Q}_{-} &= \psi^{\overline{j}} p_{\overline{j}} = i \partial^{\dagger} \end{aligned}$$

Given a nilpotent supercharge $Q^2 = 0$, construct the physical operators in a "topologically twisted" theory

$$H^*_Q(\Sigma) = \ker Q / \operatorname{im} Q$$

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A-model and B-model

The A-model and the B-model are TQFTs that can be obtained from topologically twisting a sigma model.

• A-model

$$Q_A = -i(\partial + \overline{\partial}) \cong d \implies H^*_{Q_A} \cong H^*_{\mathsf{dR}}(X)$$

• B-model

$$Q_B \cong \overline{\partial} \implies H^*_{Q_B} \cong \bigoplus_{p,q=0}^n H^{0,p}(X, \wedge^q TX) \cong H^*_{\overline{\partial}}(X)$$



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Invariants from correlation functions

The correlation functions in the A-model and the B-model are topological invariants of the Calabi-Yau.

• A-model: Gromov-Witten invariants

$$\langle \mathcal{O}_1 \dots \mathcal{O}_s \rangle = \sum_{\beta \in H_2(X,\mathbb{Z})} \int_{\mathcal{M}_{\Sigma}(X,\beta)} e^{-(\omega - iB) \cdot \beta} \mathsf{ev}_1^* \omega_{D_1} \wedge \dots \wedge \mathsf{ev}_s^* \omega_{D_s}$$
$$= \int_X \omega_1 \wedge \dots \wedge \omega_s + \sum_{\substack{\beta \in H_2(X,\mathbb{Z}) \\ \beta \neq 0}} n_{\beta,D_1\dots D_s} e^{-(\omega - iB) \cdot \beta}.$$

• B-model: Period integrals

$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \rangle = \int_X \Omega \wedge (\nabla_{\theta_1} \nabla_{\theta_2} \nabla_{\theta_3} \Omega)$$

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Gromov-Witten invariants

- GW invariants β are counts of stable maps whose image land on a fixed curve class β ∈ H₂(X, ℤ).
- Defined as integrals over moduli spaces of maps

$$D_{0,1}(\mathbb{P}_{\Delta},\beta) = \int_{[\mathcal{M}_{0,1}((\mathbb{P}_{\Delta},X),\beta)]^{virt}} \psi^{\beta \cdot X - 2} \mathsf{ev}^{\star}[\mathsf{pt}].$$



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Mirror symmetry in physics

- Mirror symmetry was first discovered by Greene and Plesser [GP90] for quintic threefold.
- Used to count rational curves by Candelas, de la Ossa, Green, Parkes [CDGP91].
- Mirror symmetry is T-duality [SYZ96] in the large complex structure limit.

Definition 1.10

Two Calabi-Yau manifolds X and \check{X} are a mirror pair (in physics) if

```
A-model(X) \cong B-model(\check{X})
A-model(\check{X}) \cong B-model(X)
```

which can be checked by comparing correlation functions.

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Mirror symmetry for Calabi-Yau manifolds

- It is a necessary but not sufficent condition that the Hodge diamonds of X and \check{X} are reflections of each other.
- - Batyrev: mirror symmetry as polar duality $\Delta \leftrightarrow \Delta^{\star}$ [Bat].
 - Can compute invariants of X in terms of invariants of X.
 - Ratios of periods are generating functions for GW invariants.
 - Deformation method for computing zeta functions [Dwo62].

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Mirror symmetry for Fano varieties

• Hori and Vafa [HV] found mirror LG model to sigma models with Fano toric variety target.

Definition 1.11

Let \mathbb{P}_{Σ} be a toric variety. Then the mirror Hori-Vafa potential is given by

$$W_{\Sigma} = \sum_{\rho \in \Sigma^{(1)}} \varphi^{\beta_{\rho}} x^{m_{\rho}}$$

where $m_{\rho} = (m_1, \ldots, m_n)$ is the primitive ray generator, $x^{m_{\rho}} = x_1^{m_1} \ldots x_n^{m_n}$, and $\varphi = (\varphi_1, \ldots, \varphi_s)$ are the complex structure moduli of the mirror $(s = h^{1,1}(\mathbb{P}_{\Sigma}))$.

W_{Σ} and curve classes

The powers of the complex structure moduli in W_{Σ} are determined by effective curve classes β_{ρ} :

• If (ρ_1, \ldots, ρ_k) is a collection of rays with $\sum_{i=1}^k m_{\rho_i} = 0$, then $\sum_{i=1}^k \varphi^{\rho} x^{m_{\rho}} = \varphi^{\beta} := \varphi_1^{d_1} \cdots \varphi_r^{d_r}$, where $\beta = d_1 \beta_1 + \ldots + d_r \beta_r$ is the effective curve class whose intersections with toric divisors are given by (ρ_1, \ldots, ρ_k) .

Example 1.12 ($W_{\Sigma_{\mathbb{D}^2}}$)



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Classical and quantum periods

Definition 1.13

The classical period of a Laurent polynomial $W\in \mathbb{C}[\varphi][x_1^{\pm 1},\ldots,x_n^{\pm 1}]$ is

$$\pi_W(\varphi) = \left(\frac{1}{2\pi i}\right)^n \int_{\Gamma_0} \frac{d\log x_1 \wedge \dots \wedge d\log x_n}{1 - W} = \sum_{k > 0} \operatorname{const}(W^k),$$

where const is with respect to the variables x_1, \ldots, x_n .

Definition 1.14

The regularized quantum period of X is

$$G_X(\varphi) = \sum_{\beta} (\beta \cdot (-K_X))! D_{0,1}(X,\beta) \varphi^{\beta}.$$

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Periods and Picard-Fuchs equations

The periods of a Calabi-Yau manifold satisfy a system of partial differential equations $\mathcal{L}_a \varpi = 0$ called the Picard-Fuchs operators

$$\mathcal{L}_a \in \mathbb{Q}\left[\varphi_b, \frac{\partial}{\partial \varphi_b}
ight]$$

Example 1.15

The local Picard-Fuchs operators [CKYZ99] associated to the canonical bundle $K_{\mathbb{P}_{\Sigma}}$ are given by the Mori vectors:

$$\mathcal{L}_{a} = \prod_{\ell_{i}^{(a)} > 0} \prod_{j=0}^{\ell_{i}^{(a)} - 1} \left(\sum_{b=1}^{s} \ell_{i}^{(b)} \theta_{b} - j \right) - \varphi_{a} \prod_{\ell_{i}^{(a)} < 0} \prod_{j=0}^{-\ell_{i}^{(a)} - 1} \left(\sum_{b=1}^{s} \ell_{i}^{(b)} \theta_{b} - j \right)$$

where $\theta_a = \varphi_a \partial_{\varphi_a}$ (no summation).

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Picard-Fuchs equations for $K_{\mathcal{F}_3^{(2)}}$

Example 1.16 (\mathcal{L}_a for $\mathcal{F}_3^{(2)}$)

The Picard-Fuchs operators for $K_{\mathcal{F}_3^{(2)}}$ are

$$\mathcal{L}_{1} = \theta_{1}(\theta_{1} - 3\theta_{2}) - \varphi_{1}(-2\theta_{1} + \theta_{2})(-2\theta_{1} + \theta_{2} - 1)$$

$$\mathcal{L}_{2} = (-2\theta_{1} + \theta_{2})\theta_{2} - \varphi_{2}(\theta_{1} - 3\theta_{2})(\theta_{1} - 3\theta_{2} - 1)(\theta_{1} - 3\theta_{2} - 2)$$

They have the solutions (obtained by recursion)

$$F_{1} = \sum_{\substack{n_{1}, n_{2} \ge 0\\n_{1} \ge 3n_{2}}} (-1)^{n_{2}} \frac{\Gamma(2n_{1} - n_{2})}{\Gamma(n_{1})\Gamma(n_{1} - 3n_{2} + 1)\Gamma^{2}(n_{2} + 1)} \varphi_{1}^{n_{1}} \varphi_{2}^{n_{2}}$$

$$F_{2} = \sum_{\substack{n_{1}, n_{2} \ge 0\\n_{1} \le 3n_{2}}} (-1)^{n_{1} + n_{2}} \frac{\Gamma(3n_{2} - n_{1})}{\Gamma(n_{1} + 1)\Gamma(n_{2} - 2n_{1} + 1)\Gamma^{2}(n_{2} + 1)} \varphi_{1}^{n_{1}} \varphi_{2}^{n_{2}}$$

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Local GW invariants for non-Fano surfaces

• A curve on a non-Fano surface can be perturbed away from the zero locus away from the canonical bundle.



• To obtain the correct GW invariants, we need to include corrections to W_{Σ} account for such curves.

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Scattering diagrams

Obtain $\vartheta_1(\mathbb{P}_{\Sigma}, X) = W_{\Sigma} + W'$ via Tim Gräfnitz's scattering.sage [Gra].



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W_{Σ} and $\vartheta_1(\mathbb{P}_{\Sigma}, X)$

Example 1.17 (Corrected superpotentials for Hirzebruch surfaces)

$$\begin{split} \vartheta_1(\mathcal{F}_0^{(2)}) &= t \cdot \left(x + y + \frac{\varphi_1}{x} + \frac{\varphi_2}{y} \right) \\ \vartheta_1(\mathcal{F}_1^{(2)}) &= t \cdot \left(x + y + \frac{\varphi_1}{x} + \frac{\varphi_1 \varphi_2}{xy} \right) \\ \vartheta_1(\mathcal{F}_2^{(2)}) &= t \cdot \left(x + y + \frac{\varphi_1}{x} + \frac{\varphi_1 \varphi_2}{x} \left(1 + \frac{\varphi_1}{xy} \right) \right) \\ \vartheta_1(\mathcal{F}_3^{(2)}) &= t \cdot \left(x + y + \frac{\varphi_1}{x} + \frac{\varphi_1 \varphi_2 y}{x} \left(1 + \frac{\varphi_1}{xy} \right)^2 \right) \\ \vartheta_1(\mathcal{F}_4^{(2)}) &= t \cdot \left(x + y + \frac{\varphi_1}{x} + \varphi_1 \varphi_2 y \left(1 + \frac{\varphi_1}{xy} \right) + \frac{\varphi_1 \varphi_2 y^2}{x} \left(1 + \frac{\varphi_1}{xy} \right)^3 \right) \end{split}$$

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GW invariants for $\mathcal{F}_3^{(2)}$



n_1, n_2	0	1	2	3	4	5	6	7	8
0	1	1	0	0	0	0	0	0	0
1	0	0	-2	-4	-6	-8	-10	-12	-14
2	0	0	0	0	5	35	135	385	910
3	0	0	0	0	0	0	-32	-400	-2592
ame as fo	or F	(2) 1	with n	$n_1 \mapsto r$	$n_1 + r$	$a_2!$ (F	$\overline{f}_m^{(n)} \cong J$	$\mathcal{F}_{m \mod n}^{(n)}$	")
				Lathw	bod	Non-Fano	Zeta		

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Mirror symmetry for non-Fano varieties

Definition 1.18

A function $W \in \mathbb{C}[\varphi_1, \ldots, \varphi_r][x_1^{-1}, \ldots, x_n^{-1}]\llbracket x_1, \ldots, x_n \rrbracket$ is called *mirror dual* to (X, D) if the classical period

$$\pi_{\mathit{W}}(arphi) = \sum_{k>0} \operatorname{const}(\mathit{W}^k) \in \mathbb{C}[arphi_1, \dots, arphi_r]$$

is equal to the regularized quantum period $G_X(\varphi)$ of X.

Theorem 1.19 ([BGL24])

For every point P inside a chamber in the scattering diagram, $t^{-1}\vartheta_1(X,D)_P$ is a mirror potential for (X,D), in the sense of Definition 1.18,

$$\pi_{t^{-1}\vartheta_1(X,D)_P}(\varphi) = G(\varphi).$$

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Mirror symmetry for non-Fano varieties

Theorem 1.20 ([BGL24])

Let (X, D) be a smooth log Calabi-Yau pair with mirror dual potential W. Then, under the change of variables $Q_i = \varphi_i (t/y)^{d_i}$, with $d_i = \beta_i \cdot D$, we have

$$\vartheta_1(t,\varphi,y)_\infty = yM_W(Q),$$

where $\vartheta_1(t, \varphi, y)_{\infty} := \vartheta_1(X, D)_{\infty}$ is the theta function at infinity and M_W is the open mirror map defined by W.

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Number Theory

Calabi-Yau manifolds over finite fields Hasse-Weil zeta functions p-adic cohomology Type IIB supergravity and the attractor mechanism

Calabi-Yau manifolds over finite fields

- Let X_φ be a one parameter family of algebraic varieties defined by a polynomial P_φ ∈ Q[x₁,..., x_{n+1}] (x_i ∈ Pⁿ).
- \bullet We clear denominators to get a defining equation over $\mathbb Z.$
- Let \mathbb{F}_q be the finite field with $q = p^k$ (p prime) elements and

$$X_{\varphi}(\mathbb{F}_q) = |\{x \in \mathbb{F}_q^{n+1} \mid P_{\varphi}(x) = 0\}|.$$

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Zeta functions

Definition 2.1

The *local zeta function* or the *Hasse-Weil zeta function* of the manifold X_{φ} at the prime p is defined as

$$\zeta_p(X_{\varphi}, T) = \exp\left(\sum_{k=1}^{\infty} \frac{X_{\varphi}(\mathbb{F}_{p^k})}{k} T^k\right)$$

Example 2.2

For $X = \{pt\}$, we have

$$\prod_{p} \zeta_p(\{\mathsf{pt}\}, p^{-s}) = \zeta(s),$$

the Riemann zeta function.

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Weil conjectures

The Weil conjectures – proved by Dwork, Grothendieck, and Deligne – are concerned with the form of local zeta functions.

Theorem 2.3 (Rationality of ζ_p)

$$\zeta_p(X_{\varphi}, T) = \frac{R_p^{(1)} R_p^{(3)} \dots R_p^{(2d-1)}}{R_p^{(0)} R_p^{(2)} R_p^{(4)} \dots R_p^{(2d)}}$$

where
$$R_p^{(i)} \in \mathbb{Z}[T]$$
 with $\deg R_p^{(i)} = \dim H^i(X_{\varphi})$.

Example 2.4

Projective space only has even dimensional cohomology, so

$$\zeta_p(\mathbb{P}^n, T) = \frac{1}{(1 - T)(1 - pT)\dots(1 - p^nT)}$$

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Weil conjectures

For a Calabi-Yau 3-fold $X_{\varphi},$ the Weil conjectures imply the local zeta function is of the form

$$\zeta_p(X_{\varphi}, T) = \frac{R_p(X_{\varphi}, T)}{(1 - T)(1 - pT)^{h^{11}}(1 - p^2T)^{h^{11}}(1 - p^3T)}$$

In other words, the problem of computing $\zeta_p(X_{\varphi}, T)$ is reducted to computing the polynomial $R_p(X_{\varphi}, T)$, called the *Frobenius polynomial*.

 \implies This can be done with *p*-adic cohomology $H^k(X_{\varphi}, \mathbb{Q}_p)$

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Frobenius map

The Frobenius map acts on coordinates as

$$\operatorname{Frob}_p : \mathbb{K}^n \longrightarrow \mathbb{K}^n$$
$$x \longmapsto (x_1^p, x_2^p, \dots, x_n^p)$$

This defines a map $\mathrm{Frob}_p:X_\varphi/\mathbb{F}_p\to X_\varphi/\mathbb{F}_p$ since

$$P_{\varphi}(x^p) = P_{\varphi}(x)^p = 0 \mod p$$

Therefore

$$|\{\mathsf{Fixed points of Frob}_p\}| = X_{\varphi}(\mathbb{F}_p)$$

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Lefschetz fixed-point theorem

The pullback of the Frobenius map gives an automorphism of $p\mbox{-}{\rm adic}$ cohomology

$$\mathrm{Fr}_p = (\mathrm{Frob}_p)_* : H^k(X_\varphi, \mathbb{Q}_p) \longrightarrow H^k(X_\varphi, \mathbb{Q}_p)$$

and applying the Lefschetz fixed-point theorem to this map gives a formula for the point counts

$$X_{\varphi}(\mathbb{F}_{p^k}) = \sum_{\ell=0}^{2n} (-1)^{\ell} \operatorname{Tr}\left(\operatorname{Fr}_{p^k} \left| H^{\ell}(X_{\varphi}, \mathbb{Q}_p)\right)\right)$$

The characteristic polynomial of the inverse Frobenius map acting on the middle cohomology $H^n(X_{\varphi}, \mathbb{Q}_p)$ is exactly the polynomial $R_p(X_{\varphi}, T)$: you've been eigenvalued.

$$R_p(X_{\varphi}, T) = \det \left(I - T \mathsf{Fr}_p^{-1} \mid H^n(X_{\varphi}, \mathbb{Q}_p) \right)$$



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Computing Frobenius trace with periods

One can compute (following [CdlOK24])

$$R_p(X_{\varphi}, T) = \det \left(I - T \operatorname{Fr}_p^{-1} | H^3(X_{\varphi}, \mathbb{Q}_p) \right) = \det \left(I - T U_p(\varphi) \right)$$

where

$$U_p(\varphi) = \widetilde{E}(\varphi^p)^{-1} \varphi^{-p\epsilon} U_p(0) \varphi^{\epsilon} \widetilde{E}(\varphi).$$

and

$$E(\varphi)_a^{\ b} = \begin{pmatrix} \theta_a \varpi^b & \theta^a \varpi^b \\ \theta_a \varpi_b & \theta^a \varpi_b \end{pmatrix}$$

is a matrix of periods $\varPi = (\varpi^0, \varpi^a, \varpi_a, \varpi_0)$ in the derivative basis.

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Attactor mechanism

- Four dimensional $\mathcal{N} = 2$ black holes can be constructed by compactifying IIB supergravity on a Calabi-Yau threefold X_{φ} with complex structure parameter φ .
- The charges of the black hole are determined by a 3-cycle

$$\gamma = q_a A^a - p^a B_a \in H_3(X_{\varphi}, \mathbb{Z}),$$

which is wrapped by D3-branes.

 The value of φ at the horizon of the black hole is an attactor point φ = φ_{*} [Moo07]. Mirror symmetry Number theory Computing zeta functions Computing Zeta f

Attactor mechanism

If $X = X_{\varphi_*}$ is an attactor variety, the middle cohomology splits as a Hodge structure. This implies Fr_p^{-1} becomes block diagonal and hence $R_p(X, T)$ factorizes.



Factorization of $R_p(X_{\varphi}, T)$ independent of $p \implies$ Rank 2 attractor point [CdIOEvS20].

Code New examples Future directions

Computing zeta functions

Code New examples Future directions

Computational methods

There exists various software to compute periods and zeta functions

- ore_algebra: periods of PF operators
- CY3Zeta: zeta functions of CY3's
- controlledreduction: modified Griffiths-Dwork reduction

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controlled reduction

- The package controlled reduction by Edgar Costa computes $\zeta_p(X_\varphi, T)$ for non-degenerate fibers X_φ using a special version of Griffiths-Dwork reduction.
- Used this to generate training data for transformer (see my String Data 2024 talk).

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Dwork pencil

Example 3.1

The Dwork pencil is the family of K3 surfaces with defined by

$$P_{\varphi} = x_1^4 + x_2^4 + x_3^4 + x_4^4 + \varphi x_1 x_2 x_3 x_4$$

Recall that for K3s dim $H^2(X_{\varphi}) = 22$.

Used controlled reduction to compute $R_p(X_{\varphi}, T)$ for all values of φ at fixed p. For example:

$$R_{11}(X_1, T) = (1 + 11T)^6 (1 - 11T)^{13} (1 + 18T + 121T^2)$$

sage: from pycontrolledreduction import controlledreduction
sage: R.<x,y,z,w> = ZZ[]
sage: controlledreduction(x⁴ + y⁴ + z⁴ + w⁴ + x*y*z*w, 11, False).factor()
(-1) * (11*T + 1)⁶ * (11*T - 1)¹³ * (121*T² + 18*T + 1)

Note that the Frobenius polynomial has degree 21 here since the factor (1 - pT) from the polarization was omitted.

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Training Dwork pencil

Example 3.2 (p = 7)

$$R_7(X_0, T) = (1 + 7T)^{10}(1 - 7T)^{11}$$

$$R_7(X_1, T) = (7T + 1)^6(1 - 7T)^{13}(49T^2 + 10T + 1)$$

$$R_7(X_2, T) = (1 - 7T)^9(7T + 1)^{10}(49T^2 - 6T + 1)$$

$$R_7(X_5, T) = (1 - 7T)^9(7T + 1)^{10}(49T^2 - 6T + 1)$$

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Lathwood Non-Fano Zeta

Code New examples Future directions

ToricZeta

- In collaboration with Pyry Kuusela at Mainz and Michael Stepniczka at Cornell, we are developing a Python package to compute zeta functions for families of Calabi-Yau hypersurfaces X_φ in toric varieties P_Δ.
- Motivation is to automate the Mathematica code CY3Zeta from arXiv:2405.08067 and port to Python, then include in CYTools.
- Our computation involves the period vector Π and linear algebra over Q_p.

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Lathwood

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New Picard-Fuchs operators

- Recent work producted new PF operators for non-Fano toric varieties [BGL24]. Compute their periods with recursion, then feed to ToricZeta to compute new zeta functions.
- Code up Griffiths-Dwork reduction in Macaulay2 to obtain a map $\Delta \mapsto \{\mathcal{L}_a\}$

Example 3.3 (New $\{\mathcal{L}_a\}$ for $\mathcal{F}_3^{(2)}$)

The new Picard-Fuchs system $\{\mathcal{L}_a^{\vartheta_1}\}$ for $\mathcal{F}_3^{(2)}$ is related to the old system by $\theta_1 \mapsto \theta_1 - \theta_2$.

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New Picard-Fuchs operators



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Hirzebruch scrolls

Apply ideas of last slide to $\mathcal{F}_3^{(4)}$, a non-Fano toric 4-fold with Calabi-Yau 3-fold as anticanonical divisor



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Ratios of weight-four special L-values are equal to an infinite series whose summands are formed out of genus-0 Gromov–Witten invariants [CdIOM25].

 \implies Apply to invariants from $\vartheta_1!$

Code New examples Future directions

Thank you!

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