

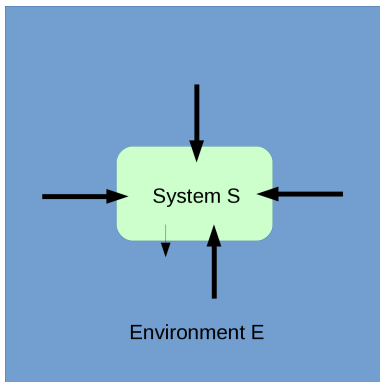
# Probing new physics with open quantum systems

Christian Käding (Atominstitut, TU Wien)

YOUNGST@RS - Quantum Sensing for Fundamental Physics,  
20 November 2024

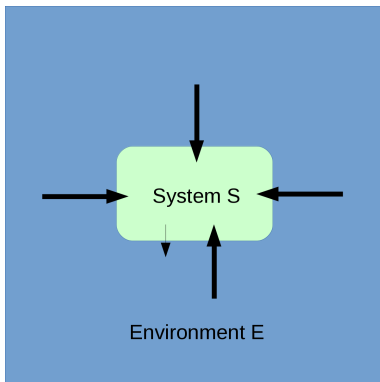


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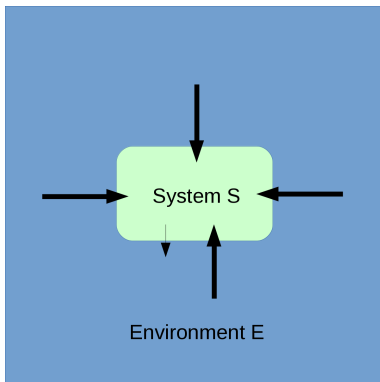
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- Examples: thermal bath, electromagnetic field, gravity, scalar field,...

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# Reduced density matrix

- Consider system scalar field  $\phi$  with environment scalar field  $\chi$
- Combination of system and environment:  $\hat{\rho}_{\phi\chi}(t)$
- Assume  $\hat{\rho}_{\phi\chi}(0) = \hat{\rho}_{\phi}(0) \otimes \hat{\rho}_{\chi}(0)$
- Reduced density matrix:  $\hat{\rho}_{\phi}(t) = \text{Tr}_{\chi} \hat{\rho}_{\phi\chi}(t)$
- Project into single-particle momentum basis:

$$\rho_{\phi}(\mathbf{p}, \mathbf{p}'; t) := \langle \mathbf{p} | \hat{\rho}_{\phi}(t) | \mathbf{p}' \rangle$$

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## Phys.Rev.D 107 (2023) 1, 016005

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 \rho_\phi(\mathbf{p}, \mathbf{p}'; t) &= \lim_{\substack{x^{0(r)} \rightarrow t^+ \\ y^{0(r)} \rightarrow 0^-}} \int d\Pi_{\mathbf{k}} d\Pi_{\mathbf{k}'} \rho_\phi(\mathbf{k}, \mathbf{k}'; 0) \\
 &\times \int_{\mathbf{x}\mathbf{x}'\mathbf{y}\mathbf{y}'} e^{-i(\mathbf{p}\cdot\mathbf{x}-\mathbf{p}'\cdot\mathbf{x}')+i(\mathbf{k}\cdot\mathbf{y}-\mathbf{k}'\cdot\mathbf{y}')} \partial_{x^0, E_{\mathbf{p}}^\phi} \partial_{x^{0r}, E_{\mathbf{p}'}^\phi}^* \partial_{y^0, E_{\mathbf{k}}^\phi}^* \partial_{y^{0r}, E_{\mathbf{k}'}^\phi} \\
 &\times \int \mathcal{D}\phi^\pm e^{i\hat{S}_\phi[\phi]} \phi^+(x) \phi^-(x') \exp \left\{ i\hat{S}_{\phi, \text{int}}[\phi; t] \right\} \hat{\mathcal{F}}[\phi; t] \phi^+(y) \phi^-(y')
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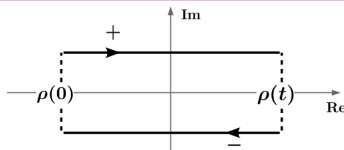
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$$\widehat{\mathcal{F}}[\phi; t] = \left\langle \exp \left\{ i(\widehat{S}_{\chi, \text{int}}[\chi; t] + \widehat{S}_{\text{int}}[\phi, \chi; t]) \right\} \right\rangle_\chi$$

Schwinger-Keldysh closed time path:



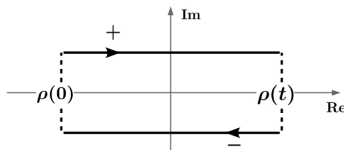
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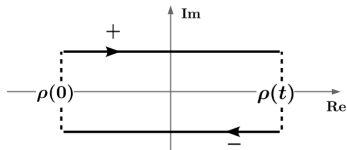
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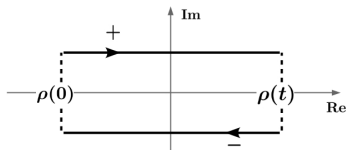
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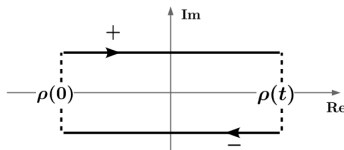
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## Example: dilatons with Damour-Polyakov mechanism

- Motivated by string theory and scalar-tensor theories
- Screening mechanism to circumvent fifth force constraints
- $\mathcal{L} = -(\partial X)^2 - V_0 e^{-\lambda X/m_{\text{Pl}}} + \frac{A_2}{2m_{\text{Pl}}^2} T^\nu{}_\nu X^2$  with  $X = X_0 + \chi$
- VEV:  $X_0 = \frac{m_{\text{Pl}}}{\lambda} W\left(-\frac{\lambda^2 V_0}{A_2 T^\nu{}_\nu}\right)$  with  $W(x) = \sum_{n=1}^{\infty} \frac{(-n)^{n-1}}{n!} x^n$
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# Eur.Phys.J.C 83 (2023) 8, 767

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- Open quantum dynamical effects could affect interference pattern in atom interferometer
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  - Real scalar  $\phi$  since dilaton only couples to mass density
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  - Restrict to single-particle subspace since atoms are cold

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# Eur.Phys.J.C 83 (2023) 8, 767

- Compute reduced density matrix elements  $\rho_\phi(\mathbf{p}, \mathbf{p}'; t)$

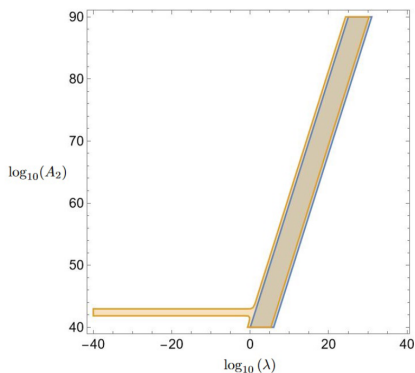
- Extract frequency shift  $\Delta u := M \frac{A_2}{4m_{\text{pl}}^2} \left[ X_0^2 + \Delta_{zz}^{\text{F}(T \neq 0)} \right] v^2$

- $M$ : cold atom's mass (rubidium-87)
- $v := \frac{\|\mathbf{p}\| - \|\mathbf{p}'\|}{M}$  non-relativistic speed difference of two atom states
- Thermal part of the dilaton Feynman propagator:

$$\Delta_{xx}^{\text{F}(T \neq 0)} = \frac{T^2}{2\pi^2} \int_{m/T}^{\infty} d\xi \frac{\sqrt{\xi^2 - (\frac{m}{T})^2}}{e^\xi - 1}$$

with dilaton mass  $m$  and temperature  $T$

## Eur.Phys.J.C 83 (2023) 8, 767



- $V_0 = 1 \text{ (MeV)}^4$
- $u(P, T_2; P, T_1)$  (blue) and  $u(P, T_3; P, T_1)$  (orange);  $P = 9.6 \times 10^{-10}$  mbar (hydrogen);  
 $T_1 = 0.5 \times 10^{-3}$  K,  $T_2 = 100$  K, and  $T_3 = 500$  K

# Conclusions

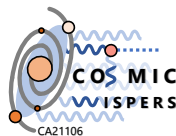
- New first principle-based formalism for computing reduced density matrices
- Open quantum dynamical effects could lead to new constraints on new physics



# Funding



This research was funded in whole or in part by the Austrian Science Fund (FWF) 10.55776/PAT8564023.



This publication is based upon work from COST Action COSMIC WISPers CA21106, supported by COST (European Cooperation in Science and Technology.)

# Eur.Phys.J.C 83 (2023) 8, 767

- Compare two measurements:

$$|\Delta u(P_a, T_b) - \Delta u(P_c, T_d)| =: u(P_a, T_b; P_c, T_d) \geq \Delta u_{\min}$$

- $P$ : Gas pressure in vacuum chamber (affects  $\varphi_0$ )
- Choose rubidium-87 atoms with  $M = 87m_u$
- $v = 50$  mm/s
- $\Delta u_{\min} \approx 10^{-8}$  Hz