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University of Science and Technology of China

Measurement of the electric dipole moment (EDM) of ^{171}Yb atoms in an optical dipole trap (ODT)

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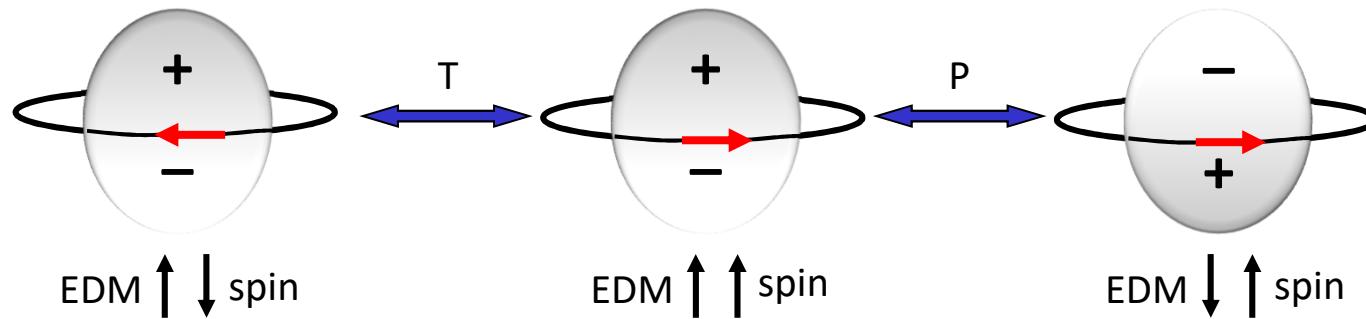
創寰宇學府
育天下英才
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一九八八年五月

MITP workshop, Nov. 2024

Permanent EDM violates P -, T - symmetry

Induced EDM: $Energy = -\frac{1}{2} \alpha E^2$

Permanent EDM: $Energy = -d \cdot E$



- CP problem in strong interaction: θ term

$$\mathcal{L}_{\text{mass}} \rightarrow -m(\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L) + \frac{\theta g^2}{32\pi^2} F_a^{\mu\nu} \tilde{F}_{a\mu\nu}$$

- SUSY: New elementary particles \rightarrow new \cancel{CP} phase
- Origins of matter – antimatter asymmetry: New physical \cancel{CP} mechanism

Searching for EDM in three different categories

PSI, US, Germany,
Japan, KAIST

UWash, PTB,
Argonne, USTC,
CENTREX, UChicago

ACME, JILA
Imperial college
ECNU

Neutron, proton

Nuclear Schiff moment
Diamagnetic Hg, Xe, Yb,
Ra, Rn, TlF, AgFr, Yb(OH)

Electron
Paramagnetic
YbF, ThO, HfF⁺

Quark EDM

Quark chromo-EDM
4 fermions, 3 gluons

Electron-EDM,
electron-quark force

New physics
beyond SM:
SUSY *et al.*

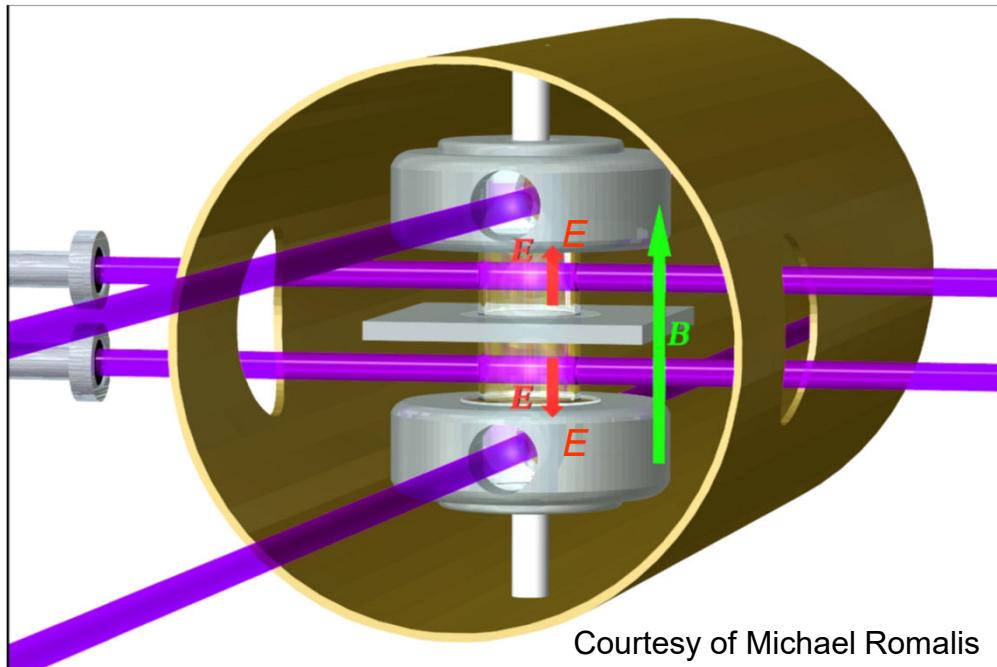
Engel, Ramsey-Musolf,
van Kolck, Prog. Part.
Nucl. Phys. 71, 21 (2013)

system	Upper limit (e-cm)	method	Value in Standard model (e-cm)
Electron	4×10^{-30}	HfF ⁺ trap, ThO beam	10^{-35}
Neutron	2×10^{-26}	Neutrons – bottle	10^{-32}
¹⁹⁹ Hg	7×10^{-30}	Atoms – vapor cell	10^{-34}
¹⁷¹ Yb	This work	Atoms – trap	10^{-34}

The Seattle EDM Measurement

^{199}Hg

stable, high Z, groundstate $^1\text{S}_0$, $I = \frac{1}{2}$, high vapor pressure



$$f_+ = \frac{2\mu B + 2dE}{h} \approx 15 \text{ Hz}$$

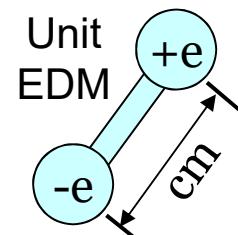
$$f_- = \frac{2\mu B - 2dE}{h} \approx 15 \text{ Hz}$$

$$|f_+ - f_-| < 25 \text{ pHz}$$

The best limit on atomic EDM

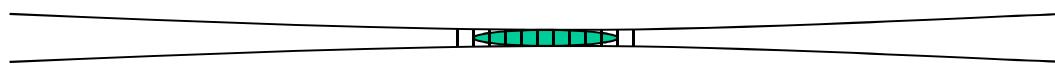
$$\text{EDM } (^{199}\text{Hg}) < 7 \times 10^{-30} \text{ e-cm}$$

Graner et al., Phys Rev Lett (2016)



Measure EDM in an Optical Dipole Trap (ODT)

M.V. Romalis and E.N. Fortson, Phys. Rev. A 59, 4547 (1999)



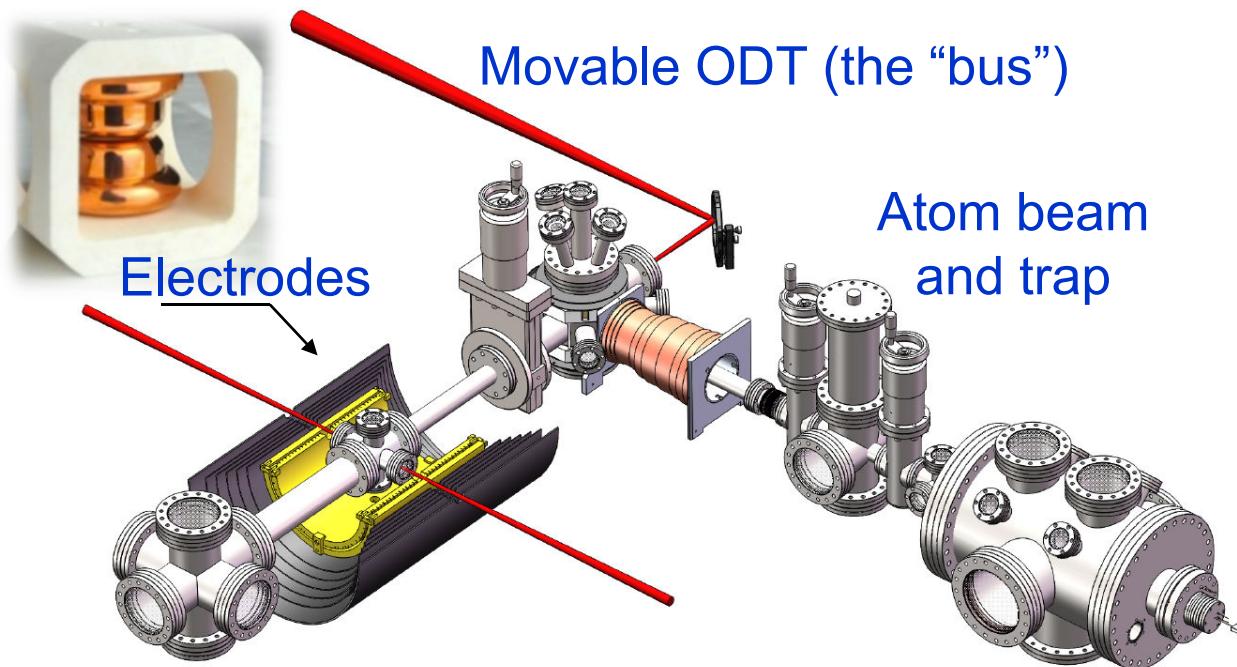
$$H = -\tilde{d}E = -\frac{1}{4}\alpha E_0^2$$

- Fiber laser: $\lambda = 1036$ nm, Power = 10 Watts
- Focused to $50 \mu\text{m}$ \rightarrow trap depth $60 \mu\text{K}$

EDM in an optical dipole trap

- $v \times E$, Berry's phase effects suppressed
- Cold scattering suppressed between cold Fermionic atoms
- Rayleigh scat. rate $\sim 10^{-1} \text{ s}^{-1}$; Raman scat. rate $\sim 10^{-12} \text{ s}^{-1}$
- Vector light shift $\sim \mu\text{Hz}$
- Parity mixing induced shift under control
- Conclusion: possible to reach 10^{-30} e cm for ^{199}Hg

^{171}Yb EDM Apparatus: Trapping + Science



$$\varpi_+ = 2\mu B + 2dE$$

E : electric field

$$\varpi_- = 2\mu B - 2dE$$

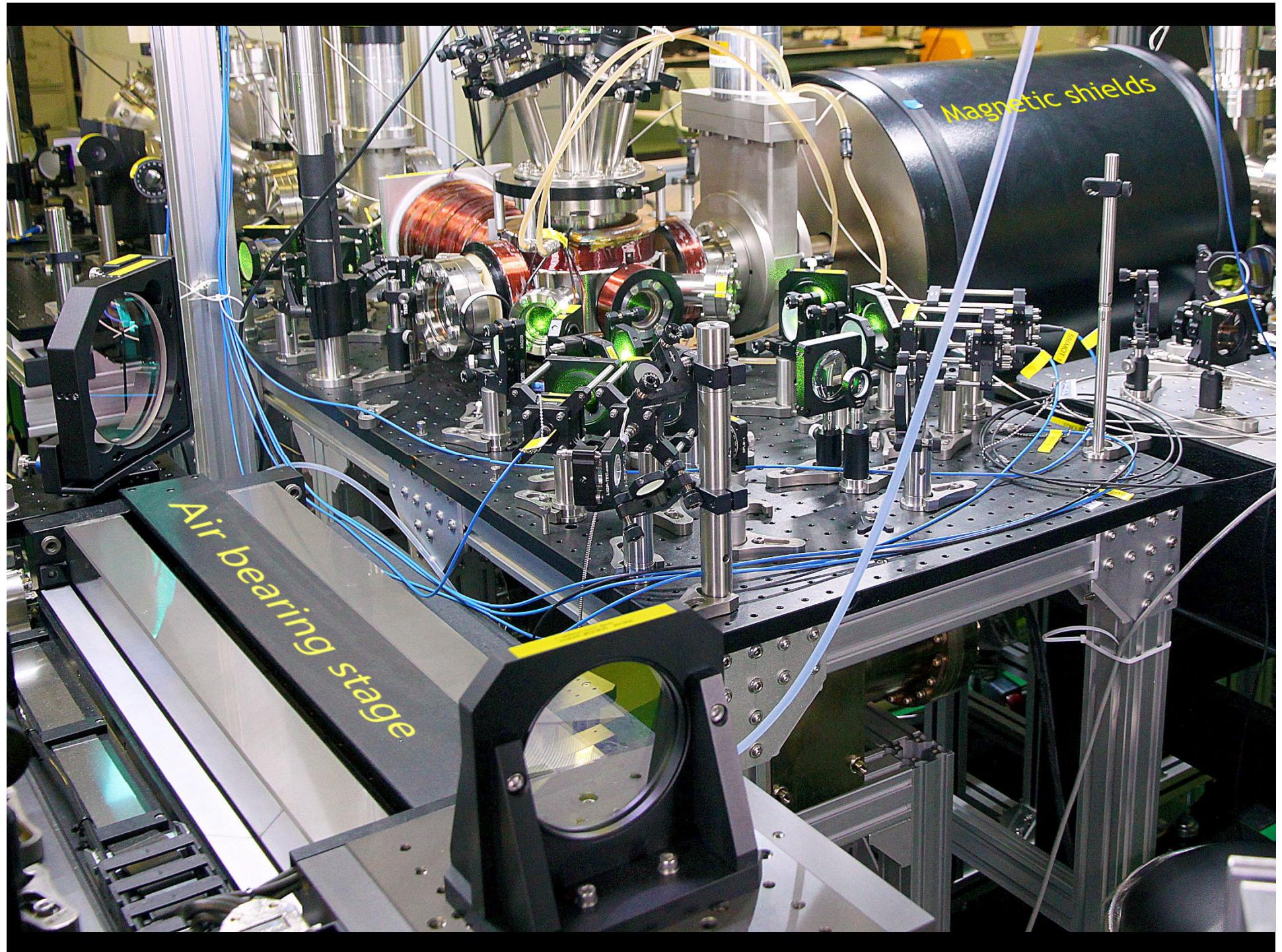
τ : precession time

$$\delta d = \frac{\hbar}{2E\sqrt{\tau N \varepsilon T}}$$

N : atom number

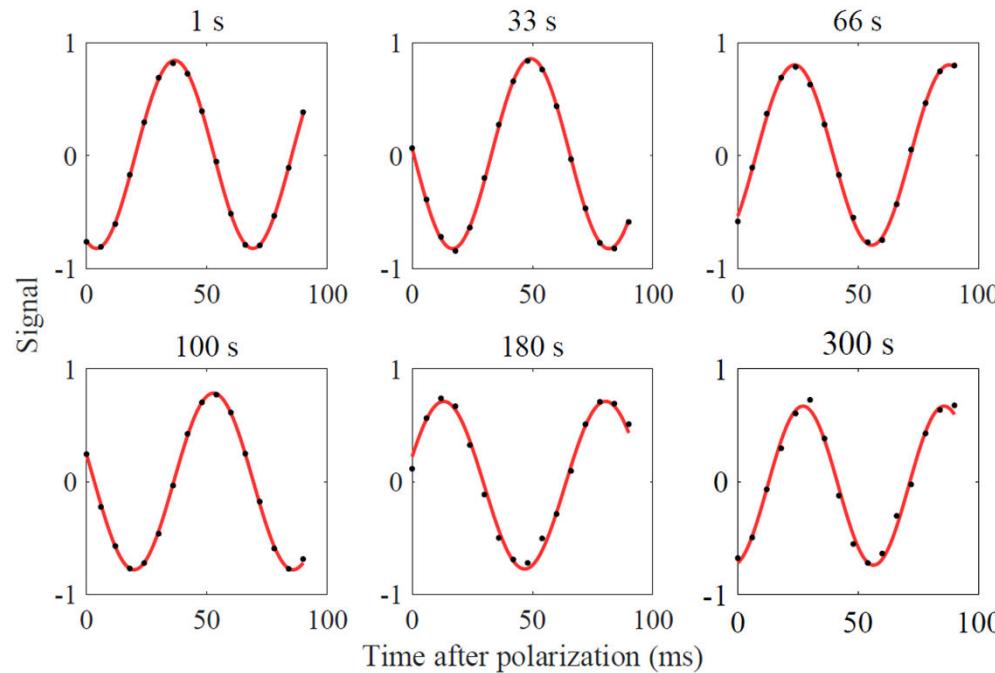
ε : spin – state detection efficiency

T : time of average



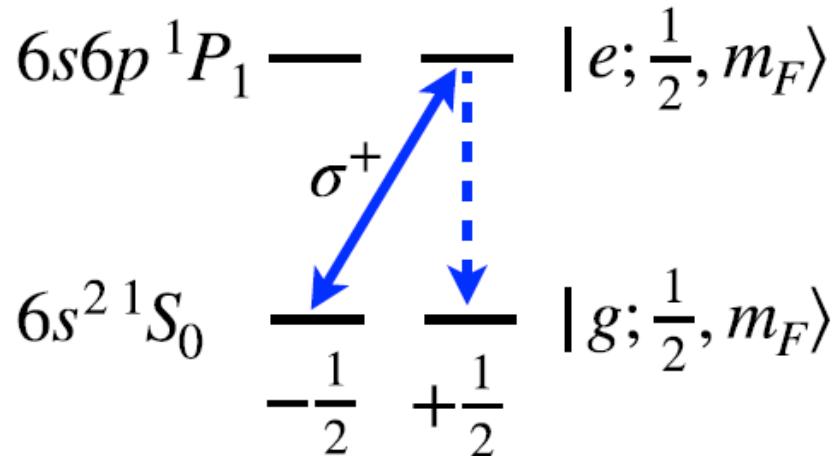
Long spin coherence time (T_2)

$$S_z = \frac{N_+ - N_-}{N_+ + N_-} = C \exp\left(-\frac{t}{T_2}\right) \cos(2\pi ft + \phi_0) + O$$



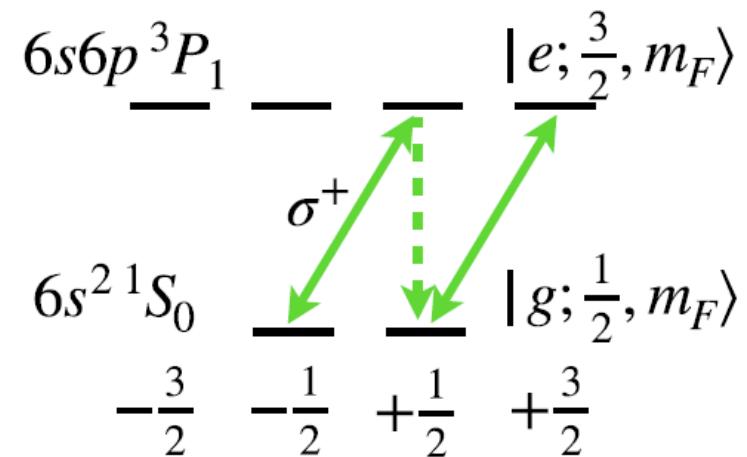
- Vacuum-limited trap lifetime: 75 s
- EDM measurement precession time: 96 s
- Observed precession to: 300 s
- T_2 : $(9 \pm 4) \times 10^3$ s (2.5 ± 1.1 hr)

Spin-state detection – conventional method



Bright state:
~ 3 photons
before state
“demolished”

Dark state:
no photons,
state preserved

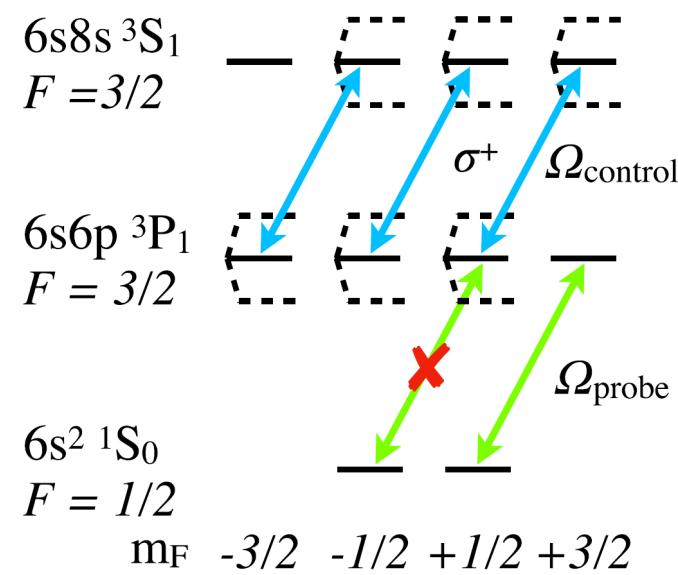
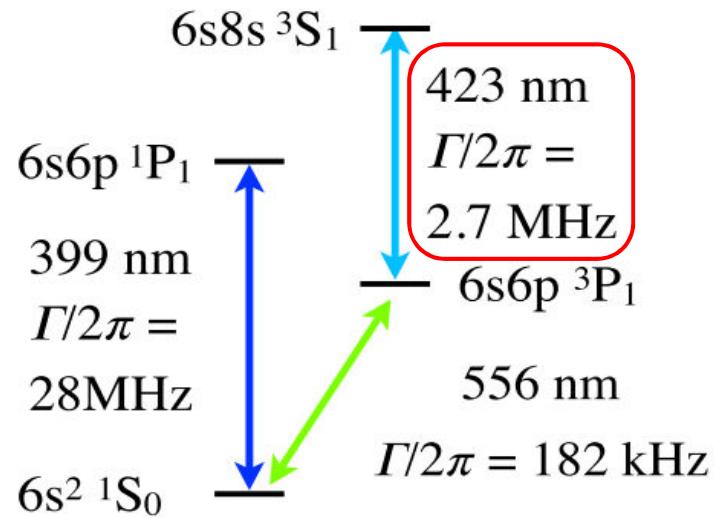


Bright state:
state
“demolished”

Brighter state:
cycling,
state preserved

Spin-state detection by a QND method

Y. Yang *et al.*, Phys Rev Applied 19, 054015 (2023)



- $\Omega_{control} \sim 2\pi \times 40 \text{ MHz}$
- $\Omega_{probe} \sim 2\pi \times 70 \text{ kHz}$
- Spin flip suppressed by: $\Omega_c^2/(\Gamma_e \Gamma_c) \sim 10^3$
- Need ODT to be at magic wavelength

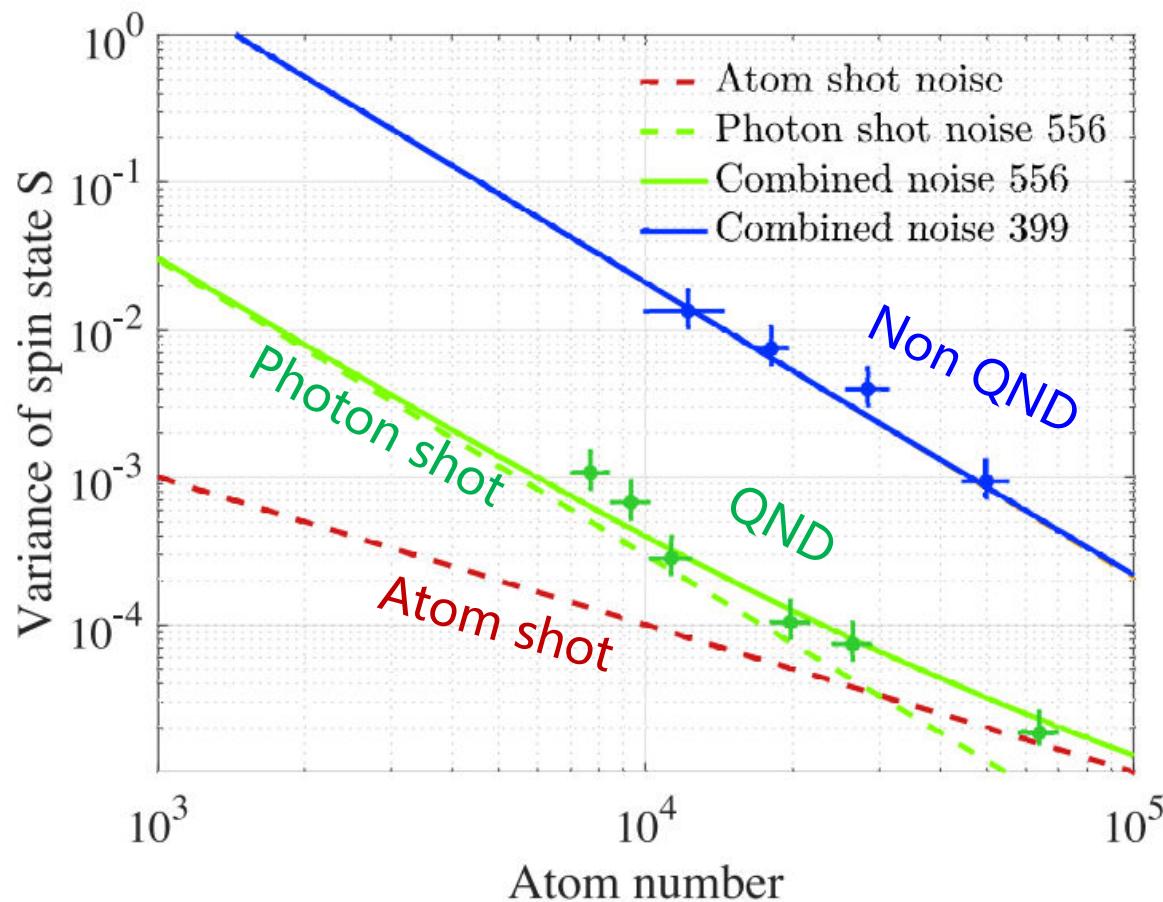
-- T. Zheng, M.S. Safronova *et al.*, PRA (2020)

Dark state:
no photons,
state preserved

Bright state:
cycling,
state preserved

Spin-state detection by a QND method

Yang et al., Phys Rev Applied 19, 054015 (2023)



- Photons scattered
non-QND: $\bar{n} \sim 2.5$
QND: $\bar{n} \sim 23$
- ~19 dB reduction of
variance

Systematic errors

- effects correlated with E-field flipping

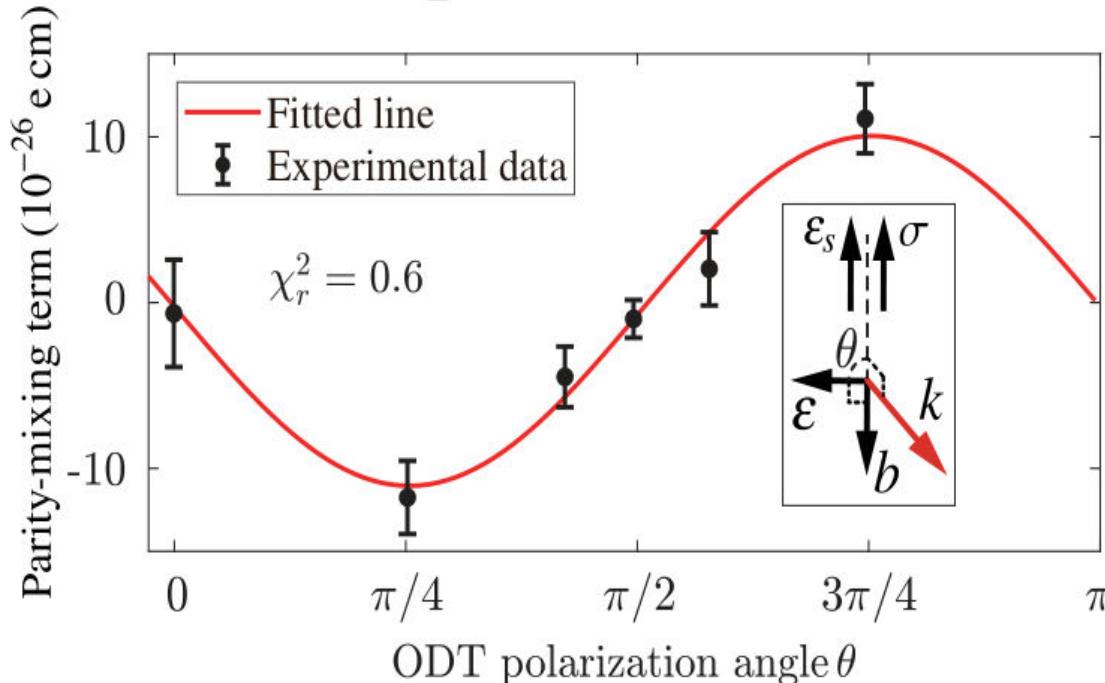
Source	Error
B-field correlations	2.5×10^{-28}
Parity mixing	2.6×10^{-28}
Leakage current	0.9×10^{-28}
ODT power effect	0.1×10^{-28}
E-squared effect	0.1×10^{-28}
Total	3.7×10^{-28}

Parity mixing

$$\Delta \nu_{F=I\pm 1/2} = \mp [\nu_{MD}^1 (\hat{b} \cdot \hat{\sigma}) (\hat{\epsilon} \cdot \hat{\epsilon}_s) + \nu_{MD}^2 (\hat{b} \cdot \hat{\epsilon}_s) (\hat{\epsilon} \cdot \hat{\sigma})] m$$

M.V. Romalis and E.N. Fortson, Phys. Rev. A 59, 4547 (1999)

$$\Delta \nu \approx -\frac{\nu_1 + \nu_2}{2} \sin(2\theta), \quad (0 \leq \theta < \pi)$$



T-even, P-even
yet EDM-like

ϵ : ac E-field of ODT

b : ac B-field of ODT

ϵ_s : Static E-field

σ : Quantization axis
Static B-field

measurement : $1.2(2) \times 10^{-25} e \text{ cm}$ ($\epsilon \cdot \epsilon_s = \epsilon \cdot \sigma = \sqrt{2}/2$)

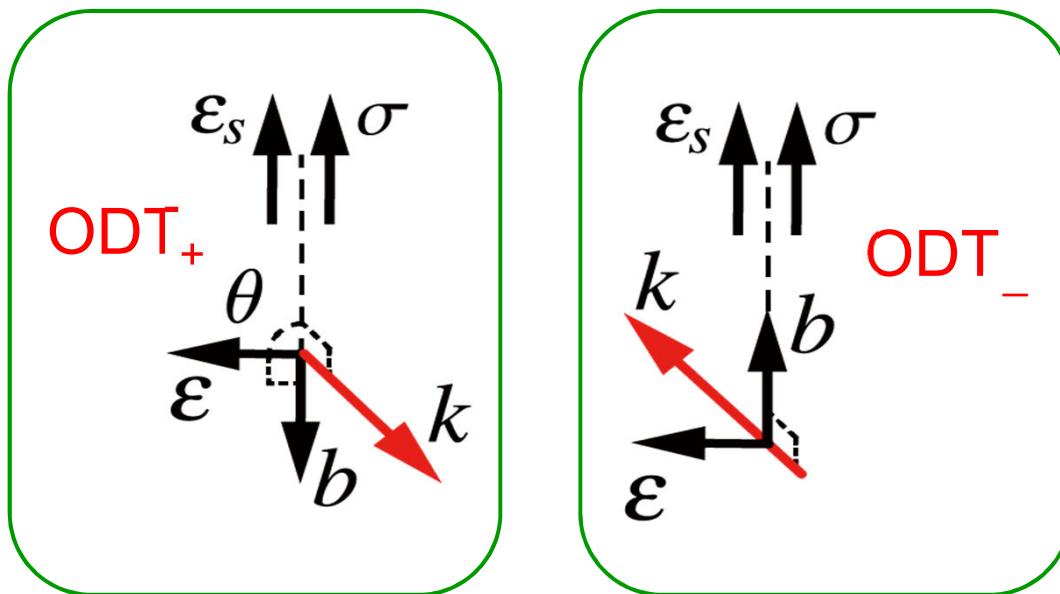
estimate : $2 \times 10^{-25} e \text{ cm}$

Calculation needed

Parity mixing

$$\Delta \nu_{F=I\pm 1/2} = \mp [\nu_{MD}^1 (\hat{b} \cdot \hat{\sigma}) (\hat{\epsilon} \cdot \hat{\epsilon}_s) + \nu_{MD}^2 (\hat{b} \cdot \hat{\epsilon}_s) (\hat{\epsilon} \cdot \hat{\sigma})] m$$

M.V. Romalis and E.N. Fortson, Phys. Rev. A 59, 4547 (1999)



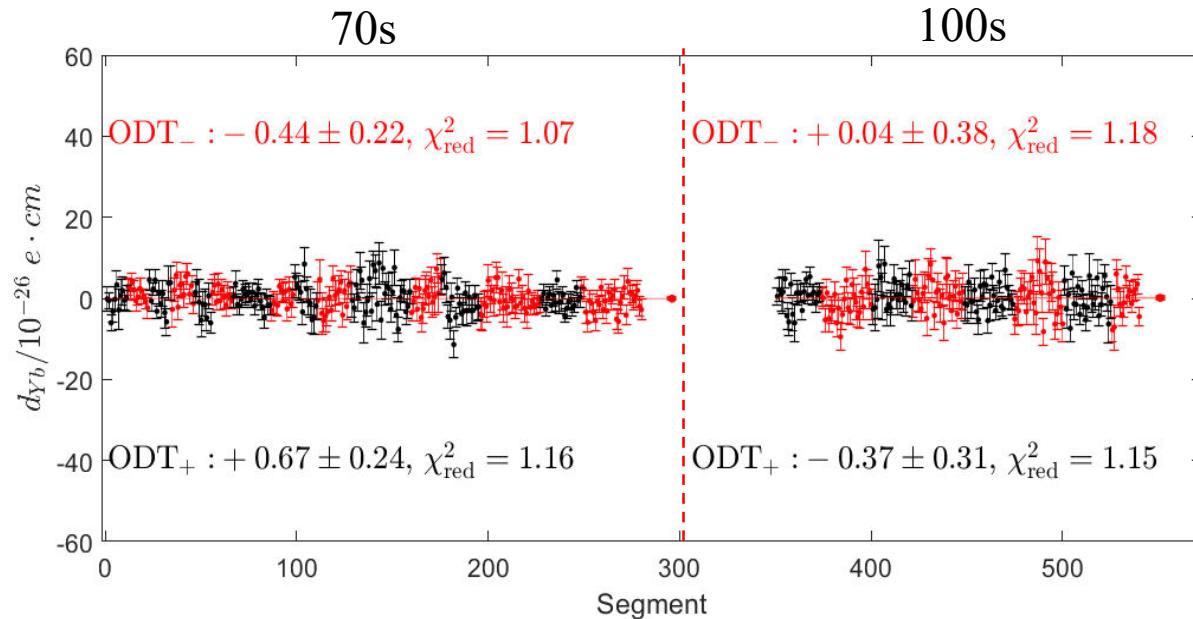
Control:

$$\theta = \pi/2 \pm 80 \text{ mrad}$$

$$(\theta_+ - \theta_-) = 0 \pm 3 \text{ mrad}$$

- Average between ODT+ and ODT-
- Residual parity mixing effect: 2.6×10^{-28} e cm
- Future: better k+ and k- balance in an optical cavity

^{171}Yb EDM measurement results



* Average ODT+ and ODT- to cancel parity-mixing systematic

$$d(^{171}\text{Yb}) = (0.3 \pm 1.4_{\text{stat}} \pm 0.4_{\text{syst}}) \times 10^{-27} \text{ e cm}$$

$$\text{Upper limit (95\%)} : |d(^{171}\text{Yb})| < 2.8 \times 10^{-27} \text{ e cm}$$

$$\text{Upper limit (95\%)} : |d(^{171}\text{Yb})| < 1.5 \times 10^{-26} \text{ e cm}$$

T.A. Zheng *et al.*, PRL **129**, 083001 (2022)

$$\delta d = \frac{\hbar}{2E\tau\sqrt{n}} \sqrt{\frac{1}{N_a \epsilon_d} + \sigma_{\phi(\delta B)}^2}$$

$$E: 192 \text{ kV/cm}$$

$$\tau: 60 - 90 \text{ s}$$

$$N: 8 \times 104$$

$$\epsilon_d: \sim 50\%$$

$$T: 840 \text{ h (40 days)}$$

$$\sigma_{\phi(\delta B)}^2 \approx 3 \times \frac{1}{N_a \epsilon_d}$$

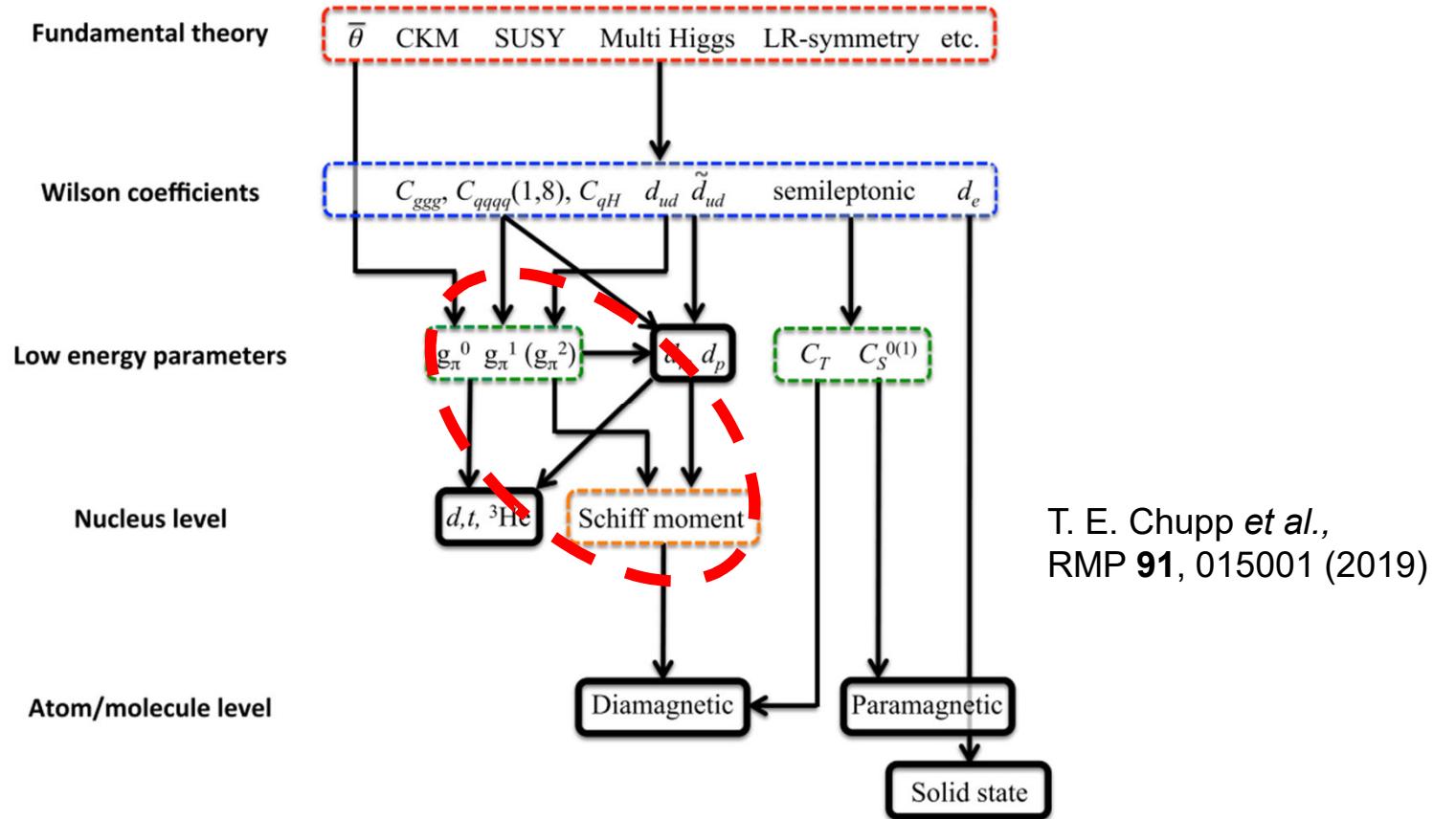
* B-field noise due to
Seebeck effect

Upper limits on Schiff moments

Diamagnetic system	^{199}Hg	$^{205}\text{Tl}^{19}\text{F}$	^{129}Xe	^{171}Yb	^{225}Ra
EDM limit (10^{-27} e cm)	7.4×10^{-3} UWash16	6.5×10^4 Yale91	1.4 PTB19	2.8 USTC23	1.4×10^4 Argonne16
Schiff moment limit (10^{-10} e fm^3)	2.6×10^{-3}	1.5	5.2	1.6	1700
Schiff moment calc (10^{-8} e fm^3)	$-1.4 \eta_{\text{np}}$ Flambaum86	$1.2 \eta_{\text{pp}} - 1.4 \eta_{\text{pn}}$ Flambaum86	$1.75 \eta_{\text{np}}$ Flambaum86	$-1.4 \eta_{\text{np}}$ Dzuba07	$300 \eta_n, 1100 \eta$ Flambaum03, Auerbach96

- ❖ First ^{171}Yb EDM result constrains BSM physics on the same order as ^{129}Xe , ^{225}Ra , $^{205}\text{Tl}^{19}\text{F}$, all lagging behind ^{199}Hg .
- ❖ Global analysis needs different systems with complementary sensitivities to BSM parameters, rather than ^{199}Hg alone.

Theories from CP violation to EDM



Calculation of Schiff moments

- ${}^{199}\text{Hg}$: J. H. de Jesus and J. Engel, PRC **72**, 045503 (2005)
- ${}^{225}\text{Ra}$: J. Dobaczewski and J. Engel, PRL **94**, 232502 (2005)
- ${}^{171}\text{Yb}$: **nuclear calculation needed !**

Outlook #1

Upgrade: ^{171}Yb EDM precision improves into E-28 e-cm

- ❖ Larger E-field
- ❖ Longer trap lifetime
- ❖ Less B-field noise
- ❖ Optical cavity

E-field

- 2024 setup: 300 kV/cm, copper
- Other works: 500 kV/cm, niobium - Ready *et al.* NIM A (2021)

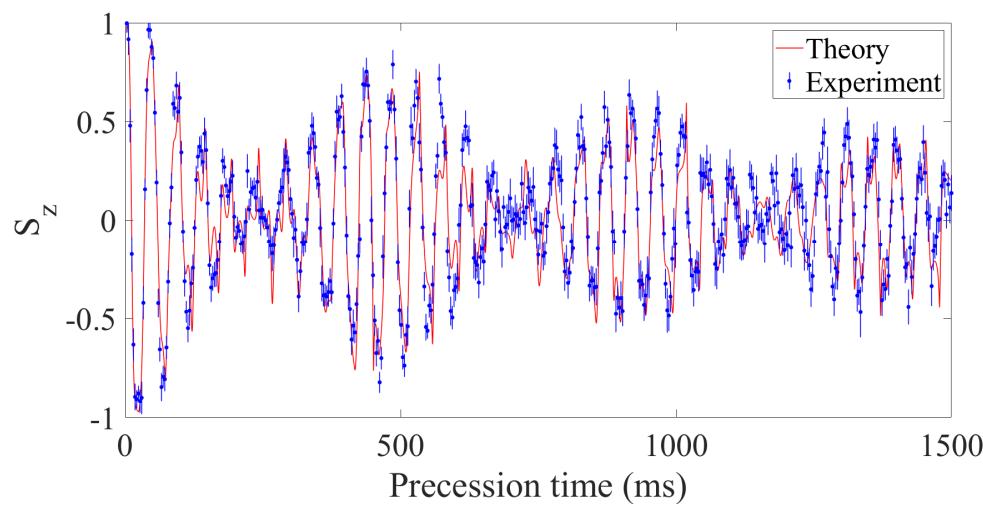
B-field

- This work: ~ 3 pT over 100 s, Seebeck effect due to ODT heating
- Environmental noise: ~ 1 pT, use magnetometers
- Johnson noise: ~ 0.6 pT, use co-magnetometer (^{173}Yb)

^{173}Yb ($I = 5/2$) co-magnetometer

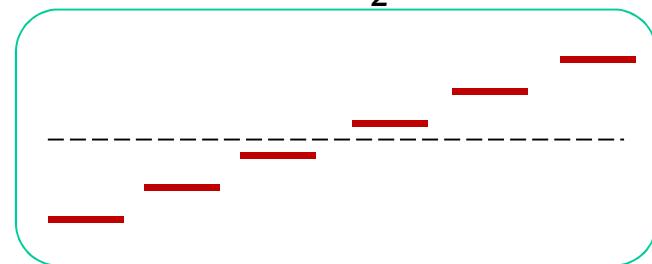
Prepare initial state: $|+5/2\rangle_x = \frac{1}{\sqrt{32}} (1 \ \sqrt{5} \ \sqrt{10} \ \sqrt{10} \ \sqrt{5} \ 1)$

Problem: decoherence due to spatially varying ODT light shift

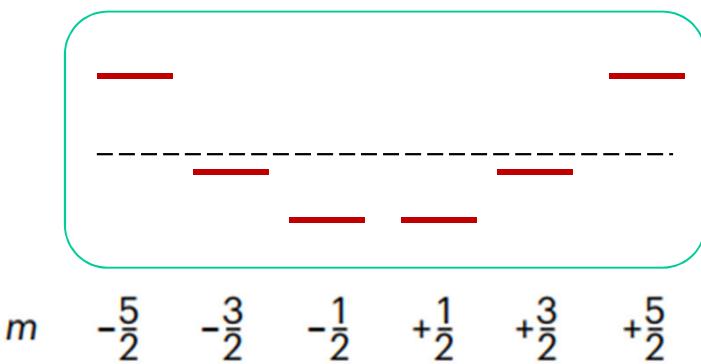


$$\text{ODT tensor shift: } H_t = \Omega(\mathbf{r}) F_z^2$$

Zeeman shift F_z



Stark shift $(F_z)^2$



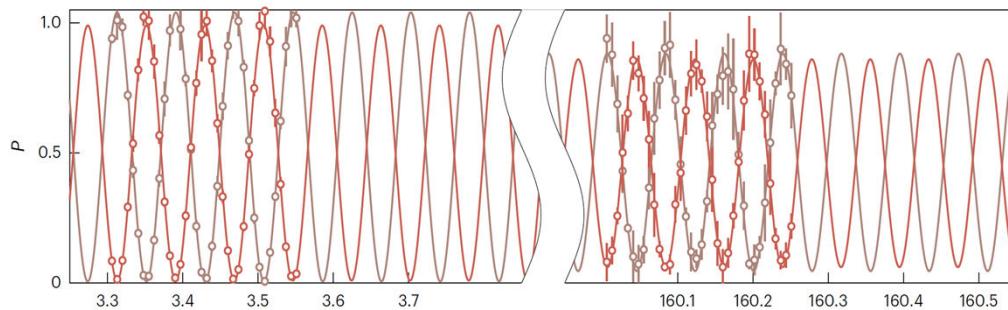
^{173}Yb ($I = 5/2$) co-magnetometer

$$|+5/2\rangle \xrightarrow{\pi/2} |\text{Cat}\rangle = \frac{1}{\sqrt{2}} (|+5/2\rangle + |-5/2\rangle)$$

High – spin $\pi/2$ pulse

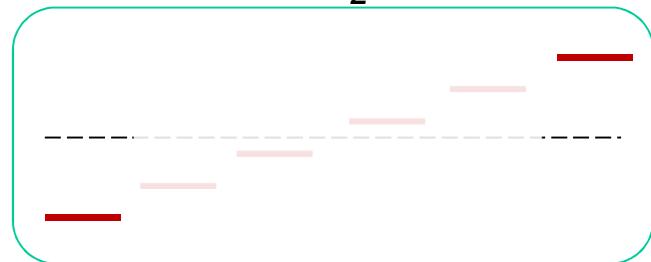
$$H = \Omega(F_x - F_x^2)$$

$$t = \frac{\pi}{2\Omega} \quad \Omega = 2\pi \cdot 0.5 \text{ kHz}$$

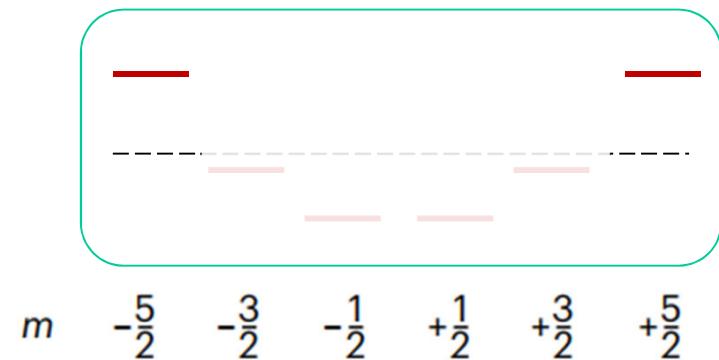


Coherence time ~ 1400 s

Zeeman shift I_z



Stark shift I_z^2



Outlook #2

^{171}Yb

2-electron system

Nuclear spin $I = \frac{1}{2}$

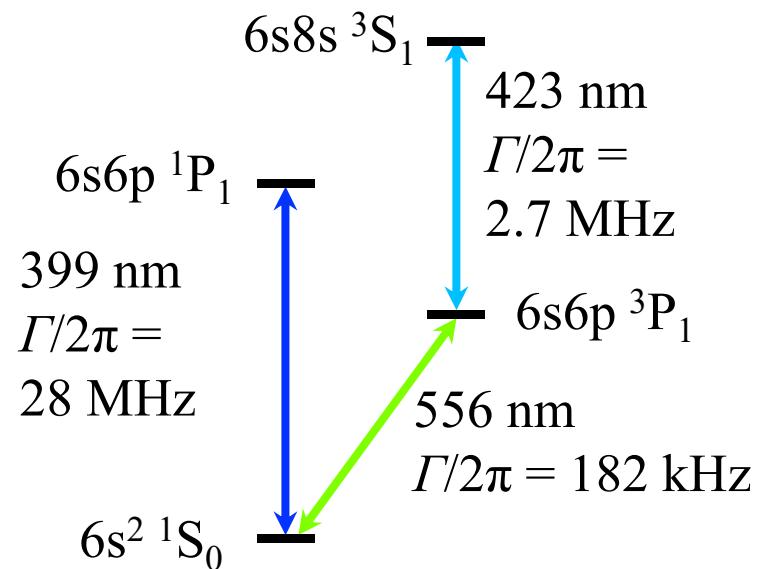
Stable and abundant

^{225}Ra :

Group II

Nuclear spin $I = \frac{1}{2}$

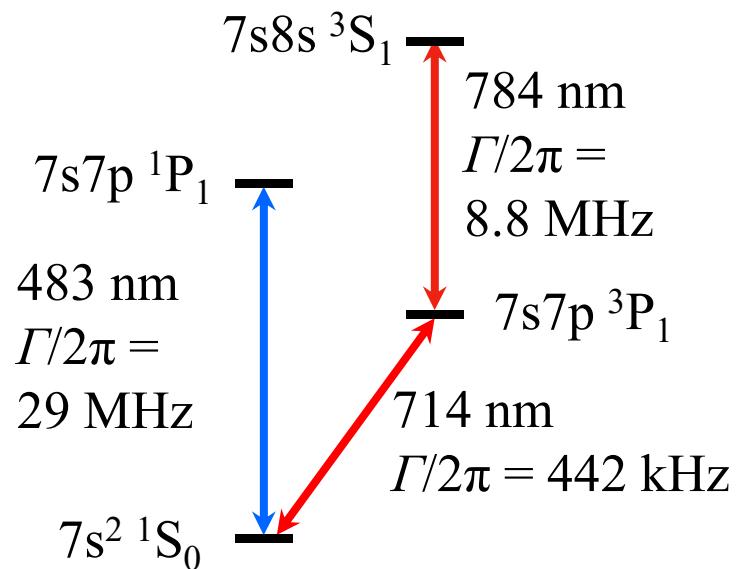
Radioactive $t_{1/2} = 15$ d



^{171}Yb
E-28 e-cm



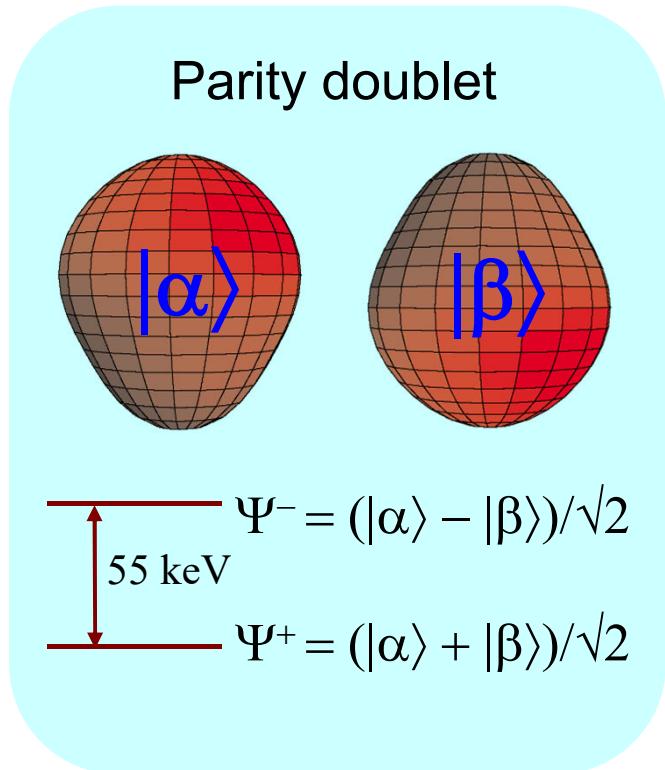
^{225}Ra
E-28 e-cm



^{199}Hg
E-31 e-cm

EDM of ^{225}Ra enhanced and more reliably calculated

- Closely spaced parity doublet – Haxton & Henley, PRL (1983)
- Large Schiff moment due to octupole deformation – Auerbach, Flambaum & Spevak, PRL (1996)
- Relativistic atomic structure ($^{225}\text{Ra} / ^{199}\text{Hg} \sim 3$) – Dzuba, Flambaum, Ginges, Kozlov, PRA (2002)



$$\text{Schiff_moment} = \sum_{i \neq 0} \frac{\langle \psi_0 | \hat{S}_z | \psi_i \rangle \langle \psi_i | \hat{H}_{PT} | \psi_0 \rangle}{E_0 - E_i} + c.c.$$

Enhancement Factor: EDM (^{225}Ra) / EDM (^{199}Hg)

	Isoscalar	Isovector
Skyrme SIII	300	4000
Skyrme SkM*	300	2000
Skyrme SLy4	700	8000

Schiff moment of ^{225}Ra , Dobaczewski, Engel, PRL (2005)

Schiff moment of ^{199}Hg , Dobaczewski, Engel et al., PRC (2010)

“[Nuclear structure] calculations in Ra are almost certainly more reliable than those in Hg.”

– Engel, Ramsey-Musolf, van Kolck, Prog. Part. Nucl. Phys. (2013)

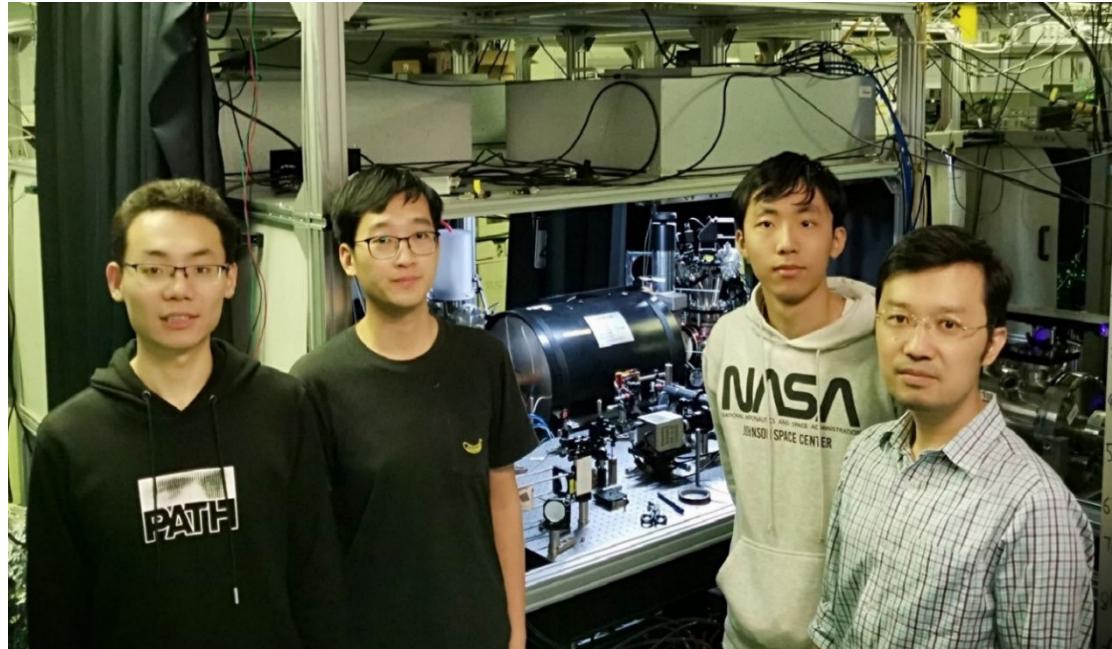
Constraining parameters in a global EDM analysis.

– Chupp, Ramsey-Musolf, PRC 91, 035502 (2015)

Collaboration and Support

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Z.-T. Lu
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