

Schwarzschild geodesics from Scattering Amplitudes to all orders in G

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LUTH



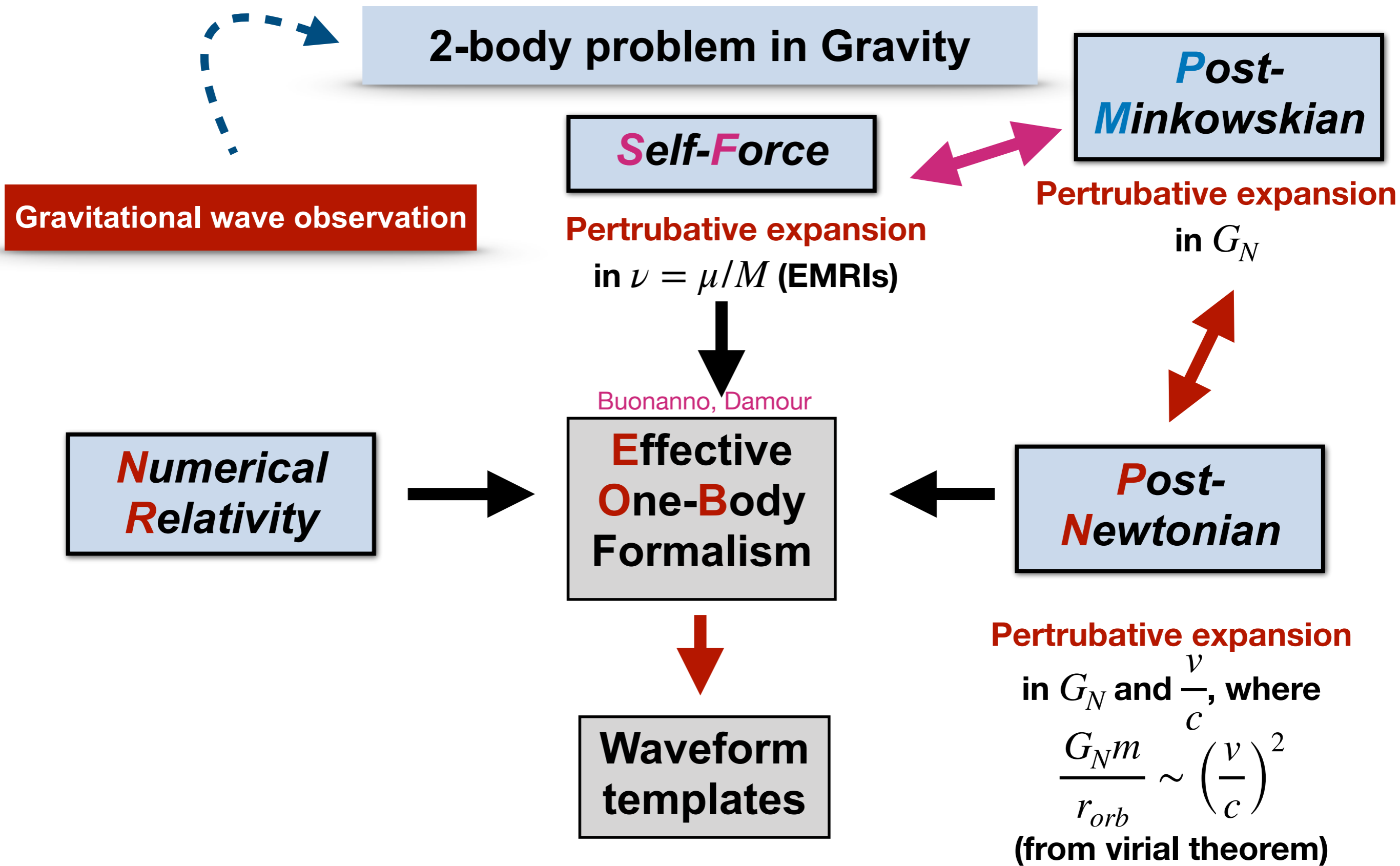
PSL 

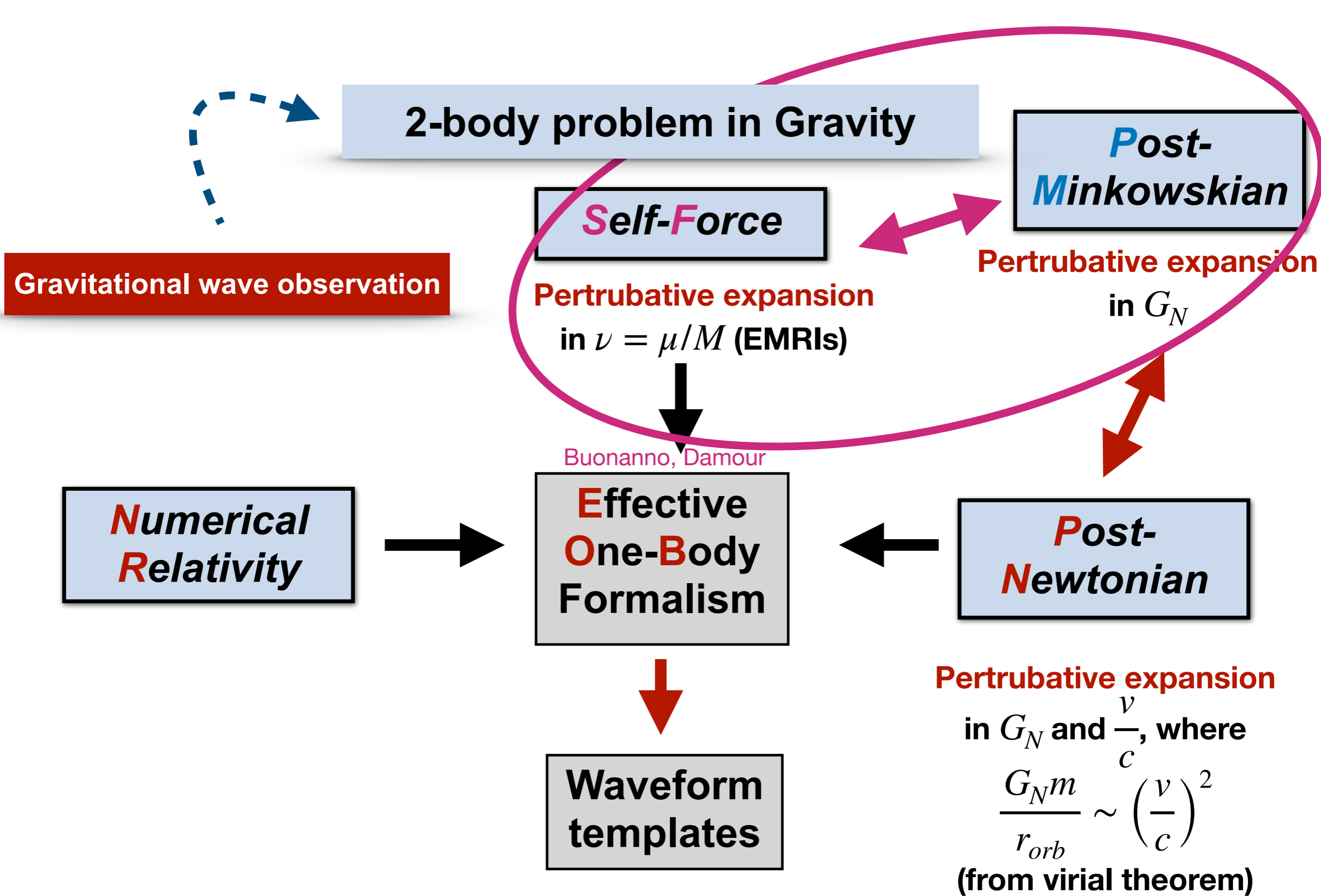
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YOUNGST@RS-Loop-the-Loop MITP, online 13/11/2024

Based on works with: P.Vanhove **[2407.09448],[2405.14421]**





2-body problem in Gravity

Post-Minkowskian

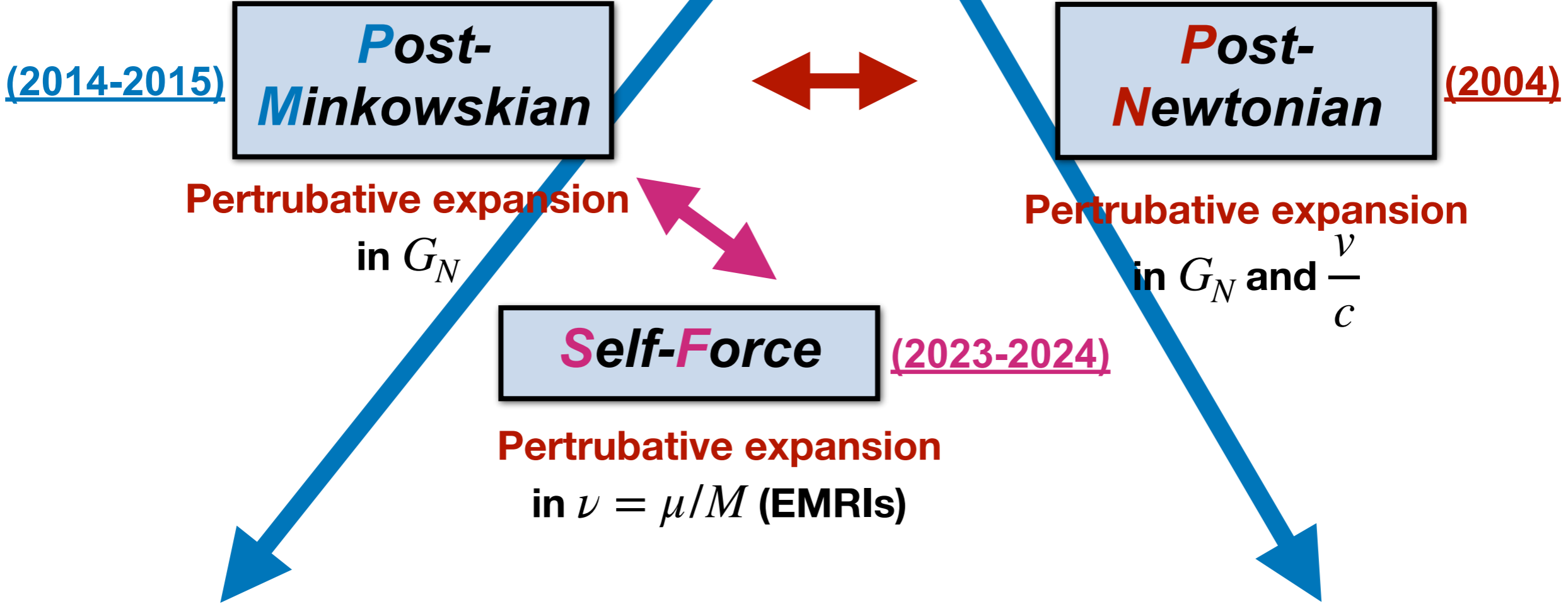
vs

Self-Force

	1PM	2PM	3PM	4PM	5PM	...
0SF	$\left[(G m_1) + (G m_1)^2 + (G m_1)^3 + (G m_1)^4 + (G m_1)^5 + \dots \right] \times m_2$					
1SF			$\left[(G m_1)^2 + (G m_1)^3 + (G m_1)^4 + \dots \right] \times G m_2^2$			
2SF					$\left[(G m_1)^3 + \dots \right] \times G^2 m_2^3$	

2-body problem in Gravity

Analytical methods/Perturbation theory



[Rothstein, Goldberger, Porto, Bern, Kosower, O'Connell, Vanhove, Damgard, Plefka et al.]

[Damour, Blanchet, Buonanno et al.]
[Poisson, Barack, Pound et al.]

2-body problem in Gravity

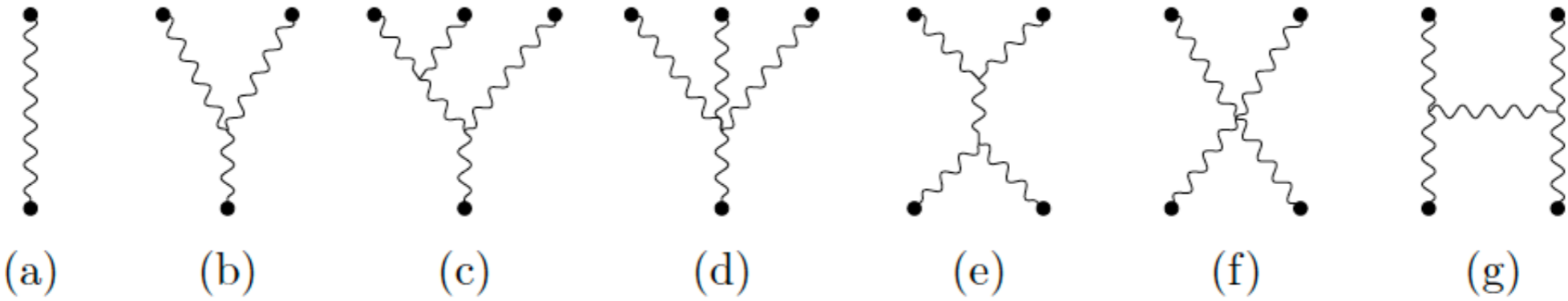
EFT+Scattering Amplitudes (in a nutshell)

$$e^{iS_{\text{eff}}[x_a]} = \int \mathcal{D}h_{\mu\nu} e^{iS_{\text{EH}}[h] + iS_{\text{GF}}[h] + iS_{\text{PP}}[x_a, h]}$$

Point particle approximation
+weak metric perturbations

Feynman rules

integrate-out graviton
(classical=**NO** grav. loops)



Schwarzschild from Amplitudes to all orders in G

[2407.09448],[2405.14421] S.M., P. Vanhove

[2308.14832,2406.14770] Cheung, Parra-Martinez et al.

[2308.15304] Kosmopoulos, Solon

Towards **SF-EFT**, it is necessary to have a
self-consistent framework to all orders in G

Simplest case: *Schwarzschild metric*

Schwarzschild from Amplitudes to all orders in G

[2407.09448],[2405.14421] S.M., P. Vanhove

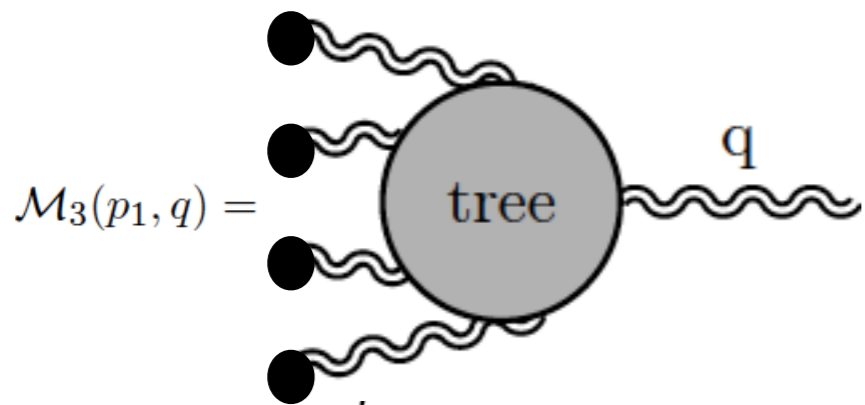
GREFT



General Relativity

Equivalence up to:

1. Choice of DOFs \longrightarrow $g_{\mu\nu} = \eta_{\mu\nu} + \sum_{n=1}^{\infty} h_{\mu\nu}^{(n)}$
2. Choice of Gauge \longrightarrow harmonic gauge



$$i\mathcal{M}_3^{(l)}(p_1, q) = -\frac{i\sqrt{32\pi G_N}}{2} \langle T^{(l)\mu\nu}(q^2) \rangle \epsilon_{\mu\nu} \quad \text{Duff (1974)}$$

$$h_{\mu\nu}^{(l+1)}(\vec{x}) = -16\pi G_N \int \frac{d^d \vec{q}}{(2\pi)^d} e^{i\vec{q}\cdot\vec{x}} \frac{1}{\vec{q}^2} \left(\langle T_{\mu\nu}^{(l)} \rangle^{\text{class.}}(q^2) - \frac{1}{d-1} \eta_{\mu\nu} \langle T^{(l)} \rangle^{\text{class.}}(q^2) \right)$$

Expectation: n -loop diagrams generate G_N^{n+1} terms of the metric

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

GREFT



General Relativity

Equivalence up to:

1. Choice of DOFs $\longrightarrow g_{\mu\nu} = \eta_{\mu\nu} + \sum_{n=1}^{\infty} h_{\mu\nu}^{(n)}$
2. Choice of Gauge \longrightarrow harmonic gauge

Too complicated!!

Main problems from previous attempts:

[2010.08882]
S.M, Vanhove

1. Infinite tower of non-minimal couplings (due to intermediate UV-divs)
2. No algorithm for higher loops (3-loops was already complicated)

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

GREFT



General Relativity

Equivalence up to:

1.

Choice of DOFs



$$g_{\mu\nu} = \eta_{\mu\nu} + \sum_{n=1}^{\infty} h_{\mu\nu}^{(n)}$$

2.

Choice of Gauge



harmonic gauge

Take a step back

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Cubic formulation of GREFT

[1705.00626] Cheung, Remmen

Choice of DOFs: 1) Gothic metric: $g^{ab} = \sqrt{-g}g^{ab} = \eta^{ab} - \sqrt{32\pi G_N}h^{ab}$

[2403.13216] Lee, Damgaard

See for a similar
derivation with
the perturbiner approach

(GR people do it, we should pay more attention)

2) Extra Auxiliary field : $A_{bc}^a = \Gamma_{bc}^a - \frac{1}{2}\delta_{(b}^a\Gamma_{c)d}^d$.

(unorthodox but necessary to constrain to 3pt vertices)

Choice of Gauge: harmonic gauge*(non unique)

+ couple to worldline $\mathcal{L}_{p.p.} = -\frac{m}{2} \int d\tau (e^{-1} g^{\mu\nu} v_\mu v_\nu + e) = -\frac{m}{2} \int d\tau \left(\frac{g^{\mu\nu} v_\mu v_\nu}{(\sqrt{-g})^{\frac{D-2}{2}}} + 1 \right)$

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Cubic formulation of GREFT

Complete ansatz:

$$\sqrt{32\pi G_N} h_{\mu\nu}^{(n)}(\mathbf{x}) = \int d^{D-1}x e^{i\mathbf{k}\mathbf{x}} J_{\mu\nu}^{(n)}(\mathbf{k}),$$

Single multi-loop Master Integral

$$\sqrt{32\pi G_N} A_{bc}^a{}^{(n)}(\mathbf{x}) = \int d^{D-1}x e^{i\mathbf{k}\mathbf{x}} Y_{bc}^a{}^{(n)}(\mathbf{k}),$$

$$J_{\mu\nu}^{(n)}(\mathbf{k}) = \rho(|\mathbf{k}|, D, n) \left(\underline{\chi_1^{(n)}} \delta_\mu^0 \delta_\nu^0 + \underline{\chi_2^{(n)}} \eta_{\mu\nu} + \underline{\chi_3^{(n)}} \frac{k_\mu k_\nu}{\mathbf{k}^2} \right)$$

$$\rho(|\mathbf{k}|, D, n) = \frac{\Gamma\left(\frac{2-(D-3)(n-1)}{2}\right)}{\Gamma\left(\frac{n(D-3)}{2}\right)} \frac{(\Gamma\left(\frac{D-3}{2}\right) G_N m)^n}{(|\mathbf{k}|/(2\sqrt{\pi}))^{2-(D-3)(n-1)}}$$

$$Y_{bc}^a{}^{(n)}(\mathbf{k}) = -i\rho(|\mathbf{k}|, D, n) \left(k_{(b} \left(\underline{\chi_7^{(n)}} \delta_c^0 \delta_0^a + \underline{\chi_8^{(n)}} \delta_c^a \right) \right. \\ \left. + k^a \left(\underline{\chi_4^{(n)}} \delta_b^0 \delta_c^0 + \underline{\chi_5^{(n)}} \eta_{bc} + \underline{\chi_6^{(n)}} \frac{k_b k_c}{\mathbf{k}^2} \right) \right).$$

form factors

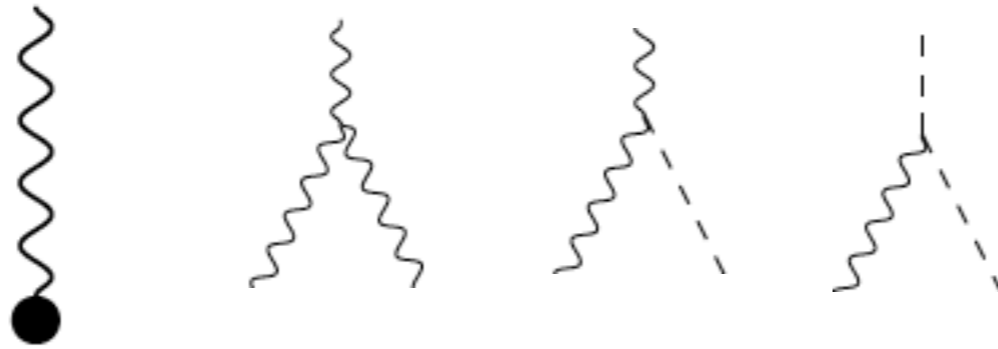
$$\chi^{(n)}(D) = (\chi_1^{(n)}, \dots, \chi_8^{(n)})$$

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Cubic formulation of GREFT

Feynman rules:



$$J_{\mu\nu}^{(3)} = \frac{1}{2} \left(\begin{array}{c} \text{wavy} \\ \text{wavy} \\ \text{wavy} \\ \bullet \\ J^{(1)} J^{(1)} J^{(1)} \end{array} + \begin{array}{c} \text{wavy} \\ \text{wavy} \\ \text{dashed} \\ \bullet \\ J^{(1)} J^{(1)} J^{(1)} \end{array} \right) = \begin{array}{c} \text{wavy} \\ \text{wavy} \\ \text{wavy} \\ \bullet \\ J^{(1)} J^{(2)} \end{array} - \begin{array}{c} \text{wavy} \\ \text{wavy} \\ \text{dashed} \\ \bullet \\ J^{(1)} Y^{(2)} \end{array}$$

$$Y_{bc}^{a(3)} = \frac{1}{2} \left(\begin{array}{c} \text{dashed} \\ \text{wavy} \\ \text{wavy} \\ \bullet \\ J^{(1)} J^{(1)} J^{(1)} \end{array} + \begin{array}{c} \text{dashed} \\ \text{wavy} \\ \text{dashed} \\ \bullet \\ J^{(1)} J^{(1)} J^{(1)} \end{array} \right) = \begin{array}{c} \text{dashed} \\ \text{wavy} \\ \text{wavy} \\ \bullet \\ J^{(1)} J^{(2)} \end{array} + \begin{array}{c} \text{dashed} \\ \text{wavy} \\ \text{dashed} \\ \bullet \\ J^{(1)} Y^{(2)} \end{array}$$

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Metric to all orders in G

Recursion relations

$$\chi_k^{(n)}(D) = \sum_{i,j=1}^8 \sum_{m=1}^{n-1} \chi_i^{(m)}(D) \chi_j^{(n-m)}(D) M_k^{ij}(D)$$

Solvable at D=4 \longrightarrow $\chi^{(n)}(4) = \left(8, 0, 0, 4 + n(-1)^n + \frac{1 + 3(-1)^n}{2n(n+2)}, \right.$

no UV-divs!!!

$$2 + \frac{1 + 3(-1)^n}{2n(n+2)}, \frac{1 + 3(-1)^n}{2n(n+2)}(n-3), \\ \left. \frac{1}{n} - 4 - \frac{1 + 3(-1)^n}{2n(n+2)}(n+1), \frac{1 + 3(-1)^n}{2n(n+2)} \right)$$

\longrightarrow **Resums to GR solution**

**GREFT computation “picks”
the simplest harmonic gauge**

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Geodesic motion (0SF)

[2308.14832,2406.14770] Cheung, Parra-Martinez, Rothstein, Shah, Wilson-Gerow

[2308.15304] Kosmopoulos, Solon

SF expansion: $\mathcal{S} = \mathcal{S}_{EH} + \mathcal{S}_l + \mathcal{S}_H$; $\mathcal{S}_H = -\frac{M}{2} \int d\tau_H \left(\frac{g^{\mu\nu} v_{H\mu} v_{H\nu}}{(\sqrt{-g})^{\frac{D-2}{2}}} + 1 \right)$ *n-point graviton vertices contrary to PM-expansion*

$$e^{i\mathcal{S}_{\text{eff}}[x_l, x_H]} = \int \mathcal{D}h \mathcal{D}A e^{i\mathcal{S}_{EH}[h, A] + i\mathcal{S}_{GF}[h] + i\mathcal{S}_l[x_l, h] + i\mathcal{S}_H[x_H, h]}$$

integrate-out via diagrams

$$\longrightarrow \mathcal{S}_{\text{eff}} = -\frac{M}{2} \int d\tau_H \eta^{\mu\nu} v_{H\mu} v_{H\nu} + M \sum_{n=0}^{\infty} \int d\tau_l \left(\frac{m}{M} \right)^{n+1} \mathcal{L}_n[x_l(\tau_l), x_H(\tau_H)]$$

+SF expand trajectories $x_l^\mu(\tau_l) \equiv x^\mu(\tau) = \sum_{n=0}^{\infty} \left(\frac{m}{M} \right)^n \delta x^{(n)\mu}(\tau)$, $x_H^\mu(\tau_H) = u_H^\mu \tau_H + \sum_{n=1}^{\infty} \left(\frac{m}{M} \right)^n \delta x_H^{(n)\mu}(\tau_H)$

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Geodesic motion (0SF)

SF expansion:
$$\mathcal{S}_{\text{eff}} = -\frac{M}{2} \int d\tau_H \eta^{\mu\nu} v_{H\mu} v_{H\nu} + M \sum_{n=0}^{\infty} \int d\tau_l \left(\frac{m}{M}\right)^{n+1} \mathcal{L}_n[x_l(\tau_l), x_H(\tau_H)]$$

$$\mathcal{L}_0[x^\mu(\tau), u_H^\mu \tau_H] = \bullet + \begin{array}{c} \bullet \\ | \\ \text{wavy} \\ | \\ \blacksquare \end{array} + \begin{array}{c} \bullet \\ / \backslash \\ | \quad | \\ \text{wavy} \quad \text{wavy} \\ | \quad | \\ \blacksquare \quad \blacksquare \end{array} + \begin{array}{c} \bullet \\ / \backslash / \backslash \\ | \quad | \quad | \quad | \\ \text{wavy} \quad \text{wavy} \quad \text{wavy} \quad \text{wavy} \\ | \quad | \quad | \quad | \\ \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \end{array} + \begin{array}{c} \bullet \\ / \backslash / \backslash / \backslash \\ | \quad | \quad | \quad | \quad | \quad | \\ \text{wavy} \quad \text{wavy} \quad \text{wavy} \quad \text{wavy} \quad \text{wavy} \quad \text{wavy} \\ | \quad | \quad | \quad | \quad | \quad | \\ \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \end{array} + \dots$$
 1) **Infinite tower of n-graviton worldline vertices**

where
$$J_{\mu\nu}(\mathbf{k}) = \begin{array}{c} \text{wavy} \\ | \\ \blacksquare \end{array} = \sum_{n=1}^{\infty} J_{\mu\nu}^{(n)}(\mathbf{k})$$
 2) Effective 1pt contains an infinite series -dressed graviton emission-, already known

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Geodesic motion (0SF)

2) Effective 1pt contains an infinite series
-dressed graviton emission-

$$\int_{\mathbb{R}^3} \frac{d^3\mathbf{k}}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}(\tau)} \int_{\mathbb{R}^{3L}} \prod_{i=1}^L \frac{d^3\mathbf{q}_i}{(2\pi)^3} \delta^{(3)}\left(\sum_{i=1}^L \mathbf{q}_i - \mathbf{k}\right) \xrightarrow{\text{decouples}} \int_{\mathbb{R}^{3L}} \prod_{i=1}^L \frac{d^3\mathbf{q}_i}{(2\pi)^3} e^{i\mathbf{x}(\tau)\cdot\mathbf{q}_i}$$

$\delta^{(3)}\left(\sum_{i=1}^L \mathbf{q}_i - \mathbf{k}\right) = \int_{\mathbb{R}^3} e^{i\mathbf{u}\cdot(\sum_{i=1}^L \mathbf{q}_i - \mathbf{k})} \frac{d^3\mathbf{u}}{(2\pi)^3}$

Fourier+ ∞ Loops

∞ decoupled Fourier

$$\begin{array}{c}
 + \text{wavy} + \text{wavy} + \text{wavy} + \text{wavy} + \text{wavy} + \dots \\
 \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare \quad \blacksquare
 \end{array}
 = \mathcal{I}_{\alpha_1\beta_1, \dots, \alpha_L\beta_L}^{(L)} = \rho^{2L} \prod_{i=1}^L \left(\delta_{\alpha_i}^0 \delta_{\beta_i}^0 \left(\frac{4(1+\rho)}{\rho(1-\rho)} - 1 \right) + n_{\alpha_i} n_{\beta_i} \right)$$

Schwarzschild from Amplitudes

[2407.09448],[2405.14421] S.M., P. Vanhove

Geodesic motion (0SF)

$$\mathcal{L}_0[x^\mu(\tau), u_H^\mu \tau_H] = \bullet + \begin{array}{c} \bullet \\ | \\ \text{---} \\ | \\ \blacksquare \end{array} + \begin{array}{c} \bullet \\ | \\ \text{---} \\ | \\ \text{---} \\ | \\ \blacksquare \end{array} + \begin{array}{c} \bullet \\ | \\ \text{---} \\ | \\ \text{---} \\ | \\ \text{---} \\ | \\ \blacksquare \end{array} + \begin{array}{c} \bullet \\ | \\ \text{---} \\ | \\ \text{---} \\ | \\ \text{---} \\ | \\ \text{---} \\ | \\ \blacksquare \end{array} + \dots = \frac{1}{2} v_\mu(\tau) v_\nu(\tau) \left(-\eta^{\mu\nu} + \sum_{L=1}^{\infty} \mathcal{I}_{\alpha_1\beta_1, \dots, \alpha_L\beta_L}^{(L)} \mathcal{T}_{(L)}^{\mu\nu\alpha_1\beta_1, \dots, \alpha_L\beta_L} \right)$$

recursion relations under control

$$\sum_{L=1}^{\infty} \mathcal{P}_{(L)}^{\alpha_1\beta_1, \dots, \alpha_L\beta_L} \mathcal{I}_{\alpha_1\beta_1, \dots, \alpha_L\beta_L}^{(L)} = \frac{1}{(1+\rho)^2} - 1.$$

resums

$$= -\frac{1}{2} v_\mu(\tau) v_\nu(\tau) g^{\mu\nu}(|\mathbf{x}|(\tau))$$

trivially gives geodesic eq.

Rediscovered America!

- 1) Trivial from GR perspective but highly non-trivial from EFT
- 2) Crucial stepping stone to go beyond 0-SF order

Conclusion

1. Derivation from flat to curved background with usual techniques. What about Kerr?
2. Interplay between Fourier transforms and loop integrals
3. **SF-EFT** setup can-in principle- be extended
4. Non-perturbative, EFT-based approach can be used for questions regarding BHs
(Love numbers, quantum effects etc.)