# Gravitational Waves from Amplitudes and EFT

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EFTs and Beyond - December 5, 2024



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## Gravitational Waves & Particle Physics



Era of gravitational-wave astronomy New insights into fundamental open questions

Present experiments



Future experiments



... and many more!



## Gravitational Waves & Particle Physics



Era of gravitational-wave astronomy

<u>Collider physics</u>: From proton-proton scattering to the Standard Model and beyond <u>Gravitational-wave physics</u>: From black hole-black hole merging to General Relativity and beyond

- New insights into fundamental open questions





## Gravitational Waves & Particle Physics



Era of gravitational-wave astronomy New insights into fundamental open questions

<u>Collider physics</u>: From proton-proton scattering to the Standard Model and beyond <u>Gravitational-wave physics</u>: From black hole-black hole merging to General Relativity and beyond

Quantum Field Theory for Gravitational-Wave Science







From the Binary's Evolution to the Observed Waveform

Merger Ringdown

Credit: ANTELIS, MORENO



### Quadrupole radiation formula

$$h \sim \frac{G}{r} \ddot{Q}$$



Credit: https://www.ligo.caltech.edu/



From the Binary's Evolution to the Observed Waveform



### Quadrupole radiation formula



Adiabatic approximation

From the Binary's Evolution to the Observed Waveform



Conservative motion via a Hamiltonian H

Energy loss due to radiation



### Quadrupole radiation formula



Adiabatic approximation

From the Binary's Evolution to the Observed Waveform



Credit: ANTELIS, MORENO

Conservative motion via a Hamiltonian(H)

Energy loss due to radiation





### Theoretical template

+ Measurement





### This talk $\longrightarrow$ Theoretical template +

Measurement











### Theoretical Gravitational-Wave Analysis Theoretical template Credit: ANTELIS, MORENO Numerical Relativity (NR) Effective One Body (EOB) [Buonanno, Damour (1998)] [Pretorius (2025)] Model built on theoretical calculations The 'truth' Verified against NR Resource intensive Resource efficient





Theoretical template

### Numerical Relativity (NR)

[Pretorius (2025)]

The 'truth' Resource intensive









[Pretorius (2025)]

This talk EOB

> LIGO-Virgo-KAGRA theoretical analysis









EOB



### From Amplitudes to the EOB Model

### EOB & NR comparison

EOB: Effective One Body NR: Numerical Relativity





### The Precision Frontier



Precision goal for future gravitational-wave detectors (e.g., Cosmic Explorer, Einstein Telescope, Lisa)

Same color = Same physical PM order

Landscape of Open Problems Relevant for Upcoming Experiments

[**DK**, Luna (2021)]

[Bern, DK, Luna, Roiban, Teng (2022)]

nPM =  $\mathcal{O}(G^a S^b)$ , with a + b = nPhysical Post-Minkowiskian counting





### From Quantum Amplitudes to Classical Hamiltonians























 $\mathscr{A}_{4} \sim \begin{array}{c} 1 \\ 2 \end{array} \begin{array}{c} 1' \\ 2 \end{array} \sim G \frac{p_{1}^{\mu} p_{1}^{\nu} \mathscr{P}_{\mu\nu\alpha\beta} p_{2}^{\alpha} p_{2}^{\beta}}{q^{2}} \\ (CoM) \end{array} \sim E_{1}E_{2} \left( \frac{Gn(\mathbf{p})}{\mathbf{q}^{2}} \right)$ 

(CoM)











$$\frac{p_1^{\nu} \mathscr{P}_{\mu\nu\alpha\beta} p_2^{\alpha} p_2^{\beta}}{q^2} \sim E_1 E_2 \left(\frac{G n(\mathbf{p})}{\mathbf{q}^2}\right) \sim E_1 E_2 V(\mathbf{p}, \mathbf{q})$$
(CoM)

(CoM)















(CoM)

















(CoM)  $p_1 = (E_1, \mathbf{p})$ 

Obtaining the  $\mathcal{O}(G)$  Potential









Full theory







Effective theory



 $p_1 = (E_1, \mathbf{p})$ 

Obtaining the  $\mathcal{O}(G)$  Potential





Full theory





The  $\mathcal{O}(G^2)$  potential

Effective theory









$$\delta V(\mathbf{p}, \mathbf{r}) = -\frac{G^2 \,\tilde{n}(\mathbf{p})}{r^2}$$

[Cheung, Rothstein, Solon (2018)]

(Non-spinning case)





## Modeling Non-Spinning Compact Objects in General Relativity





Matter (traditionally):  $S_M = -m \int d\tau$  Worldline picture

Using Quantum Field Theory to Capture the Binary's Evolution

Matter (traditionally):  $S_M$ 

<u>Effective description</u>: Replace black hole/neutron star with point particle.



Using Quantum Field Theory to Capture the Binary's Evolution

$$S_M = -m \int d\tau$$

Worldline picture

Matter (traditionally):  $S_{\Lambda}$ 

Effective description: Replace black hole/neutron star with point particle.

Matter: 
$$S_M = \int d^4x \sqrt{-g} \left(\frac{1}{2}g^{\mu\nu}\partial_{\mu}\phi\partial_{\nu}\phi - \frac{1}{2}m^2\phi^2\right)$$
 QFT picture





$$S_M = -m \int d\tau$$

Worldline picture

Using Quantum Field Theory to Capture the Binary's Evolution

Matter (traditionally):  $S_{\Lambda}$ 

<u>Effective description</u>: Replace black hole/neutron star with point particle.

Matter:

$$S_M = \int d^4x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} m^2 \phi^2 \right) \qquad \text{QFT picture}$$

QFT/Amplitudes tools:

Unitarity Double copy On-shell recursion

Using Quantum Field Theory to Capture the Binary's Evolution

+



$$S_M = -m \int d\tau$$

Worldline picture

Integration technology from collider physics





 $S_{M} = \left[ d^{4}x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi - \frac{1}{2} m^{2} \phi^{2} \right) \right]$ 



 $S_M = \left[ d^4 x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} m^2 \phi^2 \right) \right]$ 



### Computing the Amplitude

The Amplitude at  $\mathcal{O}(G^2)$  Required for Extracting the Hamiltonian



 $S_M = \int d^4x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} m^2 \phi^2 \right) \qquad S_G = -\frac{1}{16\pi G} \int d^4x \sqrt{-g} R$ 



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 $\mathscr{A}^{C} \leftrightarrow \text{Compton}$  $\mathscr{A}_{4}^{\mathrm{FT}} \leftrightarrow \mathrm{Full} \text{ theory}$ 







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#### Computing the Amplitude

 $\mathscr{A}^{C} \leftrightarrow \text{Compton}$  $\mathscr{A}^{\mathrm{FT}}_{\Lambda} \leftrightarrow \mathrm{Full} \text{ theory}$ 

#### The Amplitude at $\mathcal{O}(G^2)$ Required for Extracting the Hamiltonian





## Modeling Spinning Compact Objects in General Relativity

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Worldline picture:

QFT picture:  $S_M =$ 

$$S_M = \int d\tau \left( -m + \frac{1}{2} \Omega_{\mu\nu} S^{\mu\nu} + \dots \right)$$

$$= \int d^4x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} (\nabla_\mu \phi_s) \cdot (\nabla_\nu \phi_s) - \frac{1}{2} m^2 \phi_s \cdot \phi_s + \dots \right)$$







 $\phi_s^{\mu_1\ldots\mu_s} \leftrightarrow$ 

Worldline picture:

QFT picture:  $S_M =$ 

States in rest frame:  $\{ |s,s\rangle, |s,s-1\rangle, \dots |s,-s\rangle \}$ 

$$\begin{split} S_M &= \int d\tau \left( -m + \frac{1}{2} \Omega_{\mu\nu} S^{\mu\nu} + \dots \right) \\ &= \int d^4 x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} (\nabla_\mu \phi_s) \cdot (\nabla_\nu \phi_s) - \frac{1}{2} m^2 \phi_s \cdot \phi_s + \dots \right) \end{split}$$





 $\phi_s^{\mu_1\ldots\mu_s} \leftrightarrow$ 

Worldline picture:

QFT picture:  $S_M =$ 

States in rest frame:  $\{ |s,s\rangle, |s,s-1\rangle, \dots |s,-s\rangle \}$ 

 $\leftrightarrow$  spin magnitude fixed, spin orientation variable

$$S_{M} = \int d\tau \left( -m + \frac{1}{2} \Omega_{\mu\nu} S^{\mu\nu} + \dots \right)$$
$$= \int d^{4}x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} (\nabla_{\mu} \phi_{s}) \cdot (\nabla_{\nu} \phi_{s}) - \frac{1}{2} m^{2} \phi_{s} \cdot \phi_{s} + \dots \right)$$





 $\phi_s^{\mu_1\ldots\mu_s} \leftrightarrow$ 

Worldline picture:

QFT picture:  $S_M =$ 

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GW190412:  $m_1 \sim 30 M_{\odot}, \chi_1 \sim 0.4$ 

 $s_1 = 2\chi_1 G m_1^2 \sim 10^{79}$ 

 $\hbar = c = 1$ 





 $\phi_s^{\mu_1\ldots\mu_s} \leftrightarrow$ 

Worldline picture:

QFT picture:  $S_M =$ 

States in rest frame:

 $\{ |s,s\rangle, |s,s-1\rangle, \dots |s,-s\rangle \}$ 

↔ spin magnitude fixed,spin orientation variable

Number of indices

Using Quantum Field Theory to Capture the Binary's Evolution

$$S_{M} = \int d\tau \left( -m + \frac{1}{2} \Omega_{\mu\nu} S^{\mu\nu} + \dots \right)$$
$$= \int d^{4}x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} (\nabla_{\mu} \phi_{s}) \cdot (\nabla_{\nu} \phi_{s}) - \frac{1}{2} m^{2} \phi_{s} \cdot \phi_{s} + \dots \right)$$

GW190412:  $m_1 \sim 30 M_{\odot}, \chi_1 \sim 0.4$ 



 $\hbar = c = 1$ 





 $S_M = \left[ d^4 x \sqrt{-g} \left( \dots - \frac{C_{ES^2}}{2m^2} R_{f_1 a f_2 b} \nabla^a \phi_s \,\mathbb{S}^{(f_1} \mathbb{S}^{f_2)} \nabla^b \phi_s + \dots \right) \right]$ 

Amplitude Computation for Spinning Objects

 $\mathbb{S}^a \equiv -\frac{i}{2m} \epsilon^{abcd} M_{cd} \nabla_b$ 









$$S_{M} = \int d^{4}x \sqrt{-g} \left( \dots - \frac{C}{2} \right)$$
$$= \int d^{4}x \left( \dots + \kappa h_{\mu\nu} T^{\mu\nu} \right)$$

 $T^{\mu\nu}$ : Stress tensor  $\leftrightarrow$  Details of the body  $\leftrightarrow$  Equation of state

Amplitude Computation for Spinning Objects

 $\frac{C_{ES^2}}{2m^2} R_{f_1 a f_2 b} \nabla^a \phi_s \,\mathbb{S}^{(f_1} \mathbb{S}^{f_2)} \nabla^b \phi_s + \dots \Big) \qquad \qquad \mathbb{S}^a \equiv -\frac{i}{2m} \epsilon^{abcd} M_{cd} \nabla_b$ 

 $^{\prime\nu} + \mathcal{O}\left(h^2\right)$  $g_{\mu\nu} = \eta_{\mu\nu} + \kappa h_{\mu\nu}$ 







# Example Calculation at $O(S^2)$



 $S_{M} = \left[ d^{4}x \sqrt{-g} \left( \dots - \frac{C_{ES^{2}}}{2m^{2}} R_{f_{1}af_{2}b} \nabla^{a}\phi_{s} \mathbb{S}^{(f_{1}} \mathbb{S}^{f_{2})} \nabla^{b}\phi_{s} + \dots \right) \qquad \mathbb{S}^{a} \equiv -\frac{i}{2m} \epsilon^{abcd} M_{cd} \nabla_{b} \right]$  $= \int d^4x \left( \dots + \kappa h_{\mu\nu} T^{\mu\nu} + \mathcal{O}\left(h^2\right) \right)$ 

 $T^{\mu\nu}$ : Stress tensor  $\leftrightarrow$  Details of the body  $\leftrightarrow$  Equation of state

Amplitude Computation for Spinning Objects

 $g_{\mu\nu} = \eta_{\mu\nu} + \kappa h_{\mu\nu}$ 













 $T^{\mu\nu}$ : Stress tensor  $\leftrightarrow$  Details of the body  $\leftrightarrow$  Equation of state

 $S_{M} = \left[ d^{4}x \sqrt{-g} \left( \dots - \frac{C_{ES^{2}}}{2m^{2}} R_{f_{1}af_{2}b} \nabla^{a}\phi_{s} \mathbb{S}^{(f_{1}} \mathbb{S}^{f_{2})} \nabla^{b}\phi_{s} + \dots \right) \qquad \mathbb{S}^{a} \equiv -\frac{i}{2m} \epsilon^{abcd} M_{cd} \nabla_{b}$ 

$$g_{\mu\nu} = \eta_{\mu\nu} + \kappa h_{\mu\nu}$$

**[DK**, Luna (2021)]













 $T^{\mu\nu}$ : Stress tensor  $\leftrightarrow$  Details of the body  $\leftrightarrow$  Equation of state

 $S_M = \left[ d^4 x \sqrt{-g} \left( \dots - \frac{C_{ES^2}}{2m^2} R_{f_1 a f_2 b} \nabla^a \phi_s \,\mathbb{S}^{(f_1} \mathbb{S}^{f_2)} \nabla^b \phi_s + \dots \right) \qquad \mathbb{S}^a \equiv -\frac{i}{2m} \epsilon^{abcd} M_{cd} \nabla_b \right]$ 



**[DK**, Luna (2021)]















 $T^{\mu\nu}$ : Stress tensor  $\leftrightarrow$  Details of the body  $\leftrightarrow$  Equation of state

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**[DK**, Luna (2021)]











 $T^{\mu\nu}$ : Stress tensor  $\leftrightarrow$  Details of the body  $\leftrightarrow$  Equation of state

 $S_{M} = \int d^{4}x \sqrt{-g} \left( \dots - \frac{C_{ES^{2}}}{2m^{2}} R_{f_{1}af_{2}b} \nabla^{a}\phi_{s} \mathbb{S}^{(f_{1}}\mathbb{S}^{f_{2})} \nabla^{b}\phi_{s} + \dots \right) \qquad \mathbb{S}^{a} \equiv -\frac{i}{2m} \epsilon^{abcd} M_{cd} \nabla_{b}$  $g_{\mu\nu} = \eta_{\mu\nu} + \kappa h_{\mu\nu}$ 



**[DK**, Luna (2021)]









## Hamiltonian for Spinning Objects

$$H = \sqrt{\mathbf{p}^{2} + m_{1}^{2}} + \sqrt{\mathbf{p}^{2} + m_{2}^{2}} + V^{(0)}(\mathbf{r}^{2}, \mathbf{p}^{2}) + V^{(1,1)}(\mathbf{r}^{2}, \mathbf{p}^{2}) \frac{\mathbf{L} \cdot \mathbf{S}_{1}}{\mathbf{r}^{2}} + \dots$$
  
+  $V^{(2,4)}(\mathbf{r}^{2}, \mathbf{p}^{2}) \frac{(\mathbf{r} \cdot \mathbf{S}_{1})^{2}}{\mathbf{r}^{4}} + \dots$  [DK, Luna (20)  
 $\mathcal{O}(S^{2})$ 

$$+V^{(5,22)}(\mathbf{r}^2,\mathbf{p}^2)\frac{(\mathbf{L}\cdot\mathbf{S}_1)(\mathbf{p}\cdot\mathbf{S}_1)^2}{\mathbf{r}^6}$$



New Results for Binaries of Spinning Objects

2021)]



[Bern, **DK**, Luna, Roiban, Teng (2022)]  $\mathcal{O}(S^5)$ 

$$\mathscr{A}_{4}^{\mathrm{EFT}} = \mathscr{A}_{4}^{\mathrm{FT}} \Rightarrow H$$



### Sample Results for the Hamiltonian

 $S^{n} \qquad 1 \qquad 2 \qquad 3 \qquad 4 \qquad 5$ Number of
spin structures  $2 \qquad 9 \qquad 18 \qquad 43 \qquad 86$ 

$$H = \dots + V^{(2,4)}(\mathbf{r}^2, \mathbf{p}^2) \frac{(\mathbf{r} \cdot \mathbf{S}_1)^2}{\mathbf{r}^4} + \dots$$

$$V^{A}(\mathbf{r}^{2}, \mathbf{p}^{2}) = \frac{G}{|\mathbf{r}|}c_{1}^{A}(\mathbf{p}^{2}) + \left(\frac{G}{|\mathbf{r}|}\right)^{2}c_{2}^{A}(\mathbf{p}^{2}) + \mathcal{O}(G^{3})$$

$$E = \sqrt{\mathbf{p}^2 + m_1^2} + \sqrt{\mathbf{p}^2 + m_2^2} \qquad \sigma = \frac{\mathbf{p}^2 + \sqrt{\mathbf{p}^2 + m_1^2}}{m_1 m_2}$$

#### Hamiltonian Term at $\mathcal{O}(G^2S^2)$

$$\begin{split} & \sum_{2}^{2,1/4)} = \\ & - \left( \left[ \left( -1 + \sigma \right) \, m_2 \left( m_1^2 + 2 \, \sigma \, m_1 \, m_2 + m_2^2 \right)^5 \left( 4 \, \left( 1 + \sigma \right) \left( 3 \, \sigma^2 \left( -7 + 15 \, \sigma^2 \right) + \left( -12 - 29 \, \sigma^2 + 53 \, \sigma^4 \right) \, C_{ES^2} \right) \, m_1^8 + \\ & \left( \sigma \, \left( 1 + \sigma \right) \, \left( 3 \, \Xi \, \sigma \left( -7 + 15 \, \sigma^2 \right) + \left( -42 + 4 \, \sigma + 65 \, \sigma^2 + 161 \, \sigma^4 \right) \, m_2 \right) + C_{ES^2} \left( \Xi \, \left( -12 - 12 \, \sigma - 29 \, \sigma^2 + 53 \, \sigma^4 + 53 \, \sigma^5 \right) + \sigma \left( -142 - 130 \, \sigma - 251 \, \sigma^2 - 367 \, \sigma^3 + 529 \, \sigma^4 + 649 \, \sigma^5 \right) \, m_2 \right) \\ & m_1^6 \, m_2 \, \left( \left( 1 + \sigma \right) \, \left( 8 \, \Xi \, \sigma \left( -21 + 2 \, \sigma + 43 \, \sigma^3 + 58 \, \sigma^4 \right) + \left( -132 + 32 \, \sigma - 291 \, \sigma^2 - 208 \, \sigma^3 + 3766 \, \sigma^4 \right) \\ & 480 \, \sigma^5 - 1079 \, \sigma^6 \right) \, m_2 \right) + C_{ES^2} \left( 8 \, \Xi \, \sigma \left( -65 - 59 \, \sigma - 111 \, \sigma^2 - 169 \, \sigma^3 + 238 \, \sigma^4 + 298 \, \left( -32 + 64 \, \sigma - 3823 \, \sigma^2 - 4511 \, \sigma^3 - 346 \, \sigma^4 - 1706 \, \sigma^4 + 5659 \, \sigma^6 + 8696 \, \sigma^7 \right) \, m_2 \right) \right) + \\ & 4 \, m_2^2 \left( - \left( 1 + \sigma \right) \left( \Xi \, \sigma \left( -21 + 5 \, \sigma - 122 \, \sigma^2 + 33 \, \sigma^3 + 188 \, \sigma^4 - 90 \, \sigma^5 + 60 \, \sigma^6 \right) - \\ & \left( 4 + 40 \, \sigma^2 - 39 \, \sigma^3 + 98 \, \sigma^4 + 61 \, \sigma^5 - 176 \, \sigma^4 + 39 \, \sigma^3 \right) \, m_2 \right) \right) - \\ & m_1 \, m_2^6 \left( \left( 1 + \sigma \right) \left( \Xi \, \left( 14 + 40 \, \sigma - 1115 \, \sigma^2 + 560 \, \sigma^3 - 1716 \, \sigma^4 + 396 \, \sigma^5 + 2769 \, \sigma^5 - 720 \, \sigma^7 + 624 \, \sigma \\ & \sigma \, \left( 240 - 324 \, \sigma + 2121 \, \sigma^2 - 220 \, \sigma^3 + 66 \, \sigma^4 + 1568 \, \sigma^5 - 3083 \, \sigma^4 + 249 \, \sigma^7 \right) \, m_2 \right) + \\ & C_{ES^2} \left( \Xi \, \left( 36 + 12 \, \sigma + 813 \, \sigma^2 + 853 \, \sigma^3 - 1446 \, \sigma^4 - 578 \, \sigma^5 - 1806 \, \sigma^4 + 2851 \, \sigma^7 \right) \, m_2 \right) + \\ & C_{ES^2} \left( \Xi \, \left( -8 + 381 \, \sigma^2 + 242 \, \sigma^3 - 3290 \, \sigma^4 - 480 \, \sigma^5 + 1453 \, \sigma^6 \right) + \\ & \left( -16 + 978 \, \sigma + 428 \, \sigma^2 - 5771 \, \sigma^3 - 166 \, \sigma^4 - 757 \, \sigma^5 - 660 \, \sigma^6 + 2013 \, \sigma^7 \right) + \\ & \left( 267 + 232 \, \sigma - 788 \, \sigma^2 + 1426 \, \sigma^3 - 1446 \, \sigma^4 - 1357 \, \sigma^5 - 1184 \, \sigma^6 + 2569 \, \sigma^7 - \\ & 180 \, \sigma^4 + 183 \, \sigma^3 + 2440 \, \sigma^3 - 3483 \, \sigma^4 + 315 \, \sigma^5 + 8717 \, \sigma^5 - 660 \, \sigma^6 + 2813 \, \sigma^7 \right) + \\ & \left( 267 + 232 \, \sigma - 788 \, \sigma^2 - 1780 \, \sigma^3 - 7440 \, \sigma^4 - 13576 \, \sigma^5 - 1134 \, \sigma^6 + 2569 \, \sigma^7 - \\ & 180 \, \sigma^4 + 185 \, \sigma^9 + (1242 \, \sigma^3 + 1363 \, \sigma^4 + 105 \, \sigma^4 + 1037 \, \sigma^4 + 2356 \, \sigma^7 + 3189 \, \sigma^6 + 3283 \, \sigma$$





### Sample Results for the Hamiltonian

*S<sup>n</sup>* Number of 18 86 9 43 spin structures

$$H = \dots + V^{(2,4)}(\mathbf{r}^2, \mathbf{p}^2) \frac{(\mathbf{r} \cdot \mathbf{S}_1)^2}{\mathbf{r}^4} + \dots$$

$$V^{A}(\mathbf{r}^{2}, \mathbf{p}^{2}) = \frac{G}{|\mathbf{r}|}c_{1}^{A}(\mathbf{p}^{2}) + \left(\frac{G}{|\mathbf{r}|}\right)^{2}c_{2}^{A}(\mathbf{p}^{2}) + \mathcal{O}(G^{3})$$

$$E = \sqrt{\mathbf{p}^2 + m_1^2} + \sqrt{\mathbf{p}^2 + m_2^2} \qquad \sigma = \frac{\mathbf{p}^2 + \sqrt{\mathbf{p}^2 + m_1^2}}{m_1 m_2}$$



Hamiltonian Term at  $\mathcal{O}(G^2S^2)$ 



### Sample Results for the Hamiltonian

*S<sup>n</sup>* 3 5 Number of 9 18 86 43 spin structures

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Hamiltonian Term at  $\mathcal{O}(G^2S^2)$ 



### Quantum Field Theory around a Black Hole

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EFT in  $m/M \ll 1$  — Gravitational Self-Force (SF) expansion





EFT in  $m/M \ll 1$  — Gravitational Self-Force (SF) expansion

Motion = Motion in black-hole background +  $\mathcal{O}(m/M)$ 





EFT in  $m/M \ll 1$  — Gravitational Self-Force (SF) expansion

Motion = Motion in black-hole background +  $\mathcal{O}(m/M)$ 

Black hole in General Relativity:

$$\bar{g}_{\mu\nu}dx^{\mu}dx^{\nu} = (1 - \frac{2GM}{r})dt^2 - (1 - \frac{2GM}{r})^{-1}dr^2 - r^2d\Omega^2$$





EFT in  $m/M \ll 1$  — Gravitational Self-Force (SF) expansion

Motion = Motion in black-hole background +  $\mathcal{O}(m/M)$ 

Black hole in General Relativity:

Black hole in Quantum Field Theory:

$$\bar{g}_{\mu\nu}dx^{\mu}dx^{\nu} = (1 - \frac{2GM}{r})dt^{2} - (1 - \frac{2GM}{r})^{-1}dr^{2} - r^{2}d\Omega^{2}$$

$$+ \int \int f + \frac{1}{r} \int f + \cdots \approx \bar{g}_{\mu\nu} - \eta_{\mu\nu} \quad \text{[Duff (19)]}$$





EFT in  $m/M \ll 1$  — Gravitational Self-Force (SF) expansion

Motion = Motion in black-hole background +  $\mathcal{O}(m/M)$ 

Black hole in General Relativity:

Black hole in Quantum Field Theory:

New QFT framework:









### Manifest Power Counting



#### Tree level

Geodesic

 $0\text{SF} - \mathcal{O}\left(\left(m/M\right)^0\right)$ 

Diagram Loops Count SF order

#### 1-loop level





1SF -  $\mathcal{O}\left((m/M)^1\right)$ 

[**DK**, Solon (2023)]

• • •







Flat



Benefits: Classes of Diagrams Combined & Generically Fewer Integrals

#### Systematic Resummation



[**DK**, Solon (2023)]







Flat

Extreme-mass-ratio inspirals (EMRIs)



Benefits: Classes of Diagrams Combined & Generically Fewer Integrals

#### Systematic Resummation

Streamline calculations at high orders Enable calculations of spinning geodesics in Kerr background Model EMRIs analytically within QFT



[**DK**, Solon (2023)]







#### Outlook







#### Conclusions

EFT for Gravitational-Wave Science







#### EFT for Gravitational-Wave Science

Pushed the state-of-the-art in modeling binary systems

#### Conclusions







#### EFT for Gravitational-Wave Science

Pushed the state-of-the-art in modeling binary systems

\* Designed EFT for expanding in m/M and resumming classes of contributions to all orders in G

#### Conclusions





#### Conclusions

#### EFT for Gravitational-Wave Science

- Pushed the state-of-the-art in modeling binary systems
- $\blacksquare$  Designed EFT for expanding in m/M and resumming classes of contributions to all orders in G
- $\blacksquare$  Discovered new phenomena that may manifest in the waveform  $\Rightarrow$  Spin-magnitude change

[Bern, **DK**, Luna, Roiban, Scheopner, Teng, Vines (2023)] [Alaverdian, Bern, DK, Luna, Roiban, Scheopner, Teng (2024)] [Alaverdian, Bern, **DK**, Luna, Roiban, Scheopner, Teng (to appear)]







#### Conclusions

#### EFT for Gravitational-Wave Science

- Pushed the state-of-the-art in modeling binary systems
- $\bullet$  Designed EFT for expanding in m/M and resumming classes of contributions to all orders in G
- $\blacksquare$  Discovered new phenomena that may manifest in the waveform  $\Rightarrow$  Spin-magnitude change

Thank you!

[Bern, **DK**, Luna, Roiban, Scheopner, Teng, Vines (2023)] [Alaverdian, Bern, DK, Luna, Roiban, Scheopner, Teng (2024)] [Alaverdian, Bern, **DK**, Luna, Roiban, Scheopner, Teng (to appear)]





### Backup Slides



### Modeling Spin-Magnitude Change in Orbital Evolution



### Modeling Spin-Magnitude Change



 $S_M = \int d^4x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} (\nabla_\mu \Phi) \cdot (\nabla_\nu \Phi) - \frac{1}{2} m^2 \Phi \cdot \Phi + \dots \right)$ 

Using Quantum Field Theory to Capture the Binary's Evolution

[Bern, **DK**, Luna, Roiban, Scheopner, Teng, Vines (2023)] [Alaverdian, Bern, **DK**, Luna, Roiban, Scheopner, Teng (2024)] [Alaverdian, Bern, **DK**, Luna, Roiban, Scheopner, Teng (to appear)]

 $\Phi \sim \begin{vmatrix} \cdot \\ \phi_s \\ \phi_{s-1} \\ \cdot \end{vmatrix}$ 




## Modeling Spin-Magnitude Change



 $S_M = \int d^4x \sqrt{-g} \left(\frac{1}{2}g^{\mu\nu}(\nabla x)\right)$ 

States in rest frame:  $\{\ldots, | s, s_7 \in \{-s, \ldots, s\}\rangle,\$ 

 $\leftrightarrow$  spin magnitude variable

$$(\nabla_{\mu}\Phi)\cdot(\nabla_{\nu}\Phi)-\frac{1}{2}m^{2}\Phi\cdot\Phi+\dots$$

$$\Phi \sim \begin{pmatrix} \cdot \\ \cdot \\ \phi_s \\ \phi_{s-1} \\ \cdot \\ \cdot \\ \cdot \end{pmatrix}$$

 $|s-1, s_7 \in \{-s+1, ..., s-1\}\rangle, ...\}$ 

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#### Using Quantum Field Theory to Capture the Binary's Evolution





## Modeling Spin-Magnitude Change



 $S_M = \int d^4x \sqrt{-g} \left(\frac{1}{2}g^{\mu\nu}(\nabla_\mu)\right) d^4x \sqrt{-g} \left(\frac{1}{2}g^{\mu\nu}(\nabla_\mu$ 

#### States in rest frame:

$$\{\ldots, | s, s_z \in \{-s, \ldots, s\}\rangle,\$$

$$|s-1, s_z \in \{-s+1, ..., s-1\}\rangle, ...\}$$

#### $\leftrightarrow$ spin magnitude variable

Using Quantum Field Theory to Capture the Binary's Evolution

$$(\nabla_{\mu}\Phi)\cdot(\nabla_{\nu}\Phi)-\frac{1}{2}m^{2}\Phi\cdot\Phi+\dots$$



of the Little Group

**D** - Reducible representation

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 $\Phi \sim \begin{vmatrix} \cdot \\ \phi_s \\ \phi_{s-1} \\ \cdot \end{vmatrix}$ 





## Modeling Spin-Magnitude Change



 $S_M = \int d^4x \sqrt{-g} \left( \frac{1}{2} g^{\mu\nu} (\nabla_\mu \Phi) \cdot (\nabla_\nu \Phi) - \frac{1}{2} m^2 \Phi \cdot \Phi + \dots \right)$ 

### $\phi_s$ - Irreducible representation

#### of the Little Group

 $\Phi$  - Reducible representation

Using Quantum Field Theory to Capture the Binary's Evolution



[Bern, **DK**, Luna, Roiban, Scheopner, Teng, Vines (2023)]

 $\Phi \sim \begin{vmatrix} \cdot \\ \cdot \\ \phi_s \\ \phi_{s-1} \\ \cdot \end{vmatrix}$ 

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[Alaverdian, Bern, **DK**, Luna, Roiban, Scheopner, Teng (to appear)]





### New Structure for Compact Objects



### $H = H\left(\mathbf{r}, \mathbf{p}, \mathbf{S}_1, \mathbf{S}_2\right) + V^{(1,3)}(\mathbf{r}^2, \mathbf{p}^2) \frac{\mathbf{r} \cdot \mathbf{K}_1}{\mathbf{r}^2} + \dots$



$$V^{A}(\mathbf{r}^{2}, \mathbf{p}^{2}) = \frac{G}{|\mathbf{r}|} c_{1}^{A}(\mathbf{p}^{2}) + \left(\frac{G}{|\mathbf{r}|}\right)^{2} c_{2}^{A}(\mathbf{p}^{2}) +$$

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#### Novel Effect in Conservative Dynamics





## New Structure for Compact Objects



# $H = H\left(\mathbf{r}, \mathbf{p}, \mathbf{S}_1, \mathbf{S}_2\right) + V^{(1,3)}(\mathbf{r}^2, \mathbf{p}^2) \frac{\mathbf{r} \cdot \mathbf{K}_1}{\mathbf{r}^2} + \dots$

Additional multipolar

structure



$$V^{A}(\mathbf{r}^{2}, \mathbf{p}^{2}) = \frac{G}{|\mathbf{r}|} c_{1}^{A}(\mathbf{p}^{2}) + \left(\frac{G}{|\mathbf{r}|}\right)^{2} c_{2}^{A}(\mathbf{p}^{2}) +$$

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Additional multipolar

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$$V^{A}(\mathbf{r}^{2}, \mathbf{p}^{2}) = \frac{G}{|\mathbf{r}|}c_{1}^{A}(\mathbf{p}^{2}) + \left(\frac{G}{|\mathbf{r}|}\right)^{2}c_{2}^{A}(\mathbf{p}^{2}) + C_{2}^{A}(\mathbf{p}^{2}) + C_{2$$

$$\Delta \mathbf{S}_{1}^{2} = \frac{4GE_{1}E_{2}\left(K_{1z}^{(0)}S_{1y}^{(0)} - K_{1y}^{(0)}S_{1z}^{(0)}\right)c_{1}^{(1,3)}(p_{\infty}^{2})}{b\,p_{\infty}(E_{1} + E_{2})} + \dots$$

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### Novel Effect in Conservative Dynamics





### New Structure Manifests in the Waveform





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#### Large Modification to Observables



### New Structure Manifests in the Waveform









$$\Delta h_{+/\times}^{(1)} = h_{+/\times}^{(1)}(\dots, \mathbf{K} = \mathbf{S}) - h_{+/\times}^{(1)}(\dots, \mathbf{K} = \mathbf{0})$$

[Alaverdian, Bern, **DK**, Luna, Roiban, Scheopner, Teng (2024)] [Alaverdian, Bern, **DK**, Luna, Roiban, Scheopner, Teng (to appear)]

#### Large Modification to Observables



### New Structure Manifests in the Waveform

Spin dynamics are much **richer** than previously thought







### Large Modification to Observables

$$\Delta h_{+/\times}^{(1)} = h_{+/\times}^{(1)}(\dots, \mathbf{K} = \mathbf{S}) - h_{+/\times}^{(1)}(\dots, \mathbf{K} = \mathbf{0})$$

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