

Functional Multi-Loops Matching

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[\[2412.XXXXX\]](#)



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FTAE
High Energy Theory



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YOUNGST@ARS-EFTs and Beyond—December 3, 2024

Effective Field Theory

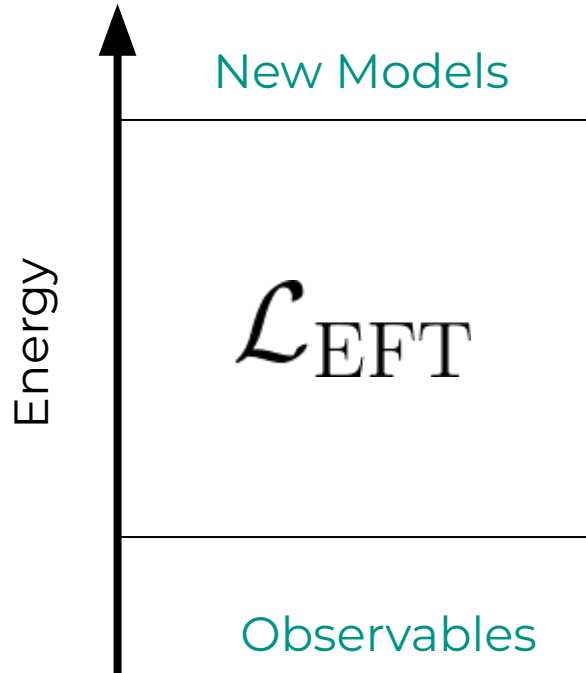
They represent our best connection between theory and experiment

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{d \leq 4} + \sum_{\ell=0}^{\infty} \sum_{n=5}^{\infty} \frac{C^{(\ell,n)}}{\Lambda^{n-4}} \mathcal{O}_n$$

Bottom up: Parameterize our lack of knowledge

Top Down: Separate scales for precision measurements

Matching



- **Connection** of UV theories to low energy observables
- Automated up to one-loop



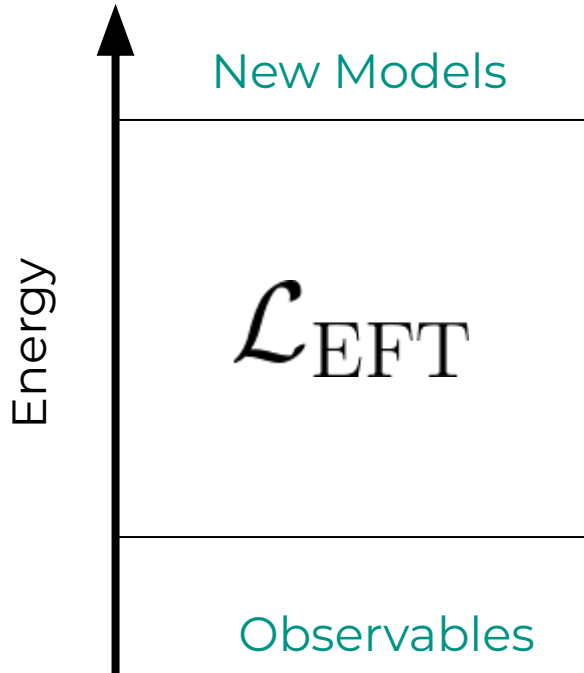
[Carmona et al-2112.10787]



[Fuentes-Martín et al-2212.04510]

- Running of the theory via RG evolution

Matching



Why **more loops**?

- There is physics we are losing at the one loop level
- Scheme independence one loop matching demands two loop running.
- Precision test of UV models are also required for future high-luminosity experiments.

Amplitude Matching

$$\mathcal{L}_{UV}(z_h, z_l) \xrightarrow{q_i \ll \Lambda} \{\mathcal{A}_{UV}(q_i)\}$$

Matching: Determining
Wilson Coefficients

$$\mathcal{L}_{EFT}(z_l) \longrightarrow \{\mathcal{A}_{EFT}(q_i)\}$$

Amplitude Matching



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$$\mathcal{L}_{UV}(z_h, z_l) \xrightarrow{q_i \ll \Lambda} \{\mathcal{A}_{UV}(q_i)\}$$

Matching: Determining
Wilson Coefficients

Feynman Diagrams

- Well-established and automated.
- Ansatz: Redundancies, redefinitions...
- Breaking of gauge symmetry in intermediate steps: additional computations

$$\mathcal{L}_{EFT}(z_l) \longrightarrow \{\mathcal{A}_{EFT}(q_i)\}$$

Functional Matching

Quantum Effective Action

We are going to “**Integrate Out**” the heavy fields

$$e^{i\Gamma_{UV}} = \int [D\Phi][D\phi] \exp \left(\int d^d x \mathcal{L}_{UV}[\Phi, \phi] \right)$$

Quantum Effective Action

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$$e^{i\Gamma_{UV}} = \int [D\Phi][D\phi] \exp \left(\int d^d x \mathcal{L}_{UV}[\Phi, \phi] \right)$$

Method of Regions: Multiscale ($m \ll M$) integrals can be separated in regions

$$\Gamma_{UV} = \Gamma_{UV} \Big|_{\text{hard}} + \Gamma_{UV} \Big|_{\text{soft}}$$

Quantum Effective Action

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Matching Condition

[Fuentes-Martín, Palavrić , Eller Thomsen-2311.13630]

$$S_{\text{EFT}} = \Gamma_{UV}[\bar{\Phi}[\phi], \phi] \Big|_{\text{hard}}$$

Quantum Effective Action

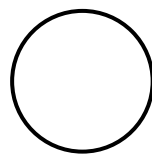
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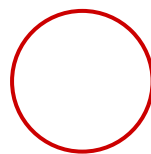
[Fuentes-Martín, Palavrić, Eller Thomsen-2311.13630]

$$S_{\text{EFT}} = \Gamma_{UV}[\bar{\Phi}[\phi], \phi] \Big|_{\text{hard}}$$



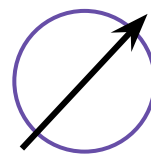
UV

=



EFT

+



Integrate
Out

Wilson Coefficients (Local)

Background Field Method

$$\phi = \bar{\phi} + \hat{\phi}$$

 $\bar{\phi}$

Classical Configuration: Tree Level

 $\hat{\phi}$

Quantum Fluctuation: Loops

Expanding the Lagrangian

$$\mathcal{L}_{UV}(\bar{\phi} + \hat{\phi}) = \mathcal{L}_{UV}(\bar{\phi}) + \frac{1}{2}\phi_i Q_{ij} \phi_j + \dots$$

Background Field Method

$$\phi = \bar{\phi} + \hat{\phi}$$

$\bar{\phi}$

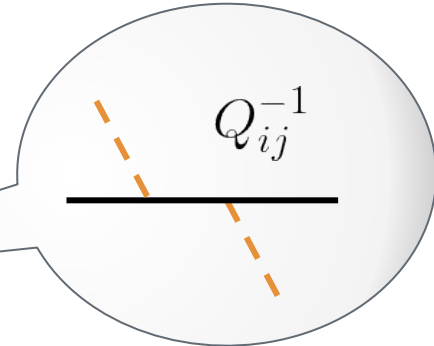
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- **One-loop**

$$\exp(i\Gamma_{\text{UV}}^{(1)}[\hat{\phi}]) = \int [D\phi] \exp \frac{1}{2} \left(\int d^d x \phi_i Q_{ij} \phi_j \right)$$

- **One-loop**

$$\exp(i\Gamma_{\text{UV}}^{(1)}[\hat{\phi}]) = \int [D\phi] \exp \frac{1}{2} \left(\int d^d x \phi_i Q_{ij} \phi_j \right)$$



Gaussian Integration

$$\Gamma_{\text{UV}}^{(1)} = \frac{i}{2} \text{STr} \ln Q$$

Differential operators under a Gauge Symmetry

$$Q_{ij}^{ab}(x, y) = Q_{ij}^{ac}(x, P_x) \delta_c^b(x, y)$$

Differential operators under a Gauge Symmetry

$$\delta_c^b(x, y) = \delta(x - y)U_c^b(x, y)$$

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Locality and gauge invariance of the action: Functional
Traces are dressed loop integrals

$$\begin{aligned}\text{Tr} \ln Q|_{\text{hard}} &= \int_{x,y} \delta_b^a(x, y) \ln Q_{ii}^{bc}(x, P_x)\delta_c^a(x, y) \\ &= \int_{x,k} \ln Q_{ii}^{ab}(x, P_x+k)U_b^a(x, y)\Big|_{x=y}\end{aligned}$$

Differential operators under a Gauge Symmetry

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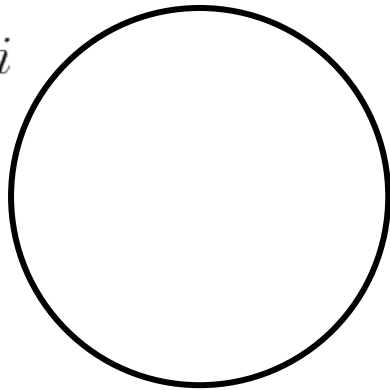
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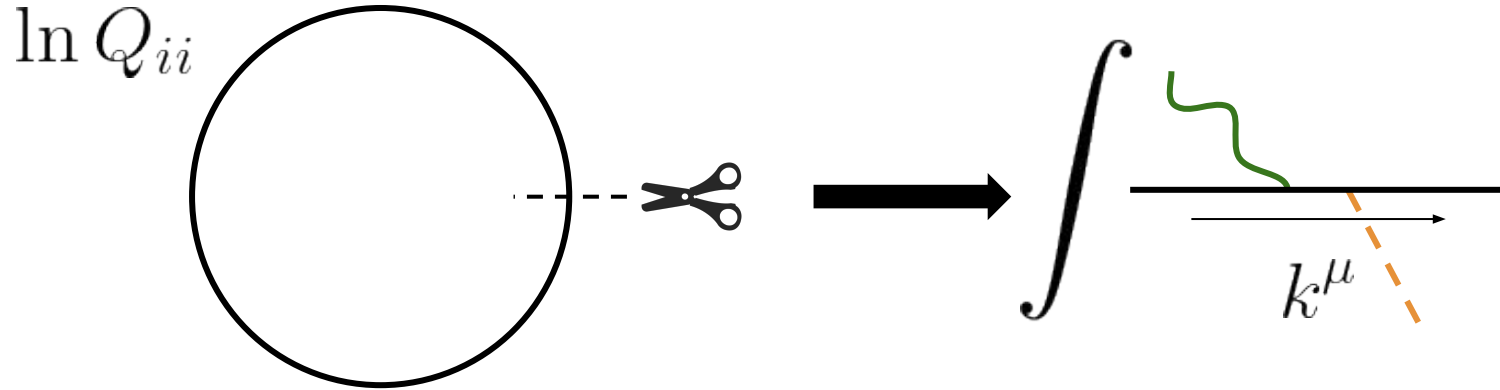
OPE around $k \sim \Lambda$
=
Explicit Gauge
Invariance

Diagrammatically,

$\ln Q_{ii}$



Diagrammatically,

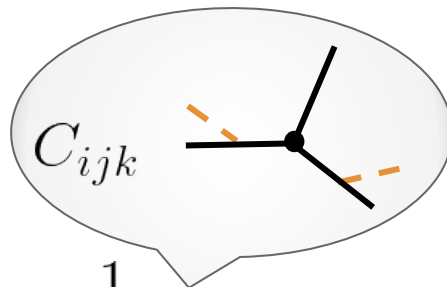
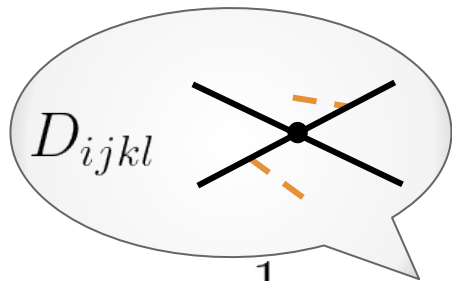


Operators traced in different points of spacetime remain local by a **momentum shift** operation

- **Two Loops:** More Topologies involved

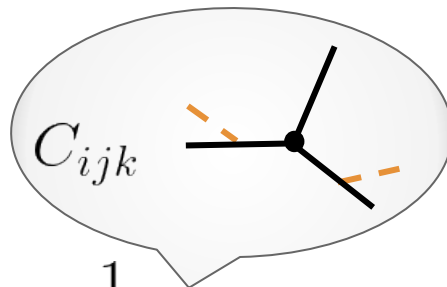
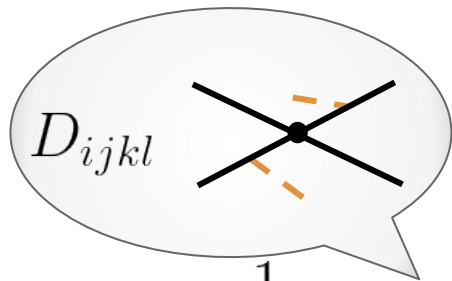
$$\Gamma_{UV}^{(2)}[\bar{\phi}] = \frac{i}{2} Q_{ij}^{-1} B_{ij} - \frac{1}{8} Q_{ij}^{-1} D_{ijkl} Q_{kl}^{-1} + \frac{1}{12} C_{ijk} Q_{il}^{-1} Q_{jm}^{-1} Q_{kn}^{-1} C_{lmn}$$

- **Two Loops:** More Topologies involved



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- **Two Loops:** More Topologies involved



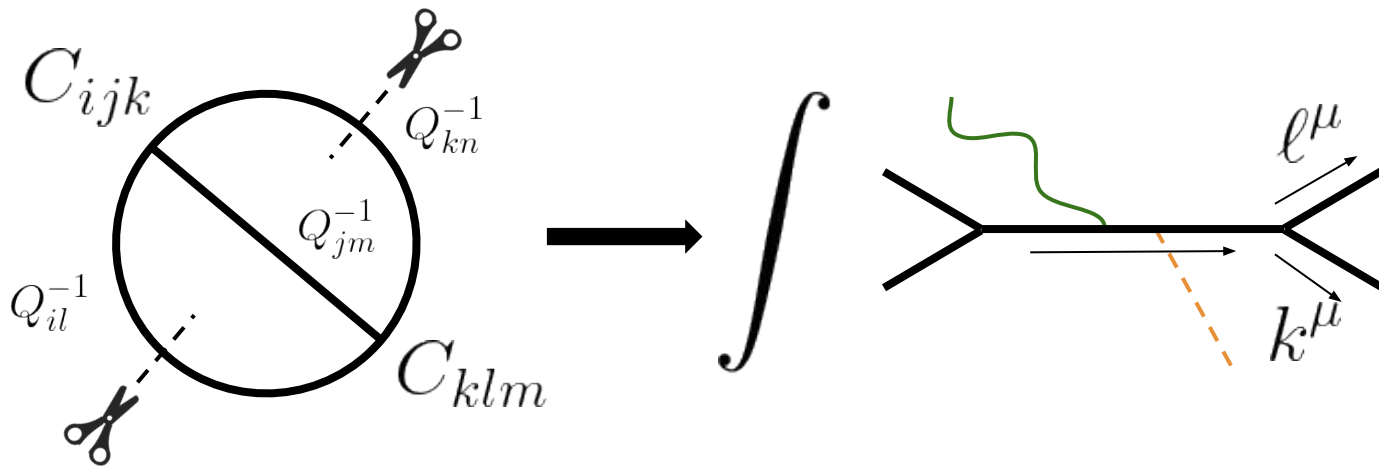
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Three Feynman diagrams representing two-loop topologies, each enclosed in a dashed circle. The first diagram is a single loop with a central black dot and a label (1) next to it, with a coefficient $\frac{i}{2}$ to its left. The second diagram is a figure-eight loop with a coefficient $\frac{1}{12}$ to its left. The third diagram is a pair of connected loops with a coefficient $\frac{1}{8}$ to its left.

$$\frac{i}{2} \text{(1)} + \frac{1}{12} \text{---} - \frac{1}{8} \text{---}$$

$$\begin{aligned}
G_{\text{ss}}|_{\text{hard}} = & \sum_{n,m,n'm'} (-1)^{n+m} \int_x \int_{\mathbf{k},\ell} C_{abc}^{(n,m)} \mathcal{Q}_{aa'}^{-1}(y, P_y - \mathbf{k} - \ell) C_{a'b'c'}^{(n',m')}(y) \\
& \times [(P_x + \mathbf{k})^m \mathcal{Q}_{be}^{-1}(x, P_x + \mathbf{k}) (P_x + \mathbf{k})^{m'} U_{b'e}(x, y)] \\
& \times [(P_x + \ell)^n \mathcal{Q}_{cf}^{-1}(x, P_x + \ell) (P_x + \ell)^{n'} U_{c'f}(x, y)] \Big|_{x=y}
\end{aligned}$$

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\end{aligned}$$



To take home



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- Functional matching is a powerful alternative for multi-loop computations.
- No need of ansatz but still we need basis reduction.
- Well suited for automation: Matchete.
- Gauge theories are fully understood: future work in other theories.

Thank you!