



### Lecture 2: Symanzik EFT — an example



# Precision Lattice Calculations

University of Illinois Urbana-Champaign

### **Frontiers and Challenges in Lattice** Gauge Theory 21 July – 1 August 2025

https://indico.mitp.uni-mainz.de/event/419/

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### Outline

- Motivation: Open Questions
- - History
  - Challenges
- Se Examples
  - Leptonic kaon, pion decay
  - Semileptonic kaon decay
  - First row CKM unitarity
  - Semileptonic D-meson decay
  - Leptonic B-meson decay

♀ muon g-2 → Lecture 3



Solution The role of (lattice) QCD in flavor physics

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# Lattice QCD Introduction: heavy quarks



 $m_b \gtrsim a^{-1} \gg \Lambda$  is leading discretization errors  $\sim (am_b)^2$ 



use EFT (HQET, NRQCD)  $\implies \Lambda/m_b$  expansion

- lattice HQET, NRQCD: use EFT to construct lattice action complicated continuum limit nontrivial matching and renormalization
- matching relativistic lattice action via HQET to continuum nontrivial matching and renormalization



 $a^{-1} > m_h \gg \Lambda + \text{highly improved light quark action}$ 

same action for all quarks simple renormalization (Ward identities)



(using same action as for light quarks)

EFTs co-developed continuum/lattice

# (few-5)% errors • relativistic heavy quarks: Fermilab (1996), also Tsukuba (2003), RHQ (2006) (1-3)% errors





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Standard Model predictions: Buras, et al [arXiv:1303.3820, JHEP 2013], Bobeth, et al [arXiv:1311.0903, PRL 2014; arXiv:2104.09521], Beneke et al [arXiv:1908.07011, JHEP 2019].

$$\overline{\mathcal{B}}(B_q \to \mu \bar{\mu}) = \frac{|\mathcal{N}_q|^2 M_{B_q}^3 f_{B_q}^2}{8\pi \Gamma_q^H} \beta_{q\mu} \left(\frac{m_\mu}{M_{B_q}}\right)^2 \left|C_{10}^{\text{eff}}\right|^2, \qquad \beta_{q\mu} \equiv \sqrt{1 - \frac{4 m_\mu^2}{M_{B_q}^2}},$$
$$\overline{\mathcal{R}}(B_q \to q^{-1}) = (2 C C + 0.14) - 10^{-9}$$

 $\mathcal{B}(B_s \to \mu \bar{\mu})_{\rm SM} = (3.66 \pm 0.14) \cdot 10^{-9}$ 





 includes structure-dependent QED corrections • dominant uncertainty due to  $|V_{cb}|$ • LQCD decay constant sub dominant source of uncertainty

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$$\mathcal{B}(\mathrm{B}^0_{\mathrm{s}} \to \mu^+\mu^-) = \left[3.83^{+0.38}_{-0.36} \text{ (stat)}^{+0.19}_{-0.16} \text{ (syst)}^{+0.14}_{-0.13} (f_{\mathrm{s}}/f_{\mathrm{s}})
ight]$$









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ight]$$







## Heavy Quark discretization errors for HISQ fermions

Tree-level HISQ action:

$$S = \sum_{x} \overline{\psi}(x) \left\{ \sum_{\mu} \gamma_{\mu} \left[ a \Delta_{\mu} - \frac{\mathbb{N}}{6} a^{3} \Delta_{\mu}^{3} \right] + a m_{0} \right\}$$

Naik term:

$$m_{1} \equiv E(0) \quad \frac{m_{1}}{m_{2}} = \lim_{p \to 0} \frac{E^{2}(p) - E(0)^{2}}{p^{2}} = 1$$
$$\implies \mathbb{N}(am_{1}) = \frac{4 - 2\sqrt{1 + 3X(0)}}{\sinh^{2}(am_{1})}$$

Normalization of heavy-light bilinears:

$$Z_{J_{hx}} = \widetilde{\mathcal{Ch}}_h^{1/2} \qquad \widetilde{\mathcal{Ch}} = \cosh am_1 \left(1 - \frac{1}{2}\mathbb{N}\sinh^2\right)$$



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Follana et al [HPQCD, hep-lat/0610092, 2007 PRD] Monahan et al [arXiv:1211.6966, 2013 PRD] Bazavov et al [FNAL/MILC, arXiv:1712.09262, 2018 PRD]

$$\psi(x) \qquad a\Delta_{\mu}\psi(x) = \frac{1}{2}[\psi(x+\hat{\mu}a) - \psi(x-\hat{\mu}a)],$$

$$am_0 = \sinh(am_1) \frac{1 + \sqrt{1 + 3X}}{3}$$

 $(am_1)$ 

$$X(am_1) = \frac{2am_1}{\sinh(2am_1)}$$

 $am_1$  '

remaining discretization errors:  $\sim (am_h)^4, \alpha_s(am_h)^2$ 





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### *B*, *D* meson decay constant results

A. Bazavov et al [FNAL/MILC, arXiv:1712.09262, 2018 PRD]





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### B, D meson decay constant results



- no renormalization (Ward identity)



