

Probing QED and nuclear physics with heavy muonic atoms

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Precision Determination of the Fine-structure Constant
Mainz, October 29, 2025



MAX-PLANCK-GESSELLSCHAFT



Outline

Introduction

Theory

Results

BSM

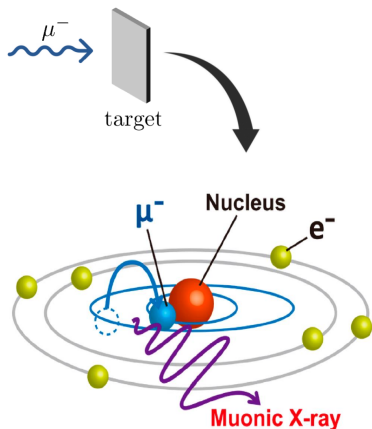
Outlook



“Live fast, die young!”

<https://www.particlezoo.net>

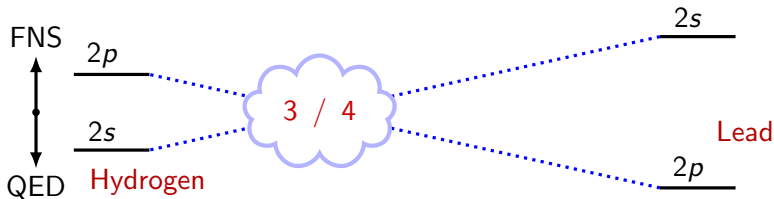
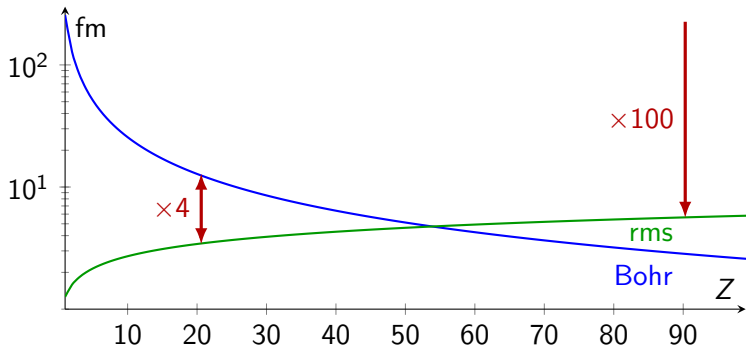
Access to muonic atoms



<http://www.mdpi.com/2412-382-X/1/1/11/htm>

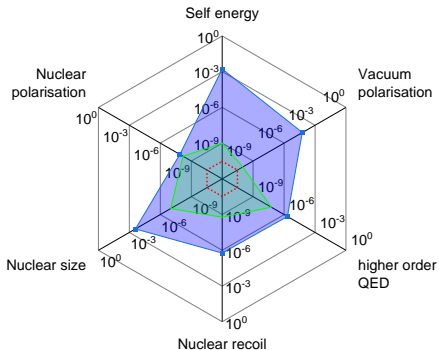
- “heavy electron”
- capture and cascade:
 $10^{-12} - 10^{-9} \text{ s}$
- lifetime: $0.1 - 2.2 \mu\text{s}$
- decay channels
 - $\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$
 - $\mu^- + p \rightarrow n + \nu_\mu$
- some electrons far away
- always H-like
- passive spectroscopy

How different are muonic atoms compared to electronic?

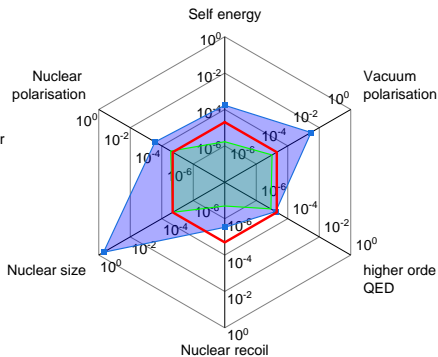


QED and nuclear effects' scaling

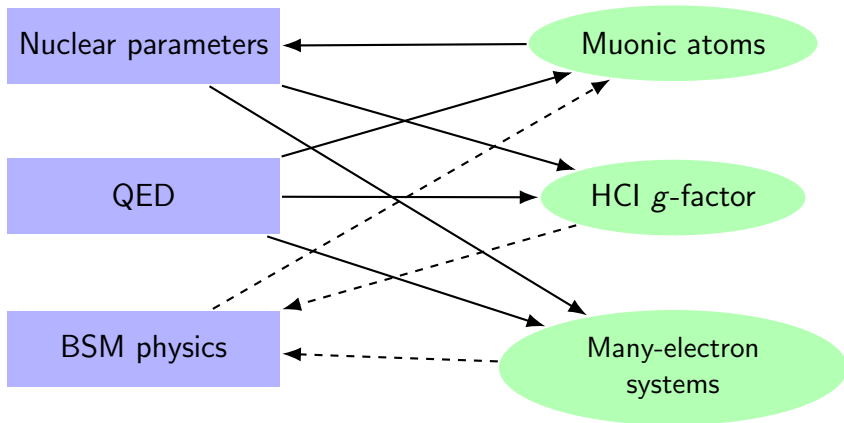
(A) g -factor hydrogenlike $^{208}\text{Pb}^{81+}$



(B) $2p_{3/2} \rightarrow 1s_{1/2}$ muonic ^{208}Pb

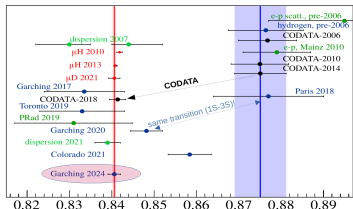


Connections



Muonic anomalies

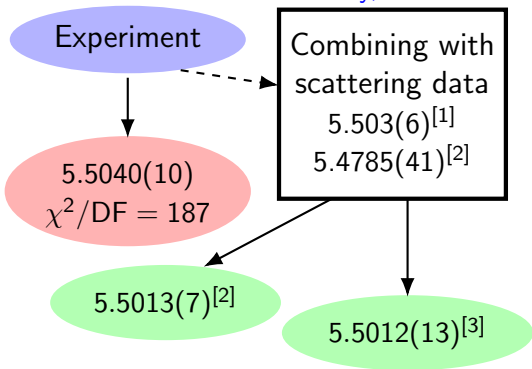
Proton radius puzzle



From: Randolf Pohl

Observed in ^{208}Pb , ^{120}Sn , and ^{90}Zr
 Blamed it on theory (NP)

Muonic fine-structure anomaly, Pb



[1] Bergem *et al.*, PRC **37**, 2821 (1988)

[2] Fricke and Bernhardt, At. Data Nucl. Data Tables **60**, 177 (1995)

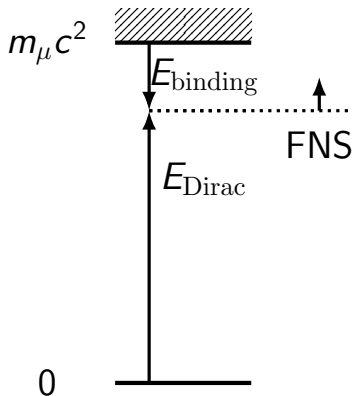
[3] Angeli and Marinova, At. Data Nucl. Data Tables **99**, 69 (2013)

Historical data

| Source | Value | Comment |
|-----------------|---|---|
| Belicard 1967 | 5.38(3), 5.39(3), 5.42(5), 5.54 | e scattering, obtained from different electron incident energies and 3pF nuclear model |
| Anderson 1969 | 5.4978(30) | muonic atoms |
| Heisenberg 1969 | 5.535, 5.539, 5.546 5.501, 5.502 | e scattering, 3pF and 4pF models e scattering + muonic data |
| Dreher 1974 | 5.498(15), 5.520(23), 5.498(10), 5.514(28) | e scattering, different energies and fitting models |
| Kessler 1975 | 5.5097(6) | muonic atoms, $\chi^2/N = 0.11^1$ |
| Euteneuer 1976 | 5.500(24) 5.494(24) | e scattering 3pF, $\chi_{\min}^2 \approx 24$ "model independent" |
| Euteneuer 1978 | 5.4927(75) 5.5032(16) | e scattering, $\overline{\chi^2} \approx 20$ e scattering + muonic, $\overline{\chi^2} \approx 22$ |
| Bergem 1988 | 5.5040(11) | muonic atoms, $\chi^2/\text{DoF} \approx 200$ |

¹all lines with 3d were excluded to improve the quality of the fit

Dirac value and nuclear size



- Muons are close to the nucleus, relativistic \rightarrow Dirac equation

- $m_\mu c^2 \approx 100$ MeV
 $E_{\text{Dirac}} \approx 80$ MeV
 $E_{\text{binding}} \approx 20$ MeV

- Extended nucleus: sphere, Fermi, deformed Fermi

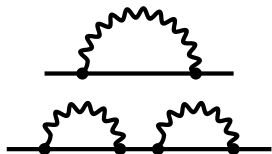
$$\rho_{a,c}^F(r_\mu) = \frac{N}{1 + e^{(r-c)/a}}$$

- $\Delta E_{\text{FNS}} \approx 10$ MeV

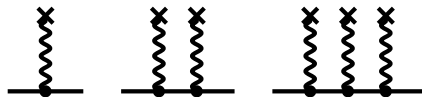
How can we do calculations?

Perturbation theory:

in α



in αZ

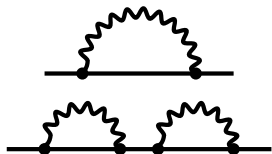


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 "Furry picture"

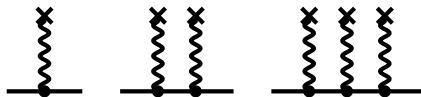
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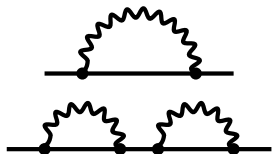


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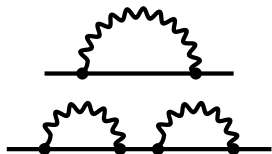


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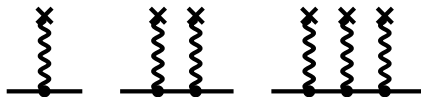
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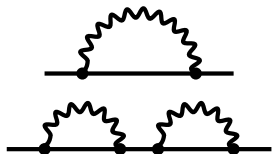


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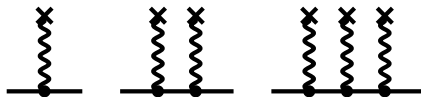
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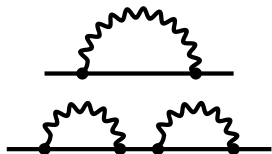


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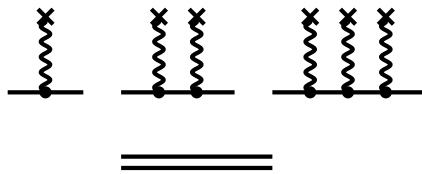
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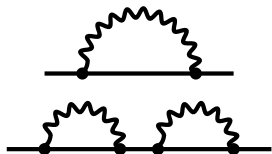


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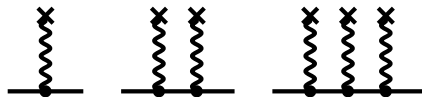
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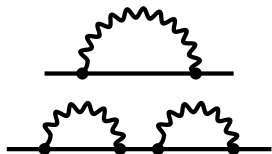


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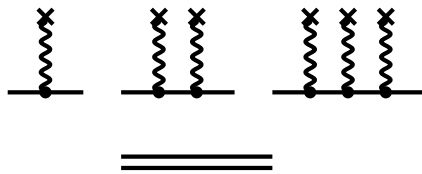
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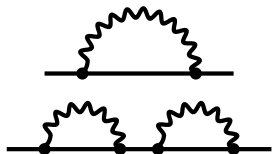


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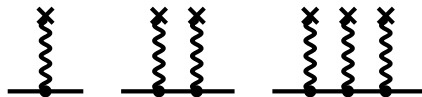
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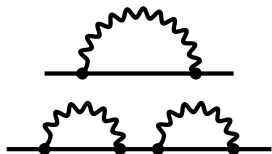


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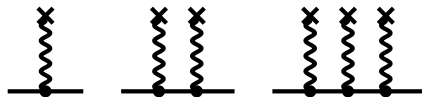
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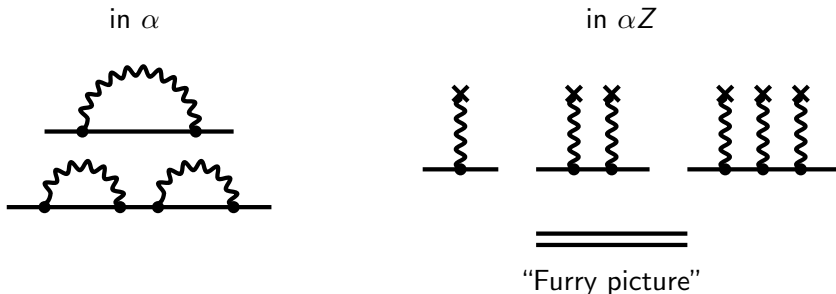
in αZ



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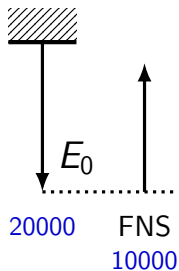
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Perturbation theory:

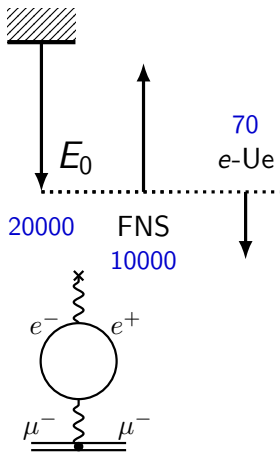


The two approaches can be independent or complimentary!

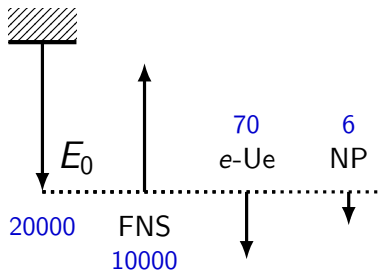
QED effects:



QED effects: Uehling



QED effects: Nuclear polarization

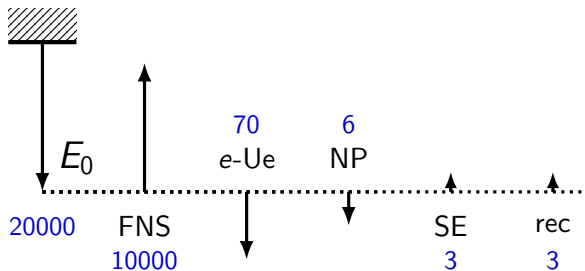


Improvements:

- field-theory approach, including transverse part
- state-of-art muonic and nuclear input, model dependence
- $0^+, 1^-, 2^+, 3^-, 4^+, 5^-$ and 1^+ excitation modes
- 4252 eV \rightarrow 5712 eV

Valuev *et al.*, PRL **128** 203001 (2022)

QED effects: Self energy and recoil



- rigorous QED calculations
- $\Delta E_{\text{SE}} = 3270(160)^{[1]}, 3373^{[2]} \text{ eV} \rightarrow 3225(15)^{[3]} \text{ eV}$
- $\Delta E_{\text{rec}} = 385^{[4]*} \text{ eV} \rightarrow 3902^{[5]} \text{ eV}$

[1] Cheng *et al.*, PRA **17**, 489 (1978)

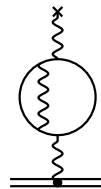
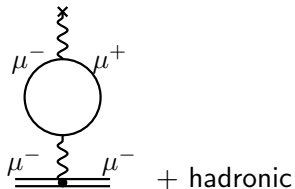
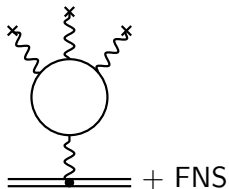
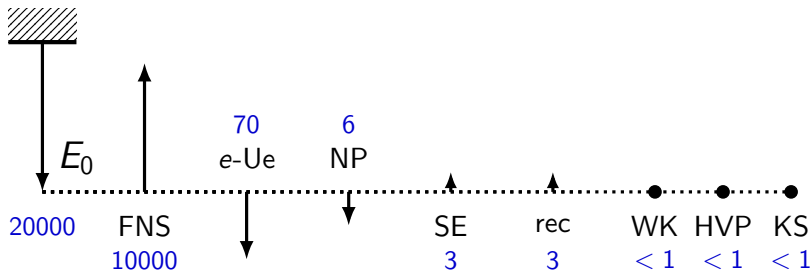
[2] Haga *et al.*, PRC **75**, 044315 (2007)

[3] Oreshkina, PRR **4**, L042040 (2022)

[4] Bergem *et al.*, PRC **37**, 2821 (1988)

[5] Yerokhin and Oreshkina, PRA **108**, 052824 (2023)

QED effects: sub-leading



Recoil with reduced mass: all orders in m/M

Schödinger equation:

$$H\psi = E\psi$$

$$\left(-\frac{\hbar^2}{2m_e} \nabla_e^2 - \frac{\hbar^2}{2m_N} \nabla_N^2 + V(|\mathbf{r}_e - \mathbf{r}_N|) \right) \psi(\mathbf{r}_e, \mathbf{r}_N) = E\psi(\mathbf{r}_e, \mathbf{r}_N)$$

$$\{\mathbf{r}_e, \mathbf{r}_N\} \rightarrow \left\{ \mathbf{r} = \mathbf{r}_e - \mathbf{r}_N, \mathbf{R} = \frac{m_e \mathbf{r}_e + m_N \mathbf{r}_N}{m_e + m_N} \right\}$$

$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right) \psi(\mathbf{r}) = E\psi(\mathbf{r})$$

Dirac equation:

$$H_D = \alpha \mathbf{p} + \beta m + V$$

Recoil with reduced mass: all orders in m/M

Schödinger equation:

$$H\psi(\mathbf{r}_1, \mathbf{r}_2) = E\psi(\mathbf{r}_1, \mathbf{r}_2)$$

$$\left(-\frac{\hbar^2}{2m_e} \nabla_e^2 - \frac{\hbar^2}{2m_N} \nabla_N^2 + V(|\mathbf{r}_e - \mathbf{r}_N|) \right) \psi(\mathbf{r}_e, \mathbf{r}_N) = E\psi(\mathbf{r}_e, \mathbf{r}_N)$$

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$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right) \psi(\mathbf{r}) = E\psi(\mathbf{r})$$


Dirac equation:

$$H_D = \alpha \mathbf{p} + \beta m + V$$


Relativistic recoil: the first order in m/M

PHYSICAL REVIEW LETTERS **130**, 053002 (2023)

QED Theory of the Nuclear Recoil with Finite Size

Krzysztof Pachucki 

Faculty of Physics, University of Warsaw, Pasteura 5, 02-093 Warsaw, Poland

Vladimir A. Yerokhin 

Peter the Great St. Petersburg Polytechnic University, Polytekhnicheskaya 29, 195251 St. Petersburg, Russia

$$\Delta E_{\text{rec,LO}} = \langle a | \frac{\mathbf{p}^2}{2m} | a \rangle$$

$$\Delta E_{\text{rec,H0}} = \text{some other integral}$$

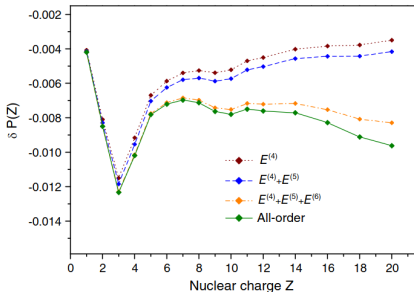


FIG. 1. Finite-size nuclear recoil correction for the $1s$ state of H-like ions, in terms of function $\delta P = E_{\text{recfs}}/[(m^2/M)(Z\alpha)^5/\pi]$.

Polarization or polarizability?

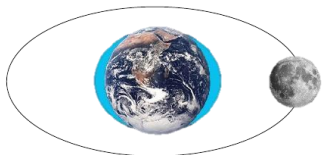


Image source: www.universetoday.com

- nucleus plus muon together
- “measure of entanglement”
- sometimes can be separated (approximately)

Nuclear polarization \approx Nuclear polarizability \times muonic factor

Nuclear polarization on a nutshell

$$V_{\text{Coul}}(r) = -\frac{\alpha Z}{r}$$

Longitudinal (Coulomb) part

$$V_{\text{ext}}(r) = -\alpha \int \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'$$

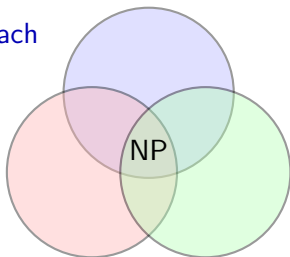
$$\Delta E_{\text{NP}} = \sum_{nN} \frac{|\langle aA | \delta V | nN \rangle|^2}{E_{aA} - E_{nN}}$$

$$V_{\text{NP}}(r) = -\alpha \sum_Z \frac{1}{|\mathbf{r} - \mathbf{r}_{N_i}|}$$

Transverse part: $H = H_N + \alpha(\mathbf{p} - e\mathbf{A}(\mathbf{r}, \mathbf{r}_{N_i})) + \beta m_\mu + V(\mathbf{r}, \mathbf{r}_{N_i})$

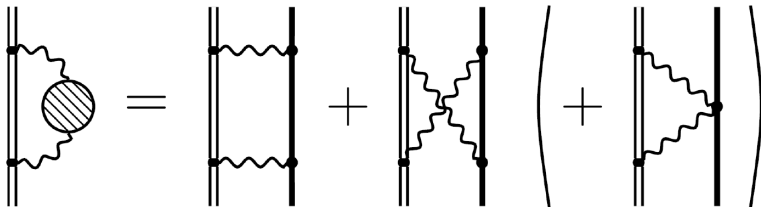
QFT approach

precise
muonic
description



state-of-art
nuclear input

Nuclear polarization: in scary formulas



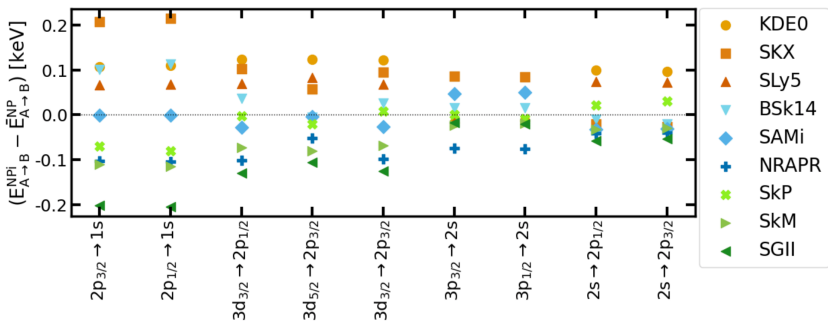
$$\Delta E_{\text{NP}}^{\text{L}} = -i(4\pi\alpha)^2 \sum_{i'I'} \iint \frac{d\mathbf{q} d\mathbf{q}'}{(2\pi)^6} \int \frac{d\omega}{2\pi} \frac{D_{\mu\xi}(\omega, \mathbf{q}) D_{\zeta\nu}(\omega, \mathbf{q}') \langle iI | j_m^\mu(-\mathbf{q}) J_N^\xi(\mathbf{q}) | i'I' \rangle \langle i'I' | J_N^\zeta(-\mathbf{q}') j_m^\nu(\mathbf{q}') | iI \rangle}{(\omega + \omega_m - iE_{i'}\epsilon)(\omega - \omega_N + i\epsilon)},$$

$$\Delta E_{\text{NP}}^{\text{X}} = +i(4\pi\alpha)^2 \sum_{i'I'} \iint \frac{d\mathbf{q} d\mathbf{q}'}{(2\pi)^6} \int \frac{d\omega}{2\pi} \frac{D_{\mu\xi}(\omega, \mathbf{q}) D_{\zeta\nu}(\omega, \mathbf{q}') \langle iI' | j_m^\mu(-\mathbf{q}) J_N^\xi(\mathbf{q}) | i'I \rangle \langle i'I | J_N^\zeta(-\mathbf{q}') j_m^\nu(\mathbf{q}') | iI' \rangle}{(\omega + \omega_m - iE_{i'}\epsilon)(\omega + \omega_N - i\epsilon)},$$

$$\Delta E_{\text{NP}}^{\text{SG}} = -i(4\pi\alpha)^2 \sum_{i'} \iint \frac{d\mathbf{q} d\mathbf{q}'}{(2\pi)^6} \int \frac{d\omega}{2\pi} \frac{D_{\mu\xi}(\omega, \mathbf{q}) \delta^{\xi\zeta} D_{\zeta\nu}(\omega, \mathbf{q}') \langle i | j_m^\mu(-\mathbf{q}) | i' \rangle \langle i' | j_m^\nu(\mathbf{q}') | i \rangle \langle I | \rho_N(\mathbf{q} - \mathbf{q}') | I \rangle}{(\omega + \omega_m - iE_{i'}\epsilon) m_p}.$$

summations over entire muonic (i') and nuclear (I') spectra

Nuclear polarization: results

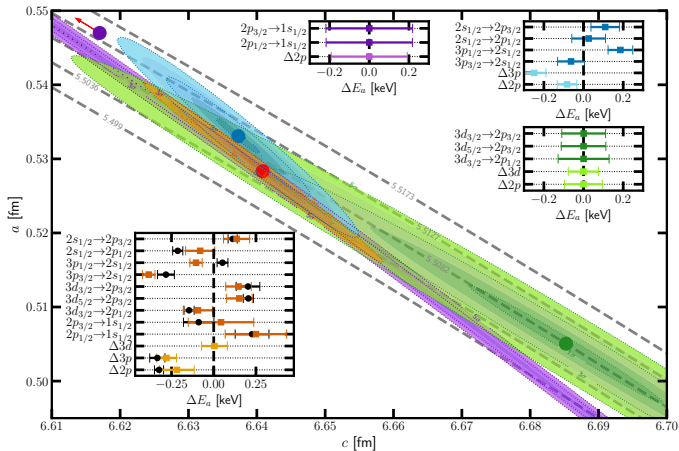


Experimental data

| Set | Transition | Energy (keV) |
|---------|-----------------------|-----------------|
| Run I | $2p_{3/2} - 1s_{1/2}$ | 5962.854(90) |
| | $2p_{1/2} - 1s_{1/2}$ | 5778.058(100) |
| Run II | $3d_{3/2} - 2p_{1/2}$ | 2642.332(30) |
| | $3d_{5/2} - 2p_{3/2}$ | 2500.590(30) |
| | $3d_{3/2} - 2p_{3/2}$ | 2457.569(70) |
| Run III | $3p_{3/2} - 2s_{1/2}$ | 1507.754(50) |
| | $3p_{1/2} - 2s_{1/2}$ | 1460.558(32) |
| | $2s_{1/2} - 2p_{1/2}$ | 1215.330(30) |
| | $2s_{1/2} - 2p_{3/2}$ | 1030.543(27) |

Bergem *et al.*, PRC **37**, 2821 (1988)

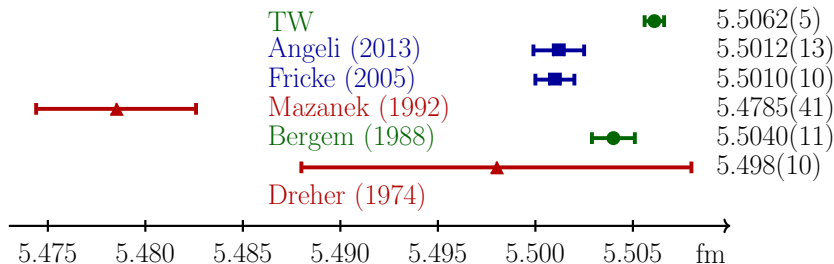
Results for ^{208}Pb



Purple: $2p_{1,3/2} - 1s_{1/2}$, green: $3d_{3,5/2} - 2p_{1,3/2}$,
 blue: $3p_{1,3/2} - 2s_{1/2}$ and $2s_{1/2} - 2p_{3,1/2}$, red: total

Sun *et al.*, *PRL* 135, 163002 (2025)

Compared with old data



- Mild tension with Bergem
- 3-4 σ with Angeli, Fricke
- $\chi^2/\text{DoF} = 9.5$: factor of 20 improvement
- Evidence against the existence of the anomaly

Normalized fit quality $\chi^2/\text{DoF} = 1$: value 5.5062(17) fm
reduced model-dependence

Sun *et al.*, *PRL* 135, 163002 (2025)

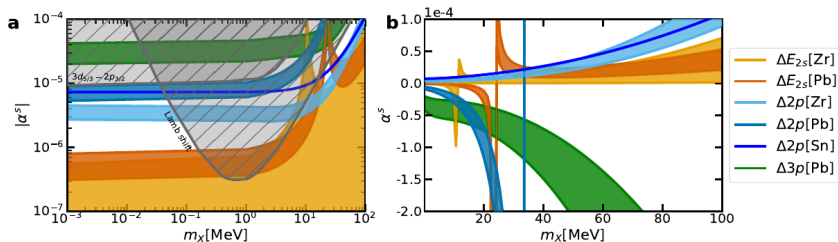
Muonic atoms anomalies: H vs Pb



BSM as (not) a solution of the anomaly

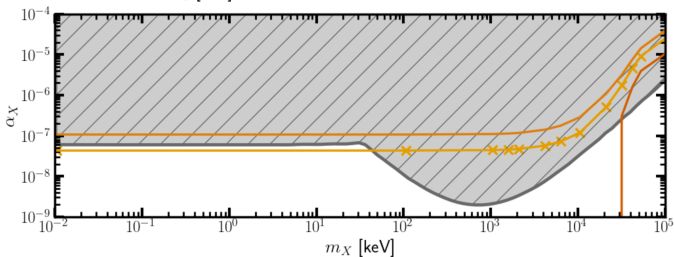
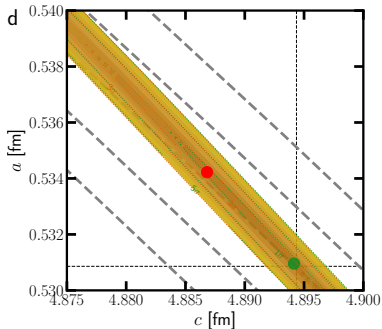
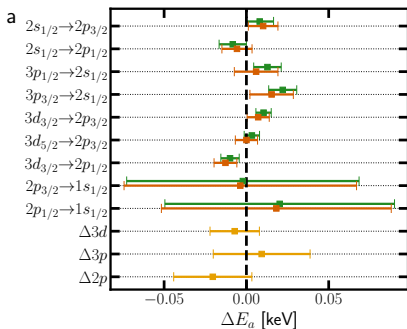
Scalar/vector new boson

$$V_{\text{BSM}} = \alpha_{\text{BSM}} \hbar c \frac{\exp[-m_{\text{BSM}} cr / \hbar]}{r} \gamma_0$$



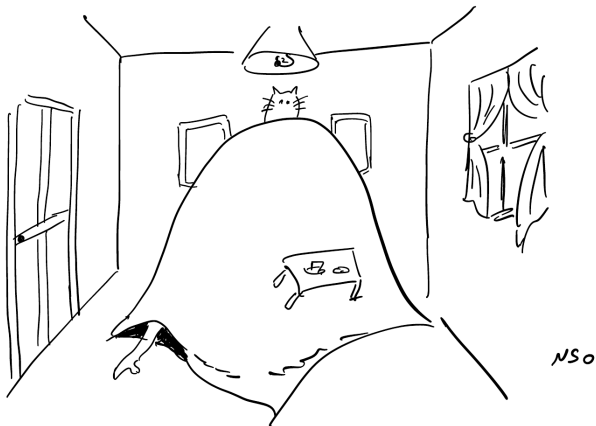
Beyer *et al.*, *Phys. Lett. B* **854**, 138746 (2024)

Self-consistent probe of BSM PRELIMINARY



Outlook

- “heavy electron” QED
- Anomaly is almost resolved
- New recommended rms for ^{208}Pb : 5.5062(17) fm
- A step forward systematic reevaluation of all rms radii



Acknowledgments

Experiments and experimentalists



My hard-working team



and you for your attention

