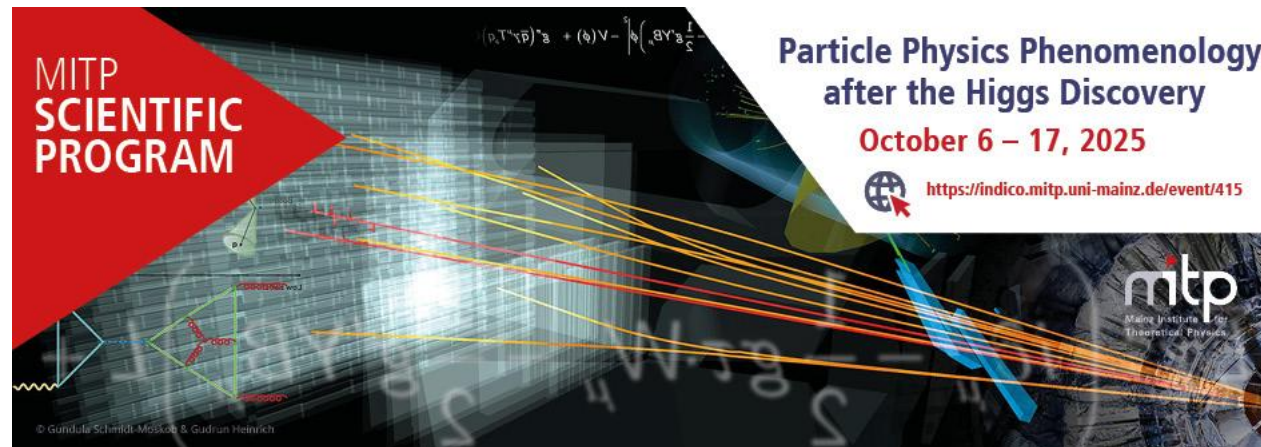


Why is SMEFT/HEFT the right approach to BSM physics and what can we expect from it?

S. Dawson, BNL

Oct. 6, 2025

MITP

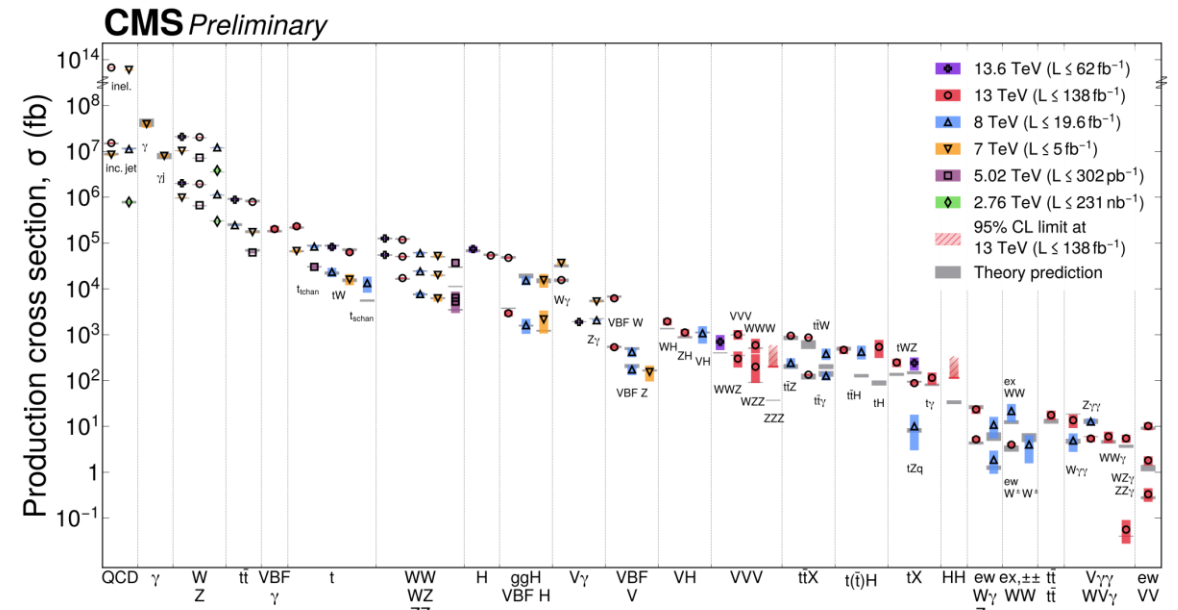
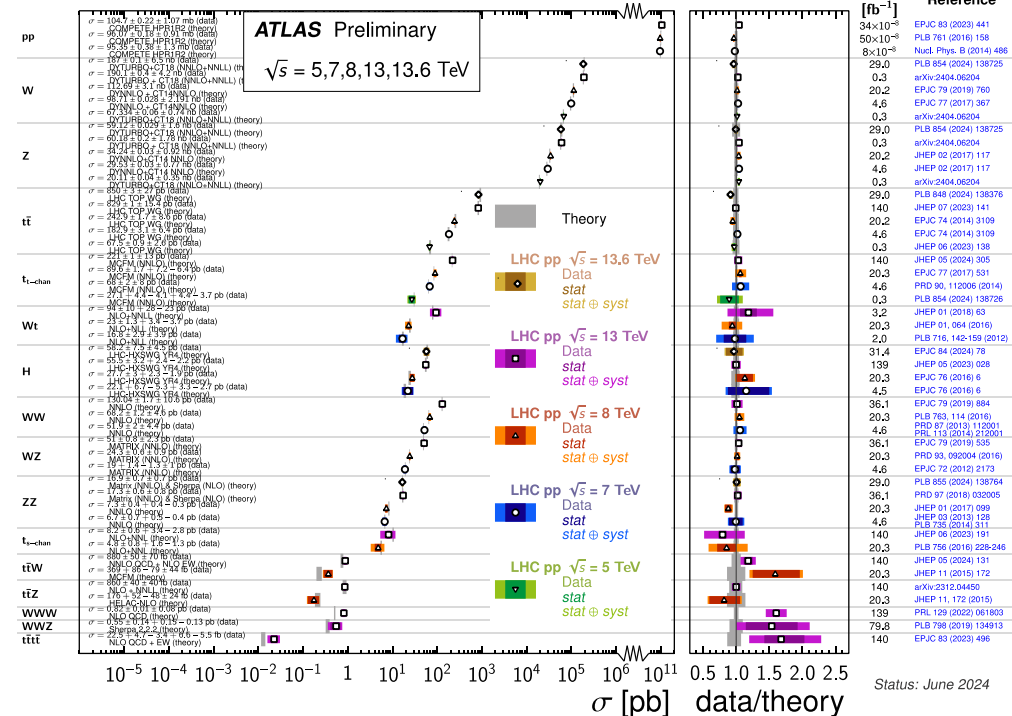


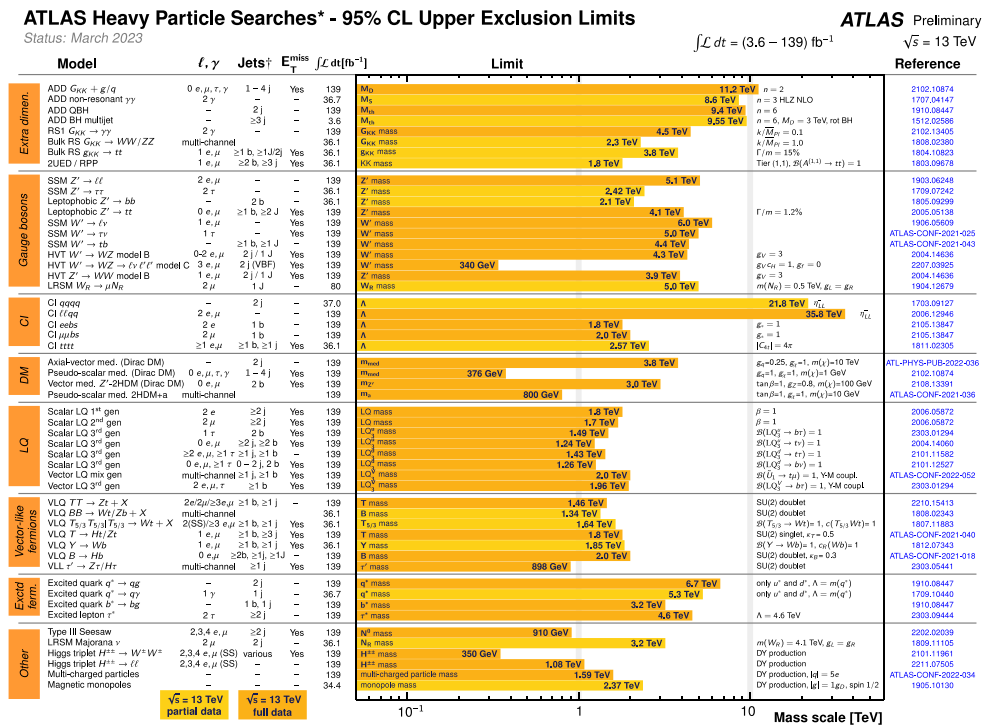
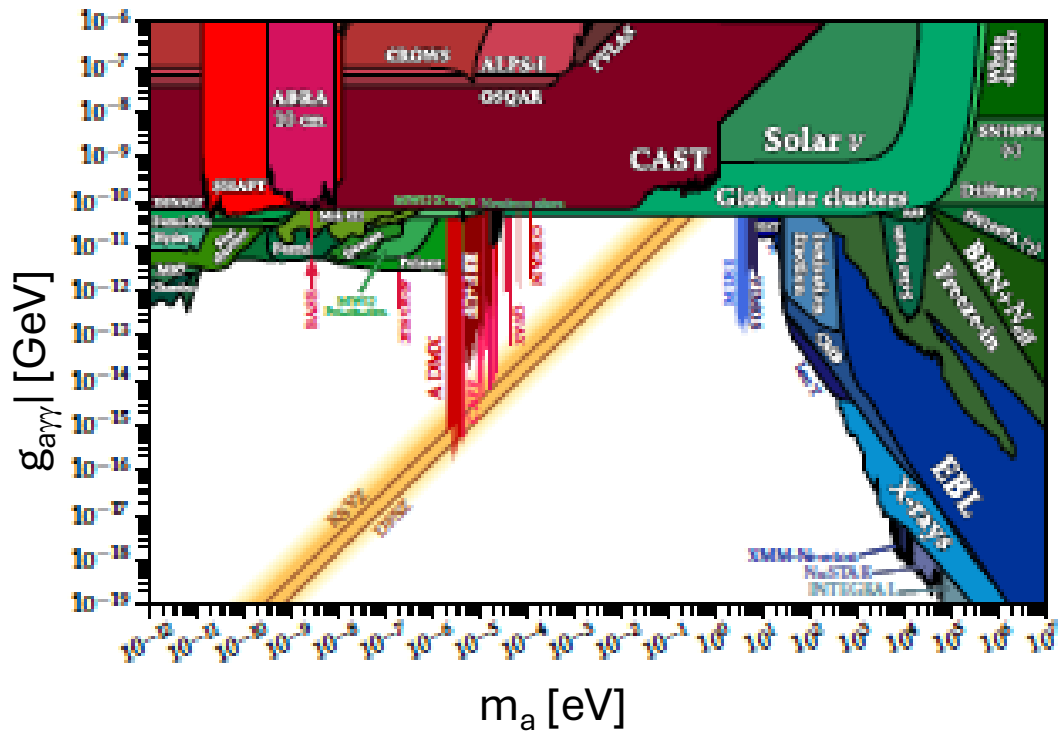
Questions for discussion at this workshop

- What has been achieved in the field after the Higgs discovery?
 - What are the perspectives of the field?
 - What needs to be developed to create further advances
 - ... regarding HL-LHC?
 - ... regarding future collider projects?
-
- Questions posed by the organizers....
 - Feel free to stop me for discussion

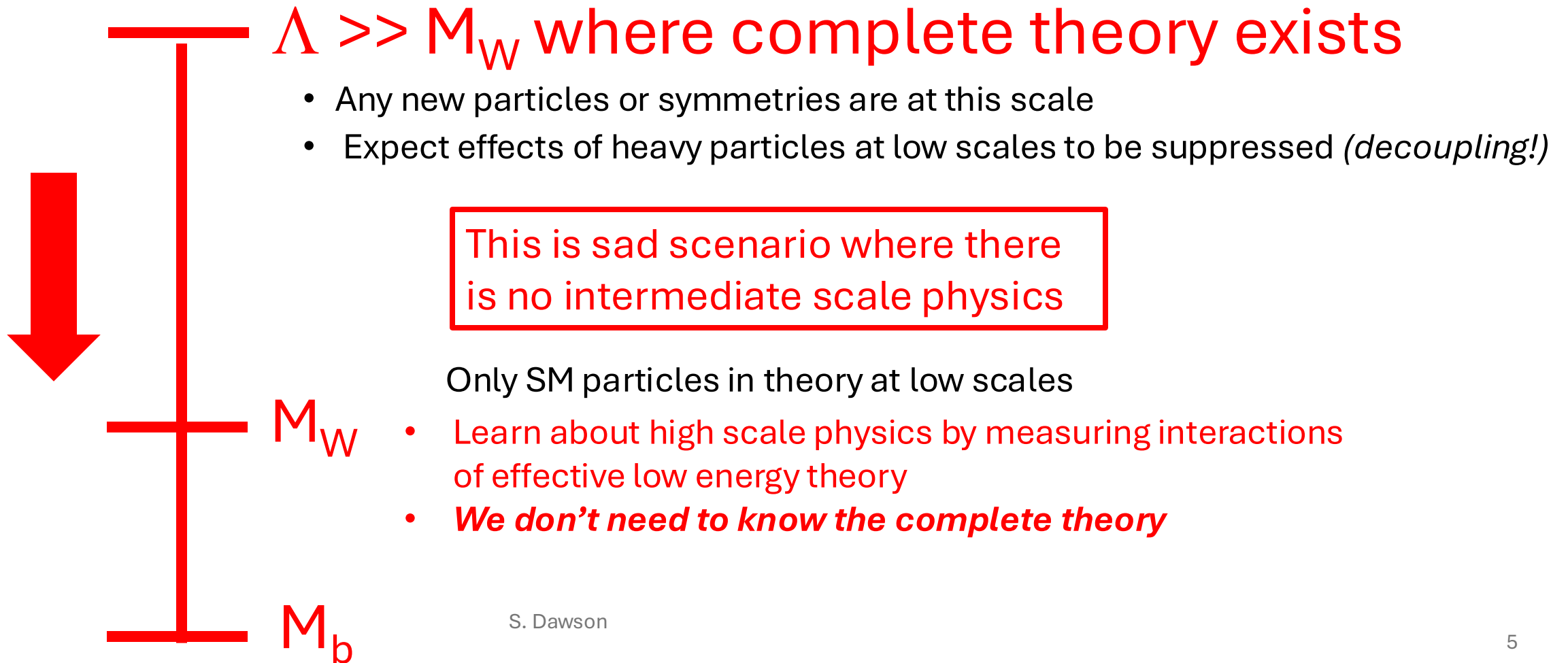
Where are we now?

- LHC results are **very, very similar** to our best SM predictions over many orders of magnitude and for many processes





Not obvious how to search for new physics: EFTs are one technique



SMEFT: SM Effective Field theory

- **Assumptions:** New physics decouples $\Lambda \gg v, E$
- At the weak scale: SM $SU(3) \times SU(2) \times U(1)$ symmetry and SM particles only
- New physics described by

$$L_{SMEFT} = L_{SM} + \frac{L_5}{\Lambda} + \frac{L_6}{\Lambda^2} + \frac{L_7}{\Lambda^3} + \frac{L_8}{\Lambda^4}$$
$$L_n = \sum_i C_i^n O_i^n$$

Assume Higgs is in
an $SU(2)$ doublet

This is the big
assumption

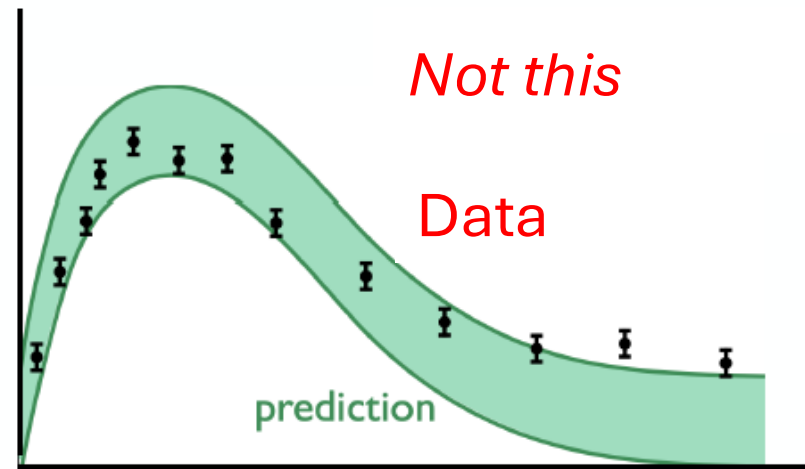
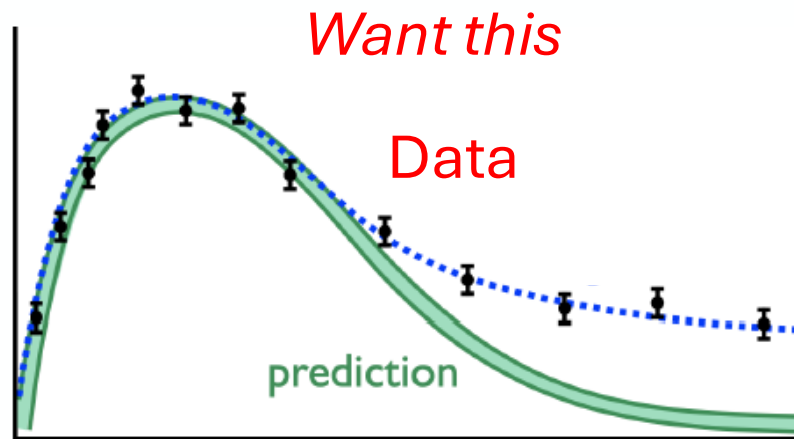
- **New physics contributions contained in coefficients C**
- Operators form a complete basis (not unique)
- L_5 and L_7 are lepton number violating

*Warsaw basis commonly used

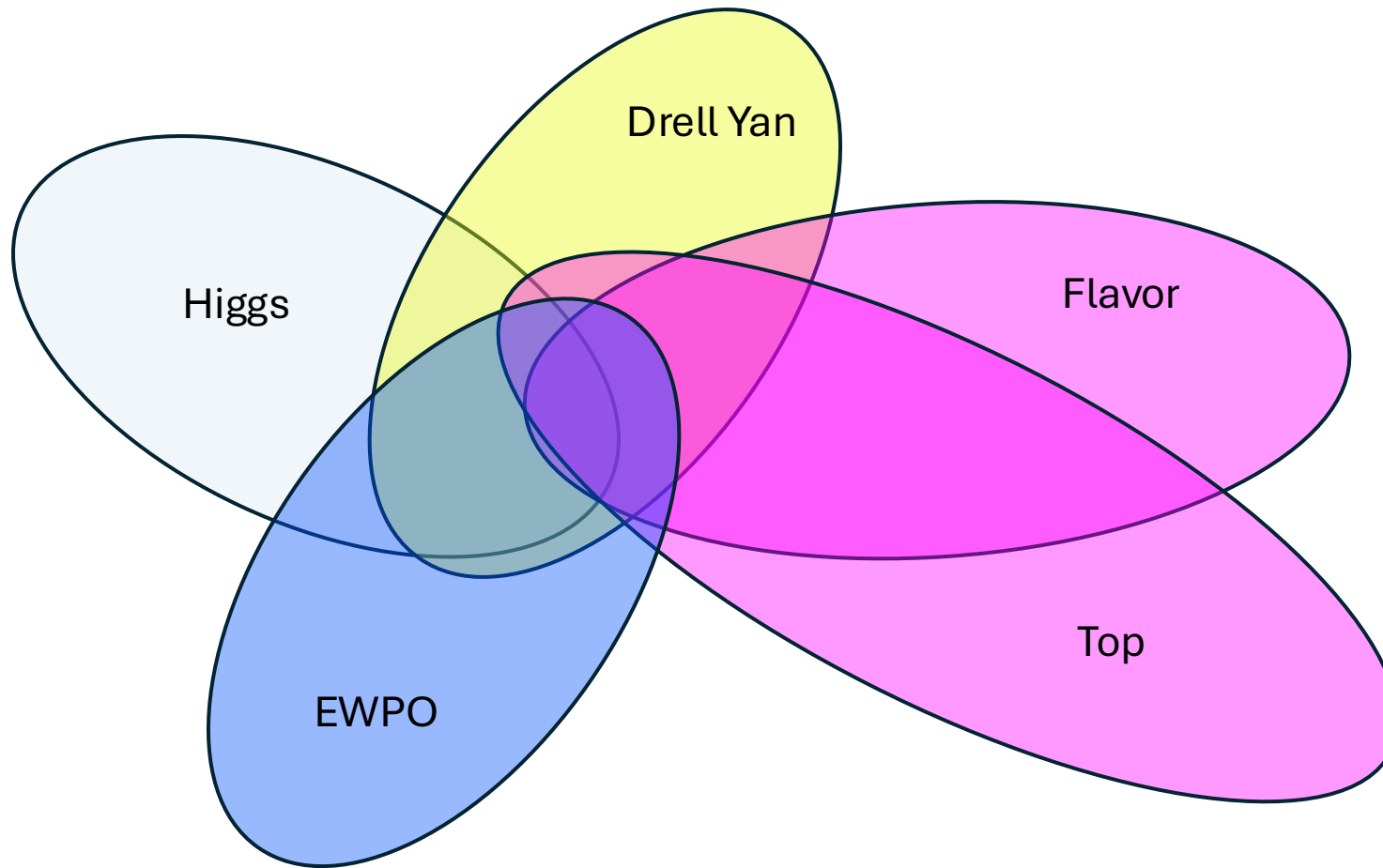
What needs to be developed to create further advances?

$$\text{Experiment} = \text{Theory}_{\text{SM}} + \sum \frac{x_i C_i^6}{\Lambda^2} + \dots$$

Precise experimental measurements Precise SM calculations Precise SMEFT calculations



Power of SMEFT is information from many observables



- SMEFT operators contribute to different types of observables
- Leads to concept of global fits

What needs to be developed to create further advances?

- Understanding of theory uncertainties on SMEFT calculations
- Can new techniques (AI/ML) improve global fits?
- Precise SMEFT calculations beyond LO
- Understanding of interpretation of SMEFT calculations
 - Do they really tell us about UV models?

Precise SMEFT calculations

- Compute an amplitude at tree level

$$A \sim A_{SM} + \frac{A_6}{\Lambda^2} + \frac{A_8}{\Lambda^4} + \dots$$

- Various possibilities for defining observables

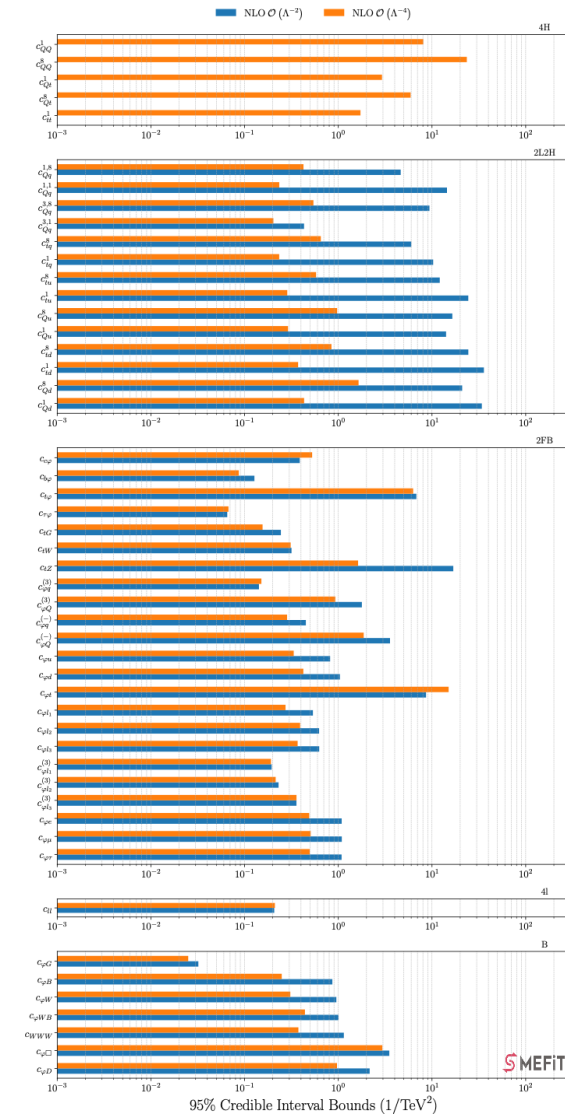
$$O \sim O_{SM} + \frac{A_6 A_{SM}}{\Lambda^2} + \left(\frac{A_6}{\Lambda^2} \right)^2 + \frac{A_{SM} A_8}{\Lambda^4} + \dots$$

Linear Quadratic

- Where to truncate?

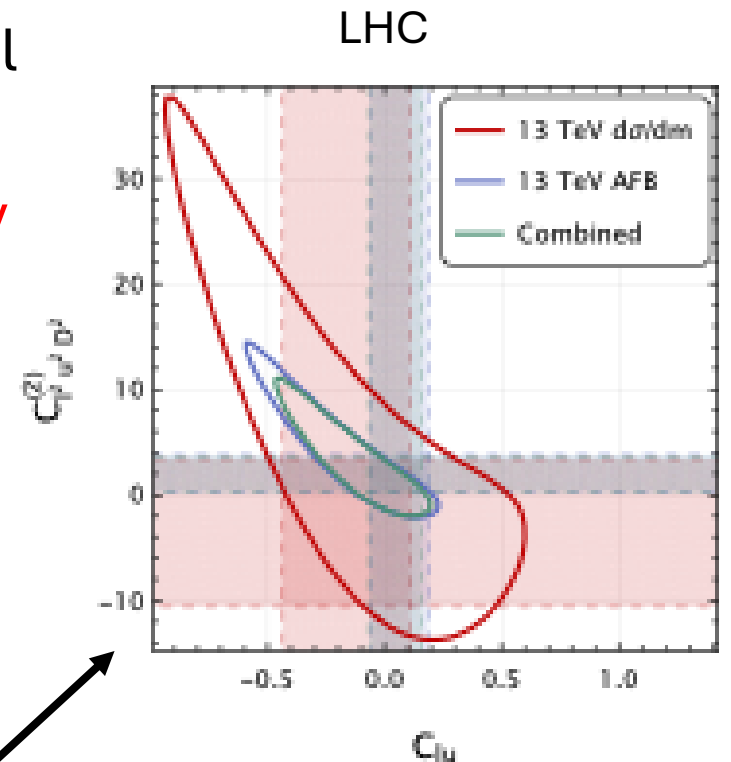
When is expansion valid?

- Expansion is both in $1/\Lambda^2$ and in loops
- Compare tree level $1/\Lambda^2$ with $1/\Lambda^4$ fits using just dimension-6 operators and including QCD@NLO
- In general, $1/\Lambda^4$ contributions change limits significantly
- (less so for observables contributing to EWPO, which are well measured)
- Is the difference useful as a handle on the theoretical uncertainty?



When does dimension-8 matter?

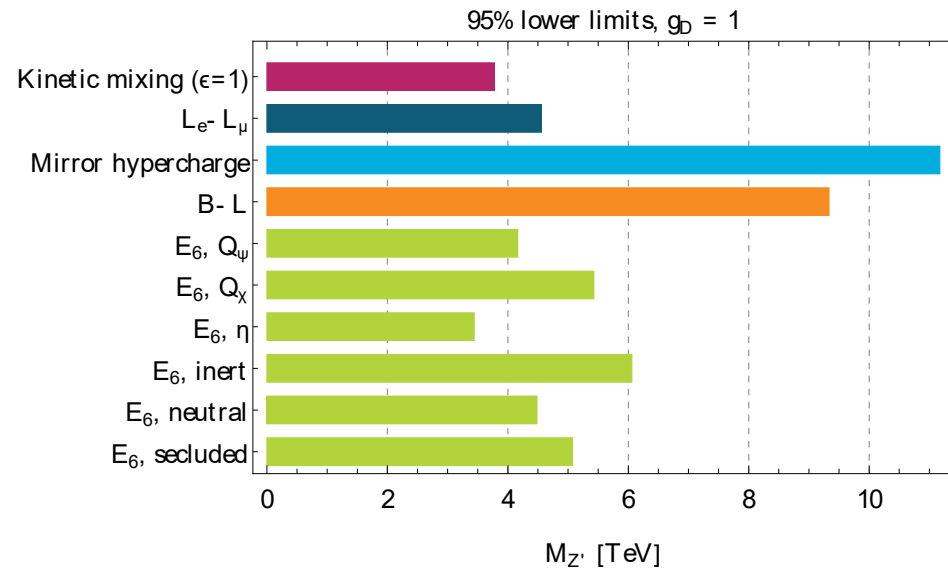
- Dimension-8 operators are $1/\Lambda^4$
- Some studies on a case- by- case basis, but no general conclusions
- Model independent approach: **Drell Yan as case study**
 - Precise SM and SMEFT calculations exist
 - 7 dim-6, 14 dim-8 operators. **Potentially a general study could be possible?**
 - Example: Left-handed lepton coupling to right-handed quark (4-quark interaction)



dim-8 changes conclusions

When does dimension-8 matter?

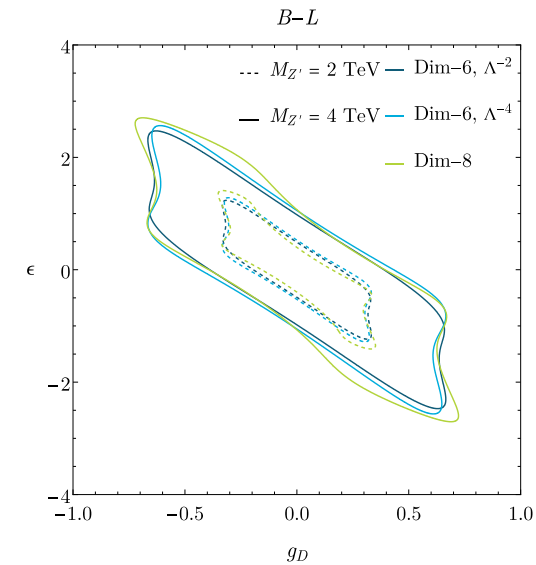
- Model dependent studies: **Z' models**



Match models to dim-6 and dim-8 SMEFT operators

$$\frac{C}{\Lambda^2} \sim \frac{g_D(..)}{M_{Z'}^2}$$

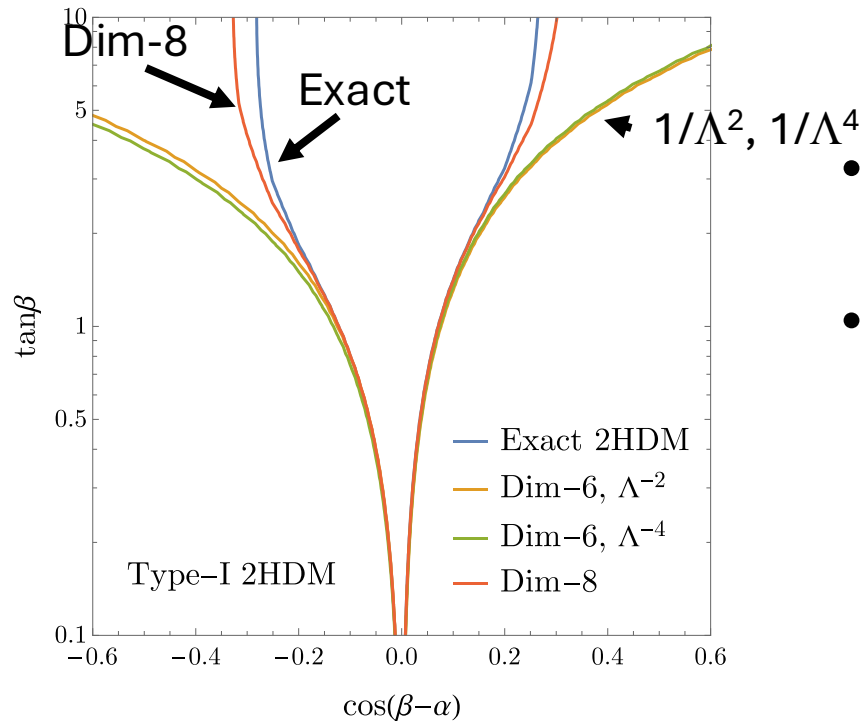
[2404.01375](#)



Dimension-8 effects small: Fit to EWPO and DY m_{ll} at LHC

When does dimension-8 matter?

- Model dependent studies: **2HDM as case study**



- VVH interactions don't occur at dimension-6 in Type-I 2HDM
- Need dimension-8 to get physics right

More case studies of the impact of dim-8 would be useful

[2205.01561](#)

Going beyond Tree Level

$$\mathcal{A} \sim a_{sm}^0 + \frac{a_{sm}^1}{16\pi^2} + \sum_i \frac{C_i^6}{\Lambda^2} \left[A_{EFT,i}^0 + \frac{B_{EFT,i}^1}{16\pi^2} + \frac{D_{EFT,i}^1}{16\pi^2} \log\left(\frac{\Lambda^2}{M_Z^2}\right) \right] + \sum_i \frac{C_i^8}{\Lambda^4} F_{EFT,i}^0 + \dots$$

- NLO QCD is automated, but NLO EW corrections done on a case- by -case basis
- At NLO, new operators contribute
 - In general, effect of more operators is to weaken many limits
- Logarithms come for free from RGE. Do they dominate?

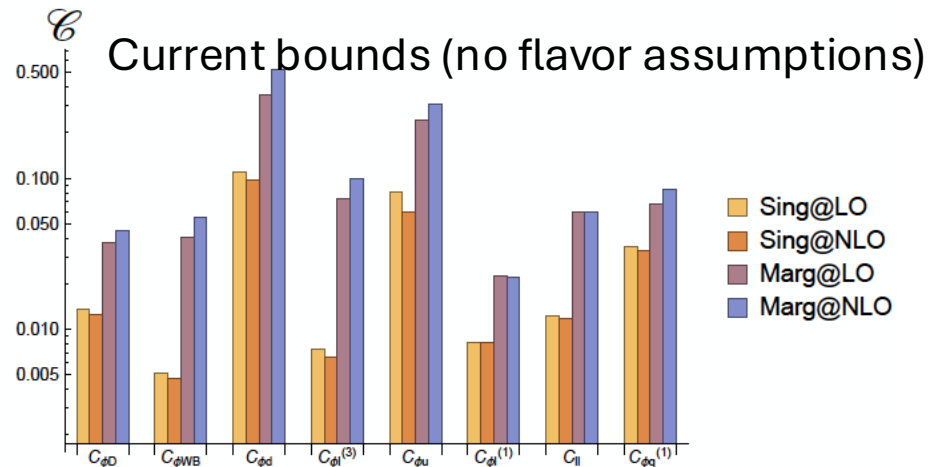
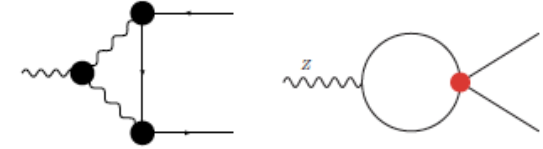
$$C_i(\mu) = \gamma_{ij} C_j \log\left(\frac{\mu}{\Lambda}\right)$$

Not diagonal

We'd really like this to be the case!

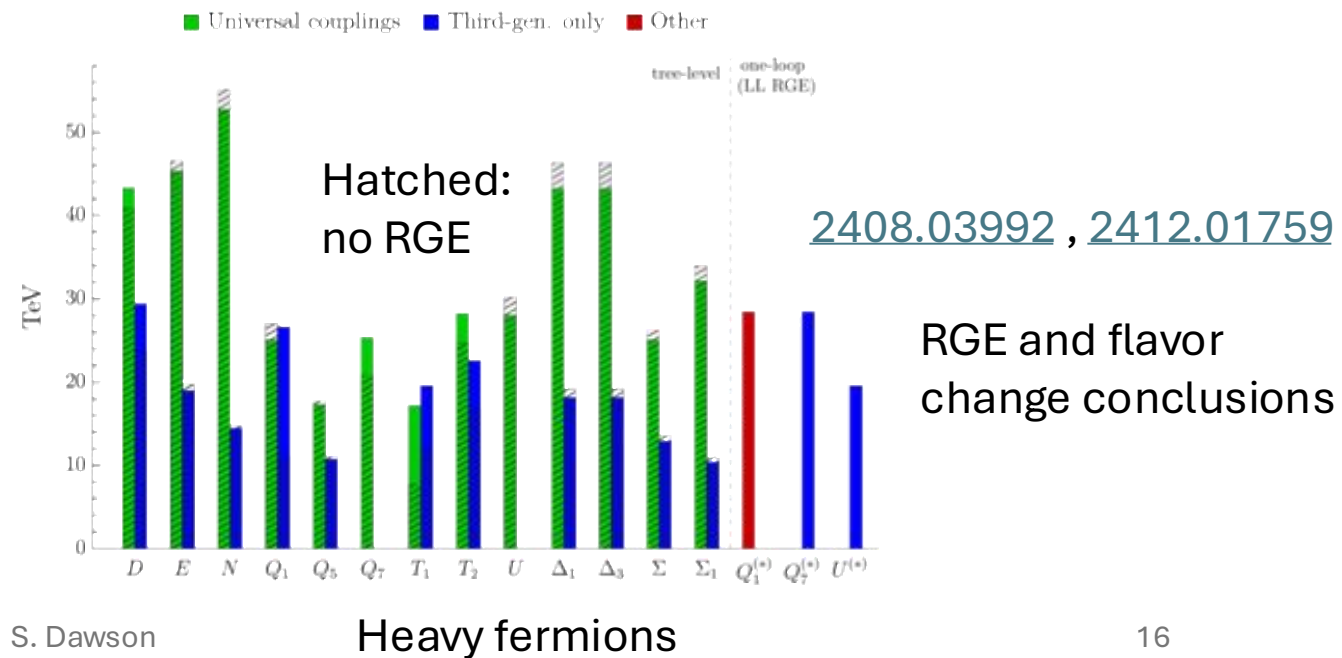
Example of loop effects

- **Z pole observables** (without flavor for now)
 - At tree level, dependence on 10 operators (2 blind directions)
 - At NLO, dependence on 32 operators (new contributions especially from 4-fermion operators)



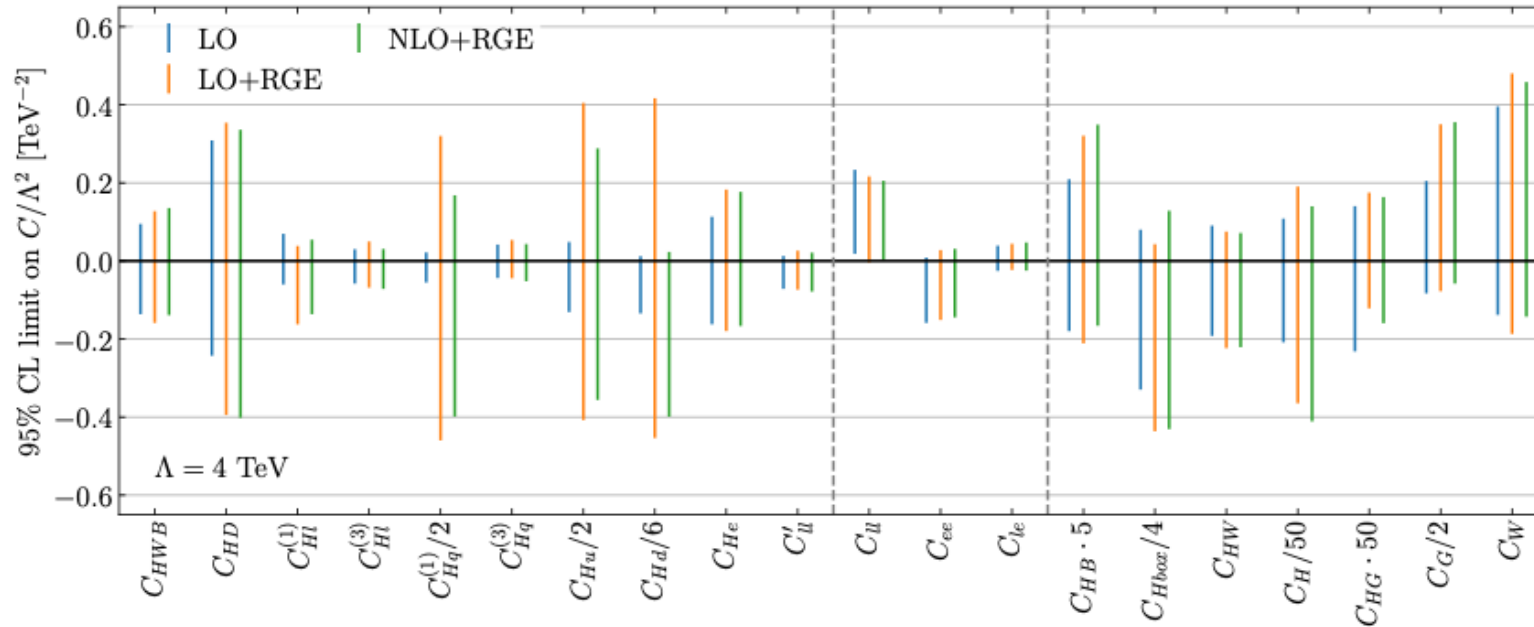
Public results for complete NLO
EW Z pole SMEFT observables,
[2304.00029](#) , [2503.07724](#)

Tera-Z: Single parameter models



Including RGE in global fits is state of the art

- RGE effects can be large
- RGE effects generate dependence on new operators



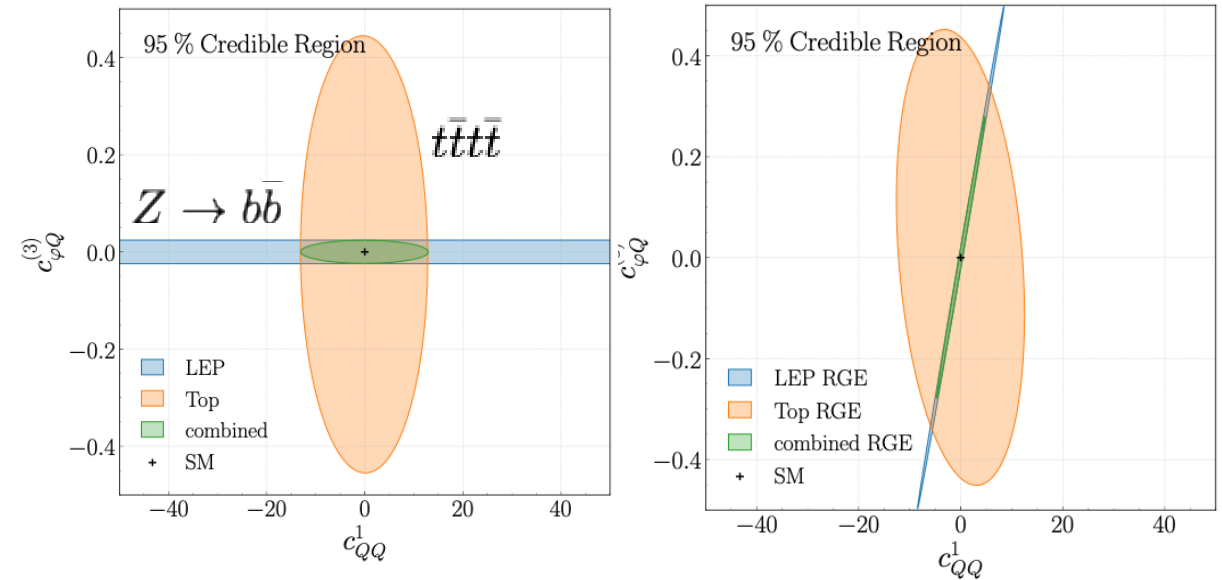
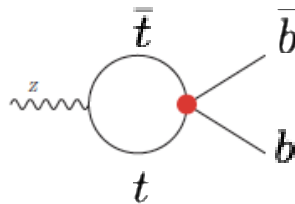
These are
observables
contributing
to EWPOs

[2412.09674](#)

See also, [2502.20453](#) [2507.06191](#)

Including RGE in global fits is state of the art

- RGE has large effects on 4-fermion heavy quark operators
 - They are poorly unconstrained at LHC
 - RGEs mix these operators with operators that contribute to well measured observables at Z pole ($Z \rightarrow b\bar{b}$)

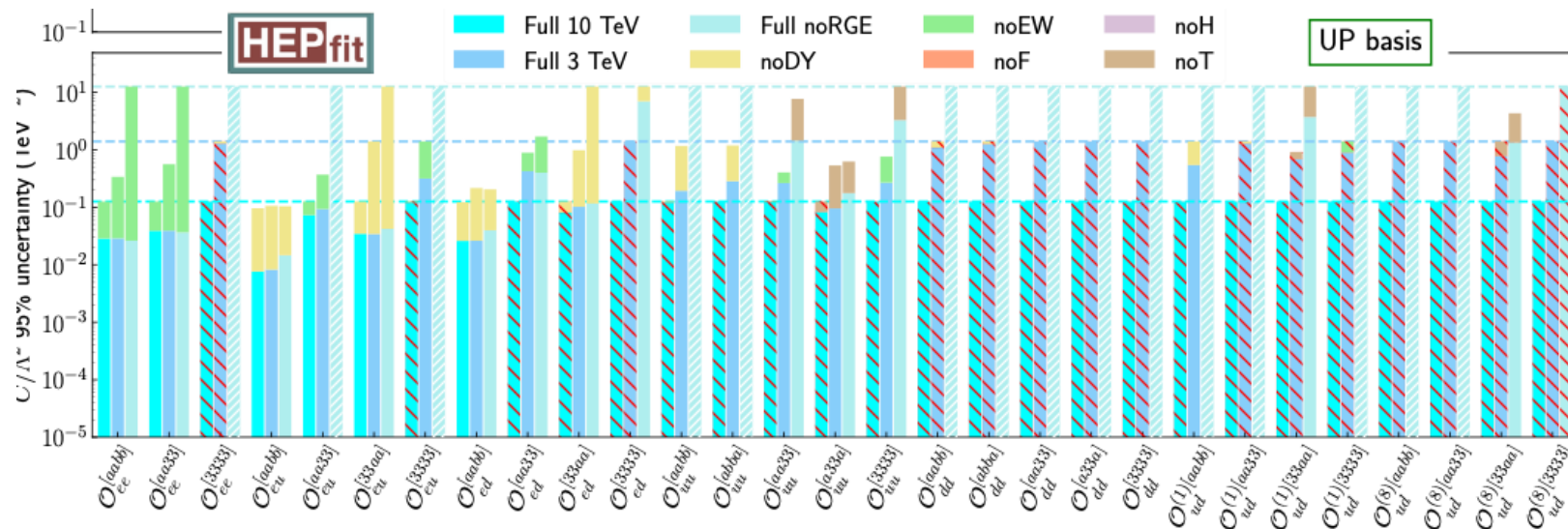


[2502.20453](#) , [2507.01137](#) , [2410.13304](#)

O_{QQ}^1 is 4-quark (LL)(LL) operator with (t,b) doublet
 $O_{\phi Q}^3$ is 2-quark LL operator with 2 fermion fields

Many global fits with different data sets

- Fits to anomalous interactions (Include Drell-Yan, EWPO, Higgs, top, B)
- Top measurements play an important role in constraining effective 4-fermion operators
- RGE effects can be important

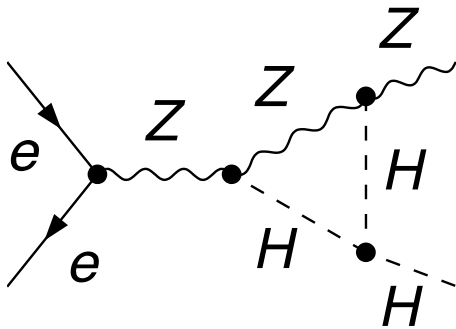


Fit includes NLO QCD, but is tree level electroweak

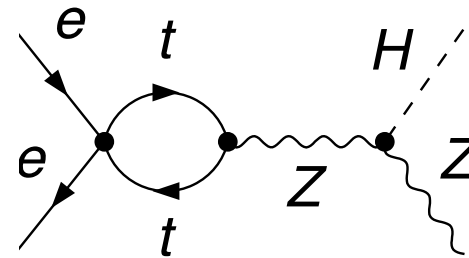
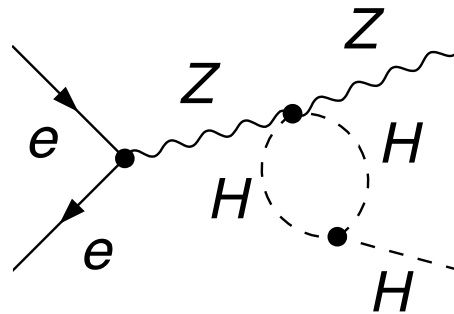
2507.06191

Higgstrahlung at NLO EW SMEFT

- Complete NLO calculation including all dimension-6 operators and no flavor assumptions
 - (~70 SMEFT operators contribute in ~ 35 combinations)
- Sensitive to poorly constrained interactions that first arise at NLO



Higgs tri-linear coupling, C_ϕ



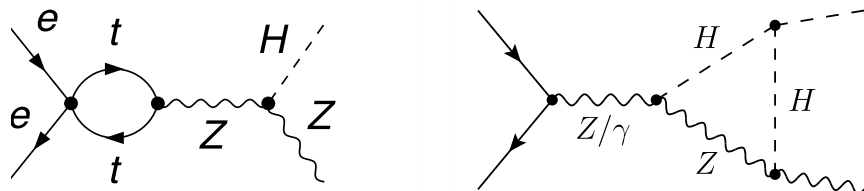
+ many more

4-fermion operators, C_{eu} [1133]

* Complete results at
<https://gitlab.com/smeft/eehz>

Note complementarity with Z-pole results:
[2304.00029](#) , [2201.09887](#) , [2412.14241](#)

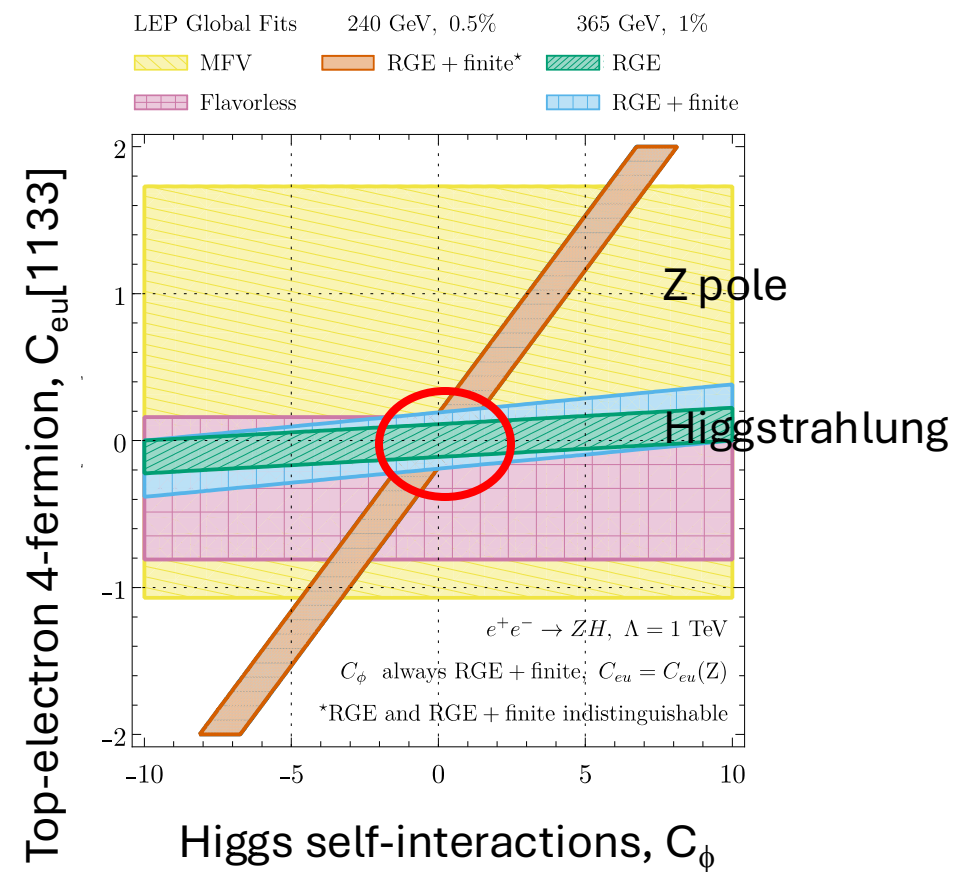
$e^+e^- \rightarrow ZH$ is window to many new interactions



- Effects of different operators is correlated
- Power of measurement at 2 different energies

Note: Z pole limits depend on flavor assumptions

[2406.03557](#)

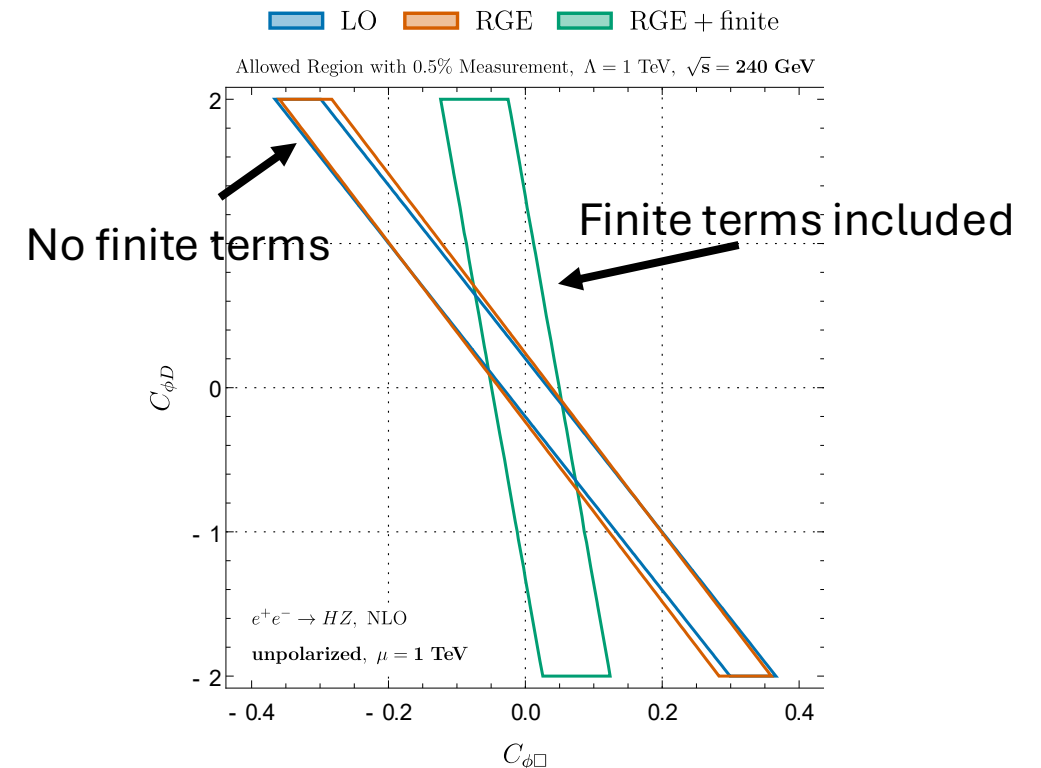


Need running at $\sqrt{s}=365$ GeV to really nail down Higgs tri-linear

NLO EW corrections in dim-6 SMEFT

- EWPOs, Drell-Yan, Higgs decays known at 1-loop EW dim-6 SMEFT
- Can test when RGE logs dominate
 - No general conclusion
- Complete set of 1-loop corrections needed for global fits that is accurate to NLO EW not available

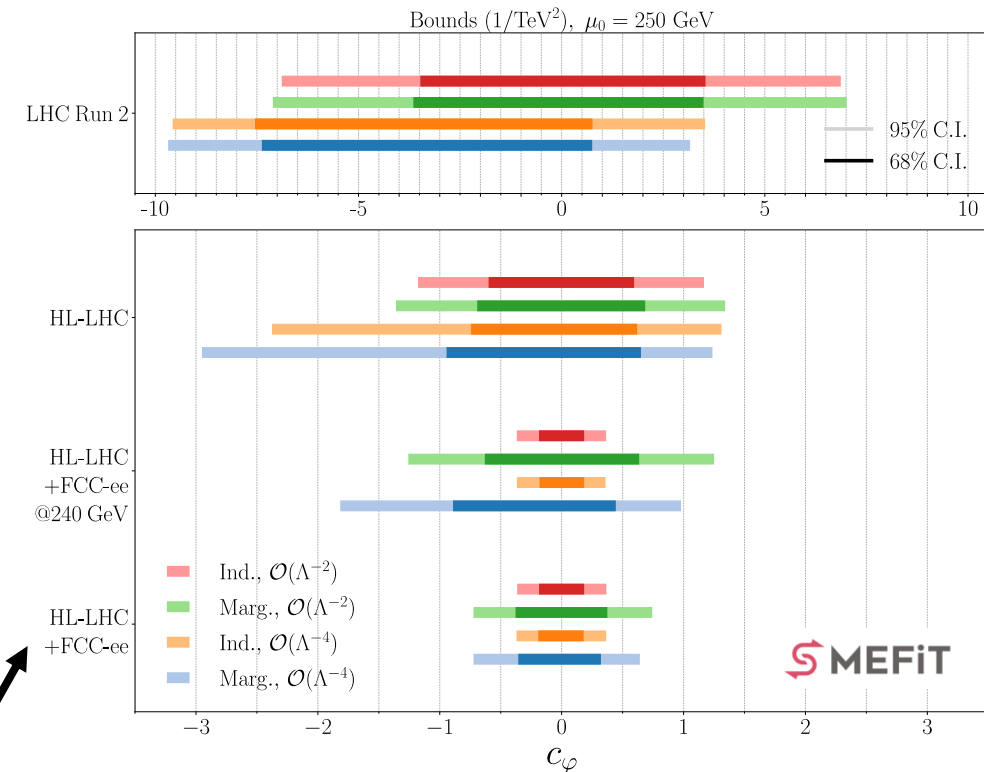
Sensitivity at FCC-ee from $e^+e^- \rightarrow ZH$



[2409.11466](#)

Global fit for Higgs tri-linear

- Include top, H, VV, HH in LHC projections
 - HL-LHC limits largely independent of contamination from other operators, (ie single parameter and marginalized fits very similar)
- Include EW loops in FCC-ee fits (don't have NLO for other pieces)
 - FCC-ee marginalized limits differ from single parameter limits
- Need $\sqrt{s}=365$ GeV @FCC-ee to improve on marginalized HL-LHC limits



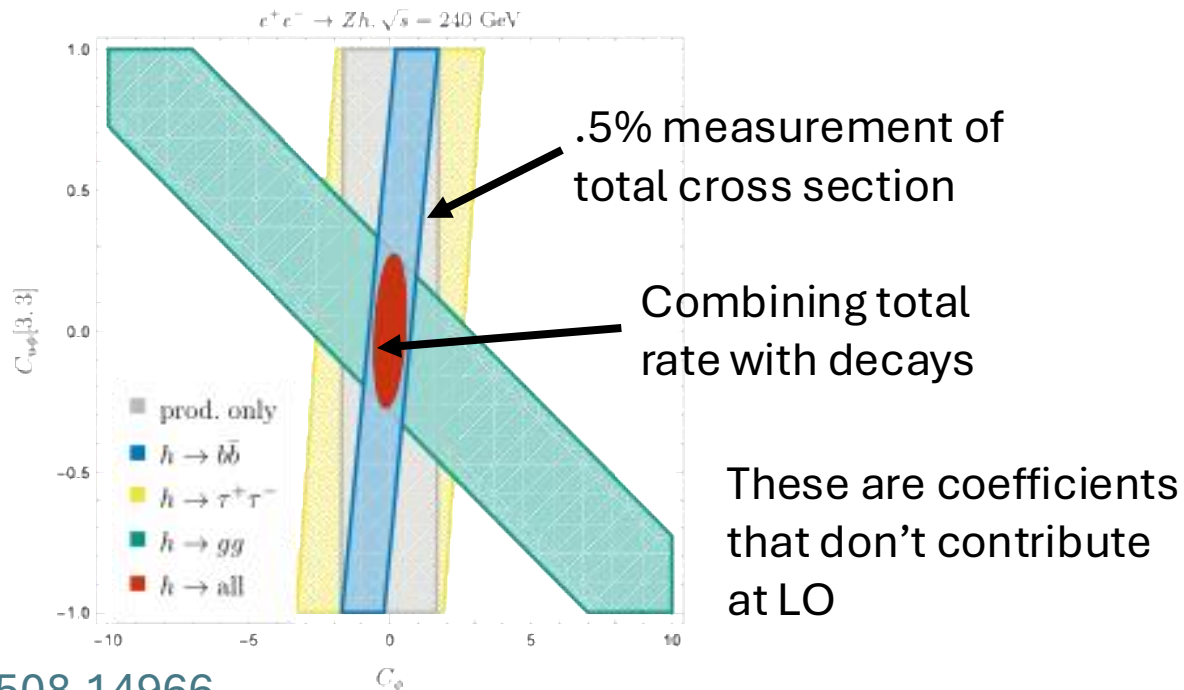
Includes $\sqrt{s}=240$ GeV
+ $\sqrt{s}=365$ GeV running

Higgs tri-linear

[2504.05974](#)

End-to-end NLO EW fit

- Example case: $e^+e^- \rightarrow ZH, H \rightarrow XX$
- Use NWA and include full NLO corrections to $e^+e^- \rightarrow ZH, Z \rightarrow ff, H \rightarrow XX$
- Significantly more information when decays are included



FCC-ee, Tera-Z

	$\sqrt{s} = 240 \text{ GeV}$	$\sqrt{s} = 365 \text{ GeV}$
$b\bar{b}$	0.21	0.38
$c\bar{c}$	1.6	2.9
$s\bar{s}$	120	350
gg	0.8	2.1
$\tau^+\tau^-$	0.58	1.2
$\mu^+\mu^-$	11	25
WW^*	0.8	1.8
ZZ^*	2.5	8.3
$\gamma\gamma$	3.6	13
$Z\gamma$	11.8	22

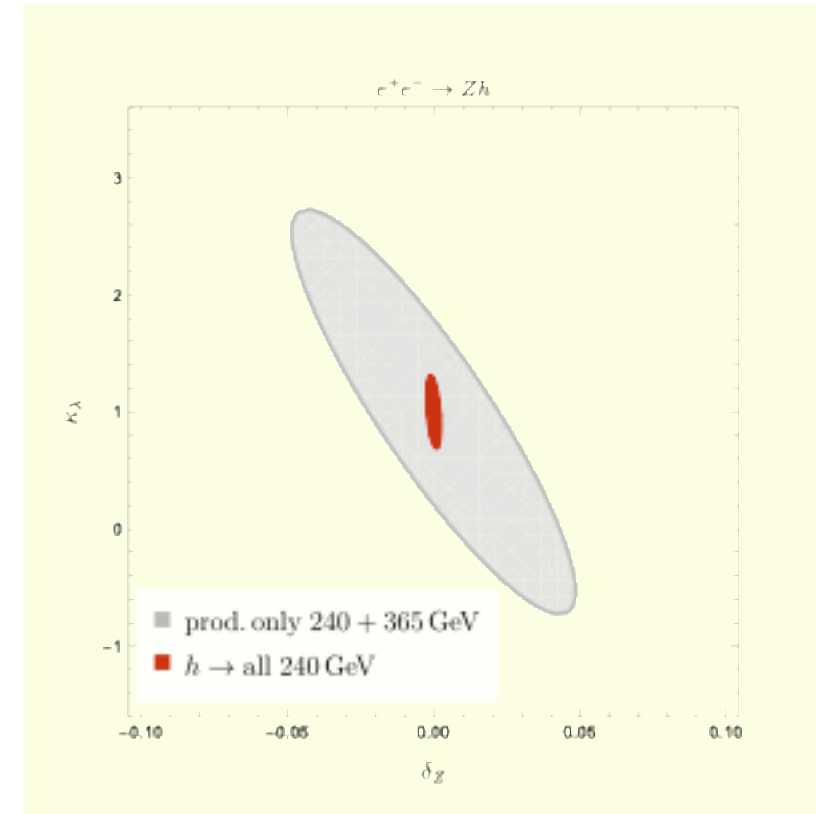
End-to-end NLO EW fit

- Translate to κ formalism

$$\delta_Z = \frac{1}{4} \frac{v^2}{\Lambda^2} \left(C_{\phi D} + 4C_{\phi\Box} \right)$$

$$\kappa_\lambda = 1 + \frac{v^2}{\Lambda^2} \left(3 \left[\frac{C_{\phi D}}{4} - C_{\phi\Box} \right] - 2 \frac{v^2}{m_H^2} C_\phi \right)$$

- Decays provide new information

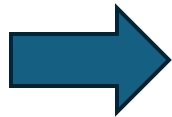


2508.14966

* $C_{\phi\Box} = 0$

Missing pieces for end-to-end NLO EW fit at LHC

- Very few of the production processes are known to NLO EW order
- Coming soon (!): public code NEWISH
 - Differential rates for ALL 2- and 3-body Higgs decays using dimension-6 SMEFT that is accurate to NLO in QCD and EW interactions
 - NWA results for ALL 4-body decays dimension-6 SMEFT that is accurate to NLO in QCD and EW interactions
 - Tree level 4-body dimension-6 SMEFT Higgs decays



Towards a global fit that is accurate to NLO QCD and EW

What about flavor?

- Much of the complexity of SMEFT studies comes from 4-fermion operators
- SM has $U(3)^5$ flavor symmetry:

$$U(3)_q \times U(3)_u \times U(3)_d \times U(3)_l \times U(3)_e$$

- Top Yukawa breaks this to:

$$U(2)^2 \times U(1) \times U(3)_d \times U(3)_l \times U(3)_e$$

- Various assumptions in the literature
 - New physics is flavor independent
 - New physics only couples to 3rd generation
 - New physics obeys a $U(3)^5$ or $U(2)^5$ symmetry
 - New physics only arises from Yukawa interactions
 -

From a global fitting point of view, these assumptions matter

Flavor assumptions reduce possibilities

Operators that contribute to EWPO at NLO

Operator	$U(3)^5$	MFV	$U(2)^5$	3 rd gen specific	3 rd gen phobic	3 rd gen phobic + $U(2)^5$	Flavorless
Class A	7	12	16	9	14	7	9
Class B	11	17	27	5	23	11	6
Class C	11	21	44	11	44	11	11
Total	29	50	87	25	81	29	26

2-fermion →

4-fermion with identical representations →

Remaining 4-fermion →

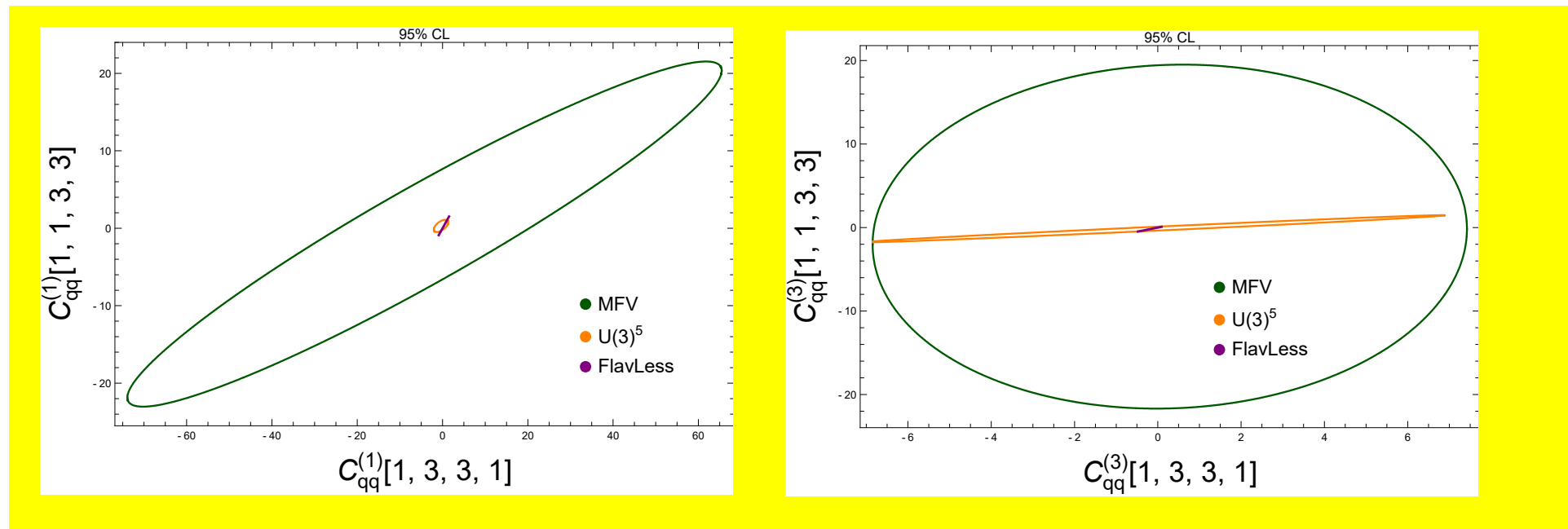
- Compare Z pole global fit results with $U(3)^5$, $U(2)^5$, MFV, only 3rd generation operators, no flavor structure

Much work to be done understanding flavor in SMEFT global fits

* In all generality, 168 operators contribute to EWPO with no flavor assumptions

Flavor matters!

Current Z pole limits on 4-fermion operators

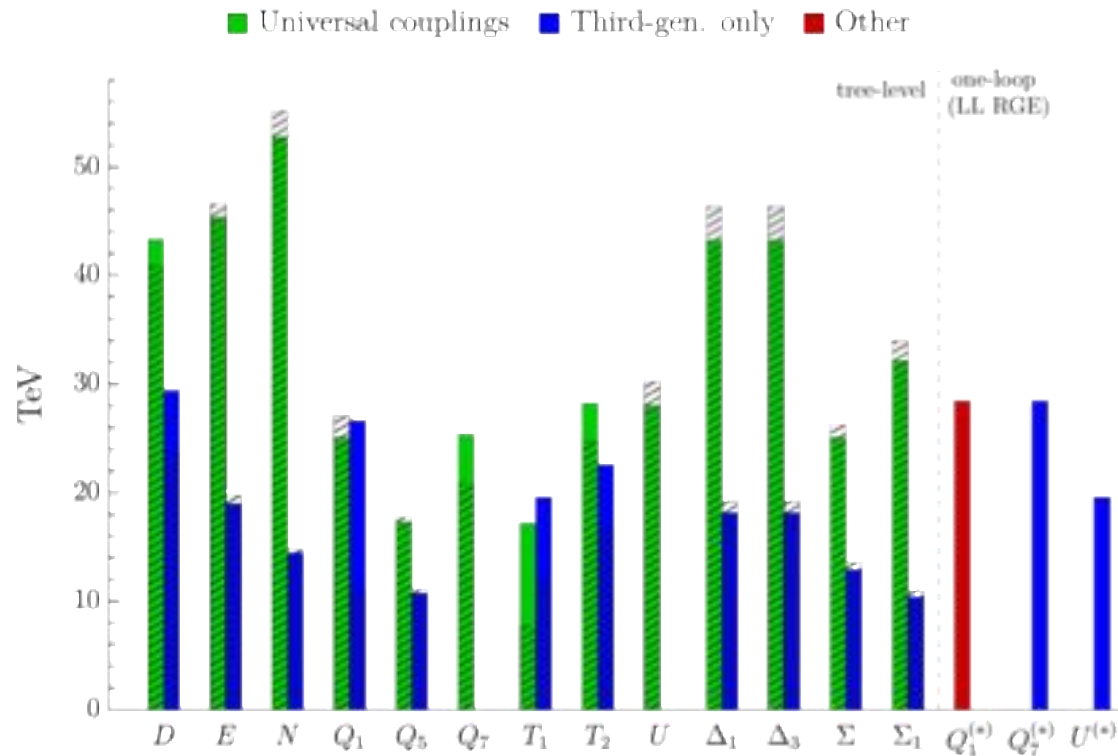


Consider 1 operator type at a time and marginalize over flavor structures not shown

[2304.00029](#)

Flavor matters

Tera-Z sensitivity to heavy fermions



[2408.03992](#)

RGE operator mixing leads to new flavor structures

- Assume all (non-SM) couplings=1, and all dimensionful parameters=M
- This “dictionary” type of analysis is perhaps oversimplified

What needs to be developed to create further advances?

- Understanding of theory uncertainties on SMEFT calculations

- When do we need dimension-8?
- Truncating dimension-6 at $1/\Lambda^2$ vs $1/\Lambda^4$?
- Flavor assumptions—when are they crucial?
- Validity of expansion
- The most straightforward-- including input scheme, scale and parametric uncertainties

Input scheme matters

$h \rightarrow b\bar{b}$		SM	$C_{H\Box}$	C_{HD}	$C_{dH_{33}}$	C_{HWB}	$C_{Hl_{jj}}^{(3)}$	$C_{ll_{1221}}$
α	NLO QCD	20.3%	20.3%	20.3%	20.3%	20.3%	-	-
	NLO EW	-5.2 %	2.1%	-11.0%	4.2%	-6.7%	-	-
	NLO correction	15.1%	22.4%	9.3%	24.5%	13.6%	-	-
α_μ	NLO QCD	20.3%	20.3%	20.3%	20.3%	-	20.3%	20.3%
	NLO EW	-0.8 %	2.1%	2.0%	1.9%	-	0.9%	-0.8%
	NLO correction	19.5%	22.4%	22.3%	22.2%	-	21.2%	19.5%
LEP	NLO QCD	20.3%	20.3%	20.3%	20.3%	-	20.3%	20.3%
	NLO EW	-0.7 %	2.1%	1.6%	1.9%	-	0.7%	-0.9%
	NLO correction	19.5%	22.3%	21.9%	22.2%	-	21.0%	19.3%

[2305.03763](#), [2312.08446](#)

α : (α , M_W , M_Z), α_μ : (G_F , M_W , M_Z), LEP: (α , G_F , M_Z)

HEFT vs SMEFT

- HEFT has Higgs (h) as a singlet, SMEFT has Higgs (H) as a doublet
- Different expansions (derivative vs $1/\Lambda$)

$$L_{HEFT} \sim \frac{v^2}{4} \left[1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right] \text{Tr} \{ D_\mu U^\dagger D_\mu U \} + \frac{1}{2} (\partial_\mu h)^2 - V(h)$$

$$V(h) = \frac{1}{2} m_h^2 h^2 \left(1 + \kappa_3 \frac{h}{v} + \frac{\kappa_4}{4} \frac{h^2}{v^2} + \dots \right)$$

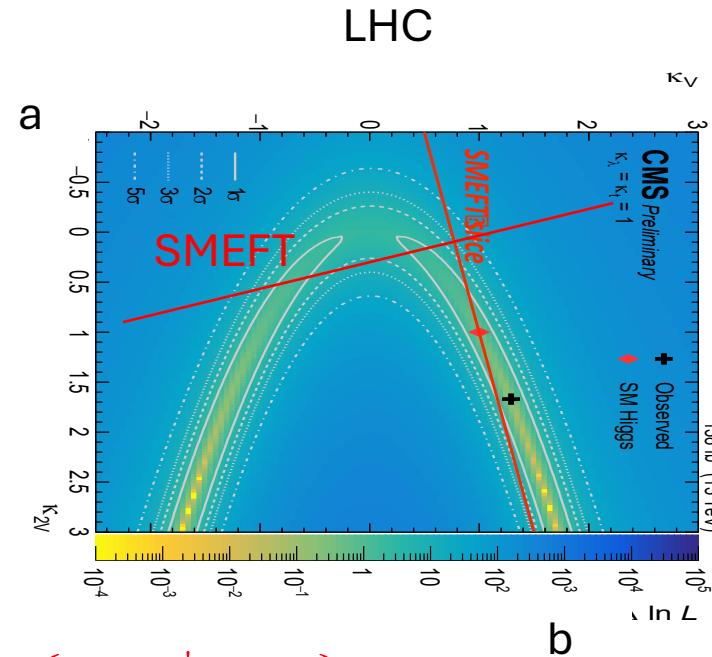
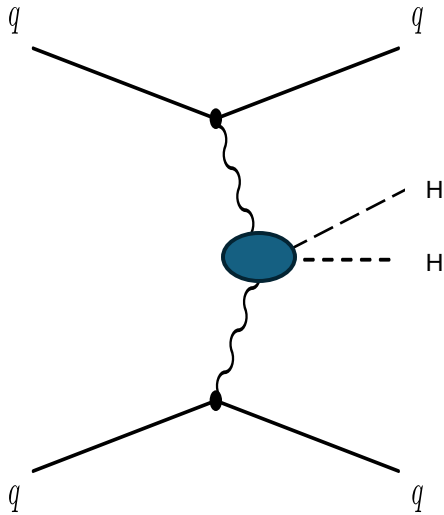
$$D_\mu U = \partial_\mu U + ig W_\mu^a \frac{\sigma^a}{2} U - ig' U \frac{\sigma^3}{2} B_\mu$$

- Unitary gauge $U=1$, suggests $WW \rightarrow hh$ is good probe of HEFT vs SMEFT
- SM: $a=b=\kappa_3=\kappa_4=1$ SMEFT: $b-a = \frac{3C_{H\Box} v^2}{\Lambda^2}$

 Correlations

SMEFT vs HEFT

HH production via VBS can potentially distinguish SMEFT from HEFT

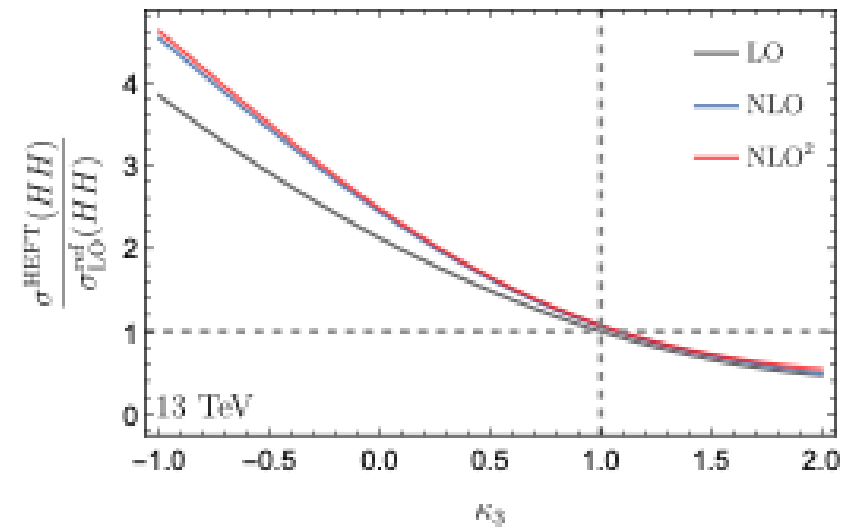
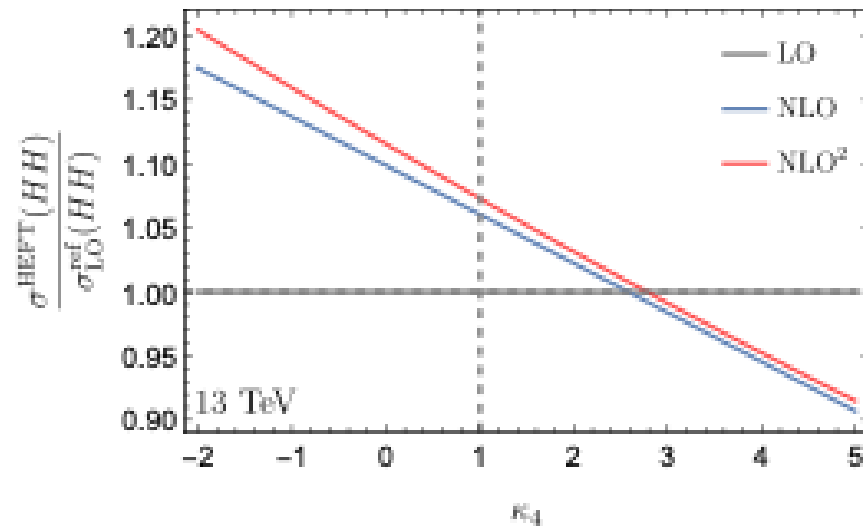
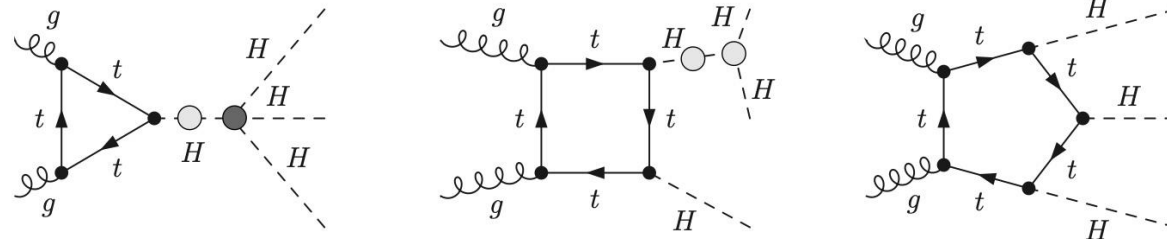


Is the Higgs in an
SU(2) doublet?

2211.09605

$$L_{HEFT} \sim \frac{v^4}{4} \left[1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right] \text{Tr} \{ D_\mu U^\dagger D_\mu U \}$$

HEFT HHH vs HH



[2405.05385](https://arxiv.org/abs/2405.05385)

Conclusions

- Much interesting work for theorists to make SMEFT useful for experiment
 - Theory uncertainties?
 - New AI/ML tools for global fits?
 - More understanding of HEFT vs SMEFT correlations
- At the top of my personal list are more dimension-8 studies and completion of NLO EW dimension-6 calculations
- Much progress that I didn't cover (in particular automated tools for 1 loop matching, progress in 2-loop SMEFT RGEs)