

Simultaneous Resonant and Broadband Detection via Cavities and Circuits

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Quantum Sensing meets

Ultra-high Frequency Gravitational Waves



THE CENTER OF GRAVITY



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Electromagnetic Detection of Ultralight Bosons and HFGW

$\vec{J}_{\text{eff}}$ in a **cavity** or **shield room**: $\square \vec{A} = \vec{J}_{\text{eff}}$.

Dark photon A' :

$$\vec{J}_{\text{eff}} = \epsilon m_{A'}^2 \vec{A}';$$

No background field.

Axion a :

$$\vec{J}_{\text{eff}} = g_{a\gamma} \vec{B}_0 \partial_t a;$$

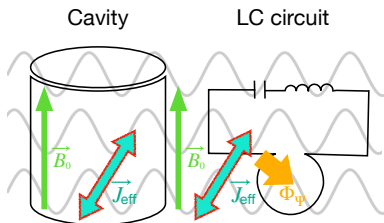
Background $\vec{B}_0$.

GW h :

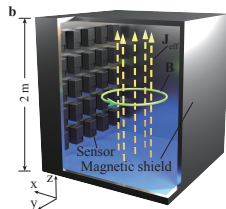
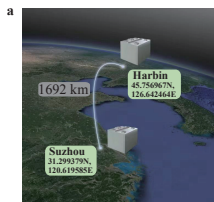
$$\vec{J}_{\text{eff}} \sim \partial(\hbar F_0);$$

Background $\vec{E}_0$ or $\vec{B}_0$.

- ▶ **Resonant cavity:** $\omega_{\text{rf}} \sim \omega_{J_{\text{eff}}}$.
- ▶ **Circuit/magnetometer:** $B \sim |\vec{J}_{\text{eff}} V^{1/3}|$ from $\vec{J}_{\text{eff}}$ where $1/\omega_{J_{\text{eff}}} \gg V^{1/3}$.



e.g. ADMX, HAYSTAC, CAPP, ORGAN, DM radio...

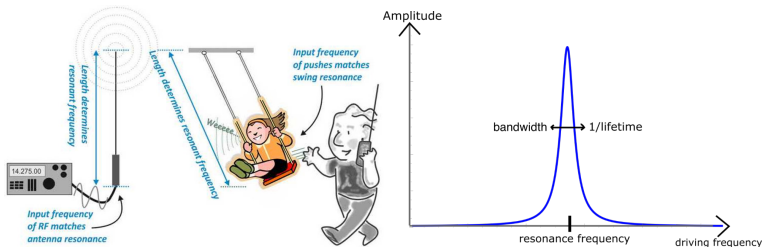


[Jiang et al, Nature Commun, 2305.00890]

Electromagnetic Single-Mode Resonator

- **Resonator:** significant response to signal at narrow bandwidth.

Trade-off between peak response and bandwidth.



e.g. cavity or circuit for capturing $\vec{J}_{\text{eff}}$.

- Require **multiple scans** to cover ultralight dark matter mass window;
- Challenge for **transient high-frequency gravitational wave** sources.

Standard Quantum Limit for Resonant Detection

- **Power law detection:**

Noise PSD: **resonant intrinsic noise** $S_{\text{int}} \propto kT$ + flat readout noise $S_r \sim \hbar\omega$.

- Responses to S_{sig} and S_{int} are the same.

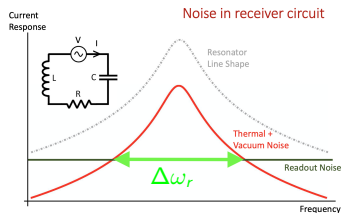
Scan rate $\propto \Delta\omega_r$ within which $S_{\text{int}} \gg S_r$.

- **Standard quantum limit:** [Chaudhuri et al 18']

Optimizing **readout coupling** Q_{cpl} :

$\beta \equiv Q_0/Q_{\text{cpl}} = 2$ for quantum-noise limited $kT \ll \hbar\omega$ [Krauss et al 85'].

$\beta \approx 2n_{\text{occ}}$ for $n_{\text{occ}} \approx kT/(\hbar\omega) \gg 1$ and $\Delta\omega_r \propto n_{\text{occ}}$ [e.g., Chaudhuri et al 18', Berlin et al 19'].



$S_{\text{int}} \propto$ Cauchy distribution

Beyond Standard Quantum Limit

► Power law detection:

Noise PSD: **resonant intrinsic noise** $S_{\text{int}} \propto kT$ + flat readout noise $S_r \sim \hbar\omega$.

► Responses to S_{sig} and S_{int} are the same.

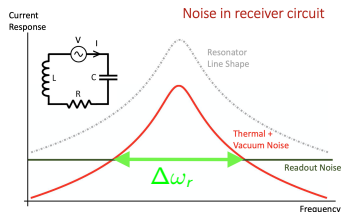
Scan rate $\propto \Delta\omega_r$ within which $S_{\text{int}} \gg S_r$.

► Beyond standard quantum limit:

Squeezing S_r , e.g., HAYSTAC Phase II.

Increasing the sensitivity to S_{sig} , e.g., white light cavity in optomechanics/GW detection [Miao et al 15].

Qubit.



$S_{\text{int}} \propto$ Cauchy distribution

Full length article

White-light cavities, atomic phase coherence, and gravitational wave detectors

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Abstract

We propose a new concept to realize optical cavities with **large buildup but broadband response** (*white-light cavities*) using atomic phase coherence. We demonstrate that strongly driven double- Λ systems can show **negative dispersion** without absorption, which is needed in order to compensate for the variation of the wavelength with frequency. Internal buildup profiles and the cavity bandwidth of standard devices and *white-light cavities* will be briefly compared. These devices may be useful to improve the bandwidth and sensitivity of future generations of laser interferometric gravitational wave detectors.

PACS: 07.60.Ly; 42.50.Ar; 95.30.SF

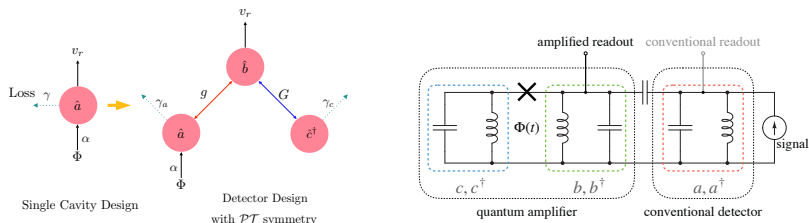
Keywords: Atomic phase coherence; Gravitational wave detector; White-light cavity; Negative dispersion; Vanishing absorption; Four-level system

- **Negative dispersion medium cancels the frequency-dependent phase shift:**

Achieves **simultaneously large response** and **broad bandwidth**.

- **Microwave cavity?**

White Light Haloscope [Li et al 20']



Probing mode:
 $\hbar\alpha(\hat{a} + \hat{a}^\dagger)\Psi$

- ▶ **Beam-splitting:** $\hbar g(\hat{a}\hat{b}^\dagger + \hat{a}^\dagger\hat{b})$.
- ▶ **Non-degenerate parametric interaction:** $\hbar G(\hat{b}\hat{c} + \hat{b}^\dagger\hat{c}^\dagger)$.
- ▶ **$\mathcal{PT}$ -symmetry ($\hat{a} \leftrightarrow \hat{c}^\dagger$) emerges** when $g = G$.

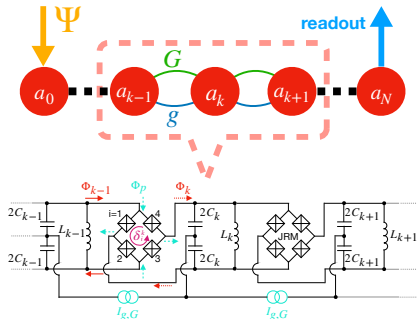
$$\begin{aligned} (\dot{\hat{a}} + \dot{\hat{c}}^\dagger) &= -i(g - G)\hat{b} - i\alpha\Psi + \dots; \\ \dot{\hat{b}} &= -\gamma_r\hat{b} - ig(\hat{a} + \hat{c}^\dagger) + \dots. \end{aligned}$$
- ▶ **Double resonance** when $g \gg$ **intrinsic dissipation** γ :

$$S_{\text{sig}}^{\text{WLC}}(\Omega) = \frac{2\gamma_r\alpha^2 S_\Psi(\Omega)}{(\gamma + \gamma_r)^2 + \Omega^2} \left(\frac{g^2}{\gamma^2 + \Omega^2} \right).$$

Readout coupling γ_r

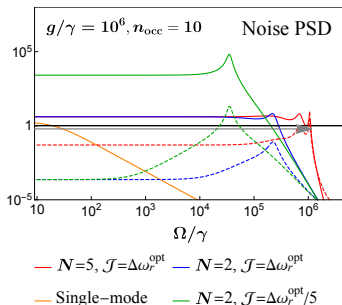
Response Width for Multi-mode Resonators

Broadened response in multi-mode resonators:



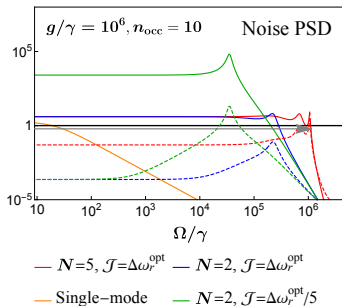
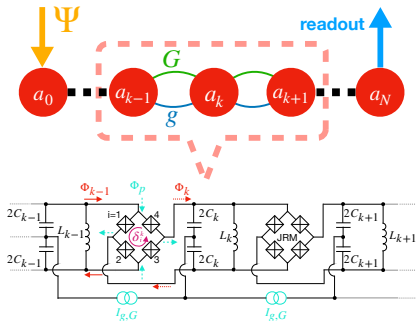
$$H_{\text{ch}} = \sum_{k=0}^{N-1} \left(\underset{\text{flatten}}{ig} \hat{a}_k \hat{a}_{k+1}^\dagger + \underset{\text{amplify}}{iG} \hat{a}_k \hat{a}_{k+1} + h.c. \right).$$

- ▶ $|g| > |G|$ to make the system stable.
- ▶ Large $\mathcal{J} \equiv \sqrt{|g|^2 - |G|^2}$ to **flatten** the response as required by **dynamic range** and **robustness**.



Response Width for Multi-mode Resonators

Broadened response in multi-mode resonators:



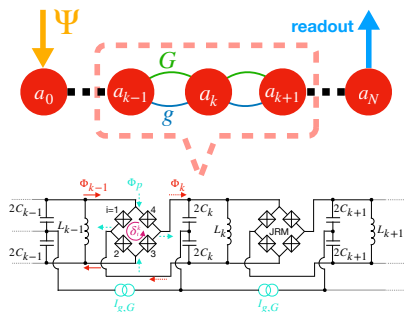
$$H_{\text{ch}} = \sum_{k=0}^{N-1} \left(i g \hat{a}_k \hat{a}_{k+1}^\dagger + i G \hat{a}_k \hat{a}_{k+1} + h.c. \right).$$

► New sensitivity limit for multi-mode resonators, optimal for SRF cavity:

$$\frac{\Delta\omega_r^{\text{MM}}}{\Delta\omega_r^{\text{SM}}} \propto \left(\frac{gQ_0}{\omega_{\text{rf}} n_{\text{occ}}} \right)^{\frac{2N}{2N+1}} \rightarrow \frac{Q_0}{n_{\text{occ}}} \text{ as } N \gg 1, g \rightarrow \omega_{\text{rf}}, \Delta\omega_r^{\text{MM}} \rightarrow \omega_{\text{rf}}.$$

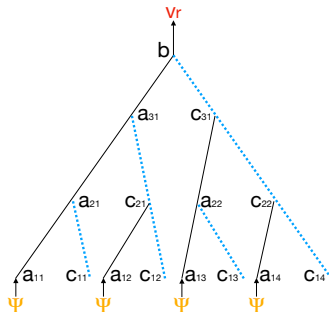
Chain and Binary Tree

Non-hermitian chain [McDonald et al 18',
YC et al Rept.Prog.Phys. 2309.12387]



- **Single-quadrature** amplification
- **Phase fluctuations & nonlinearities?**

Binary tree [YC et al, PRR 2103.12085,
Rept.Prog.Phys. 2309.12387]



- **Both quadratures** experience gain
- Compatible with **sensor networks**

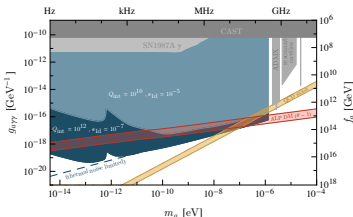
- **HAYSTAC Phase IV, CEASEFIRE: prototype calibration** for $N = 1$ **non-hermitian chain** with $g \sim \text{MHz}$ [Wurtz et al 21', Jiang et al 23'].

Heterodyne Upconversion

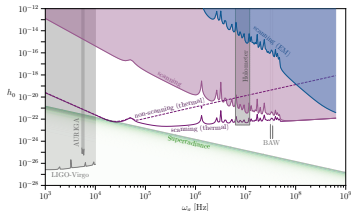
- **Heterodyne upconversion** [Goryachev et al 19', Berlin et al 19', UPLOAD 19' 23']:

injecting **AC pump mode**: $\omega_{rf} - \omega_0 \approx \omega_a$ or ω_h .

- Resonant scan: tune $\omega_{rf} - \omega_0$ from Hz to GHz:



[Berlin et al 19']



[Berlin et al 23']

- Both **EM** and **mechanical** coupling from **GW** [Berlin et al 21' 23'].
- **Jing Shu's** talk for **SHANHE** collaboration's **experimental progress**.

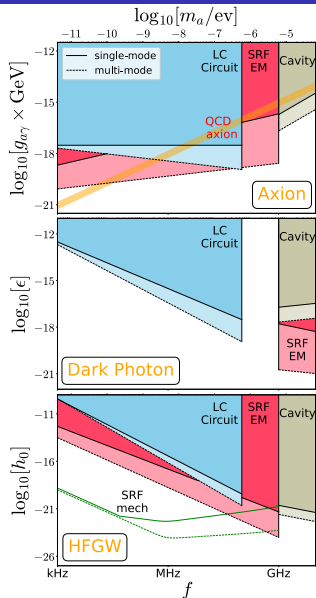
Simultaneous Resonant and Broadband Detection

- ▶ e-fold time: 10^7 s.
- ▶ Solid lines: **single-mode resonators**
with $\Delta\omega_r \sim n_{\text{occ}}\omega_{\text{rf}}/Q_0$.
- ▶ Dashed lines: **multi-mode resonators**
with $\Delta\omega_r \rightarrow \omega_{\text{rf}}$.

- ▶ DC cavity and LC circuits:

$$\frac{\text{SNR}_{\text{MM}}^2}{\text{SNR}_{\text{SM}}^2} \approx \frac{Q_0}{n_{\text{occ}}}.$$

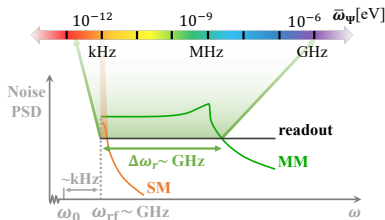
High n_{occ} of LC circuits at low frequency made enhancement ineffective.



[YC et al, Rept.Prog.Phys. 2309.12387]

Simultaneous Resonant and Broadband Detection

► SRF heterodyne upconversion:

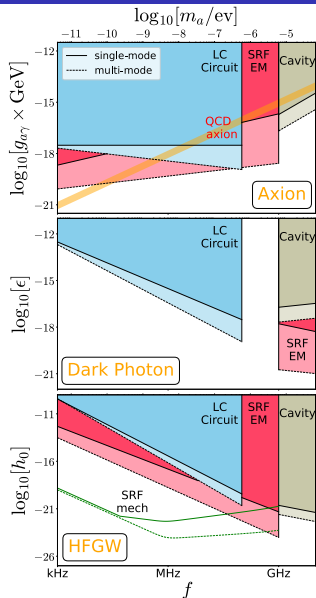


Simultaneous scan $N_e = 6$ orders of ω_ψ

with **significant response**:

$$\frac{\text{SNR}_{\text{MM}}^2}{\text{SNR}_{\text{SM}}^2} \simeq N_e \frac{\bar{\omega}_\psi Q_0}{\omega_{\text{rf}} n_{\text{occ}}^{\text{eff}}}$$

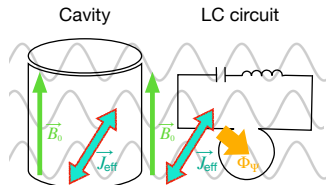
- Ineffective enhancement for **GW mechanical resonance** near kHz due to high $n_{\text{occ}}^{\text{eff}}$.



[YC et al, Rept.Prog.Phys. 2309.12387]

Summary

- ▶ **Resonant cavities or circuits** are powerful detectors for **ultralight bosons and HFGW**.



- ▶ **Multi-mode resonantors** have **broadened response**.
→ **Simultaneous resonant and broadband** detection for **SRF-based upconversion**.

Thank you!