

# Flavour Physics: Current Status and Future Directions

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**Flavour for New Physics at Present and Future Colliders**

**MITP**

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## Outline:

1. The problem of flavour
2. Open problems in hadronic physics
3. Beyond the LHC and Belle II

# Motivation

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there are open problems:

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**Hierarchy problem**

**dark matter/dark energy**

**flavour hierarchies**

**neutrino masses**

**gravity**

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SM(EFT)

$\Lambda_{EW}$

Energy

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UV theory

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$\Lambda_{UV}$

$\Lambda_{EW}$

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# Motivation

Despite the SM successes,  
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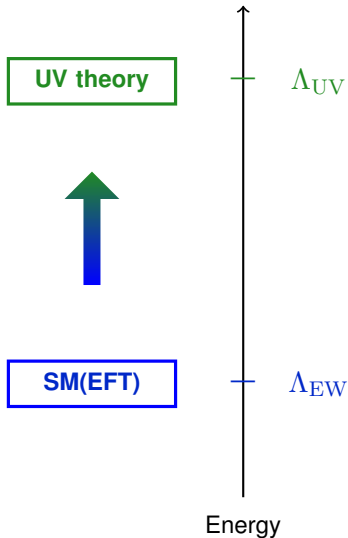
Hierarchy problem

dark matter/dark energy

flavour hierarchies

neutrino masses

gravity



# The (two) flavour problems

1. **The SM flavour problem:** The measured Yukawa pattern doesn't seem accidental

⇒ Is there any deeper reason for that?

2. **The NP flavour problem:** If we regard the SM as an EFT valid below a certain energy cutoff  $\Lambda$ , why don't we see any deviations in flavour changing processes?

⇒ Which is the flavour structure of BSM physics?



# The SM flavour problem

$$\mathcal{L}_{\text{Yukawa}} \supset Y_u^{ij} \bar{Q}_L^i H u_R^j$$

$$Y_u \sim y_t \begin{pmatrix} \text{light green circle} & \text{light green circle} & \text{dark green circle with } 0.003 \\ & \text{dark green circle} & \text{dark green circle with } 0.04 \\ & & 1 \end{pmatrix}$$

# The SM flavour problem

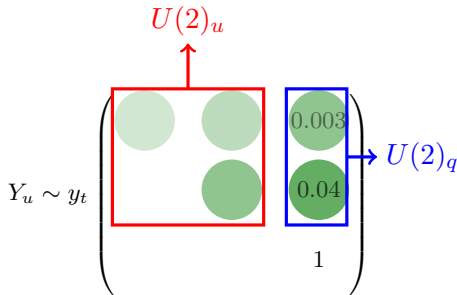
$$\mathcal{L}_{\text{Yukawa}} \supset Y_u^{ij} \bar{Q}_L^i H u_R^j$$

$$Y_u \sim y_t \begin{pmatrix} & & \\ & & \\ & & 1 \end{pmatrix}$$

Exact  $U(2)^n$  limit

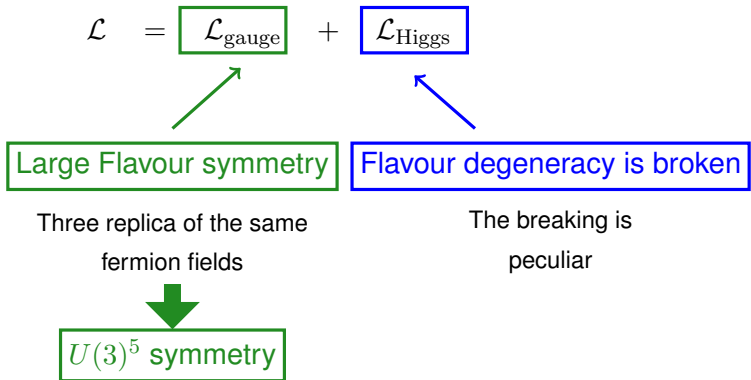
# The SM flavour problem

$$\mathcal{L}_{\text{Yukawa}} \supset Y_u^{ij} \bar{Q}_L^i H u_R^j$$



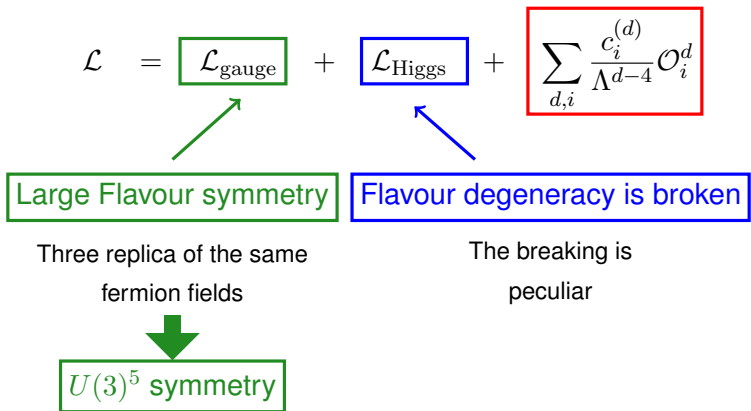
An approximate  $U(2)^n$  is acting  
on the light families!

# The NP flavour problem



- In the SM: accidental  $U(3)^5 \rightarrow \text{approx } U(2)^n$

# The NP flavour problem

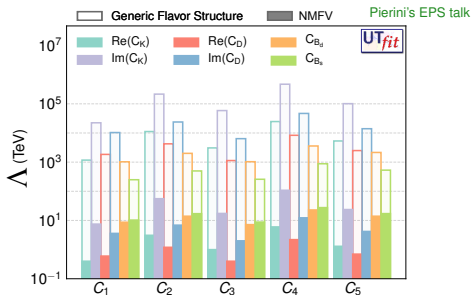
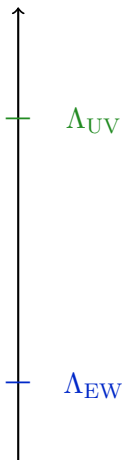


- In the SM: accidental  $U(3)^5 \rightarrow \text{approx } U(2)^n$
- **What happens when we switch on NP?**

# The NP flavour problem

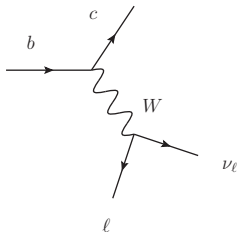
$$\mathcal{L} = \boxed{\mathcal{L}_{\text{gauge}}} + \boxed{\mathcal{L}_{\text{Higgs}}} + \boxed{\sum_{d,i} \frac{c_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^d}$$

- What is the energy scale of NP?
- Why haven't observed any violation of accidental symmetries yet?

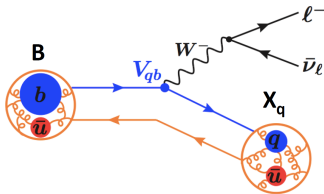


**no breaking of the  $U(2)^n$  flavour symmetry at low energies**

# Partonic vs Hadronic



$$\mu_{\text{partonic}} = m_b$$



$$\mu_{\text{hadronic}} = \Lambda_{\text{QCD}}$$

**Fundamental challenge to match  
partonic and hadronic descriptions**

# Open problems in hadronic physics



# What are the open themes in hadronic physics?

1. Extraction of the CKM elements  $V_{cb}$  and  $V_{ub}$
2. Non-local effects in  $b \rightarrow s \ell \ell$
3. Non-leptonic decays

# The $V_{cb} - V_{ub}$ puzzle

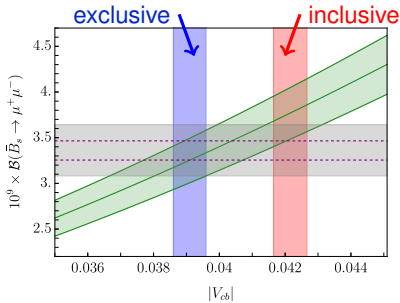
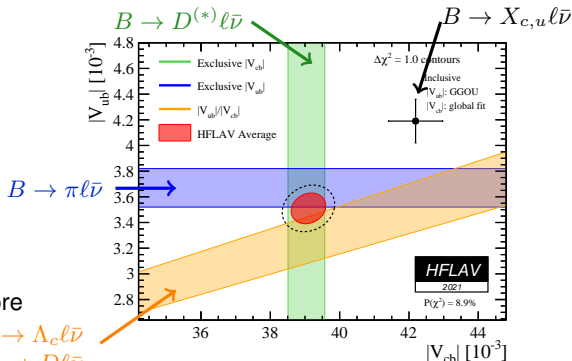
- Large discrepancies between inclusive and exclusive determinations
- Recent work mostly on  $B \rightarrow D^*$  due to new lattice QCD form factors determinations
- When precision increases, more puzzles arise

$$\Lambda_b \rightarrow p \ell \bar{\nu} / \Lambda_b \rightarrow \Lambda_c \ell \bar{\nu}$$

$$B_s \rightarrow K \ell \bar{\nu} / B \rightarrow D \ell \bar{\nu}$$

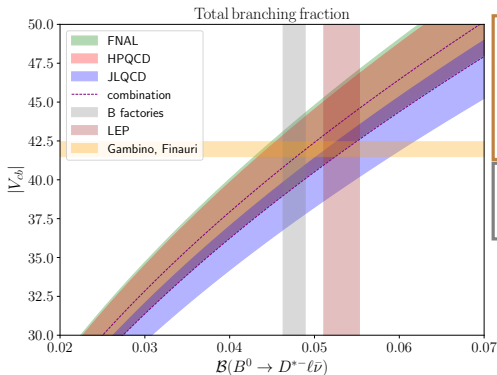
- $|V_{cb}|$  is a fundamental parameter

- $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) \sim |V_{cb}|^2$
- $\mathcal{B}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \sim |V_{cb}|^4$



# Exclusive $V_{cb}$ from $B \rightarrow D^*$

MB, A. Jüttner, '23 + WIP



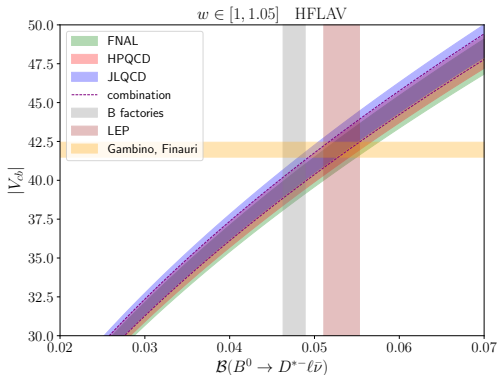
| Experiment     | BF (rescaled) [%]      | Parameters                                       |
|----------------|------------------------|--|
| ALEPH          | 5.56 +/- 0.27 +/- 0.33 | <a href="#">input</a> <a href="#">parameters</a> |
| OPAL incl      | 6.13 +/- 0.28 +/- 0.57 | <a href="#">input</a> <a href="#">parameters</a> |
| OPAL excl      | 5.17 +/- 0.20 +/- 0.36 | <a href="#">input</a> <a href="#">parameters</a> |
| DELPHI incl    | 4.96 +/- 0.14 +/- 0.35 | <a href="#">input</a> <a href="#">parameters</a> |
| DELPHI excl    | 5.23 +/- 0.20 +/- 0.42 | <a href="#">input</a> <a href="#">parameters</a> |
| CLEO           | 6.17 +/- 0.19 +/- 0.37 | <a href="#">input</a> <a href="#">parameters</a> |
| BELLE untagged | 4.90 +/- 0.02 +/- 0.16 | <a href="#">input</a> <a href="#">parameters</a> |
| BELLE tagged   | 4.95 +/- 0.11 +/- 0.22 | <a href="#">input</a> <a href="#">parameters</a> |
| BABAR untagged | 4.52 +/- 0.04 +/- 0.33 | <a href="#">input</a> <a href="#">parameters</a> |
| BABAR tagged   | 5.26 +/- 0.16 +/- 0.31 | <a href="#">input</a> <a href="#">parameters</a> |
| Average        | 5.06 +/- 0.02 +/- 0.12 | chi2/dof = 16.0/9 (CL = 0.0661)                  |

- Shape information shifts the total branching fraction prediction

Thanks to C. Schwanda  
for the averages!

# Exclusive $V_{cb}$ close to zero-recoil

MB, A. Jüttner, '23 + WIP



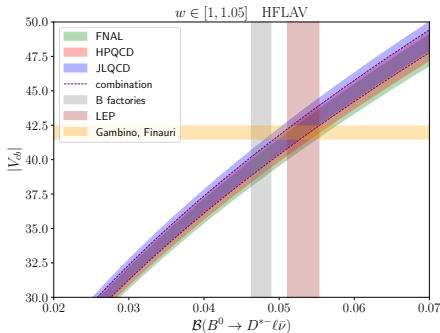
- Combination of Belle and Belle II angular distribution data
  - Only the zero-recoil bin
  - In that region, one form factor dominates
- Which branching ratio should we use?

B factories:  $|V_{cb}| = 40.07 \pm 0.86$

LEP:  $|V_{cb}| = 42.37 \pm 1.09$

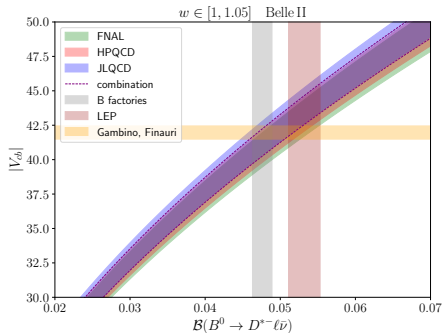
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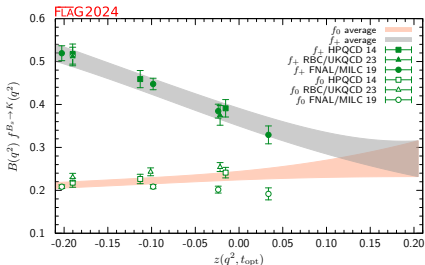
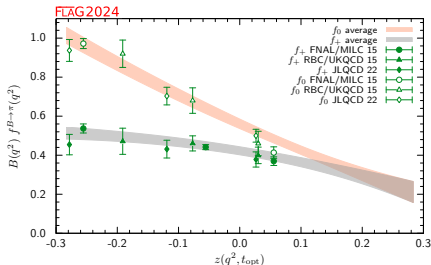
LEP:  $|V_{cb}| = 42.37 \pm 1.09$



B factories:  $|V_{cb}| = 41.24 \pm 1.15$

LEP:  $|V_{cb}| = 43.60 \pm 1.35$

# A few words on exclusive $V_{ub}$

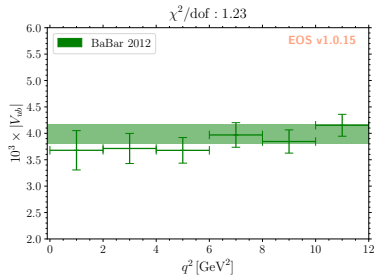
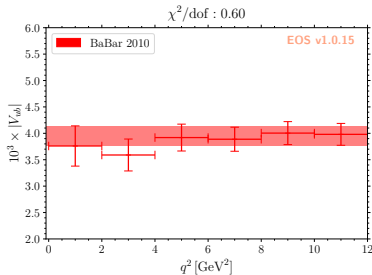
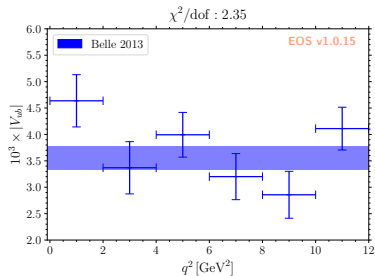
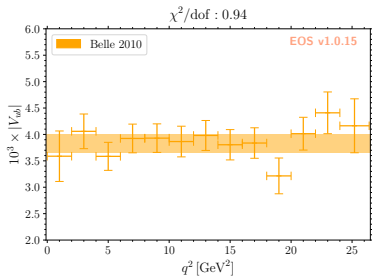


- There are tensions in the lattice determinations of  $f_0$
- $f_+$  and  $f_0$  are correlated through the kinematic constraint

$$f_+(q^2 = 0) = f_0(q^2 = 0)$$

- The chiral continuum extrapolation using the helicity base or the "lattice" base has some tricky points and checks are ongoing

# $V_{ub}$ from $B \rightarrow \pi \ell \bar{\nu}$



**Not inflating uncertainties reveals a few tensions**

# Charm-loop effects in $b \rightarrow s \ell^+ \ell^-$

$$\mathcal{H}_{\text{eff}} = -4 \frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* [-\mathcal{C}_1 \mathcal{O}_1 - \mathcal{C}_2 \mathcal{O}_2 + \mathcal{C}_7 \mathcal{O}_7 + \mathcal{C}_9 \mathcal{O}_9 + \mathcal{C}_{10} \mathcal{O}_{10}]$$

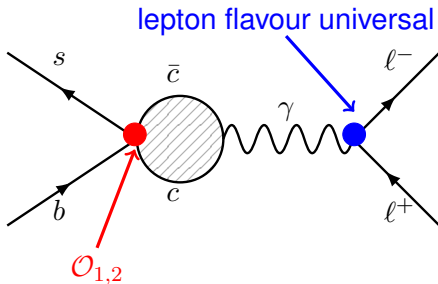
$$\mathcal{O}_1 = (\bar{s} \gamma^\mu P_L b) (\bar{c} \gamma_\mu c)$$

$$\mathcal{O}_2 = (\bar{s} \gamma^\mu T^a P_L b) (\bar{c} \gamma_\mu T^a c)$$

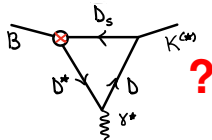
$$\mathcal{O}_9 = (\bar{s} \gamma^\mu P_L b) (\bar{\ell} \gamma_\mu \ell)$$

$$\mathcal{O}_{10} = (\bar{s} \gamma^\mu P_L b) (\bar{\ell} \gamma_\mu \gamma_5 \ell)$$

$$\mathcal{O}_7 = (\bar{s} \sigma^{\mu\nu} P_R b) F_{\mu\nu}$$



$$C_9 \rightarrow C_9^{\text{eff}}(q^2) = C_9 + C_9^{\text{LD}}(q^2)$$

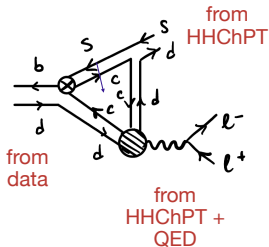


[Ciuchini, Fedele, Franco, Paul, Silvestrini, Valli, '22]

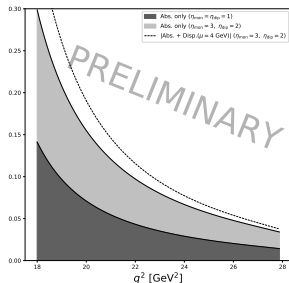
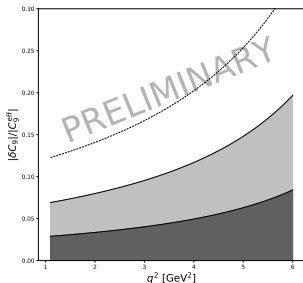


# Models for the rescattering

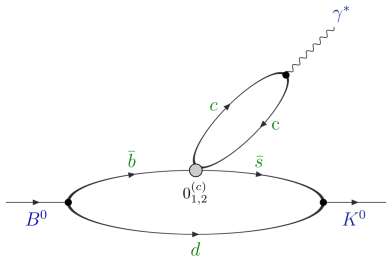
G. isidori, Z. Polonsky, A.Tinari, '24 + WIP



- Use of heavy hadron chiral perturbation theory + data + QED to determine the various vertices
- Better reliability at maximum  $q^2$
- The tower of all possible combinations of states is accounted for by multiplicity factors based on the respective branching ratio
- Relative phases and how various contributions add up depend on unknown relative phases



# News from the lattice



$$H^\mu(\mathbf{q}) \sim \int_{E_0}^{\infty} \frac{dE}{2\pi} \frac{\rho^\mu(E, \mathbf{q})}{E - m_B - i\epsilon}$$

⇒ Spectral densities are connected to correlation functions

$$C^\mu(t, \mathbf{q}) \sim \int_{E_0}^{\infty} \frac{dE}{2\pi} e^{-Et} \rho^\mu(E, \mathbf{q})$$

**One has to solve a very complicated inverse problem**

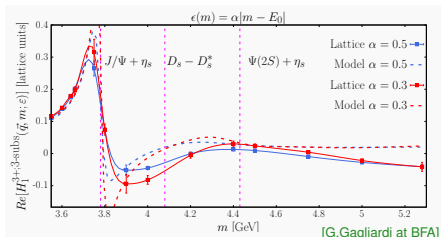
# News from the lattice

Solution: **smeared correlators**

$$H^\mu(\mathbf{q}) \sim \int_{E_0}^{\infty} \frac{dE}{2\pi} \frac{\rho^\mu(E, \mathbf{q})}{E - m_B - i\epsilon} = \lim_{\epsilon \rightarrow 0} \int_{E_0}^{\infty} \frac{dE}{2\pi} \mathcal{K}_\epsilon(E, \mathbf{q}) \rho^\mu(E, \mathbf{q})$$

- The smearing makes the problem solvable in a finite volume
- This technique has been successfully applied to various problems like  $g - 2$  computation, inclusive semileptonic  $D$  decays

$$B_s \rightarrow \eta_s \ell^+ \ell^-$$



- Extend to  $B \rightarrow K \ell^+ \ell^-$
- Attempt the  $\epsilon \rightarrow 0$  limit
- Include more ensembles
- Perform non-perturbative renormalisation

# The $B_{(s)} \rightarrow D_{(s)}^* K(\pi)$ decays

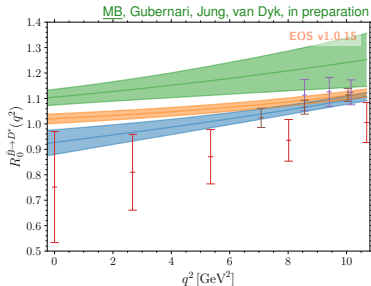
## Why?

⇒ crucial inputs to extract fragmentation fractions from data

$$\mathcal{A}(\bar{B}_q^0 \rightarrow D_q^{*+} L^-) = -i \frac{G_F}{\sqrt{2}} V_{uq2}^* V_{cb} a_1(D_q^{*+} L^-) f_L$$

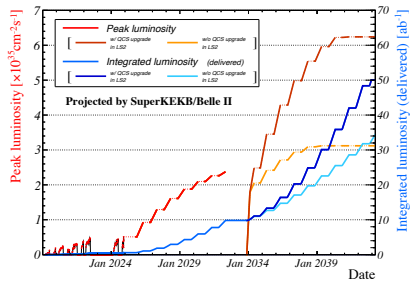
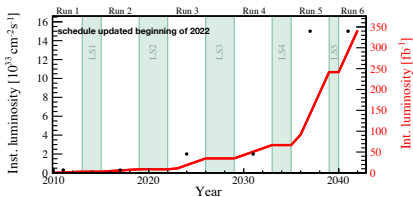
$$\times A_0^{\bar{B}_q \rightarrow D_q^*}(M_L^2) 2M_{D_q^*} \varepsilon^*(\lambda=0) \cdot q$$

- Currently a  $4\sigma$  discrepancy  
MB, Huber, Gubernari, Jung, van Dyk, '20
- Initiated further works on power correction  
M.L. Piscopo, A. Rusov, '23



- New form factor analyses suggest a 15% downward shift in the scalar form factor
- This would reduce the tension with experimental data
- This is just for  $B \rightarrow D^*$ ,  $B \rightarrow D$  is unchanged at the current status

# **Beyond the LHC and Belle II**



- Both the LHCb and Belle II programs are supported by not only unprecedented datasets, but also by software and analysis improvements
- What can be yet improved?

# FCC-ee flavour opportunities

FCC: combines some of the best features of  $pp$  colliders and  $B$  factories

- High statistics
- Relatively clean environment due to high boost of  $b$  quarks

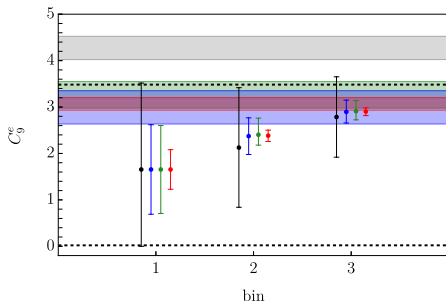
| Particle production ( $10^9$ ) | $B^0/\overline{B}^0$ | $B^+/B^-$ | $B_s^0/\overline{B}_s^0$ | $B_c^+/\overline{B}_c^-$ | $\Lambda_b/\overline{\Lambda}_b$ | $c\overline{c}$ | $\tau^+\tau^-$ |
|--------------------------------|----------------------|-----------|--------------------------|--------------------------|----------------------------------|-----------------|----------------|
| Belle II                       | 27.5                 | 27.5      | n/a                      | n/a                      | n/a                              | 65              | 45             |
| FCC-ee                         | 620                  | 620       | 150                      | 4                        | 130                              | 600             | 170            |

- ⇒ unique opportunity to study flavour under all possible angles
- ⇒ all possible decay modes are accessible

# $B \rightarrow K^* e^+ e^-$ at FCC-ee

MB, C. Cornella, J. Davighi, '25

- Reconstruction is as good as for muons  
⇒ estimated around  $\sim 80\%$ \* in all phase space
- The statistical error is below 1% and half of what is expected from HL-LHC



— current  
— HL-LHC+P<sub>1</sub>

— FCC+P<sub>1</sub>  
— FCC+P<sub>2</sub>

- Only  $\mathcal{B}(B \rightarrow K^* e^+ e^-)$
- High Lummy projections:  $\sigma_{\text{stat}} = \sigma_{\text{syst}}$
- Theory projections:

$$P_1 : \sigma_{F_i} \rightarrow \sigma_{F_i}/2,$$

$$P_2 : \sigma_{F_i} \rightarrow \sigma_{F_i}/5,$$

- Only with P<sub>2</sub> we can fully exploit the FCC-ee statistics

\*estimates from  $B \rightarrow K^* \tau^+ \tau^-$  studies confirmed by ad hoc simulations by S. Monteil



# Conclusions

- Flavour physics is a powerful test for new physics living at different energy scales
- We have a lot of puzzles to solve, but this is just a sign of the advancements in both theory and experiments
- One of the most urgent puzzle to resolve concerns  $V_{cb}$  and  $V_{ub}$
- Precise SM predictions are the key to unveiling potential NP signals
- The excellent experimental prospects, on the short and long term, combining with theory and experimental advancements

# Appendix