



Universidad de Granada

FTAE
High Energy Theory

Efficient on-shell matching

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[2411.12798]

Matching a theory: *off-shell* vs *on-shell*

<i>OFF-SHELL</i>	<i>VS</i>	<i>ON-SHELL</i>
Small number of diagrams (one-light-particle-irreducible)	1 — 0	All diagrams (light bridges too)
Large set of operators (Green's basis)	1 — 1	Smaller set of operators (physical basis)
Reduction to the physical basis	1 — 2	No reduction is needed
Contribution directly local in momenta		Delicate cancelation of non-local contributions

Matching a theory: *off-shell* vs *on-shell*

4-pt:

$$\text{Tree-level exchange} + \text{Contact} = \lambda + \frac{\kappa^2}{M^2} = C_4$$

6-pt:

$$\text{Tree-level exchange} + \text{Contact} + \dots = C_6 + \frac{i C_4^2}{(p_{456}^2 - m^2)} + \dots$$

$$\frac{\eta \kappa^2}{M^4} + \frac{\lambda \kappa^2}{M^2} \frac{i}{(p_{456}^2 - m^2)} + \dots = C_6 + \frac{i C_4^2}{(p_{456}^2 - m^2)} + \dots$$

~~$\frac{i C_4^2}{(p_{456}^2 - m^2)}$~~
cancel

Matching a theory: *off-shell* vs *on-shell*

OFF-SHELL

VS

ON-SHELL

Small number of diagrams
(one-light-particle-irreducible)

1 — 0

All diagrams (light bridges too)

Large set of operators
(Green's basis)

1 — 1

Smaller set of operators
(physical basis)

Reduction to the physical basis

1 — 2

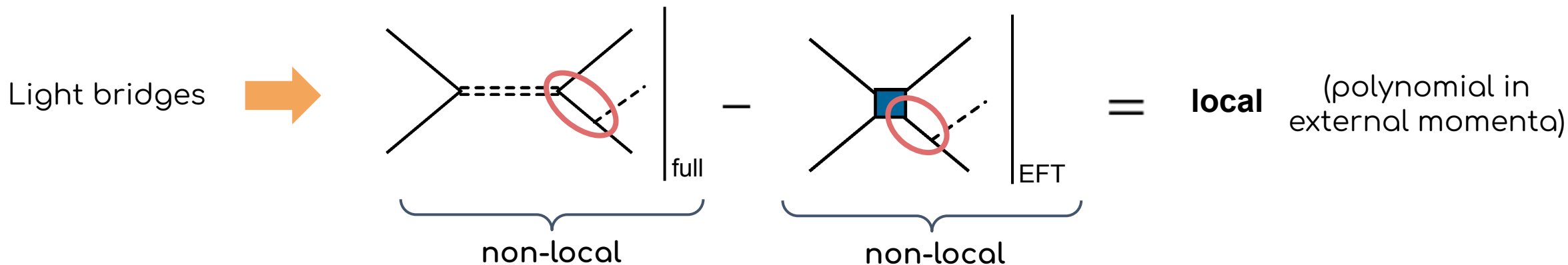
No reduction is needed

Contribution directly local
in momenta

Delicate cancellation of
non-log terms



On-shell matching: *non-localities*



■ Difficult to follow this cancelation analytically

■ The procedure is to be **numerical** but exact



Substitution of **random-generated kinematics**



Rational kinematics

Spinor Helicity Formalism

[arXiv:2304.01589, arXiv:2202.02681]



$$\mathcal{M} = \alpha p_2 \cdot p_3 + \frac{\beta^2}{(p_1 + p_4)^2 - m^2} = \frac{9044503}{1681920} m^2 \alpha - \frac{840960}{8203543} \frac{\beta^2}{m^2}$$

Rational values for $\left\{ \begin{array}{l} \text{momenta} \\ \text{polarizations} \\ \text{spinors} \end{array} \right\}$ with symbolic masses m_i



Satisfying...

- Momentum conservation
- On-shell condition
- Dirac equation
- Transversality

On-shell matching: *the matching condition*

$$\frac{p^2}{M_\Phi^2} \ll 1$$

» Tree-level matching:

$$\left. \text{Diagram with } \mathcal{M}^{\text{tree}} \right|_{\text{EFT}} = \left. \text{Diagram with } \mathcal{M}^{\text{tree}} \right|_{\text{full}} + \mathcal{O}\left[\left(\frac{p^2}{M_\Phi^2}\right)^n\right]$$

» One-loop matching:
(method of regions)

hard: $k \sim M_\Phi \gg p, m$

soft: $M_\Phi \gg k \sim p, m$

$$\left. \text{Diagram with } \mathcal{M}^{\text{tree}} \right|_{\text{EFT renorm}} = \left. \text{Diagram with } \mathcal{M}^{\text{hard}} \right|_{\text{full renorm}} + \left. \text{Diagram with } \mathcal{M}^{\text{soft}} \right|_{\text{full renorm}} - \left. \text{Diagram with } \mathcal{M}^{\text{soft}} \right|_{\text{EFT renorm}}$$

On-shell condition: $p^2 = m_{\text{phys}}^2$

Wavefunction factors: $\mathcal{M} = (\sqrt{Z})^n \mathcal{M}_{\text{amp}}$

Almost cancel
evanescent contributions



On-shell matching: *evanescent shifts*

» One-loop matching:
(method of regions)

hard: $k \sim M_\Phi \gg p, m$

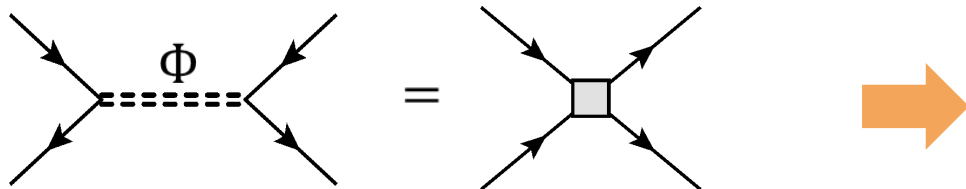
soft: $M_\Phi \gg k \sim p, m$

$$\left. \text{Diagram with } \mathcal{M}^{\text{tree}} \right|_{\text{EFT}} = \left. \text{Diagram with } \mathcal{M}^{\text{hard}} \right|_{\text{full}} + \underbrace{\left. \text{Diagram with } \mathcal{M}^{\text{soft}} \right|_{\text{full}} - \left. \text{Diagram with } \mathcal{M}^{\text{soft}} \right|_{\text{EFT}}}_{\text{Almost cancel evanescent contributions}}$$

Almost cancel
evanescent contributions



Copy of SM Higgs: $\mathcal{L} = \mathcal{L}_{\text{SM}} + D_\mu \Phi^\dagger D^\mu \Phi - M^2 \Phi^\dagger \Phi - (\mathcal{Y}^{pr} \bar{\ell}^p \Phi e^r + \text{h.c.})$



$$(\bar{\ell}^p e^r)(\bar{e}^s \ell^t) = -\frac{1}{2} (\bar{\ell}^p \gamma^\mu \ell^r)(\bar{e}^s \gamma_\mu e^t) \quad (d=4)$$

$$\left(\text{Diagram with red square vertex} - \text{Diagram with blue square vertex} \right) \neq 0$$

Only finite parts
coming from $\epsilon \times \frac{1}{\epsilon_{\text{UV}}}$

ΔC_{eB}

An example: *finite one-loop matching to SMEFT*

We solve for effective couplings perturbatively in the EFT order:

$$\lambda \rightarrow \lambda + \frac{g_2^4 m_H^2}{960\pi^2 M^2},$$

$$y_e^{pr} \rightarrow y_e^{pr} - \frac{1}{128\pi^2} (\mathcal{Y}\mathcal{Y}^\dagger y_e + 2y_e \mathcal{Y}^\dagger \mathcal{Y})^{pr} + \frac{m_H^2}{32\pi^2 M^2} \mathcal{Y}^{pr} (\mathcal{Y}^\dagger)^{st} y_e^{ts},$$

$$c_{HD} \rightarrow -\frac{g_1^4}{1920\pi^2 M^2},$$

$$c_{H\Box} \rightarrow -\frac{1}{7680\pi^2 M^2} (g_1^4 + 3g_2^4),$$

$$c_{He}^{pr} \rightarrow \frac{g_1^4}{1920\pi^2 M^2} \delta^{pr} + \frac{7g_1^2}{576\pi^2 M^2} (\mathcal{Y}^\dagger \mathcal{Y})^{pr} + \frac{1}{192\pi^2 M^2} (6\mathcal{Y}^\dagger y_e y_e^\dagger \mathcal{Y} + y_e^\dagger \mathcal{Y} \mathcal{Y}^\dagger y_e)^{pr},$$

$$c_{Hl}^{(1)pr} \rightarrow \frac{g_1^4}{3840\pi^2 M^2} \delta^{pr} + \frac{17g_1^2}{1152\pi^2 M^2} (\mathcal{Y}\mathcal{Y}^\dagger)^{pr} - \frac{1}{192\pi^2 M^2} (6\mathcal{Y} y_e^\dagger y_e \mathcal{Y}^\dagger + y_e \mathcal{Y}^\dagger \mathcal{Y} y_e^\dagger)^{pr},$$

$$c_{Hl}^{(3)pr} \rightarrow -\frac{g_2^4}{3840\pi^2 M^2} \delta^{pr} + \frac{g_2^2}{1152\pi^2 M^2} (\mathcal{Y}\mathcal{Y}^\dagger)^{pr} - \frac{1}{192\pi^2 M^2} (y_e \mathcal{Y}^\dagger \mathcal{Y} y_e^\dagger)^{pr},$$

$$c_{eB}^{pr} \rightarrow -\frac{g_1}{768\pi^2 M^2} (5\mathcal{Y}\mathcal{Y}^\dagger y_e + 2y_e \mathcal{Y}^\dagger \mathcal{Y})^{pr} + \frac{3g_1}{128\pi^2 M^2} \mathcal{Y}^{pr} (\mathcal{Y}^\dagger)^{st} y_e^{ts},$$

$$c_{eW}^{pr} \rightarrow -\frac{5g_2}{768\pi^2 M^2} (\mathcal{Y}\mathcal{Y}^\dagger y_e)^{pr} - \frac{g_2}{128\pi^2 M^2} \mathcal{Y}^{pr} (\mathcal{Y}^\dagger)^{st} y_e^{ts},$$

$$c_{eH}^{pr} \rightarrow -\frac{g_2^4}{3840\pi^2 M^2} y_e^{pr} + \frac{1}{192\pi^2 M^2} (y_e y_e^\dagger \mathcal{Y}\mathcal{Y}^\dagger y_e + 2y_e \mathcal{Y}^\dagger \mathcal{Y} y_e^\dagger y_e - 3\mathcal{Y} y_e^\dagger y_e \mathcal{Y}^\dagger y_e - 3y_e \mathcal{Y}^\dagger y_e y_e^\dagger \mathcal{Y})^{pr} + \frac{1}{16\pi^2 M^2} \mathcal{Y}^{pr} y_e^{st} (y_e^\dagger)^{tu} y_e^{uv} (\mathcal{Y}^\dagger)^{vs} - \frac{\lambda}{32\pi^2 M^2} \mathcal{Y}^{pr} (\mathcal{Y}^\dagger)^{st} y_e^{ts}.$$



Cross-check with

Matchete [arXiv:2212.04510]

and

MatchMakerEFT [arXiv:2112.10787]

A second example: *reduction of Green's basis*

Green's basis: minimal set of operators to match amplitudes off-shell.

Physical basis: minimal set of operators to match amplitudes on-shell.

Field
redefinitions

Tedious
Non systematic
Hard to automate

$\mathcal{L}_{\text{full}} \supset$ operators in Green's basis

$\mathcal{L}_{\text{phys}} \supset$ operators in physical basis

Match both theories on-shell at tree level



» We find the couplings of $\mathcal{L}_{\text{phys}}$ in terms of those in $\mathcal{L}_{\text{full}}$.

**REDUCTION OF
GREEN'S BASIS**

■ **Example I:** Real scalar singlet with Z_2 symmetry up to dimension 8

■ **Example II:** Bosonic sector in the SMEFT up to dimension 8

A second example: *reduction of Green's basis*

BOSONIC SECTOR



Cross-check with [\[2003.12525v5\]](#)

$X^2 H^2$

$$\begin{aligned} c_{HB} &\rightarrow c_{HB} - 2m_0^2 c_{HB} r_{DH} \\ c_{H\tilde{B}} &\rightarrow c_{H\tilde{B}} - 2m_0^2 c_{H\tilde{B}} r_{DH} \end{aligned}$$

$H^4 D^2$

$$\begin{aligned} c_{H\Box} &\rightarrow c_{H\Box} - \frac{1}{8}g'^2 r_{2B} + \frac{1}{2}g' r_{BDH} - m_0^2(4c_{H\Box} r_{DH} + g' r_{BDH} r_{DH} + 2r_{DH} r'_{HD}) \\ c_{HD} &\rightarrow c_{HD} - \frac{1}{2}g'^2 r_{2B} + 2g' r_{BDH} - m_0^2(4c_{HD} r_{DH} + 4g' r_{BDH} r_{DH}) \end{aligned}$$

H^6

$$\begin{aligned} c_H &\rightarrow c_H + \lambda^2 r_{DH} + \lambda r'_{HD} + m_0^2 \left(\frac{1}{4}g'^2 c_{HD} r_{2B} - \frac{1}{16}g'^4 r_{2B}^2 - \frac{1}{2}g' c_{HD} r_{BDH} + \frac{1}{2}g'^3 r_{2B} r_{BDH} \right. \\ &\quad \left. - \frac{3}{4}g'^2 r_{BDH}^2 - 6c_H r_{DH} - \lambda c_{HD} r_{DH} + 8\lambda c_{H\Box} r_{DH} + g' \lambda r_{BDH} r_{DH} - 11\lambda^2 r_{DH}^2 \right. \\ &\quad \left. - \frac{1}{2}c_{HD} r'_{HD} + 4c_{H\Box} r'_{HD} + \frac{1}{2}g' r_{BDH} r'_{HD} - 9\lambda r_{DH} r'_{HD} - \frac{1}{4}r_{HD}^2 - r_{HD}''^2 \right) \end{aligned}$$

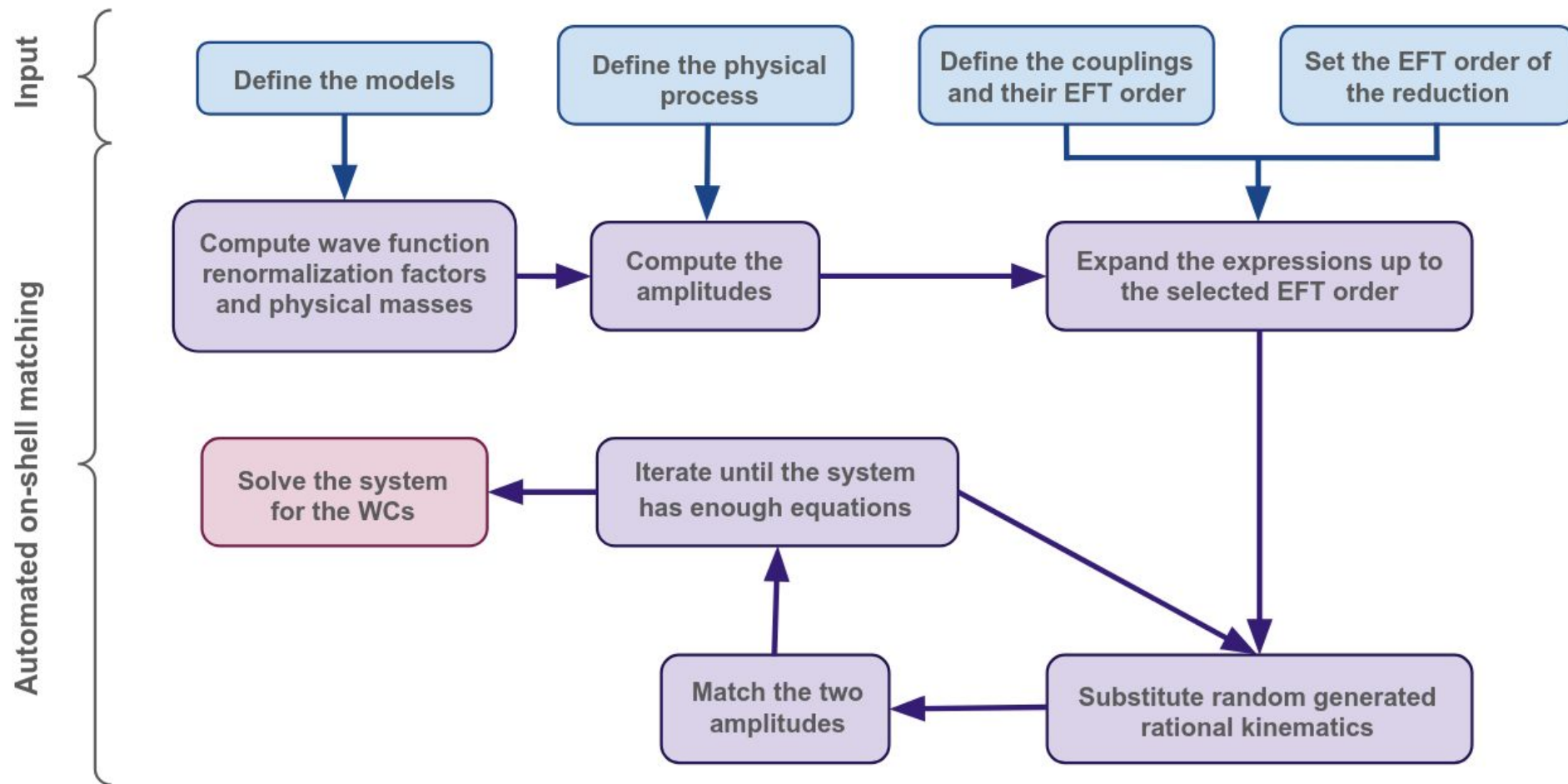
X^3		$X^2 H^2$		$H^2 D^4$	
$\mathcal{O}_{3G}$	$f^{ABC} G_{\mu}^{A\nu} G_{\nu}^{B\rho} G_{\rho}^{C\mu}$	$\mathcal{O}_{HG}$	$G_{\mu\nu}^A G^{A\mu\nu} (H^\dagger H)$	$\mathcal{O}_{DH}$	$(D_\mu D^\mu H)^\dagger (D_\nu D^\nu H)$
$\mathcal{O}_{3\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu}^{A\nu} G_{\nu}^{B\rho} G_{\rho}^{C\mu}$	$\mathcal{O}_{H\tilde{G}}$	$\tilde{G}_{\mu\nu}^A G^{A\mu\nu} (H^\dagger H)$	$H^4 D^2$	
$\mathcal{O}_{3W}$	$\epsilon^{IJK} W_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$	$\mathcal{O}_{HW}$	$W_{\mu\nu}^I W^{I\mu\nu} (H^\dagger H)$	$\mathcal{O}_{H\Box}$	$(H^\dagger H) \Box (H^\dagger H)$
$\mathcal{O}_{3\tilde{W}}$	$\epsilon^{IJK} \tilde{W}_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$	$\mathcal{O}_{H\tilde{W}}$	$\tilde{W}_{\mu\nu}^I W^{I\mu\nu} (H^\dagger H)$	$\mathcal{O}_{HD}$	$(H^\dagger D^\mu H)^\dagger (H^\dagger D_\mu H)$
$X^2 D^2$		$\mathcal{O}_{HB}$	$B_{\mu\nu} B^{\mu\nu} (H^\dagger H)$	$\mathcal{O}'_{HD}$	$(H^\dagger H) (D_\mu H)^\dagger (\overleftrightarrow{D}^\mu H)$
$\mathcal{O}_{2G}$	$-\frac{1}{2}(D_\mu G^{A\mu\nu})(D^\rho G_{\rho\nu}^A)$	$\mathcal{O}_{H\tilde{B}}$	$\tilde{B}_{\mu\nu} B^{\mu\nu} (H^\dagger H)$	$\mathcal{O}''_{HD}$	$(H^\dagger H) D_\mu (H^\dagger \overleftrightarrow{D}^\mu H)$
$\mathcal{O}_{2W}$	$-\frac{1}{2}(D_\mu W^{I\mu\nu})(D^\rho W_{\rho\nu}^I)$	$\mathcal{O}_{HWB}$	$W_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$	H^6	
$\mathcal{O}_{2B}$	$-\frac{1}{2}(\partial_\mu B^{\mu\nu})(\partial^\rho B_{\rho\nu})$	$\mathcal{O}_{H\tilde{W}B}$	$\tilde{W}_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$	$\mathcal{O}_H$	$(H^\dagger H)^3$
		$H^2 X D^2$			
		$\mathcal{O}_{WDH}$	$D_\nu W^{I\mu\nu} (H^\dagger \overleftrightarrow{D}_\mu^I H)$		
		$\mathcal{O}_{BDH}$	$\partial_\nu B^{\mu\nu} (H^\dagger \overleftrightarrow{D}_\mu H)$		

V. Gherardi, D. Marzocca y E. Venturini (2021) [\[2003.12525v5\]](#)

Reduction of dim. 6 operators up to dim. 8!!

Creating a tool: *MOSCA*

MOSCA: A **M**ATCHING **O**N-SHELL **C**ALCULATOR



Creating a tool: *MOSCA*

MOSCA: A **M**ATCHING **O**N-**S**HELL **C**ALCULATOR

Let us consider a Green's basis for a Z_2 -symmetric scalar theory up to dimension 8. A suitable choice is given by

$$\mathcal{L} = \mathcal{L}_4 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \frac{1}{\Lambda^4} \mathcal{L}_8, \quad (10)$$

where

$$\mathcal{L}_4 = -\frac{1}{2} \phi (\partial^2 + m^2) \phi - \lambda \phi^4, \quad (11)$$

$$\mathcal{L}_6 = \alpha_{61} \phi^6 + \beta_{61} \partial^2 \phi \partial^2 \phi + \beta_{62} \phi^3 \partial^2 \phi, \quad (12)$$

$$\mathcal{L}_8 = \alpha_{81} \phi^8 + \alpha_{82} \phi^2 \partial_\mu \partial_\nu \phi \partial^\mu \partial^\nu \phi + \beta_{81} \phi \partial^6 \phi + \beta_{82} \phi^3 \partial^4 \phi + \beta_{83} \phi^2 \partial^2 \phi \partial^2 \phi + \beta_{84} \phi^5 \partial^2 \phi. \quad (13)$$

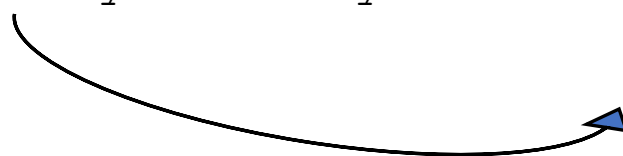
β are the redundant WCs (they are in the Green's basis but not in the physical one)

Conclusions

- ★ On-shell matching is a **diagrammatic alternative** that allows to compute the **coefficients directly into the physical basis**.
- ★ **More diagrams** to be computed in a single amplitude, but **several operators can be matched** at the **same time**.
- ★ **No** need of defining **evanescent operators** or using **background field method** for finite matching.
- ★ Useful for **renormalization** and computing **beta functions** (see [arXiv:2409.15408]).
- ★ Can be used to compute the **reduction to any physical basis** from any redundant basis.



Code in Mathematica in progress (**MOSCA**)
based in FeynRules+FeynArts+FeynCalc
(with F. Vilches)



FIXED BUGS

- $(p_i \cdot p_j)^n$ structures
- $\epsilon^{\mu\nu\rho\sigma}\gamma_\mu$ structures
- Duplicate components in 'FFSS'



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THANKS FOR YOUR ATTENTION !

Green's basis vs physical basis

EFT Lagrangian :
$$\mathcal{L} = \mathcal{L}^{(4)} + \sum_{d>4} \sum_i \frac{c_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^{(d)}$$

Valid operators

Local operators

Preserve the symmetries of the Lagrangian

Finite number of operators

Integration by parts, Fierz identities (4D), etc.

Green's basis

Some operators are still **redundant** on-shell

Physical basis

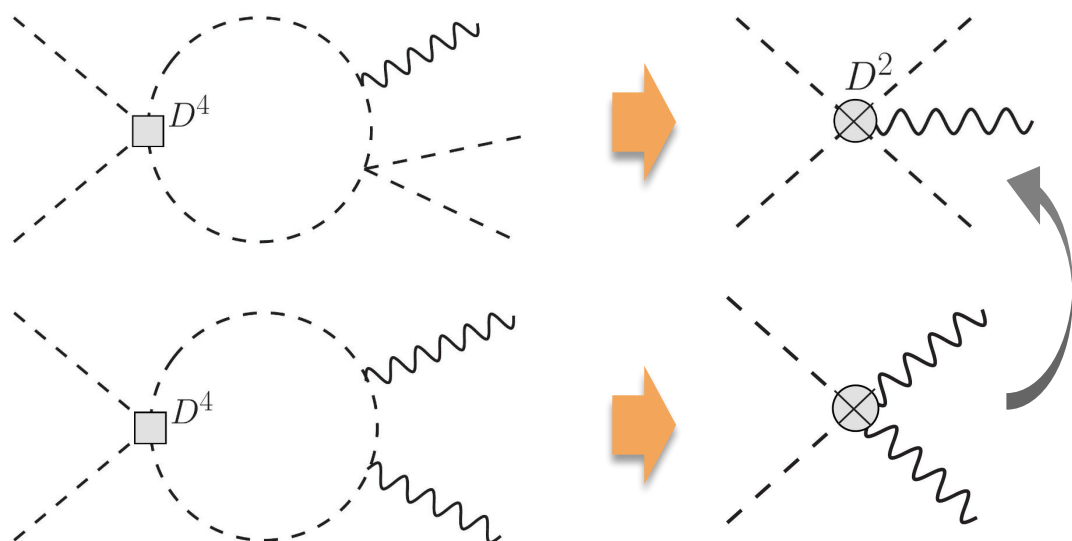
Applications: on-shell RGEs

A theory expressed in terms of the physical basis can still generate redundant operators at the loop level through RGEs.

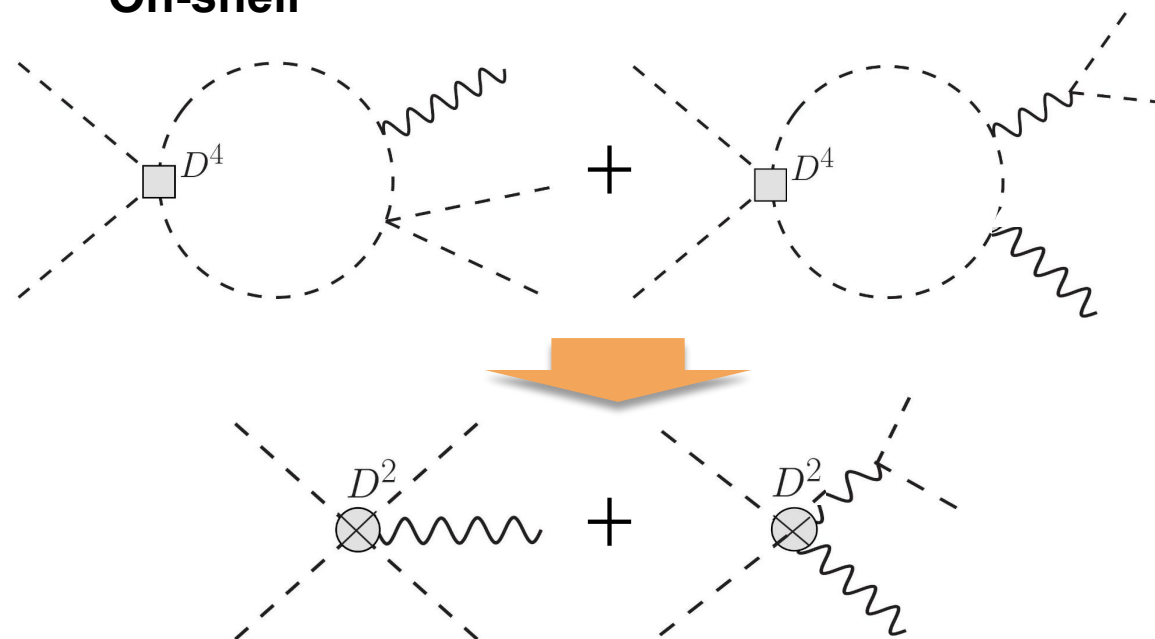
$$\mathcal{L}^{\text{div}} = \sum_j \frac{1}{16\pi^2\epsilon} c'_j(c) \mathcal{O}_j \quad \longrightarrow \quad 16\pi^2\mu \frac{dc_i}{d\mu} = -2c'_i \quad (\text{RGEs})$$

Mixing of $H^4 D^4$ into $H^4 B D^2$

Off-shell



On-shell



Some results in the SMEFT

Dimension 4 up to dimension 8 (H and B)

$$m_0^2 \rightarrow m_0^2 - m_0^4 r_{DH} + 2m_0^6 r_{DH}^2$$

$$\lambda \rightarrow \lambda - m_0^2(4\lambda r_{DH} + 2r'_{HD}) + m_0^4(16\lambda r_{DH}^2 + 10r_{DH}r'_{HD})$$

$$y_E \rightarrow y_E(1 - m_0^2 r_{DH} + \frac{5}{2}m_0^4 r_{DH}^2)$$

Some results in the SMEFT

Dimension 6 up to dimension 8 (H and B)

$$X^2 H^2 \quad \begin{aligned} c_{HB} &\rightarrow c_{HB} - 2m_0^2 c_{HB} r_{DH} \\ c_{H\tilde{B}} &\rightarrow c_{H\tilde{B}} - 2m_0^2 c_{H\tilde{B}} r_{DH} \end{aligned}$$

$$H^4 D^2 \quad \begin{aligned} c_{H\Box} &\rightarrow c_{H\Box} - \frac{1}{8}g'^2 r_{2B} + \frac{1}{2}g' r_{BDH} - m_0^2(4c_{H\Box} r_{DH} + g' r_{BDH} r_{DH} + 2r_{DH} r'_{HD}) \\ c_{HD} &\rightarrow c_{HD} - \frac{1}{2}g'^2 r_{2B} + 2g' r_{BDH} - m_0^2(4c_{HD} r_{DH} + 4g' r_{BDH} r_{DH}) \end{aligned}$$

$$H^6 \quad \begin{aligned} c_H &\rightarrow c_H + \lambda^2 r_{DH} + \lambda r'_{HD} + m_0^2 \left(\frac{1}{4}g'^2 c_{HD} r_{2B} - \frac{1}{16}g'^4 r_{2B}^2 - \frac{1}{2}g' c_{HD} r_{BDH} + \frac{1}{2}g'^3 r_{2B} r_{BDH} \right. \\ &\quad - \frac{3}{4}g'^2 r_{BDH}^2 - 6c_H r_{DH} - \lambda c_{HD} r_{DH} + 8\lambda c_{H\Box} r_{DH} + g' \lambda r_{BDH} r_{DH} - 11\lambda^2 r_{DH}^2 \\ &\quad \left. - \frac{1}{2}c_{HD} r'_{HD} + 4c_{H\Box} r'_{HD} + \frac{1}{2}g' r_{BDH} r'_{HD} - 9\lambda r_{DH} r'_{HD} - \frac{1}{4}r_{HD}^2 - r_{HD}^{\prime 2} \right) \end{aligned}$$

Some results in the SMEFT

Dimension 8 (H and B)

$$XH^4D^2$$

$$c_{BH^4D^2}^{(1)} \rightarrow c_{BH^4D^2}^{(1)} - 4g'c_{HB}r_{2B} + \frac{1}{2}g'^3r_{2B}^2 + 8c_{HB}r_{BDH} - 2g'^2r_{2B}r_{BDH} + 2g'r_{BDH}^2$$

$$c_{BH^4D^2}^{(2)} \rightarrow -4g'c_{H\tilde{B}}r_{2B} + 8c_{H\tilde{B}}r_{BDH}$$

$$X^2H^2D^2$$

$$c_{B^2H^2D^2}^{(1)} \rightarrow 0$$

$$c_{B^2H^2D^2}^{(2)} \rightarrow 0$$

$$c_{B^2H^2D^2}^{(3)} \rightarrow 0$$

$$X^4$$

$$c_{B^4}^{(1)} \rightarrow 0$$

$$c_{B^4}^{(2)} \rightarrow 0$$

$$c_{B^4}^{(3)} \rightarrow 0$$

$$X^2H^4$$

$$c_{B^2H^4}^{(1)} \rightarrow -c_{HB}g'^2r_{2B} + \frac{1}{16}g'^4r_{2B}^2 + 2c_{HB}g'r_{BDH} - \frac{1}{4}g'^3r_{2B}r_{BDH} + \frac{1}{4}g'^2r_{BDH}^2 - 2c_{HB}\lambda r_{DH} - c_{HB}r'_{HD}$$

$$c_{B^2H^4}^{(2)} \rightarrow -g'^2c_{H\tilde{B}}r_{2B} + 2g'c_{H\tilde{B}}r_{BDH} - 2\lambda c_{H\tilde{B}}r_{BDH} - c_{H\tilde{B}}r'_{HD}$$

Some results in the SMEFT

Dimension 8 (H and B)

$H^4 D^4$

$$c_{H^4}^{(1)} \rightarrow \frac{1}{2}g'^2 r_{2B}^2 - 2g' r_{2B} r_{BDH} + 2r_{BDH}^2$$

$$c_{H^4}^{(2)} \rightarrow -\frac{1}{2}g'^2 r_{2B}^2 + 2g' r_{2B} r_{BDH} - 2r_{BDH}^2$$

$$c_{H^4}^{(3)} \rightarrow 0$$

$H^6 D^2$

$$\begin{aligned} c_{H^6}^{(1)} \rightarrow & -\frac{3}{4}g'^2 c_{HD} r_{2B} + \frac{3}{16}g'^4 r_{2B}^2 + \frac{3}{2}g' c_{HD} r_{BDH} - \frac{3}{2}g'^3 r_{2B} r_{BDH} + \frac{9}{4}g'^2 r_{BDH}^2 - \lambda c_{HD} r_{DH} \\ & - 8\lambda c_{H\Box} r_{DH} - 3g' \lambda r_{BDH} r_{DH} + \lambda^2 r_{DH}^2 - \frac{1}{2}c_{HD} r'_{HD} - 4c_{H\Box} r'_{HD} - \frac{3}{2}g' r_{BDH} r'_{HD} \\ & - 3\lambda r_{DH} r'_{HD} - \frac{7}{4}r_{HD}^{\prime 2} + r_{HD}^{\prime\prime 2} \end{aligned}$$

$$\begin{aligned} c_{H^6}^{(2)} \rightarrow & -\frac{1}{2}g'^2 c_{HD} r_{2B} + \frac{1}{8}g'^4 r_{2B}^2 + g' c_{HD} r_{BDH} - g'^3 r_{2B} r_{BDH} + \frac{3}{2}g'^2 r_{BDH}^2 - 2\lambda c_{HD} r_{DH} \\ & - 2g' \lambda r_{BDH} r_{DH} - c_{HD} r'_{HD} - g' r_{BDH} r'_{HD} \end{aligned}$$

Generation of random momenta

$$SL(2, \mathbb{C}) \cong SU(2)_L \times SU(2)_R \left\{ \begin{array}{l} \lambda \in SU(2)_L \\ \tilde{\lambda} \in SU(2)_R \end{array} \right\} \quad \begin{array}{l} \lambda^\alpha = \varepsilon^{\alpha\beta} \lambda_\beta \\ \tilde{\lambda}_{\dot{\alpha}} = \varepsilon_{\dot{\alpha}\dot{\beta}} \tilde{\lambda}^{\dot{\beta}} \end{array}$$

**Massless
momenta :**

$$P_{\alpha\dot{\alpha}} = \lambda_\alpha \tilde{\lambda}_{\dot{\alpha}} \quad \Rightarrow \quad P = p_\mu \sigma^\mu = \begin{pmatrix} p_0 + p_3 & p_1 - ip_2 \\ p_1 + ip_2 & p_0 - p_3 \end{pmatrix}$$

**Massive
momenta :**

$$P^\mu := q^\mu + \frac{m^2}{2q \cdot k} k^\mu \quad \left| \begin{array}{l} q^2, k^2 = 0 \\ q_{\alpha\dot{\alpha}} = \lambda_\alpha \tilde{\lambda}_{\dot{\alpha}} \\ k_{\alpha\dot{\alpha}} = \mu_\alpha \tilde{\mu}_{\dot{\alpha}} \end{array} \right.$$

Evanescent operators

$$\mathcal{R} = \alpha \mathcal{O} \quad \xrightarrow{d = 4 - 2\epsilon} \quad \mathcal{R} = \alpha \mathcal{O} + \mathcal{E}$$

$$IR^{(0)} + IR^{(1)} = UV^{(0)} + UV^{(1)}$$

$$\mathcal{O}(\epsilon)$$

Additional finite local contributions in loop amplitudes

$$IR^{(0)} + \boxed{IR_{soft}^{(1)}} = UV^{(0)} + UV_{hard}^{(1)} + \boxed{UV_{soft}^{(1)}}$$

We take the hard region

$$\int \mathcal{O} = \frac{1}{\epsilon} (a + b_{\mathcal{O}} \epsilon)$$

$$\int \mathcal{R} = \frac{1}{\epsilon} (a + b_{\mathcal{R}} \epsilon)$$

$$\int \mathcal{R} - \mathcal{O} = \frac{1}{\epsilon} (b_{\mathcal{R}} \epsilon - b_{\mathcal{O}} \epsilon) = b$$

Matching a theory: *off-shell* vs *on-shell*

OFF-SHELL

VS

ON-SHELL

Small number of diagrams
(1lPI in UV, 1PI in IR)

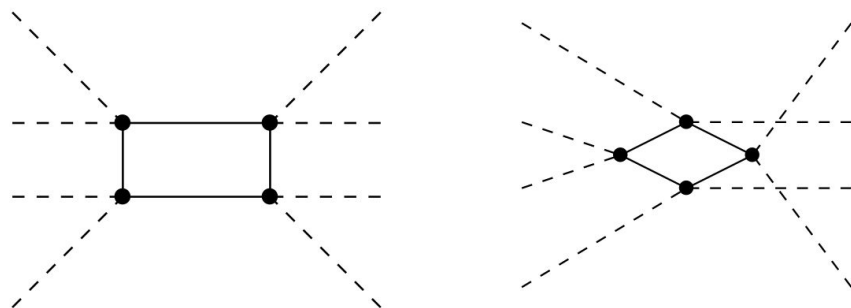
All diagrams (light bridges too)

Matching a theory: *off-shell* vs *on-shell*

$$\mathcal{L}_{\text{int}} = \eta \phi^2 \Phi^2 + \lambda \phi^4$$

OFF-SHELL

$\phi \quad \phi \quad \phi \quad \phi \rightarrow \phi \quad \phi \quad \phi \quad \phi$

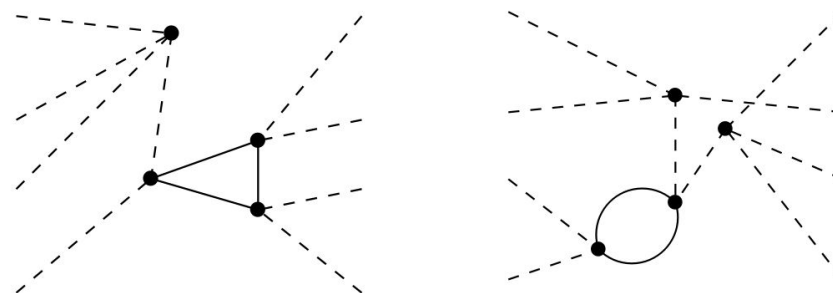


```
In[66]:= Length[diags1lPIHeavy]
```

```
Out[66]= 315
```

ON-SHELL

$\phi \quad \phi \quad \phi \quad \phi \rightarrow \phi \quad \phi \quad \phi \quad \phi$



```
In[65]:= Length[diagsHeavy]
```

```
Out[65]= 6475
```

Matching a theory: *off-shell* vs *on-shell*

OFF-SHELL

vs

ON-SHELL

Small number of diagrams
(one-light-particle-irreducible)

1 — 0

All diagrams (light bridges too)

Matching a theory: *off-shell* vs *on-shell*

OFF-SHELL

VS

ON-SHELL

Small number of diagrams
(one-light-particle-irreducible)

1 — 0

All diagrams (light bridges too)

Large set of operators
(Green's basis)

Smaller set of operators
(physical basis)

Matching a theory: off-shell vs on-shell

[2003.12525]

$\psi^2 D^3$		$\psi^2 X D$		$\psi^2 D H^2$	
$\mathcal{O}_{qD}$	$\frac{i}{2} \bar{q} \{D_\mu D^\mu, \not{D}\} q$	$\mathcal{O}_{Gq}$	$(\bar{q} T^A \gamma^\mu q) D^\nu G_{\mu\nu}^A$	$\mathcal{O}_{Hq}^{(1)}$	$(\bar{q} \gamma^\mu q) (H^\dagger i \overleftrightarrow{D}_\mu H)$
$\mathcal{O}_{uD}$	$\frac{i}{2} \bar{u} \{D_\mu D^\mu, \not{D}\} u$	$\mathcal{O}'_{Gq}$	$\frac{1}{2} (\bar{q} T^A \gamma^\mu i \overleftrightarrow{D}^\nu q) G_{\mu\nu}^A$	$\mathcal{O}'_{Hq}^{(1)}$	$(\bar{q} i \not{D} q) (H^\dagger H)$
$\mathcal{O}_{dD}$	$\frac{i}{2} \bar{d} \{D_\mu D^\mu, \not{D}\} d$	$\mathcal{O}''_{Gq}$	$\frac{1}{2} (\bar{q} T^A \gamma^\mu i \overleftrightarrow{D}^\nu q) \tilde{G}_{\mu\nu}^A$	$\mathcal{O}''_{Hq}^{(1)}$	$(\bar{q} \gamma^\mu q) \partial_\mu (H^\dagger H)$
$\mathcal{O}_{\ell D}$	$\frac{i}{2} \bar{\ell} \{D_\mu D^\mu, \not{D}\} \ell$	$\mathcal{O}_{Wq}$	$(\bar{q} \sigma^I \gamma^\mu q) D^\nu W_{\mu\nu}^I$	$\mathcal{O}_{Hq}^{(3)}$	$(\bar{q} \sigma^I \gamma^\mu q) (H^\dagger i \overleftrightarrow{D}_\mu^I H)$
$\mathcal{O}_{eD}$	$\frac{i}{2} \bar{e} \{D_\mu D^\mu, \not{D}\} e$	$\mathcal{O}'_{Wq}$	$\frac{1}{2} (\bar{q} \sigma^I \gamma^\mu i \overleftrightarrow{D}^\nu q) W_{\mu\nu}^I$	$\mathcal{O}''_{Hq}^{(3)}$	$(\bar{q} i \not{D}^I q) (H^\dagger \sigma^I H)$
$\psi^2 H D^2 + \text{h.c.}$		$\mathcal{O}'_{\tilde{W}q}$	$\frac{1}{2} (\bar{q} \sigma^I \gamma^\mu i \overleftrightarrow{D}^\nu q) \tilde{W}_{\mu\nu}^I$	$\mathcal{O}'''_{Hq}^{(3)}$	$(\bar{q} \sigma^I \gamma^\mu q) D_\mu (H^\dagger \sigma^I H)$
$\mathcal{O}_{uHD1}$	$(\bar{q} u) D_\mu D^\mu \tilde{H}$	$\mathcal{O}_{Bq}$	$(\bar{q} \gamma^\mu q) \partial^\nu B_{\mu\nu}$	$\mathcal{O}_{Hu}$	$(\bar{u} \gamma^\mu u) (H^\dagger i \overleftrightarrow{D}_\mu H)$
$\mathcal{O}_{uHD2}$	$(\bar{q} i \sigma_{\mu\nu} D^\mu u) D^\nu \tilde{H}$	$\mathcal{O}'_{Bq}$	$\frac{1}{2} (\bar{q} \gamma^\mu i \overleftrightarrow{D}^\nu q) B_{\mu\nu}$	$\mathcal{O}'_{Hu}$	$(\bar{u} i \not{D} u) (H^\dagger H)$
$\mathcal{O}_{uHD3}$	$(\bar{q} D_\mu D^\mu u) \tilde{H}$	$\mathcal{O}''_{Bq}$	$\frac{1}{2} (\bar{q} \gamma^\mu i \overleftrightarrow{D}^\nu q) \tilde{B}_{\mu\nu}$	$\mathcal{O}''_{Hu}$	$(\bar{u} \gamma^\mu u) \partial_\mu (H^\dagger H)$
$\mathcal{O}_{uHD4}$	$(\bar{q} D_\mu u) D^\mu \tilde{H}$	$\mathcal{O}_{Gu}$	$(\bar{u} T^A \gamma^\mu u) D^\nu G_{\mu\nu}^A$	$\mathcal{O}_{Hd}$	$(\bar{d} \gamma^\mu d) (H^\dagger i \overleftrightarrow{D}_\mu H)$
$\mathcal{O}_{dHD1}$	$(\bar{q} d) D_\mu D^\mu H$	$\mathcal{O}'_{Gu}$	$\frac{1}{2} (\bar{u} T^A \gamma^\mu i \overleftrightarrow{D}^\nu u) G_{\mu\nu}^A$	$\mathcal{O}'_{Hd}$	$(\bar{d} i \not{D} d) (H^\dagger H)$
$\mathcal{O}_{dHD2}$	$(\bar{q} i \sigma_{\mu\nu} D^\mu d) D^\nu H$	$\mathcal{O}'_{\tilde{Gu}}$	$\frac{1}{2} (\bar{u} T^A \gamma^\mu i \overleftrightarrow{D}^\nu u) \tilde{G}_{\mu\nu}^A$	$\mathcal{O}''_{Hd}$	$(\bar{d} \gamma^\mu d) \partial_\mu (H^\dagger H)$
$\mathcal{O}_{dHD3}$	$(\bar{q} D_\mu D^\mu d) H$	$\mathcal{O}_{Bu}$	$(\bar{u} \gamma^\mu u) \partial^\nu B_{\mu\nu}$	$\mathcal{O}_{Hud}$	$(\bar{u} \gamma^\mu d) (H^\dagger i D_\mu H)$
$\mathcal{O}_{dHD4}$	$(\bar{q} D_\mu d) D^\mu H$	$\mathcal{O}'_{Bu}$	$\frac{1}{2} (\bar{u} \gamma^\mu i \overleftrightarrow{D}^\nu u) B_{\mu\nu}$	$\mathcal{O}_{H\ell}^{(1)}$	$(\bar{\ell} \gamma^\mu \ell) (H^\dagger i \overleftrightarrow{D}_\mu H)$
$\mathcal{O}_{eHD1}$	$(\bar{\ell} e) D_\mu D^\mu H$	$\mathcal{O}'_{\tilde{Bu}}$	$\frac{1}{2} (\bar{u} \gamma^\mu i \overleftrightarrow{D}^\nu u) \tilde{B}_{\mu\nu}$	$\mathcal{O}'_{H\ell}^{(1)}$	$(\bar{\ell} i \not{D} \ell) (H^\dagger H)$
$\mathcal{O}_{eHD2}$	$(\bar{\ell} i \sigma_{\mu\nu} D^\mu e) D^\nu H$	$\mathcal{O}_{Gd}$	$(\bar{d} T^A \gamma^\mu d) D^\nu G_{\mu\nu}^A$	$\mathcal{O}''_{H\ell}^{(1)}$	$(\bar{\ell} \gamma^\mu \ell) \partial_\mu (H^\dagger H)$
$\mathcal{O}_{eHD3}$	$(\bar{\ell} D_\mu D^\mu e) H$	$\mathcal{O}'_{Gd}$	$\frac{1}{2} (\bar{d} T^A \gamma^\mu i \overleftrightarrow{D}^\nu d) G_{\mu\nu}^A$	$\mathcal{O}_{H\ell}^{(3)}$	$(\bar{\ell} \sigma^I \gamma^\mu \ell) (H^\dagger i \overleftrightarrow{D}_\mu^I H)$
$\mathcal{O}_{eHD4}$	$(\bar{\ell} D_\mu e) D^\mu H$	$\mathcal{O}'_{\tilde{Gd}}$	$\frac{1}{2} (\bar{d} T^A \gamma^\mu i \overleftrightarrow{D}^\nu d) \tilde{G}_{\mu\nu}^A$	$\mathcal{O}'_{H\ell}^{(3)}$	$(\bar{\ell} i \not{D}^I \ell) (H^\dagger \sigma^I H)$
$\psi^2 X H + \text{h.c.}$		$\mathcal{O}_{Bd}$	$(\bar{d} \gamma^\mu d) \partial^\nu B_{\mu\nu}$	$\mathcal{O}''_{H\ell}^{(3)}$	$(\bar{\ell} \sigma^I \gamma^\mu \ell) D_\mu (H^\dagger \sigma^I H)$
$\mathcal{O}_{uG}$	$(\bar{q} T^A \sigma^{\mu\nu} u) \tilde{H} G_{\mu\nu}^A$	$\mathcal{O}'_{Bd}$	$\frac{1}{2} (\bar{d} \gamma^\mu i \overleftrightarrow{D}^\nu d) B_{\mu\nu}$	$\mathcal{O}_{He}$	$(\bar{e} \gamma^\mu e) (H^\dagger i \overleftrightarrow{D}_\mu H)$
$\mathcal{O}_{uW}$	$(\bar{q} \sigma^{\mu\nu} u) \sigma^I \tilde{H} W_{\mu\nu}^I$	$\mathcal{O}'_{\tilde{Bd}}$	$\frac{1}{2} (\bar{d} \gamma^\mu i \overleftrightarrow{D}^\nu d) \tilde{B}_{\mu\nu}$	$\mathcal{O}'_{He}$	$(\bar{e} i \not{D} e) (H^\dagger H)$
$\mathcal{O}_{uB}$	$(\bar{q} \sigma^{\mu\nu} u) \tilde{H} B_{\mu\nu}$	$\mathcal{O}_{W\ell}$	$(\bar{\ell} \sigma^I \gamma^\mu \ell) D^\nu W_{\mu\nu}^I$	$\mathcal{O}''_{He}$	$(\bar{e} \gamma^\mu e) \partial_\mu (H^\dagger H)$
$\mathcal{O}_{dG}$	$(\bar{q} T^A \sigma^{\mu\nu} d) H G_{\mu\nu}^A$	$\mathcal{O}'_{W\ell}$	$\frac{1}{2} (\bar{\ell} \sigma^I \gamma^\mu i \overleftrightarrow{D}^\nu \ell) W_{\mu\nu}^I$	$\psi^2 H^3 + \text{h.c.}$	
$\mathcal{O}_{dW}$	$(\bar{q} \sigma^{\mu\nu} d) \sigma^I H W_{\mu\nu}^I$	$\mathcal{O}'_{\tilde{W}\ell}$	$\frac{1}{2} (\bar{\ell} \sigma^I \gamma^\mu i \overleftrightarrow{D}^\nu \ell) \tilde{W}_{\mu\nu}^I$	$\mathcal{O}_{uH}$	$(H^\dagger H) \bar{q} \tilde{H} u$
$\mathcal{O}_{dB}$	$(\bar{q} \sigma^{\mu\nu} d) H B_{\mu\nu}$	$\mathcal{O}_{B\ell}$	$(\bar{\ell} \gamma^\mu \ell) \partial^\nu B_{\mu\nu}$	$\mathcal{O}_{dH}$	$(H^\dagger H) \bar{q} H d$
$\mathcal{O}_{eW}$	$(\bar{\ell} \sigma^{\mu\nu} e) \sigma^I H W_{\mu\nu}^I$	$\mathcal{O}'_{B\ell}$	$\frac{1}{2} (\bar{\ell} \gamma^\mu i \overleftrightarrow{D}^\nu \ell) B_{\mu\nu}$	$\mathcal{O}_{eH}$	$(H^\dagger H) \bar{\ell} H e$
$\mathcal{O}_{eB}$	$(\bar{\ell} \sigma^{\mu\nu} e) H B_{\mu\nu}$	$\mathcal{O}'_{\tilde{B}\ell}$	$\frac{1}{2} (\bar{\ell} \gamma^\mu i \overleftrightarrow{D}^\nu \ell) \tilde{B}_{\mu\nu}$		
		$\mathcal{O}_{Be}$	$(\bar{e} \gamma^\mu e) \partial^\nu B_{\mu\nu}$		
		$\mathcal{O}'_{Be}$	$\frac{1}{2} (\bar{e} \gamma^\mu i \overleftrightarrow{D}^\nu e) B_{\mu\nu}$		
		$\mathcal{O}'_{\tilde{Be}}$	$\frac{1}{2} (\bar{e} \gamma^\mu i \overleftrightarrow{D}^\nu e) \tilde{B}_{\mu\nu}$		

X^3		$X^2 H^2$		$H^2 D^4$	
$\mathcal{O}_{3G}$	$f^{ABC} G_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	$\mathcal{O}_{HG}$	$G_{\mu\nu}^A G^{A\mu\nu} (H^\dagger H)$	$\mathcal{O}_{DH}$	$(D_\mu D^\mu H)^\dagger (D_\nu D^\nu H)$
$\mathcal{O}_{3\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	$\mathcal{O}_{H\tilde{G}}$	$\tilde{G}_{\mu\nu}^A G^{A\mu\nu} (H^\dagger H)$	$H^4 D^2$	
$\mathcal{O}_{3W}$	$\epsilon^{IJK} W_{\mu\nu}^I W_{\nu\rho}^J W_{\rho\mu}^K$	$\mathcal{O}_{HW}$	$W_{\mu\nu}^I W^{I\mu\nu} (H^\dagger H)$	$\mathcal{O}_{H\Box}$	$(H^\dagger H) \Box (H^\dagger H)$
$\mathcal{O}_{3\tilde{W}}$	$\epsilon^{IJK} \tilde{W}_{\mu\nu}^I W_{\nu\rho}^J W_{\rho\mu}^K$	$\mathcal{O}_{H\tilde{W}}$	$\tilde{W}_{\mu\nu}^I W^{I\mu\nu} (H^\dagger H)$	$\mathcal{O}_{HD}$	$(H^\dagger D^\mu H)^\dagger (H^\dagger D_\mu H)$
$X^2 D^2$		$\mathcal{O}_{HB}$	$B_{\mu\nu} B^{\mu\nu} (H^\dagger H)$	$\mathcal{O}'_{HD}$	$(H^\dagger H) (D_\mu H)^\dagger (D^\mu H)$
$\mathcal{O}_{2G}$	$-\frac{1}{2} (D_\mu G^{A\mu\nu}) (D^\rho G_{\rho\nu}^A)$	$\mathcal{O}_{H\tilde{B}}$	$\tilde{B}_{\mu\nu} B^{\mu\nu} (H^\dagger H)$	$\mathcal{O}''_{HD}$	$(H^\dagger H) D_\mu (H^\dagger i \overleftrightarrow{D}^\mu H)$
$\mathcal{O}_{2W}$	$-\frac{1}{2} (D_\mu W^{I\mu\nu}) (D^\rho W_{\rho\nu}^I)$	$\mathcal{O}_{HWB}$	$W_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$	H^6	
$\mathcal{O}_{2B}$	$-\frac{1}{2} (\partial_\mu B^{\mu\nu}) (\partial^\rho B_{\rho\nu})$	$\mathcal{O}_{H\tilde{W}B}$	$\tilde{W}_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$	$\mathcal{O}_H$	$(H^\dagger H)^3$
		$H^2 X D^2$			
		$\mathcal{O}_{WDH}$	$D_\nu W^{I\mu\nu} (H^\dagger i \overleftrightarrow{D}_\mu^I H)$		
		$\mathcal{O}_{BDH}$	$\partial_\nu B^{\mu\nu} (H^\dagger i \overleftrightarrow{D}_\mu H)$		

Four-quark		Four-lepton		Semileptonic	
$\mathcal{O}_{qq}^{(1)}$	$(\bar{q} \gamma^\mu q) (\bar{q} \gamma_\mu q)$	$\mathcal{O}_{\ell\ell}$	$(\bar{\ell} \gamma^\mu \ell) (\bar{\ell} \gamma_\mu \ell)$	$\mathcal{O}_{\ell q}^{(1)}$	$(\bar{\ell} \gamma^\mu \ell) (\bar{q} \gamma_\mu q)$
$\mathcal{O}_{qq}^{(3)}$	$(\bar{q} \gamma^\mu \sigma^I q) (\bar{q} \gamma_\mu \sigma^I q)$	$\mathcal{O}_{ee}$	$(\bar{e} \gamma^\mu e) (\bar{e} \gamma_\mu e)$	$\mathcal{O}_{\ell q}^{(3)}$	$(\bar{\ell} \gamma^\mu \sigma^I \ell) (\bar{q} \gamma_\mu \sigma^I q)$
$\mathcal{O}_{uu}$	$(\bar{u} \gamma^\mu u) (\bar{u} \gamma_\mu u)$	$\mathcal{O}_{\ell e}$	$(\bar{\ell} \gamma^\mu \ell) (\bar{e} \gamma_\mu e)$	$\mathcal{O}_{eu}$	$(\bar{e} \gamma^\mu e) (\bar{u} \gamma_\mu u)$
$\mathcal{O}_{dd}$	$(\bar{d} \gamma^\mu d) (\bar{d} \gamma_\mu d)$			$\mathcal{O}_{ed}$	$(\bar{e} \gamma^\mu e) (\bar{d} \gamma_\mu d)$
$\mathcal{O}_{ud}^{(1)}$	$(\bar{u} \gamma^\mu u) (\bar{d} \gamma_\mu d)$			$\mathcal{O}_{qe}$	$(\bar{q} \gamma^\mu q) (\bar{e} \gamma_\mu e)$
$\mathcal{O}_{ud}^{(8)}$	$(\bar{u} \gamma^\mu T^A u) (\bar{d} \gamma_\mu T^A d)$			$\mathcal{O}_{\ell u}$	$(\bar{\ell} \gamma^\mu \ell) (\bar{u} \gamma_\mu u)$
$\mathcal{O}_{qu}^{(1)}$	$(\bar{q} \gamma^\mu q) (\bar{u} \gamma_\mu u)$			$\mathcal{O}_{\ell d}$	$(\bar{\ell} \gamma^\mu \ell) (\bar{d} \gamma_\mu d)$
$\mathcal{O}_{qu}^{(8)}$	$(\bar{q} \gamma^\mu T^A q) (\bar{u} \gamma_\mu T^A u)$			$\mathcal{O}_{\ell edq}$	$(\bar{\ell} e) (\bar{d} q)$
$\mathcal{O}_{qd}^{(1)}$	$(\bar{q} \gamma^\mu q) (\bar{d} \gamma_\mu d)$			$\mathcal{O}_{\ell e qu}^{(1)}$	$(\bar{\ell} e) \epsilon_{rs} (\bar{q}^s u)$
$\mathcal{O}_{qd}^{(8)}$	$(\bar{q} \gamma^\mu T^A q) (\bar{d} \gamma_\mu T^A d)$			$\mathcal{O}_{\ell e qu}^{(3)}$	$(\bar{\ell}^r \sigma^{\mu\nu} e) \epsilon_{rs} (\bar{q}^s \sigma_{\mu\nu} u)$
$\mathcal{O}_{quqd}^{(1)}$	$(\bar{q}^r u) \epsilon_{rs} (\bar{q}^s d)$				
$\mathcal{O}_{quqd}^{(8)}$	$(\bar{q}^r T^A u) \epsilon_{rs} (\bar{q}^s T^A d)$				

Physical basis ops: 59

Green's basis ops: 129

Matching a theory: *off-shell* vs *on-shell*

OFF-SHELL

VS

ON-SHELL

Small number of diagrams
(one-light-particle-irreducible)

1 — 0

All diagrams (light bridges too)

Large set of operators
(Green's basis)

1 — 1

Smaller set of operators
(physical basis)

Matching a theory: *off-shell* vs *on-shell*

OFF-SHELL

VS

ON-SHELL

Small number of diagrams
(one-light-particle-irreducible)

1 — 0

All diagrams (light bridges too)

Large set of operators
(Green's basis)

1 — 1

Smaller set of operators
(physical basis)

Reduction to the physical basis

No reduction is needed

Matching a theory: *off-shell* vs *on-shell*

Elimination of redundant operators in Green's basis via:

EOMs:

$$\begin{aligned}(D_\mu D^\mu H)^a &= m^2 H^a - \lambda(H^\dagger H)H^a - \bar{e}y_E^\dagger \ell^a + i(\sigma_2)^{ab}\bar{q}^b y_U u - \bar{d}y_D^\dagger q^a, & i\not{D}\ell &= y_E e H, \\(D^\nu G_{\mu\nu})^A &= g_s(\bar{q}_i \gamma_\mu T^A q_i + \bar{u}_i \gamma_\mu T^A u_i + \bar{d}_i \gamma_\mu T^A d_i), & i\not{D}e &= y_E^\dagger H^\dagger \ell, \\(D^\nu W_{\mu\nu})^I &= \frac{g}{2}(H^\dagger i \overleftrightarrow{D}_\mu H + \bar{\ell}_\alpha \gamma_\mu \sigma^I \ell_\alpha + \bar{q}_i \gamma_\mu \sigma^I q_i), & i\not{D}q &= y_U u \tilde{H} + y_D d H, \\(\partial^\nu B_{\mu\nu}) &= g'(Y_H H^\dagger i \overleftrightarrow{D}_\mu H + \sum_f Y_f \bar{f} \gamma_\mu f), & i\not{D}u &= y_U^\dagger \tilde{H}^\dagger q, \\ & & i\not{D}d &= y_D^\dagger H^\dagger q.\end{aligned}$$

Straightfoward *but ...*
only valid up to linear
order!
[1811.09413]

Modified EOMs: [2210.14761]

Field redefinitions



Non-trivial process

Hard to program it in a
systematic way

Matching a theory: *off-shell* vs *on-shell*

<i>OFF-SHELL</i>	<i>VS</i>	<i>ON-SHELL</i>
Small number of diagrams (one-light-particle-irreducible)	<i>1</i> — <i>0</i>	All diagrams (light bridges too)
Large set of operators (Green's basis)	<i>1</i> — <i>1</i>	Smaller set of operators (physical basis)
Reduction to the physical basis	<i>1</i> — <i>2</i>	No reduction is needed