GRAVITATIONAL WAVE – BACKGROUNDS –

4th Lecture DANIEL G. FIGUEROA IFIC, Valencia

ICTP's Summer School on Cosmology, 17-28 June 2024, Trieste, Italy

EARLY UNIVERSE



Universe expands, temperature decreases: phase transition triggered !

1

* Potential barrier separates true and false vacua



$$V(\Phi, T) = \begin{cases} Bare potential + \\ thermal corrections \end{cases}$$

Universe expands, temperature decreases: phase transition triggered !

* Potential barrier separates **true** and **false** vacua



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Thermal Fluctuations Or Quantum Tunnelling

Universe expands, temperature decreases: phase transition triggered !

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"bubble" formation





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Real Simulation (Helsinki Group: Hindmarsh, Rummukainen, Weir, ...)







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source: \prod_{ij} tensor anisotropic stress



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• collisions of bubble walls

source: \prod_{ij} tensor anisotropic stress

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source: Π_{ij} tensor anisotropic stress • collisions of bubble walls

• sound waves and turbulence in the fluid

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- collisions of bubble walls
- sound waves and turbulence in the fluid
- primordial magnetic fields (MHD)

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- **sound waves** and turbulence in the fluid
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Technically

source: Π_{ij} tensor anisotropic stress $\begin{aligned} \Pi_{ij} \sim \partial_i \phi \, \partial_j \phi \\ \Pi_{ij} \sim \gamma^2 (\rho + p) \, v_i v_j + \text{Scalar-Fluid coupling} \\ \Pi_{ij} \sim \frac{(E^2 + B^2)}{3} - E^i E^j - B^i B^j + \text{MHD} \end{aligned}$

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Sounds waves (relativistic fluids) + Magneto-hydrodynamic (MHD) effects + Turbulence



M. Hindmarsh *et al*, Phys.Rev.D 92 (2015) 12, 123009; **Arxiv:** 1504.03291 (Series of papers 2012-2023)

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* **GW causal source**: cannot 'operate' beyond the **horizon**

$$f_* = \frac{H(T_*)}{\epsilon_*}$$



parameter characterising source

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$$f_c = f_* \frac{a_*}{a_0} = \frac{2 \cdot 10^{-5}}{\epsilon_*} \frac{T_*}{1 \text{ TeV}} \text{ Hz}$$

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GW generation <--> bubbles properties

$$\beta^{-1}$$
: duration of PhT
 $v_b \leq 1$: speed of bubble walls $\rightarrow R_* = v_b \beta^{-1}$ size of bubbles at collision

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Freq.
(today)

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 Parameter List
(not independent)
 $\epsilon \simeq \frac{H_*}{\beta}, H_* R_*$ \longrightarrow $\frac{\beta}{H_*}, v_b, T_*$

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Amplitude
(today)
 $\Omega_{\text{GW}} \sim \Omega_{\text{rad}} \epsilon_*^2 \left(\frac{\rho_s^*}{\rho_{\text{tot}}^*}\right)^2$ \longrightarrow $\alpha = \frac{\rho_{\text{vac}}}{\rho_{\text{rad}}^*}$
 $\frac{\rho_s^*}{\rho_{\text{tot}}^*} = \frac{\kappa \alpha}{1 + \alpha}$ $\kappa = \frac{\rho_{\text{kin}}}{\rho_{\text{vac}}}$

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 $\sim 10^{-4} \times (10^{-2})^2 \times 10^{-X}$ $\alpha = \frac{\rho_{\text{kin}}}{\rho_{\text{vac}}}$

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 $\sim 10^{-8} \times 10^{-X}$ $\kappa = \frac{\rho_{\text{kin}}}{\rho_{\text{vac}}}$





(From Collisions + Sound Waves + Turbulence)



peak of bubble collisions







Evaluation of the signal

• **bubble collisions**: **analytical** and **numerical** simulations

Huber, Konstandin '08 Cutting, Hindmarsh et al 2018, ...

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 sound waves: numerical simulations of scalar field and fluid Hindmarsh, Weir et al 2012 - 2019, analytical Hindmarsh 2016, 2019,
Evaluation of the signal

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• MDH turbulence: analytical evaluation Kosowsky et al '07, Caprini et al '09, Niksa et al '18 numerical Pol et al 2019

 LISA sensitive to energy scale 10 GeV - 100 TeV ! (mHZ)

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- LISA can probe the EWPT in BSM models ...
 - singlet extensions of (MS)SM (e.g. Huber et al 2015)
 - direct coupling of Higgs to scalars (e.g. Kozackuz et al 2013)
 - SM + dimension six operator (e.g. Grojean et al 2004)

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 - Warped extra dimensions : PT from the dilaton/radion stabilisation in RS-like models (Randall and Servant 2015)

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 LISA can problem: LHC is putting great pressure over BSM scenarios
 hterplay! *ar 2015 ckuz et al 2013*) *t al 2004*)
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Can we really detect a 1st-O Ph-T?

* LISA/ET can, but LHC pressures typical BSM extensions to promote EW-PhT into First Order

* Assuming LHC does not rule out models, LISA/ET can detect/constrain significant fraction of Param Space

Can we really detect a 1st-O Ph-T ?

* LISA/ET can, but LHC pressures typical BSM extensions to promote EW-PhT into First Order

- * Assuming LHC does not rule out models, LISA/ET can detect/constrain significant fraction of Param Space
- * Predictions depend on many assumptions (particularly in sound waves), so is our modelling correct?
 - * Even if we detect it, then we infer α and β , but what BSM model is behind? not univocal !



Gravitational Wave Backgrounds

OUTLINE

Early Universe Sources 2) GWs from Inflation ✓

1) Grav. Waves (GWs)

- 3) GWs from Preheating 🗸
- 4) GWs from Phase Transitions 🗸

1st Topic

Core

Topics

- 5) GWs from Cosmic Defects
- 6) Astrophysical Background(s)
- 7) Observational Constraints/Prospects

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Cosmic Defects

Aftermath product of a Ph.T.

Topology of cosmic domains and strings

T W B Kibble

Blackett Laboratory, Imperial College, Prince Consort Road, London SW7 2BZ, UK

Received 11 March 1976

Abstract. The possible domain structures which can arise in the universe in a spontaneously broken gauge theory are studied. It is shown that the formation of domain walls, strings or monopoles depends on the homotopy groups of the manifold of degenerate vacua. The subsequent evolution of these structures is investigated. It is argued that while theories generating domain walls can probably be eliminated (because of their unacceptable gravitational effects), a cosmic network of strings may well have been formed and may have had important cosmological effects.

Kibble pioneered the study of topological defect generation in the early universe.



 $\mathbf{M} = \mathbf{G}/\mathbf{H}$



M = G/H

Spatial Dist.





M = G/H

Spatial Dist.























Example: Cosmic Strings

(e.g. From PhT after Hybrid Inflation)



Dufaux et al PRD 2010



Dufaux et al PRD 2010

DEFECTS: Aftermath of PhT \rightarrow $\begin{cases}
Domain Walls \\
Cosmic Strings \\
Cosmic Monopoles \\
Non - Topological
\end{cases}$

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CAUSALITY & MICROPHYSICS \Rightarrow Corr. Length: $\xi(t) = \lambda(t) H^{-1}(t)$


Introduction to Cosmic Defects

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Domain Walls Cosmic Strings Cosmic Monopoles Non - Topological
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CAUSALITY & MICROPHYSICS \Rightarrow Corr. Length: $\xi(t) = \lambda(t) H^{-1}(t)$

(Kibble' 76)

SCALING:
$$\lambda(t) = \text{const.} \rightarrow \lambda \sim 1 \implies k/\mathcal{H} = kt$$

comoving momentum conformal time

.

Introduction to Cosmic Defects

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Introduction to Cosmic Defects









DEFECTS: GW Source $\rightarrow \{T_{ij}\}^{TT} \propto \{\partial_i \phi \partial_j \phi, E_i E_j, B_i B_j\}^{TT}$

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UTC: $\langle T_{ij}^{TT}(\mathbf{k},t)T_{ij}^{TT}(\mathbf{k}',t')\rangle = (2\pi)^3 \Pi^2(\mathbf{k},t_1,t_2) \ \delta^3(\mathbf{k}-\mathbf{k}')$ (Unequal Time Correlator)

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GW spectrum:ExpansionUTC $\frac{d\rho_{\rm GW}}{d\log k}(k,t) \propto \frac{k^3}{M_p^2 a^4(t)} \int dt_1 dt_2 \ a(t_1)a(t_2) \ \cos(k(t_1-t_2)) \ \Pi^2(k,t_1,t_2)$

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GW spectrum:
$$(x_i \equiv kt_i)$$
ExpansionUTC $\frac{d\rho_{\rm GW}}{d\log k}(k,t) \propto \left(\frac{V}{M_p}\right)^4 \frac{M_p^2}{a^4(t)} \left[\int dx_1 dx_2 \sqrt{x_1 x_2} \cos(x_1 - x_2) \ U(x_1, x_2)\right]$ Rad. DomSCALING

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 $F_U \sim \text{Const.}$ (Dimensionless)



DGF, Hindmarsh, Urrestilla, PRL 2013



 \forall PhT (1st, 2nd, ...), \forall Defects (top. or non-top.)

DGF, Hindmarsh, Urrestilla, PRL 2013



$$\mathsf{RD} \qquad F_U^{(\mathrm{R})} \equiv \frac{32}{3} \int_0^x dx_1 dx_2 \, (x_1 x_2)^{1/2} \cos(x_1 - x_2) \, U_{\mathrm{RD}}(x_1, x_2)$$
$$\mathsf{MD} \qquad F_U^{(\mathrm{M})} \equiv \frac{32}{3} \frac{(\sqrt{2} - 1)^2}{2} \int_{x_{\mathrm{eq}}}^x dx_1 dx_2 \, (x_1 x_2)^{3/2} \cos(x_1 - x_2) \, U_{\mathrm{MD}}(x_1, x_2)$$

DGF, Hindmarsh, Lizarraga, Urrestilla, PRD 2020



However this assumes exact scaling !

DGF, Hindmarsh, Lizarraga, Urrestilla, PRD 2020





Gorghetto, Hardy, Villadoro, JHEP 2018



Global strings (e.g. axion DM) → Log corrections argued & found

Gorghetto, Hardy, Villadoro, JHEP 2018



Not scale invariant due to Log enhancements !



e.g. Gorghetto et al



Not scale invariant due to Log enhancements !





* Scaling dynamics (exact)

* Infinitely thin

* Inter-commutation

* Scaling dynamics

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* Inter-commutation

Infinitely thin: $H^{-1} \gg m^{-1}$



* Scaling dynamics

* Infinitely thin

* Inter-commutation

Intercommutation





Loops !

Loops !



Loops are formed !



Loops are formed ! Vibrate under their tension !



Periodic Oscillations

Loops are formed !

Vibrate under their tension !



Gravitational Waves (GW) are emitted !

Loops are formed ! V

Vibrate under their tension !



Gravitational Waves (GW) are emitted !

Superposition from many loop signals



Gravitational Wave Background

Traditional picture --- Nambu-Goto approximation (zero width)

String networks = Infinite strings + Loops

Traditional picture --- Nambu-Goto approximation (zero width)

String networks = Infinite strings + Loops
Decay to loops
Traditional picture --- Nambu-Goto approximation (zero width)

► Decay to GWs (Vilenkin '81)

String networks = Infinite strings + Loops
Decay to loops

Traditional picture --- Nambu-Goto approximation (zero width)

Decay to GWs
 String networks = Infinite strings + Loops
 Decay to loops

> Loops decay via GWs radiated in all harmonic frequencies ν_j

$$P_j = \Gamma G \mu^2 \frac{j^{-q}}{\zeta(q)} \longrightarrow P_{GW} = \dot{E}_{GW} = \sum_{j=1}^{\infty} P_j = \Gamma G \mu^2$$

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But ...

Field-theory strings can also decay via particle emission

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Following ... we assume GW emission dominates !

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But ...

Field-theory strings can also decay via particle emission

Following ... we assume GW emission dominates^{*}! (*Lattice calculations show opposite... not published yet)

Cosmic string loop (length *l*) oscillates under tension μ
 emits GWs in a series of harmonic modes

Cosmic string loop (length *l***) <u>oscillates</u> under tension µ emits GWs in a series of harmonic modes**

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> Original emission of GWs ! (Vilenkin '81) and many others !

"extra" emission on top of Irreducible background (only for strings)

Cosmic string loop (length *l***) <u>oscillates</u> under tension µ emits GWs in a series of harmonic modes**

$$\frac{d\rho^{(o)}}{df} \equiv \Gamma G \mu^2 \int_{t_*}^{t_o} dt \left(\frac{a(t)}{a_o}\right)^3 \int_0^{\alpha/H(t)} dl \ln(l,t) \mathcal{P}((a_o/a(t))fl)$$

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expansion
history

Cosmic string loop (length *l*) oscillates under tension μ emits GWs in a series of harmonic modes

$$\begin{split} \frac{d\rho^{(\mathrm{o})}}{df} &\equiv \Gamma G \mu^2 \int_{t_*}^{t_o} dt \left(\frac{a(t)}{a_o}\right)^3 \int_{0}^{\alpha/H(t)} \frac{dlln(l,t) \,\mathcal{P}((a_o/a(t))fl)}{\sqrt{\sqrt{\frac{1}{1}}}} \\ & \text{expansion} \\ & \text{history} \\ & \text{length} \\ & \text{number} \\ & \text{density} \\ & \left(\frac{\text{Nambu-Goto}}{\text{simulations}}\right) \end{split}$$

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expansion
history
length
lengt

Cosmic strings loops: GW background



e.g. Sanidas et al 2012

Cosmic strings loops: GW background



Sanidas et al 2012



Model II (BOS) from LISA paper



LISA paper: 1909.00819 [astro-ph.CO]

Model II (BOS) from LISA paper



Model II (BOS) from LISA paper



@ LISA: Very large parameter space !

GW background constrained by LISA

 $G\mu \gtrsim 10^{-17} ~(v \gtrsim 10^{10} {
m GeV})$

CMBPTA (today)PTA (future) $G\mu \sim 10^{-7}$ $G\mu \sim 10^{-11}$ $G\mu \sim 10^{-14}$

GW background constrained by LISA





GW background constrained by LISA



LISA * Best constraints on Comic Strings * (actually only way to obtain them) * Discovery, or stringent constraints

EARLY UNIVERSE in GWs 🗸



EARLY UNIVERSE in GWs \checkmark



Gravitational Wave Backgrounds

OUTLINE

Early Universe Sources 2) GWs from Inflation 🗸

1) Grav. Waves (GWs)

- 3) GWs from Preheating 🗸
- 4) GWs from Phase Transitions 🗸

| 1st Topic

Core

Topics

- 5) GWs from Cosmic Defects 💊
- 6) Astrophysical Background(s)
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Gravitational Wave Backgrounds

OUTLINE

1) Grav. Waves (GWs) 1st Topic

Early Universe Sources

> Late Universe + Observations

2) GWs from Inflation ✓
3) GWs from Preheating √

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Late Universe + Observations

Late Universe $(0 \le z \le 10)$



LIGO/VIRGO 2015-now



Late Universe $(0 \le z \le 10)$













$(0 \le z \le 10)$











$(0 \le z \le 10)$







 $(0 \le z \le 10)$

$$\Omega_{\rm GW}(f) = \frac{1}{\rho_c} \, \frac{d\rho_{\rm GW}}{d \log f} \,,$$



Late Universe $(0 \le z \le 10)$

$$\Omega_{\rm GW}(f) = \frac{1}{\rho_c} \frac{d\rho_{\rm GW}}{d{\log}f} = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f) ,$$

$$\uparrow$$
Characteristic strain



Late Universe $(0 \le z \le 10)$

$$\Omega_{\rm GW}(f) = \frac{1}{\rho_c} \frac{d\rho_{\rm GW}}{d\log f} = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f) \,,$$

For binary population:
$$\frac{dn}{dz} \begin{pmatrix} \text{comoving} \\ \text{number} \\ \text{density} \end{pmatrix}$$

/


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$$h_c^2(f) = \frac{4}{\pi} \frac{G}{c^2} f^{-2} \int_0^\infty dz \frac{dn}{dz} \frac{1}{1+z} f_r \frac{dE_{\rm GW}}{df_r} \bigg|_{f_r = f(1+z)}$$



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$$h_c^2(f) = \frac{4}{\pi} \frac{G}{c^2} f^{-2} \int_0^\infty dz \frac{dn}{dz} \frac{1}{1+z} f_r \frac{dE_{\rm GW}}{df_r} \bigg|_{f_r = f(1+z)}$$

E.S. Phinney
astro-ph/0108028
Source-frame

energy spectrum



a



$$\begin{split} \Omega_{\rm GW}(f) &= \frac{1}{\rho_c} \left. \frac{d\rho_{\rm GW}}{d\log f} = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f) \,, \\ h_c^2(f) &= \frac{4}{\pi} \frac{G}{c^2} f^{-2} \int_0^\infty dz \frac{dn}{dz} \frac{1}{1+z} f_r \frac{dE_{\rm GW}}{df_r} \bigg|_{f_r = f(1+z)} \\ &\left. \frac{dE_{\rm GW}}{df_r} = \frac{\pi}{3} \frac{1}{G} \frac{(G\mathcal{M})^{5/3}}{\pi^{1/3} f_r^{1/3}} \right|_{f_r = f(1+z)} \end{split}$$

Chirp mass $\mathcal{M} = (m_1 + m_2)^{2/5} \left(\frac{m_1 m_2}{m_1 + m_2}\right)^{3/5}$





$$\Omega_{\rm GW}(f) = \frac{1}{\rho_c} \frac{d\rho_{\rm GW}}{d\log f} = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f) ,$$

$$f) = \frac{4}{\pi} \frac{G}{c^2} f^{-2} \int_0^\infty dz \frac{dn}{dz} \frac{1}{1+z} f_r \frac{dE_{\rm GW}}{df_r} \bigg|_{f_r = f(1+z)}$$
Ising:
$$\frac{dn}{dz} \equiv \frac{d\mathcal{N}}{dz dm_1 dm_1 dV} \equiv R(z) \ p(m_1, m_2) \frac{dt_r}{dz}$$
Mass
Merging function
Rate (distribution)

$$\Omega_{\rm GW}(f) = \frac{1}{\rho_c} \frac{d\rho_{\rm GW}}{d\log f} = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f) ,$$

$$h_c^2(f) = \frac{4}{\pi} \frac{G}{c^2} f^{-2} \int_0^\infty dz \frac{dn}{dz} \frac{1}{1+z} f_r \frac{dE_{\rm GW}}{df_r} \bigg|_{f_r = f(1+z)}$$

$$h_c^2(f) = \frac{4G^{5/3} f^{-2/3}}{3\pi^{1/2} c^2} \int dm_1 dm_2 p(m_1, m_2) \int_0^{z_{\rm max}} dz' R(z') \frac{\mathcal{M}^{5/3}(m_1, m_2)}{(1+z)^{1/3}} \frac{dt_r}{dz}$$





Mass Function Spin distribution

(Black Holes, Neutron Stars & White Dwarfs)

Late Universe

$(0 \le z \le 10)$





$(0 \le z \le 10)$



Calibrated Mass Function Spin distribution (Black Holes, Neutron Stars & White Dwarfs)







Calibrated Mass Function Spin distribution (Black Holes, Neutron Stars & White Dwarfs)

Computation of Solar-Mass Binary BH Grav. Wave Background

(LISA Collaboration

with

S. Babak, C. Caprini, N. Karnesis, P. Marcoccia, M. Pieroni A. Ricciardone, A. Sesana, J. Torrado

Late Universe

$(0 \le z \le 10)$



Calibrated Mass Function Spin distribution (Black Holes, Neutron Stars & White Dwarfs)



Late Universe

$(0 \le z \le 10)$



Calibrated Mass Function Spin distribution (Black Holes, Neutron Stars & White Dwarfs)



Babak et al 2023

Lehoucq et al 2023

Gravitational Wave Backgrounds

OUTLINE

1) Grav. Waves (GWs) 1st Topic

Early Universe Sources

> Late Universe + Observations

2) GWs from Inflation
3) GWs from Preheating
4) GWs from Phase Transitions
5) GWs from Cosmic Defects
6) Astrophysical Background(s)

7) Observational Constraints/Prospects

Gravitational Wave Backgrounds

OUTLINE

1) Grav. Waves (GWs) 1st Topic

Early Universe Sources 2) GWs from Inflation
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5) GWs from Cosmic Defects
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7) Observational Constraints/Prospects

Late Universe + Observations

GWB Observations

O1. Data AnalysisO2. ConstraintsO3. Evidence (PTA)

01. Data Analysis

O1. Data Analysis O2. Constraints

O3. Evidence (PTA)

01. Data Analysis

O1. Data Analysis O2. Constraints O3. Evidence (PTA)







(Caprini et al <u>1906.09244</u>)

O1. Data Analysis 01. Data Analysis O2. Constraints **O3. Evidence (PTA) POWER SPECTRUM RECONSTRUCTION Machine Learning Analysis** 10^{-8} **GW Spectrum** --- central noise **GW Spectrum** theoretical signal 10^{-9} reconstructed signal 68% CI 95% CI 10-10 $h^2\Omega$ 10-11 10^{-12} 10^{-13} 10-3 10-2 Frequency [Hz] **Code** GWBackFinder (Dimitriou et al 2309.08430) **Code** SGWBinner Code saqqara (Alvey et al 2309.07954)

O1. Data Analysis 01. Data Analysis O2. Constraints **O3. Evidence (PTA) POWER SPECTRUM RECONSTRUCTION Machine Learning Analysis** 10^{-8} **GW Spectrum** --- central noise **GW Spectrum** --- theoretical signal 10^{-9} reconstructed signal 68% CI 95% CI 10-10 $h^2\Omega$ 10-11 10^{-12} 10^{-13} 10-3 10-2 Frequency [Hz] **Code** GWBackFinder (Dimitriou et al <u>2309.08430</u> **Code** SGWBinner Can be adapted to any experimental Code saggar Noise (parametrised)



O1. Data AnalysisO2. ConstraintsO3. Evidence (PTA)

O1. Data Analysis O2. Constraints O3. Evidence (PTA)

(Before June 2023)



O1. Data Analysis O2. Constraints O3. Evidence (PTA)

Future ~ 15-20 yrs



O1. Data Analysis O2. Constraints O3. Evidence (PTA)

Future ~ 30-40 yrs



CMB Latest Analysis



(from Copeland et al 2004.11396)

BACK CURRENT

PLANNED

PLANNED II

ALL

CMB

PTA-detection

O3. Evidence (PTA) O1. Data Analysis O2. Constraints O3. Evidence (PTA)



NanoGrav, EPTA+IPTA, PPTA, CPTA (Pulsar Timing Array Collaborations)



(June 2023)

NanoGrav, EPTA+IPTA, PPTA, CPTA (Pulsar Timing Array Collaborations)



[Hellings, Downs: Astrophys. J. 265 (1983) L39]

Hellings - Downs



"Strong" Evidence

2306.16213: NANOGrav



⁶⁸ pulsars, 16 yr of data, HD at $\sim 3 \cdots 4 \sigma$

2306.16215: PPTA



32 pulsars, 18 yr of data, HD at $\sim 2 \sigma$



²⁵ pulsars, 25 yr of data, HD at \sim 3 σ

2306.16216: CPTA



57 pulsars, 3.5 yr of data, HD at \sim 4.6 σ

"Strong" Evidence



Interpretation

Super Massive BHB (SMBHB, expected)



NanoG: 2306.16220 [astro-ph.HE] EPTA: 2306.16227 [astro-ph.CO]

Interpretation

or

Super Massive BHB (SMBHB, expected)



NanoG: 2306.16220 [astro-ph.HE] EPTA: 2306.16227 [astro-ph.CO]

Cosmological GWB (more speculative)



EPTA: 2306.16227 [astro-ph.CO] NanoG: 2306.16219 [astro-ph.HE]

Interpretation

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O3. Evidence (PTA) O3. Detection (PTA) O3. Detection (PTA)

Interpretation

or

Super Massive BHB (SMBHB, expected)



Key: Anisotropies ?

Cosmological GWB (more speculative)


O3. Evidence (PTA) O3. Data Analysis O3. Detection (PTA)

Interpretation

or

Super Massive BHB (SMBHB, expected)



Key: Anisotropies ?

Cosmological GWB (more speculative)



Key: Spectrum

Inflation, Domain Walls, Cosmic Strings 1stO-PhT, audible Axions, 2ndOI-GWB, ...

Summary & Outlook

Cosmological

Early Universe





Universe

Universe



Universe

Universe



Early Universe





Early Universe



Early Universe







C5. Astrophysical Background

UNIQUE OPPORTUNITY Probe high energy physics (otherwise inaccessible) Complementary/Beyond Particle Colliders ! Late Universe [C5]



C5. Astrophysical Background

Late Universe

[**C**5]







Nature/observability of Inflation Eq. of State of the early universe Primordial perturbations at all scales Viability of Axion species Did phase transitions take place? Are there primordial black holes ? Are cosmic strings in the Universe ? Foreground role Population properties Anisotropies Cross-correlations Gravity properties



















Early Universe GWB Program



Early Universe GWB Program

But before we conclude ...

Final Remarks

Expertise



Expertise



Expertise



Review on Gravitational Waves from the Early Universe

Caprini & Figueroa arXiv:1801.04268

For you early universe numerics ...





For you early universe numerics ...



For you early universe numerics ...



For you early universe numerics ...



2nd CL School 2023: Sept 25-29

https://www.youtube.com /@CosmoLattice/videos



CosmoLattice School 2023, Day 4: Practice 3 (Simulating Gravitational Waves)

17 views · 4 months ago



CosmoLattice School 2023, Day 4: Lecture 8 (Plotting Features of CosmoLattice)



CosmoLattice School 2023, Day 3: Lecture 7 [SU(2) Scalar-Gauge Theory Lattice...

Evolution of GWs modes
non-local operations are computationally expensive!
Solution: we define a set of unphysical tensor modes u's
(1) Evolve equation of motion of u's)

$$u_{ij} + 3H\dot{u}_{ij} - \frac{\nabla^2}{a^2}u_{ij} = \frac{2}{m_{ij}^2a^2}\Pi_{ij}^{-1}$$

(2) When needed, (compute power spectrum energy
density) we apply transformation
 $h_{ij}(k, t) = \Lambda_{ij,kl}(k)u_{kl}(k, t)$
1:19:39

CosmoLattice School 2023, Day 3: Lecture 6 (Creation and Propagation of Grav. Waves)

12 views • 4 months ago

For you early universe numerics ...



Details for 3rd CL School TBA at:

https://cosmolattice.net



Thanks for your attention

Vielen Dank für Ihre Aufmerksamkeit







CMB defects Kibble Mechanism

Lattice Strings

Modeling

What about lattice simulations ?



(Image: David Daverio)
Modeling

What about lattice simulations ?

Abelian-Higgs Simulations

- * Loops formed ! ... but decay into scalar/gauge fields
- * If loops disappear... then no GW?
- * There is an **irreducible GW emission** from the long string network, but negligible vs NG loop GW emission

Modeling

What about lattice simulations ?

Abelian-Higgs Simulations

- * Loops formed ! ... but decay into scalar/gauge fields
- * If loops disappear... then no GW?
- * There is an **irreducible GW emission** from the long string network, but negligible vs NG loop GW emission

Open debate in the ArXiv !

Vachaspati et al 2019/20 Copeland et al 2023/24

Modeling

What about lattice simulations ?

Abelian-Higgs Simulations

- * Loops formed ! ... but decay into scalar/gauge fields
- * If loops disappear... then no GW?
- * There is an **irreducible GW emission** from the long string network, but negligible vs NG loop GW emission

So ... LISA results based on Nambu-Goto strings !

Kibble Mechanism

Introduction to Cosmic Defects

Kibble as recall the more general situation. In a model with symmetry group G, the vacuum expectation value $\langle \phi \rangle$ will be restricted to lie on some orbit of G. If H is the isotropy subgroup of G at one point $\langle \phi \rangle$, i.e. the subgroup of transformations leaving $\langle \phi \rangle$ unaltered, then the orbit may be identified with the coset space M = G/H. Physically H is the subgroup of unbroken symmetries, and M is the manifold of degenerate vacua. As we shall see, the topological properties of M (specifically its homotopy groups) largely determine the geometry of possible domain structures.





Introduction to Cosmic Defects

6. Conclusions and discussion

On this basis we showed that a domain structure can be expected to arise. The topological character of this structure depends on the homotopy groups $\pi_k(M)$ of the manifold M of degenerate vacua. Domain walls can form if $\pi_0(M)$ is nontrivial, i.e. if M is non-connected. If it has n connected components we find an n-phase emulsion. The formation of cosmic strings requires that $\pi_1(M)$ be nontrivial, i.e. that M is not formed of simply connected components. Finally, 'monopoles' can form if $\pi_2(M)$ is nontrivial.



M = G/H



Introduction to Cosmic Defects





 $\mathbf{M} = \mathbf{G}/\mathbf{H}$



CMB Defects (Back) SLIDES





Durrer et al, JCAP 2014



Durrer et al, JCAP 2014

