

## 1. Horizons

The *particle horizon* is the largest comoving distance from which light emitted in the past could reach the observer at a given time. One could also define a *cosmological event horizon* in an expanding Universe, which corresponds to the largest comoving distance from which light emitted now can ever reach the observer in the future.

- a) Assuming the scale factor obeys  $a(t) = (t/t_0)^n$ , which horizon exists?
- b) And what if the scale factor is given by  $a(t) = \exp(Ht)$ ?
- c) Are there FLRW metrics that exhibit neither horizon? Or both horizons?

## 2. Einstein's static universe

The cosmological constant was originally introduced by Einstein to make the universe static, i.e., not expanding.

- a) Consider a universe filled with a perfect fluid with  $\rho > 0$  and  $P \geq 0$ . Show that no static solution to the Friedmann equations exists in this case.
- b) Now consider a universe filled with pressureless matter ( $P_m = 0$ ) and a cosmological constant ( $P_\Lambda = -\rho_\Lambda$ ). Show that in this case it is possible to find a static solution and find the ratio  $\xi_{\text{static}} \equiv \rho_m / \rho_\Lambda$ .
- c) Now imagine there were small density fluctuations in the matter component of this universe. Describe qualitatively what would happen for
  - (i) a small matter overdensity,  $\xi > \xi_{\text{static}}$ ,
  - (ii) and a small matter underdensity,  $\xi < \xi_{\text{static}}$ .

## 3. Conformal time

Consider a flat FLRW universe. Define the conformal time  $\eta$

$$d\eta \equiv \frac{dt}{a}.$$

What is the functional dependence of the scale factor  $a$  on  $\eta$  during

- a) matter domination?
- b) radiation domination?
- c)  $\Lambda$  domination?

## 4. Neutrinos

Standard model neutrinos decouple from cosmic plasma at temperatures of  $\sim 1$  MeV. Since laboratory experiments have shown that neutrino masses are sub-eV, it is clear that these neutrinos must be ultra-relativistic when decoupling happens. If the decoupling is instantaneous, the phase space density of neutrinos at the time of decoupling takes the relativistic Fermi-Dirac form to a very good approximation, i.e.,

$$f(p) = \frac{1}{\exp(p/T_\nu) + 1} \quad (1)$$

- a) Argue why the phase space distribution  $f(p)$  must remain the same form, with  $T_\nu \propto a^{-1}$ , *even after* the temperature  $T_\nu$  drops below the neutrino rest mass.

- b) Suppose one species of neutrinos has a mass of  $m$ , where  $m \gg T_{\nu,0} \simeq 10^{-5}$  eV and  $T_{\nu,0} \simeq 1.95$  K is the present-day neutrino temperature. Show that the present-day reduced energy density of these neutrinos is given by

$$\Omega_\nu h^2 = \frac{m}{93 \text{ eV}}. \quad (2)$$

- c) The present-day neutrino temperature is related to the present-day CMB temperature by

$$T_{\nu,0} = \left(\frac{4}{11}\right)^{1/3} T_{\gamma,0}, \quad (3)$$

where  $T_{\gamma,0} = 2.725$ . If the Fermi coupling constant  $G_F$  were a factor of 100 smaller, at what temperature would you expect the neutrinos to decouple? Would the relation given in equation (3) still apply to the present-day neutrino temperature in this case?

## 5. Light thermal relics

Suppose we extend the standard model (SM) of particle physics with one extra massless scalar particle  $\phi$ . This particle interacts with the SM particles and was once in a state of thermodynamic equilibrium in the early universe. At some temperature  $T_{\text{dec}(\phi)}$ ,  $\phi$  decouples from the rest of the cosmic plasma.

- a) Suppose  $\phi$  decouples before neutrino decoupling, i.e.,  $T_{\text{dec}(\phi)} \gg O(1)$  MeV. Assuming as usual instantaneous decoupling, show using entropy conservation arguments that the present-day temperature of  $\phi$ ,  $T_{\phi,0}$ , is related to the present-day neutrino temperature  $T_{\nu,0}$  via

$$T_{\phi,0} = \left[ \frac{43}{4} \frac{1}{g_{*s}(T_{\text{dec}(\phi)}) - 1} \right]^{1/3} T_{\nu,0}, \quad (4)$$

where  $g_{*s}$  is the effective massless entropy degrees of freedom, and three generations of neutrinos have been assumed. How does  $T_{\phi,0}$  relate to the present-day photon temperature  $T_{\gamma,0}$ ?

- b) Now suppose  $\phi$  decoupling happens after  $e^+e^-$  annihilation ( $T \sim 0.5$  MeV). How does  $T_{\phi,0}$  now relate to  $T_{\nu,0}$  and to  $T_{\gamma,0}$ ?
- c) Compared with the standard model, how does the presence of this hypothetical particle  $\phi$  affect the expansion rate and hence the neutrino decoupling temperature?

## 6. Boltzmann equation

The integrated Boltzmann equation for the physical number density of a particle species  $i$  is given by

$$\frac{dn}{dt} + 3\frac{\dot{a}}{a}n = \frac{g}{(2\pi)^3} \int d^3p \frac{1}{E} C[f], \quad (5)$$

where  $C[f]$  is the collision integral. Assuming an interaction

$$i + j \leftrightarrow k + l \quad (6)$$

and  $CP$  invariance, verify that the r.h.s. of equation (5) evaluates to exactly zero when the interaction (6) is in equilibrium.

## 7. Scalar fields

Begin with the Lagrangian density for a real scalar field  $\phi(x^\alpha)$ ,

$$\mathcal{L}_\phi = -\frac{1}{2}g^{\mu\nu}\partial_\mu\phi\partial_\nu\phi - V(\phi). \quad (7)$$

(metric convention:  $(-, +, +, +)$ ).

a) Using

$$T^{\mu\nu} = -\frac{2}{\sqrt{-g}}\frac{\delta(\sqrt{-g}\mathcal{L})}{\delta g_{\mu\nu}} = g^{\mu\nu}\mathcal{L} - 2\frac{\delta\mathcal{L}}{\delta g_{\mu\nu}}, \quad (8)$$

where  $g$  is the determinant of the metric, derive the stress-energy tensor  $T_{\mu\nu}$  for  $\phi$ .

b) Suppose the field is homogeneous, i.e.,  $\phi = \phi(t)$ . Show that the energy density  $\rho_\phi$  and the pressure  $P_\phi$  of  $\phi(t)$  are given by

$$\begin{aligned} \rho_\phi(t) &= \frac{1}{2}\left(\frac{d\phi}{dt}\right)^2 + V(\phi), \\ P_\phi(t) &= \frac{1}{2}\left(\frac{d\phi}{dt}\right)^2 - V(\phi). \end{aligned} \quad (9)$$

c) Consider the scalar field  $\phi$  described by equation (7), and suppose again it is homogeneous, i.e.,  $\phi = \phi(t)$ . Show that in an FLRW universe, the time evolution of  $\phi(t)$  is governed by the Klein-Gordon equation:

$$\ddot{\phi} + 3\frac{\dot{a}}{a}\dot{\phi} + \frac{\partial V}{\partial\phi} = 0. \quad (10)$$

You can start the derivation from first principles, i.e., minimise the Lagrangian. But you can also arrive at the same point via the continuity equation,  $d\rho/dt + 3H(\rho + P) = 0$ . You should try both.