

# THE ANOMALOUS MAGNETIC MOMENT OF THE MUON

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OPEN QUESTIONS AND FUTURE DIRECTIONS IN FLAVOUR PHYSICS  
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# THE MUON $g-2$ : A PROBE FOR NEW PHYSICS

- Magnetic moment of charged leptons  $l \in \{e, \mu, \tau\}$ :

$$\vec{\mu}_l = g_l \cdot \frac{e}{2m_l} \cdot \vec{s}$$

- Quantum corrections lead to deviations from the classical value  $g = 2$  (Dirac), the anomalous magnetic moment,

$$a_l = \frac{g_l - 2}{2} = \frac{\alpha}{2\pi} + O(\alpha^2) \quad (\text{Schwinger})$$



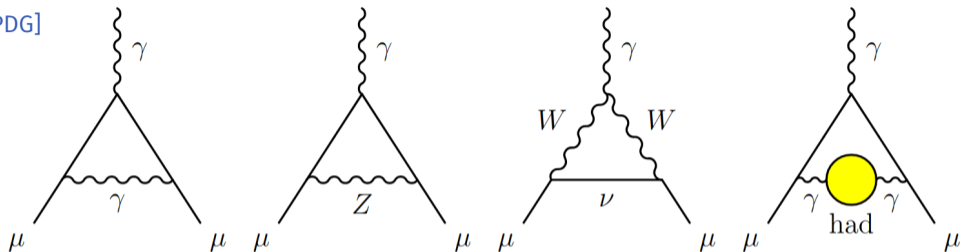
- Contributions from new physics at the scale  $\Lambda_{\text{NP}}$  enter  $a_l$  via

$$a_l - a_l^{\text{SM}} \propto \frac{m_l^2}{\Lambda_{\text{NP}}^2}$$

with  $m_\mu/m_e \approx 207$ .

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Figure: [PDG]



- SM prediction from QED, electroweak and hadronic contributions:

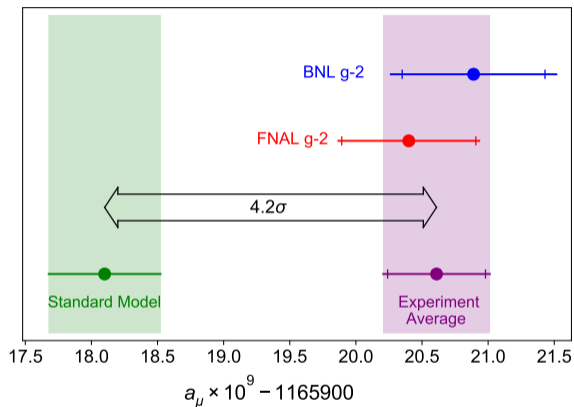
$$a_l^{\text{SM}} = a_l^{\text{QED}} + a_l^{\text{EW}} + a_l^{\text{had}} \quad \text{where} \quad a_l^{\text{had}} = a_l^{\text{hvp}} + a_l^{\text{HLbL}}.$$

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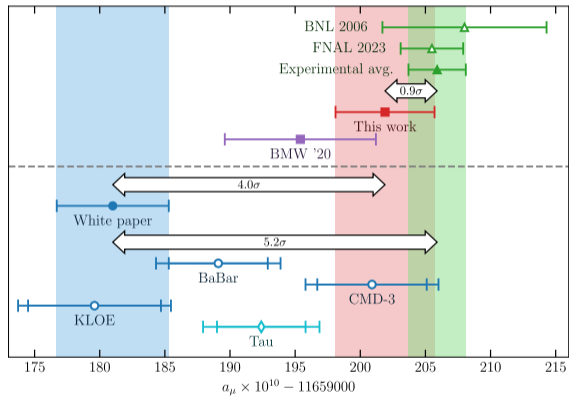
where  $a_\tau$  is inaccessible for experiment and  $m_\mu/m_e \approx 207$ .

# THE MUON $g - 2$ : A PROBE FOR NEW PHYSICS



- Comparison of Standard Model prediction and experimental average [Muon  $g - 2$ , 2104.03281]
- After Run-1 results of the Fermilab  $g - 2$  experiment.
- Standard Model prediction based on the White Paper of the Muon  $g-2$  Theory Initiative [Aoyama et al., 2006.04822]

# THE MUON $g - 2$ – 2: A PROBE FOR NEW PHYSICS



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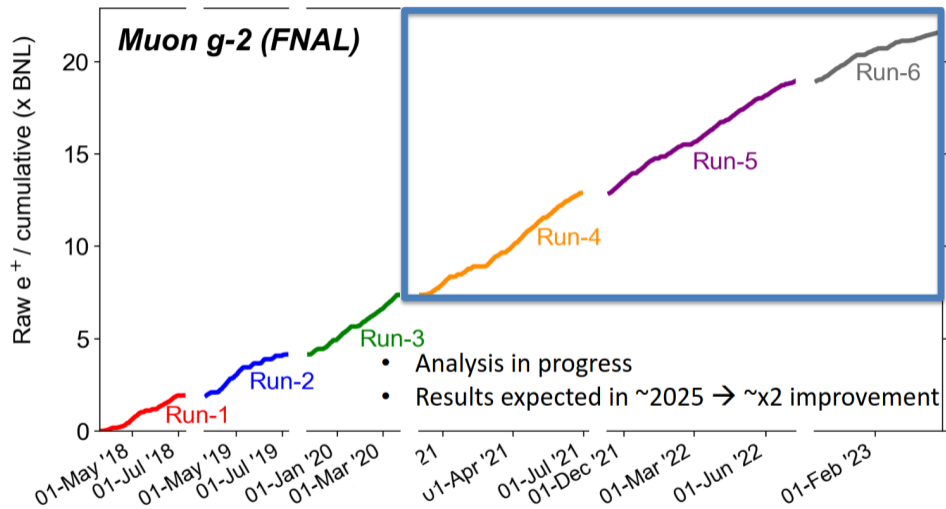
[Aoyama et al., 2006.04822]

← Everything is much more complicated now

[Boccaletti et al., 2407.10913]

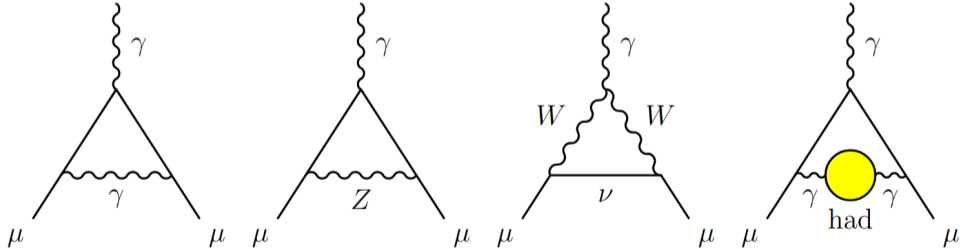
[Muon  $g - 2$ , 2308.06230].

# THE MUON $g - 2$ : EXPERIMENT AT FNAL

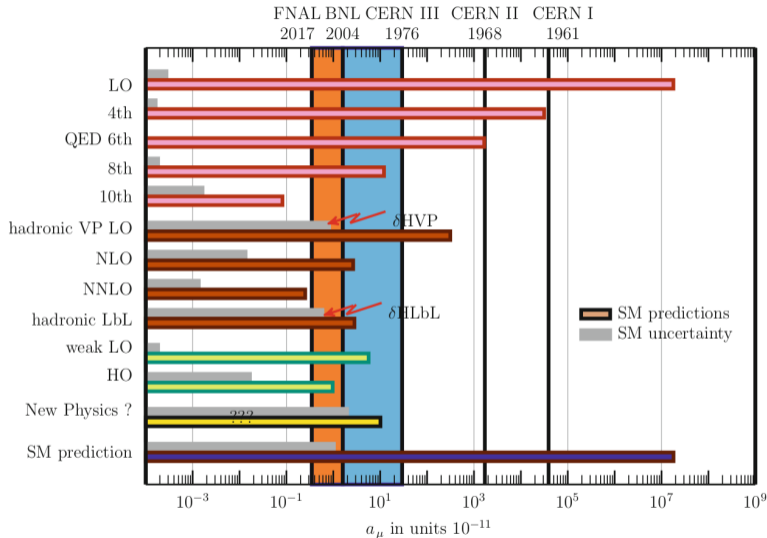


The  $g - 2$  experiment will surpass their initial target precision! [Venanzoni]

# THE STANDARD MODEL PREDICTION FOR $a_\mu$



# THE MUON $g - 2$ : STANDARD MODEL PREDICTION



Muon  $g - 2$  experiments and the sensitivity to various contributions. [Jegerlehner, 2017]



Contribution	Value $\times 10^{11}$
Experiment (E821 + E989)	116 592 059(22)
HVP LO ( $e^+e^-$ )	6931(40)
HVP NLO ( $e^+e^-$ )	-98.3(7)
HVP NNLO ( $e^+e^-$ )	12.4(1)
HVP LO (lattice, $udsc$ )	7116(184)
HLbL (phenomenology)	92(19)
HLbL NLO (phenomenology)	2(1)
HLbL (lattice, $uds$ )	79(35)
HLbL (phenomenology + lattice)	90(17)
QED	116 584 718.931(104)
Electroweak	153.6(1.0)
HVP ( $e^+e^-$ , LO + NLO + NNLO)	6845(40)
HLbL (phenomenology + lattice + NLO)	92(18)
Total SM Value	116 591 810(43)
Difference: $\Delta a_\mu := a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	249(48)

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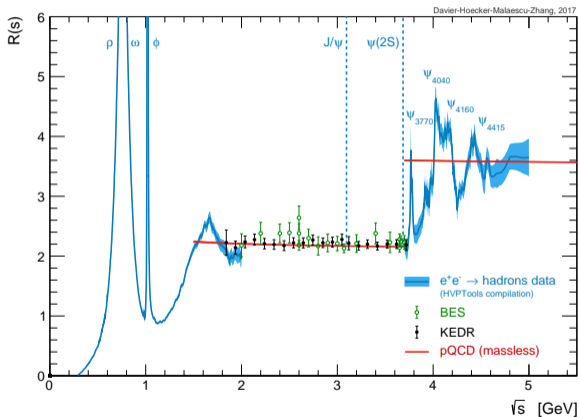
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- 5.1  $\sigma$  discrepancy between **experiment** and prediction?
- Uncertainty from  $a_\mu^{\text{hvp,LO}}$  dominates  $a_\mu$ .
- A number of new inputs for the **upcoming second White Paper...**

- The QED contribution completely dominates the SM prediction.
- Recent developments [Makiko Nio at KEK]:
  - ▶ Small inconsistencies in the 10th order contributions are resolved.
  - ▶ Uncertainty halved due to improved measurement of the fine-structure constant.
- Improved precision for the electroweak contribution [Martin Hoferichter at KEK].
- Both changes are completely irrelevant for the final uncertainty.

# THE STANDARD MODEL PREDICTION FOR $a_\mu$

**HADRONIC CONTRIBUTIONS TO  $a_\mu$**

R-ratio:  $R(s) = \frac{\sigma^0(e^+e^- \rightarrow \text{hadrons}(+\gamma))}{\sigma_{\text{pt}}}$ ,  $\sigma_{\text{pt}} = \frac{4\pi\alpha^2}{3s}$

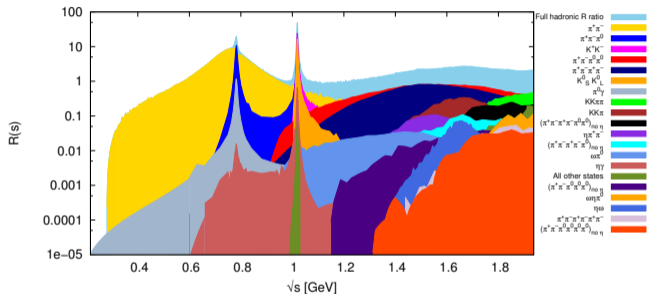


- Data-driven extraction of the HVP contribution via dispersion integral

$$a_{\mu}^{\text{HVP,LO}} = \frac{\alpha^2}{3\pi^2} \int_{M_{\pi}^2}^{\infty} \frac{K(s)}{s} R(s) ds$$

# $a_\mu^{\text{hvp}}$ : THE DISPERSIVE APPROACH

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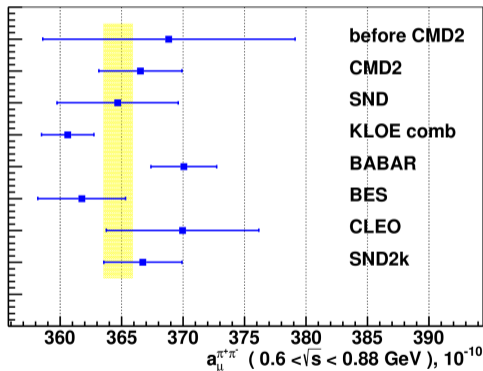
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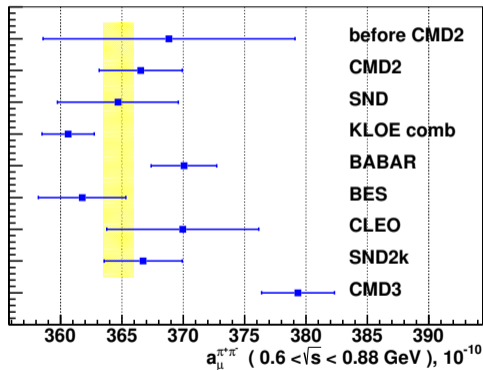
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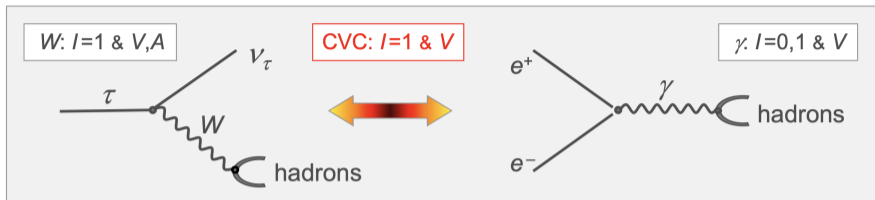


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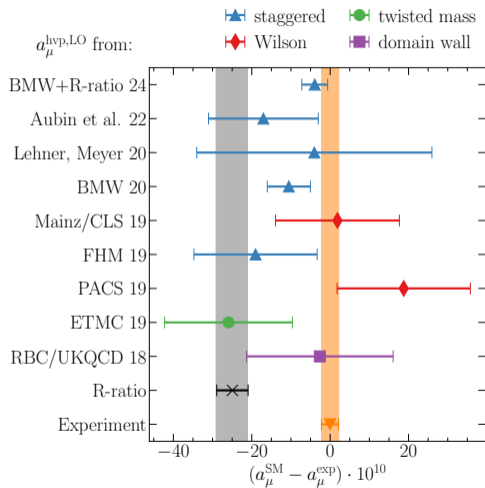
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- $R$ -ratio constructed from exclusive channels  
→ source of systematic uncertainty.
- The discrepancies are **not understood**.

- Can missing NNLO terms in the MC generators (e.g. PHOKHARA) that are used to remove radiative corrections affect experimental results significantly?
- First, partly preliminary, analyses: Not the case for [BaBar](#), [KLOE](#), [BES III](#).
- [The RadioMonteCarlow 2 Effort](#) assesses different Monte Carlo codes.
- New data (KLOE, BaBar, Belle II, BES III, SND) will come in eventually
- All of this will be too late for the WP update.



- Can use  $\tau$  spectral functions to evaluate the LO HVP  
[\[Alemany et al., hep-ph/9703220\]](#) [\[Zhang at KEK\]](#).
- Requires (model-dependent) isospin-breaking corrections.
- Model-independent evaluations on the way
  - ▶ on the lattice [\[Bruno et al.\]](#)
  - ▶ with dispersive methods [\[Cottini et al.\]](#)
 and will allow to include the  $\tau$  data in a data-driven evaluation.

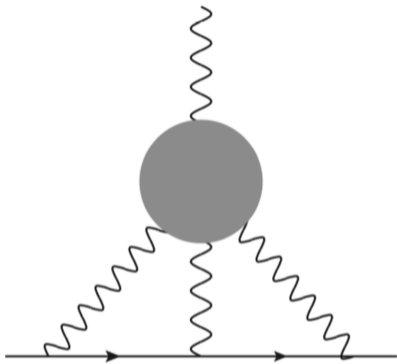


- One sub-percent determination of  $a_\mu^{\text{hvp}}$  from the lattice [BMWc, 2002.12347]: In tension with the dispersive result.
- Several new lattice results at the percent level in the (very) near future to come.
- Consistency would allow to replace the data-driven estimate.

## Goal

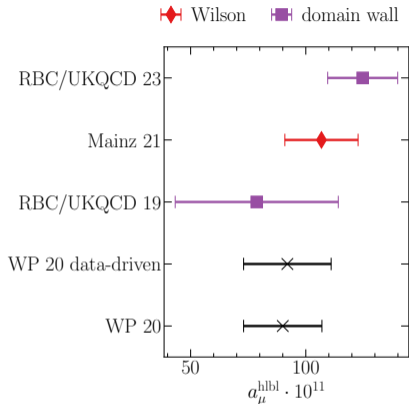
Several lattice results at  $< 0.5\%$  precision.

# HADRONIC LIGHT-BY-LIGHT SCATTERING



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 $O(\alpha^3)$ , target precision: 10%.

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- White paper recommended value:

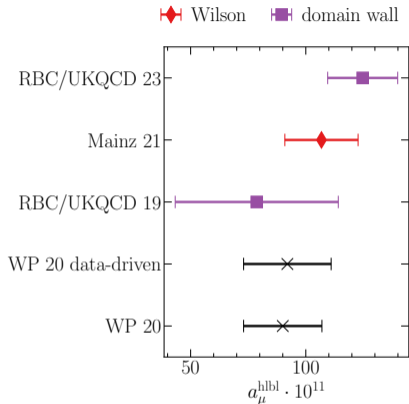
$$a_\mu^{\text{hlbl}} = (92 \pm 18) \cdot 10^{-11}$$

- Two lattice calculations since then, [Mainz 21, 2104.02632, 2204.08844] and [RBC/UKQCD 23, 2304.04423].

- Preliminary results by ETMc and BMW.

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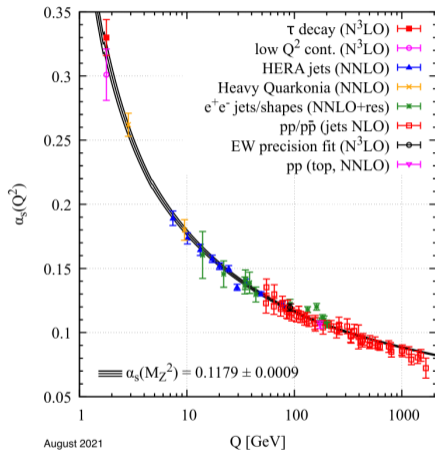
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- Lattice also enters data-driven determination via transition form factors.
- Lattice and data-driven computations are an outstanding success.
- No obvious tension at the current level of uncertainty.

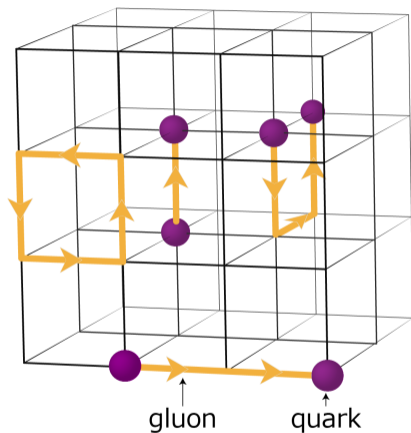


$a_{\mu}^{\text{hvp}}$  **ON THE LATTICE**



- QCD is a strongly coupled theory in the hadronic regime at  $Q \sim 300$  MeV.
- Perturbative expansion fails below 1 GeV.

<sup>1</sup>[PDG, PTEP **2022** (2022), o83Co1]



- QCD is a strongly coupled theory in the hadronic regime at  $Q \sim 300 \text{ MeV}$ .
- Perturbative expansion fails below  $1 \text{ GeV}$ .
- Formulate the theory
  - ▶ on a finite grid  $\rightarrow$  regulator  $\Lambda_{UV}$ .
  - ▶ in finite volume  $\rightarrow \Lambda_{IR}$ .
  - ▶ in Euclidean space-time
  - ▶ as a Boltzmann distribution
- Compute expectation values  $\langle O \rangle$  by sampling the QCD path integral with Markov Chain Monte Carlo methods.

<sup>2</sup><http://www.jicfus.jp/en/promotion/pr/mj/guido-cossu/>

## The QCD Lagrange density

$$\mathcal{L}_{\text{QCD}} = \sum_{f=1}^{N_f} \bar{\psi}_f (\not{D} + m_f) \psi_f + \frac{1}{4} \text{tr} F_{\mu\nu} F^{\mu\nu}$$

- Contains  $N_f + 1$  **bare** parameters (gauge coupling and  $N_f$  quark masses)
- Renormalize the theory from hadronic input, e.g.,  $m_\Omega, m_\pi, m_K, m_{D_s}, m_{B_s}$ .  
→ All other observables are **predictions**.
- Freedom of choice on how to discretize  $\mathcal{L}_{\text{QCD}}$ : Wilson, twisted mass, staggered, domain wall, overlap, ...
- *Ab initio* predictions after lifting the cutoffs:
  - ▶  $\Lambda_{\text{IR}}$ : Infinite-volume limit.
  - ▶  $\Lambda_{\text{UV}}$ : Continuum limit.

- Compute  $a_\mu^{\text{hvp}}$  via [Laurup et al.] [Blum, hep-lat/0212018]

$$a_\mu^{\text{hvp}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dQ^2 f(Q^2) \hat{\Pi}(Q^2), \quad \text{with} \quad \hat{\Pi}(Q^2) = 4\pi^2 [\Pi(Q^2) - \Pi(0)]$$

from a known QED kernel function  $f(Q^2)$  and the polarization tensor

$$\Pi_{\mu\nu}(Q) = \int d^4x e^{iQ \cdot x} \langle j_\mu^{\text{em}}(x) j_\nu^{\text{em}}(0) \rangle = (Q_\mu Q_\nu - \delta_{\mu\nu} Q^2) \Pi(Q^2).$$

- $a_\mu^{\text{hvp}}$  in the time-momentum representation (TMR) [Bernecker, Meyer, 1107.4388],

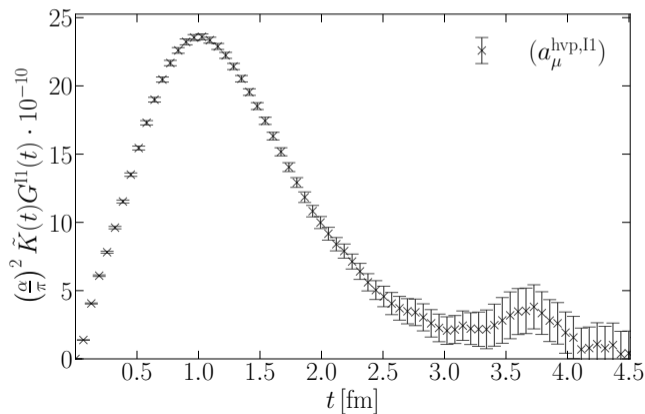
$$a_\mu^{\text{hvp}} := \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt G(t) \tilde{K}(t) \quad \text{with the known QED kernel function } \tilde{K}(t),$$

in terms of the zero-momentum vector correlator  $G(t)$  (de facto standard).

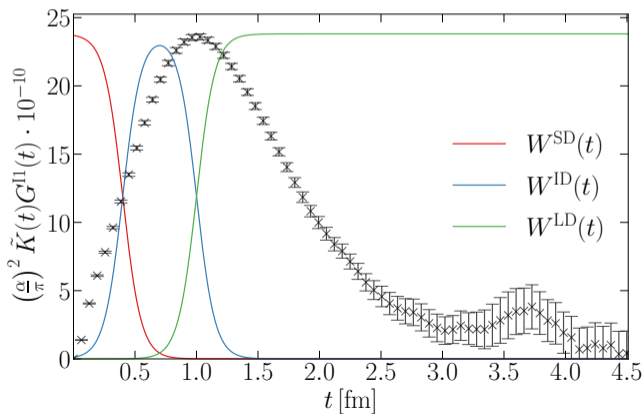
- Alternative: coordinate space method [Meyer, 1706.01139] [Chao et al., 2211.15581].

$$(a_\mu^{\text{hvp}}) = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt G(t) \tilde{K}(t),$$

$$G(t) = -\frac{a^3}{3} \sum_{k=1}^3 \sum_{\vec{x}} \langle j_k^{\text{em}}(t, \vec{x}) j_k^{\text{em}}(0) \rangle$$



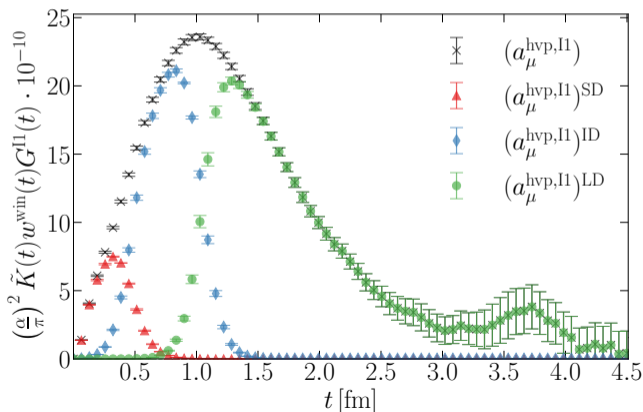
$$(a_\mu^{\text{hvp}})^i = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt G(t) \tilde{K}(t) W^i(t; t_0; t_1), \quad G(t) = -\frac{a^3}{3} \sum_{k=1}^3 \sum_{\vec{x}} \langle j_k^{\text{em}}(t, \vec{x}) j_k^{\text{em}}(0) \rangle$$



- Windows in the TMR: separate short- from long-distance effects [RBC/UKQCD, 1801.07224].

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- Windows in the TMR: separate short- from long-distance effects [RBC/UKQCD, 1801.07224].
- Intermediate window  $a_\mu^{\text{win}}$ :
  - ▶ Cutoff effects suppressed.
  - ▶ No signal-to-noise problem.
  - ▶ Finite-volume effects small.

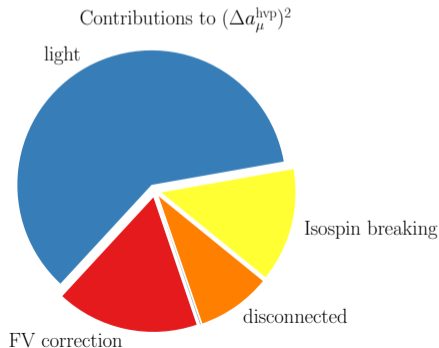
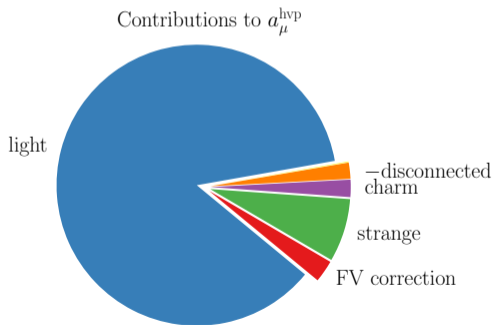


The electromagnetic current

$$j_\mu^{\text{em}} = \frac{2}{3}\bar{u}\gamma_\mu u - \frac{1}{3}\bar{d}\gamma_\mu d - \frac{1}{3}\bar{s}\gamma_\mu s + \frac{2}{3}\bar{c}\gamma_\mu c + \dots = j_\mu^{I=1} + j_\mu^{I=0}$$

from zero-momentum vector-vector correlation functions

$$G^{\text{isoQCD}}(t) = \frac{5}{9}G^{\text{light}}(t) + \frac{1}{9}G^{\text{strange}}(t) + \frac{4}{9}G^{\text{charm}}(t) + G^{\text{disc}}(t) + \dots$$

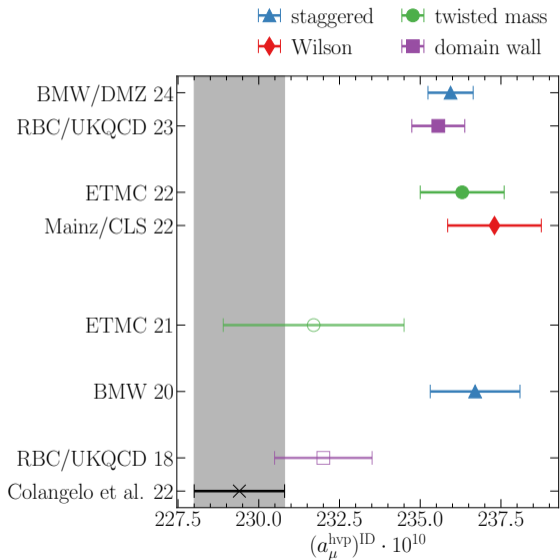


Based on [BMWc, 2002.12347]:  $a_\mu^{\text{hvp}} = 707.5 (5.5) \cdot 10^{-10}$

# $a_{\mu}^{\text{hvp}}$ ON THE LATTICE

## WINDOW OBSERVABLES

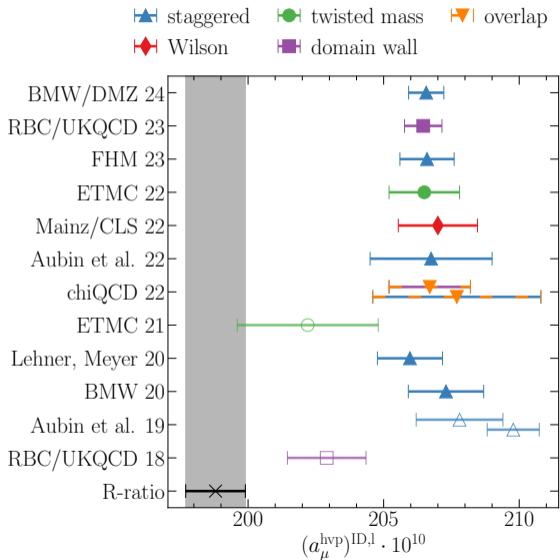
# THE INTERMEDIATE-DISTANCE WINDOW



■  $3.8\sigma$  tension between lattice QCD and data-driven evaluation [Colangelo et al., 2205.12963].

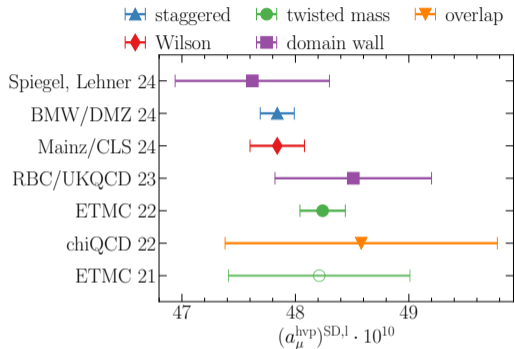
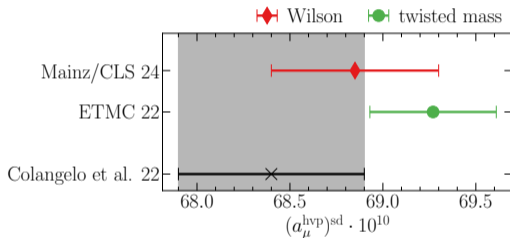
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- This accounts for **50% of the difference** between BMW 20 and the White Paper average for  $a_\mu^{\text{hvp}}$ .
- Agreement across many actions for the light-connected contribution (87% of  $(a_\mu^{\text{hvp}})^{\text{ID}}$ ).
- Data-driven estimate: [Benton et al., 2306.16808]

# THE SHORT-DISTANCE WINDOW

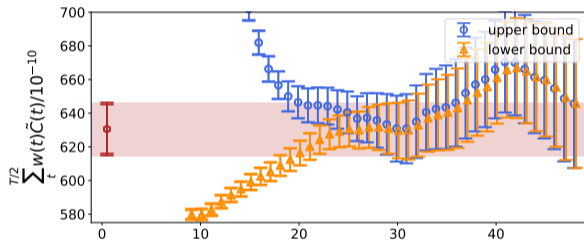


- Continuum extrapolation is the major difficulty for the short-distance window.
- However: Small uncertainties w.r.t. the full HVP.
- No significant difference between lattice and R-ratio - could expect about 1 unit (1.44%) based on what is seen in the intermediate window [SK at al., 2401.11895].

# $a_{\mu}^{\text{hvp}}$ ON THE LATTICE

**DOMINANT SOURCES OF UNCERTAINTY FOR  $a_{\mu}^{\text{hvp}}$**

# CONTROLLING THE LONG-DISTANCE TAIL



Exponential deterioration of the signal-to-noise ratio.

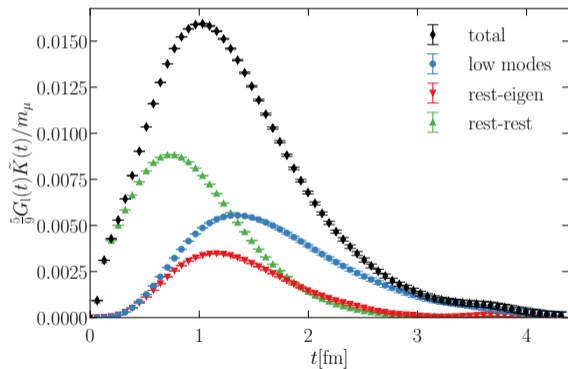
Improve the signal at large  $t$  via:

- **Bounds on the correlator.**
- Noise reduction methods:
  - ▶ Truncated Solver Method
  - ▶ Low Mode Averaging
  - ▶ All Mode Averaging
- Spectral reconstruction of the  $\pi\pi$  contributions.
- Multi-level integration.  
[Dalla Brida et al., 2007.02973]

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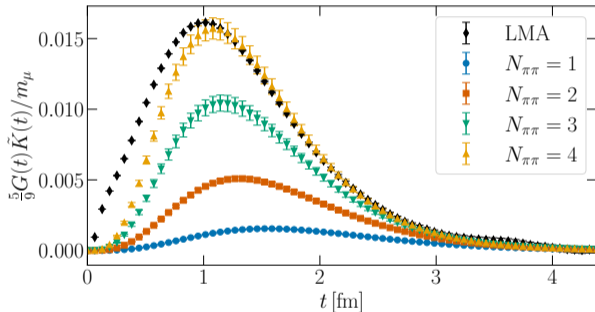
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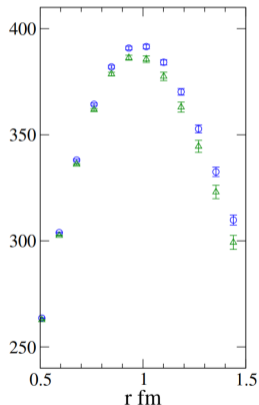
3% finite- $L$  corrections for  $a_\mu^{\text{hvp}}$  at  $m_\pi L = 4$ , mostly in the **isovector channel**.

## ■ EFT and model calculations.

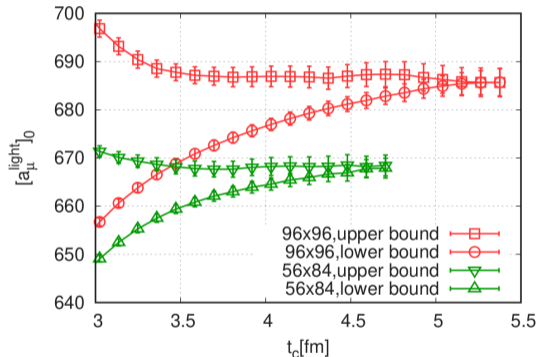
- ▶ NNLO  $\chi$ PT
- ▶ Two-pion spectrum in finite-volume and the timelike pion form factor [Meyer, 1105.1892] [Lellouch and Lüscher, hep-lat/0003023] [Giusti et al., 1808.00887].
- ▶ Pions winding around the torus and the electromagnetic pion form factor [Hansen, Patella, 1904.10010, 2004.03935].
- ▶ Rho-pion-gamma model [Sakurai] [Jegerlehner, Szafron, 1101.2872] [HPQCD, 1601.03071].

## ■ Simulations at $L > 10$ fm [PACS, 1902.00885] [BMWc, 2002.12347].

- ▶ Uncertainty statistics dominated.
- ▶ Show good consistency with models.



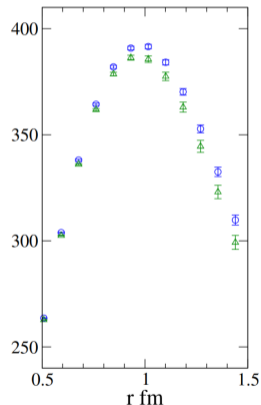
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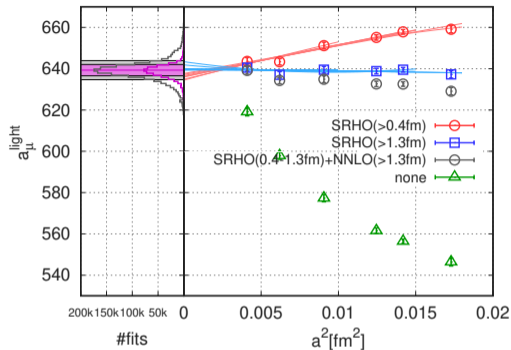
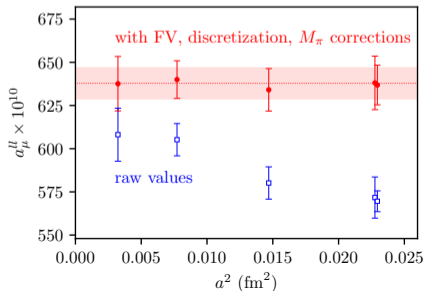
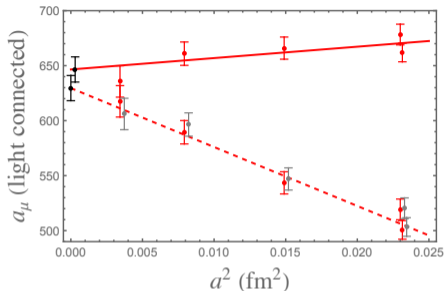
].



Systematic uncertainties from the continuum extrapolation may be dominant.

- Extrapolation to the continuum limit guided by Symanzik effective theory.
  - Cutoff effects start at  $O(a^2)$  in modern lattice calculations.
  - Mandatory to
    - ▶ include  $\geq 4$  resolutions to constrain higher order cutoff effects.
    - ▶ include fine resolutions  $a \leq 0.05$  fm for per-mil uncertainties.
  - Staggered quarks: taste violations distort the pion spectrum.
    - ▶ This is a cutoff effect: Vanishes in the continuum limit.
    - ▶ Taste breaking may introduce non-linear effects (in  $a^2$ ).
- Corrections applied at finite lattice spacing.

# THE CONTINUUM LIMIT: STAGGERED QUARKS

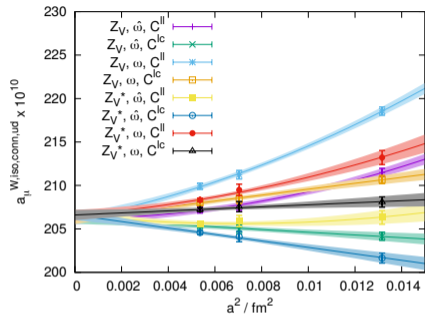
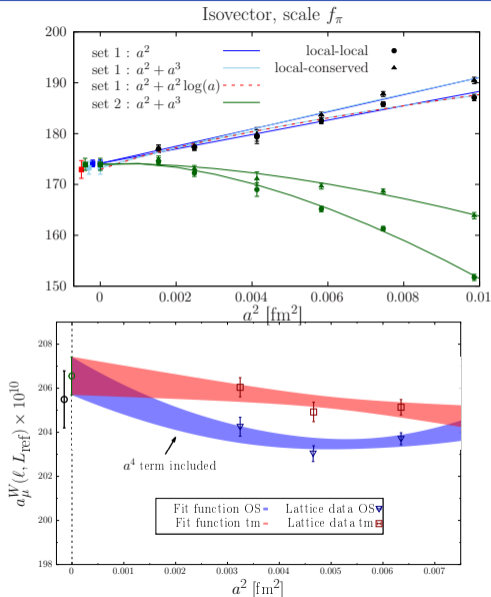


- Continuum extrapolations of  $a_\mu^{\text{hvp}}$  computed with staggered quarks.
- Compare raw and corrected data.

[Aubin et al., 2204.12256] [BMWc, 2002.12347]

[Fermilab, HPQCD, MILC, 1902.04223]

# THE CONTINUUM LIMIT: INTERMEDIATE WINDOW



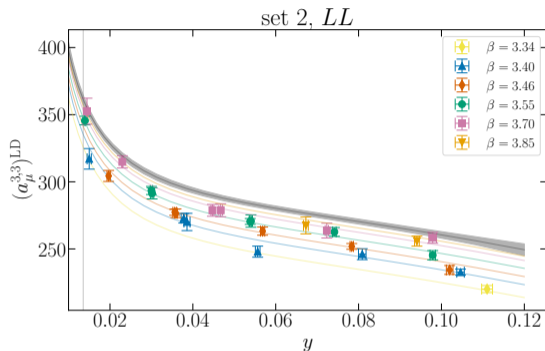
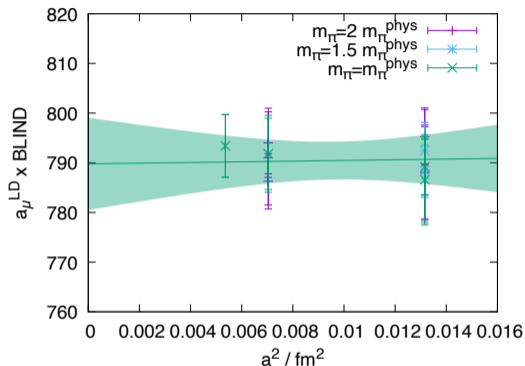
- Different discretization prescriptions have to agree in the continuum.
- Strong cross-check for **valence** cutoff effects.

[Mainz, 2206.06582] [RBC/UKQCD, 2301.08696]  
 [ETMC, 2206.15084]

# $a_{\mu}^{\text{hvp}}$ ON THE LATTICE

## THE LONG-DISTANCE CONTRIBUTION

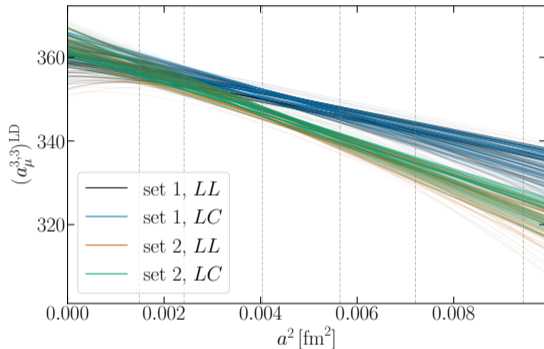
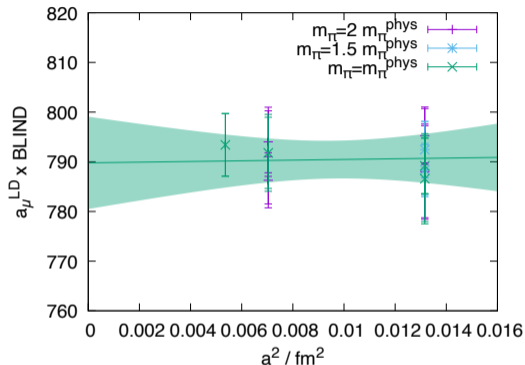
# THE LONG-DISTANCE CONTRIBUTION



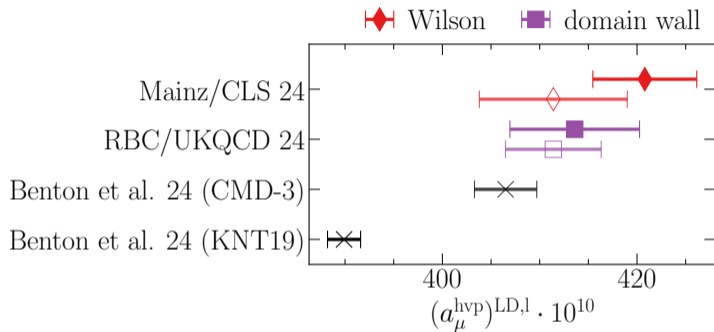
- Two new high-precision evaluations of the long-distance light-connected contribution [RBC/UKQCD, 2410.20590][Mainz, 2411.07969][SK at KEK].
- The main stepping stone towards  $a_\mu^{\text{hvp}}$  at sub-percent precision.



# THE LONG-DISTANCE CONTRIBUTION

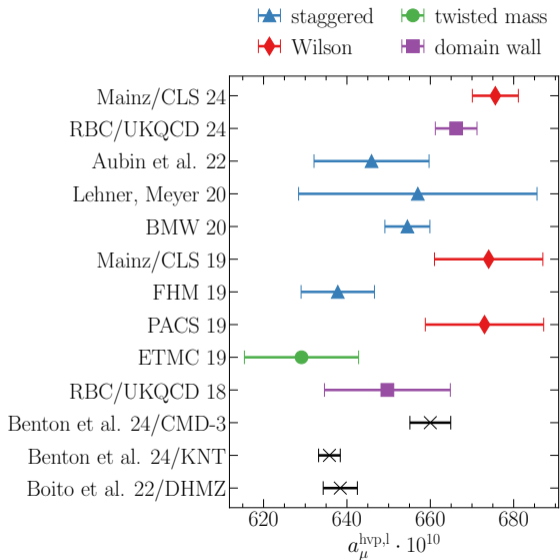


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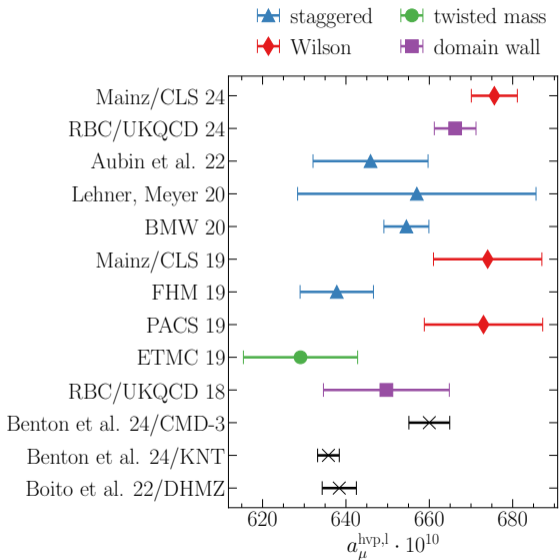
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- The main stepping stone towards  $a_\mu^{\text{hvp}}$  at sub-percent precision.
- Significant scheme dependence! Agreement in the same scheme for isoQCD.

# THE LIGHT-CONNECTED CONTRIBUTION



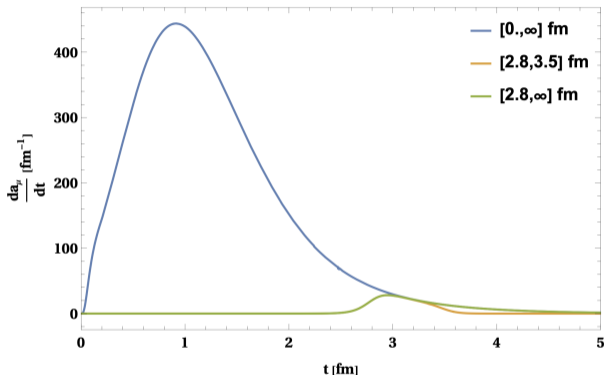
- The light-connected contribution makes up for  $> 90\%$  of  $a_\mu^{\text{hvp}}$ .
- No consistent isoQCD scheme in this comparison!
- Unfortunately no update by BMW in [BMWc, 2407.10913]...

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- Unfortunately no update by BMW in [BMWc, 2407.10913]...
- No clear tension when using CMD-3 for the data-driven estimate [Benton et al., 2411.06637].

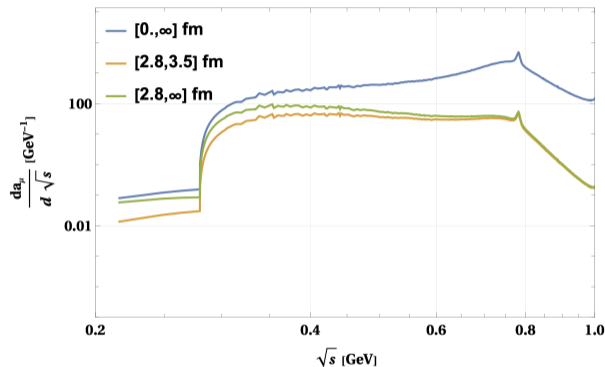
# THE LIGHT-CONNECTED CONTRIBUTION



[Lellouch at KEK]

- [BMWc, 2407.10913] uses a data-driven of the contribution beyond 2.8 fm.
  - ▶ About 4% of  $a_\mu^{\text{hvp}}$ .
  - ▶ Smaller statistical uncertainties.
  - ▶ Smaller finite-volume effects.
  - ▶ Smaller cutoff effects with staggered fermions.

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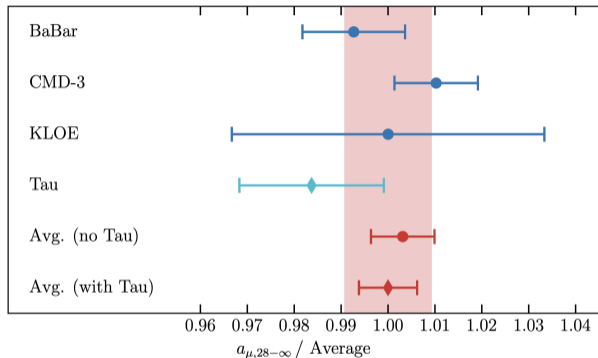


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  - ▶ Smaller statistical uncertainties.
  - ▶ Smaller finite-volume effects.
  - ▶ Smaller cutoff effects with staggered fermions.
- Contribution below the  $\rho$  peak seems to be compatible across experiments.
- No updated pure lattice result for  $a_{\mu}^{\text{hVP}}$  or  $(a_{\mu}^{\text{hVP}})^{\text{(LD)}}$ .

# $a_{\mu}^{\text{hvp}}$ ON THE LATTICE

## ISOSPIN BREAKING EFFECTS



Need to include  $O(\frac{m_u - m_d}{\Lambda_{\text{QCD}}})$  and  $O(\alpha)$  effects for per-mil precision.

- Various ways to compute isospin breaking corrections:
  - ▶ **Perturbative expansion around isospin symmetric QCD** [RM123, 1303.4896].
  - ▶ Simulation of dynamical QCD+QED [CSSM/QCDSF/UKQCD] [RC\*, 2212.11551].
  - ▶ Infinite volume QED [RBC/UKQCD, 1801.07224] [Biloshytskyi et al., 2209.02149]
- Major challenge: Formulation of QED in a finite box.
- $\text{QED}_L$ : Finite-volume corrections scale as  $O(1/L^3)$  [Bijnens et al., 1903.10591]  
→ sufficient for the precision goal.

# QED AND STRONG ISOSPIN BREAKING: RESULTS

Overview of published results - contributions to  $a_\mu \times 10^{10}$

- Strong isospin breaking:  
Five groups agree within  $1\sigma$ .



6.60(63)(53)		BMW
10.6(4.3)(6.8)		RBC/UKQCD
6.0(2.3)		ETM
7.7(3.7)	9.0(2.3)	FHM
9.0(0.8)(1.2)		LM

BMW [Nature 593 (2021) 7857, 51-55]  
RBC/UKQCD [Phys.Rev.Lett. 121 (2018) 2, 022003]  
ETM [Phys. Rev. D 99, 114502 (2019)]  
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Adapted from [V. Gülpers @ Lattice HVP workshop 2020]

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Overview of published results - contributions to  $a_\mu \times 10^{10}$



BMW	-1.23(40)(31)
RBC/UKQCD	5.9(5.7)(1.7)
ETM	1.1(1.0)



	-0.55(15)(10)	BMW
	-6.9(2.1)(2.0)	RBC/UKQCD

- Strong isospin breaking:  
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- QED: agreement on the total valence contribution.



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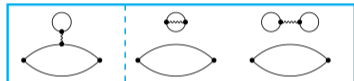
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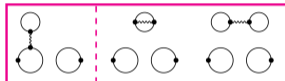
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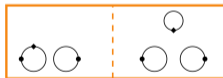
$-0.0093(86)(95)$   $0.37(21)(24)$  BMW



$0.011(24)(14)$   $-0.040(33)(21)$  BMW



$6.60(63)(53)$  BMW  
 $10.6(4.3)(6.8)$  RBC/UKQCD  
 $6.0(2.3)$  ETM  
 $7.7(3.7)$   $9.0(2.3)$  FHM  
 $9.0(0.8)(1.2)$  LM



$-4.67(54)(69)$  BMW

BMW [Nature 593 (2021) 7857, 51-55]  
 RBC/UKQCD [Phys.Rev.Lett. 121 (2018) 2, 022003]  
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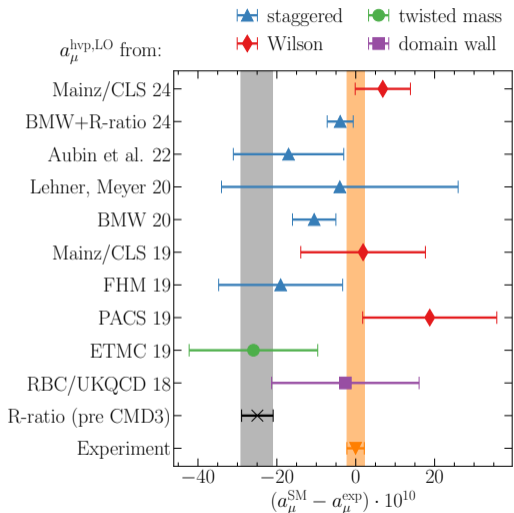
■ QED: agreement on the total  
 valence contribution.

■ One complete calculation  
 [BMWc, 2002.12347]:  
 $\delta a_\mu^{\text{hvp}} = 0.5(1.4) \cdot 10^{-10}$

■ Work in progress:  
 [Mainz, 2206.06582]  
 [RBC/UKQCD, Lattice 2022]  
 [BMWc, Lattice 2022]  
 [FHM, 2212.12031]  
 [Harris et al., 2301.03995]  
 [Mainz, 2411.07969]

Adapted from [v. Gülpers @ Lattice HVP workshop 2020]

# HADRONIC VACUUM POLARIZATION CONTRIBUTION TO THE MUON $g - 2$



- Adding further contributions and isospin breaking corrections allows to add Mainz/CLS 24 to the overview for  $a_\mu^{\text{hvp}}$  [Mainz, 2411.07969].
- More independent data points to follow.
- Going significantly below 1% uncertainty remains a challenge.
- Combination with data-driven estimates and MUonE in a few years?

# CONCLUSIONS

- Final result of the FNAL  $g - 2$  experiment in spring 2025.
- New White Paper will provide an updated SM prediction before that.
- The main difference compared to 2020 will be  $a_\mu^{\text{hvp}}$ 
  - ▶ Currently no reliable estimate from data-driven dispersive methods.
  - ▶ Multiple high-precision lattice QCD results.
  - ▶ An *ab initio* lattice QCD result for  $a_\mu^{\text{hvp}}$  can replace the old estimate.
- There is little hope for the anomaly in  $(g - 2)_\mu$  to survive.
- More work on the SM prediction is required to match the expected experimental precision.