

PDFs &

New developments on GPDs from lattice QCD and potential synergies

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Temple University

MITP
SCIENTIFIC
PROGRAM

Next Generation Perturbative QCD for Hadron Structure:
Preparing for the Electron-Ion Collider
Oktober 21 – 31, 2024

<https://indico.mitp.uni-mainz.de/event/370>

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Theoretical Physics

The banner features a central illustration of a hadron structure with a central core and surrounding green and black spheres, with a wavy line representing a photon or gluon. The background is black with a red triangle on the left.

Next Generation Perturbative QCD for Hadron Structure: Preparing for the Electron-Ion Collider

Disclaimer

The field of GPDs is still developing and sources of systematic uncertainties have not been fully addressed

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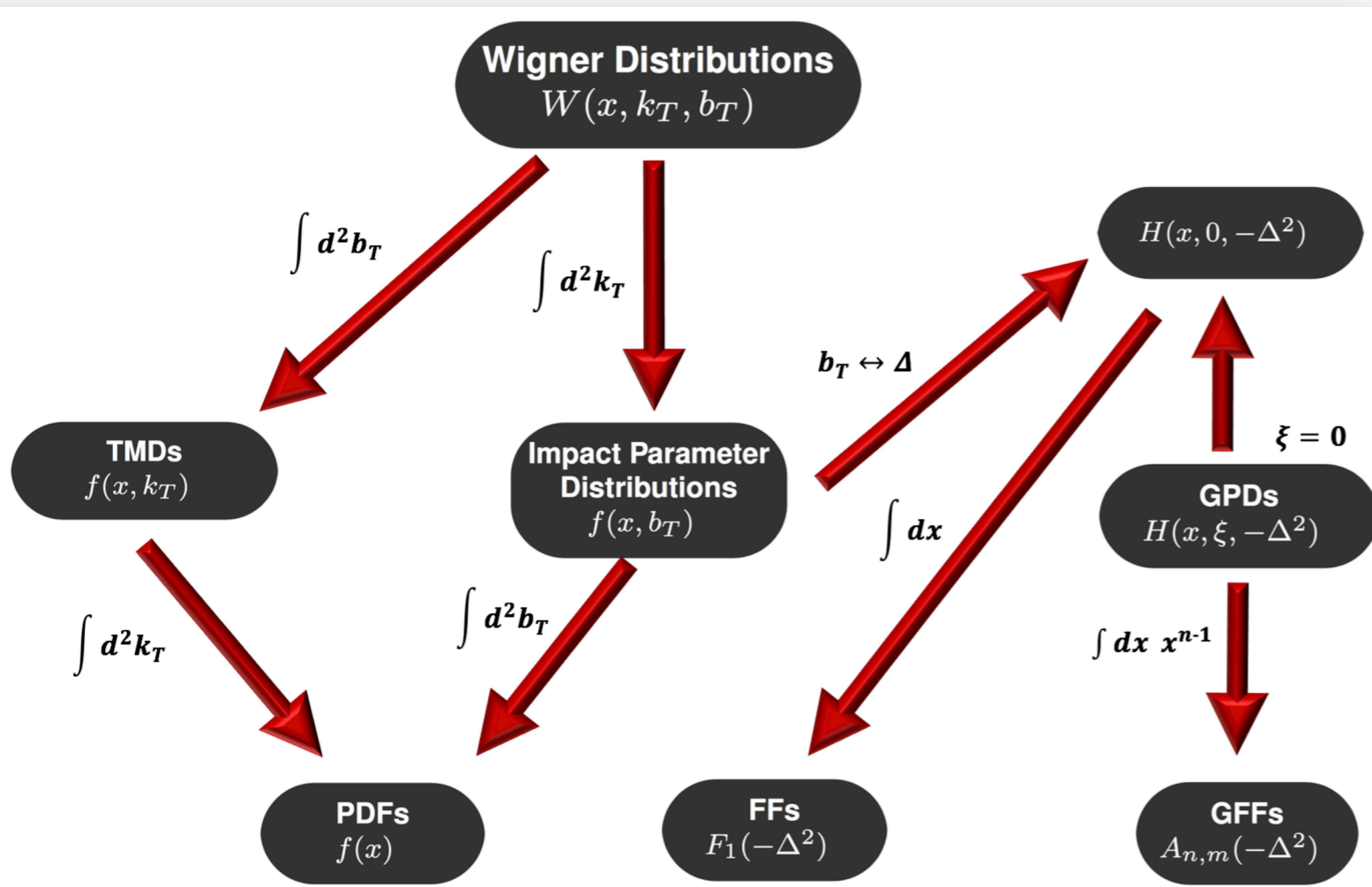
- **Discretization effects**
- **physical pion mass**
- **Volume effects**
- **inverse problem**
- **matching formalism**
- **connection to light-cone, higher twist contaminations, ...**
- **...**

Accessing PDFs/GPDs from lattice QCD

Nucleon Characterization

Wigner distributions

- ★ Fully characterize partonic structure of hadrons
- ★ Provide multi-dim images of the parton distributions in phase space

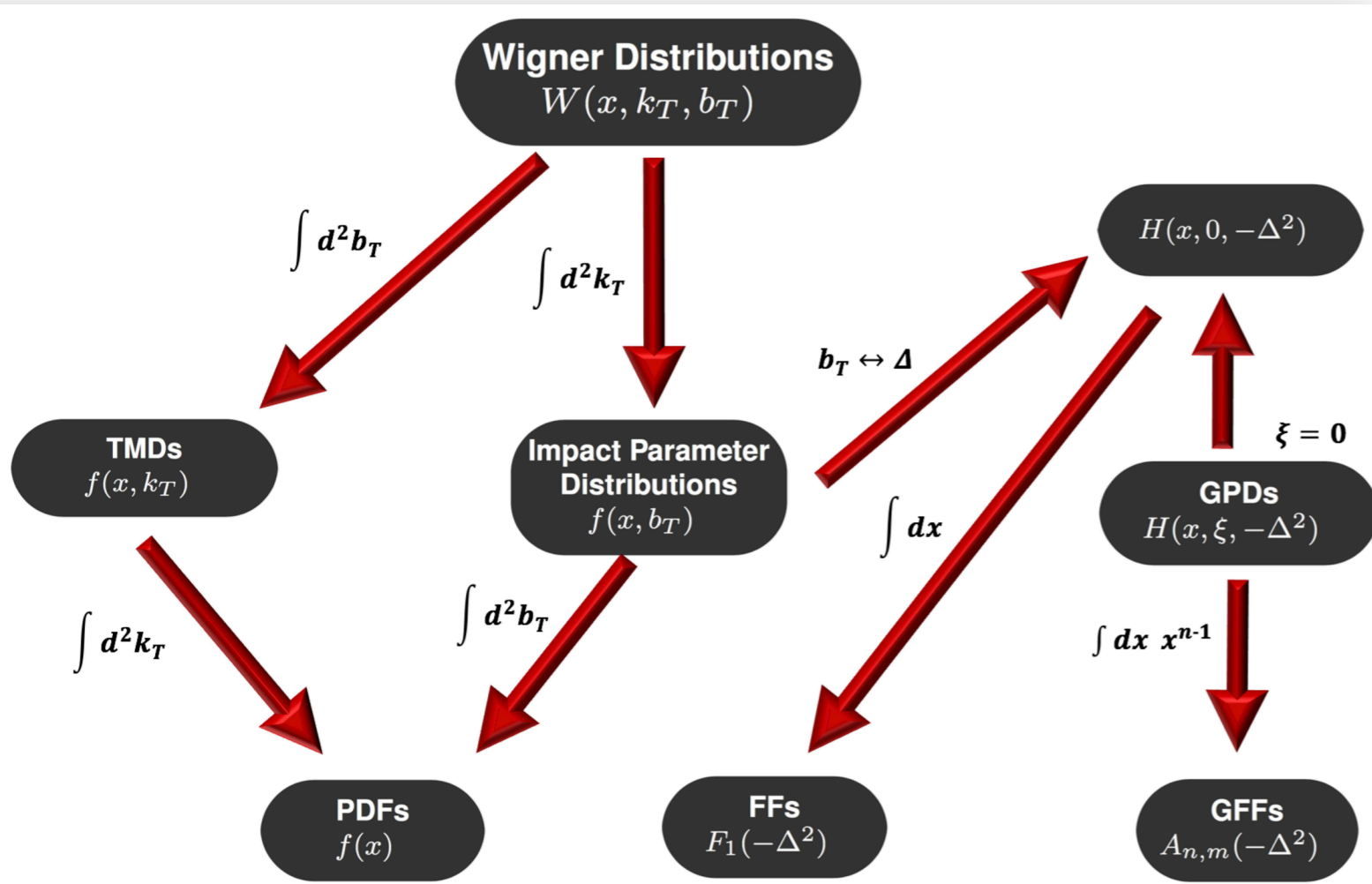


- ★ Correlations between momenta, positions, spins
- ★ Information on the hadron's mechanical properties (OAM, pressure, etc.)

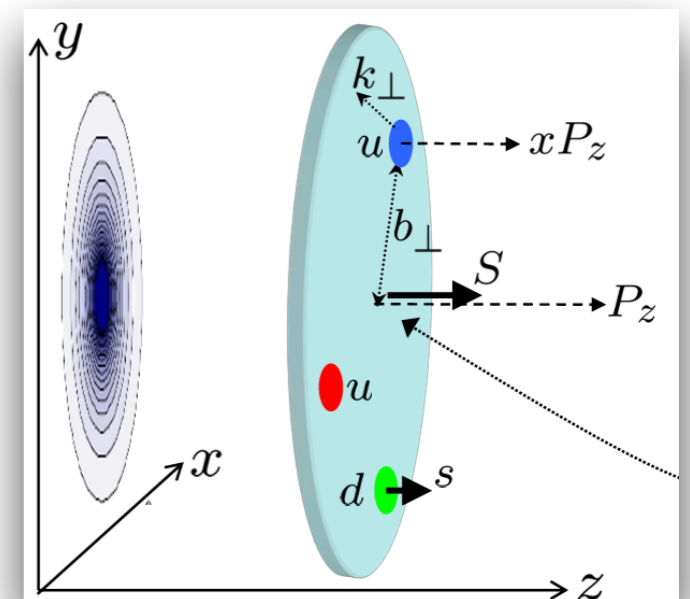
Nucleon Characterization

Wigner distributions

- ★ Fully characterize partonic structure of hadrons
- ★ Provide multi-dim images of the parton distributions in phase space



- ★ Partons contain information on
 - x : longitudinal momentum fraction
 - k_T : transverse momentum
 - b_\perp : impact parameter



- ★ Correlations between momenta, positions, spins
- ★ Information on the hadron's mechanical properties (OAM, pressure, etc.)

Accessing information on PDFs/GPDs

- ★ In parton model, physical picture valid for infinite momentum frame
[R. P. Feynman, Phys. Rev. Lett. 23, 1415 (1969)]
- ★ PDFs parameterized via matrix elements of nonlocal light-cone operators

$$f(x) = \frac{1}{4\pi} \int dy^- e^{-ixP^+y^-} \langle P, S | \bar{\psi}_f \gamma^+ \mathcal{W} \psi_f | P, S \rangle$$

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- ★ **Mellin moments**
(local OPE expansion)

$$\bar{q}\left(-\frac{1}{2}z\right) \gamma^\sigma W\left[-\frac{1}{2}z, \frac{1}{2}z\right] q\left(\frac{1}{2}z\right) = \sum_{n=0}^{\infty} \frac{1}{n!} z_{\alpha_1} \dots z_{\alpha_n} \left[\bar{q} \gamma^\sigma \overleftrightarrow{D}^{\alpha_1} \dots \overleftrightarrow{D}^{\alpha_n} q \right]$$

local operators

$$\langle N(P') | \mathcal{O}_V^{\mu\mu_1 \dots \mu_{n-1}} | N(P) \rangle \sim \sum_{\substack{i=0 \\ \text{even}}}^{n-1} \left\{ \gamma^{\mu} \Delta^{\mu_1} \dots \Delta^{\mu_i} \bar{P}^{\mu_{i+1}} \dots \bar{P}^{\mu_{n-1}} \left[A_{n,i}(t) - i \frac{\Delta_\alpha \sigma^{\alpha\mu}}{2m_N} \Delta^{\mu_1} \dots \Delta^{\mu_i} \bar{P}^{\mu_{i+1}} \dots \bar{P}^{\mu_{n-1}} B_{n,i}(t) \right] + \frac{\Delta^\mu \Delta^{\mu_1} \dots \Delta^{\mu_{n-1}}}{m_N} C_{n,0}(\Delta^2) \Big|_{n \text{ even}} \right\} + \frac{\Delta^\mu \Delta^{\mu_1} \dots \Delta^{\mu_{n-1}}}{m_N} C_{n,0}(\Delta^2) \Big|_{n \text{ even}} U(P)$$

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Reconstruction of PDFs/GPDs very challenging

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- ★ Matrix elements of nonlocal operators
(quasi-GPDs, pseudo-GPDs, ...)

$$\langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z,0) \Psi(0) | N(P_i) \rangle_\mu$$

Nonlocal operator with Wilson line

$$\langle N(P') | O_V^\mu(x) | N(P) \rangle = \bar{U}(P') \left\{ \gamma^\mu H(x, \xi, t) + \frac{i\sigma^{\mu\nu} \Delta_\nu}{2m_N} E(x, \xi, t) \right\} U(P) + \text{ht},$$

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This talk

Novel Approaches

- ★ Hadronic tensor [K.F. Liu, S.J. Dong, PRL 72 (1994) 1790, K.F. Liu, PoS(LATTICE 2015) 115]
- Auxiliary scalar quark [U. Aglietti et al., Phys. Lett. B441, 371 (1998), arXiv:hep-ph/9806277]
- Fictitious heavy quark [W. Detmold, C. J. D, Lin, Phys. Rev. D73, 014501 (2006)]
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- Higher moments [Z. Davoudi, M. Savage, Phys. Rev. D86, 054505 (2012)]
- Quasi-distributions (LaMET) [X. Ji, PRL 110 (2013) 262002, arXiv:1305.1539; Sci. China PPMA. 57, 1407 (2014)]
- Compton amplitude and OPE [A. Chambers et al. (QCDSF), PRL 118, 242001 (2017), arXiv:1703.01153]
- Pseudo-distributions [A. Radyushkin, Phys. Rev. D 96, 034025 (2017), arXiv:1705.01488]
- Good lattice cross sections [Y-Q Ma & J. Qiu, Phys. Rev. Lett. 120, 022003 (2018), arXiv:1709.03018]
- PDFs without Wilson line [Y. Zhao Phys.Rev.D 109 (2024) 9, 094506, arXiv:2306.14960]
- Moments of PDFs of any order [A. Shindler, Phys.Rev.D 110 (2024) 5, L051503, arXiv:2311.18704]

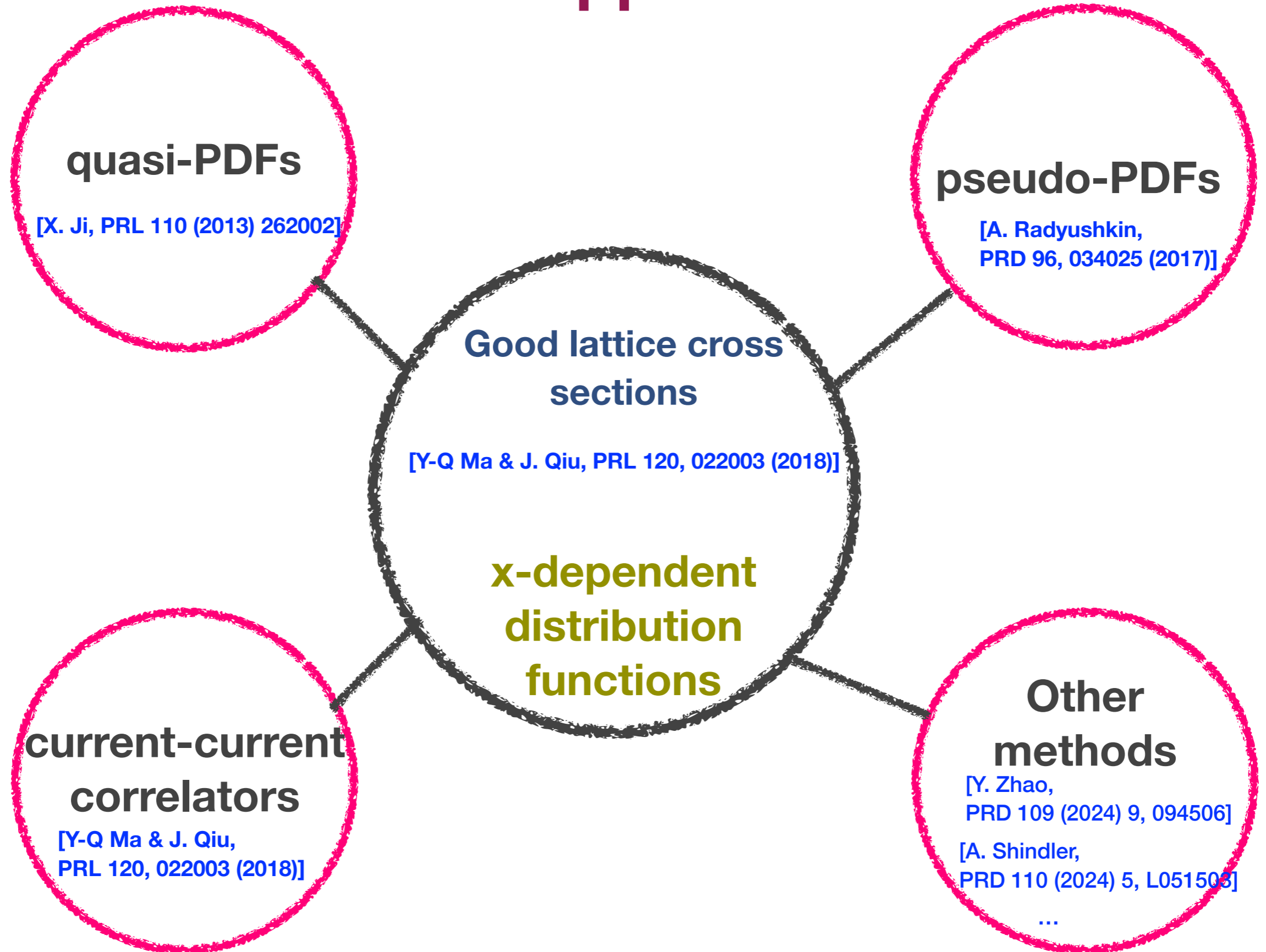
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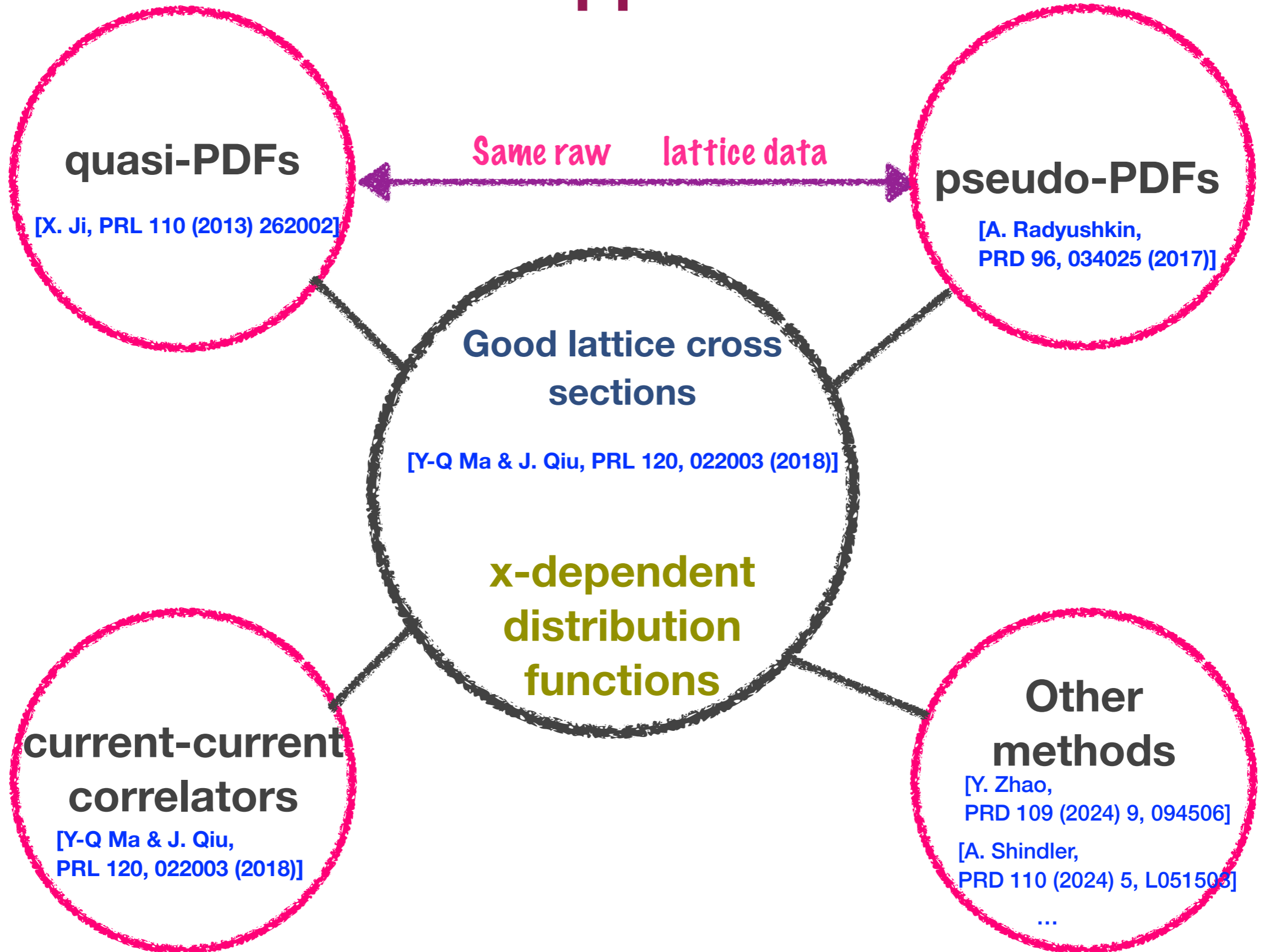
★ Reviews of methods and applications

- *A guide to light-cone PDFs from Lattice QCD: an overview of approaches, techniques and results*
K. Cichy & M. Constantinou (invited review) Advances in HEP 2019, 3036904, arXiv:1811.07248
- *Large Momentum Effective Theory*
X. Ji, Y.-S. Liu, Y. Liu, J.-H. Zhang, and Y. Zhao (2020), 2004.03543
- *The x-dependence of hadronic parton distributions: A review on the progress of lattice QCD*
M. Constantinou (invited review) Eur. Phys. J. A 57 (2021) 2, 77, arXiv:2010.02445

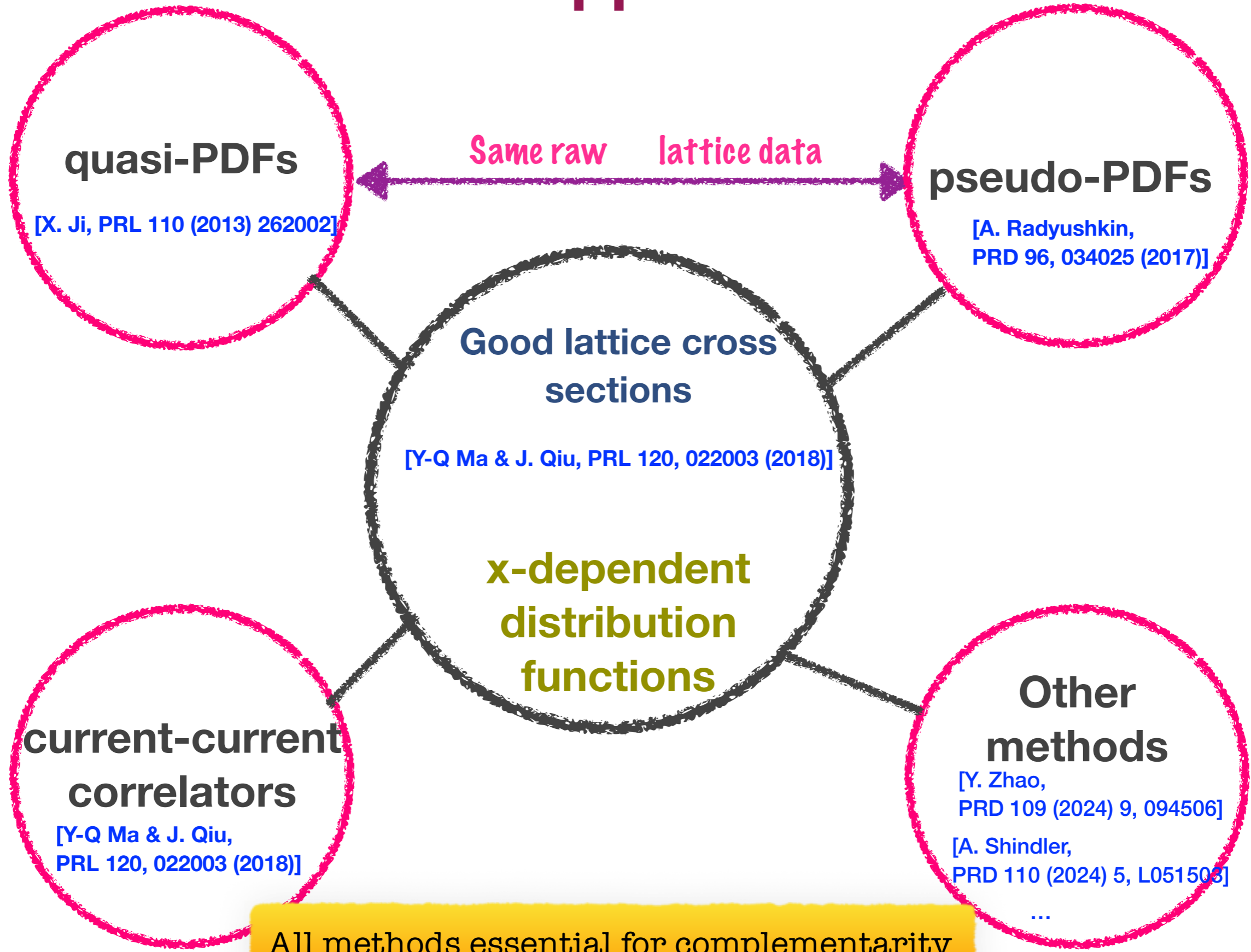
Novel Approaches



Novel Approaches



Novel Approaches



All methods essential for complementarity



Well-studied “novel” methods for PDFs/GPDs in LQCD

Matrix elements of non-local operators (space-like separated fields) with **boosted hadrons**

$$\mathcal{M}(P_f, P_i, z) = \langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_\mu$$

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quasi-PDFs

pseudo-ITD

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$$\mathfrak{M}(\nu, \xi, t; z_3^2) \equiv \frac{\mathcal{M}(\nu, \xi, t; z_3^2)}{\mathcal{M}(0,0,0; z^2)} \quad (\nu = z \cdot p)$$

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Matching in momentum space
(Large Momentum
Effective Theory)

Matching in ν space

$$Q(\nu, \mu^2) = \int_{-1}^1 dx e^{i\nu x} q(x, \mu^2)$$

Light-cone PDFs & GPDs

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Light-cone PDFs & GPDs

Calculation very taxing!

- length of the Wilson line (z)
- nucleon momentum boost (P_3)
- momentum transfer (t)
- skewness (ξ)

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- } PDFs, GPDs
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PDFs:

The unpolarized case

See also QCD@LHC 2024

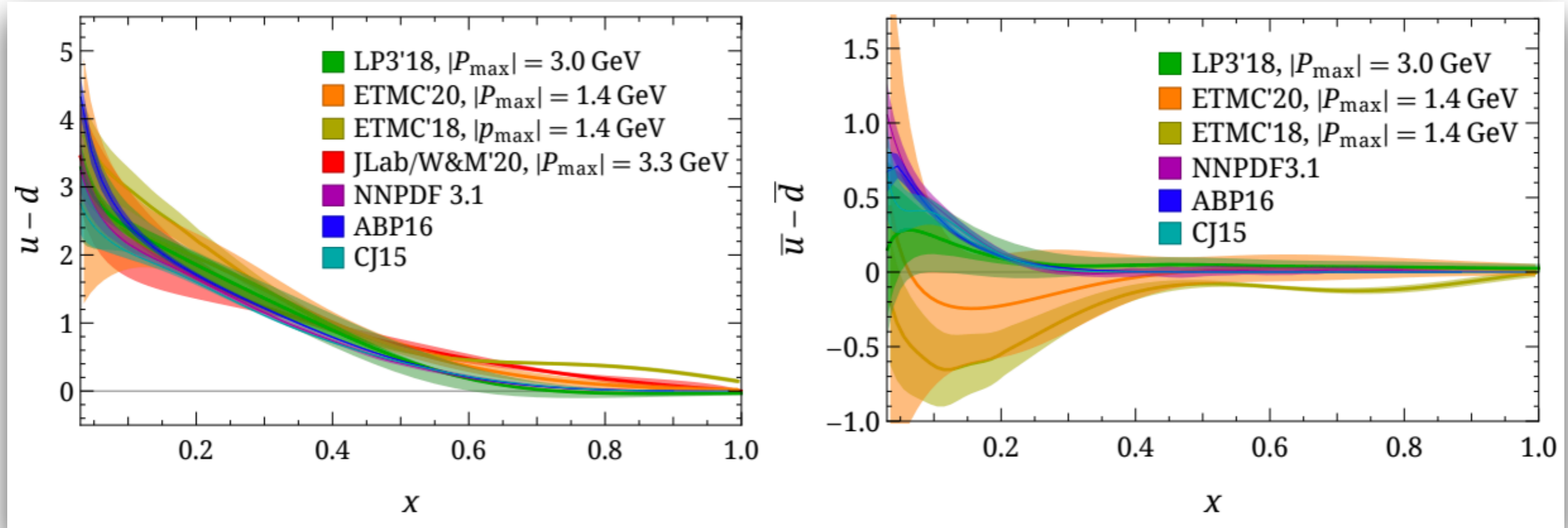
PDFs:

The unpolarized case

See also QCD@LHC 2024

- Results presented in $\overline{\text{MS}}$ scheme at 2 GeV
- Errors reported are statistical (in majority of cases)

Collection of results



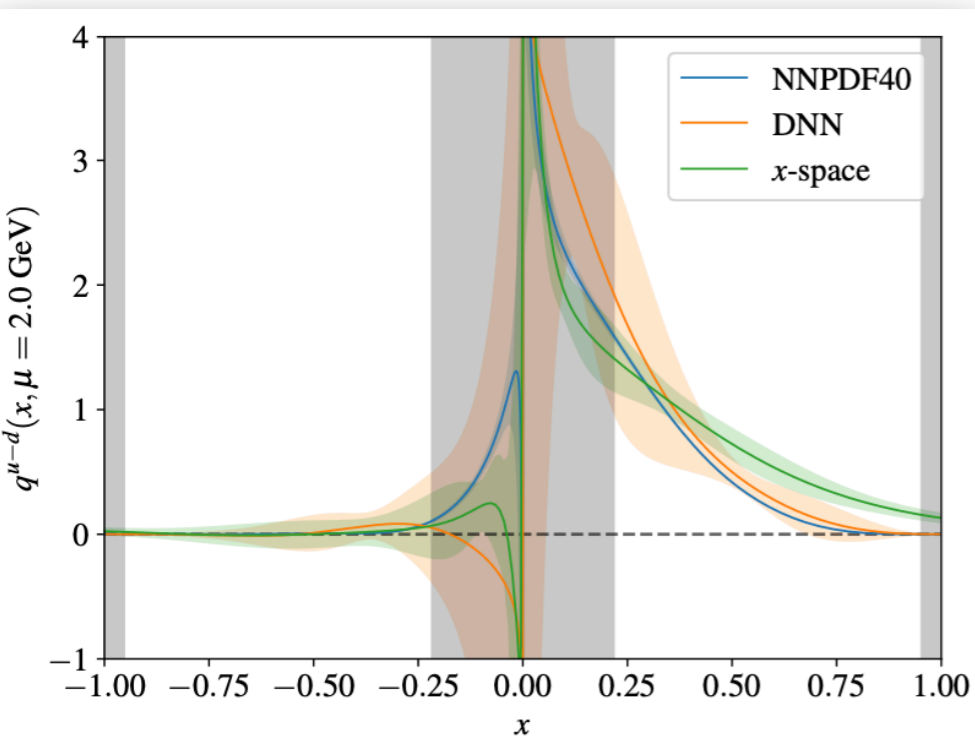
[M. Constantinou et al. (2020 PDFLattice Report), Prog.Part.Nucl.Phys. 121 (2021) 103908]

- ★ Several improvements:
 - More calculations at physical quark masses
 - Ensembles at various lattice spacings
 - Addressing systematic uncertainties due to methodologies

Refining the unpolarized proton PDF (u-d)

★ Physical quark masses

- HISQ, $a=0.076$ fm, $P\sim 1.5$ GeV
- Deep Neural Network for inverse problem
- NNLO for matching

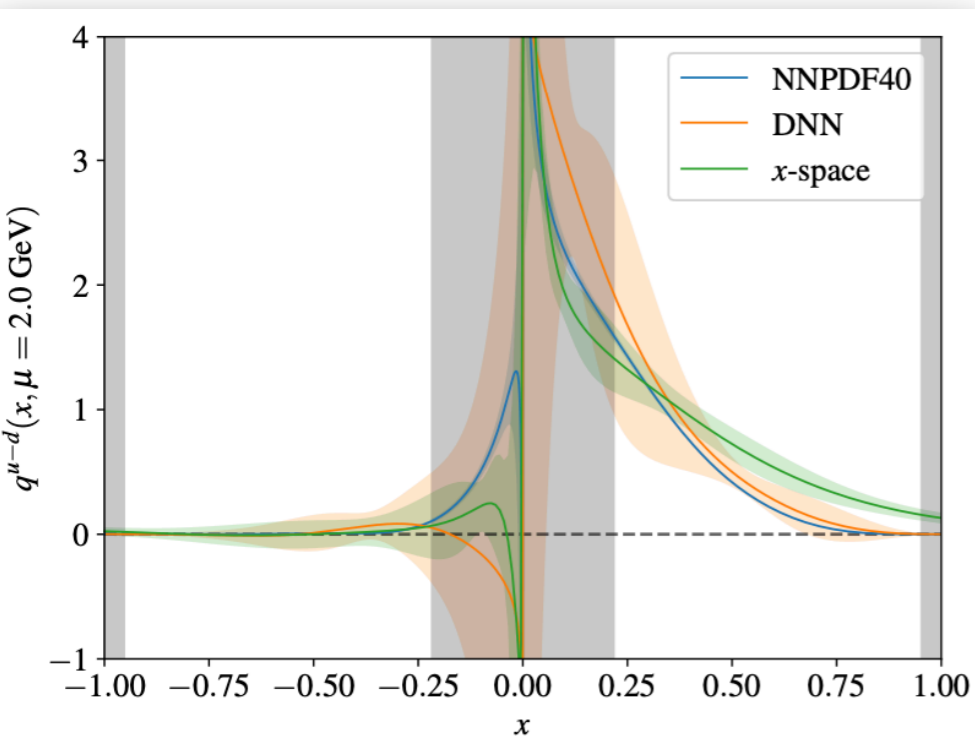


[X. Gao et al., PRD 107 (2023) 7, 074509]

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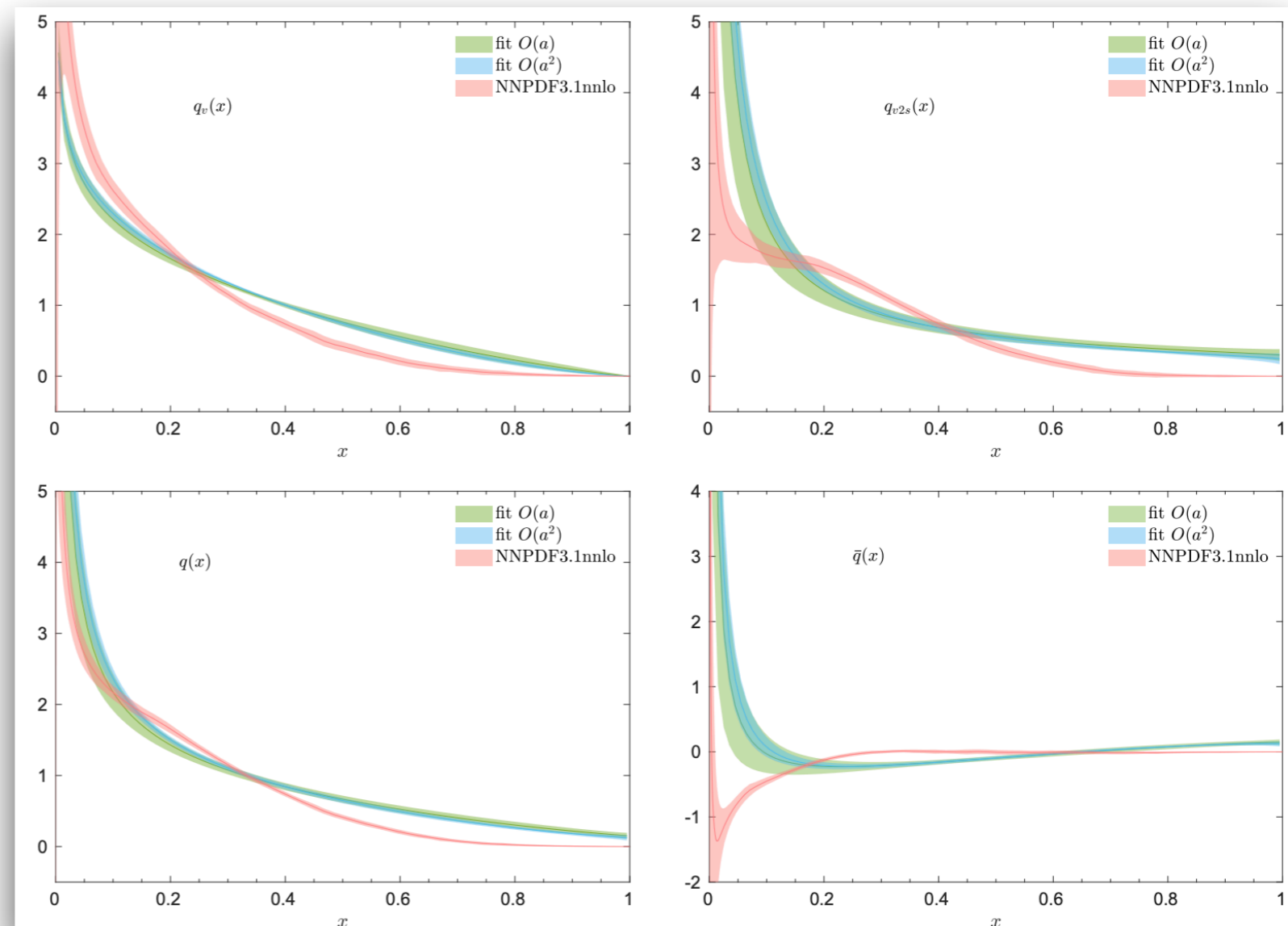
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★ Continuum limit

- TM&clover, $a=0.09$ fm, $m_\pi=350$ MeV
- $P\sim 1.8$ GeV
- NNLO for matching

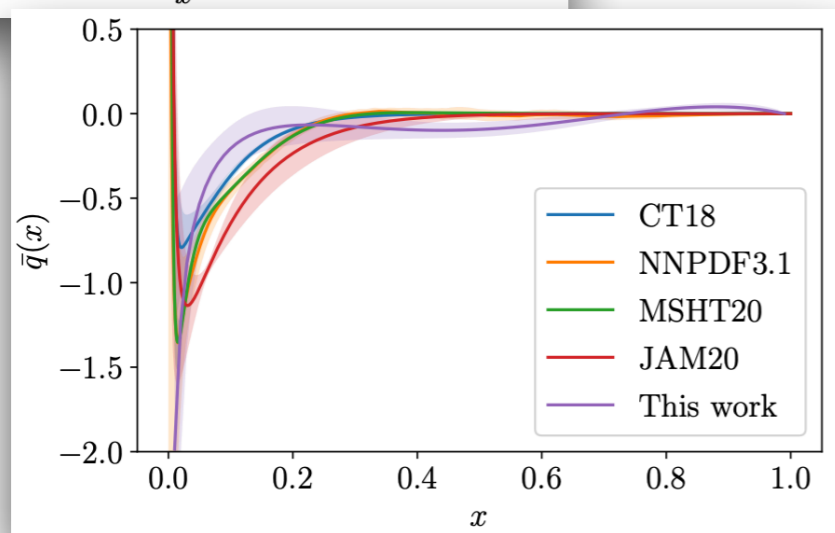
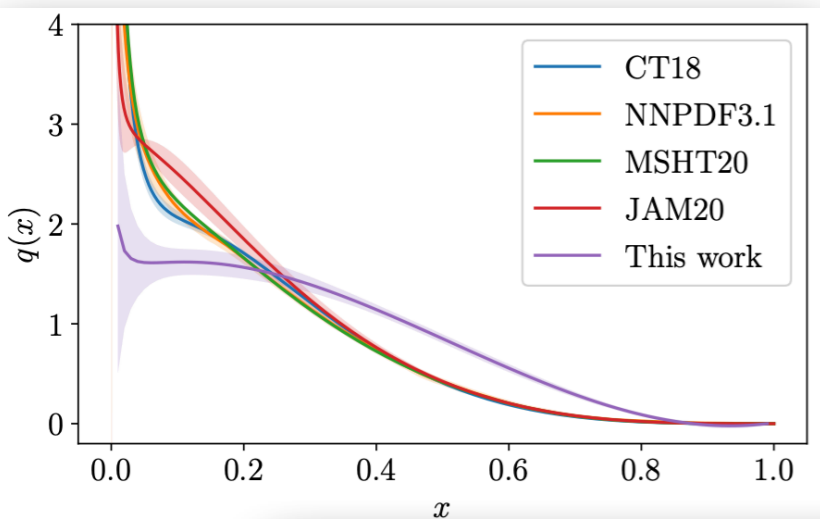


[Bhat et al., PRD 106 (2022) 5, 054504]

Improving evolution of PDFs

★ Continuum limit - higher twist effects

- Clover, $a=0.075, 0.065, 0.048$ fm
- $m_\pi=440$ MeV
- Jacobi polynomials for controlling finite- a & higher twist

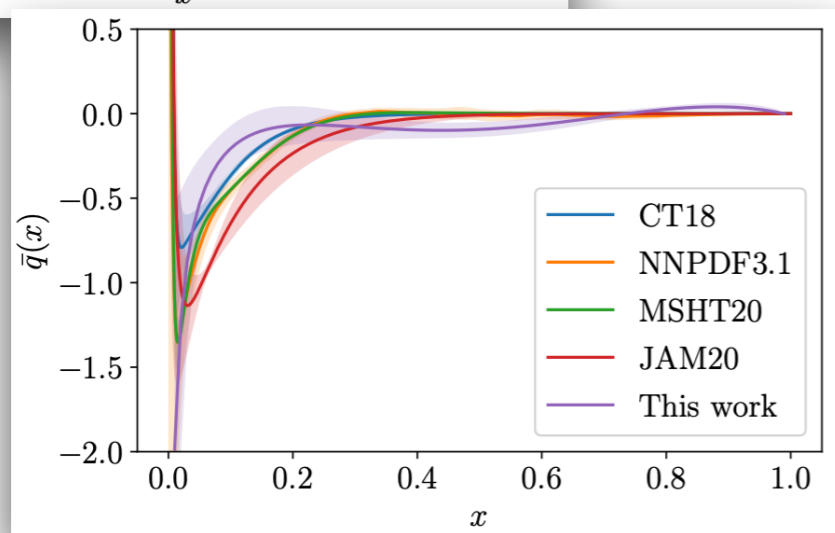
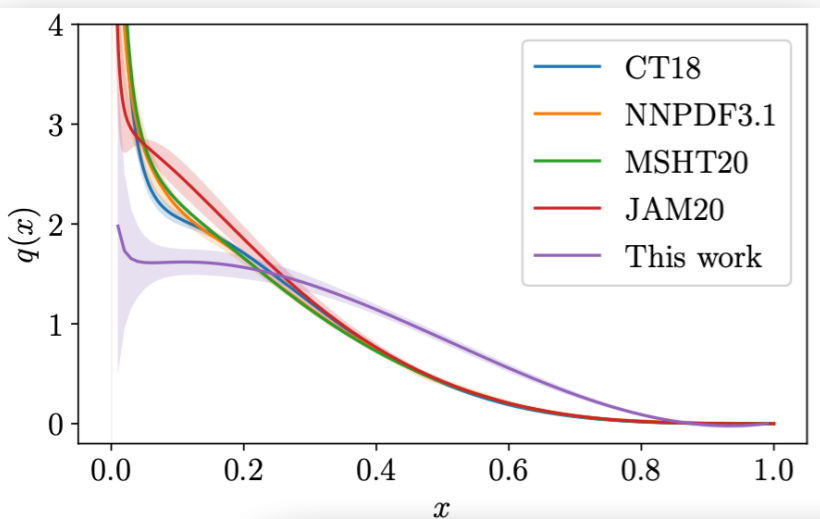


[Karpie et al., JHEP 11 (2021) 024]

Improving evolution of PDFs

★ Continuum limit - higher twist effects

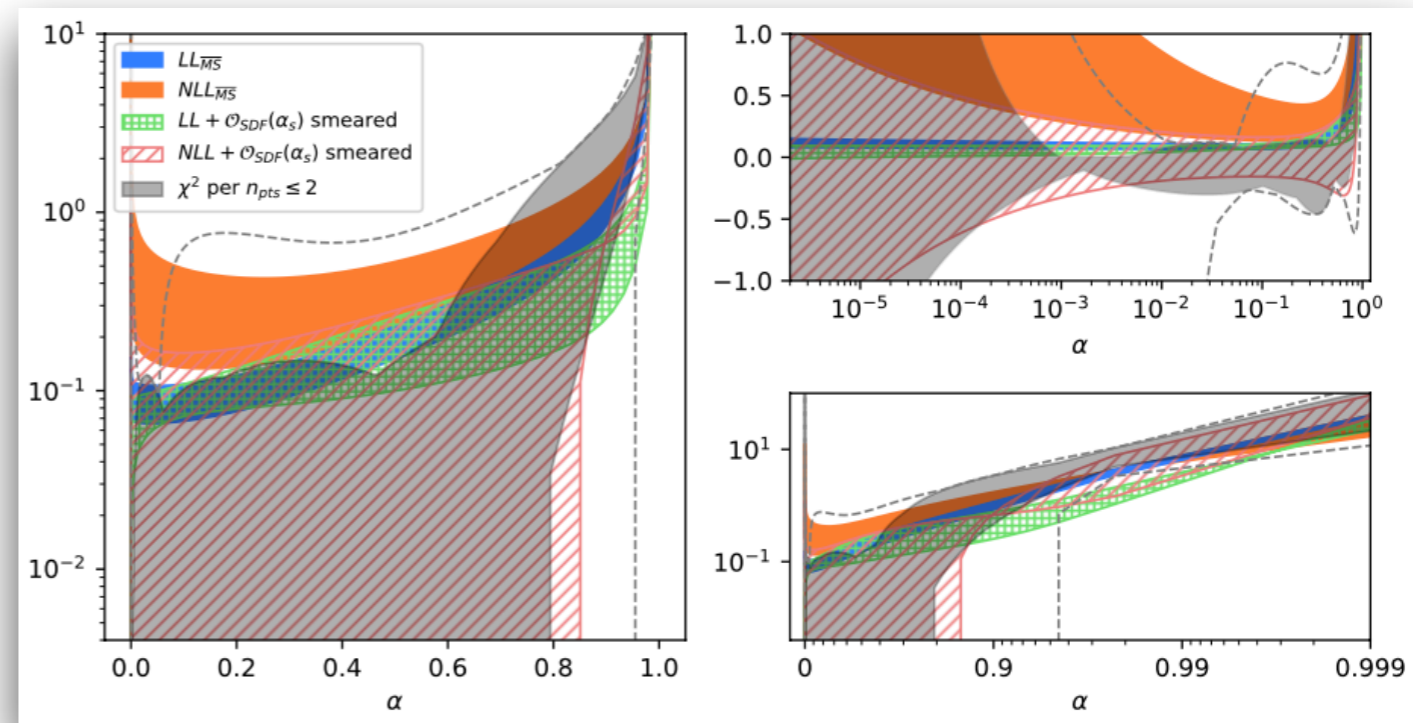
- Clover, $a=0.075, 0.065, 0.048$ fm
- $m_\pi=440$ MeV
- Jacobi polynomials for controlling finite- a & higher twist



[Karpie et al., JHEP 11 (2021) 024]

★ Non-perturbative scale evolution of pseudo distributions:

- lattice scale much different than scale for light-cone PDFs
- addresses the subtle z^2 behavior of matrix elements
- reduces fluctuation of lattice data

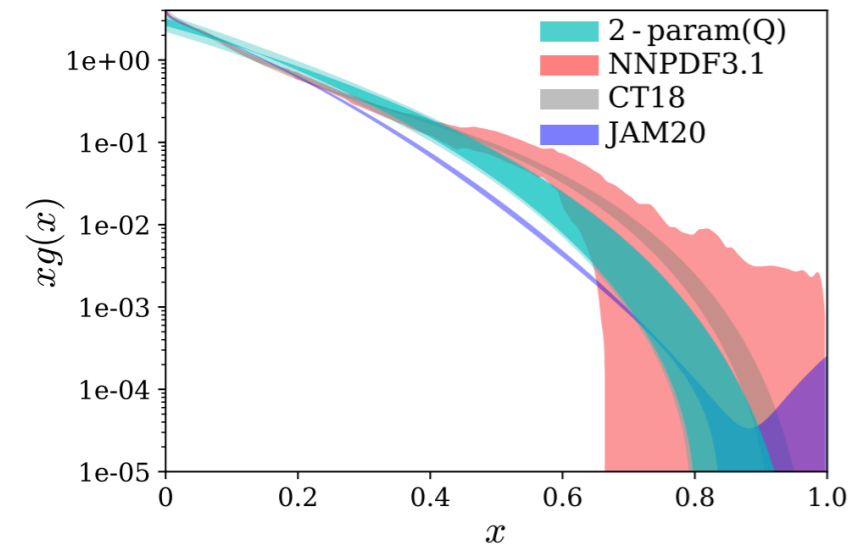


[H. Dutrieux et al. (HadStruc), JHEP 04 (2024) 061]

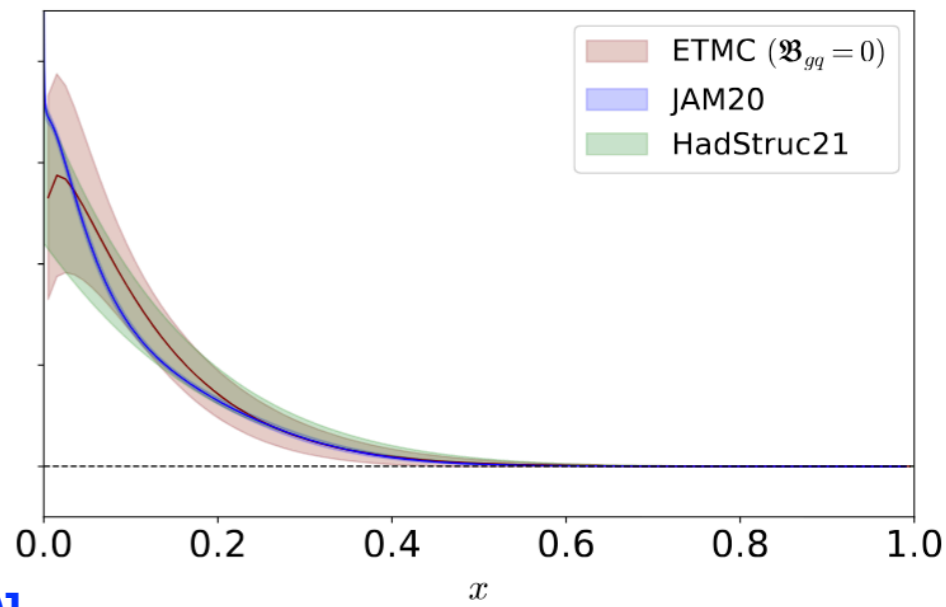
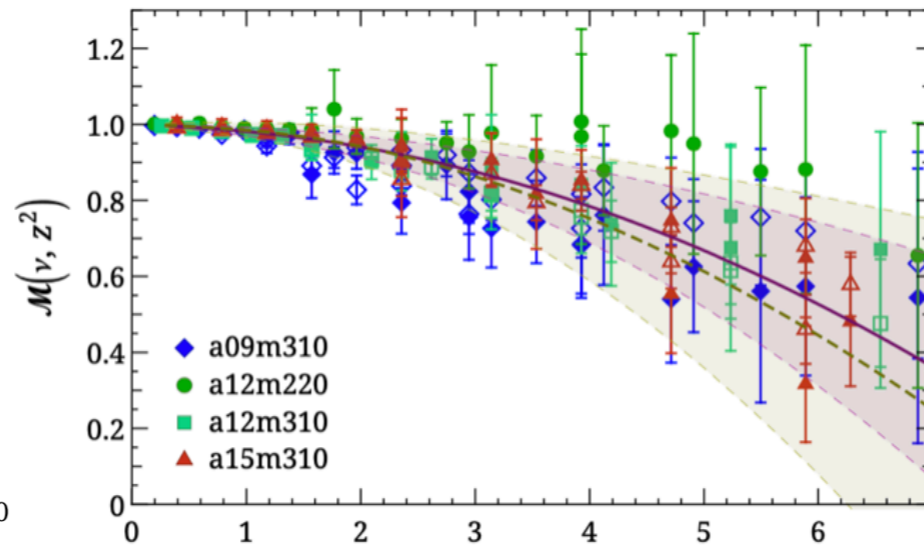
Evolution of vector operator much larger than anticipated

Gluon PDF

[Khan et al., JHEP 11, 148 (2021)]



[Delmar et al., PRD 108 (2023) 9, 094515]

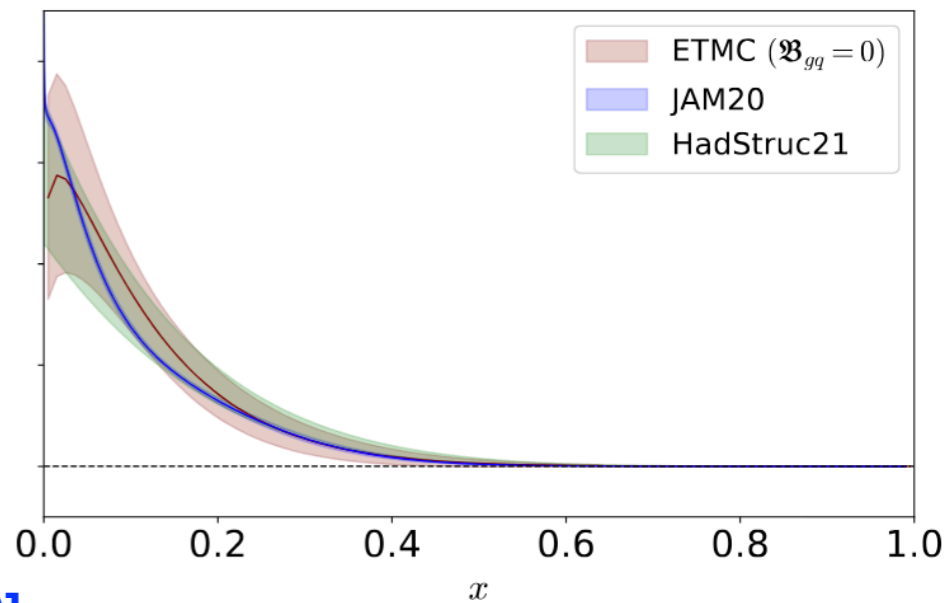
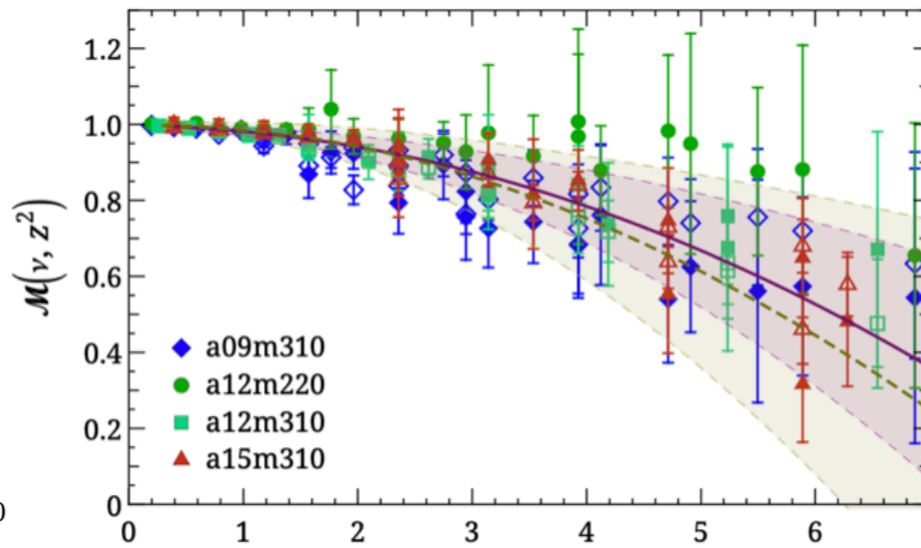
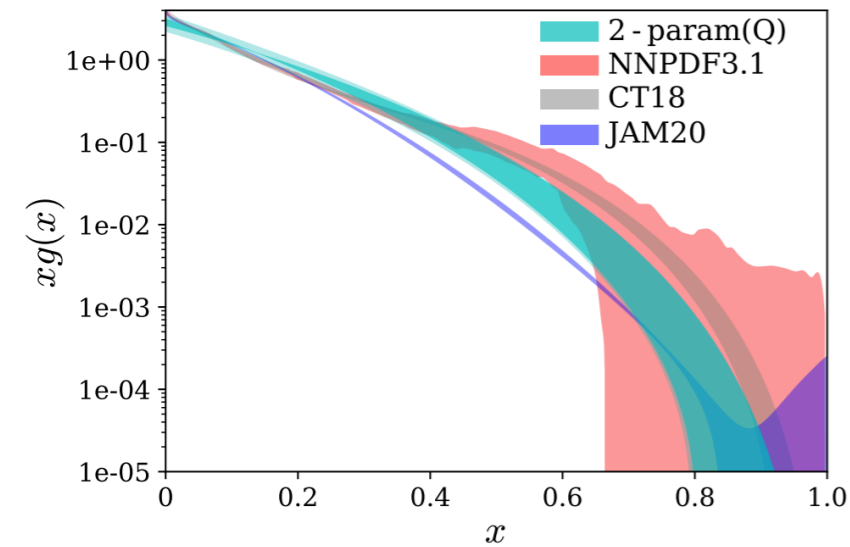


[Fan et al., Phys. Rev. D 108, 014508 (2023)]

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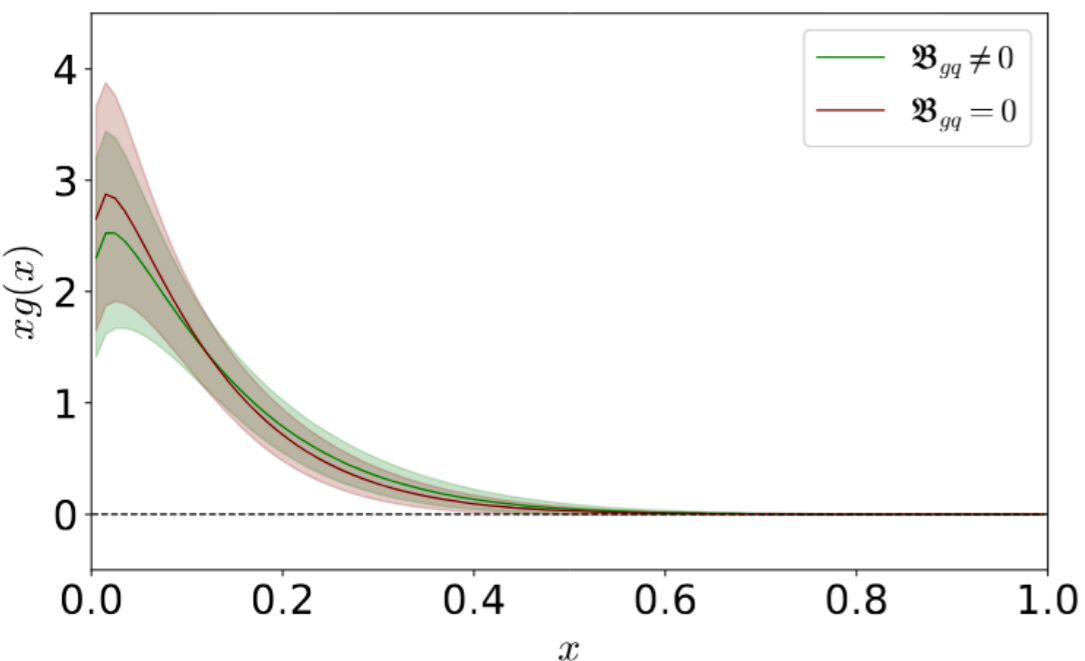
[Delmar et al., PRD 108 (2023) 9, 094515]



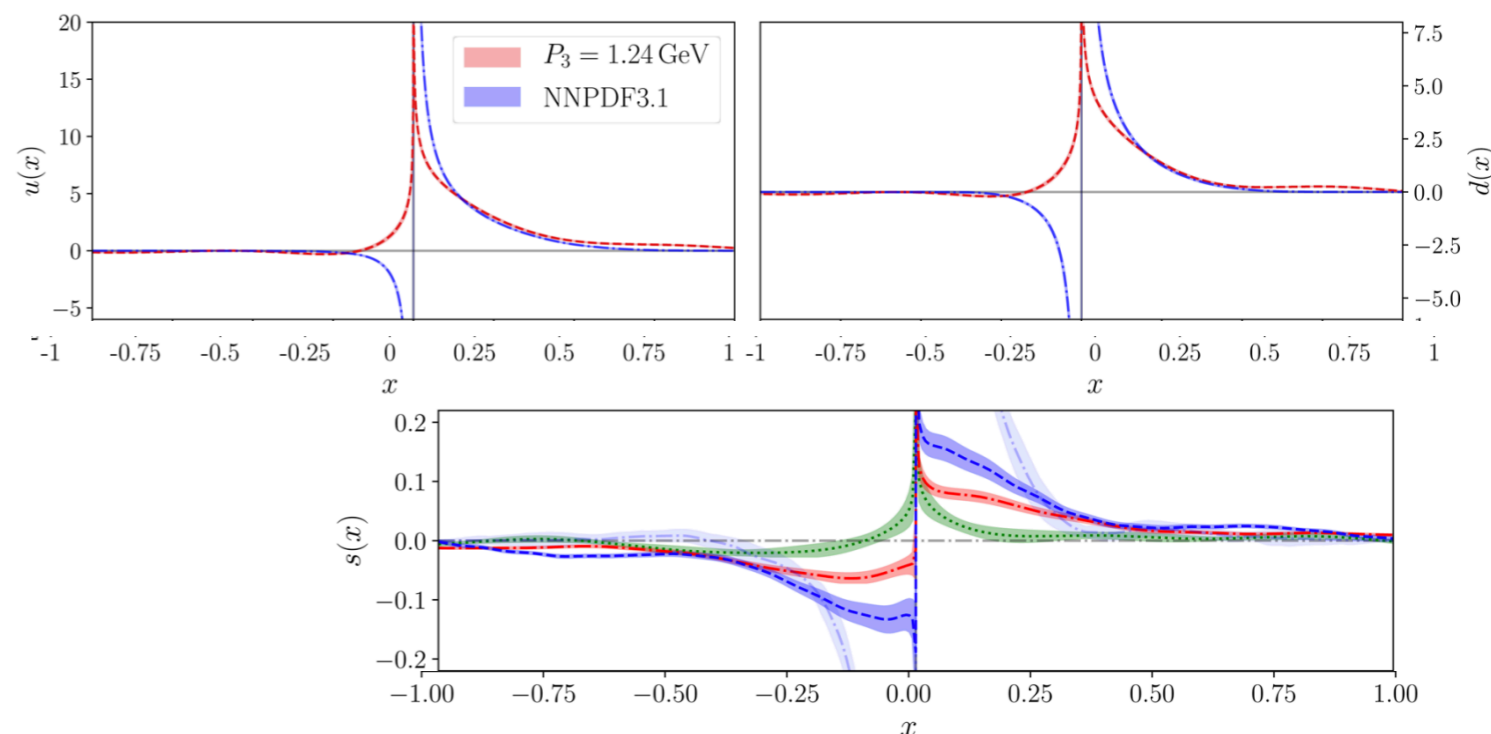
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★ Elimination of Mixing with Quark Singlet PDFs

[J. Delmar et al., PRD 108 (2023) 9, 094515]



Dedicated calculation of quark-singlet PDFs

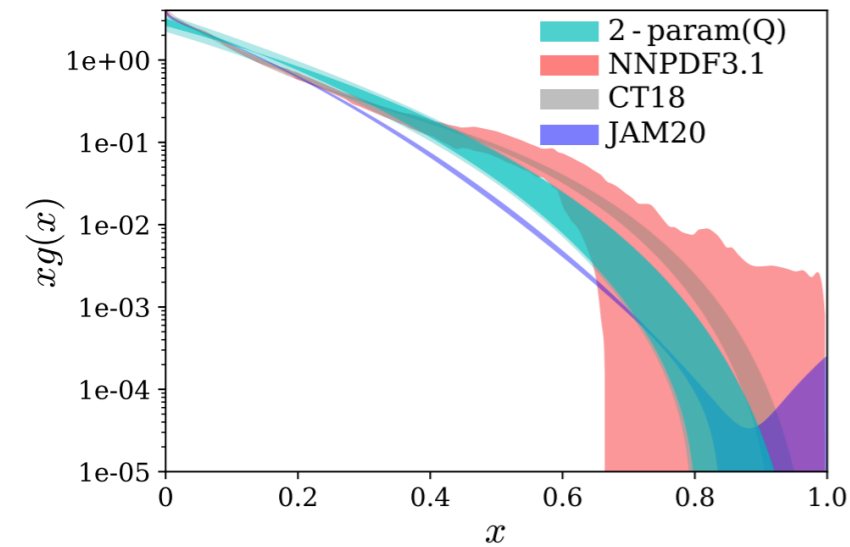


Lattice QCD can provide key information

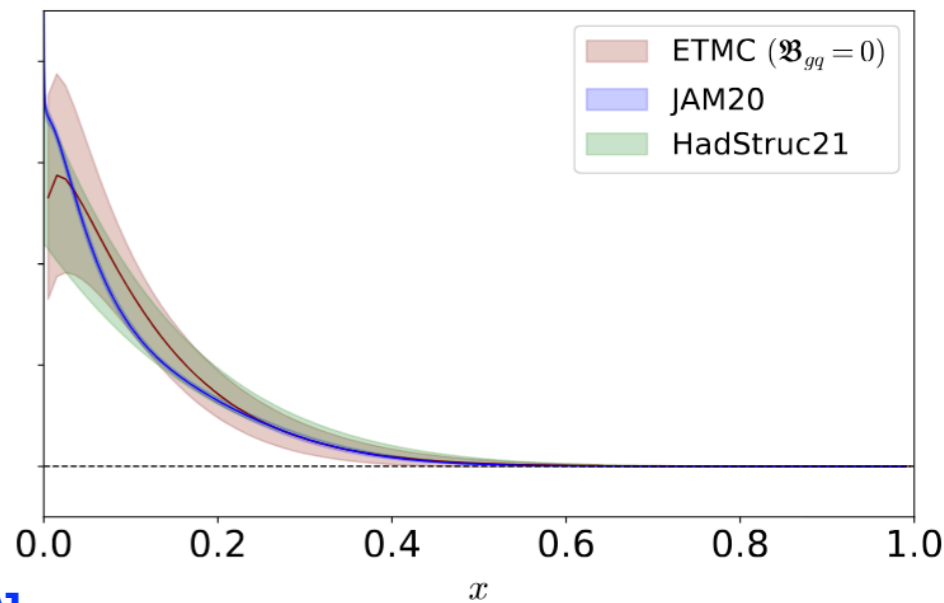
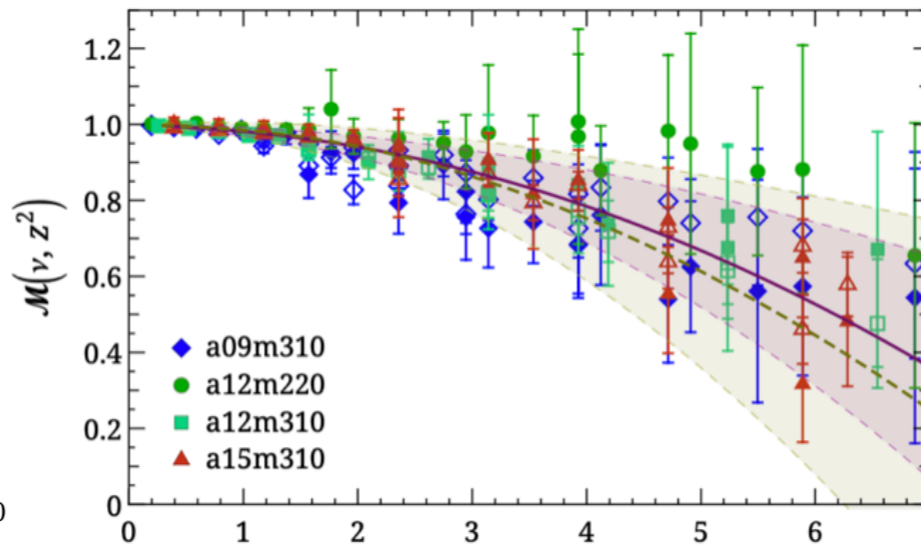
[C. Alexandrou, PRD 104 (2021) 5, 054503]

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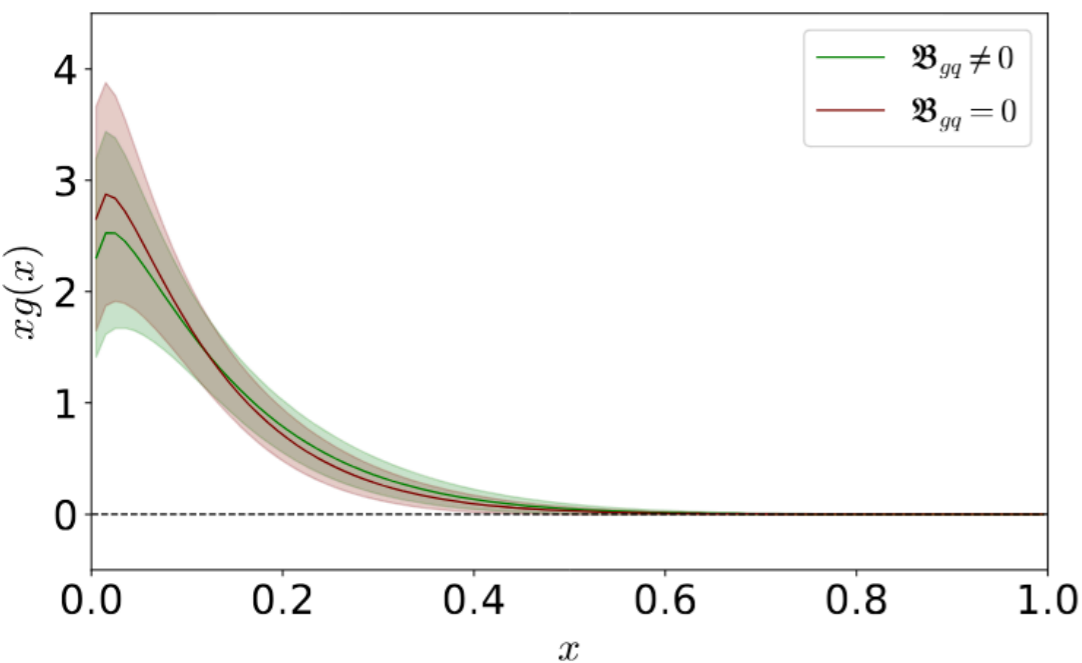
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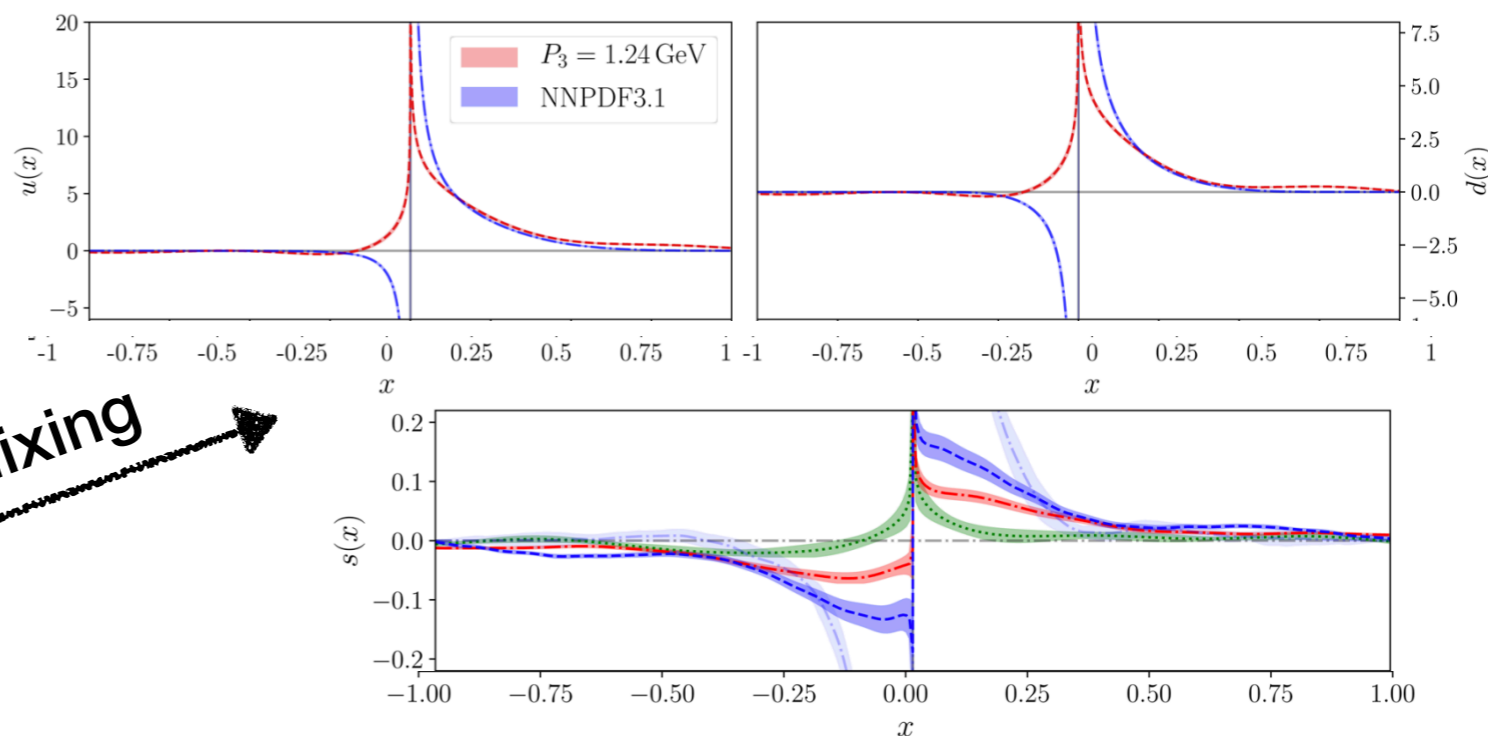
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GPDs

leading twist

GPDs on the lattice

- ★ Off-forward matrix elements of non-local light-cone operators

$$F^{[\gamma^+]}(x, \Delta; \lambda, \lambda') = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ik \cdot z} \langle p'; \lambda' | \bar{\psi}(-\frac{z}{2}) \gamma^+ \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p; \lambda \rangle \Big|_{z^+=0, \vec{z}_\perp=\vec{0}_\perp}$$

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[Constantinou & Panagopoulos (2017)]

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γ^0 ideal for PDFs

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γ^0 ideal for PDFs

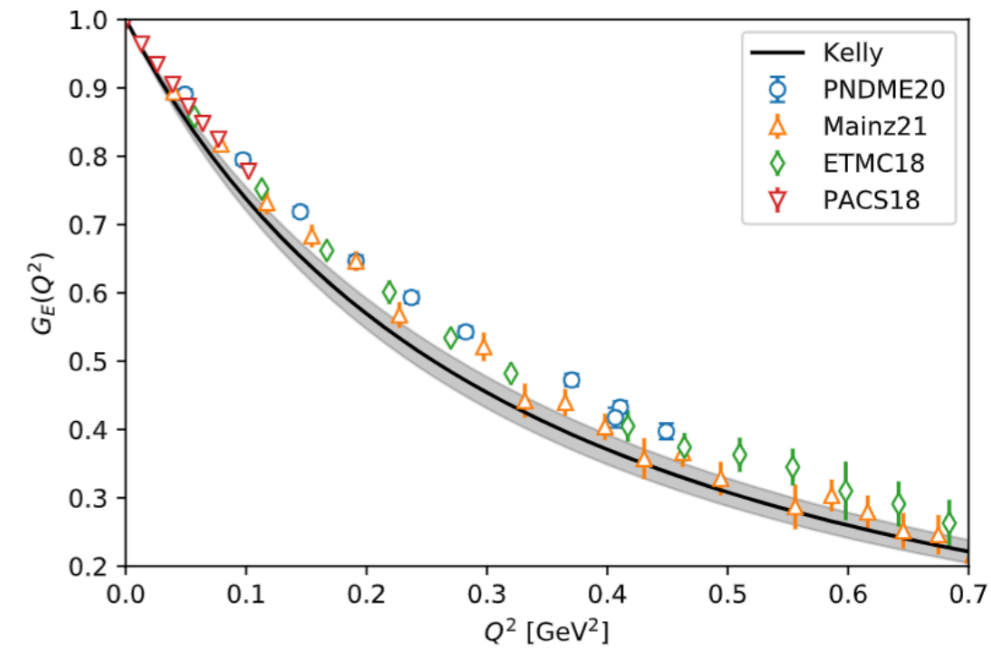
γ^0 parametrization is prohibitively expensive

Definition of GPDs in Euclidean lattice

- ★ Calculation expected to be performed in symmetric frame to extract “standard” GPDs
- ★ Symmetric frame requires separate calculations at each t

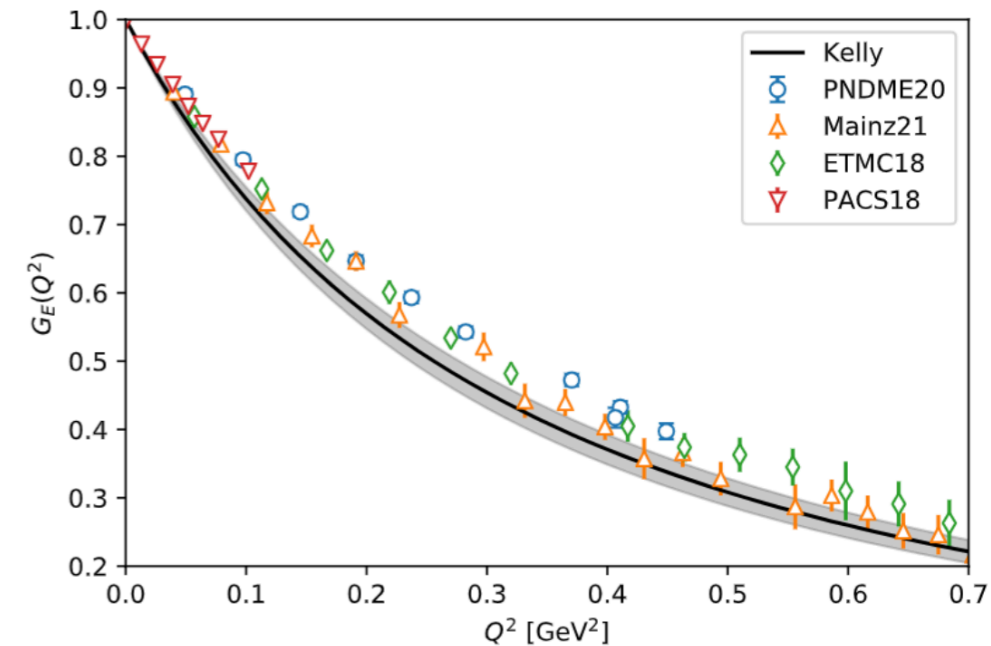
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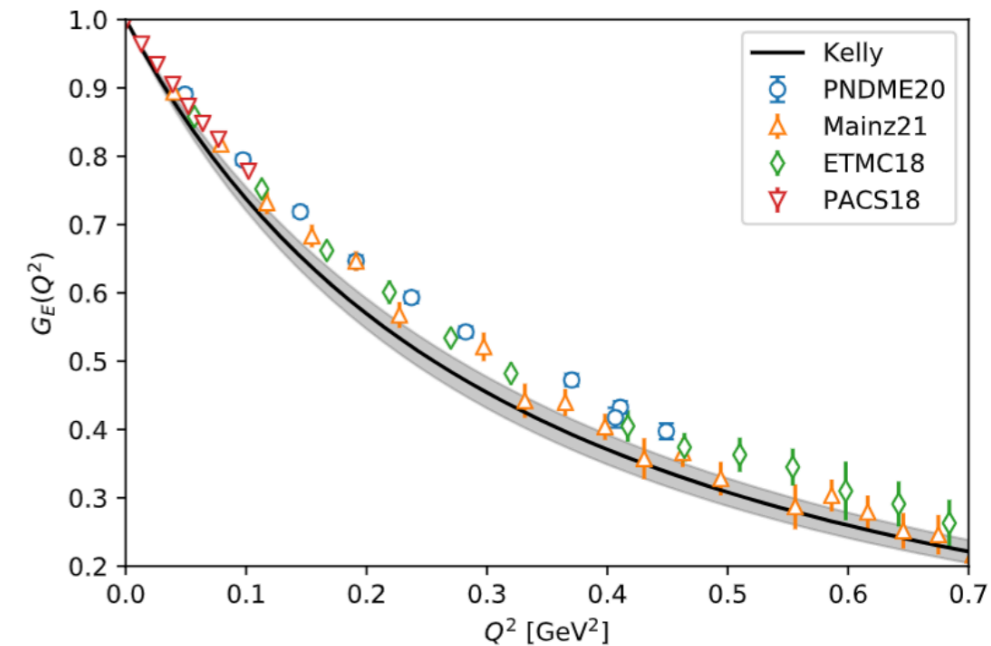
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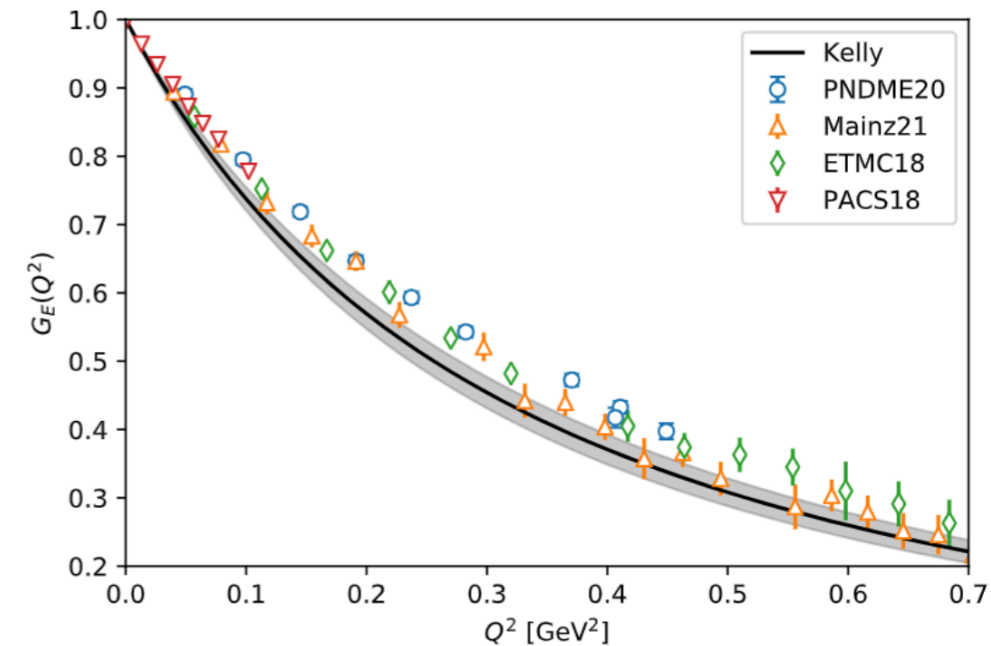
[S. Bhattacharya et al., PRD 106 (2022) 11, 114512]

Advantages

- Applicable to any kinematic frame and have definite symmetries
- Lorentz invariant amplitudes A_i can be related to the standard H, E GPDs
- Quasi H, E may be redefined (Lorentz covariant) to eliminate $1/P_3$ contributions:

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→ Proof-of-concept calculation ($\xi = 0$):

- symmetric frame: $\vec{p}_f^s = \vec{P} + \frac{\vec{Q}}{2}, \quad \vec{p}_i^s = \vec{P} - \frac{\vec{Q}}{2} \quad -t^s = \vec{Q}^2 = 0.69 \text{ GeV}^2$

- asymmetric frame: $\vec{p}_f^a = \vec{P}, \quad \vec{p}_i^a = \vec{P} - \vec{Q} \quad t^a = -\vec{Q}^2 + (E_f - E_i)^2 = 0.65 \text{ GeV}^2$

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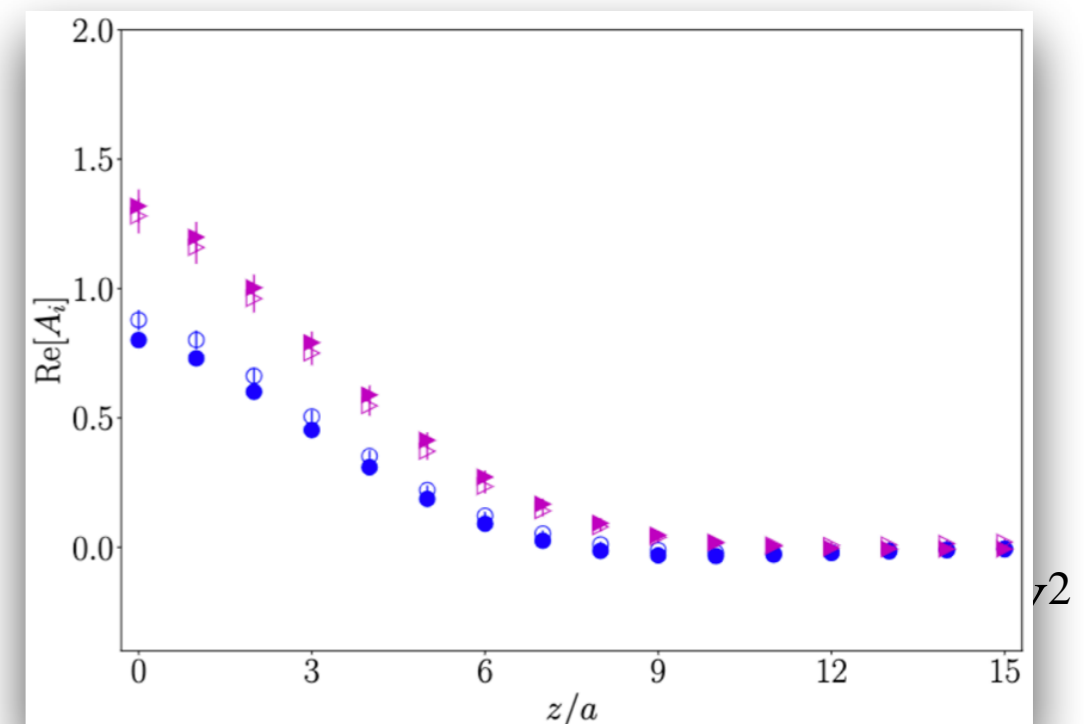
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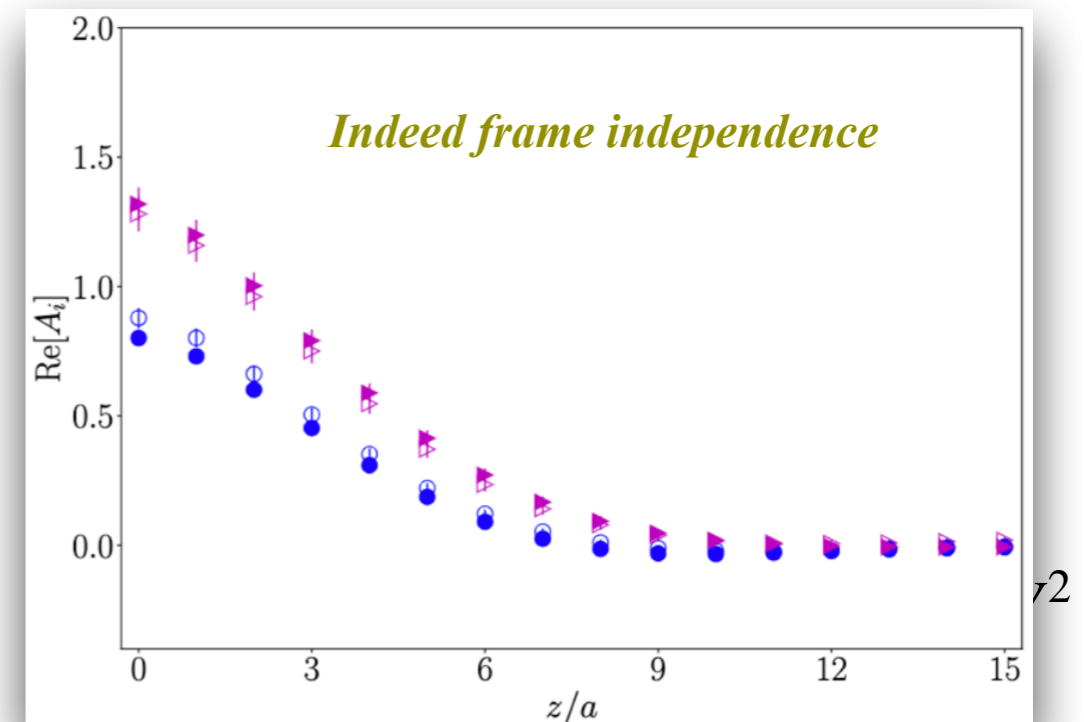
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- asymmetric frame: $\vec{p}_f^a = \vec{P}, \quad \vec{p}_i^a = \vec{P} - \vec{Q}$



Light-cone GPDs

- ★ Nf=2+1+1 twisted mass fermions with a clover term

Name	β	N_f	$L^3 \times T$	a [fm]	M_π	$m_\pi L$
cA211.32	1.726	u, d, s, c	$32^3 \times 64$	0.093	260 MeV	4

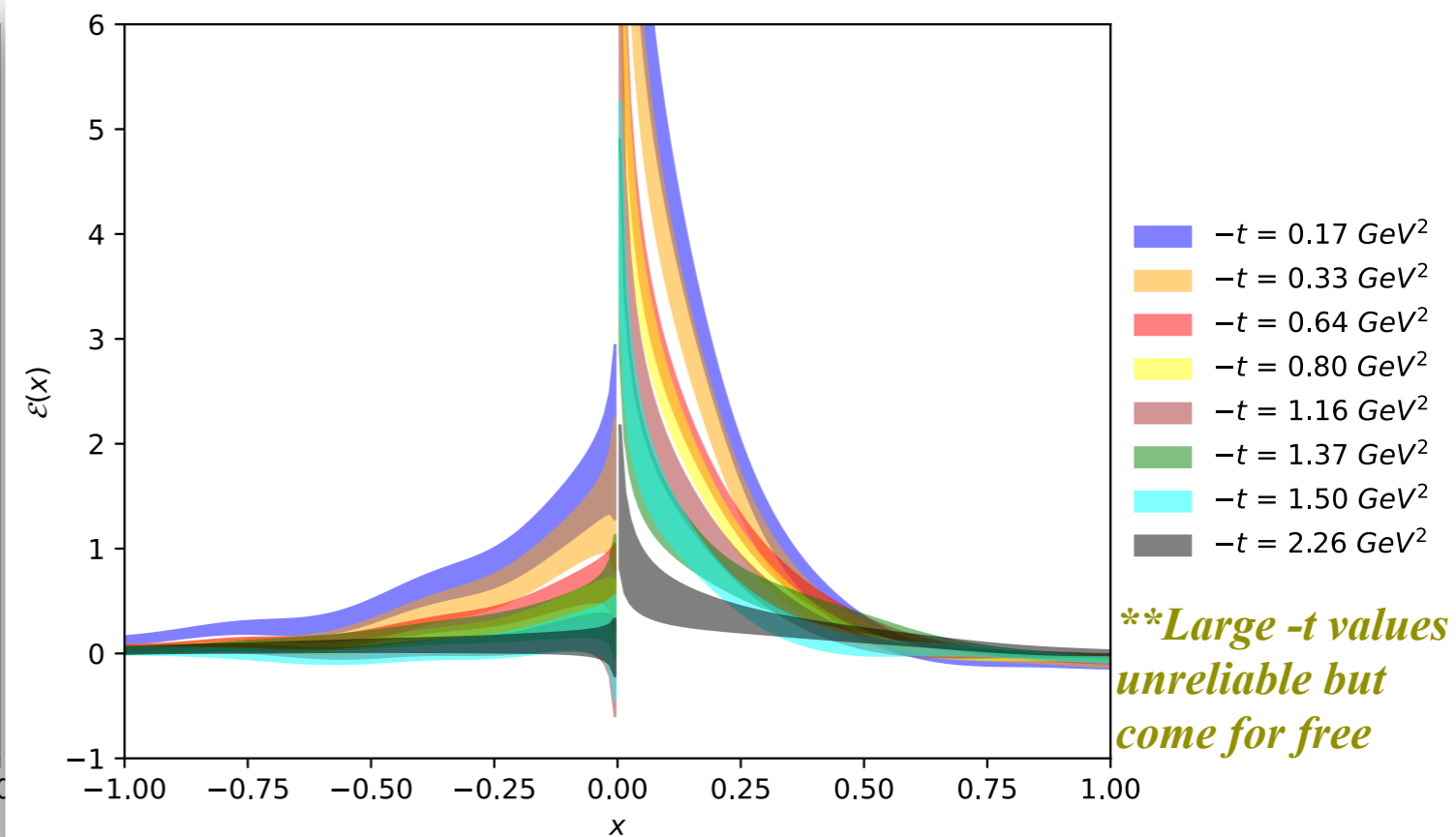
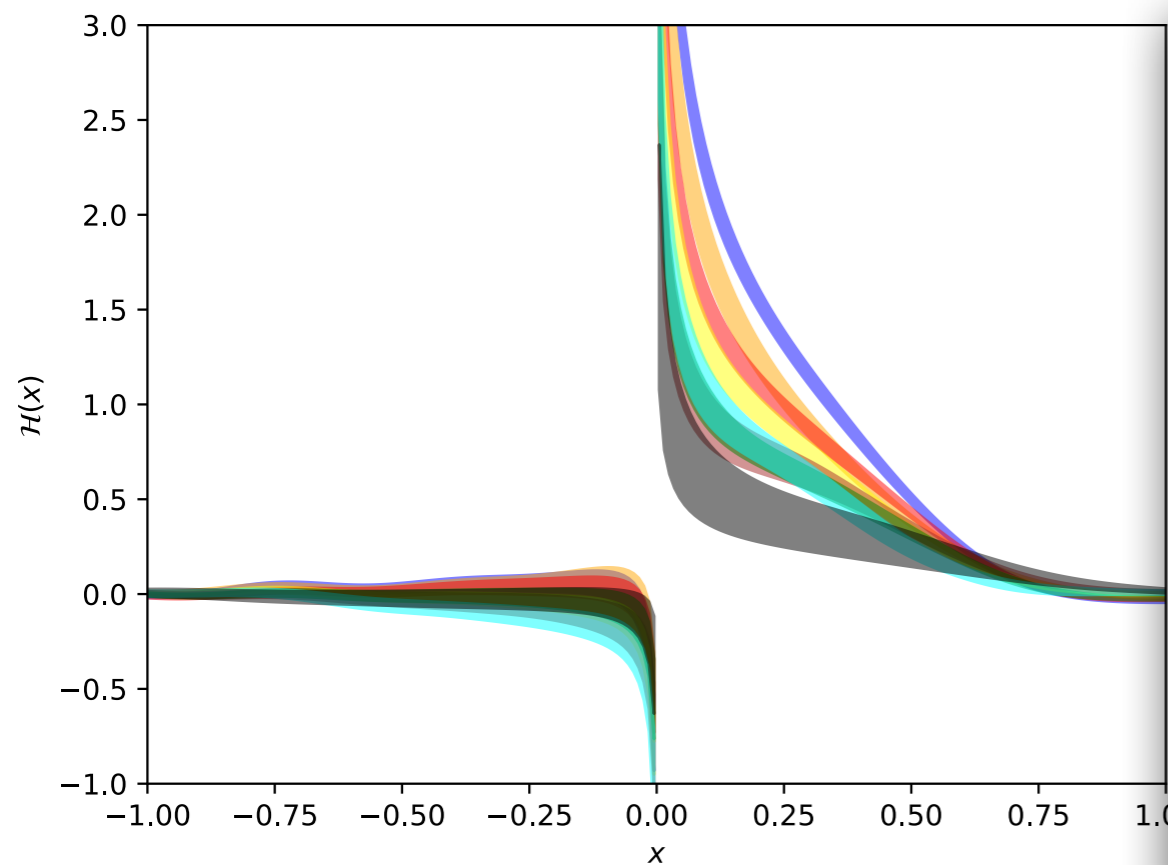
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- ★ $N_f=2+1+1$ twisted mass fermions with a clover term

H – GPD

E – GPD



- ★ $\pm x$ correspond to quark and anti-quark region
- ★ anti-quark region susceptible to systematic uncertainties
- ★ small- x region not reliably extracted
- ★ perturbative matching breaks down at $\xi = x$

Helicity GPDs

$$F^{[\gamma^3 \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^0} \bar{u}(p_f, \lambda') \left[\gamma^3 \gamma_5 \tilde{\mathcal{H}}_3(x, \xi, t; P^3) + \frac{\Delta^3 \gamma_5}{2m} \tilde{\mathcal{E}}_3(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

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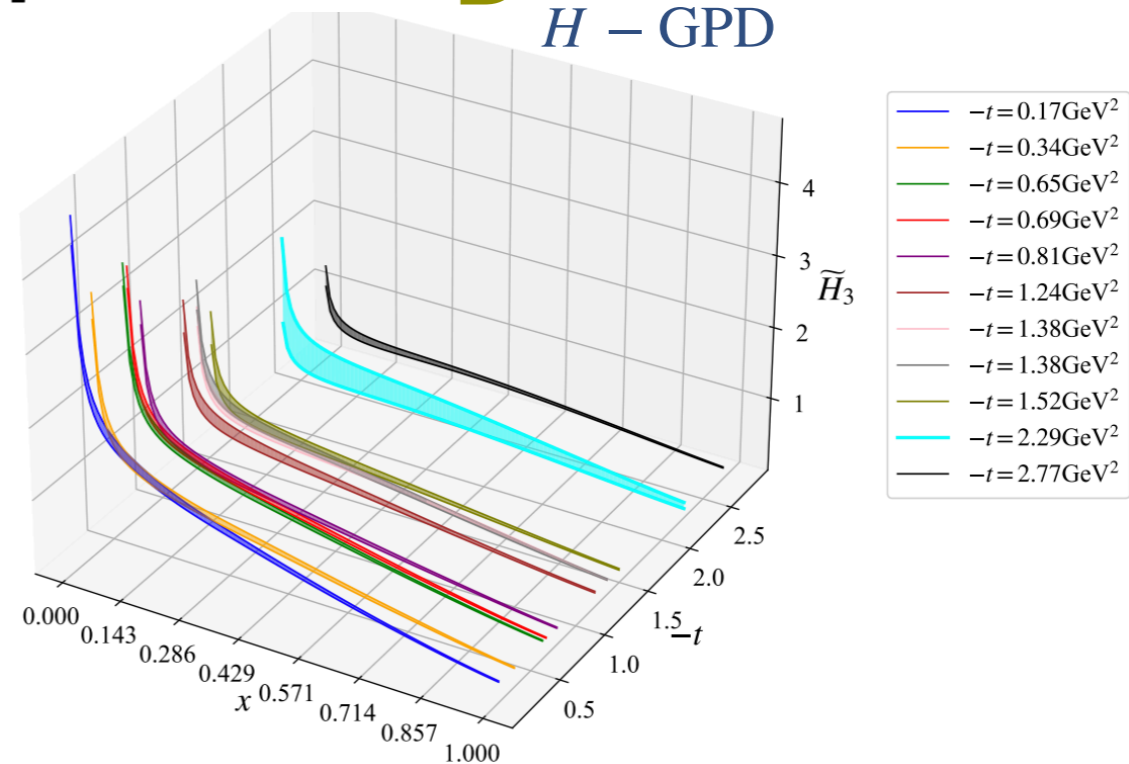
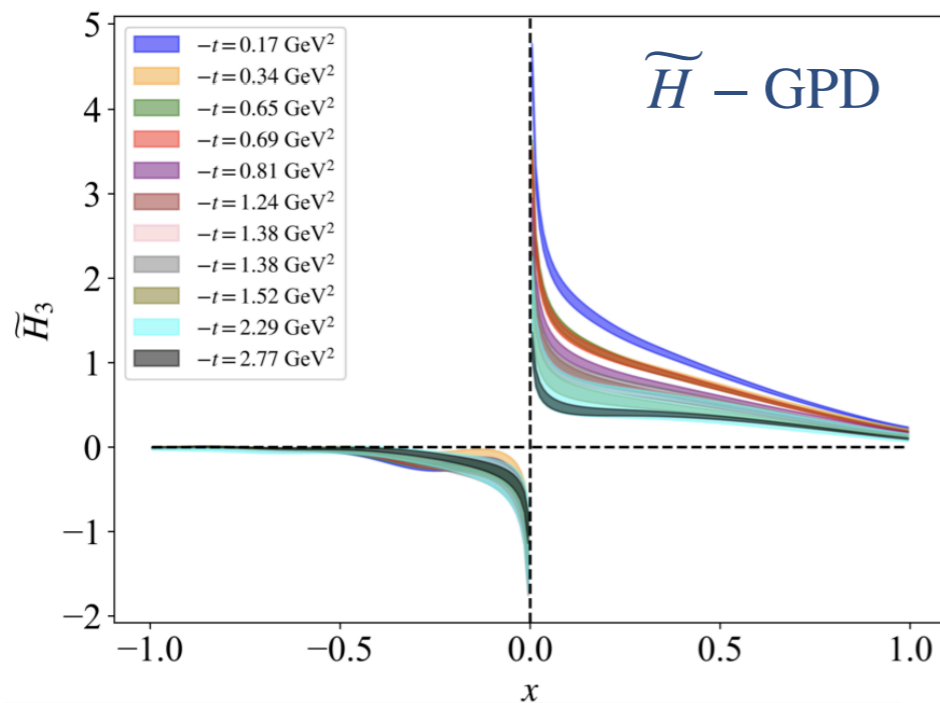
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\tilde{H} - GPD



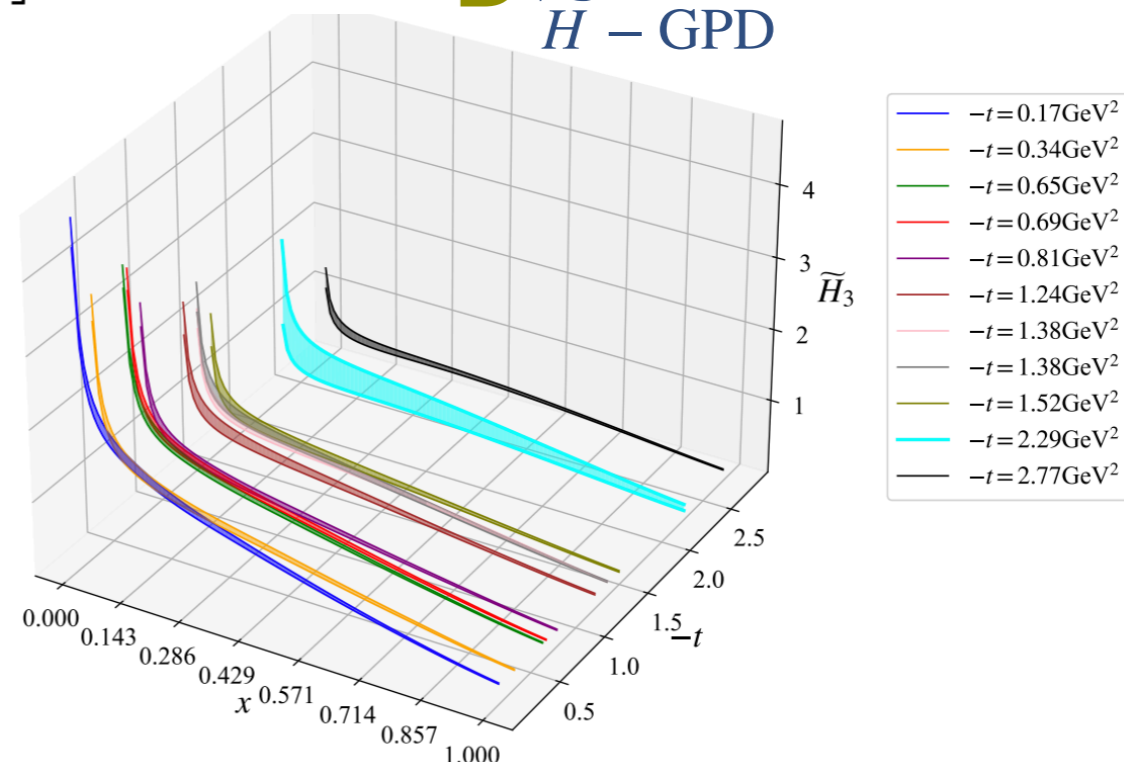
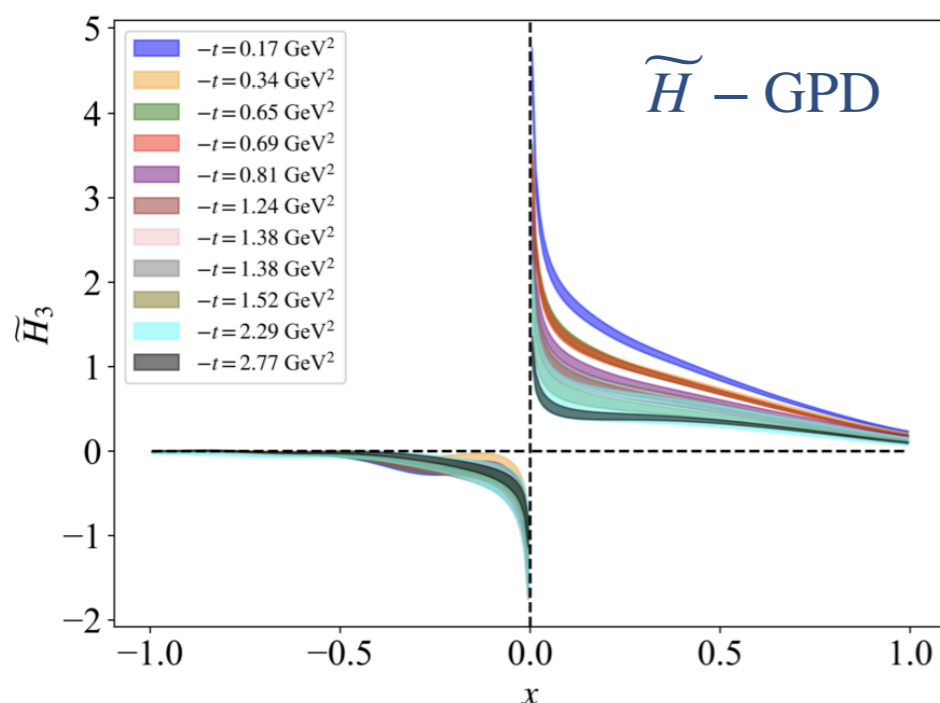
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★ \tilde{E} -GPD cannot be extracted directly at $\xi = 0$

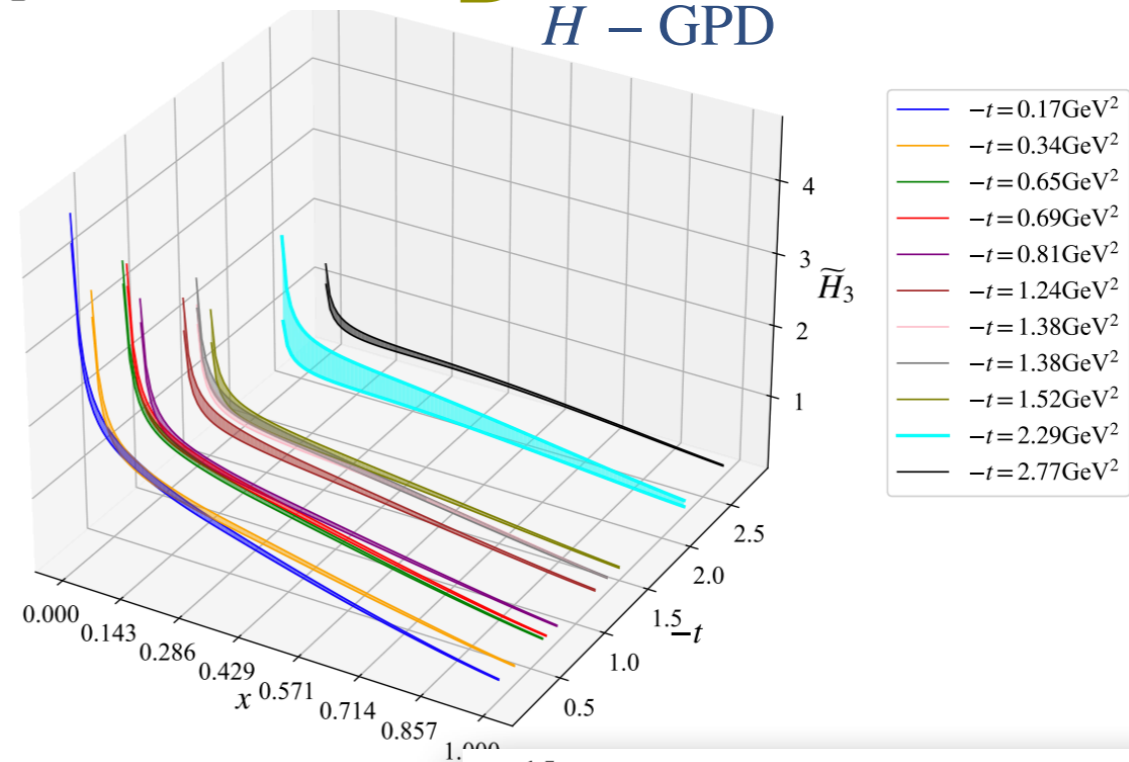
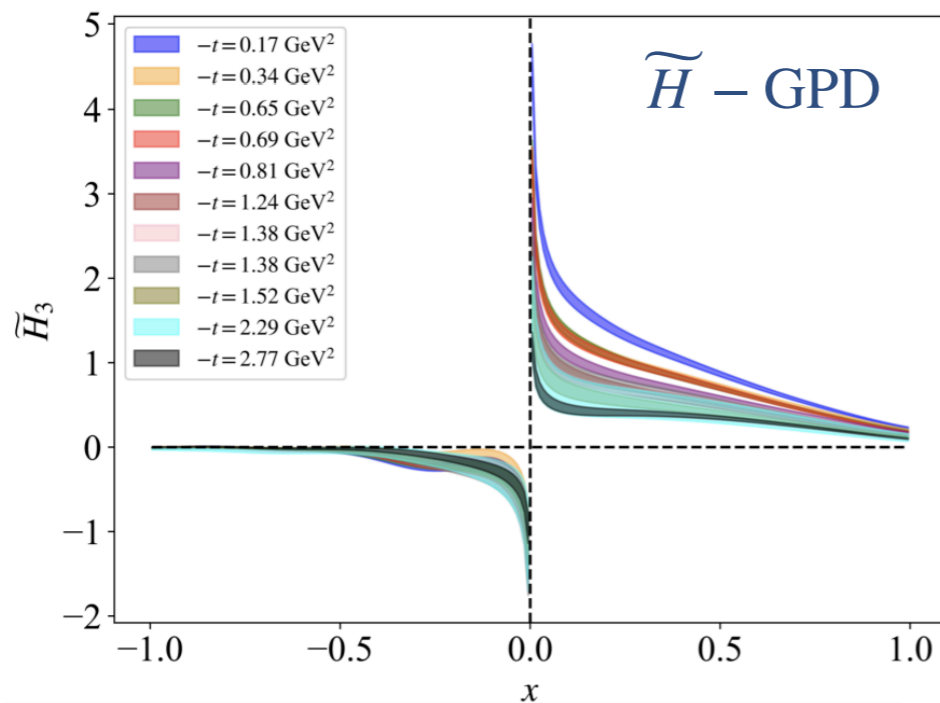
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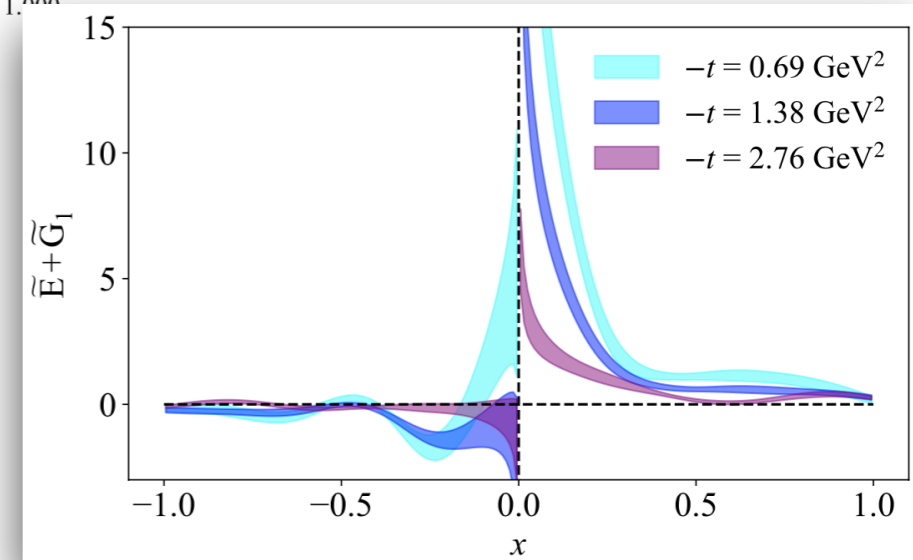
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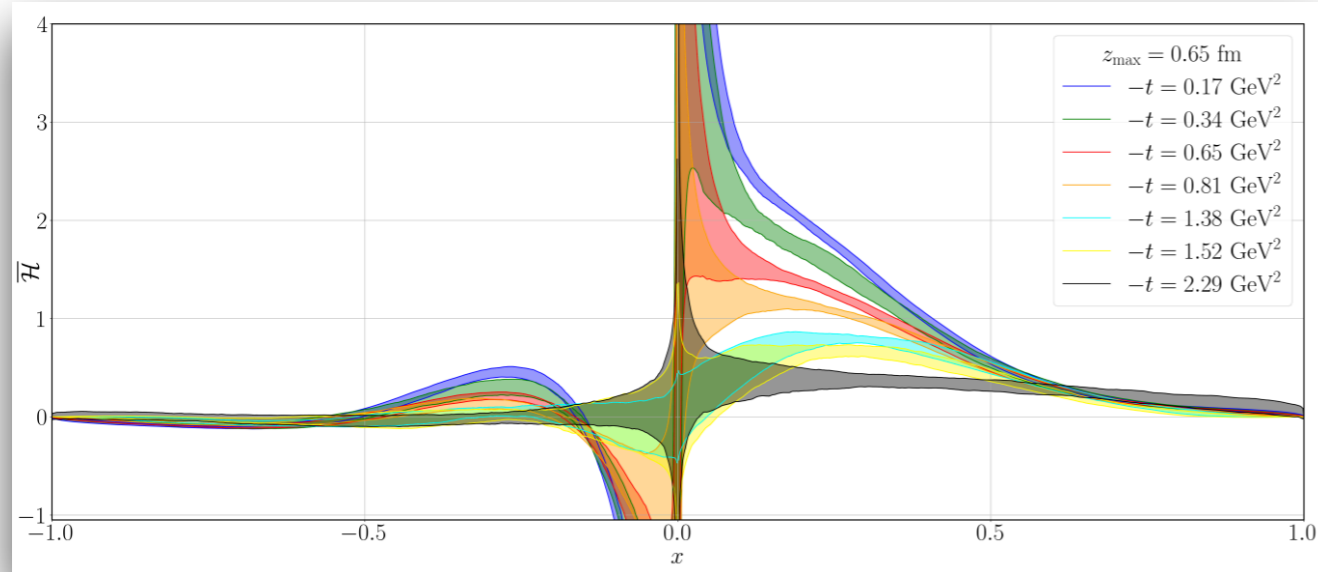
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Glimpse of \tilde{E} from twist-3 GPDs



Alternative approach: pseudo-ITD



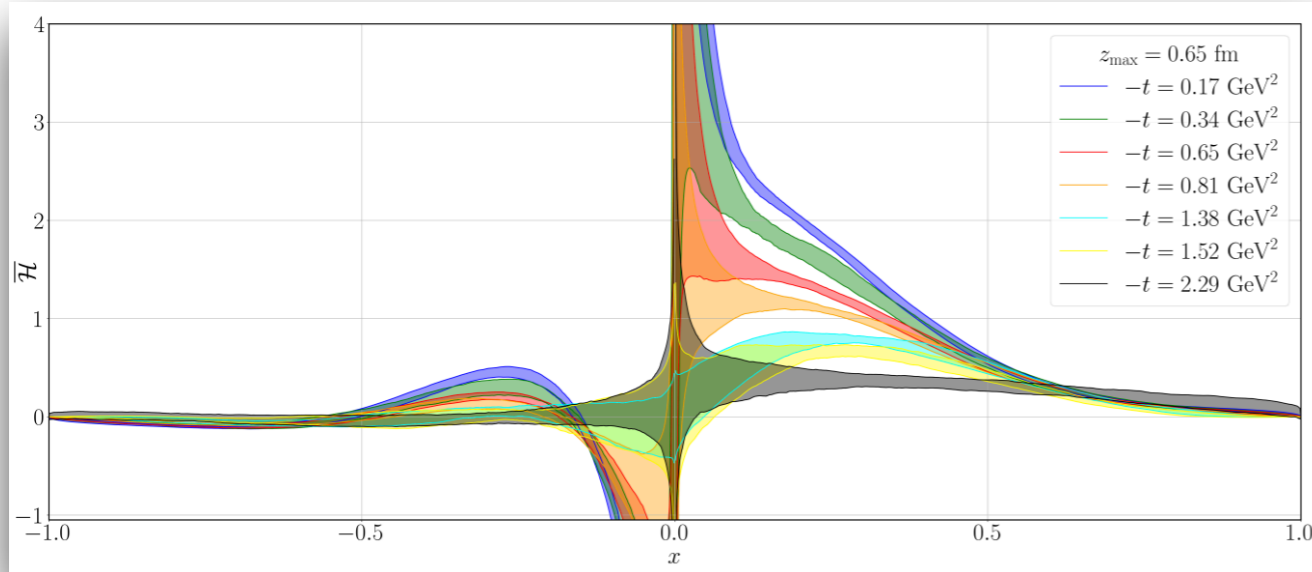
[Battacharya et al., PRD 110 (2024) 5, 054502]

Different steps between approaches:

- *renormalization*
- *x-dependence reconstruction*
- *matching formalism*

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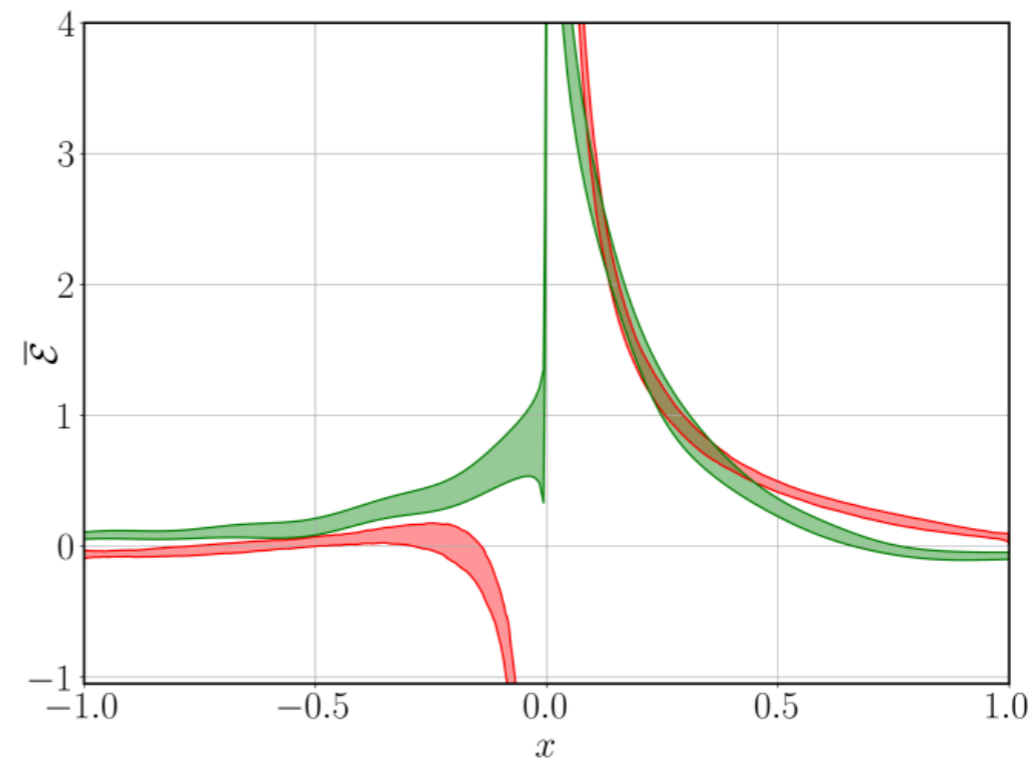
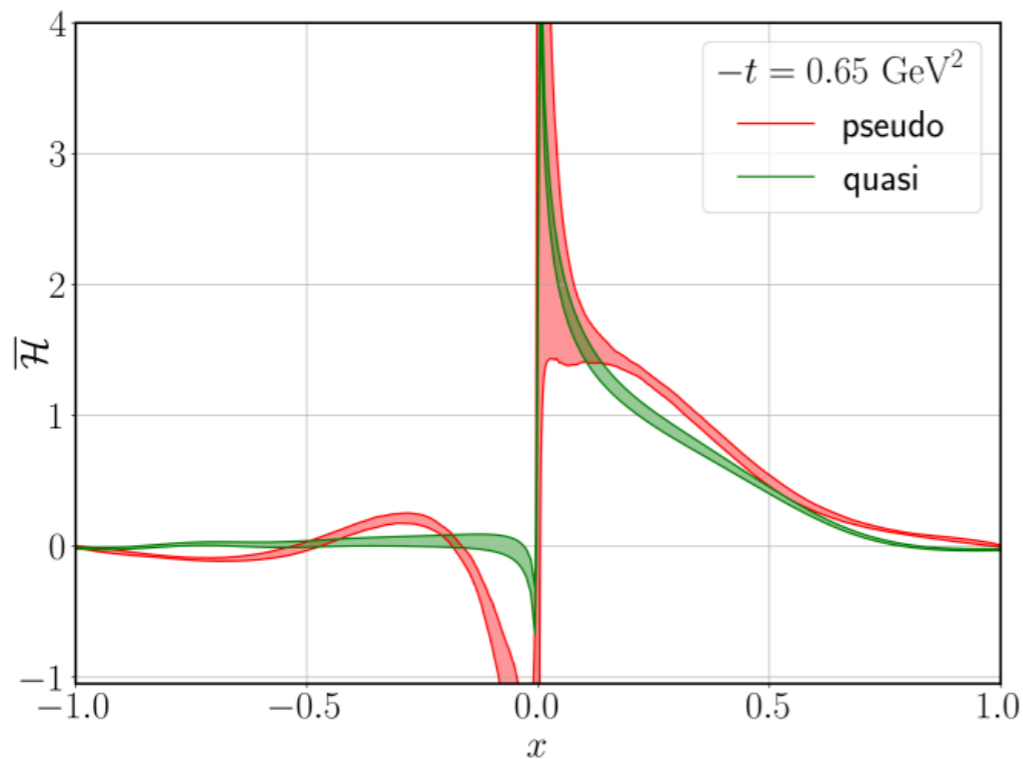
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Different steps between approaches:

- renormalization
- x -dependence reconstruction
- matching formalism

★ Comparison between methods helps assess systematic effects



- ★ $x < 0$ and small- x regions susceptible to systematic effects
- ★ Comparison only includes systematic uncertainties

Mellin moments from non-local operators

[S. Bhattacharya et al., PRD 108 (2023) 1, 014507; arXiv:2410.03539]

★ Leading-twist factorization formula

$$\mathcal{M}(z, P, \Delta) \equiv \frac{\mathcal{F}(z, P, \Delta)}{\mathcal{F}(z, P=0, \Delta=0)} = \sum_{n=0} \frac{(-izP)^n}{n!} \frac{C_n^{\overline{\text{MS}}}(\mu^2 z^2)}{C_0^{\overline{\text{MS}}}(\mu^2 z^2)} \langle x^n \rangle + \mathcal{O}(\Lambda_{\text{QCD}}^2 z^2)$$

★ Avoid power-divergent mixing of multi-derivative operators

★ Wilson coefficients known to NLO (or NNLO)

★ Both isovector and isoscalar (ignores disconnected; found to be tiny)

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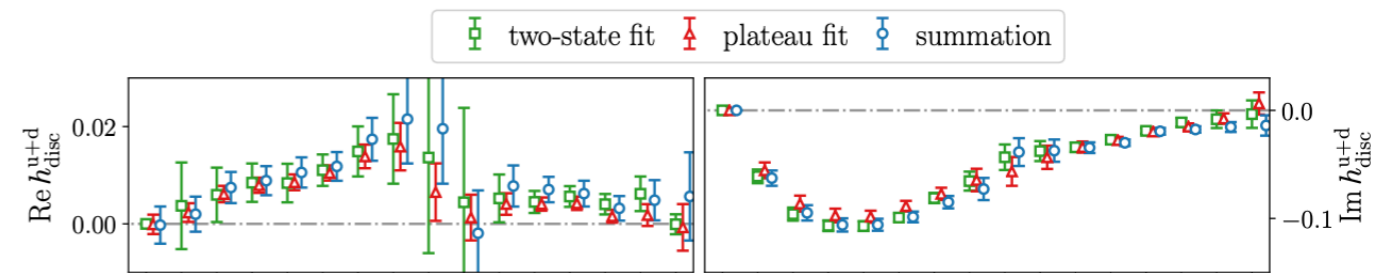
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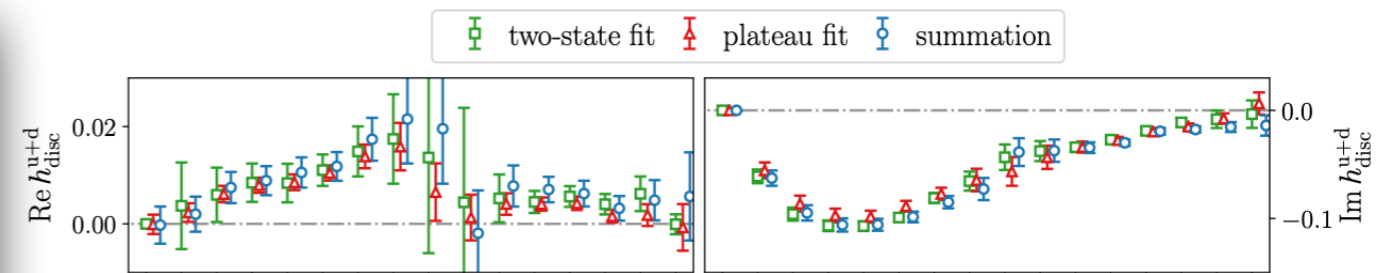
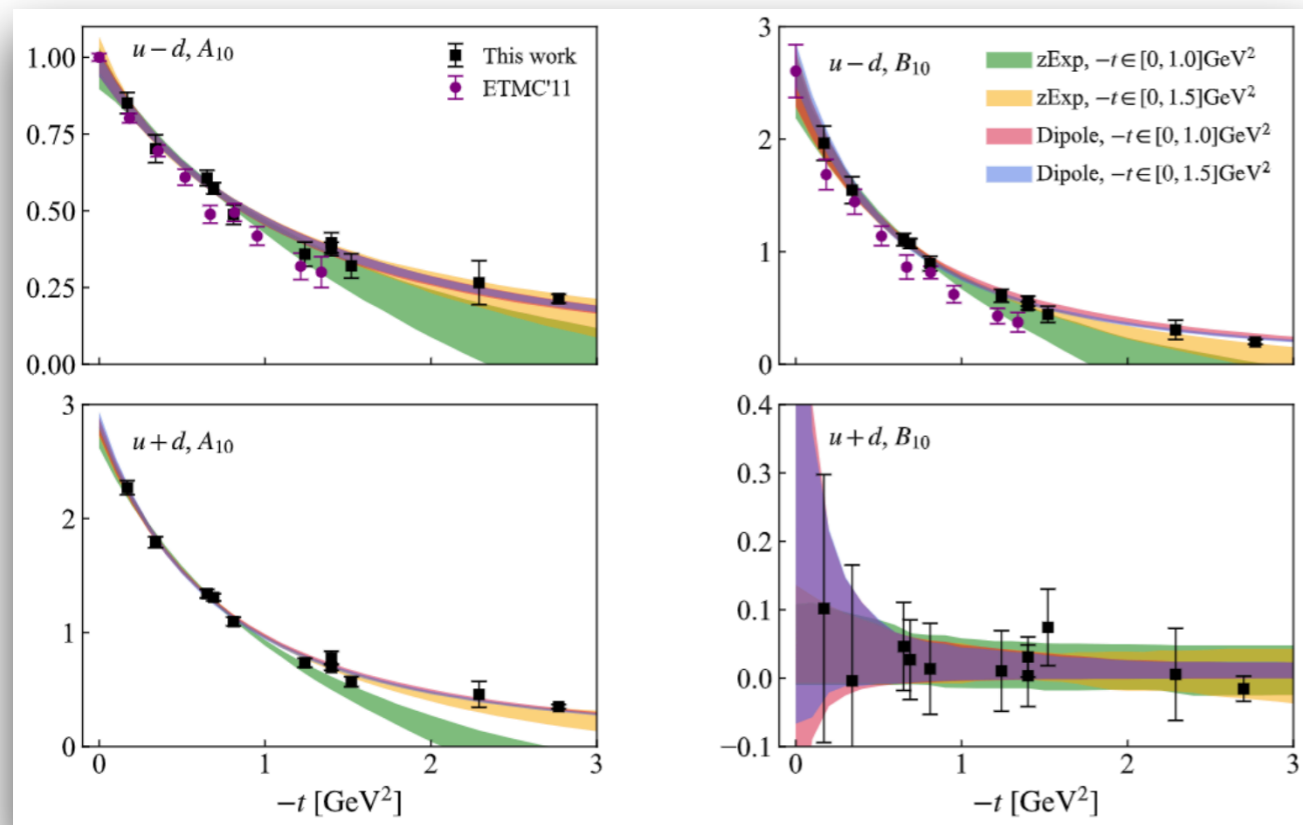
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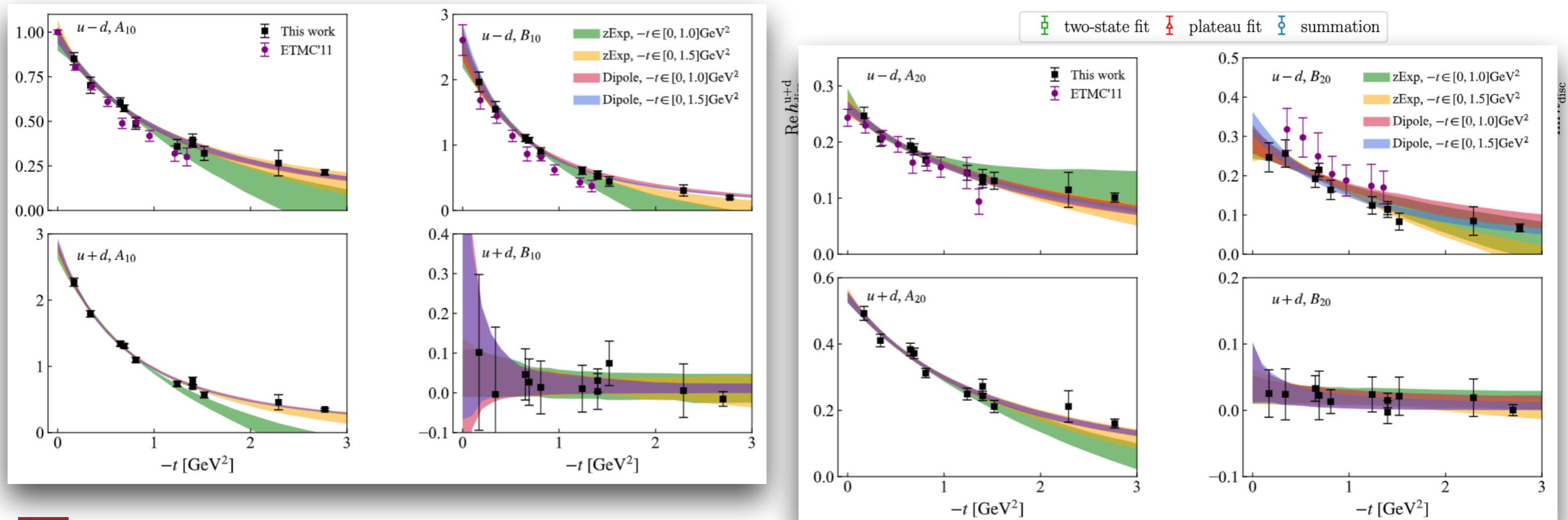
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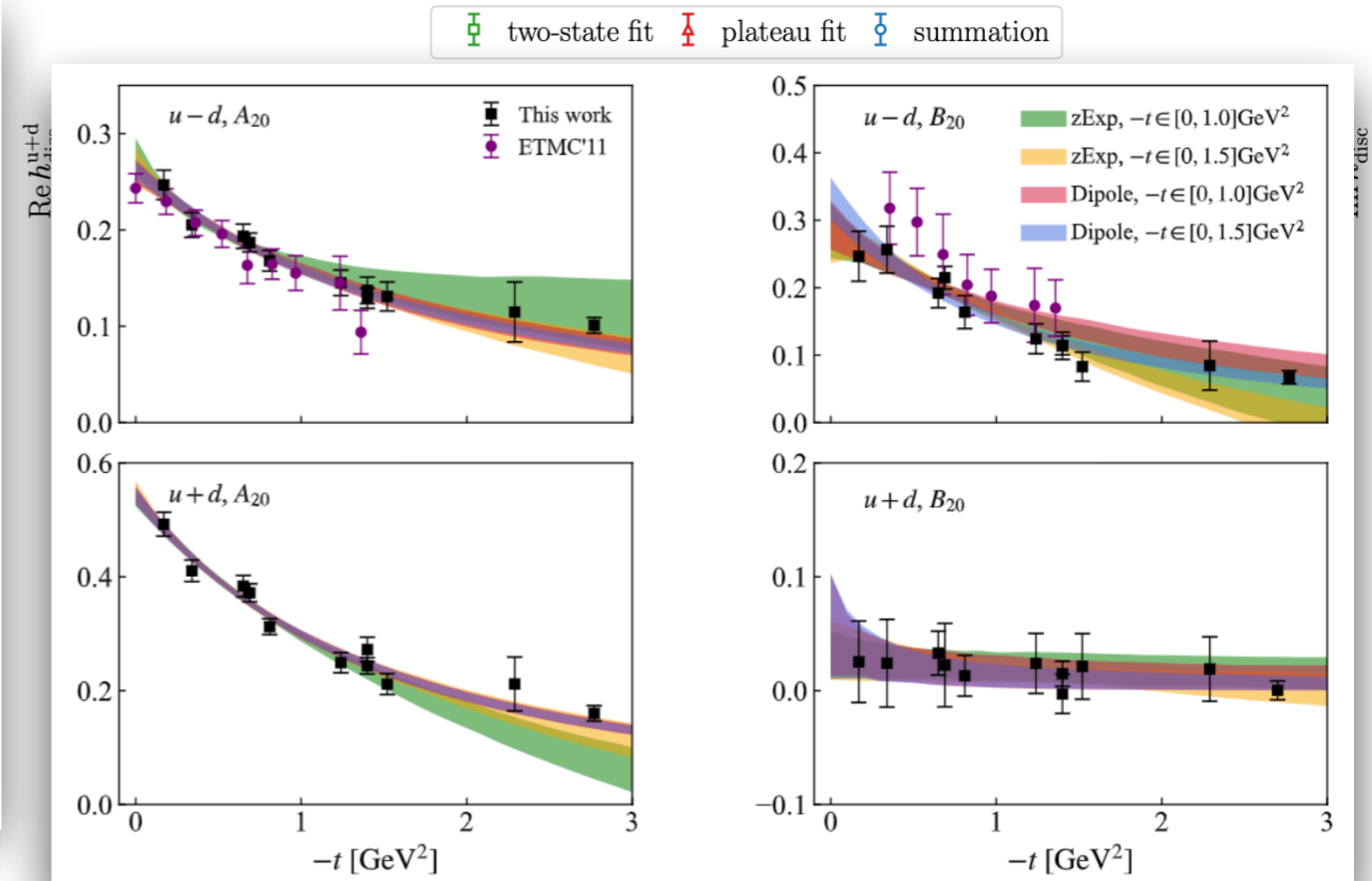
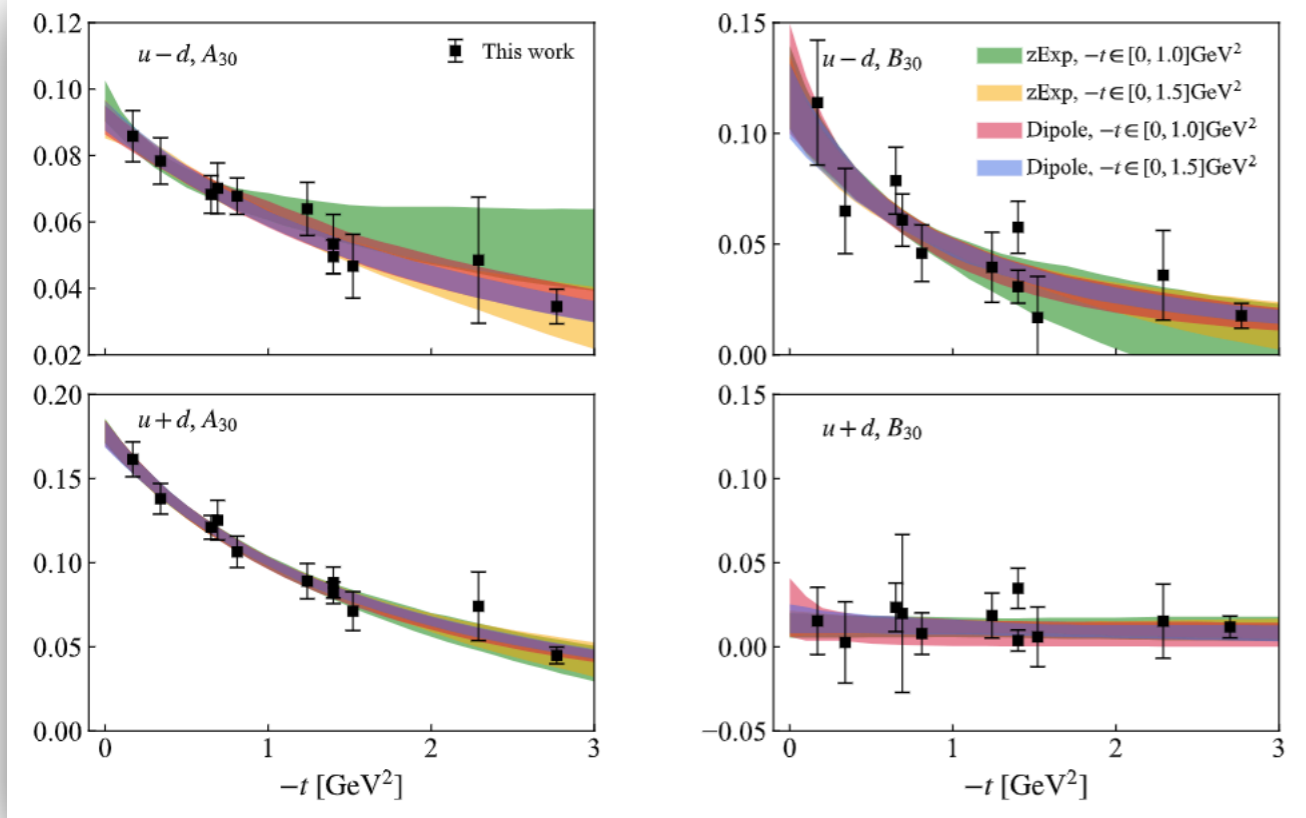
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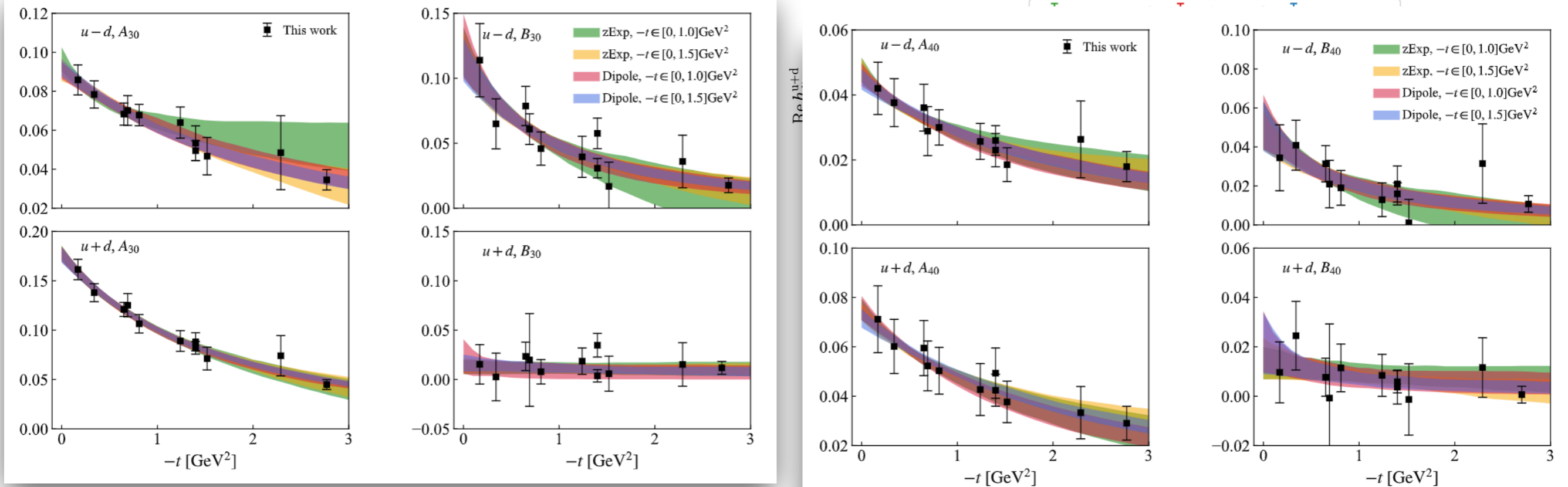
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UNSC

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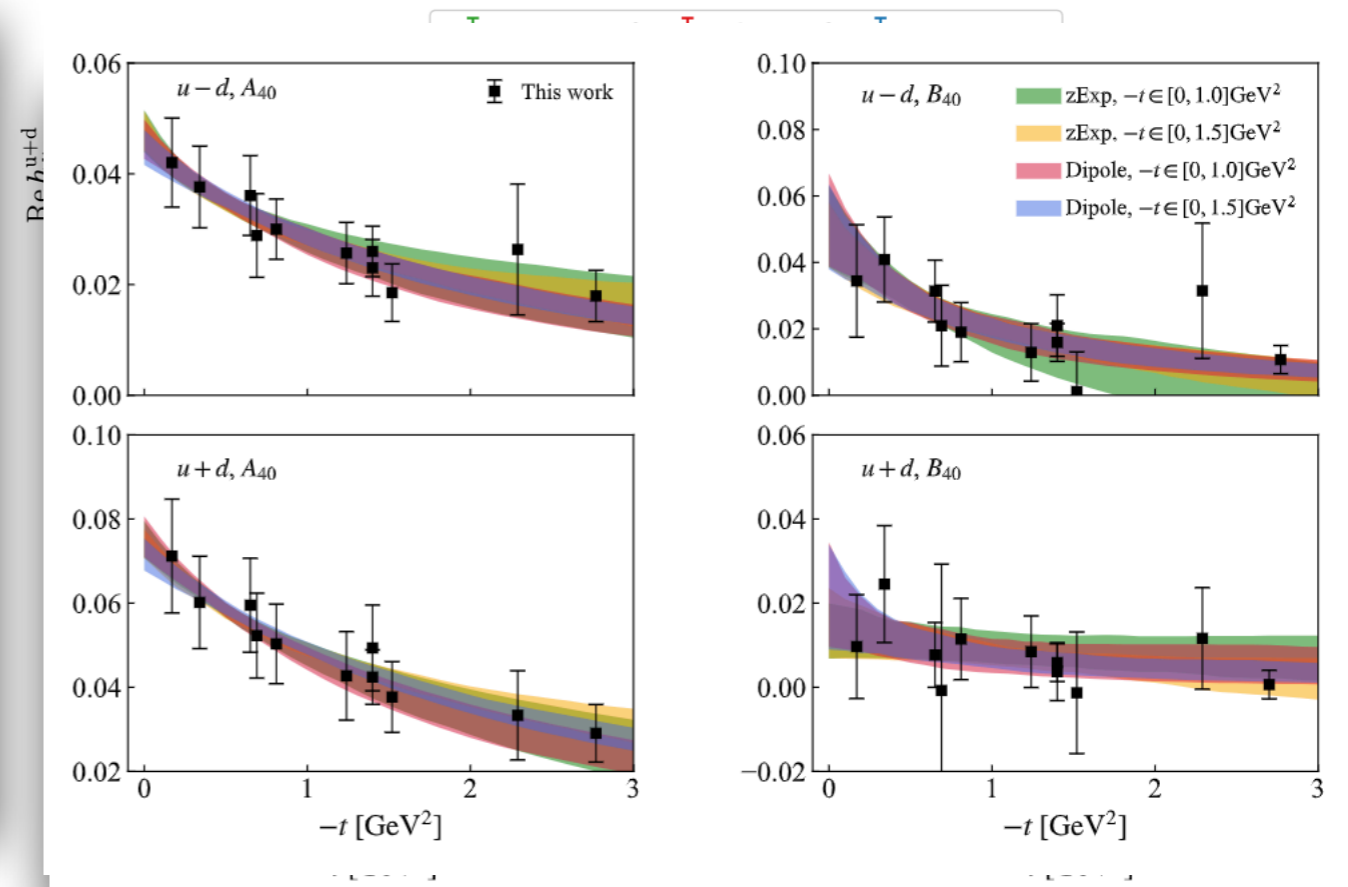
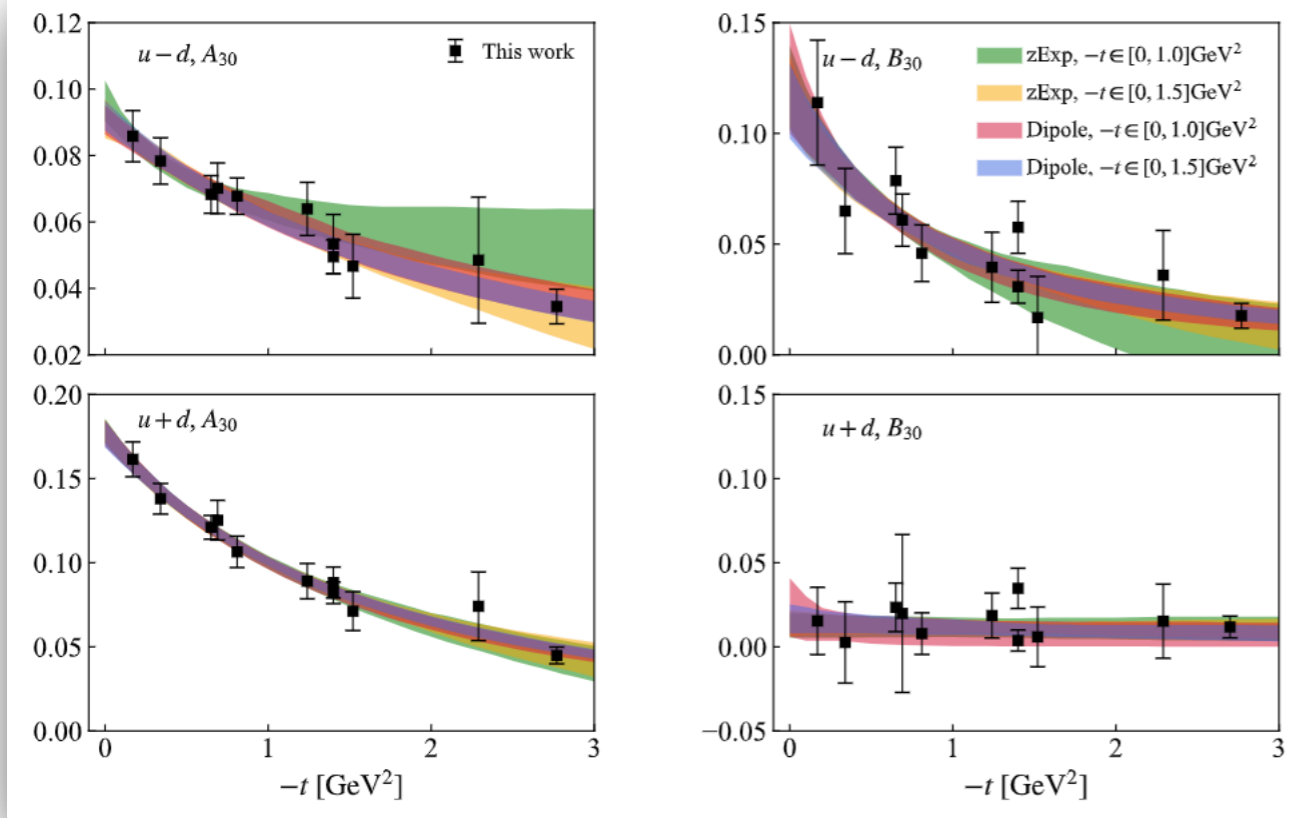
Access to Mellin moments beyond local operators

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UNSC



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beyond leading twist

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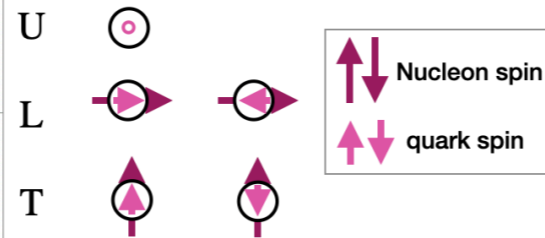
Twist-classification of PDFs, GPDs, TMDs

★ Twist: The order in Q^{-1} entering factorization

$$f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \dots$$

Twist-2 ($f_i^{(0)}$)

Quark \ Nucleon	U (γ^+)	L ($\gamma^+\gamma^5$)	T (σ^{+j})
U	$H(x, \xi, t)$ $E(x, \xi, t)$ unpolarized		
L		$\widetilde{H}(x, \xi, t)$ $\widetilde{E}(x, \xi, t)$ helicity	
T			H_T, E_T $\widetilde{H}_T, \widetilde{E}_T$ transversity



(Selected) Twist-3 ($f_i^{(1)}$)

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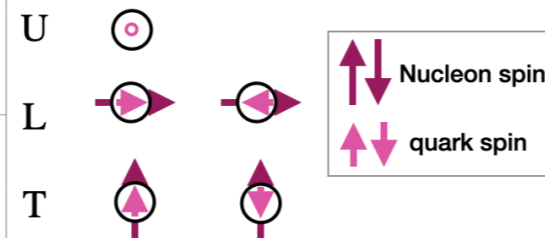
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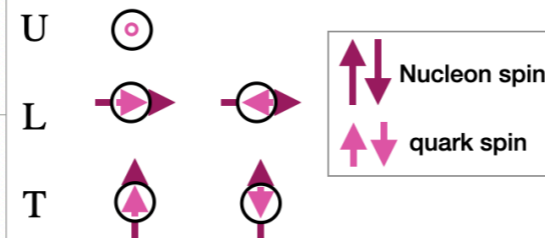
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The extraction of twist-3 is very challenges both experimentally and theoretically

Theoretical setup

★ Correlation functions in coordinate space

$$F^{[\Gamma]}(x, \Delta; P^3) = \frac{1}{2} \int \frac{dz^3}{2\pi} e^{ik \cdot z} \langle p_f, \lambda' | \bar{\psi}(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p_i, \lambda \rangle \Big|_{z^0=0, \vec{z}_\perp=\vec{0}_\perp}$$

★ Parametrization of coordinate-space correlation functions

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018)]

$$F^{[\gamma^\mu \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^3} \bar{u}(p_f, \lambda') \left[P^\mu \frac{\gamma^3 \gamma_5}{P^0} F_{\tilde{H}}(x, \xi, t; P^3) + P^\mu \frac{\Delta^3 \gamma_5}{2m P^0} F_{\tilde{E}}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma_5}{2m} F_{\tilde{E}+\tilde{G}_1}(x, \xi, t; P^3) + \gamma_\perp^\mu \gamma_5 F_{\tilde{H}+\tilde{G}_2}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma^3 \gamma_5}{P^3} F_{\tilde{G}_3}(x, \xi, t; P^3) + i \varepsilon_\perp^{\mu\nu} \Delta_\nu \frac{\gamma^3}{P^3} F_{\tilde{G}_4}(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

Theoretical setup

★ Correlation functions in coordinate space

$$F^{[\Gamma]}(x, \Delta; P^3) = \frac{1}{2} \int \frac{dz^3}{2\pi} e^{ik \cdot z} \langle p_f, \lambda' | \bar{\psi}(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p_i, \lambda \rangle \Big|_{z^0=0, \vec{z}_\perp=\vec{0}_\perp}$$

★ Parametrization of coordinate-space correlation functions

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018)]

$$F^{[\gamma^\mu \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^3} \bar{u}(p_f, \lambda') \left[P^\mu \frac{\gamma^3 \gamma_5}{P^0} F_{\tilde{H}}(x, \xi, t; P^3) + P^\mu \frac{\Delta^3 \gamma_5}{2m P^0} F_{\tilde{E}}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma_5}{2m} F_{\tilde{E}+\tilde{G}_1}(x, \xi, t; P^3) + \gamma_\perp^\mu \gamma_5 F_{\tilde{H}+\tilde{G}_2}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma^3 \gamma_5}{P^3} F_{\tilde{G}_3}(x, \xi, t; P^3) + i \varepsilon_\perp^{\mu\nu} \Delta_\nu \frac{\gamma^3}{P^3} F_{\tilde{G}_4}(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

Theoretical setup

★ Correlation functions in coordinate space

$$F^{[\Gamma]}(x, \Delta; P^3) = \frac{1}{2} \int \frac{dz^3}{2\pi} e^{ik \cdot z} \langle p_f, \lambda' | \bar{\psi}(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p_i, \lambda \rangle \Big|_{z^0=0, \vec{z}_\perp=\vec{0}_\perp}$$

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★ Twist-3 contributions to helicity GPDs: $\gamma^{1,2} \gamma_5$

Theoretical setup

★ Correlation functions in coordinate space

$$F^{[\Gamma]}(x, \Delta; P^3) = \frac{1}{2} \int \frac{dz^3}{2\pi} e^{ik \cdot z} \langle p_f, \lambda' | \bar{\psi}(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p_i, \lambda \rangle \Big|_{z^0=0, \vec{z}_\perp=\vec{0}_\perp}$$

★ Parametrization of coordinate-space correlation functions

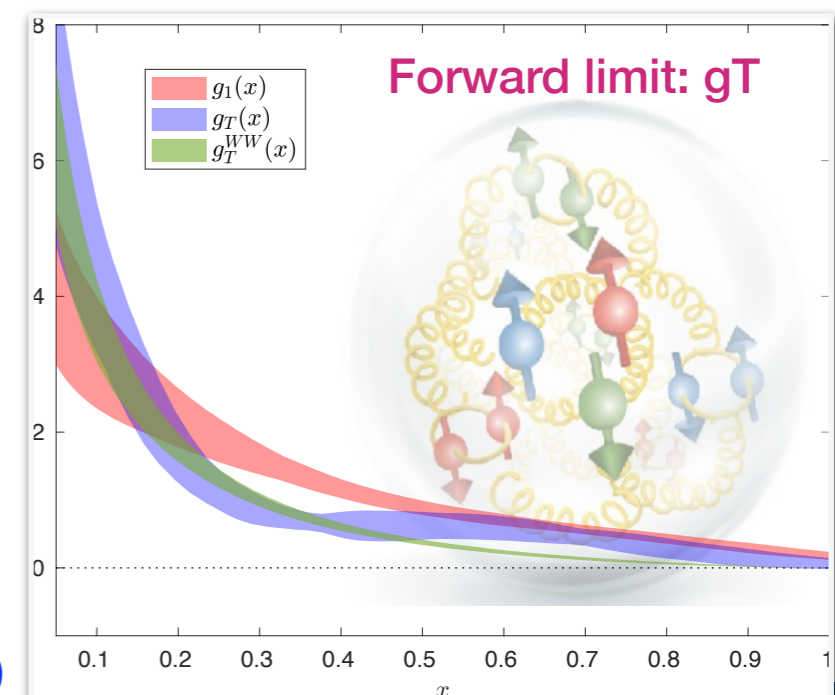
[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018)]

$$F^{[\gamma^\mu \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^3} \bar{u}(p_f, \lambda') \left[P^\mu \frac{\gamma^3 \gamma_5}{P^0} F_{\tilde{H}}(x, \xi, t; P^3) + P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} F_{\tilde{E}}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma_5}{2m} F_{\tilde{E}+\tilde{G}_1}(x, \xi, t; P^3) + \gamma_\perp^\mu \gamma_5 F_{\tilde{H}+\tilde{G}_2}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma^3 \gamma_5}{P^3} F_{\tilde{G}_3}(x, \xi, t; P^3) + i\varepsilon_\perp^{\mu\nu} \Delta_\nu \frac{\gamma^3}{P^3} F_{\tilde{G}_4}(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

★ Twist-3 contributions to helicity GPDs: $\gamma^{1,2} \gamma_5$

[S. Bhattacharya et al., PRD 102 (2020) 11] (Editors Highlight)



Theoretical setup

★ Correlation functions in coordinate space

$$F^{[\Gamma]}(x, \Delta; P^3) = \frac{1}{2} \int \frac{dz^3}{2\pi} e^{ik \cdot z} \langle p_f, \lambda' | \bar{\psi}(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p_i, \lambda \rangle \Big|_{z^0=0, \vec{z}_\perp=\vec{0}_\perp}$$

★ Parametrization of coordinate-space correlation functions

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018)]

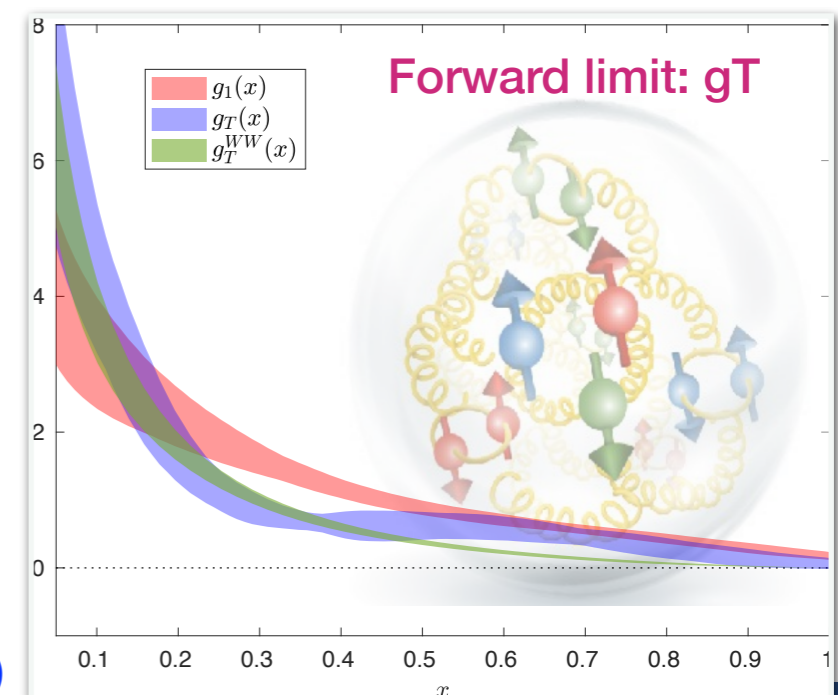
$$F^{[\gamma^\mu \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^3} \bar{u}(p_f, \lambda') \left[P^\mu \frac{\gamma^3 \gamma_5}{P^0} F_{\tilde{H}}(x, \xi, t; P^3) + P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} F_{\tilde{E}}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma_5}{2m} F_{\tilde{E}+\tilde{G}_1}(x, \xi, t; P^3) + \gamma_\perp^\mu \gamma_5 F_{\tilde{H}+\tilde{G}_2}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma^3 \gamma_5}{P^3} F_{\tilde{G}_3}(x, \xi, t; P^3) + i\varepsilon_\perp^{\mu\nu} \Delta_\nu \frac{\gamma^3}{P^3} F_{\tilde{G}_4}(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

★ Twist-3 contributions to helicity GPDs: $\gamma^{1,2} \gamma_5$

★ Kinematic twist-3 contributions to pseudo- and quasi-GPDs to restore translation invariance

[V. Braun et al., JHEP 10 (2023) 134]

[S. Bhattacharya et al., PRD 102 (2020) 11] (Editors Highlight)



$$F^{[\gamma^\mu \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^3} \bar{u}(p_f, \lambda') \left[P^\mu \frac{\gamma^3 \gamma_5}{P^0} F_{\tilde{H}}(x, \xi, t; P^3) + P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} F_{\tilde{E}}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma_5}{2m} F_{\tilde{E}+\tilde{G}_1}(x, \xi, t; P^3) + \gamma_\perp^\mu \gamma_5 F_{\tilde{H}+\tilde{G}_2}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma^3 \gamma_5}{P^3} F_{\tilde{G}_3}(x, \xi, t; P^3) + i\varepsilon_\perp^{\mu\nu} \Delta_\nu \frac{\gamma^3}{P^3} F_{\tilde{G}_4}(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

$$\tilde{F}^\mu(z, P, \Delta) \equiv \langle p_f; \lambda' | \bar{\psi}(-\frac{z}{2}) \gamma^\mu \gamma_5 \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p_i; \lambda \rangle \\ = \bar{u}(p_f, \lambda') \left[\frac{i\varepsilon^{\mu P z \Delta}}{m} \tilde{A}_1 + \gamma^\mu \gamma_5 \tilde{A}_2 + \gamma_5 \left(\frac{P^\mu}{m} \tilde{A}_3 + m z^\mu \tilde{A}_4 + \frac{\Delta^\mu}{m} \tilde{A}_5 \right) \right. \\ \left. + m \not{z} \gamma_5 \left(\frac{P^\mu}{m} \tilde{A}_6 + m z^\mu \tilde{A}_7 + \frac{\Delta^\mu}{m} \tilde{A}_8 \right) \right] u(p_i, \lambda),$$

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

[S. Bhattacharya et al., 109 (2024) 3, 034508]

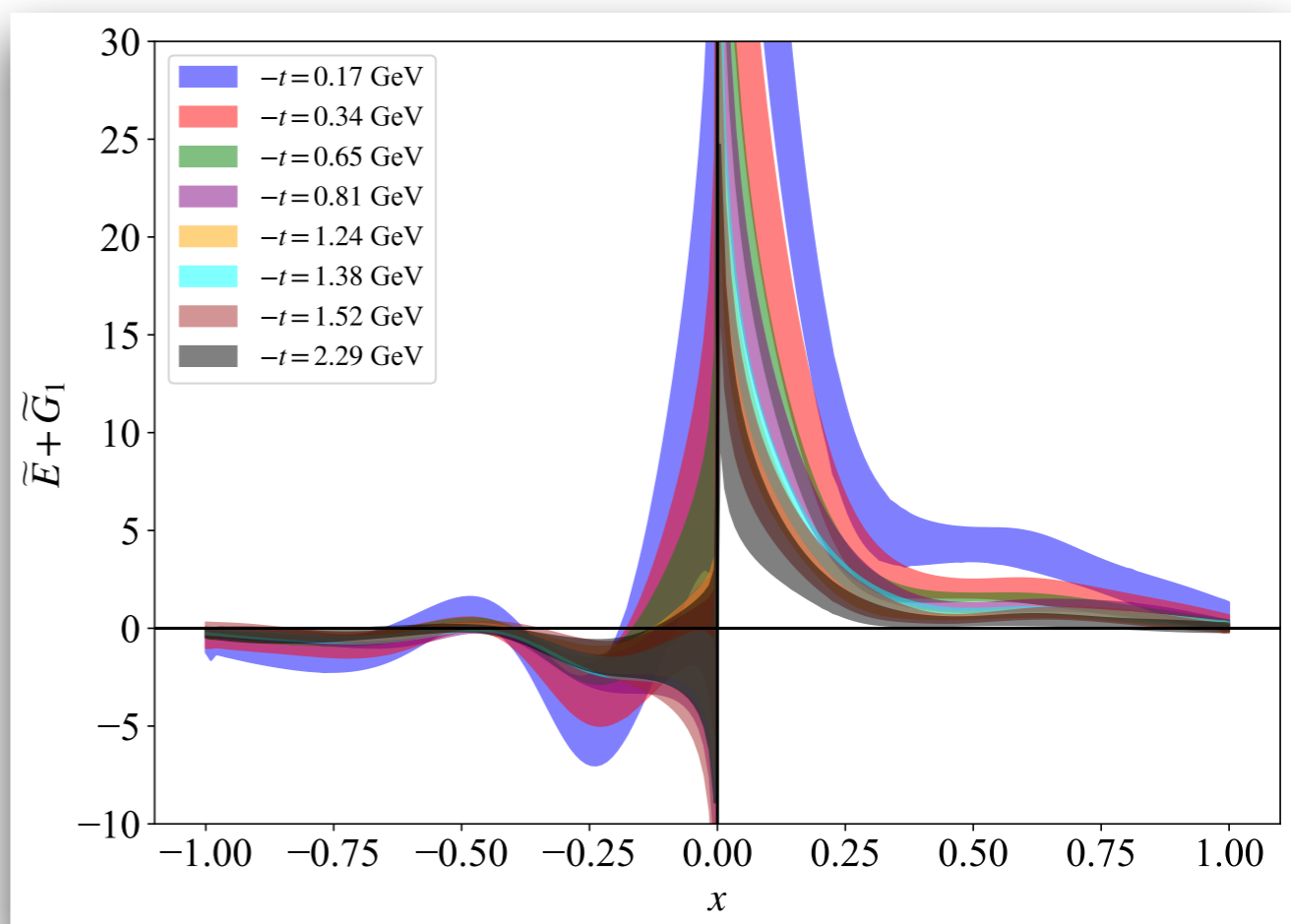
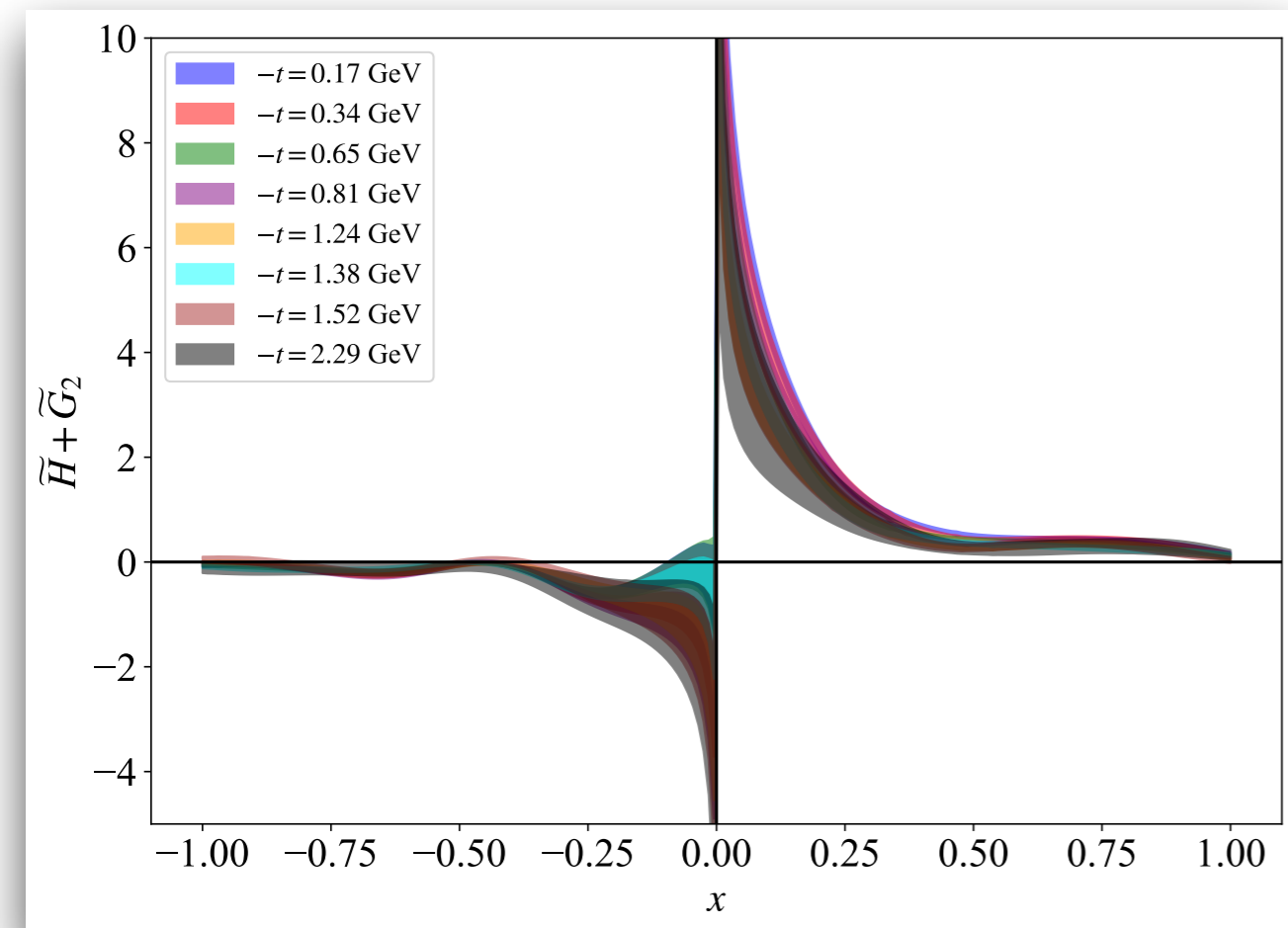
$$F^{[\gamma^\mu \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^3} \bar{u}(p_f, \lambda') \left[P^\mu \frac{\gamma^3 \gamma_5}{P^0} F_{\tilde{H}}(x, \xi, t; P^3) + P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} F_{\tilde{E}}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma_5}{2m} F_{\tilde{E}+\tilde{G}_1}(x, \xi, t; P^3) + \gamma_\perp^\mu \gamma_5 F_{\tilde{H}+\tilde{G}_2}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma^3 \gamma_5}{P^3} F_{\tilde{G}_3}(x, \xi, t; P^3) + i\varepsilon_\perp^{\mu\nu} \Delta_\nu \frac{\gamma^3}{P^3} F_{\tilde{G}_4}(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

$$\tilde{F}^\mu(z, P, \Delta) \equiv \langle p_f; \lambda' | \bar{\psi}(-\frac{z}{2}) \gamma^\mu \gamma_5 \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p_i; \lambda \rangle$$

$$= \bar{u}(p_f, \lambda') \left[\frac{i\varepsilon^{\mu P z \Delta}}{m} \tilde{A}_1 + \gamma^\mu \gamma_5 \tilde{A}_2 + \gamma_5 \left(\frac{P^\mu}{m} \tilde{A}_3 + m z^\mu \tilde{A}_4 + \frac{\Delta^\mu}{m} \tilde{A}_5 \right) \right. \\ \left. + m \not{z} \gamma_5 \left(\frac{P^\mu}{m} \tilde{A}_6 + m z^\mu \tilde{A}_7 + \frac{\Delta^\mu}{m} \tilde{A}_8 \right) \right] u(p_i, \lambda),$$

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

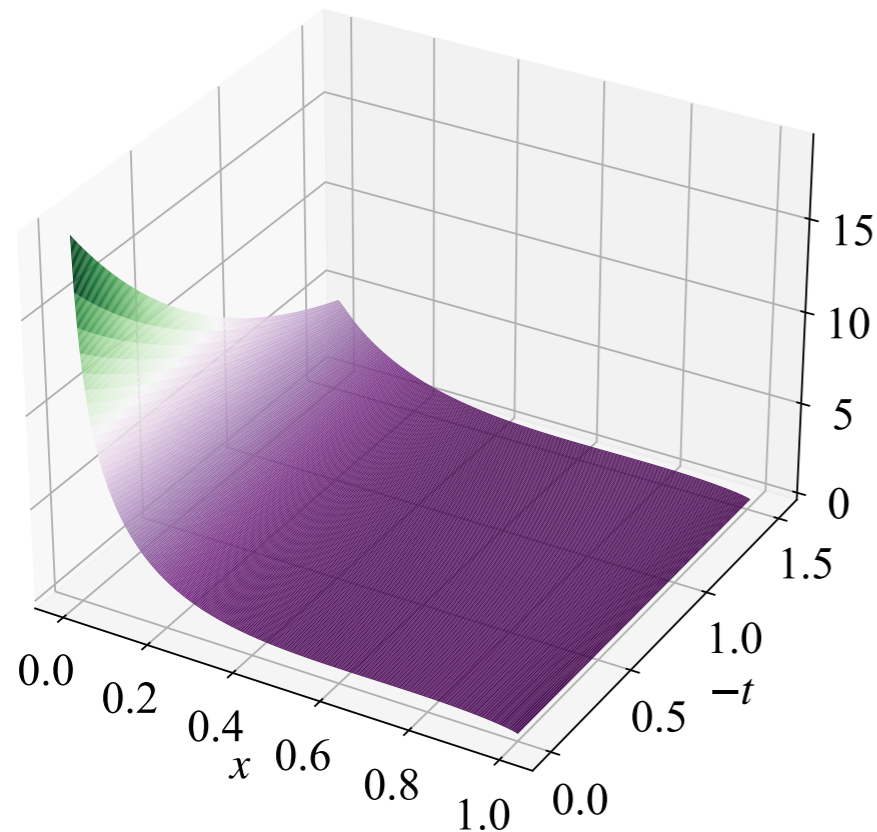
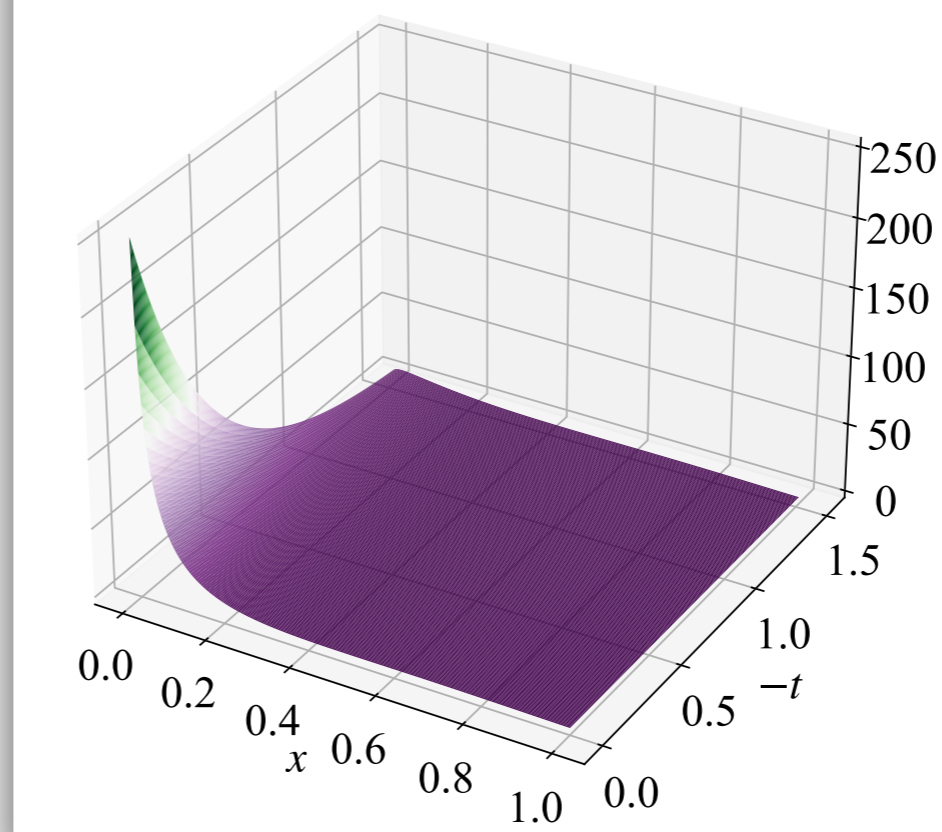
[S. Bhattacharya et al., 109 (2024) 3, 034508]



★ Parametrization of -t dependence

$$\text{GPD}(x, -t, 0) = A x^{\alpha_0 - \alpha_1 t} (1 - x)^\beta$$

Ademollo & Del Giudice Gatto & Preparata

$\tilde{H} + \tilde{G}_2$  $\tilde{E} + \tilde{G}_1$ 

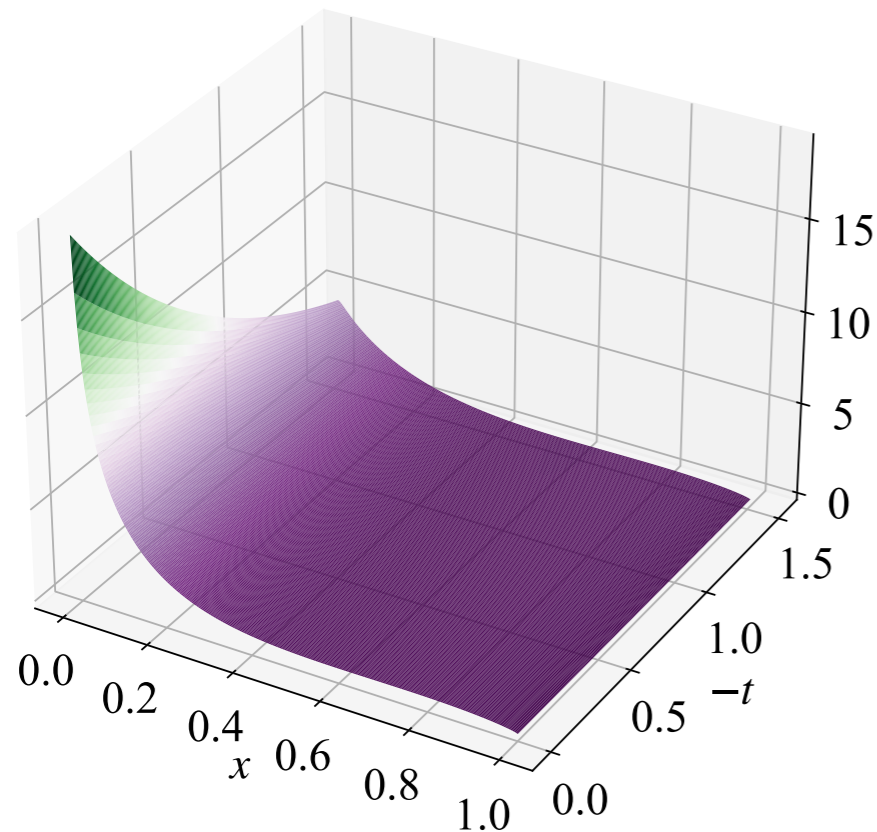
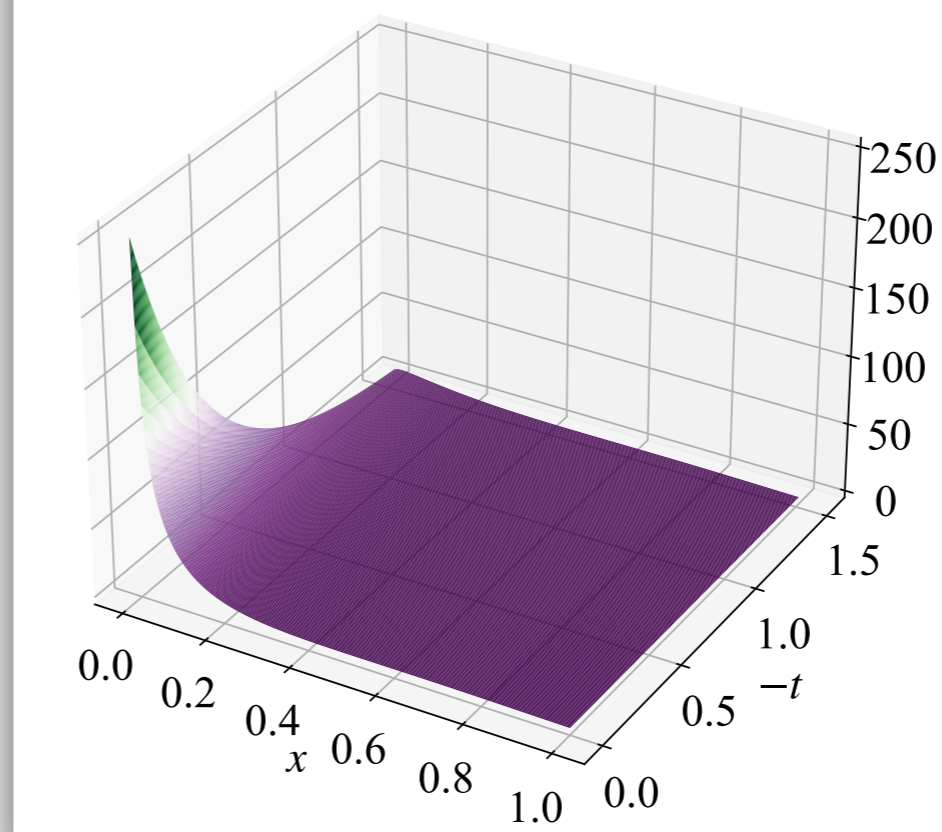
- ★ Direct access to \tilde{E} -GPD not possible for zero skewness

$$P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} \underline{F_{\tilde{E}}}(x, \xi, t; P^3)$$

- ★ Glimpse into \tilde{E} -GPD through twist-3 :

$$\int_{-1}^1 dx \tilde{E}(x, \xi, t) = G_P(t)$$

$$\int_{-1}^1 dx \tilde{G}_i(x, \xi, t) = 0, \quad i = 1, 2, 3, 4$$

$\tilde{H} + \tilde{G}_2$  $\tilde{E} + \tilde{G}_1$ 

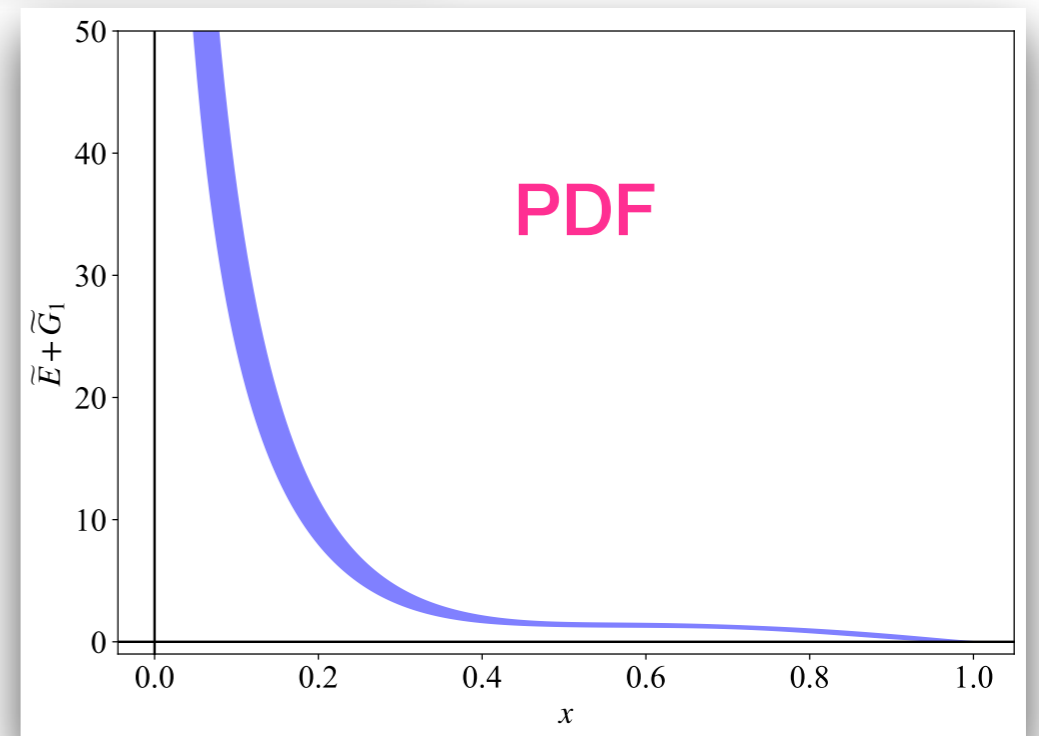
- ★ Direct access to \tilde{E} -GPD not possible for zero skewness

$$P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} \underline{F_{\tilde{E}}}(x, \xi, t; P^3)$$

- ★ Glimpse into \tilde{E} -GPD through twist-3 :

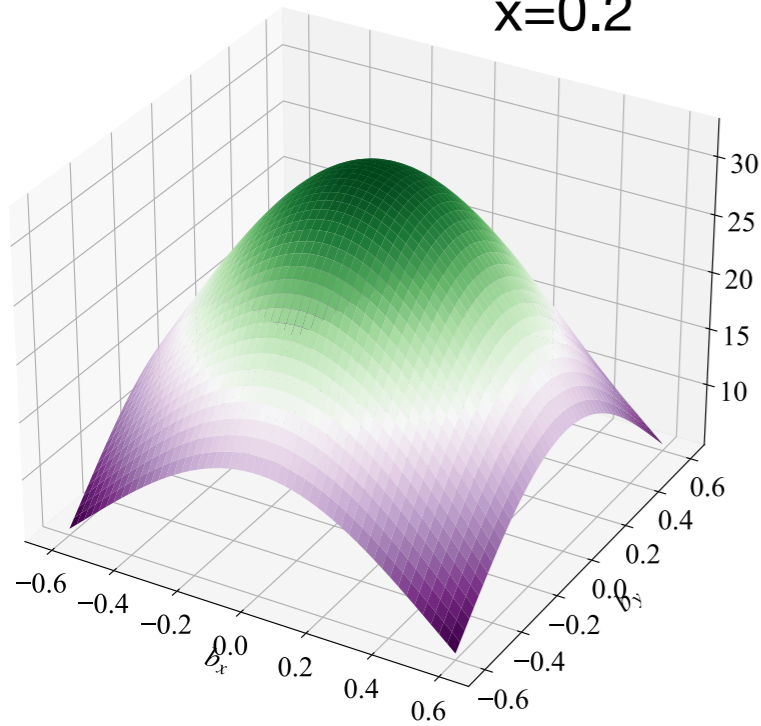
$$\int_{-1}^1 dx \tilde{E}(x, \xi, t) = G_P(t)$$

$$\int_{-1}^1 dx \tilde{G}_i(x, \xi, t) = 0, \quad i = 1, 2, 3, 4$$

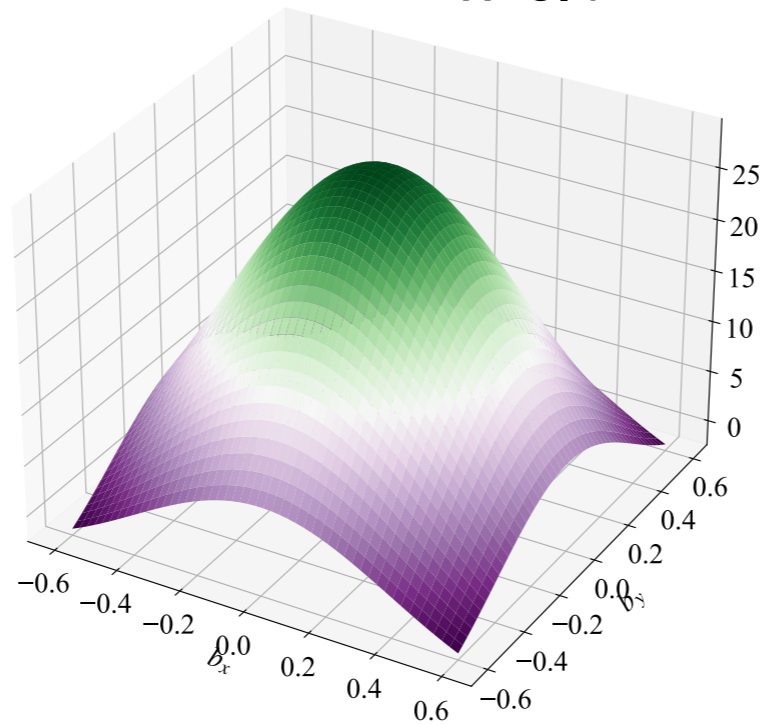


Impact parameter space $\widetilde{H} + \widetilde{G}_2$

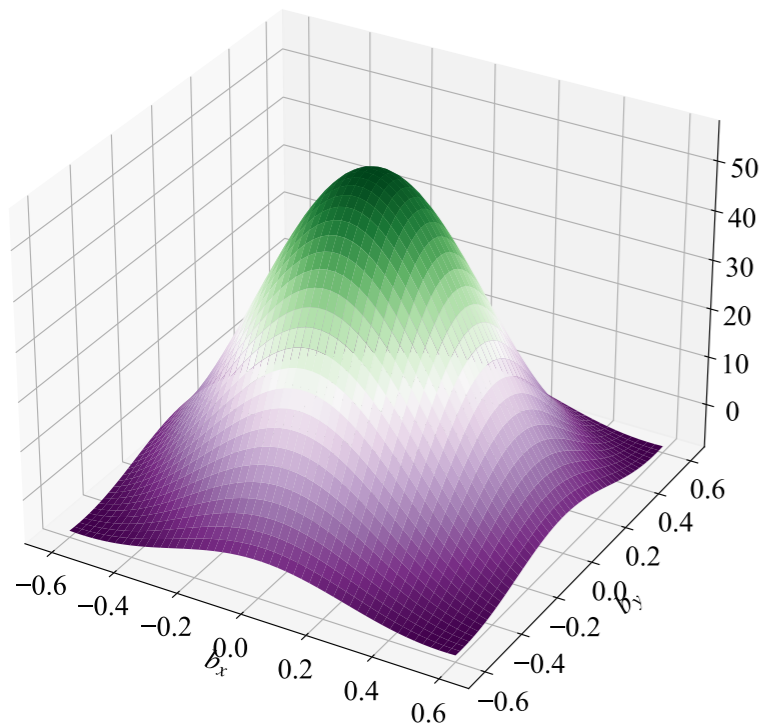
$x=0.2$



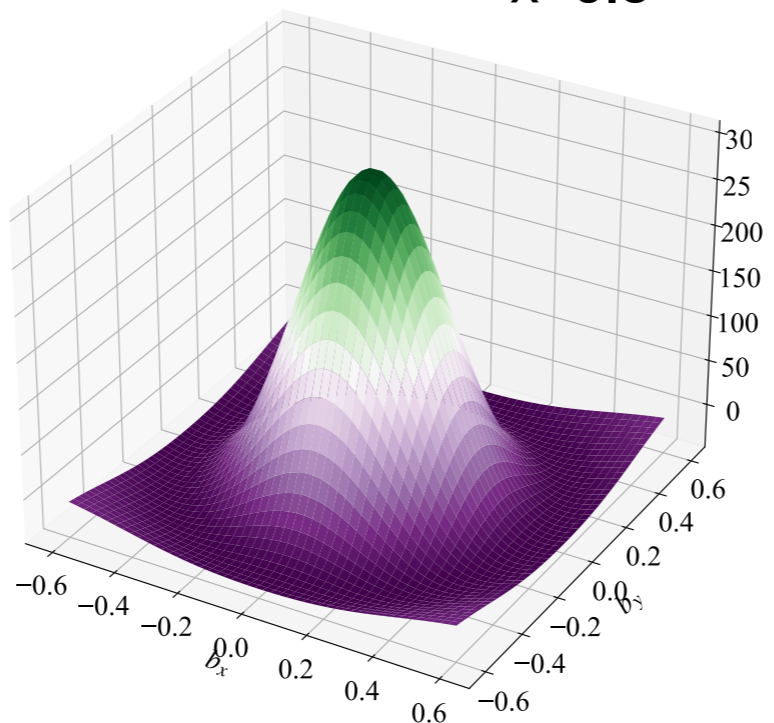
$x=0.4$



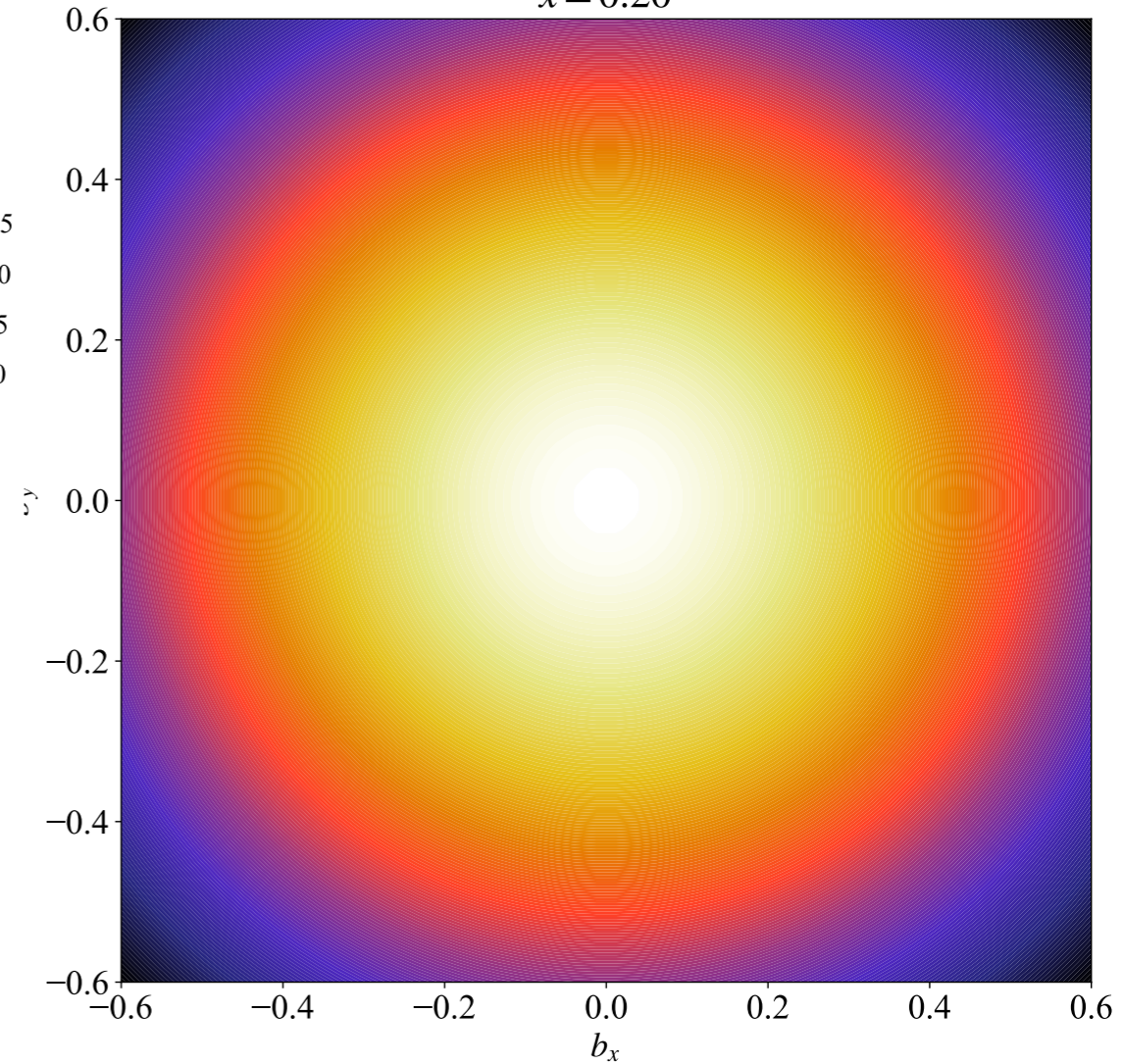
$x=0.6$



$x=0.8$



$x=0.20$



★ GPDs in transverse plane

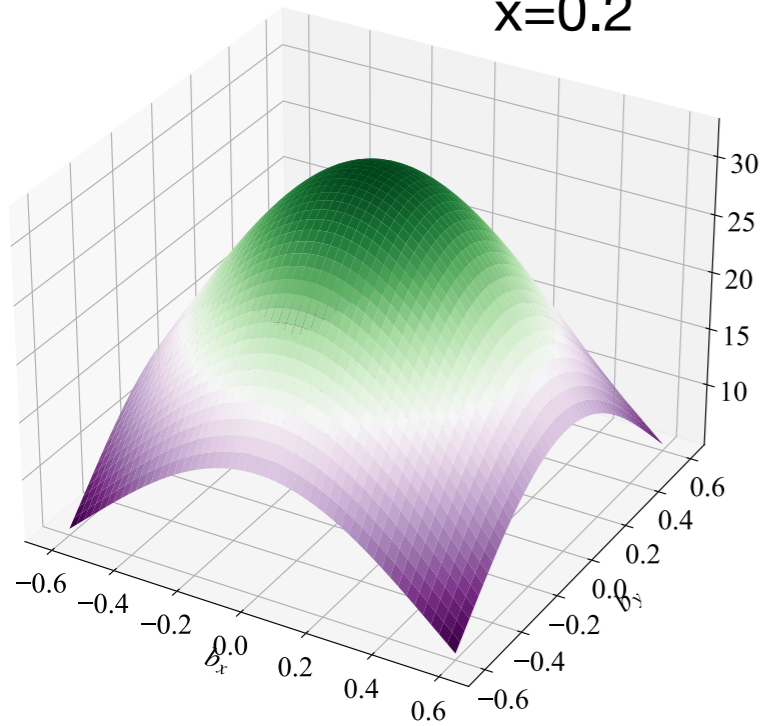
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

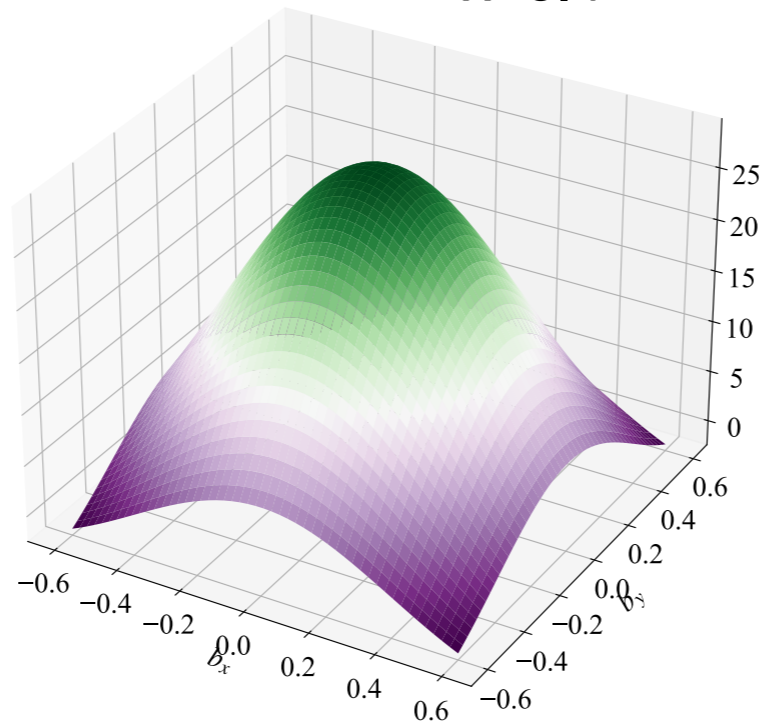
\mathbf{b}_\perp : transverse distance from the
(transverse) center of momentum

Impact parameter space $\widetilde{H} + \widetilde{G}_2$

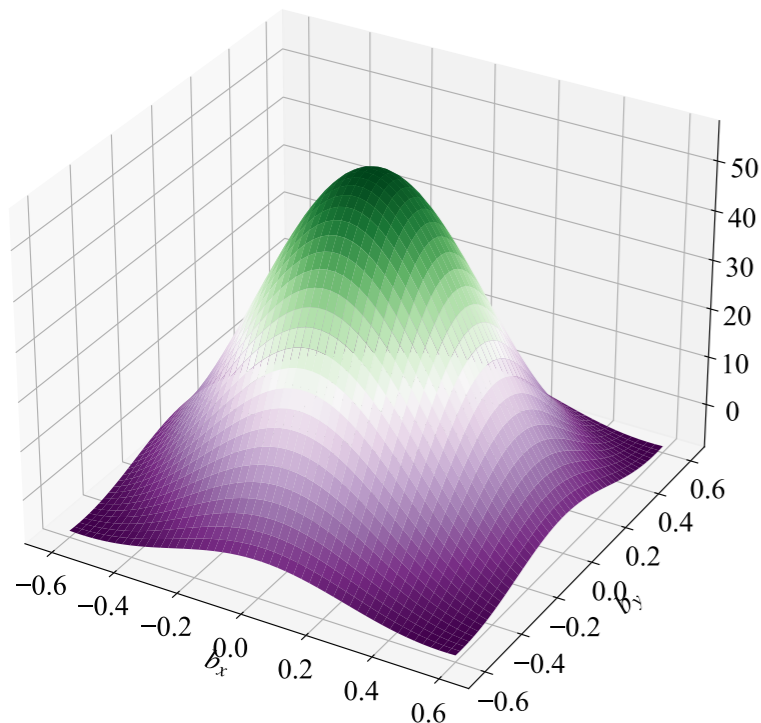
$x=0.2$



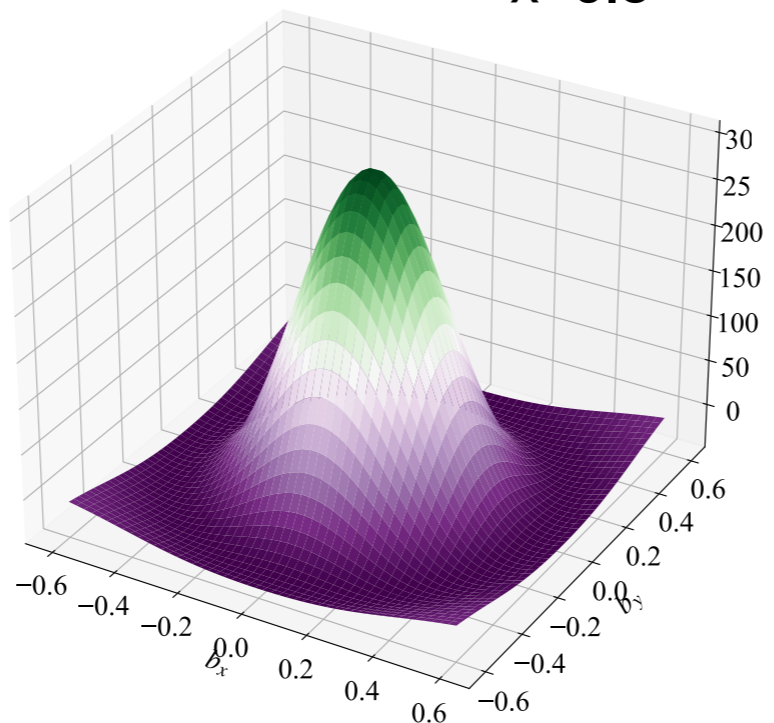
$x=0.4$



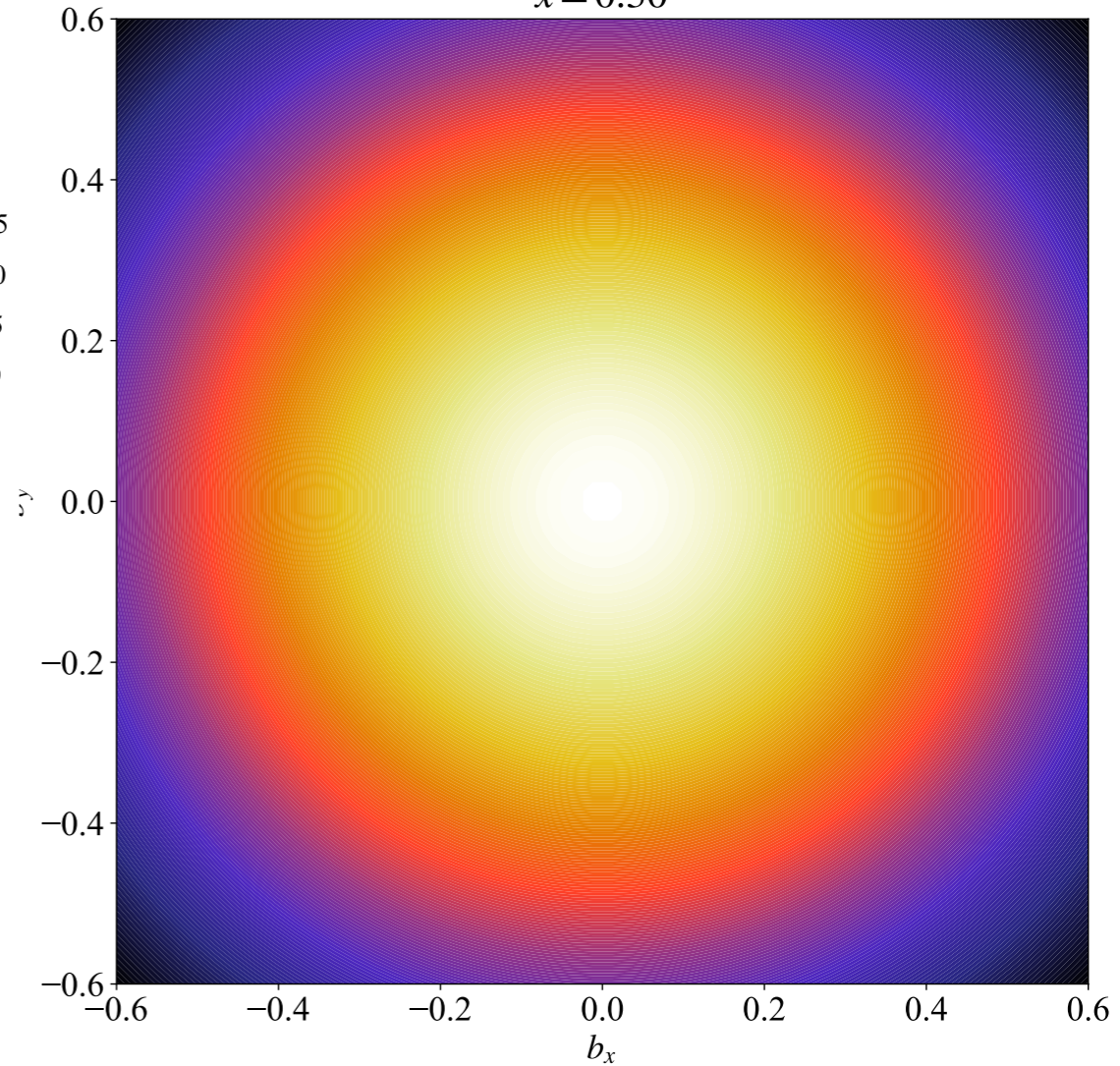
$x=0.6$



$x=0.8$



$x=0.30$



★ GPDs in transverse plane

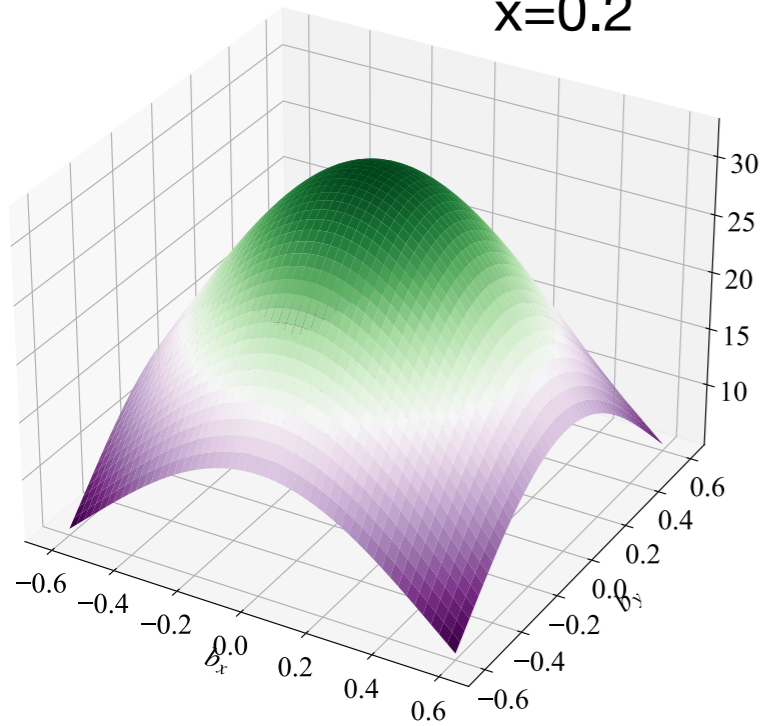
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

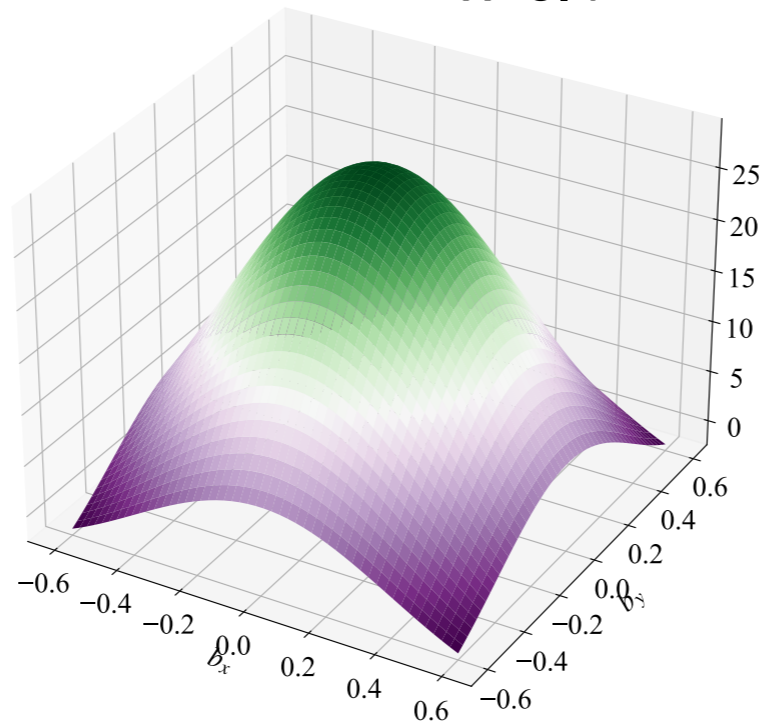
\mathbf{b}_\perp : transverse distance from the
(transverse) center of momentum

Impact parameter space $\widetilde{H} + \widetilde{G}_2$

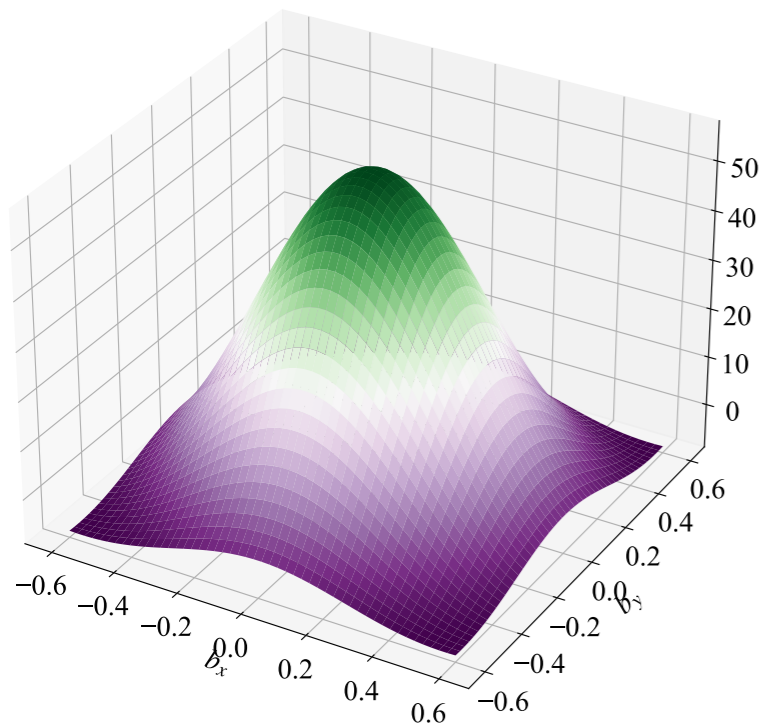
$x=0.2$



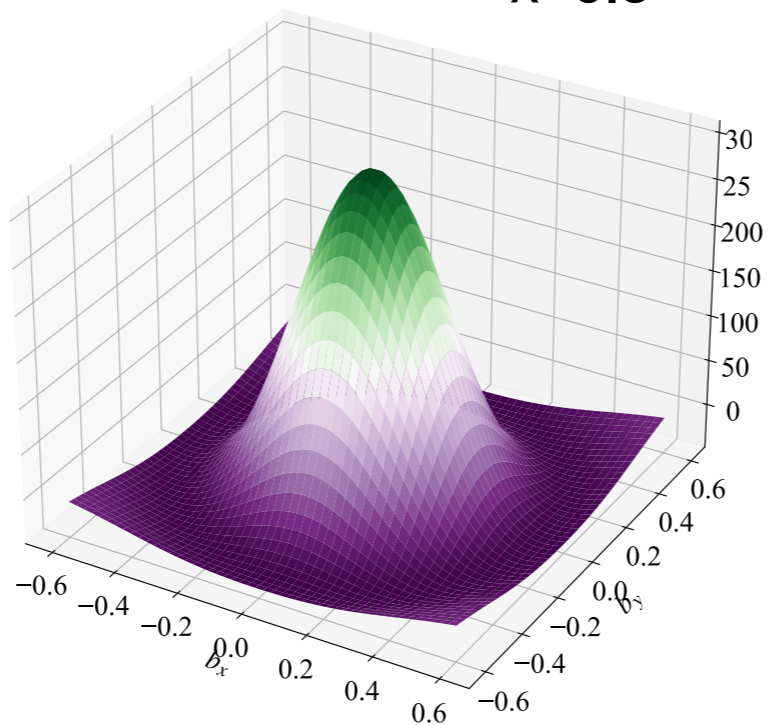
$x=0.4$



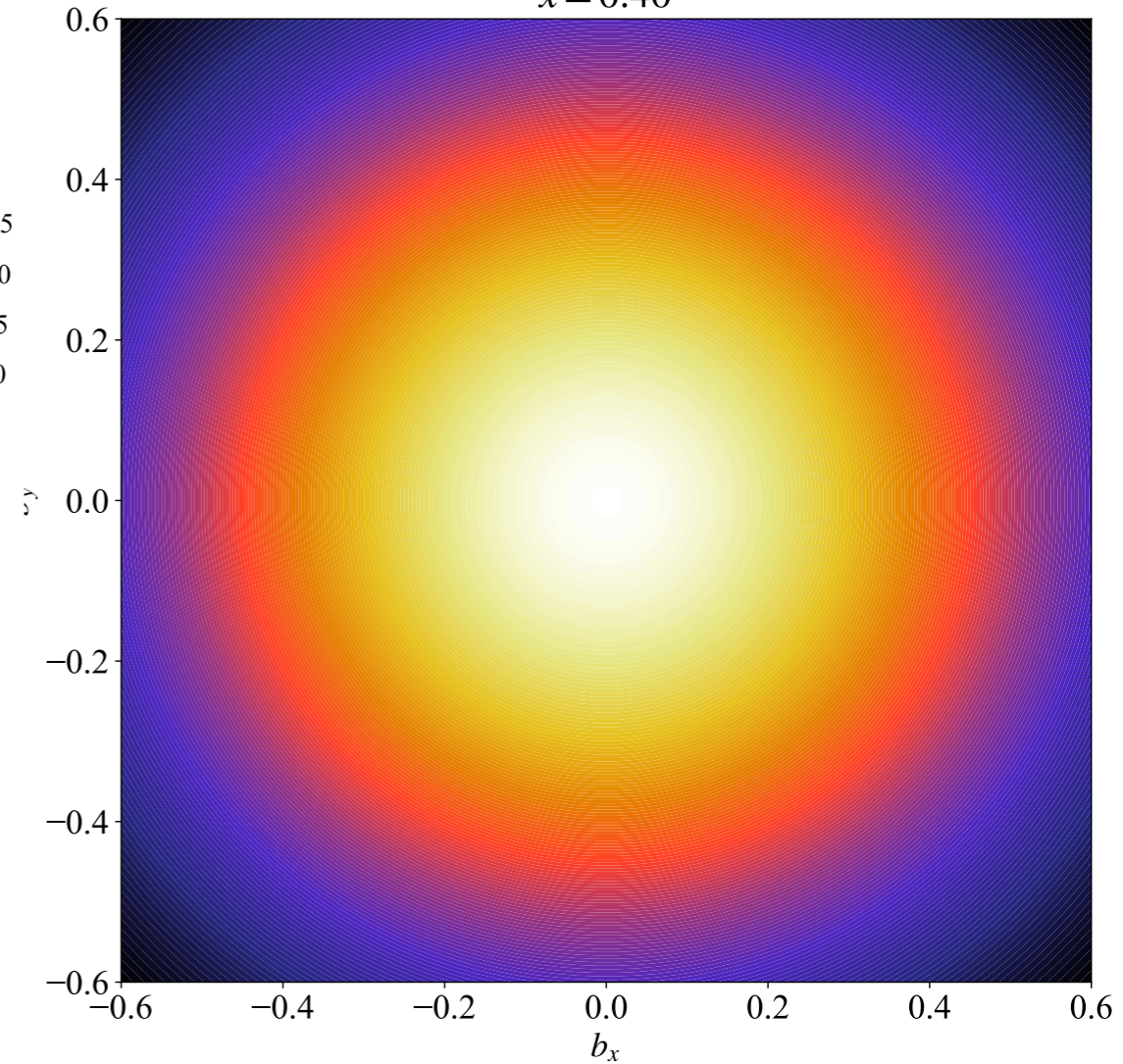
$x=0.6$



$x=0.8$



$x=0.40$



★ GPDs in transverse plane

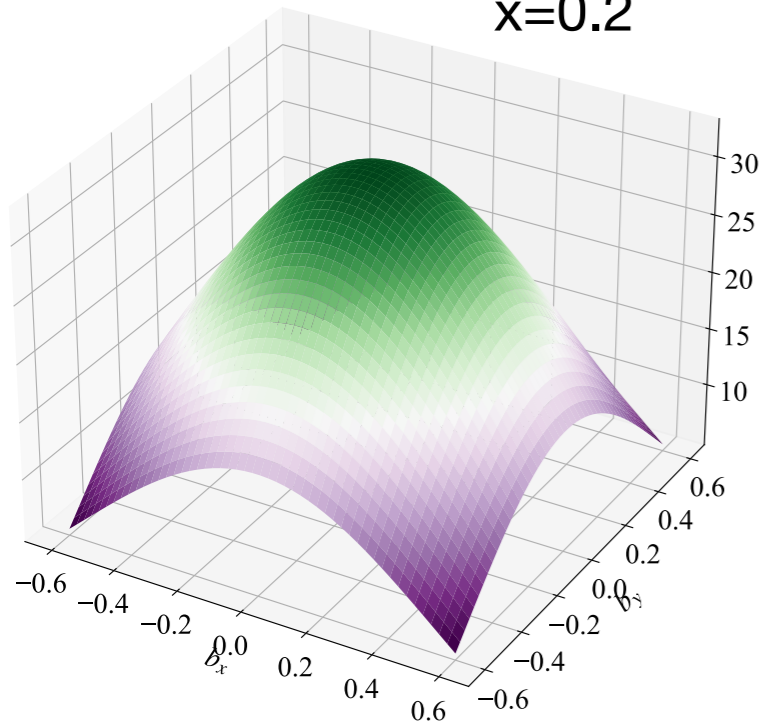
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

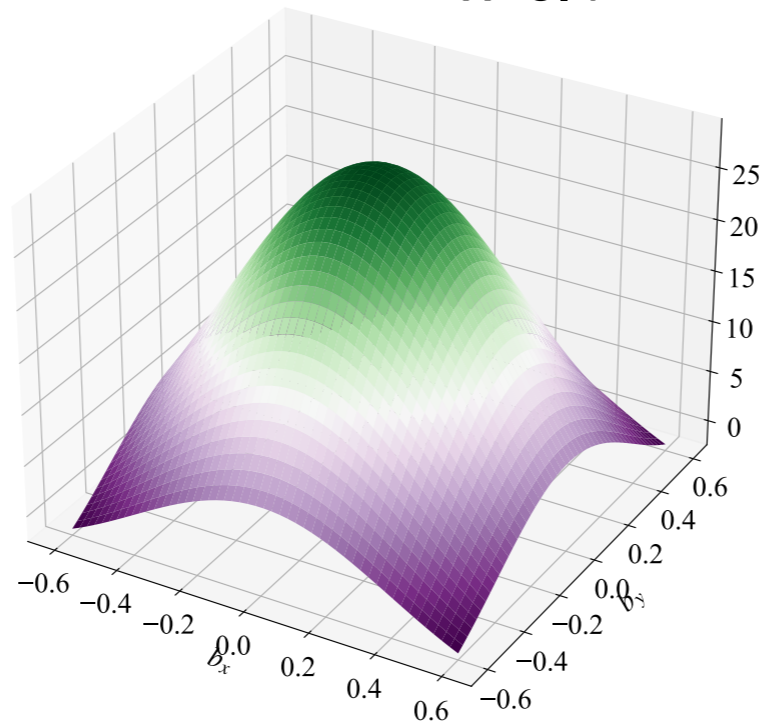
\mathbf{b}_\perp : transverse distance from the (transverse) center of momentum

Impact parameter space $\widetilde{H} + \widetilde{G}_2$

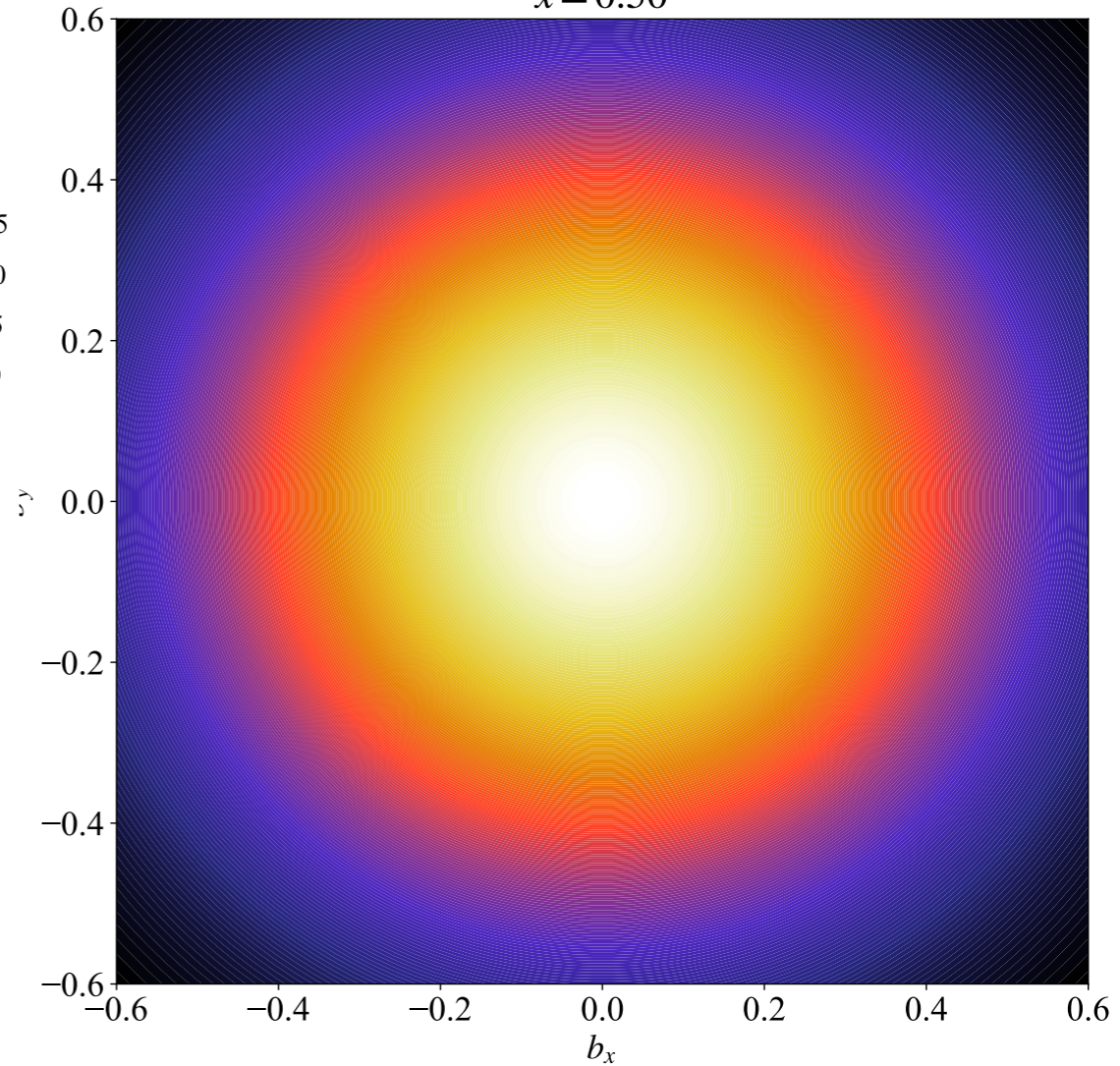
$x=0.2$



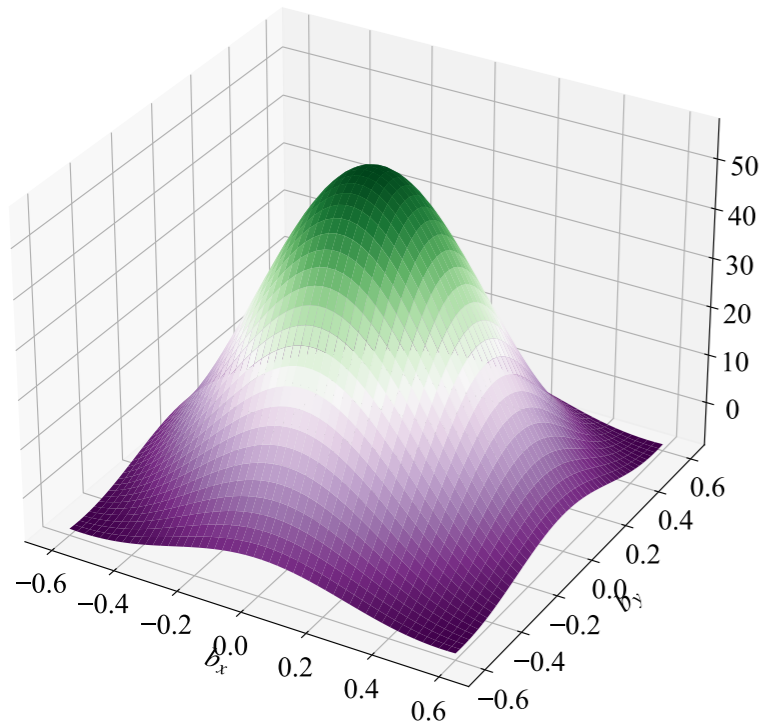
$x=0.4$



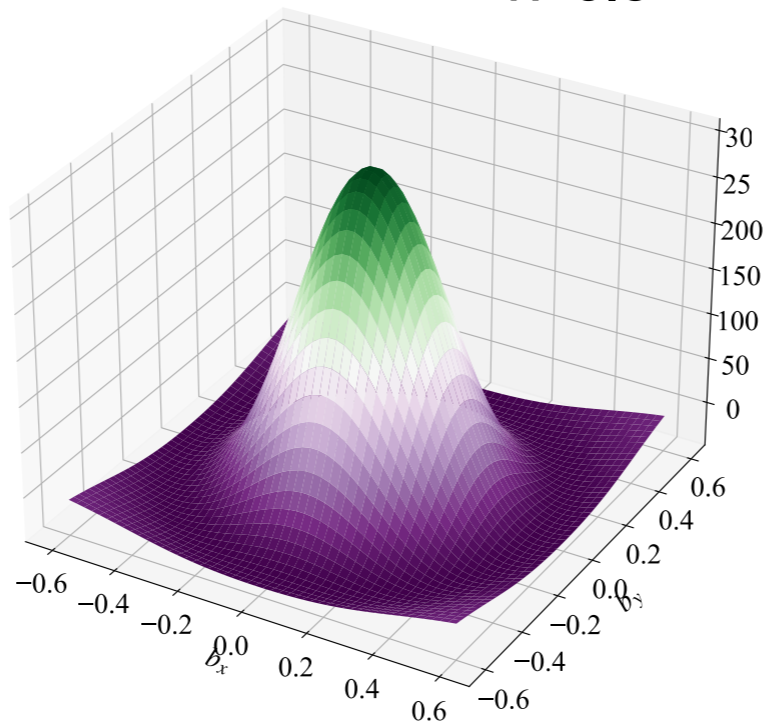
$x=0.50$



$x=0.6$



$x=0.8$



★ GPDs in transverse plane

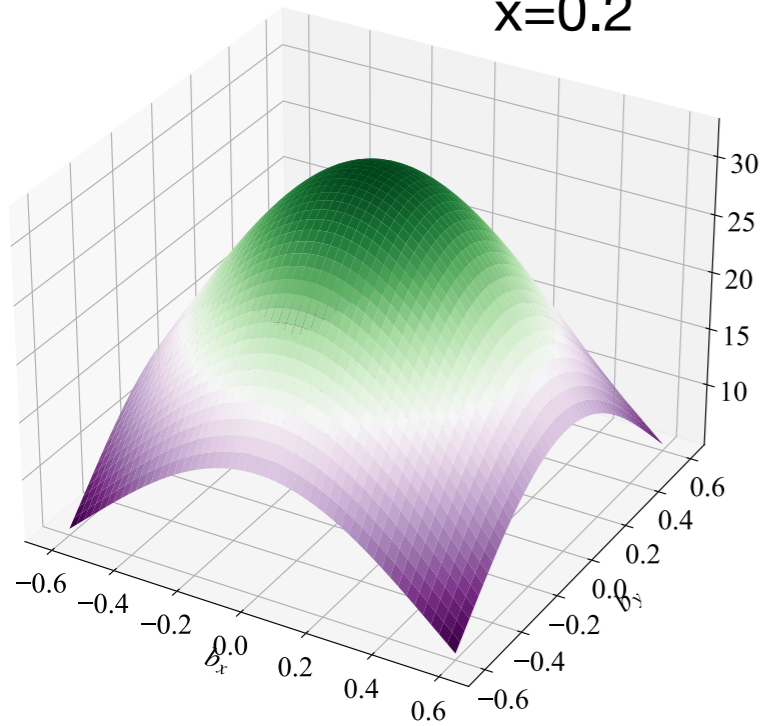
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

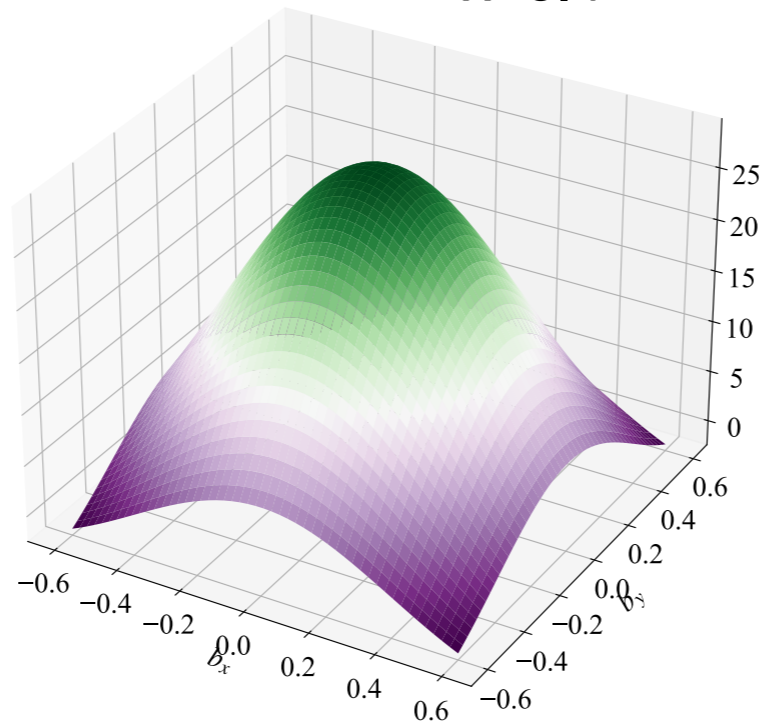
\mathbf{b}_\perp : transverse distance from the
(transverse) center of momentum

Impact parameter space $\widetilde{H} + \widetilde{G}_2$

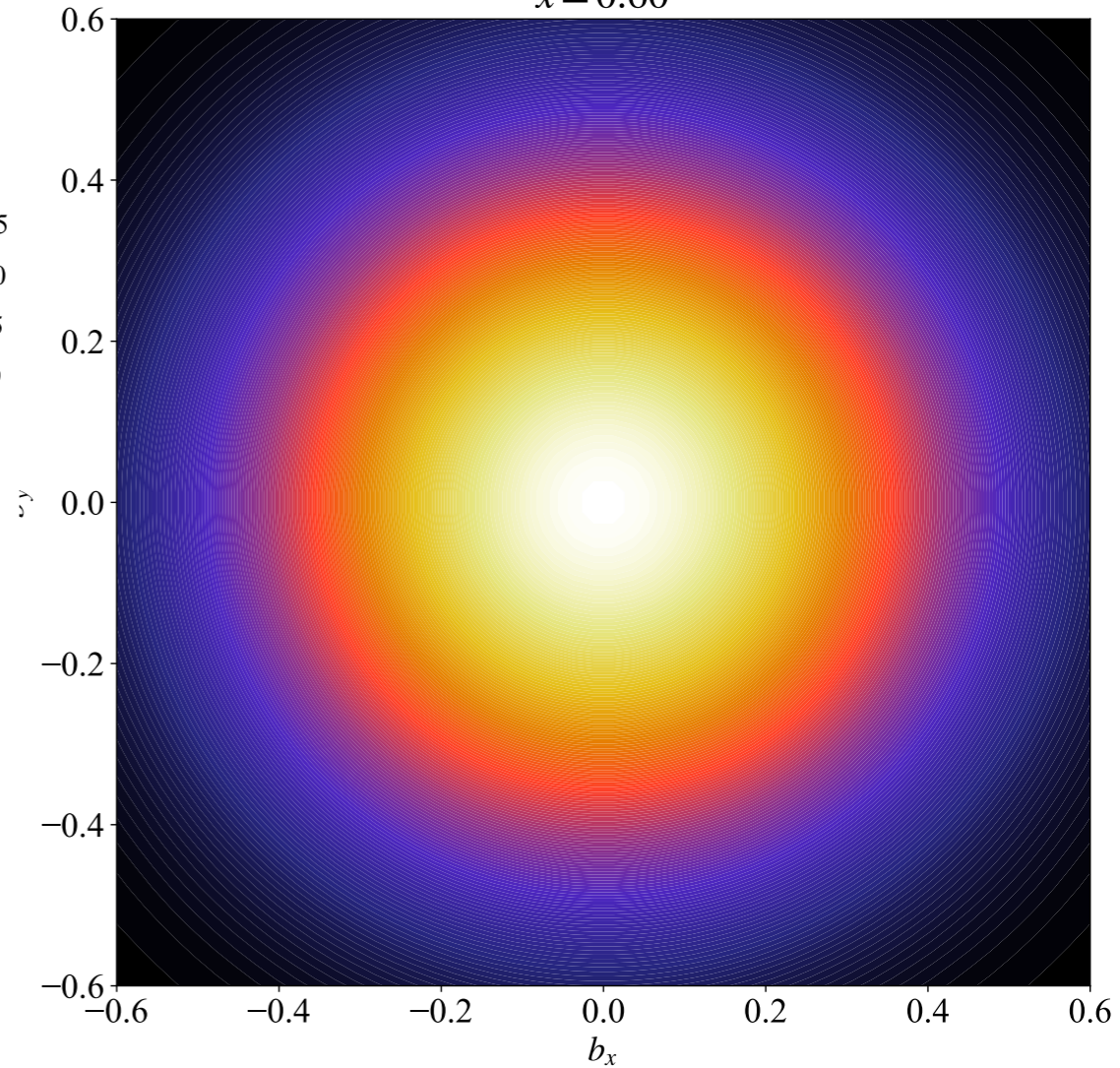
$x=0.2$



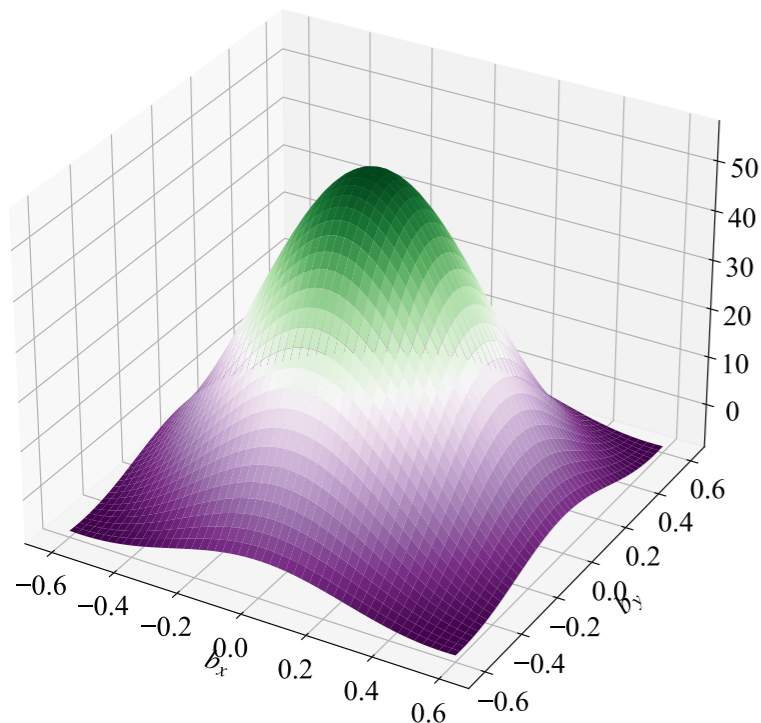
$x=0.4$



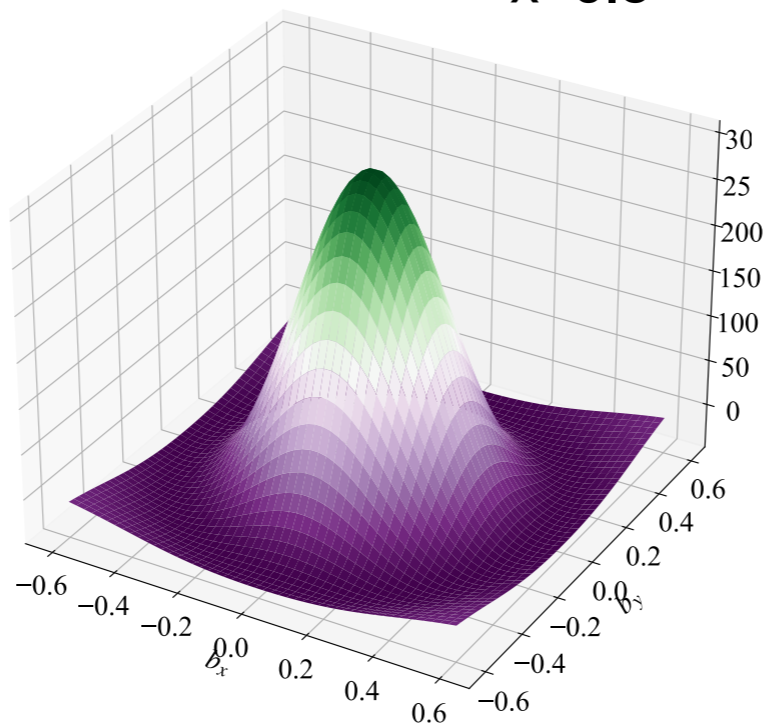
$x=0.60$



$x=0.6$



$x=0.8$



★ GPDs in transverse plane

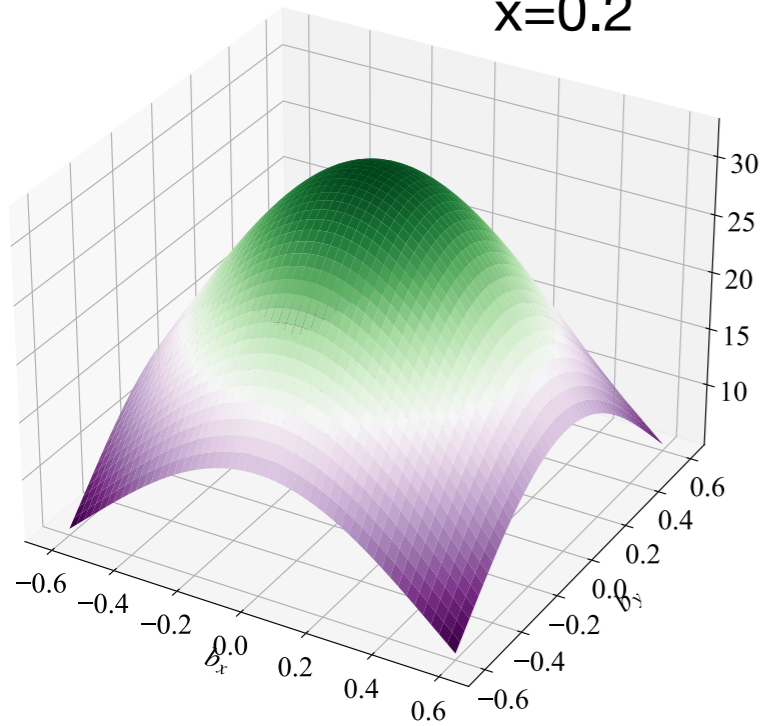
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

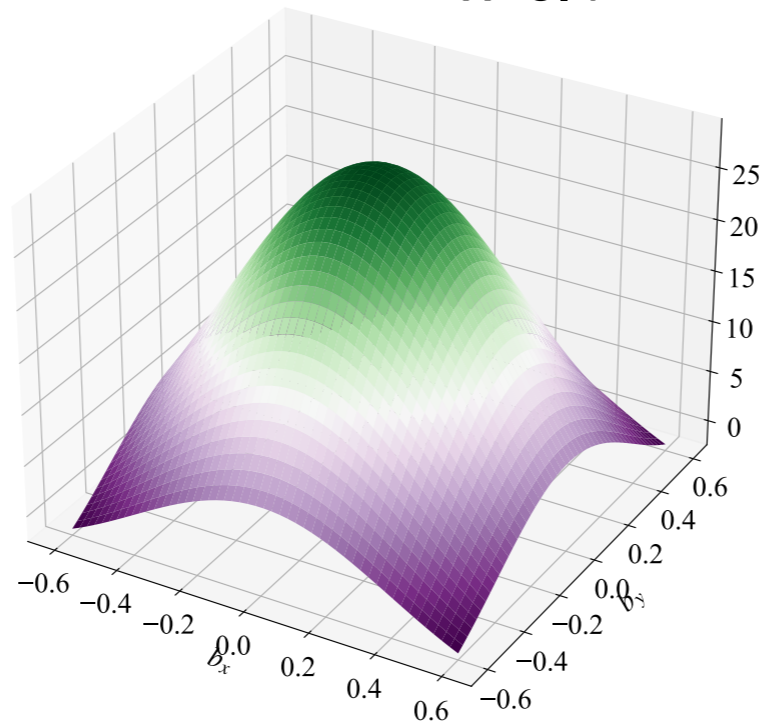
b_\perp : transverse distance from the (transverse) center of momentum

Impact parameter space $\widetilde{H} + \widetilde{G}_2$

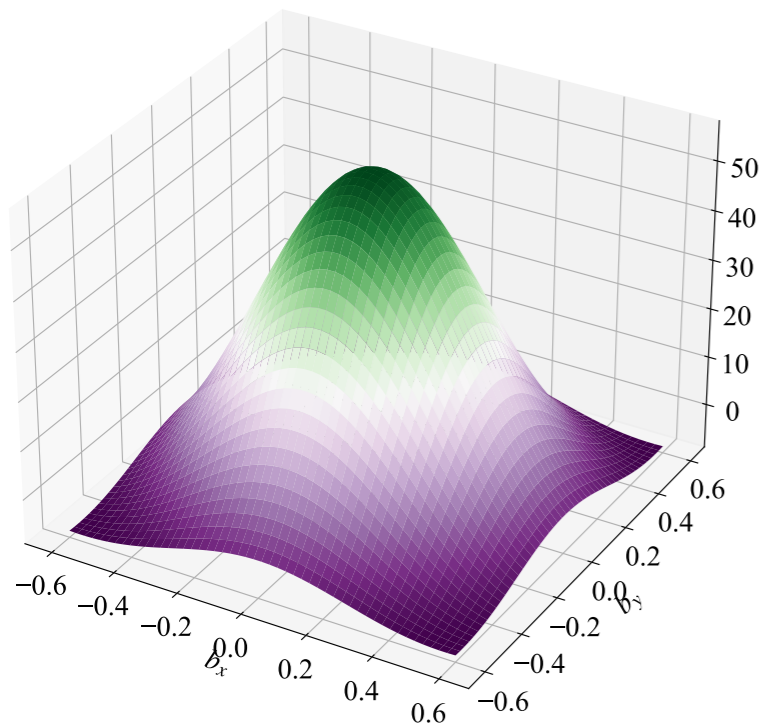
$x=0.2$



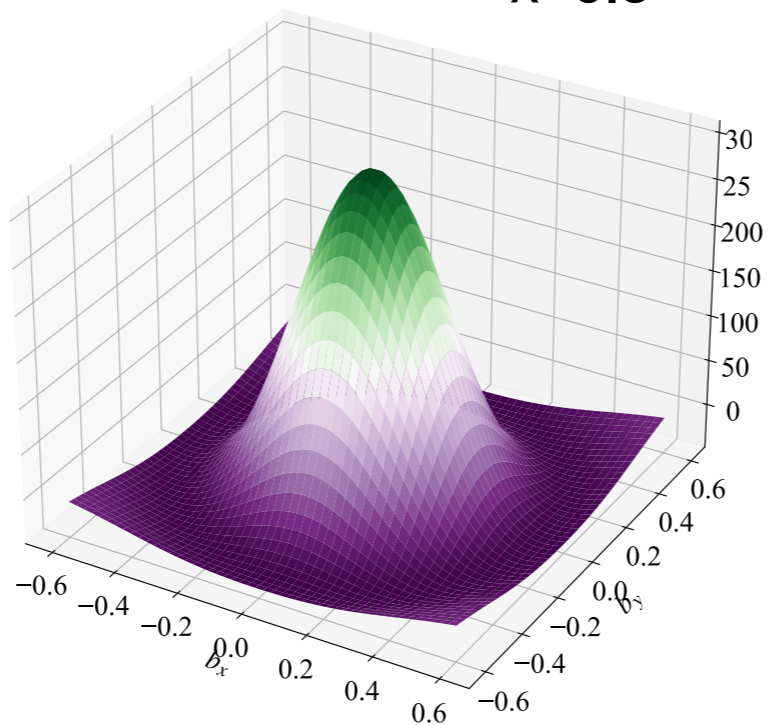
$x=0.4$



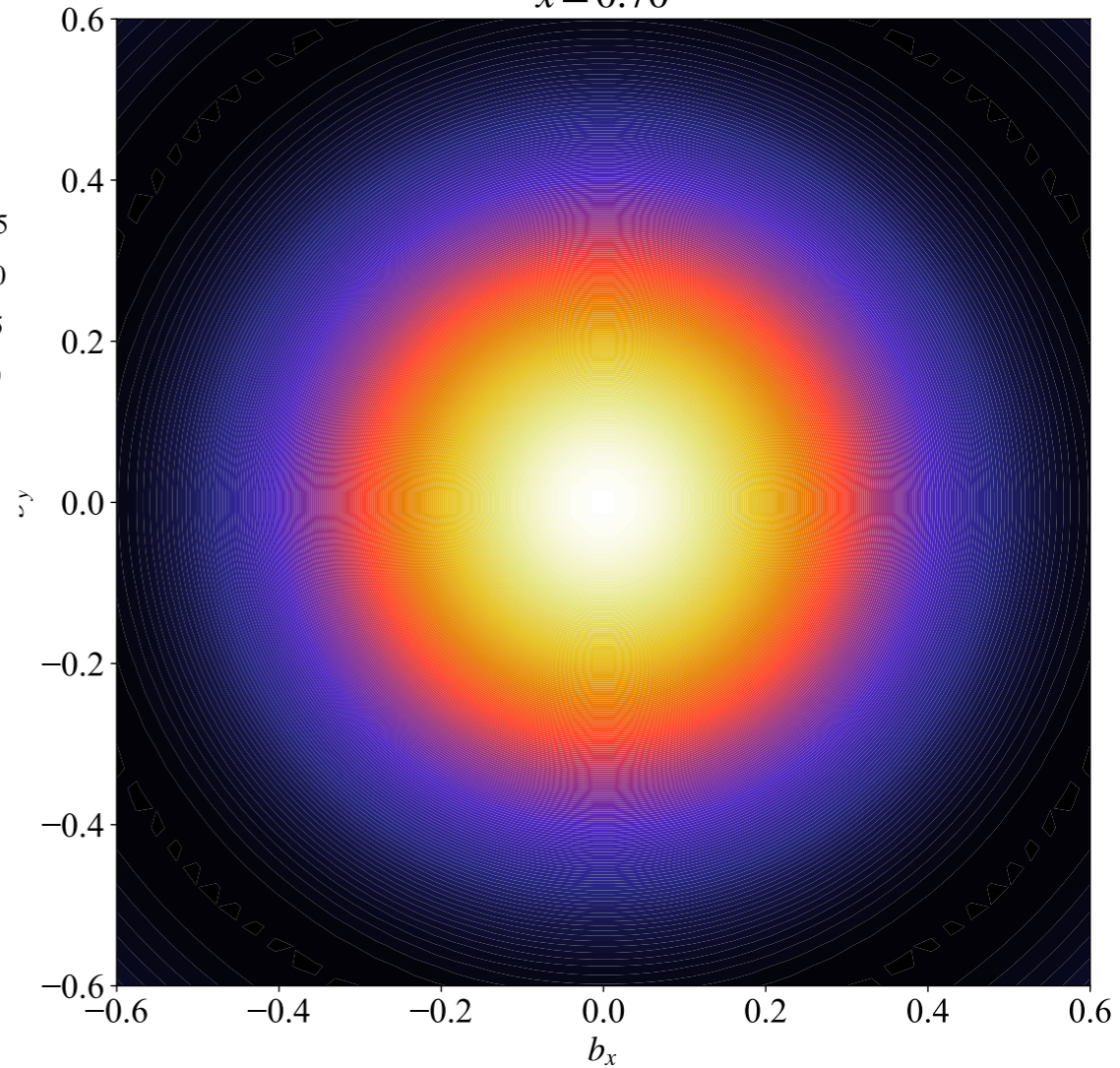
$x=0.6$



$x=0.8$



$x=0.70$



★ GPDs in transverse plane

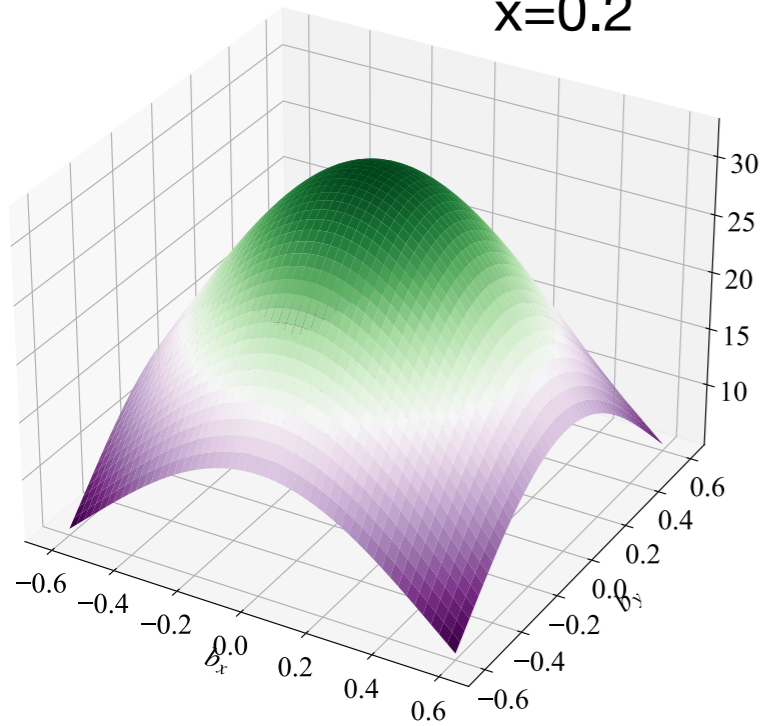
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

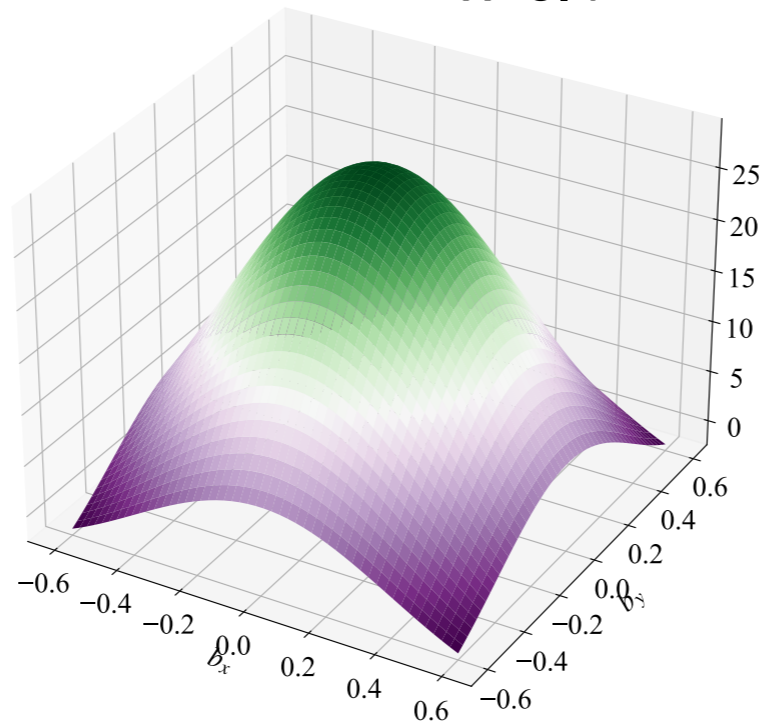
b_\perp : transverse distance from the (transverse) center of momentum

Impact parameter space $\widetilde{H} + \widetilde{G}_2$

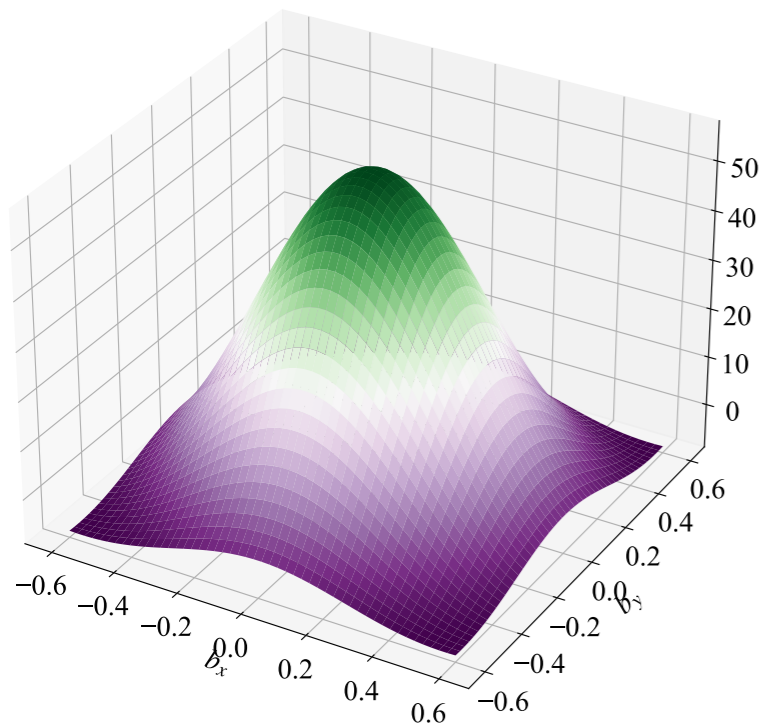
$x=0.2$



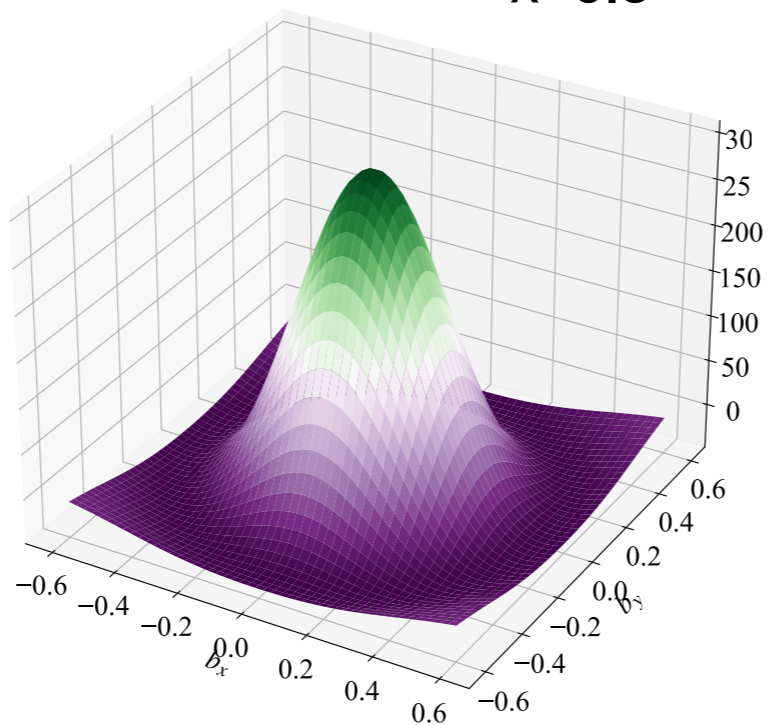
$x=0.4$



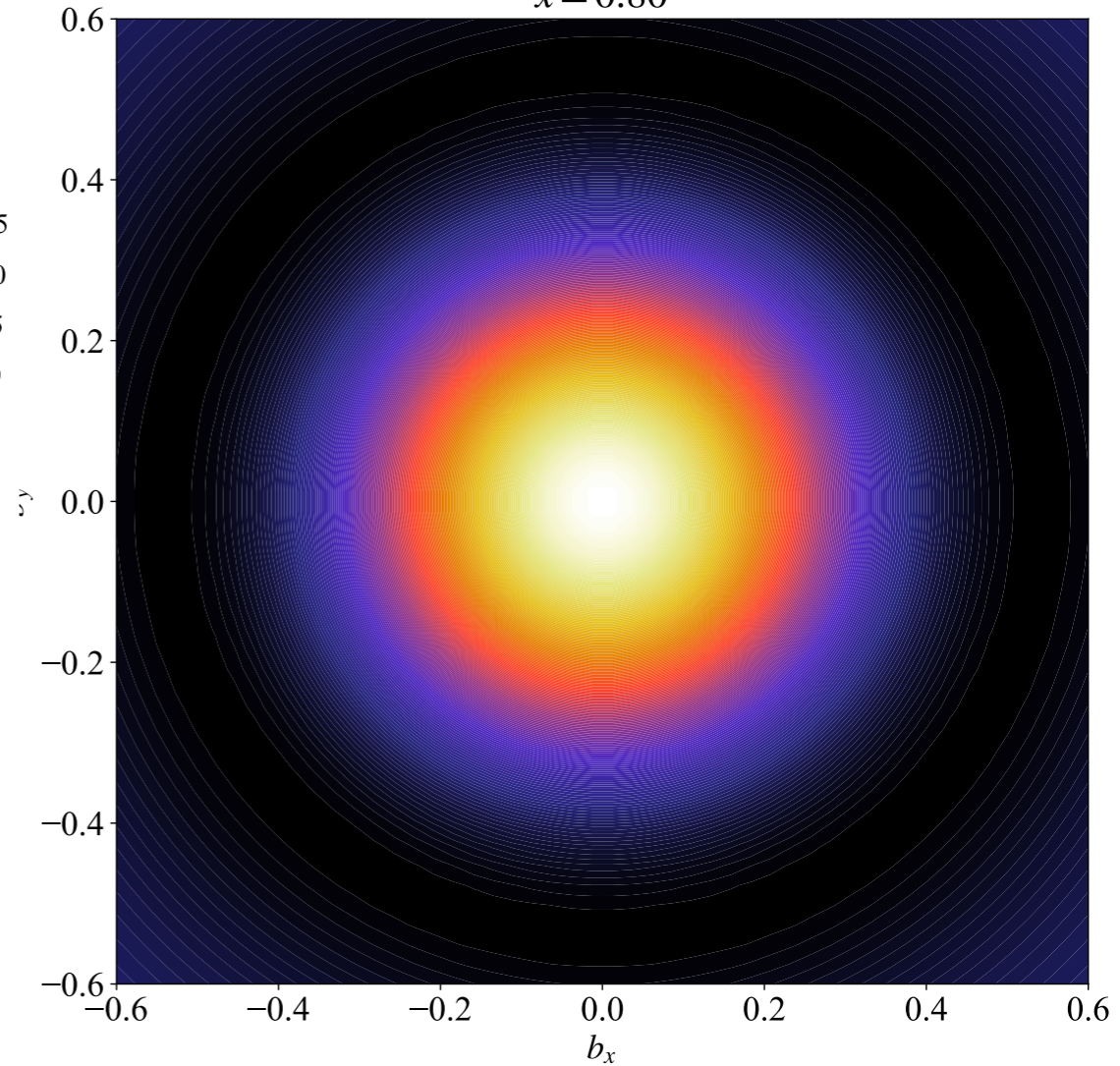
$x=0.6$



$x=0.8$



$x=0.80$



★ GPDs in transverse plane

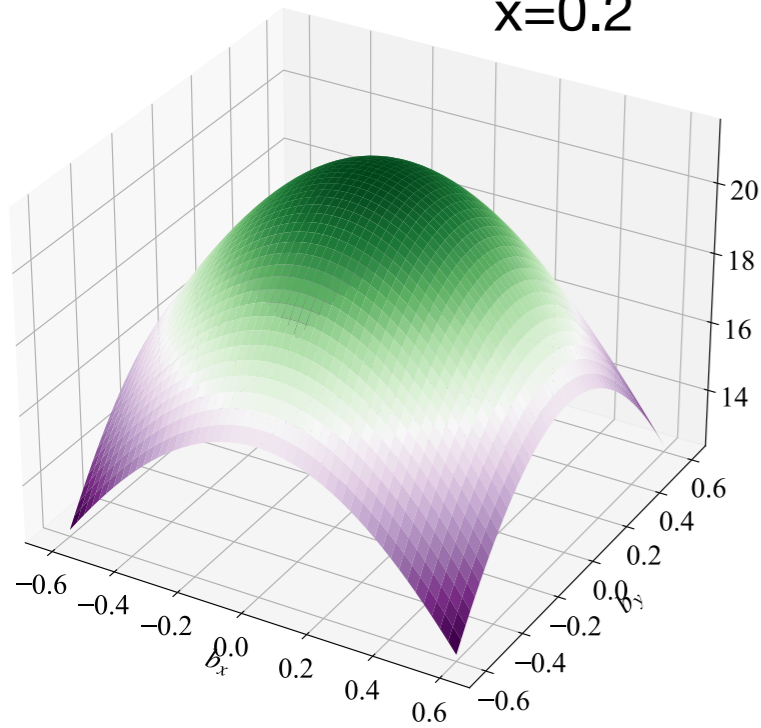
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

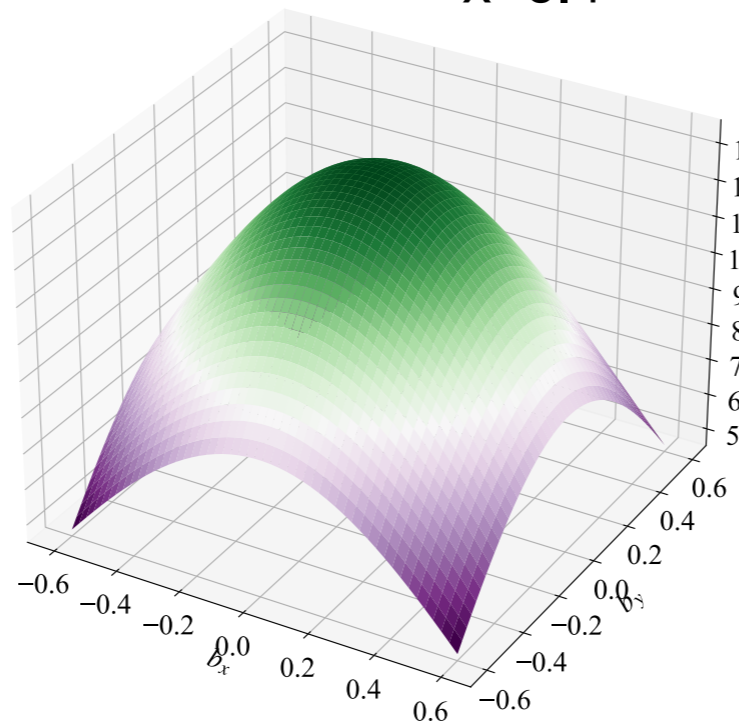
b_\perp : transverse distance from the (transverse) center of momentum

Impact parameter space $\widetilde{E} + \widetilde{G}_1$

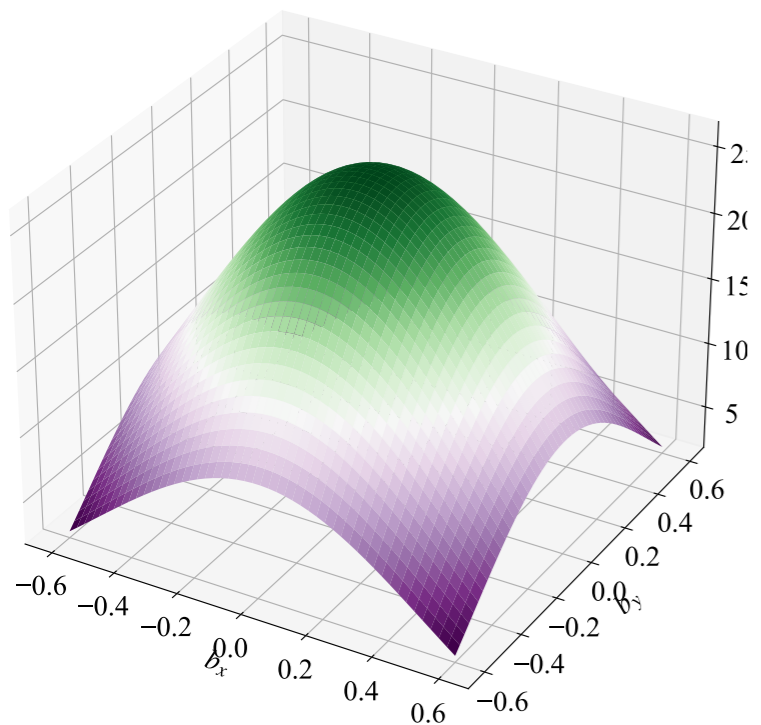
$x=0.2$



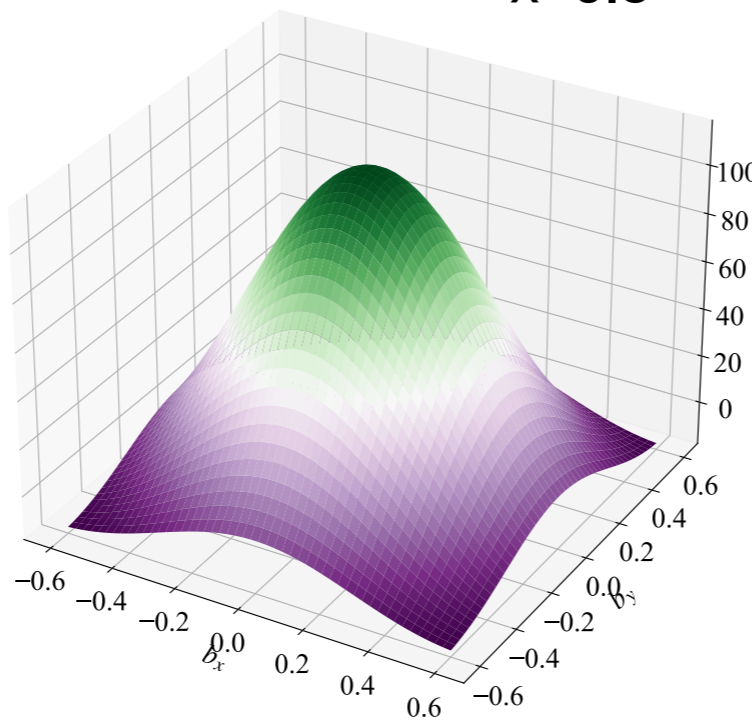
$x=0.4$



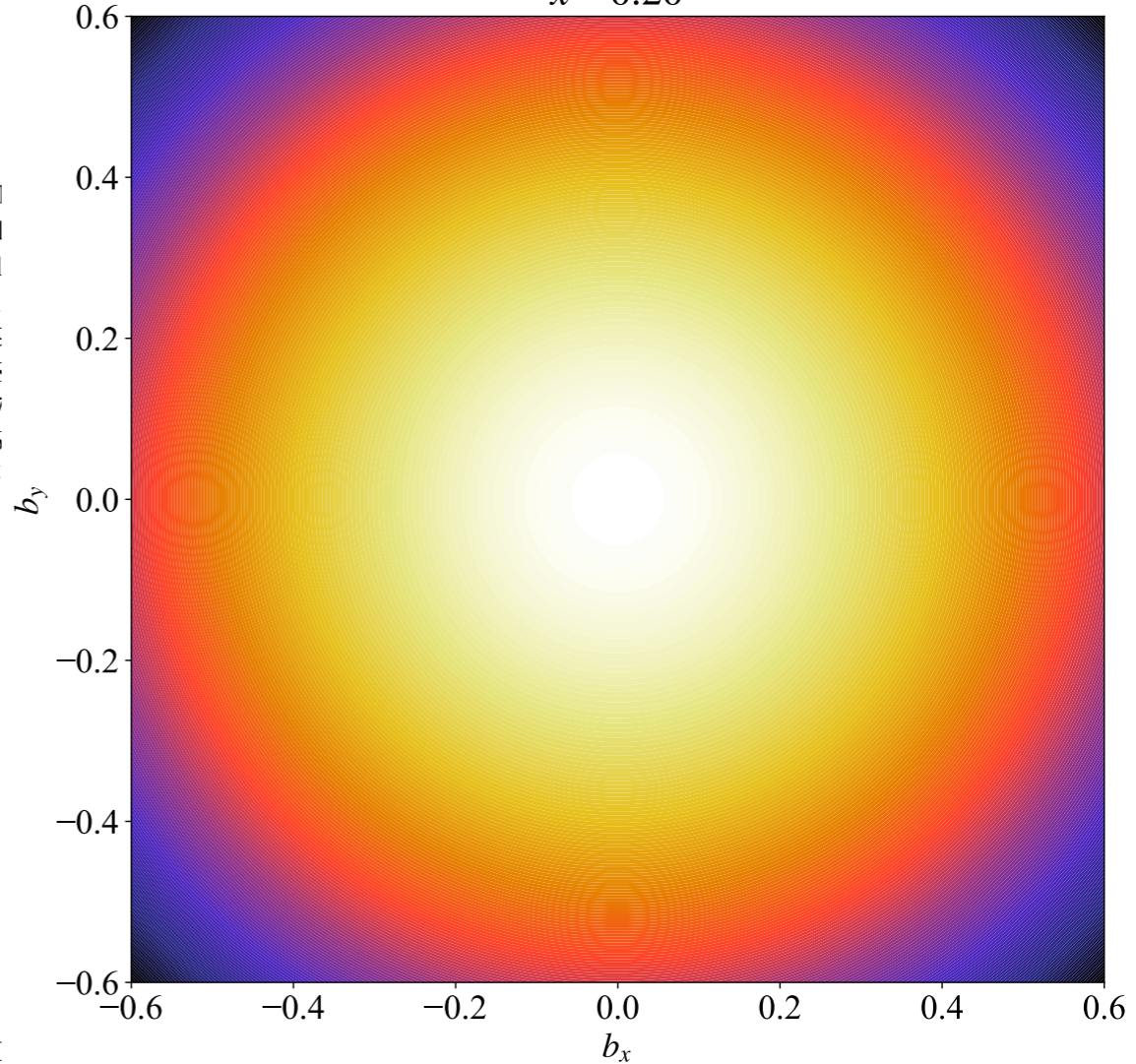
$x=0.6$



$x=0.8$



$x=0.20$



★ GPDs in transverse plane

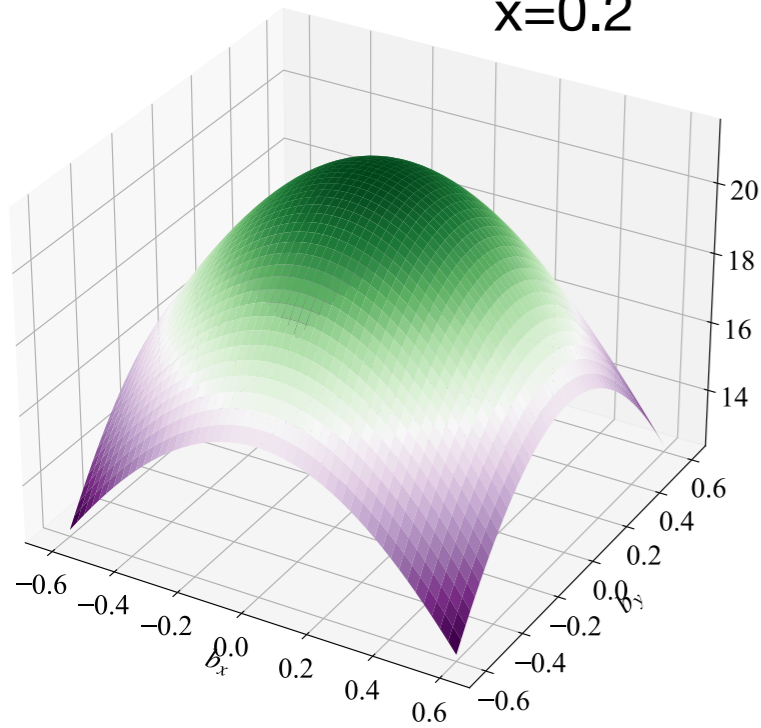
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

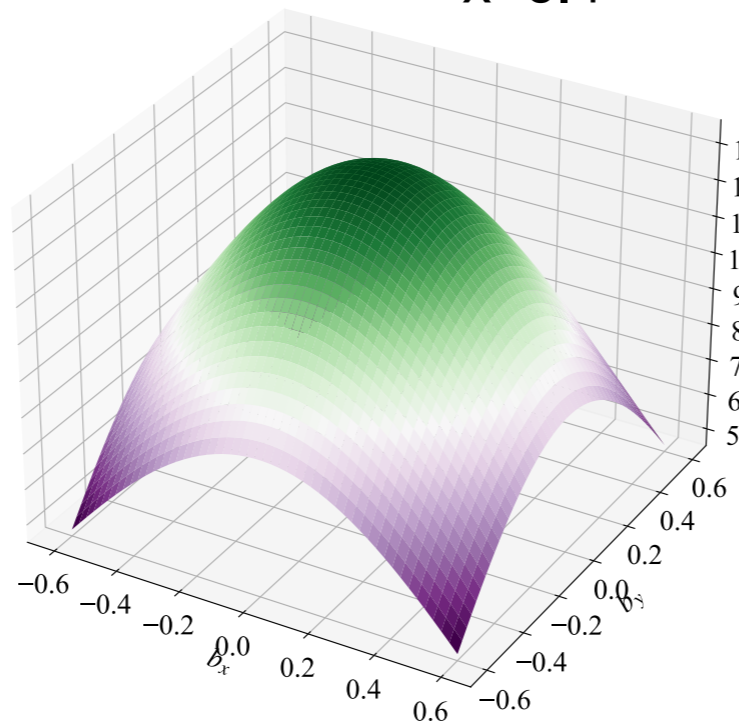
\mathbf{b}_\perp : transverse distance from the (transverse) center of momentum

Impact parameter space $\widetilde{E} + \widetilde{G}_1$

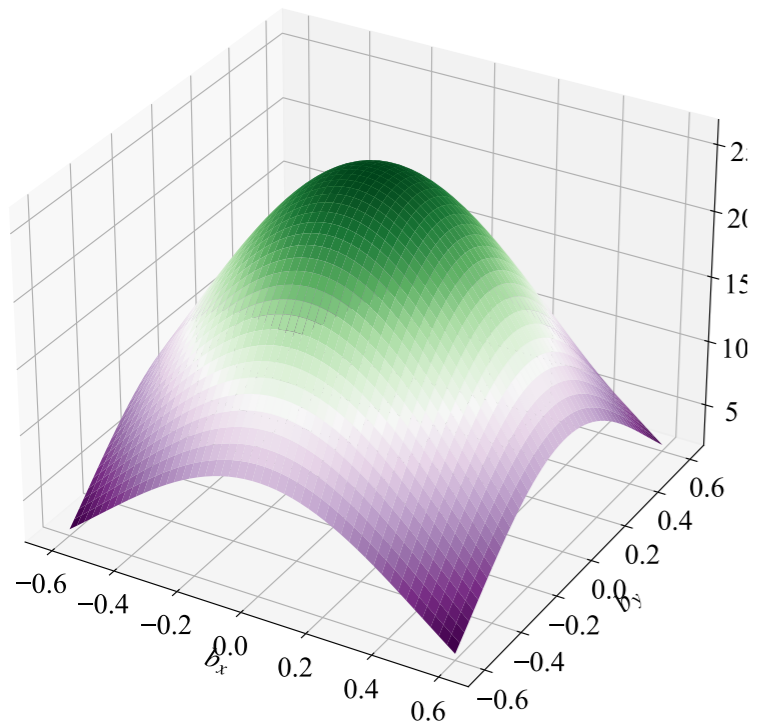
$x=0.2$



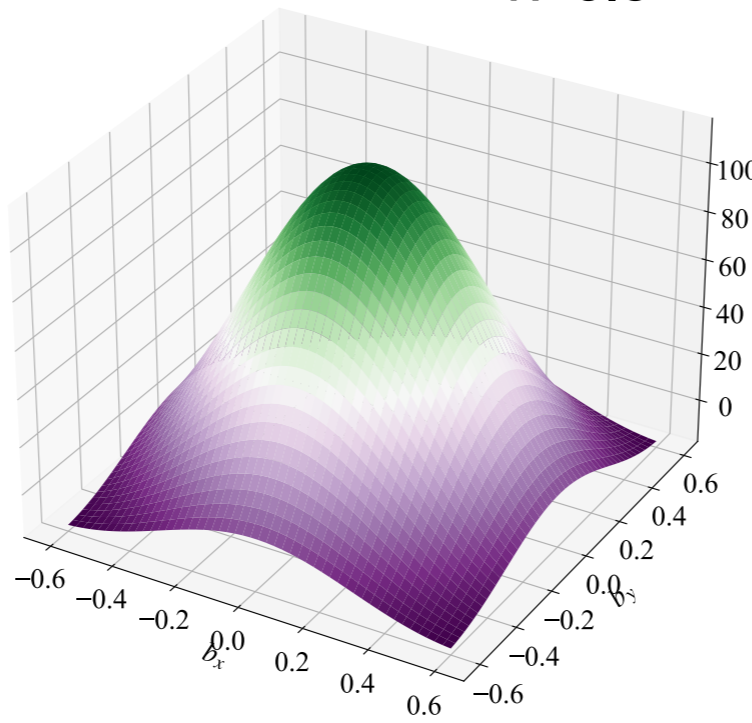
$x=0.4$



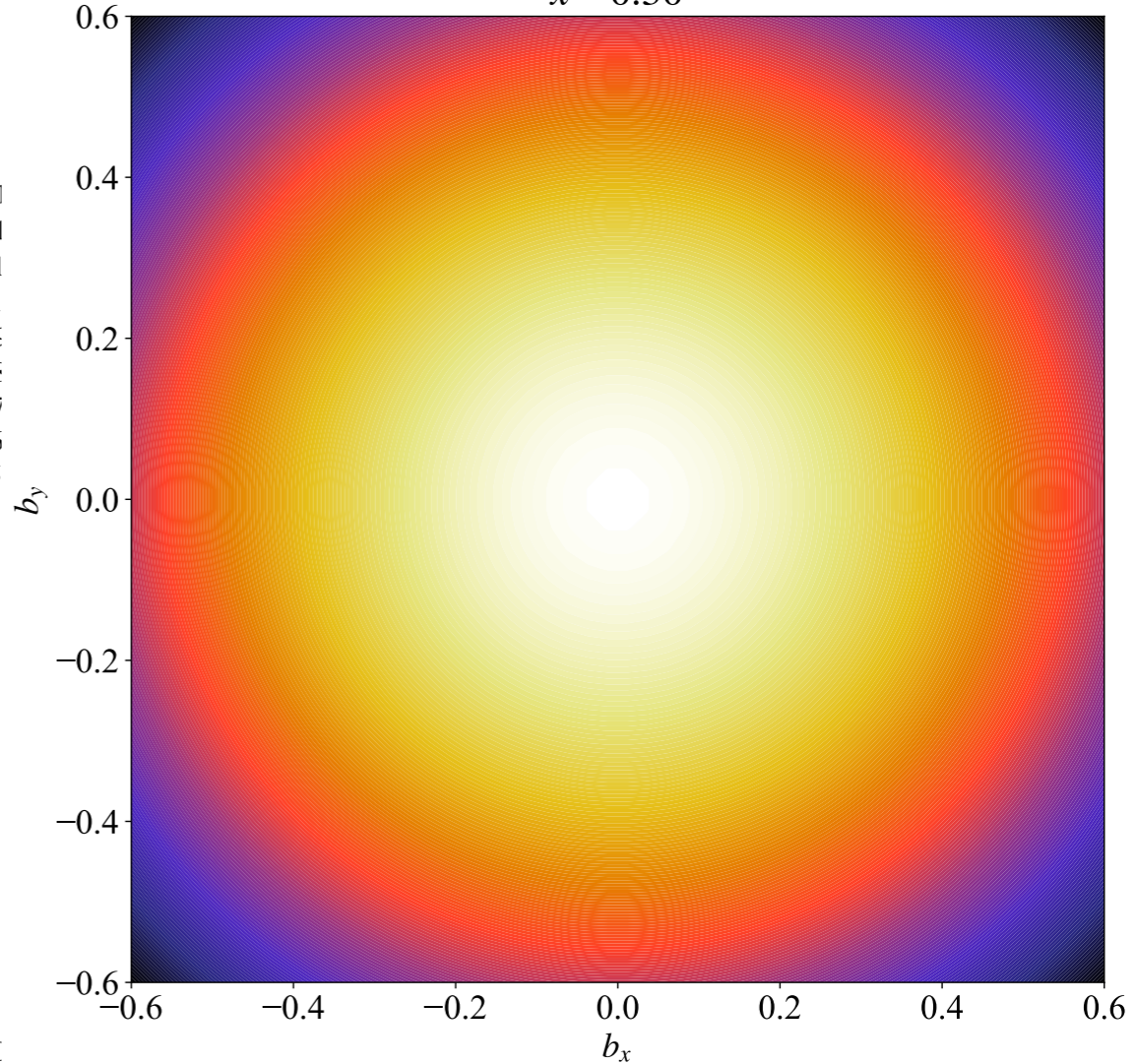
$x=0.6$



$x=0.8$



$x=0.30$



★ GPDs in transverse plane

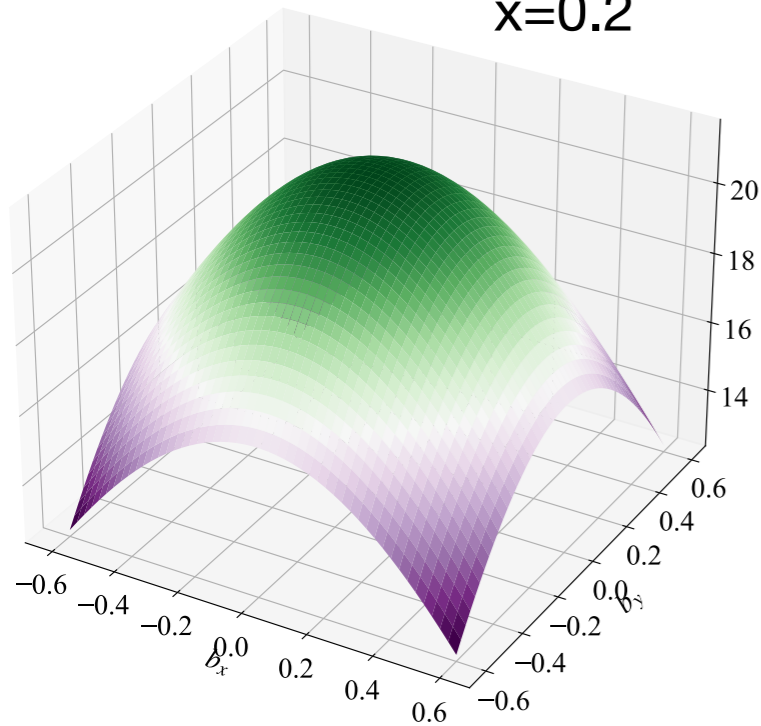
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

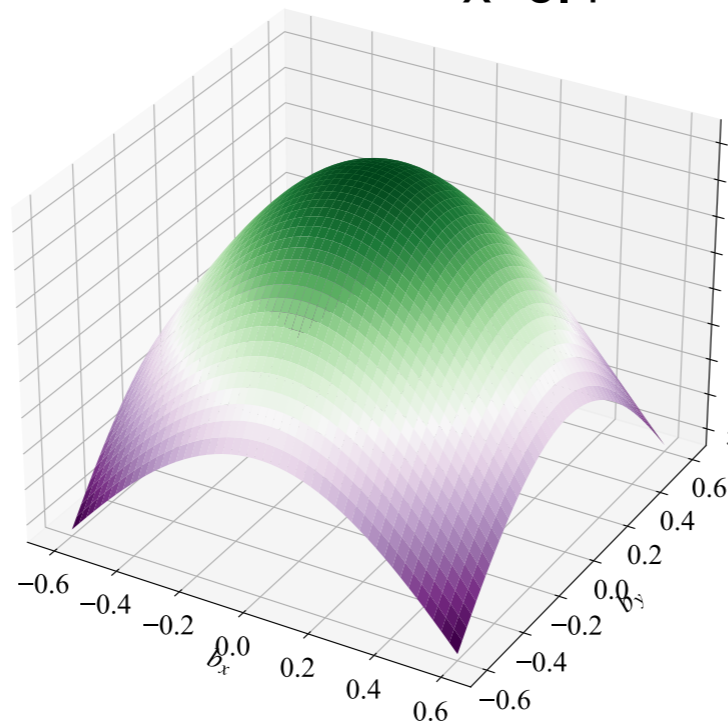
\mathbf{b}_\perp : transverse distance from the (transverse) center of momentum

Impact parameter space $\widetilde{E} + \widetilde{G}_1$

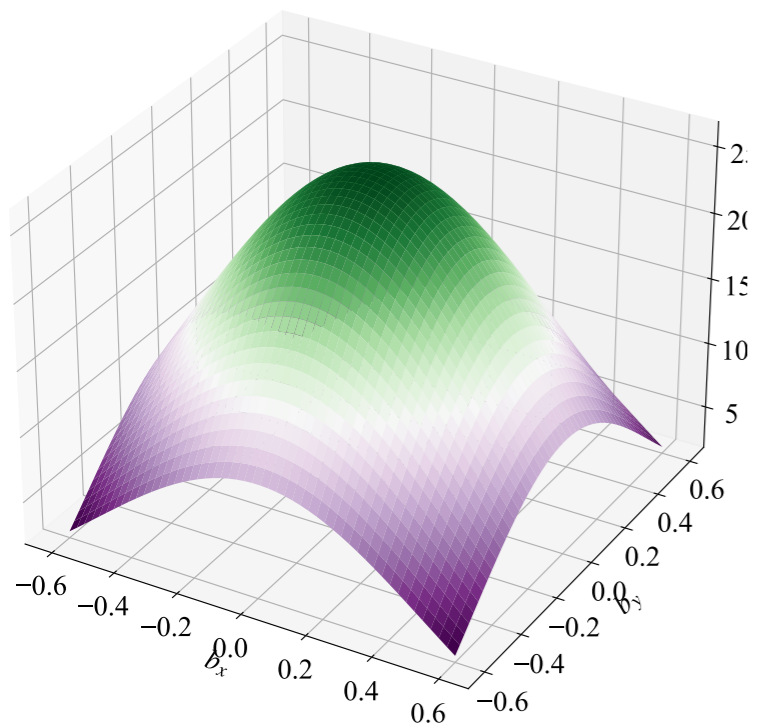
$x=0.2$



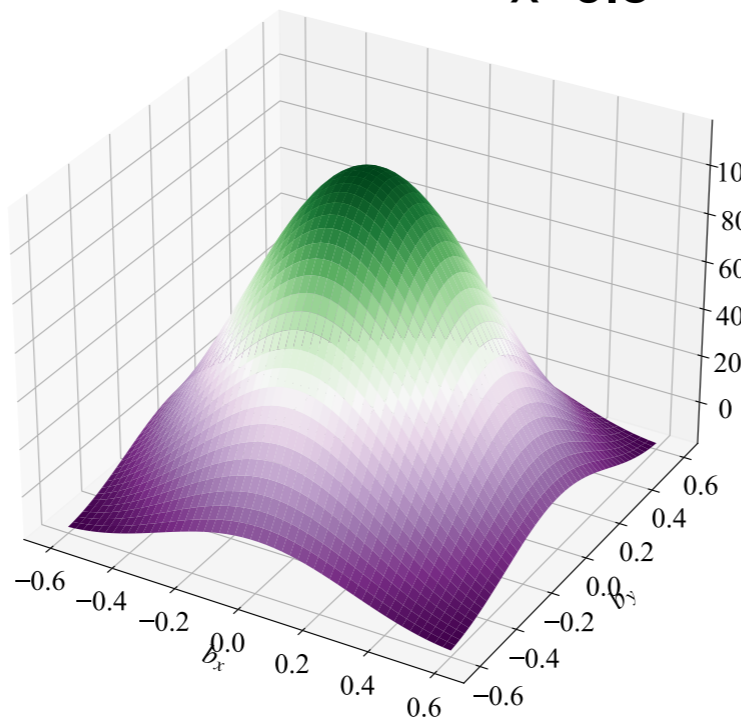
$x=0.4$



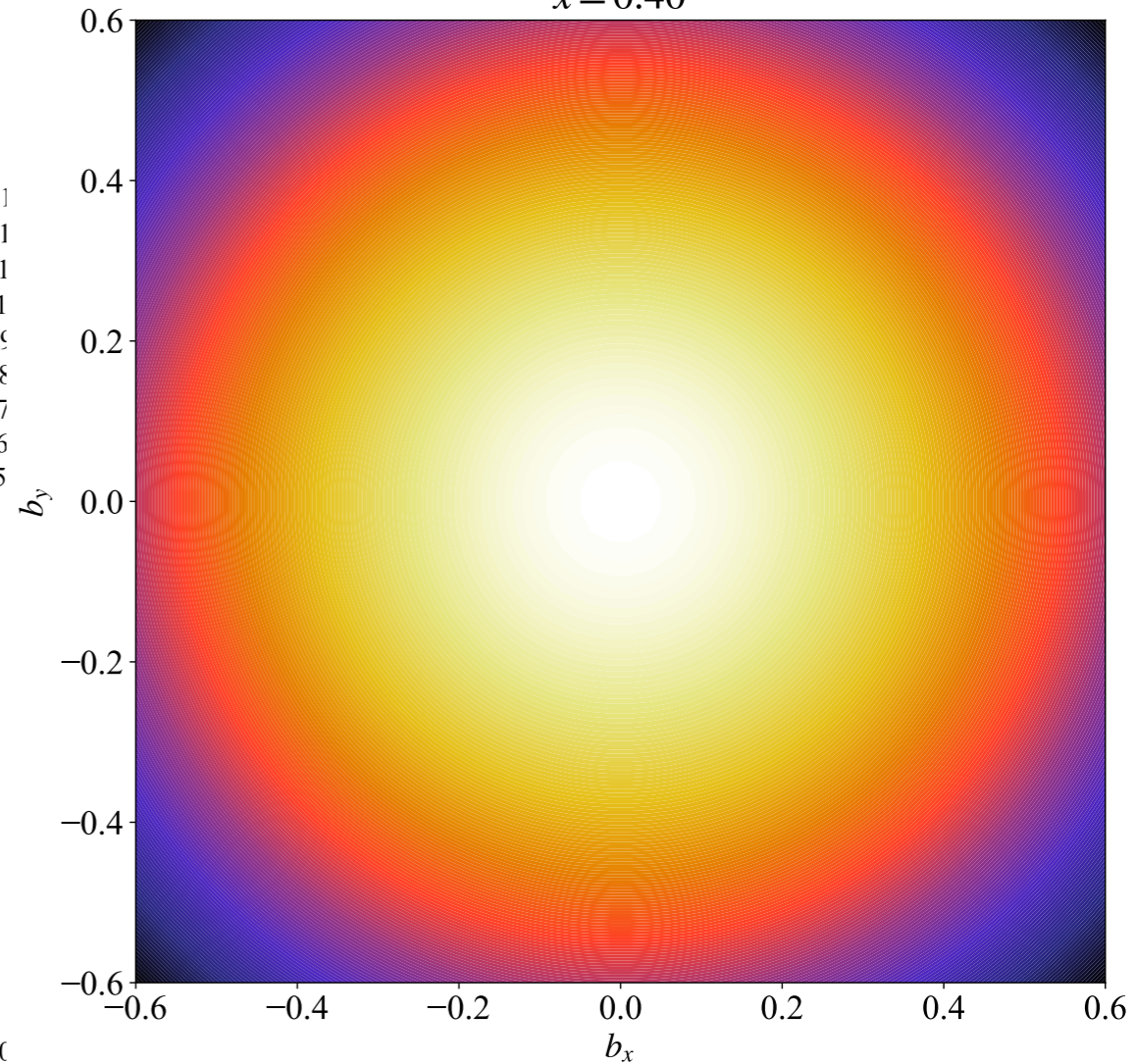
$x=0.6$



$x=0.8$



$x=0.40$



★ GPDs in transverse plane

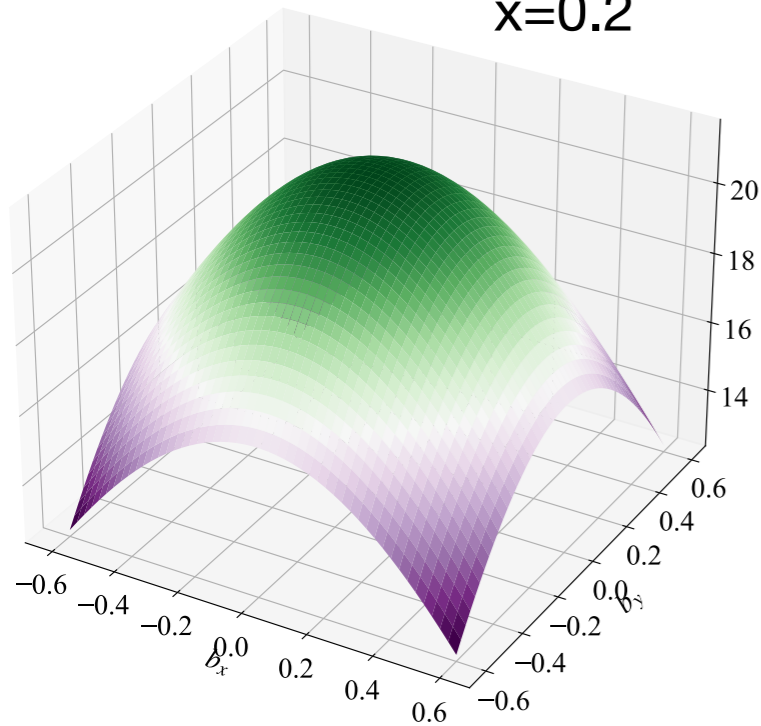
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

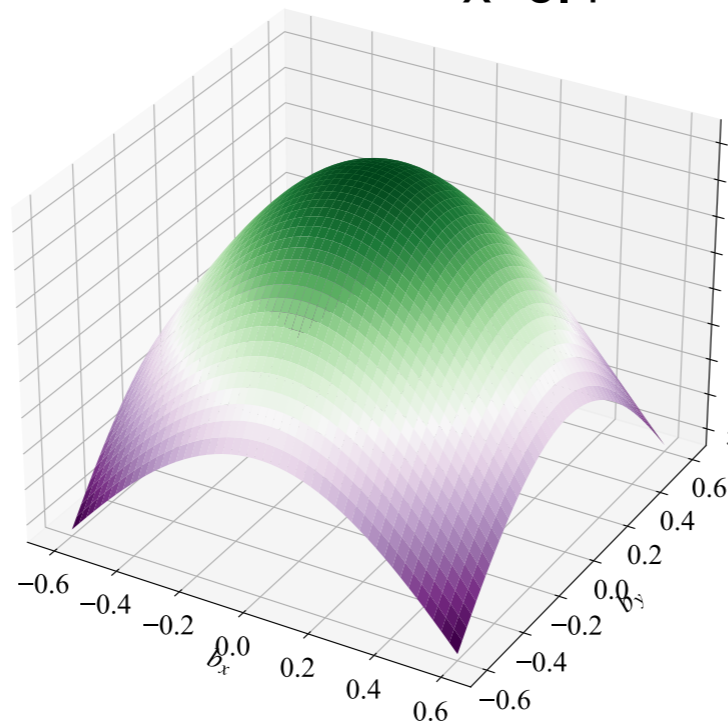
b_\perp : transverse distance from the
(transverse) center of momentum

Impact parameter space $\widetilde{E} + \widetilde{G}_1$

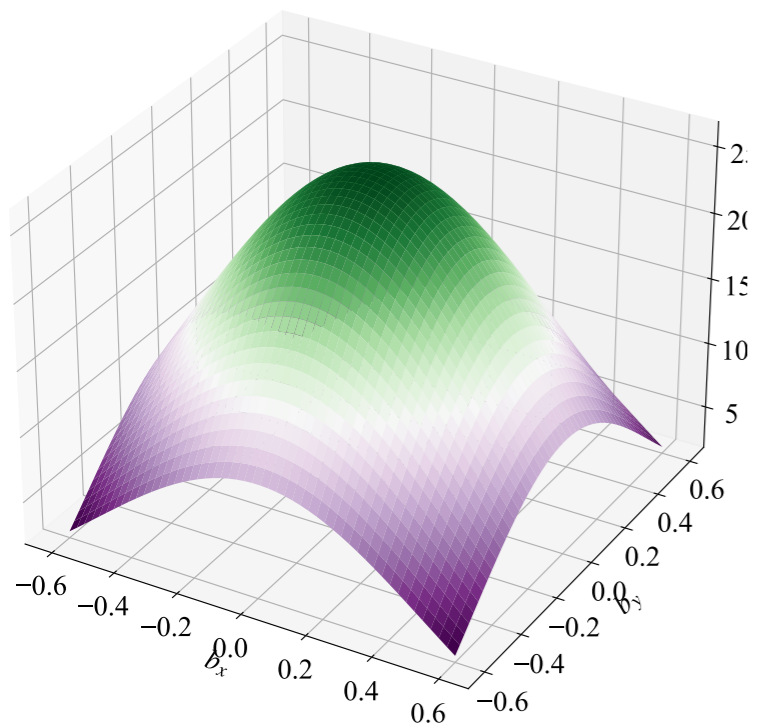
$x=0.2$



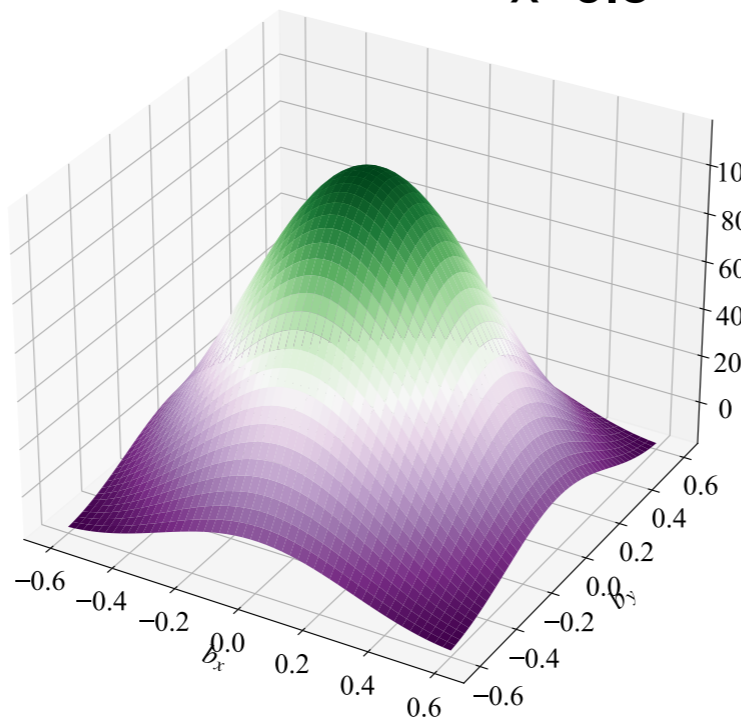
$x=0.4$



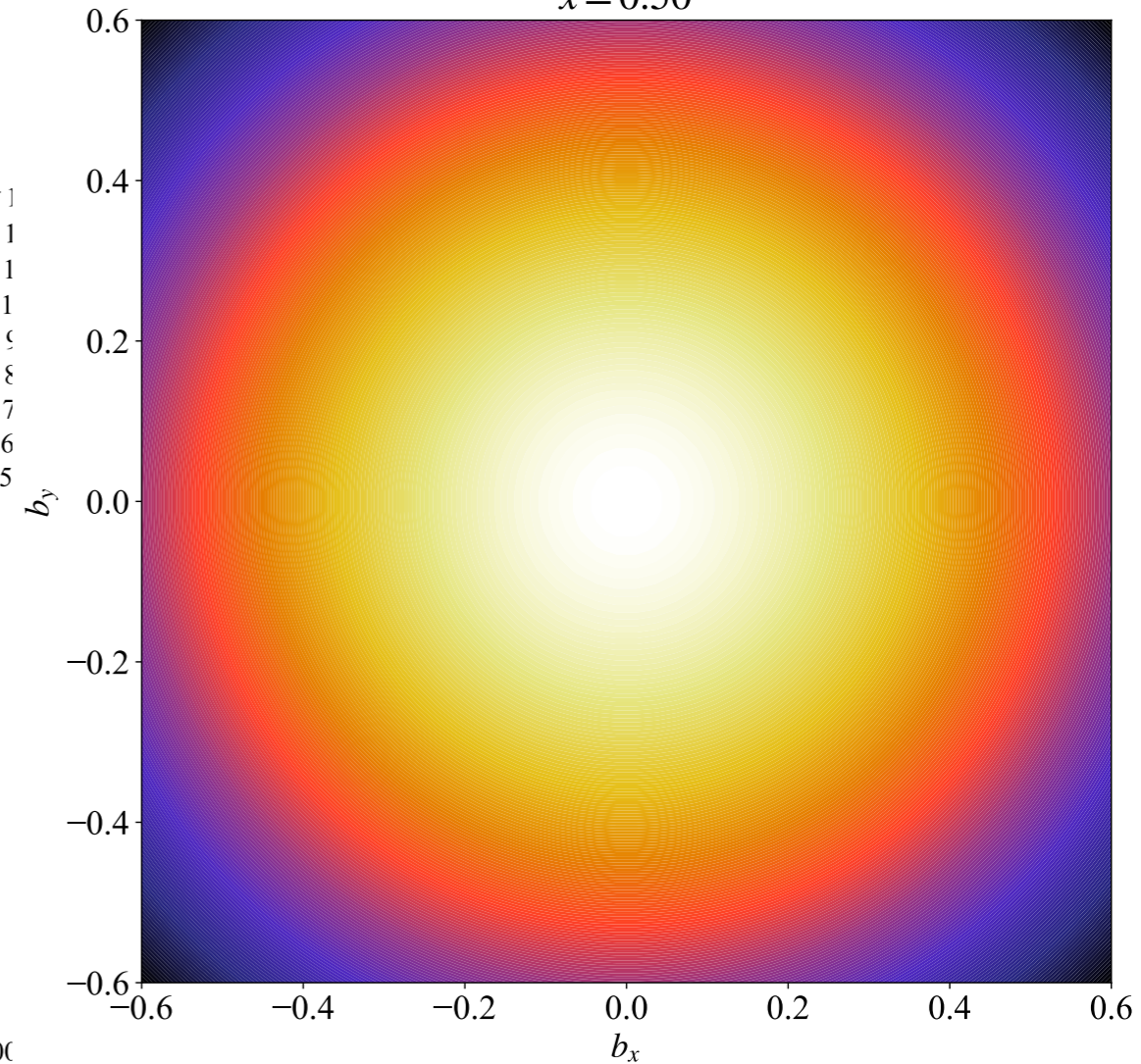
$x=0.6$



$x=0.8$



$x=0.50$



★ GPDs in transverse plane

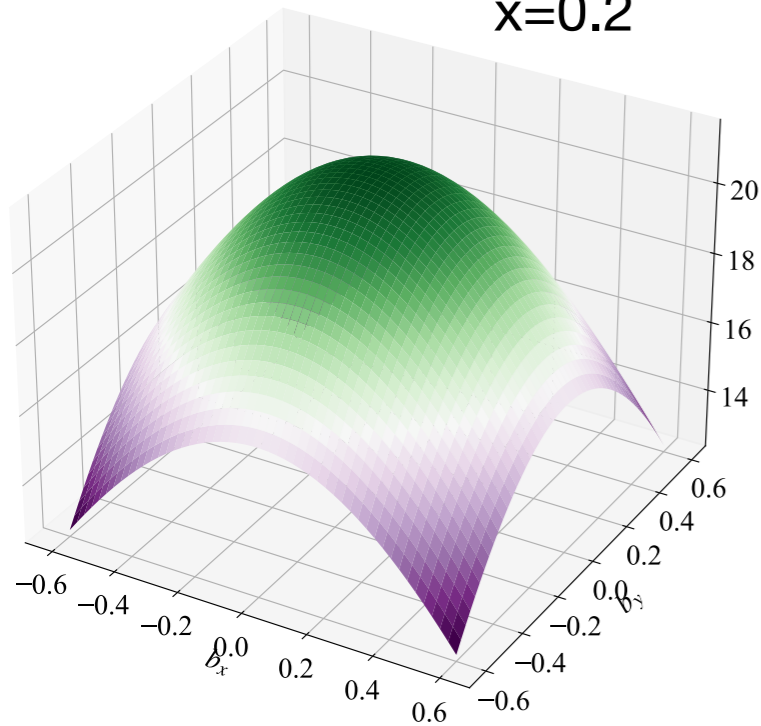
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

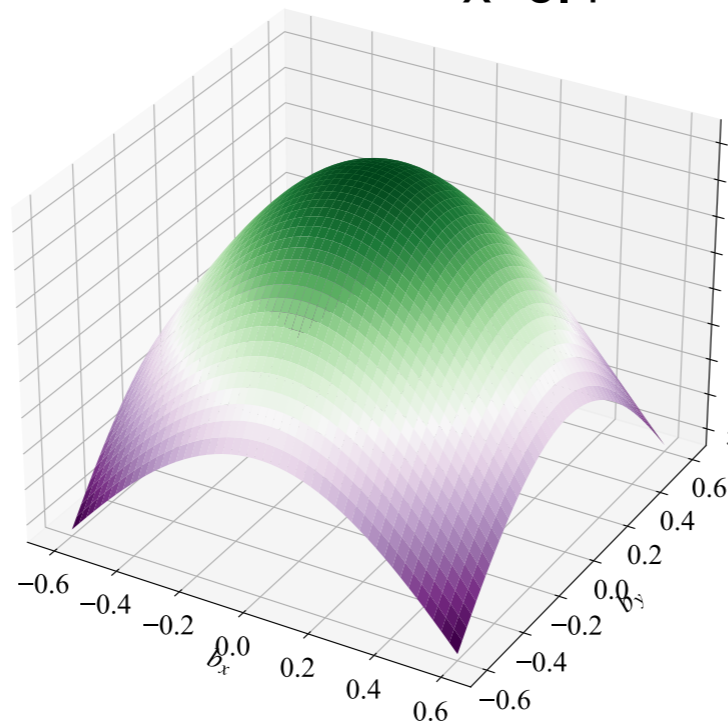
\mathbf{b}_\perp : transverse distance from the
(transverse) center of momentum

Impact parameter space $\widetilde{E} + \widetilde{G}_1$

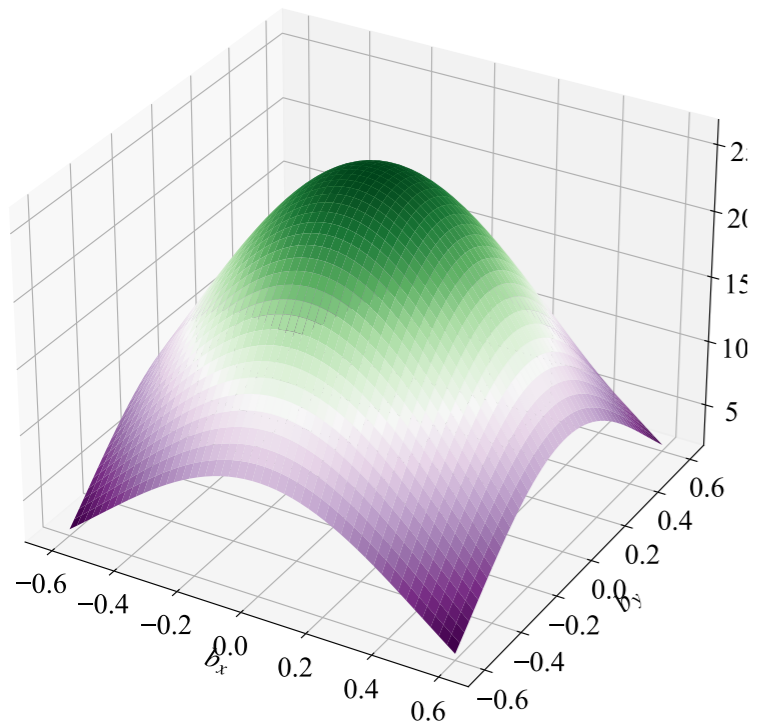
$x=0.2$



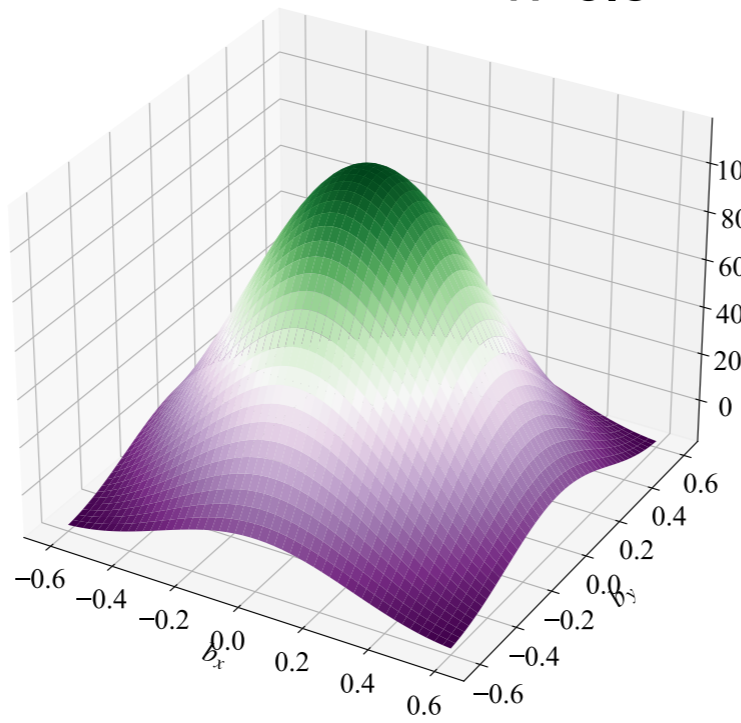
$x=0.4$



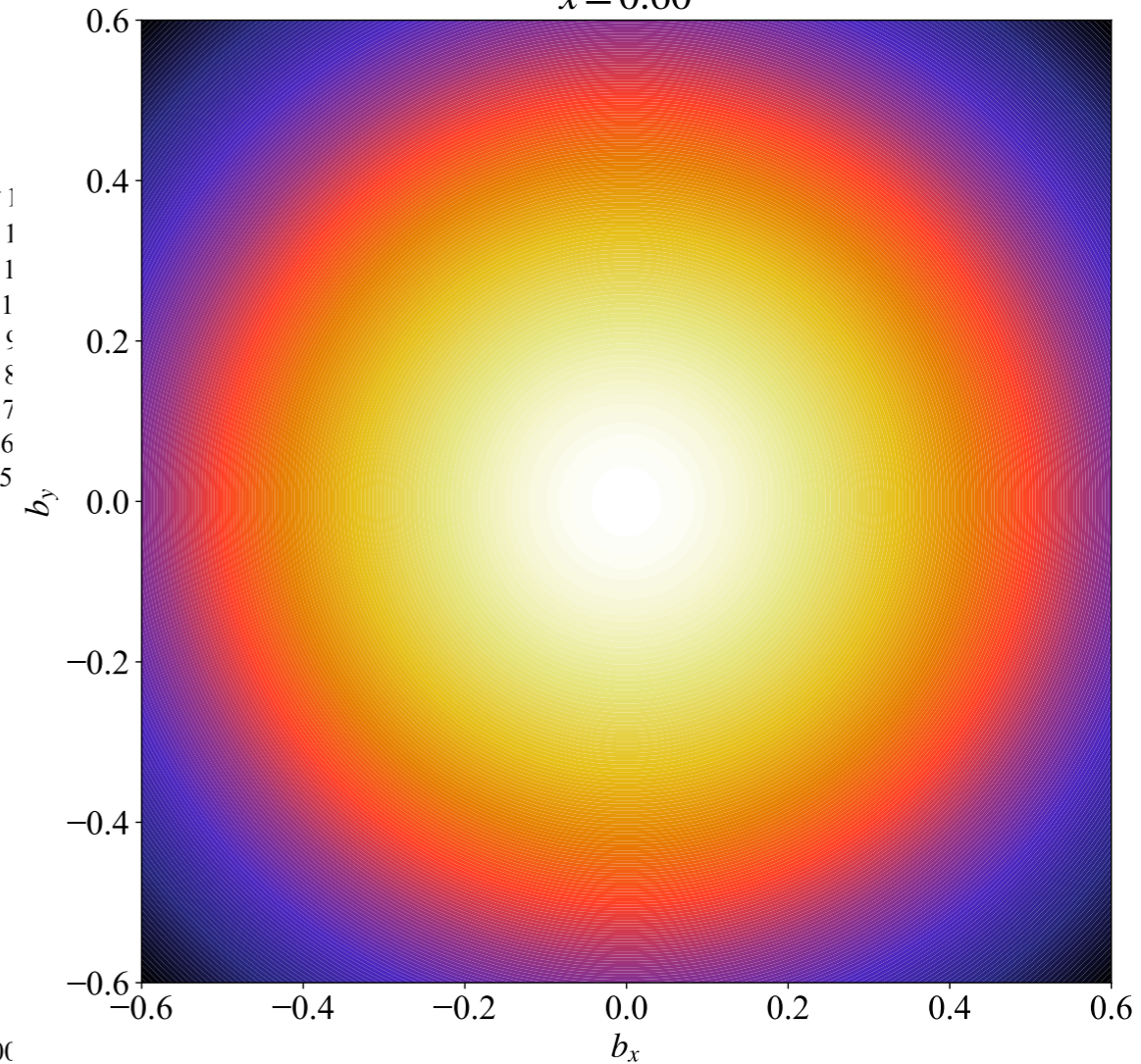
$x=0.6$



$x=0.8$



$x=0.60$



★ GPDs in transverse plane

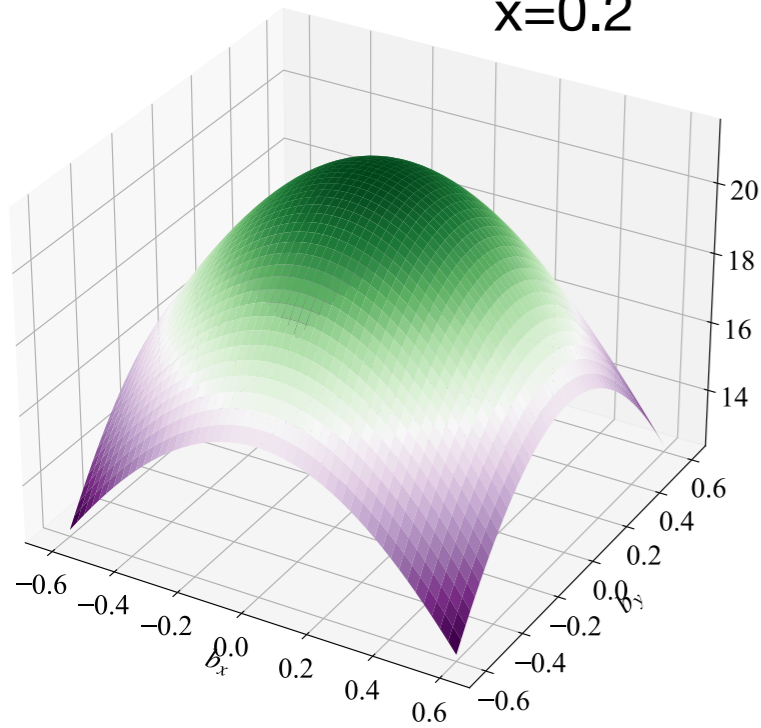
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

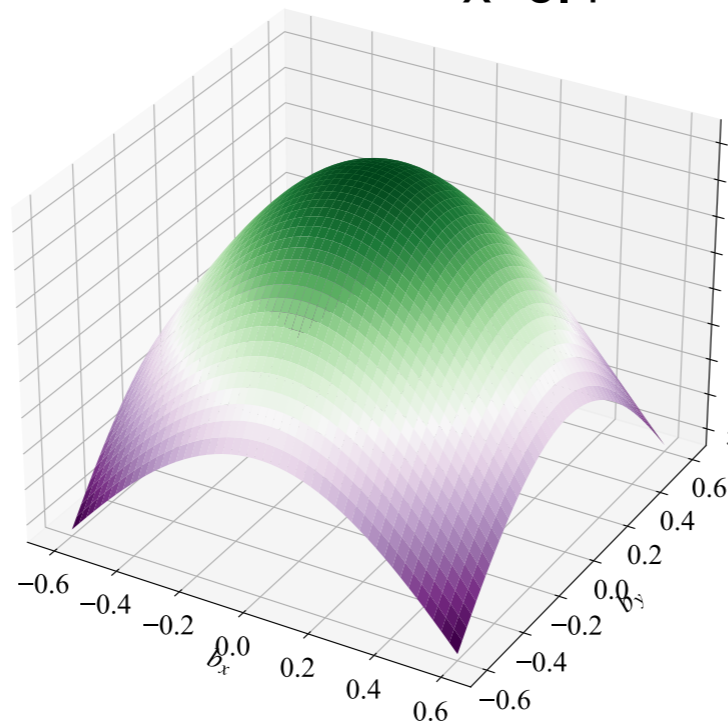
b_\perp : transverse distance from the
(transverse) center of momentum

Impact parameter space $\widetilde{E} + \widetilde{G}_1$

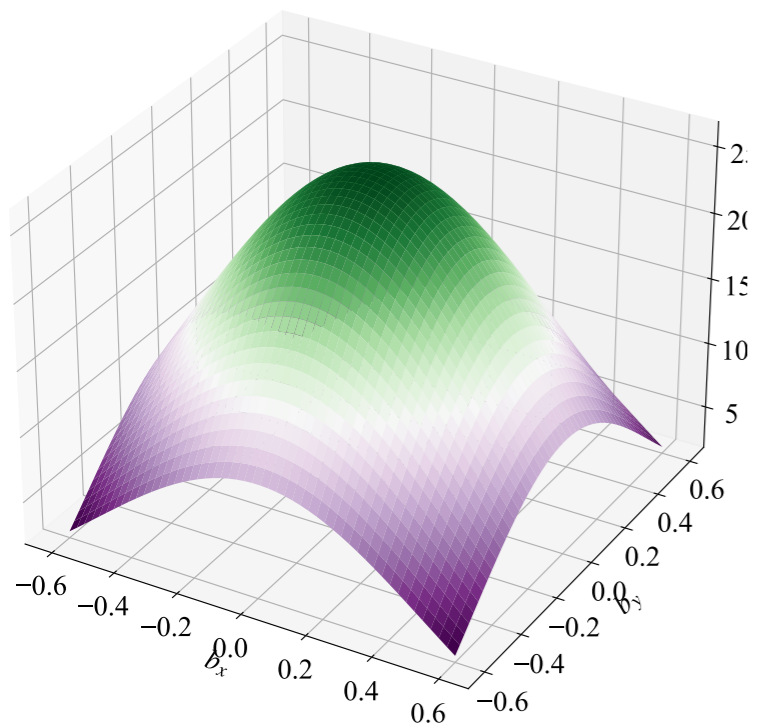
$x=0.2$



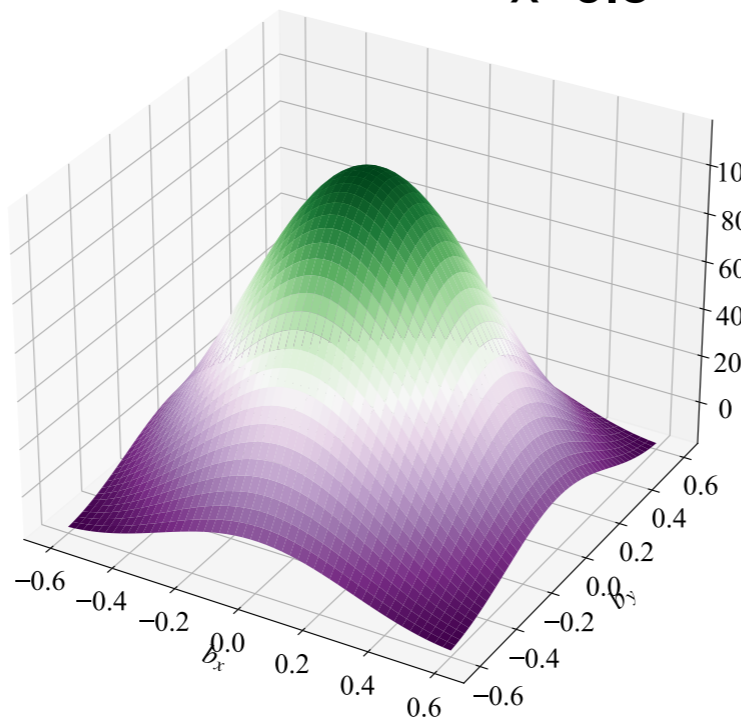
$x=0.4$



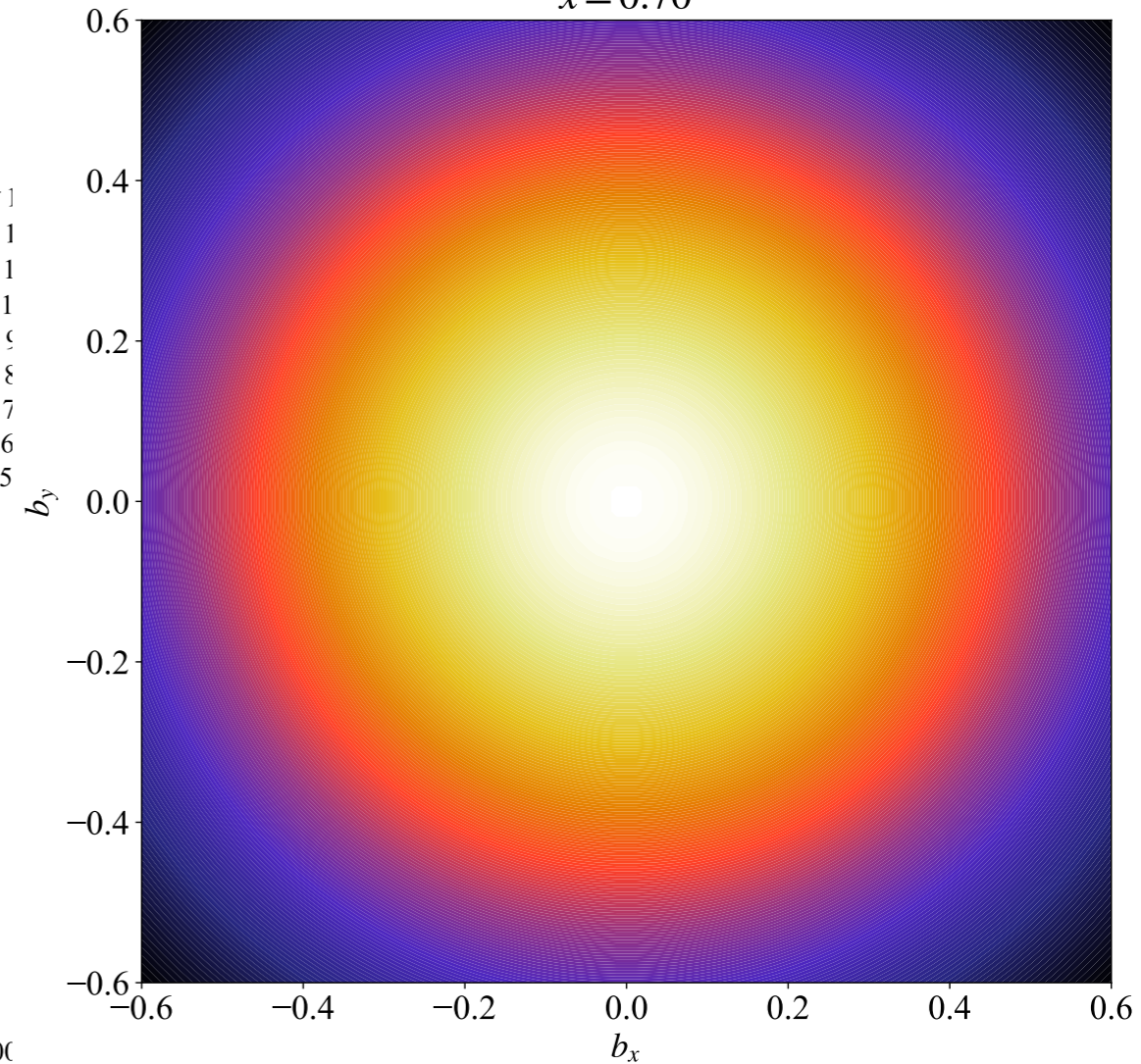
$x=0.6$



$x=0.8$



$x=0.70$



★ GPDs in transverse plane

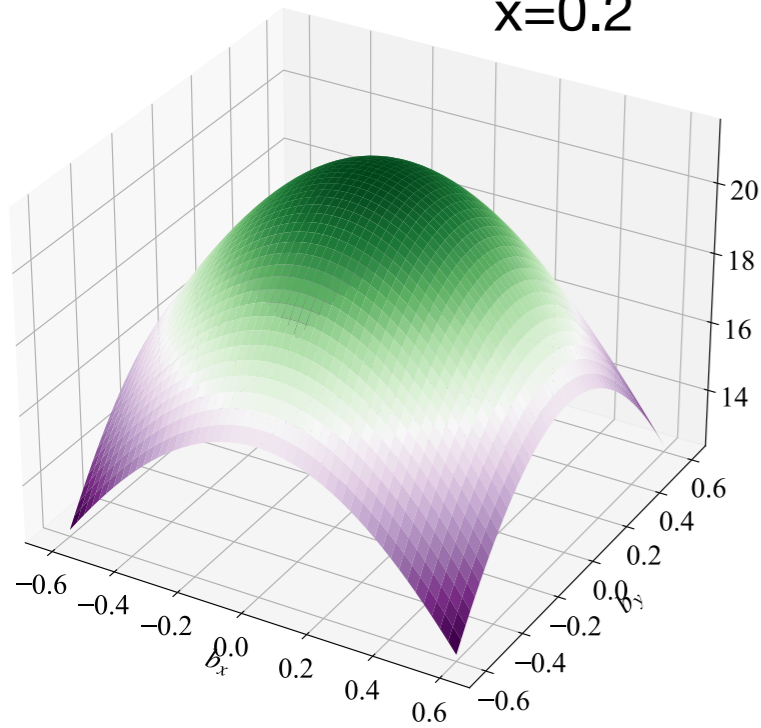
$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

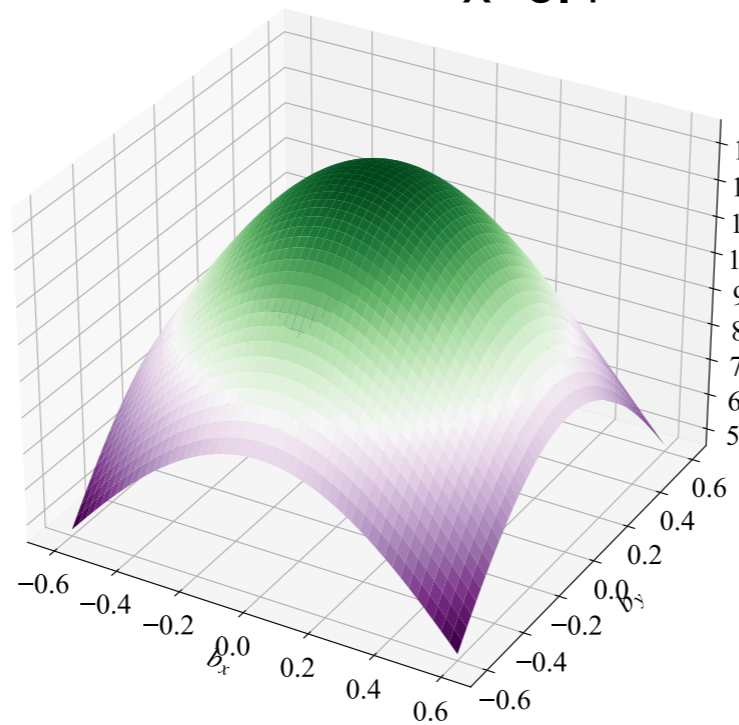
\mathbf{b}_\perp : transverse distance from the
(transverse) center of momentum

Impact parameter space $\widetilde{E} + \widetilde{G}_1$

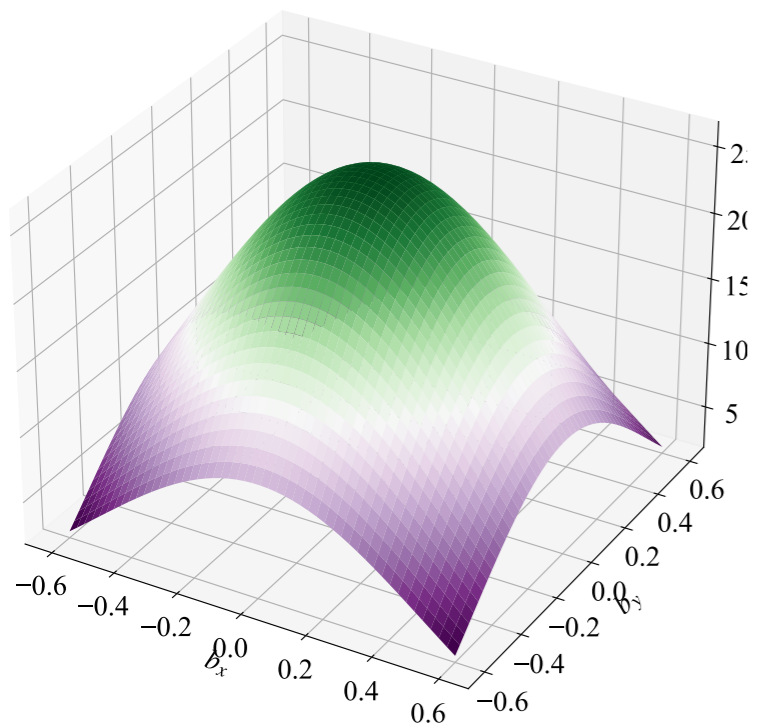
$x=0.2$



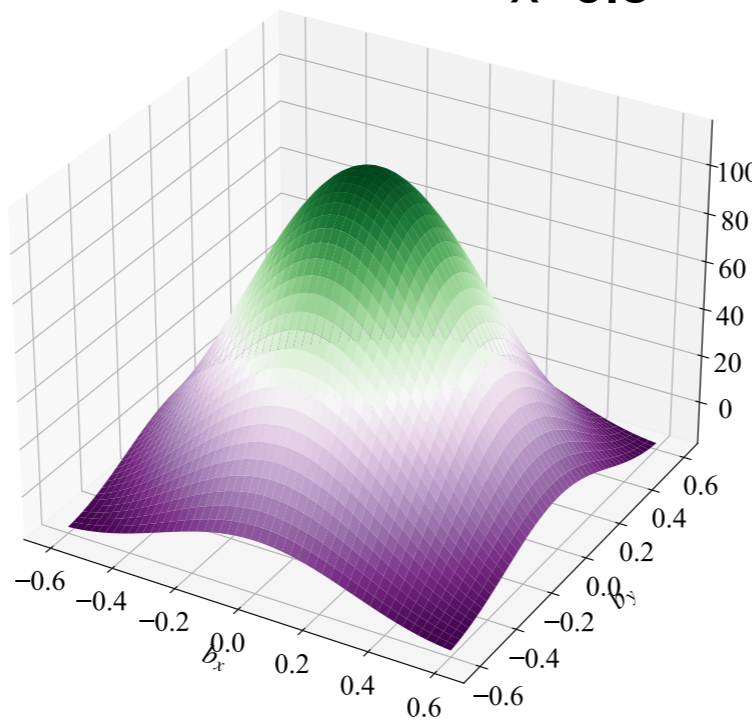
$x=0.4$



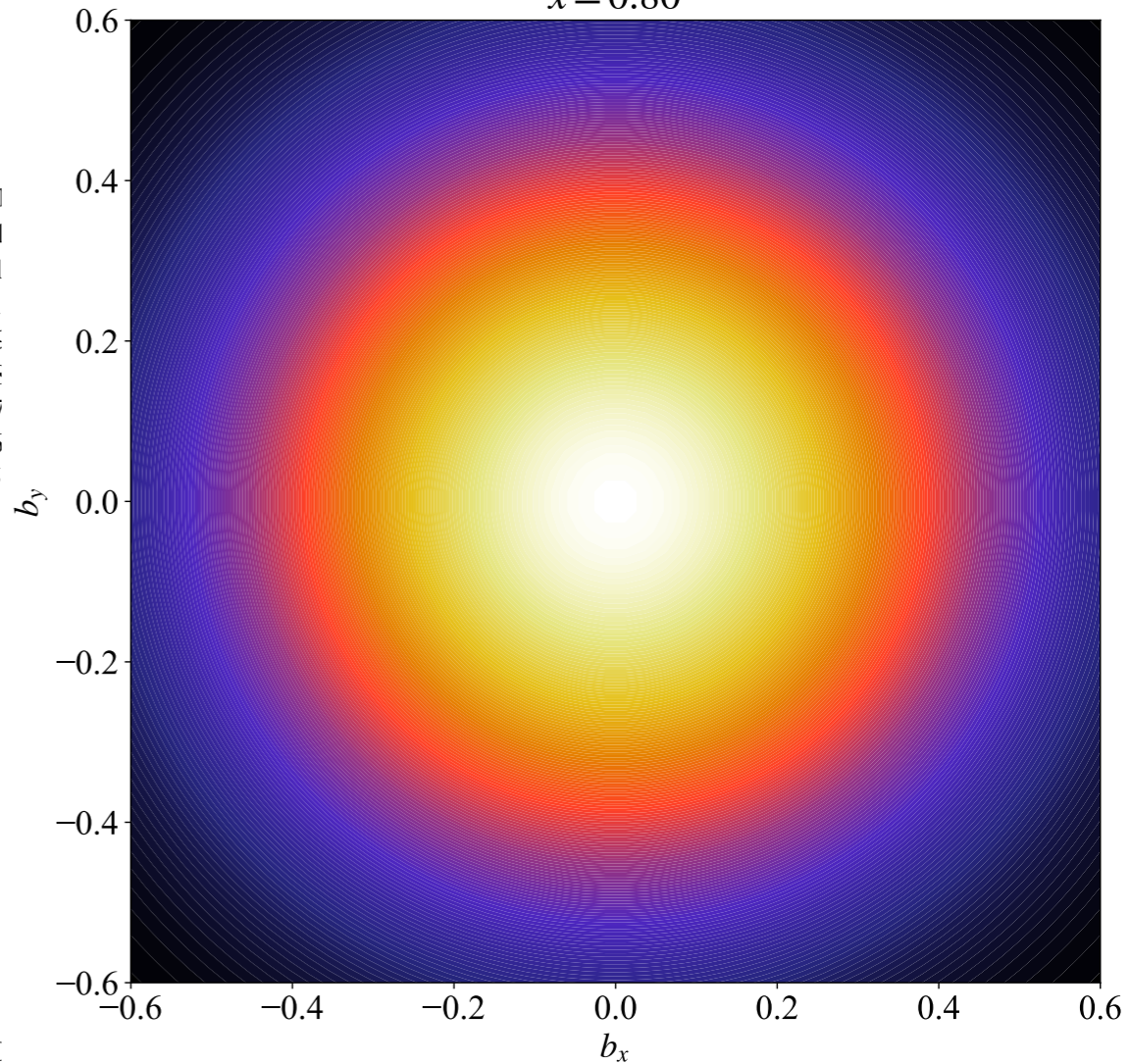
$x=0.6$



$x=0.8$



$x=0.80$



★ GPDs in transverse plane

$$q(x, \mathbf{b}_\perp) = |\mathcal{N}|^2 \int \frac{d^2 \mathbf{p}_\perp}{(2\pi)^2} \int \frac{d^2 \mathbf{p}'_\perp}{(2\pi)^2} H_q(x, -(\mathbf{p}_\perp - \mathbf{p}'_\perp)^2) e^{i\mathbf{b}_\perp \cdot (\mathbf{p}_\perp - \mathbf{p}'_\perp)}$$

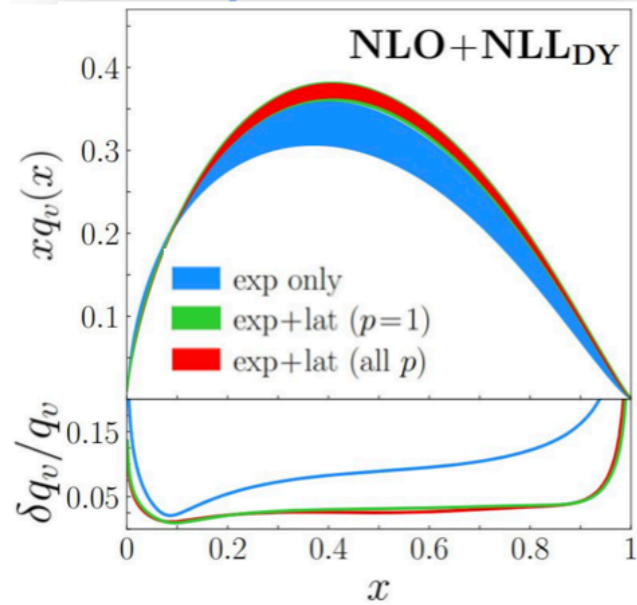
$$= \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp},$$

b_\perp : transverse distance from the
(transverse) center of momentum

Synergy/Complementarity of lattice and phenomenology

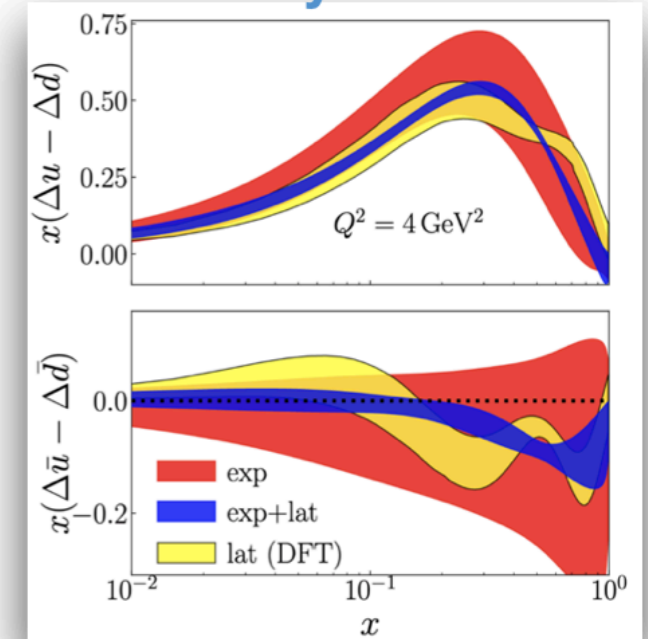
Synergies: constraints & predictive power of lattice QCD

pion PDF

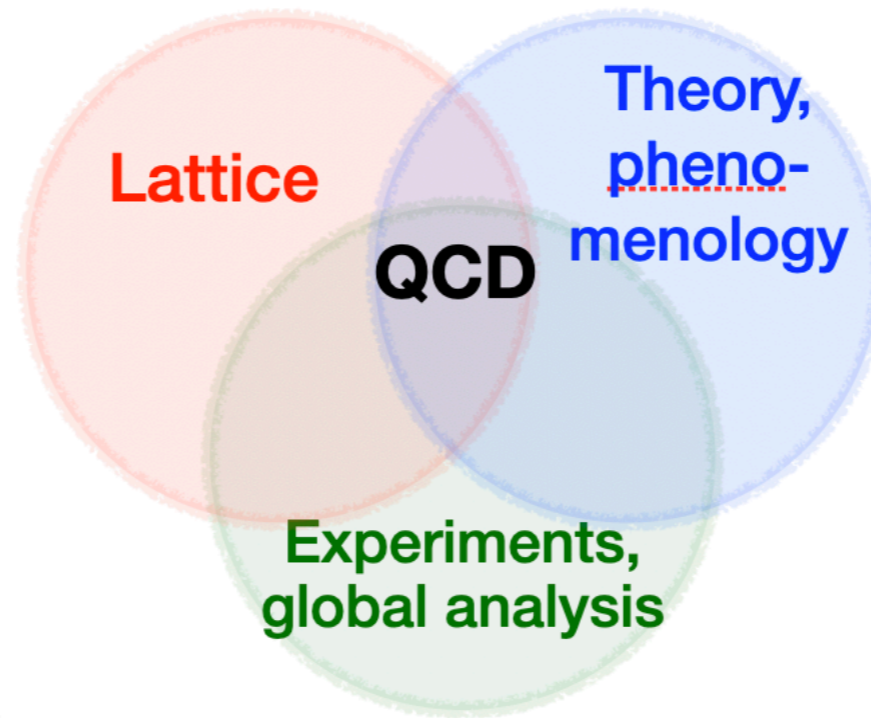


[JAM/HadStruc, PRD105 (2022) 114051]

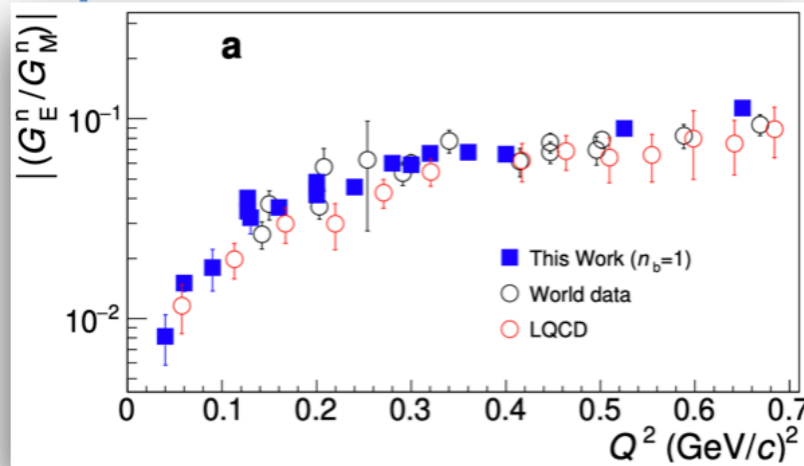
helicity PDF



[JAM & ETMC, PRD 103 (2021) 016003]

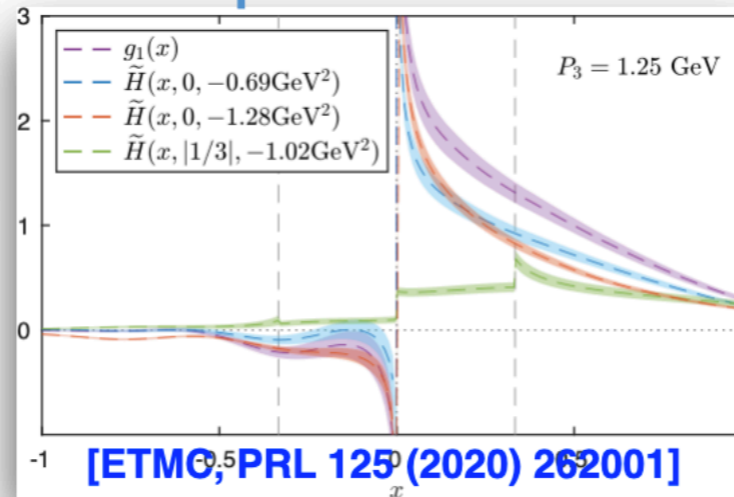


proton & neutron radius



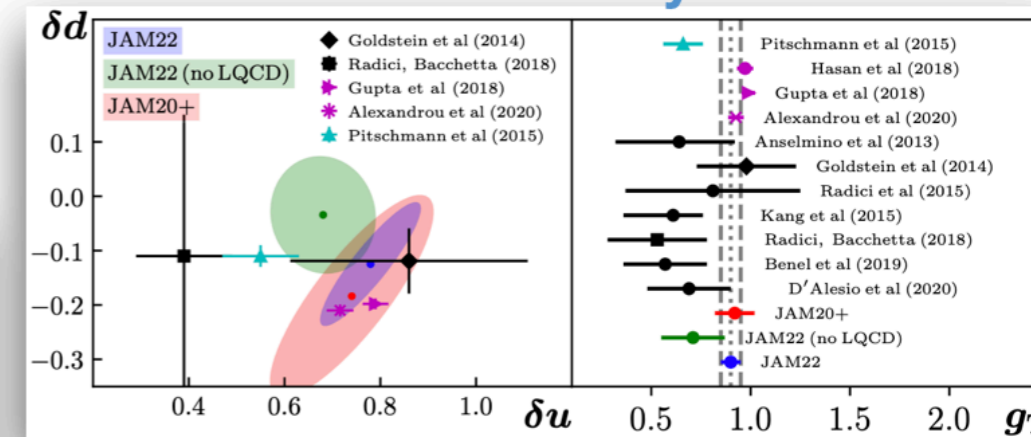
[Atac et al., Nature Comm. 12, 1759 (2021)]

proton GPDs



[ETMC, PRL 125 (2020) 262001]

transversity PDF



[JAM, PRD 106 (2022) 3, 034014]

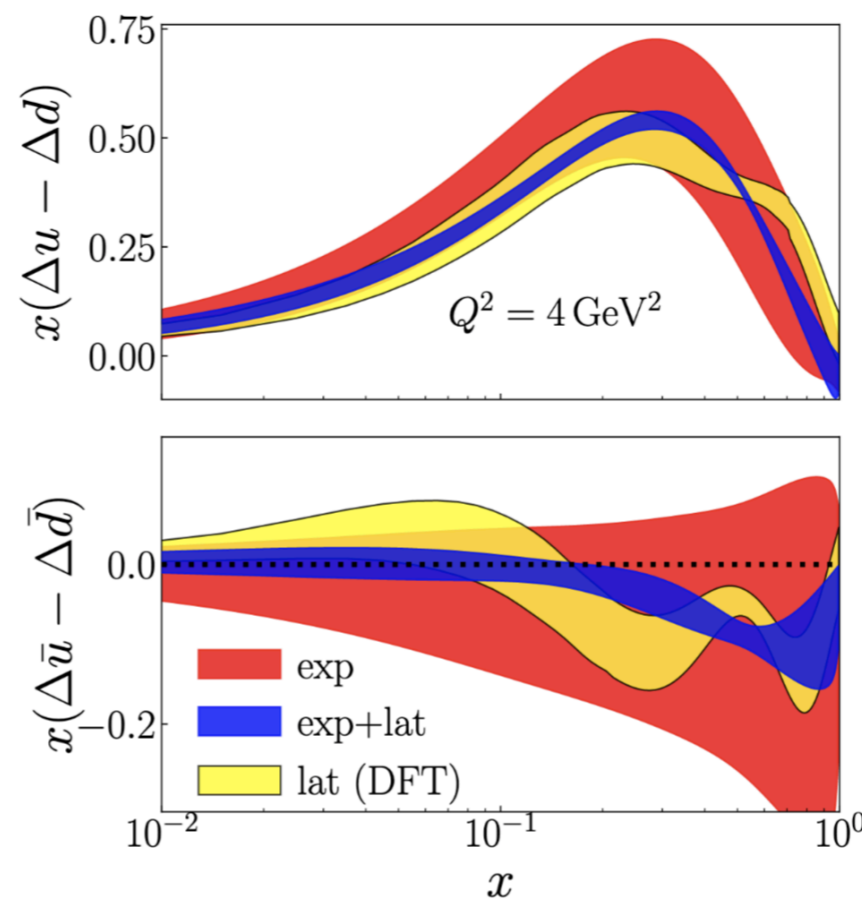
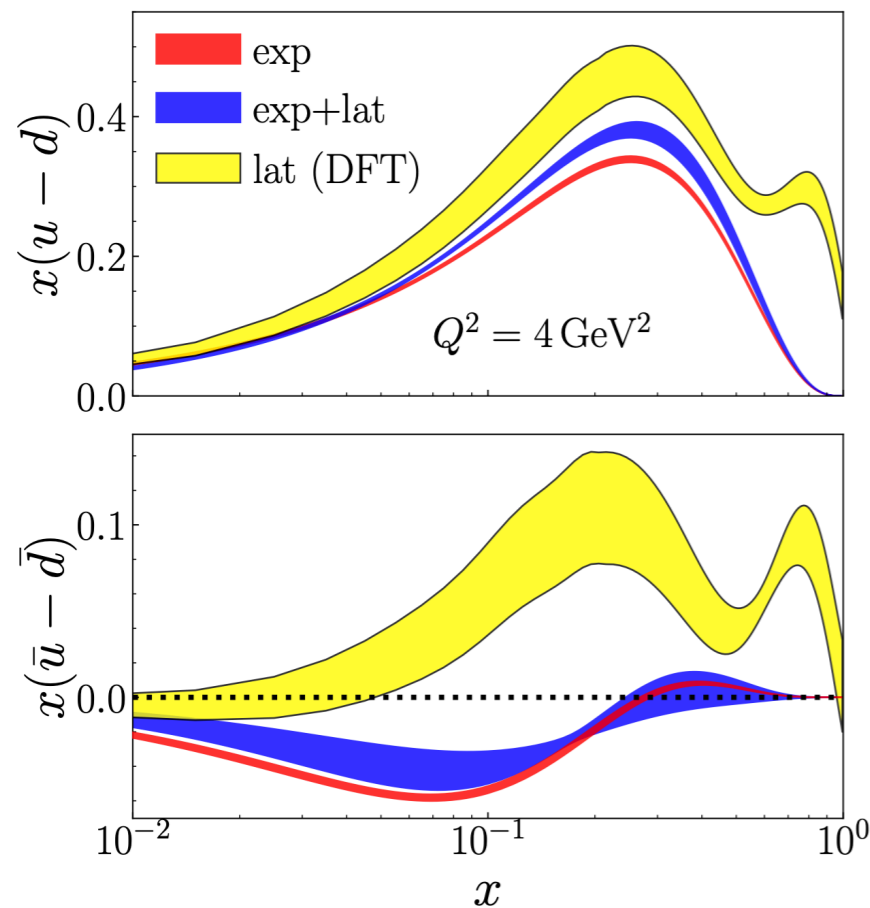
And many more!

Incorporating lattice PDFs in global analyses

Synergy between lattice and phenomenology

- ★ Lattice and experimental data sets data within the same global analysis (JAM framework)

[J. Bringewatt et al., PRD 103 (2021) 016003, arXiv:2010.00548]



- Consistent picture with JAM for unpolarized PDF

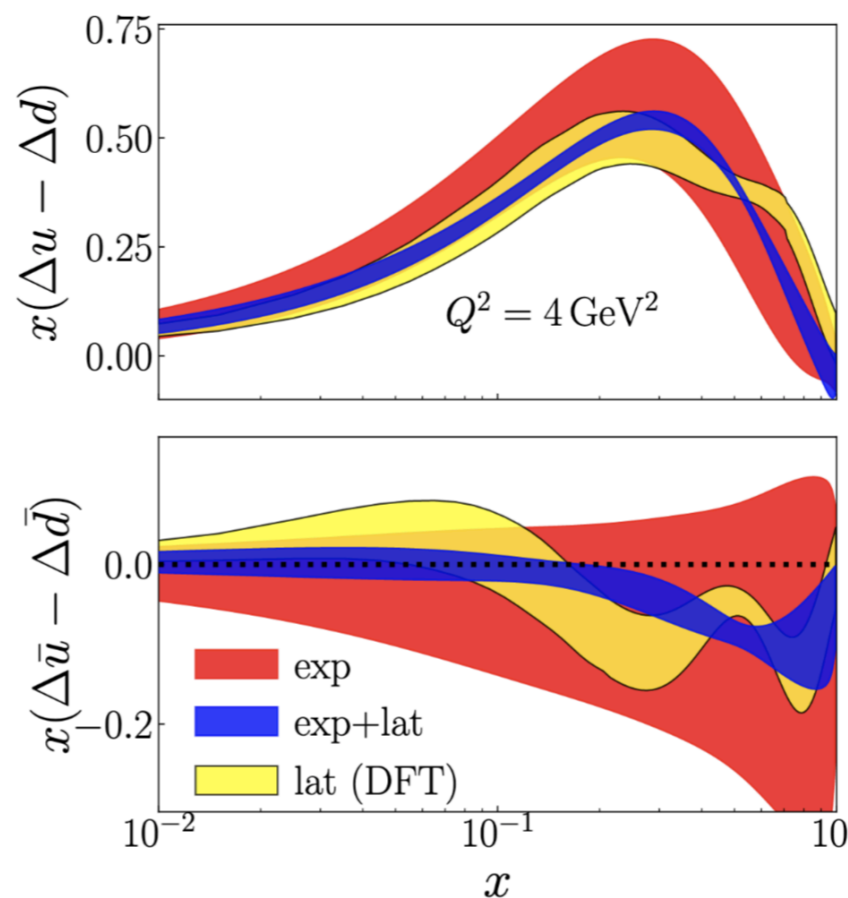
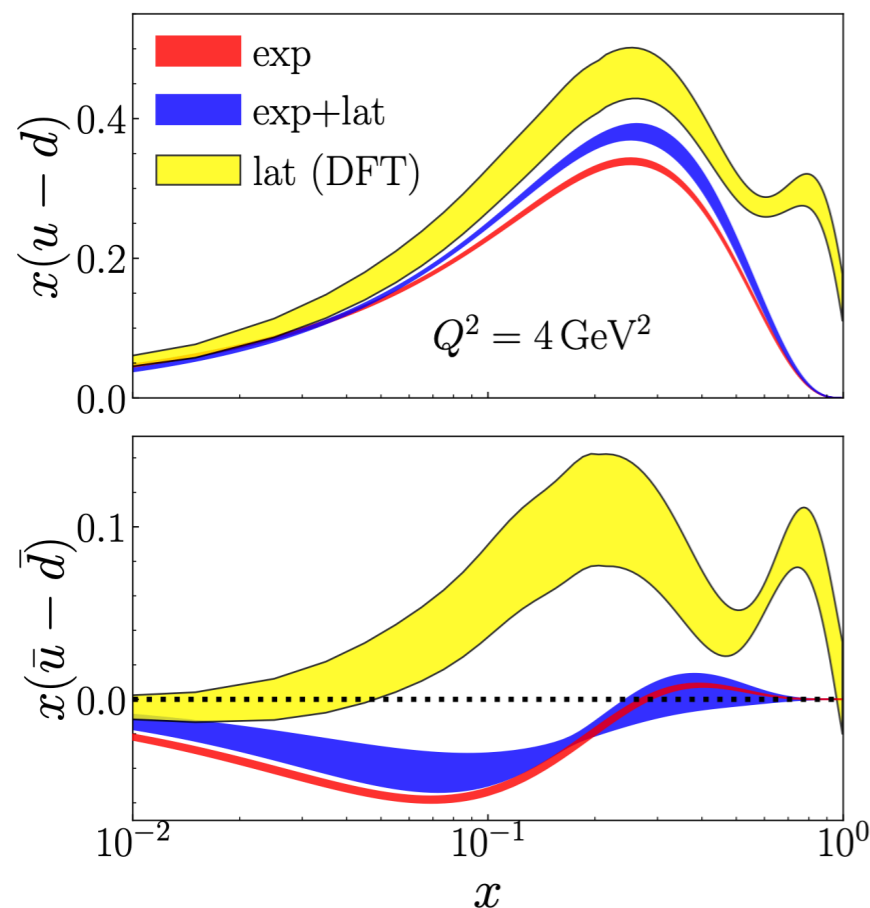
- Significant impact for helicity PDF

Incorporating lattice PDFs in global analyses

Synergy between lattice and phenomenology

- ★ Lattice and experimental data sets data within the same global analysis (JAM framework)

[J. Bringewatt et al., PRD 103 (2021) 016003, arXiv:2010.00548]



- Consistent picture with JAM for unpolarized PDF

- Significant impact for helicity PDF

- ★ Other efforts within NNPDF framework

[K. Cichy et al., JHEP 10 (2019) 137, arXiv:1907.06037]
[L. Del Debbio et al., JHEP 02 (2021) 138, 2010.03996]

- ★ Interest in applying similar approach to quantities that are more challenging to extract experimentally (GPDs, twist-3 distributions, ...)

Toward synergy for GPDs

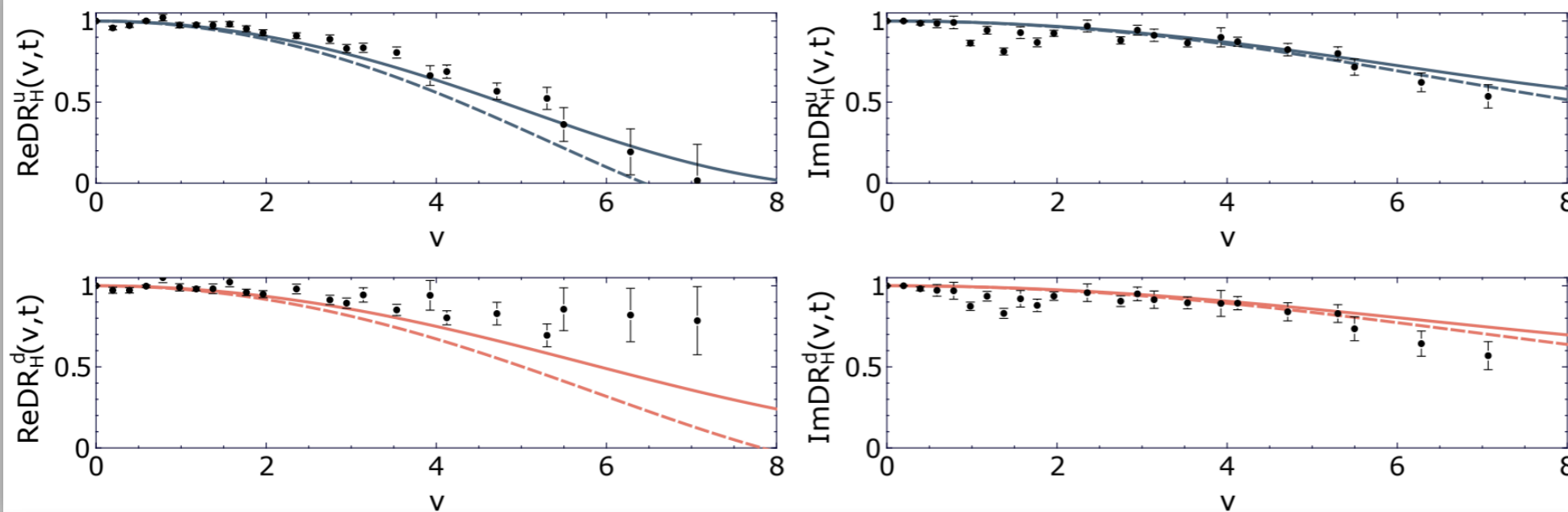
- ★ Forming ratios of GPDs seems to suppress systematic uncertainties

[K. Cichy et al., arXiv:2409.17955]

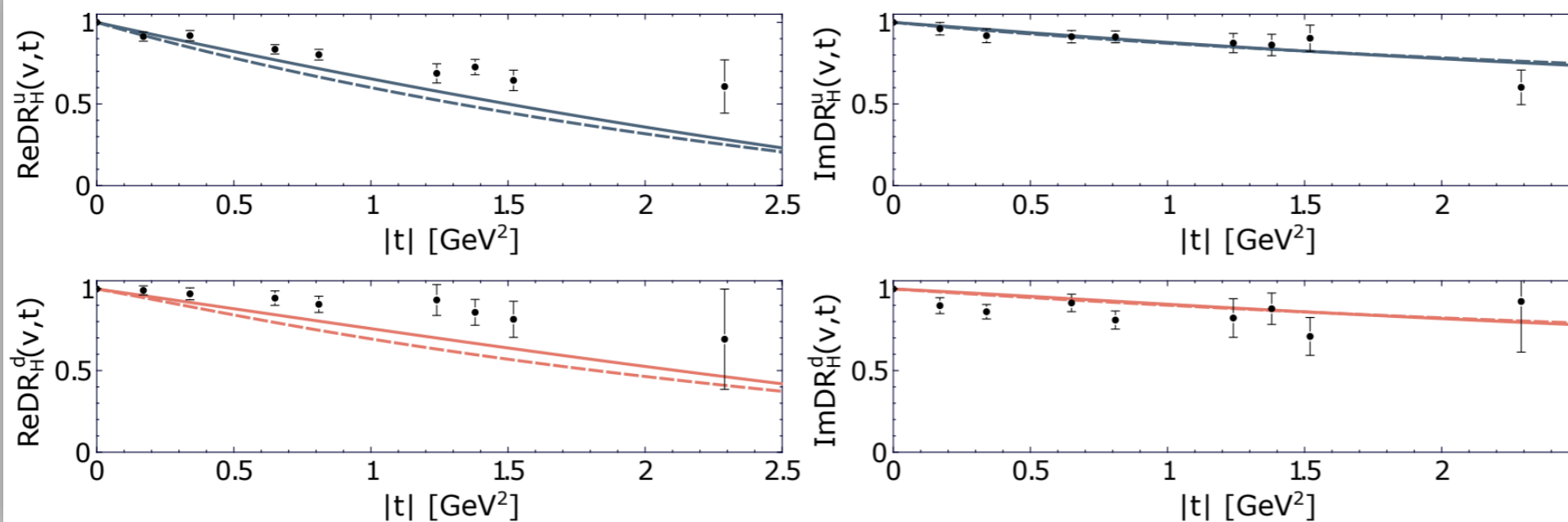
$$\text{DR}_{\text{Re}}^{\hat{H}^q}(\nu, t) = \frac{\text{Re}\hat{H}^q(\nu, t) \text{Re}\hat{H}^q(0, 0)}{\text{Re}\hat{H}^q(\nu, 0) \text{Re}\hat{H}^q(0, t)},$$

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(a) As a function of ν for $|t| = 0.65 \text{ GeV}^2$.



(b) As a function of $|t|$ for $\nu = 3.14$.



- GK (solid curve)
- VGG (dashed curve)
- Good agreement for up quark
- Reasonable agreement for down quark

Toward synergy for GPDs

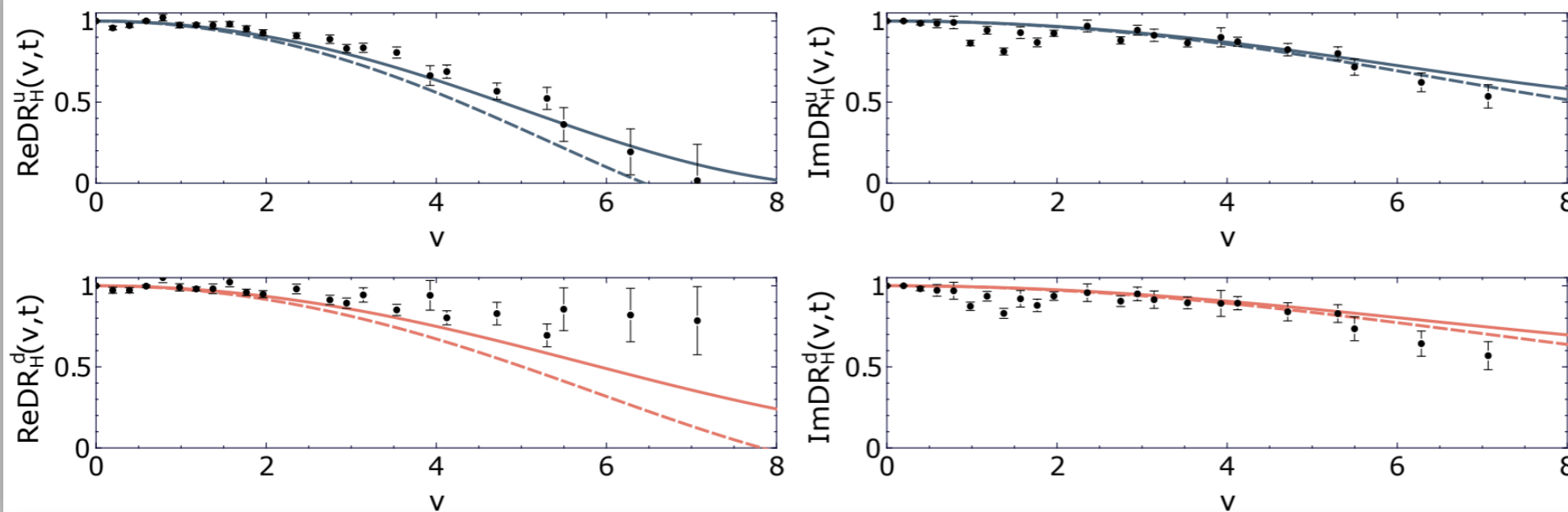
- ★ Forming ratios of GPDs seems to suppress systematic uncertainties

[K. Cichy et al., arXiv:2409.17955]

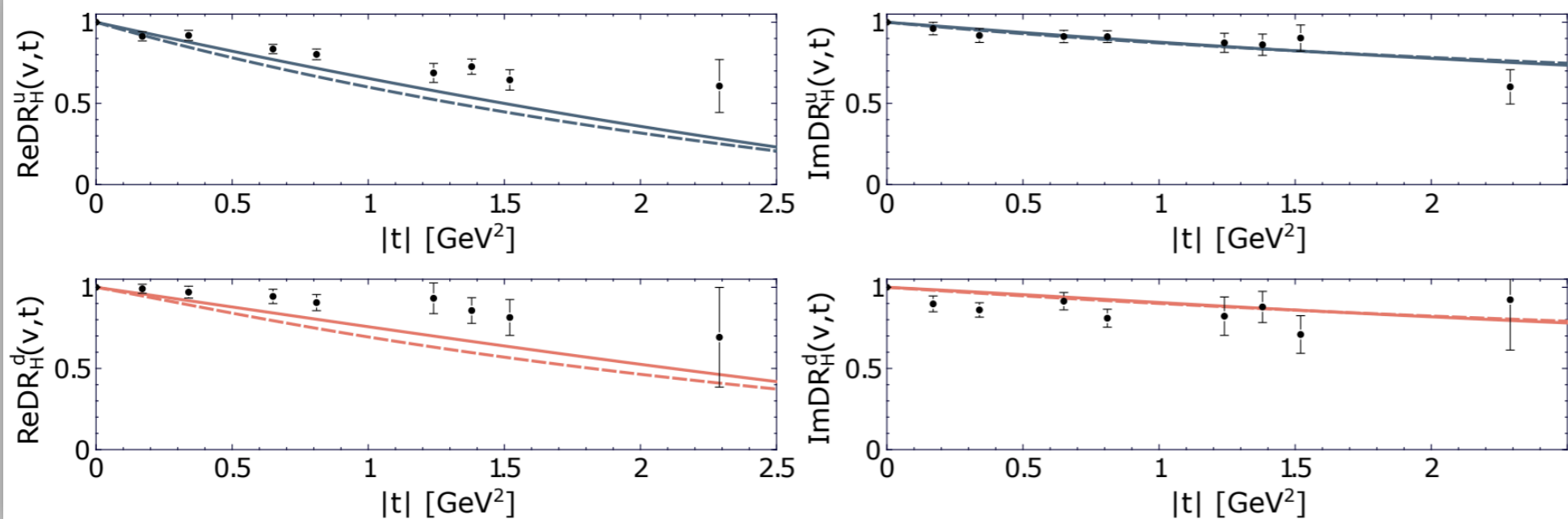
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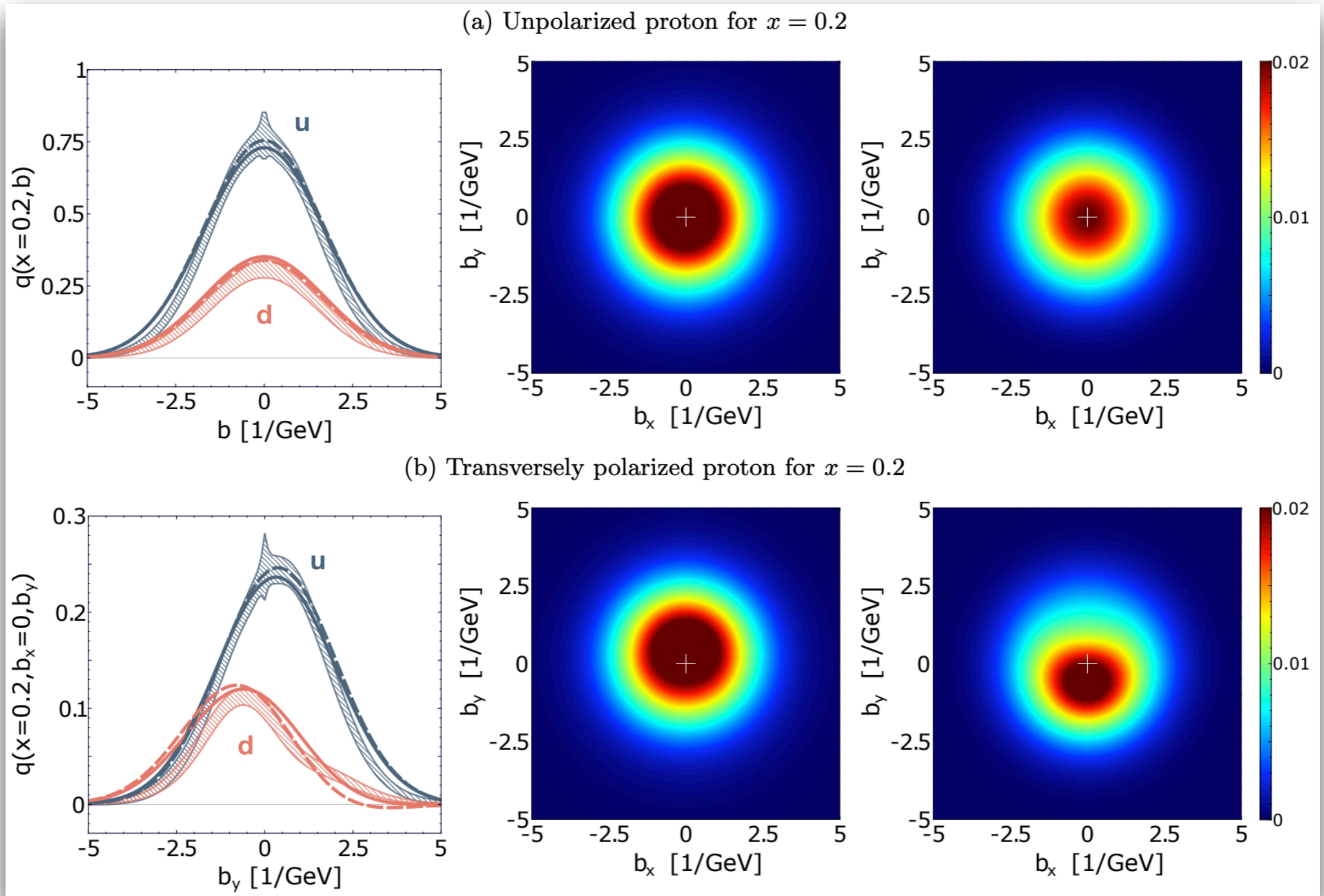


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- Good agreement for up quark
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- Further study needed on how to combine lattice results with data

Tomographic Images



- GK (solid line), VGG (dashed line)

[K. Cichy et al., arXiv:2409.17955]

How to lattice QCD data fit into the overall effort for hadron tomography

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QUARK-GLUON TOMOGRAPHY COLLABORATION



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Award Number:
DE-SC0023646

1. **Theoretical studies** of high-momentum transfer processes using perturbative QCD methods and study of GPDs properties
2. **Lattice QCD** calculations of GPDs and related structures
3. **Global analysis** of GPDs based on experimental data using modern data analysis techniques for inference and uncertainty quantification

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Other GPD global analysis efforts:

- Gepard [<https://gepard.phy.hr/>]
- PARTONS [<https://partons.cea.fr>]
- EXCLAIM [<https://exclaimcollab.github.io/web.github.io/#/>]

Concluding remarks

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- ★ Extensive programs in Gluon PDFs
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Thank you



DOE Early Career Award (NP)
Grant No. DE-SC0020405



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WW approximation

[S. Bhattacharya et al., PRD 102 (2020) 11 (Editors Selection)]

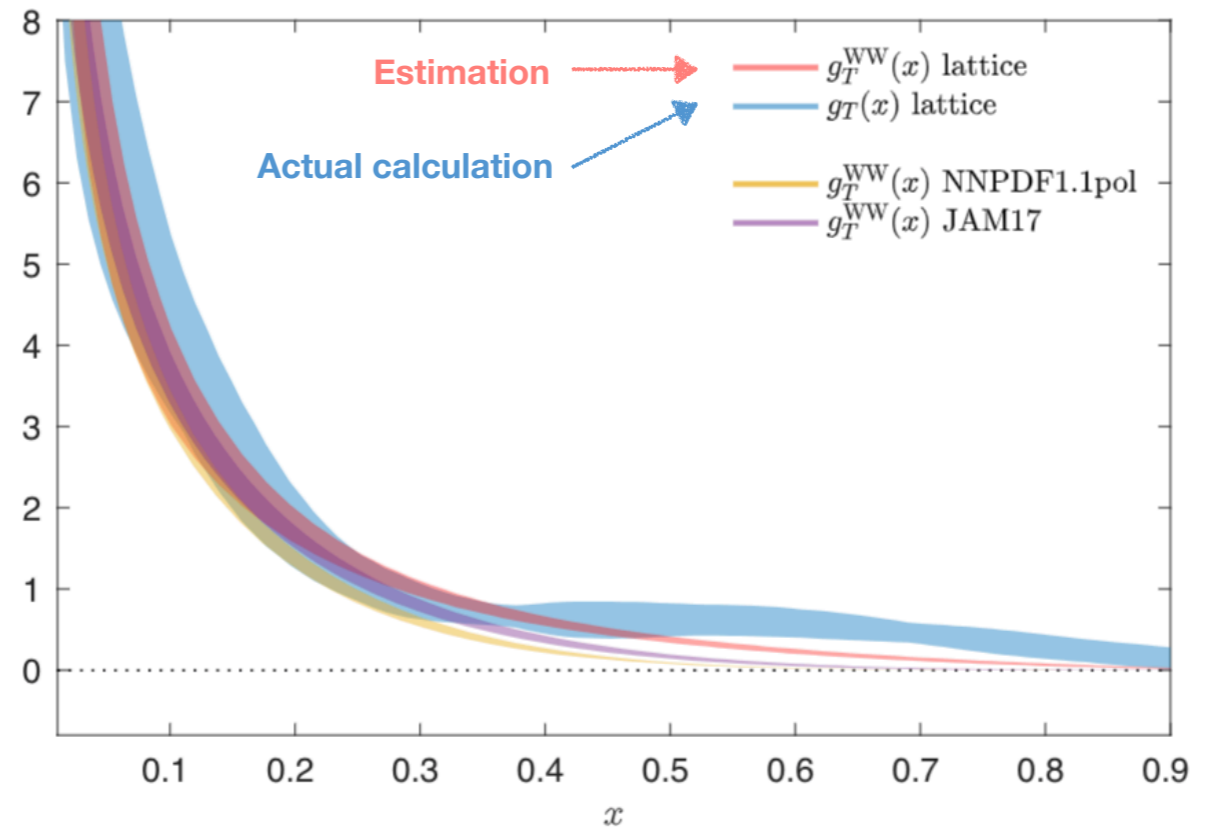
WW approximation:

$$g_T^{\text{WW}}(x) = \int_x^1 \frac{dy}{y} g_1(y)$$

twist-3 $g_T(x)$ determined by the twist-2 $g_1(x)$

- $g_T(x)$ agrees with $g_T^{\text{WW}}(x)$ for $x < 0.5$ (violations up to 30-40% possible)
- Violations of 15-40% expected from experimental data

[A. Accardi et al., JHEP 11 (2009) 093]



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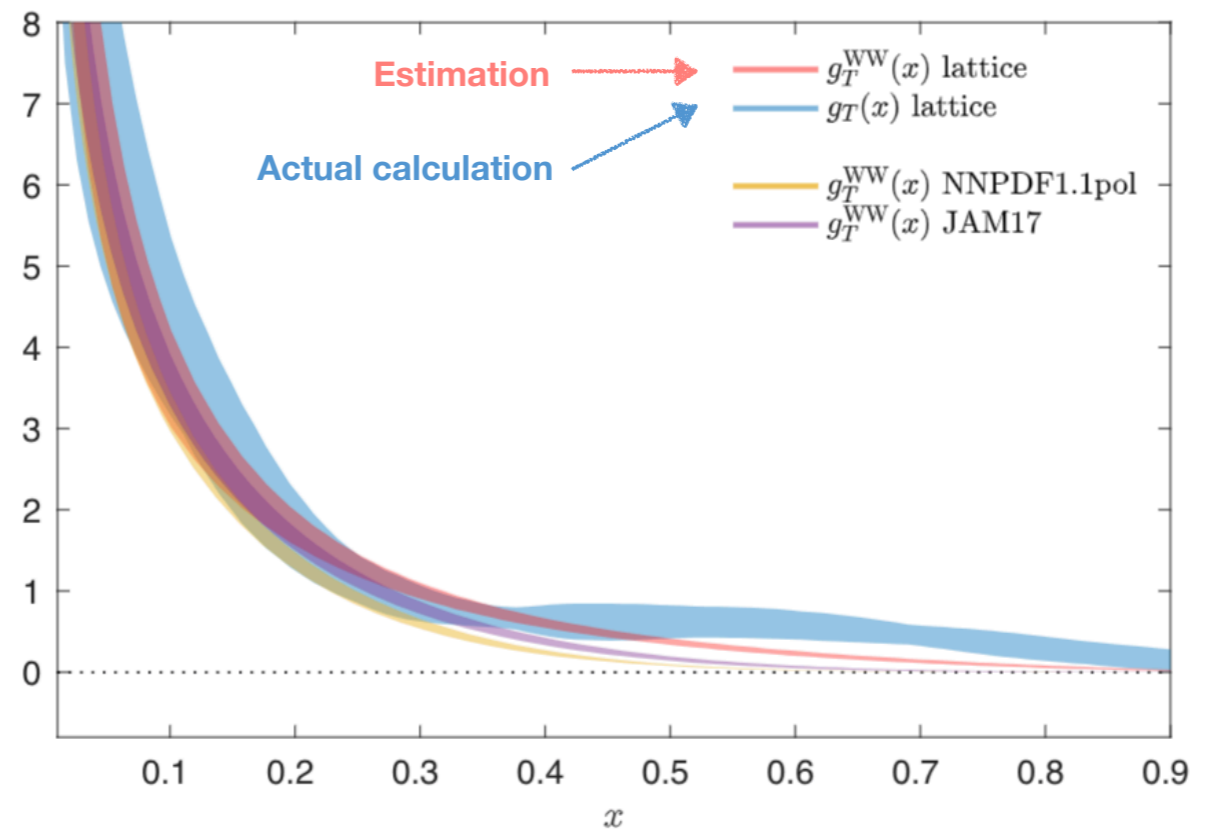
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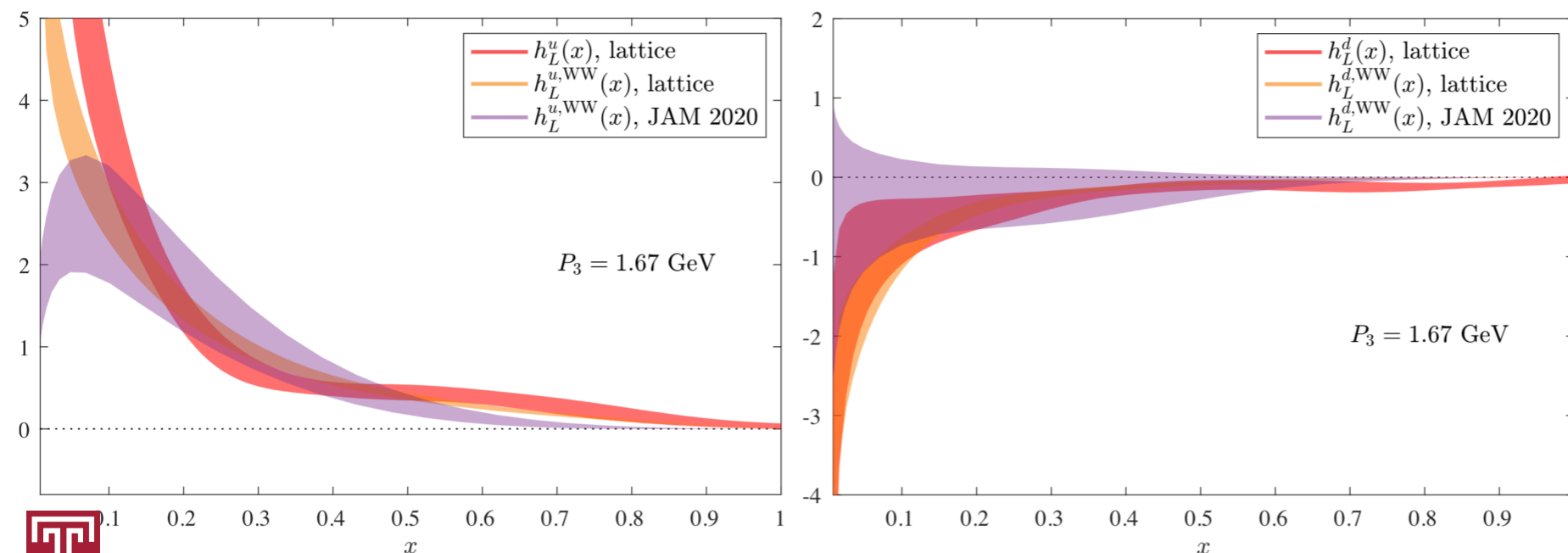
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Twist-3 $h_L(x)$ PDF

[S. Bhattacharya et al., PRD 104 (2021) 11, 114510]

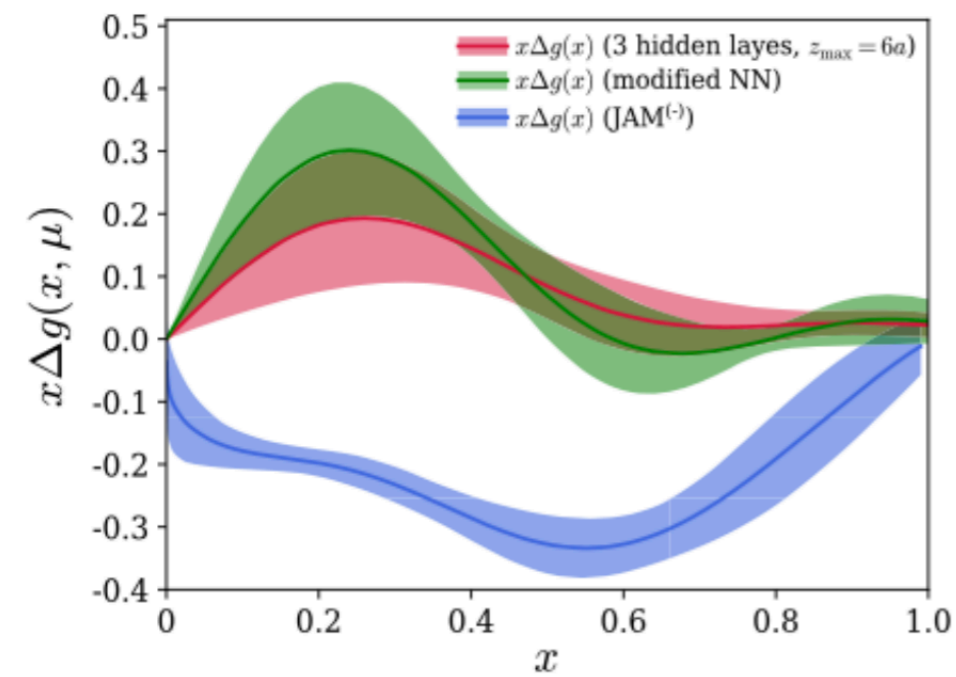


- h_L^u dominant - tension between h_L & h_L^{WW}
- $h_L^d < 0$ and decays faster than h_L^u



Gluon Helicity PDF

- Neural network analysis of lattice calculation disfavors negative gluon polarizability



[T. Khan et al., PRD 108, 074502]

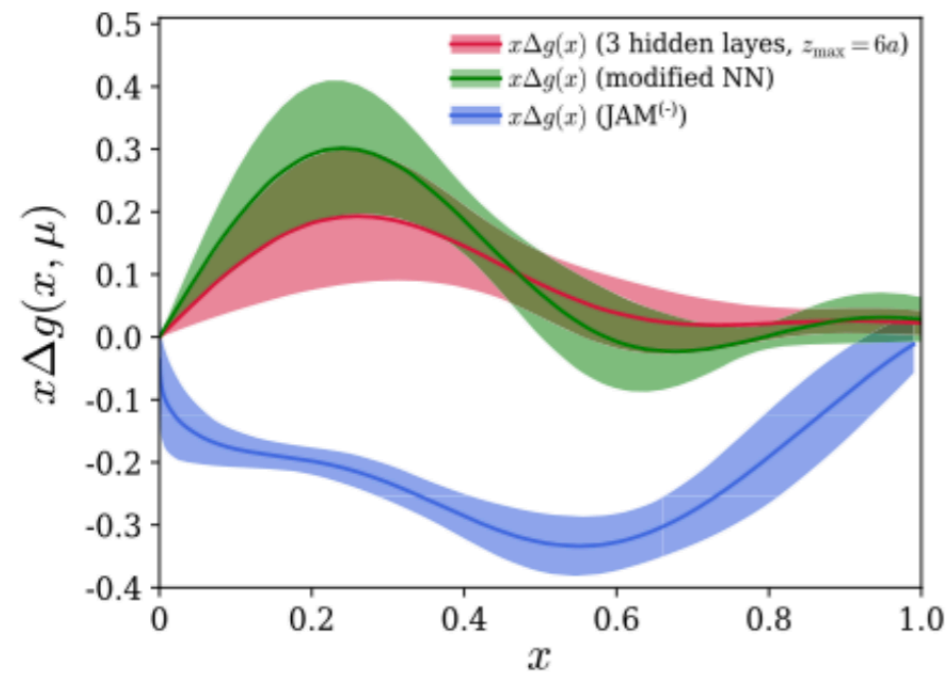
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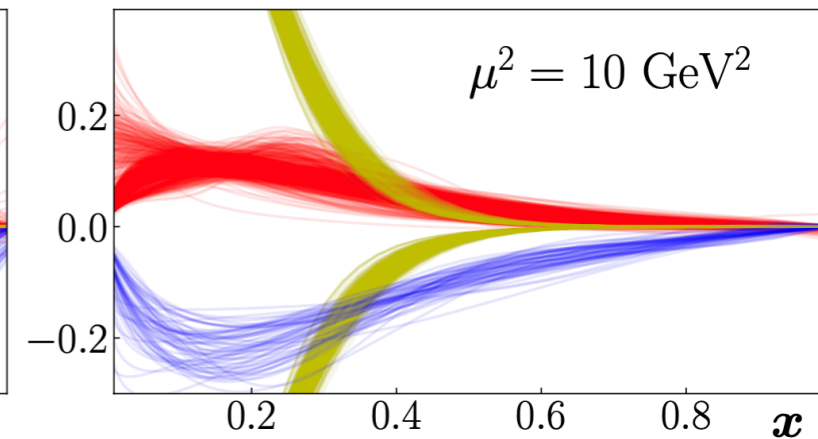
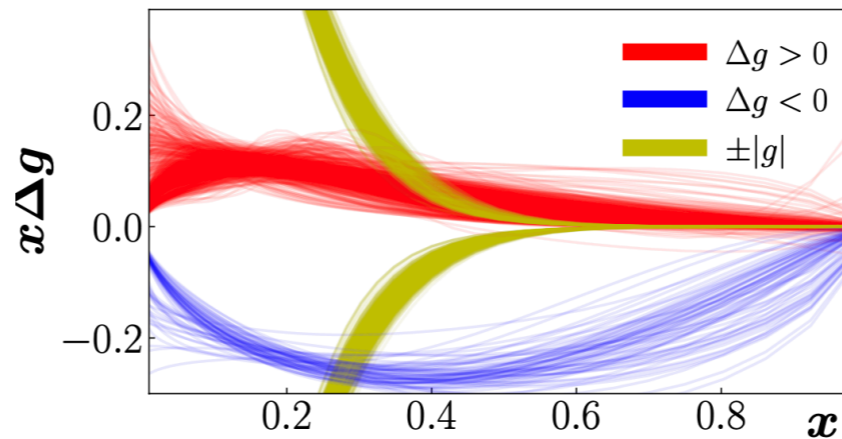
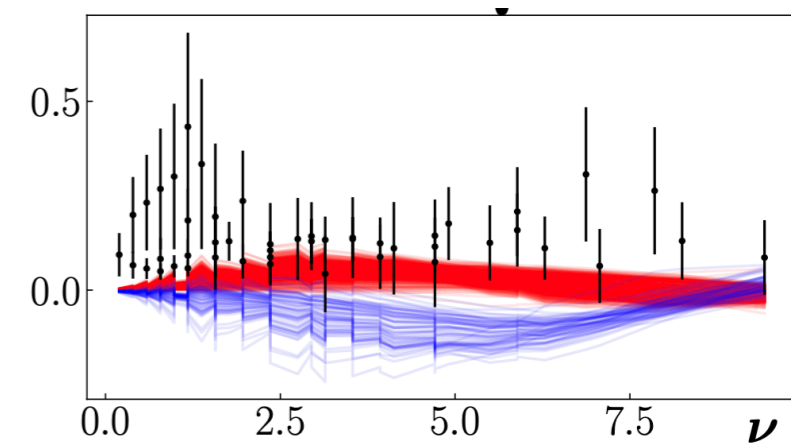
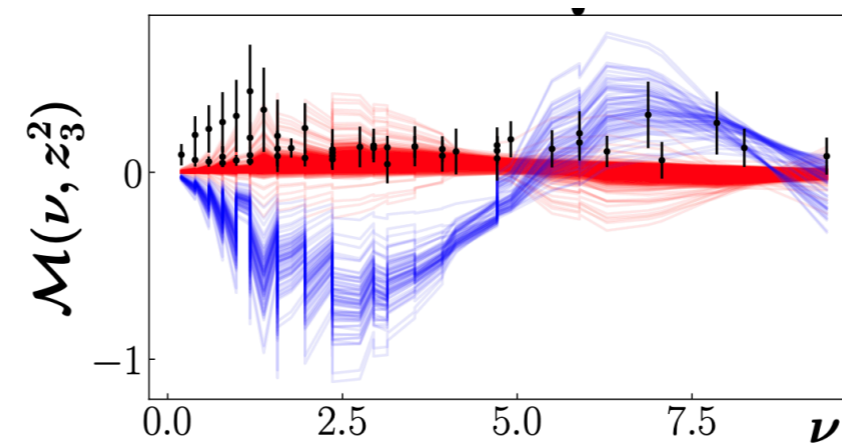
[J. Karpie et al., PRD 109 (2024) 3, 036031]

Without Lattice

Including Lattice



[T. Khan et al., PRD 108, 074502]



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 JAM analysis: No positivity constraint
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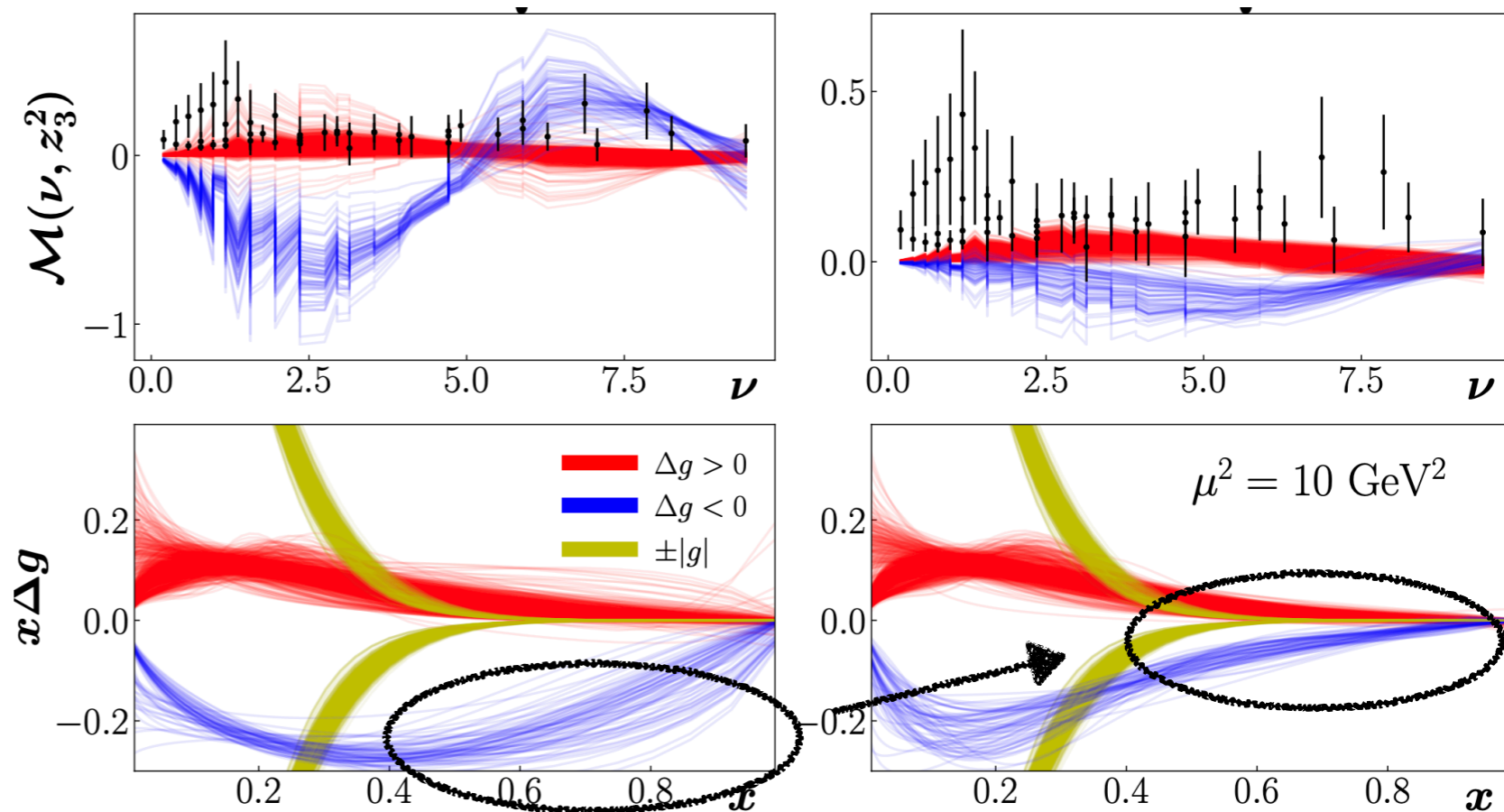
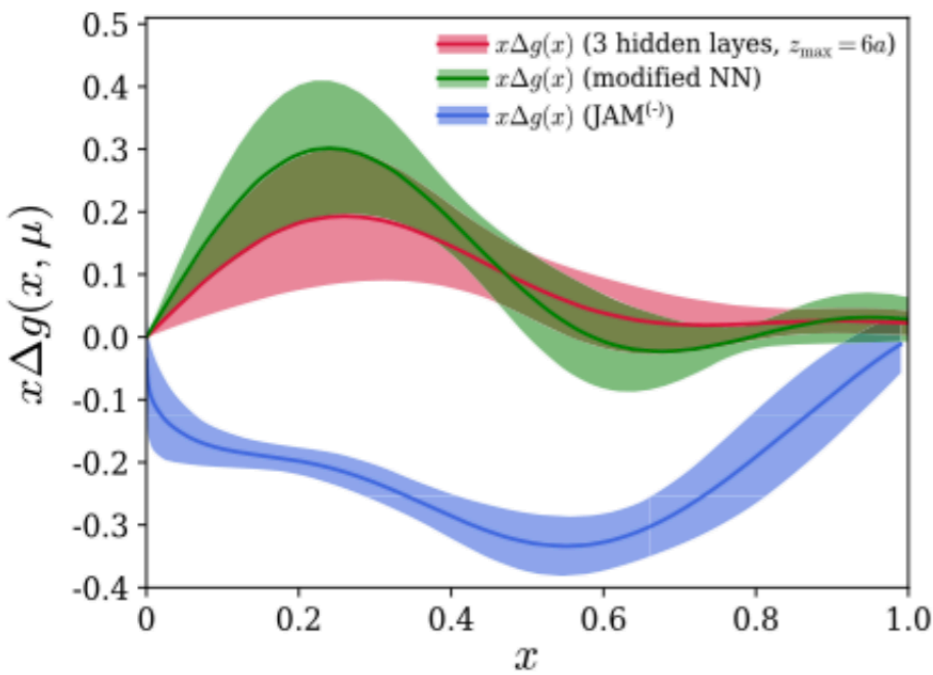
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