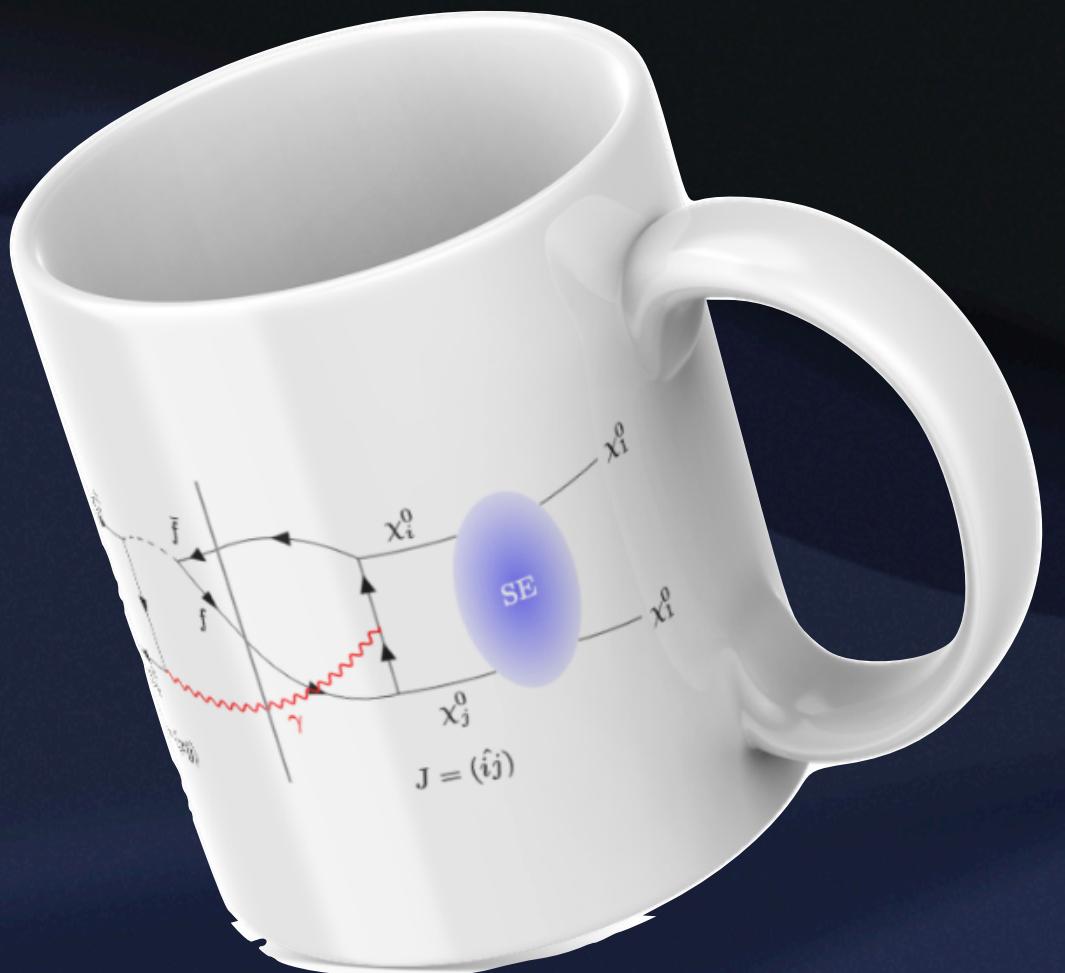


TBA

The Dark Matter Landscape
MITP 2024

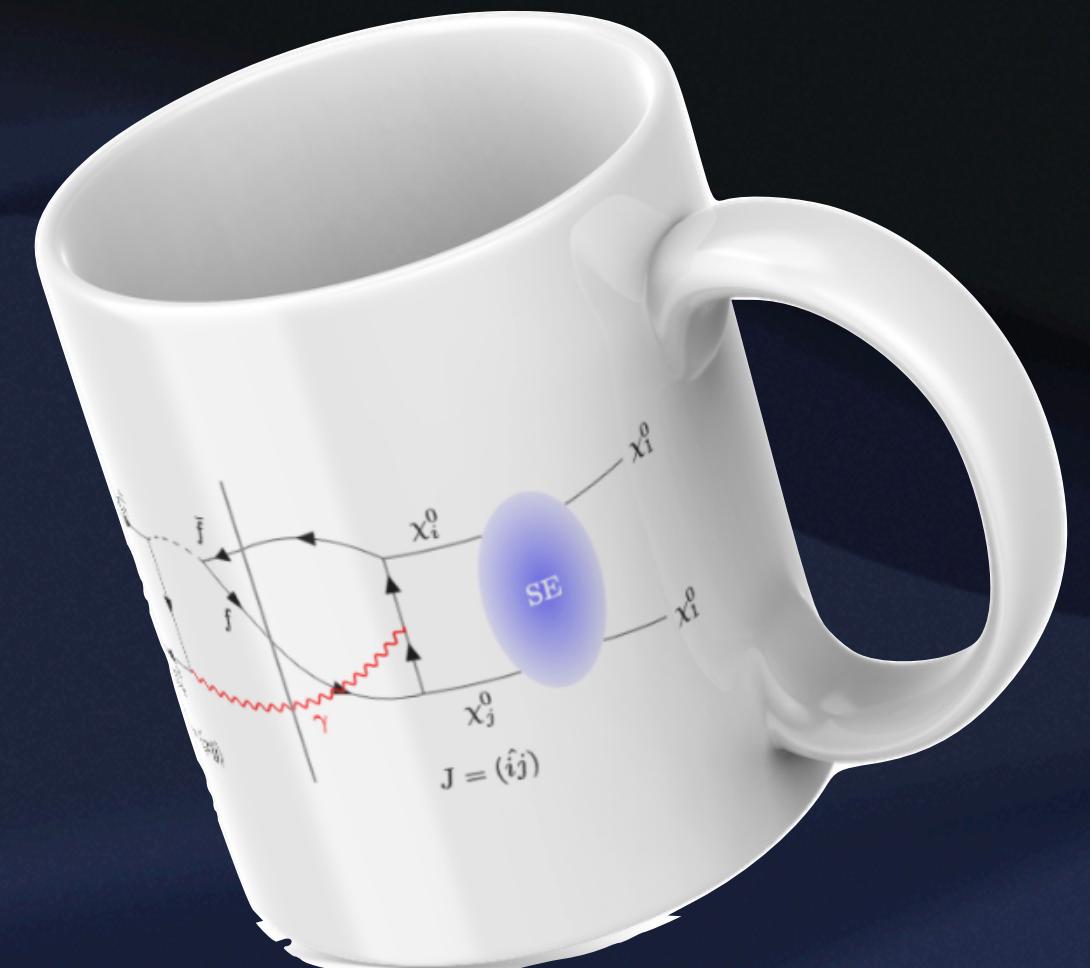
Martin Vollmann – Uni Tuebingen – 13.9.2024



The Basics of Annihilation*

The Dark Matter Landscape
MITP 2024

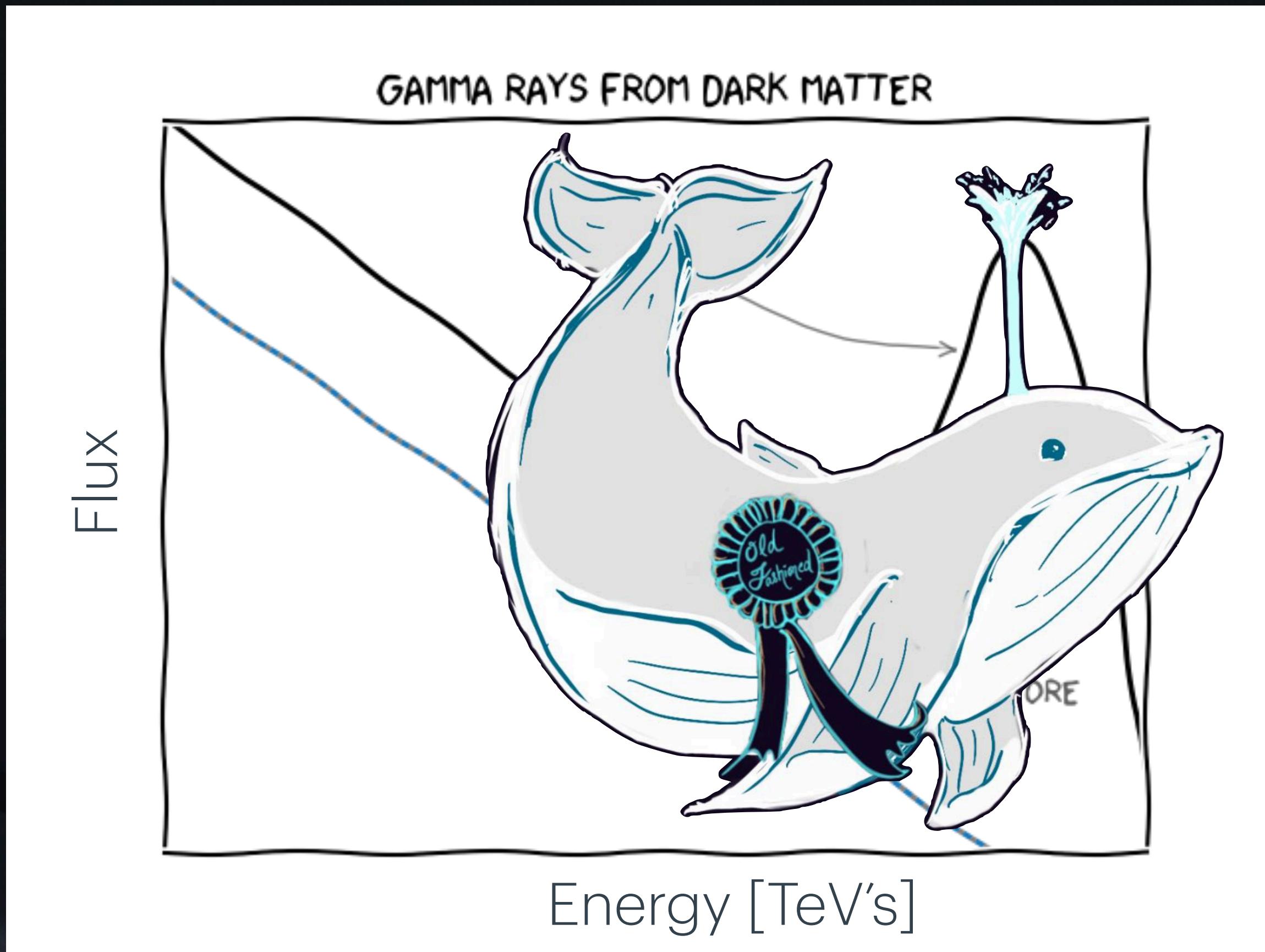
Martin Vollmann – Uni Tuebingen – 13.9.2024

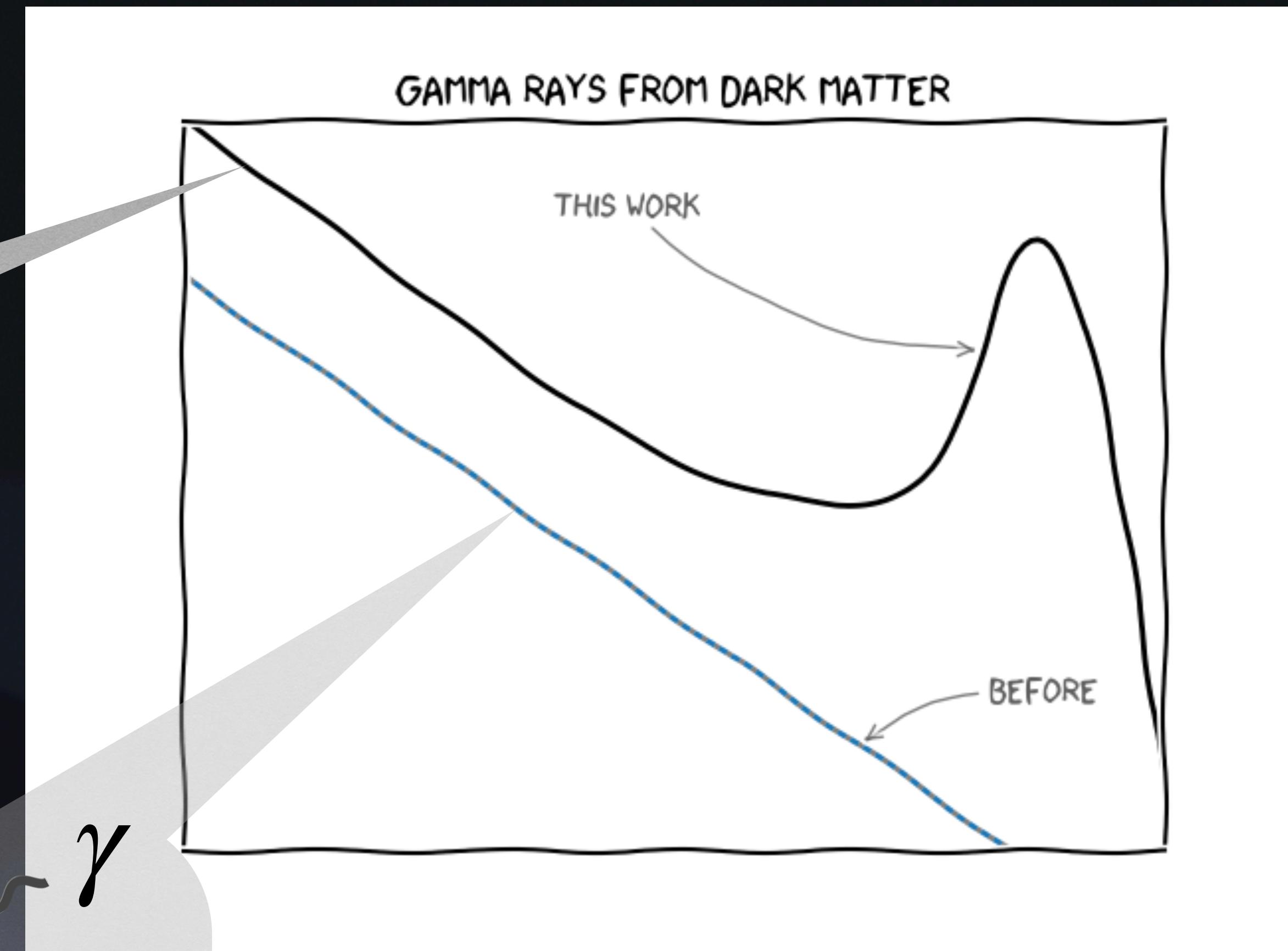
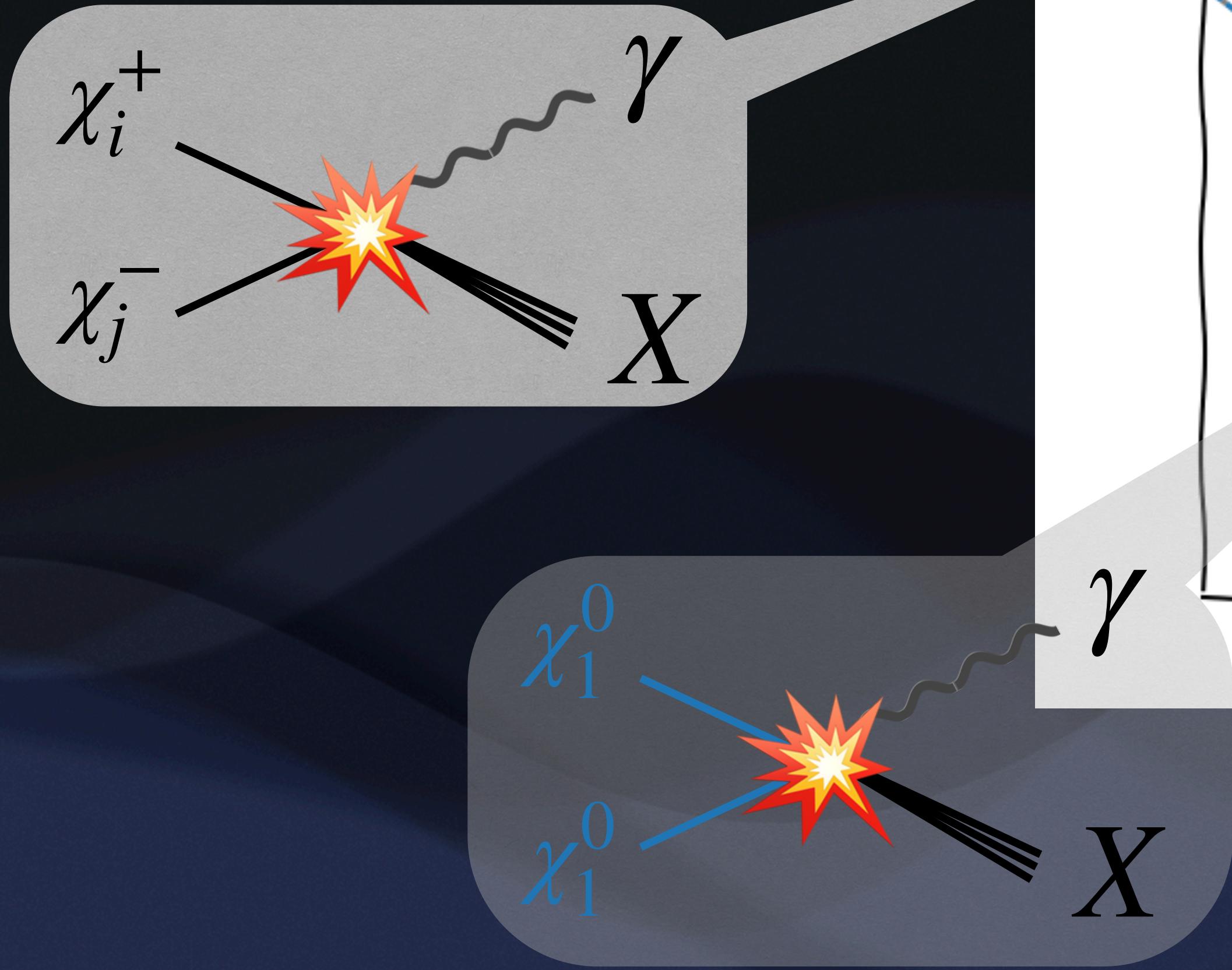


Credit: S. Sevillano/M. Fujiwara

Intro

Very rough and (deliberately) misleading picture





2410.XXXXX

Resummed photon spectrum from ~~MOTOKINO~~ interacting spin 1 dark matter annihilation

Motoko Fujiwara^{a*}, Martin Vollmann^{b†}

^aPhysik-Department, Technische Universität München,
James-Franck-Straße, 85748 Garching, Germany

^bInstitute for Theoretical Physics, University of Tübingen, Auf der Morgenstelle 14, 72076
Tübingen, Germany

Sommerfeld effect for continuum gamma-ray spectra from
Dark Matter annihilation

Barbara Jäger^{ID}, Martin Vollmann^{ID}

Institute for Theoretical Physics, University of Tübingen,
Auf der Morgenstelle 14, 72076 Tübingen, Germany

October 18, 2023

October 18, 2023

Sommerfeld effect for continuum gamma-ray spectra from
Dark Matter annihilation

Outline

Motivation

Indirect detection

Sommerfeld factor

Endpoint resummations

Plots

Conclusions

Motivation

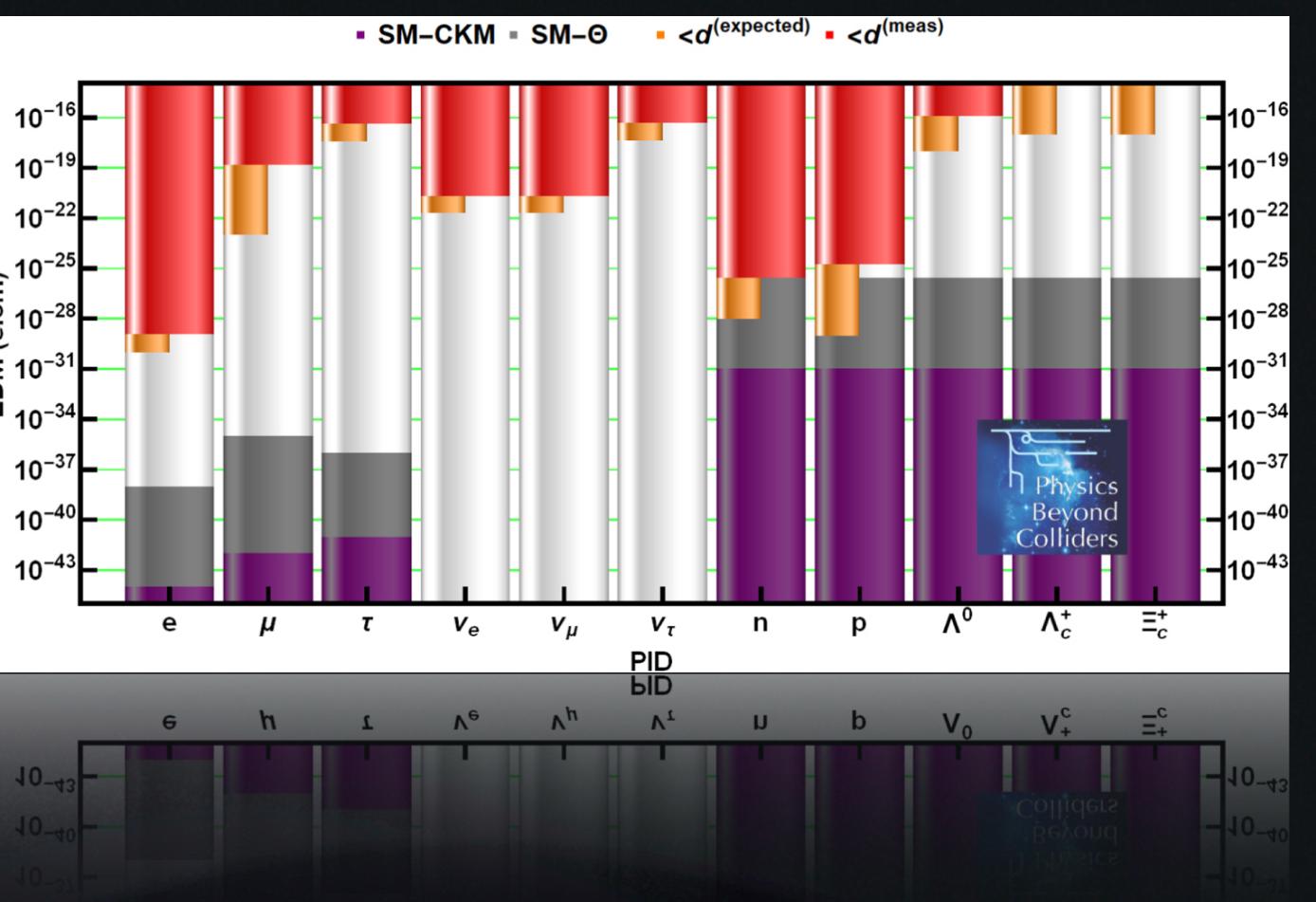
WIMP: the devil that you know

Did you know that...

- There is experimental evidence for **Dark Matter**
- There is experimental evidence for **electroweak interactions**
- Although there is no experimental evidence for the **freeze-out mechanism**
 - ▶ still the simplest!

High-scale SUSY

- Unification of gauge couplings improved
- No SUSY CP problem (EDMs)
- Pure wino/higgsino \rightarrow very predictive scenarios
 - Maybe detectable with next-generation γ -ray telescopes (e.g. Rodd, Safdi, Xu
[2405.13104](#))
- Naturalness



But really... Why?



And why you should still care!

- WIMPs are **not** dead yet!
- The miracle is still a miracle
- If you want to rule them out better be sure that you didn't overlook anything

NEED ACCURATE THEORETICAL PREDICTIONS

- The physics is fun and the “spare parts” are useful as long as there are electroweak interactions in nature

Outline

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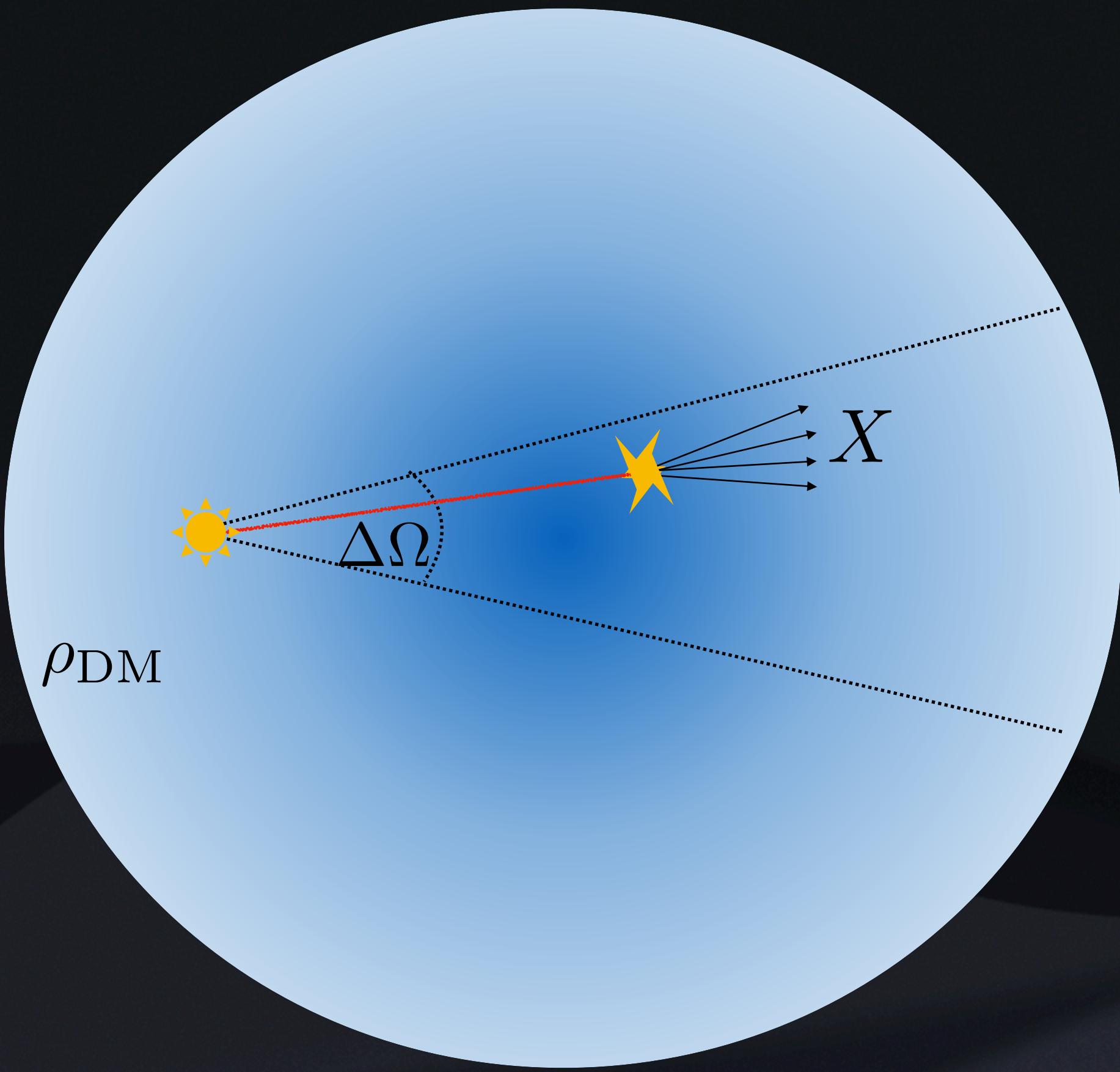
Conclusions

Indirect detection

Gamma-ray flux formula

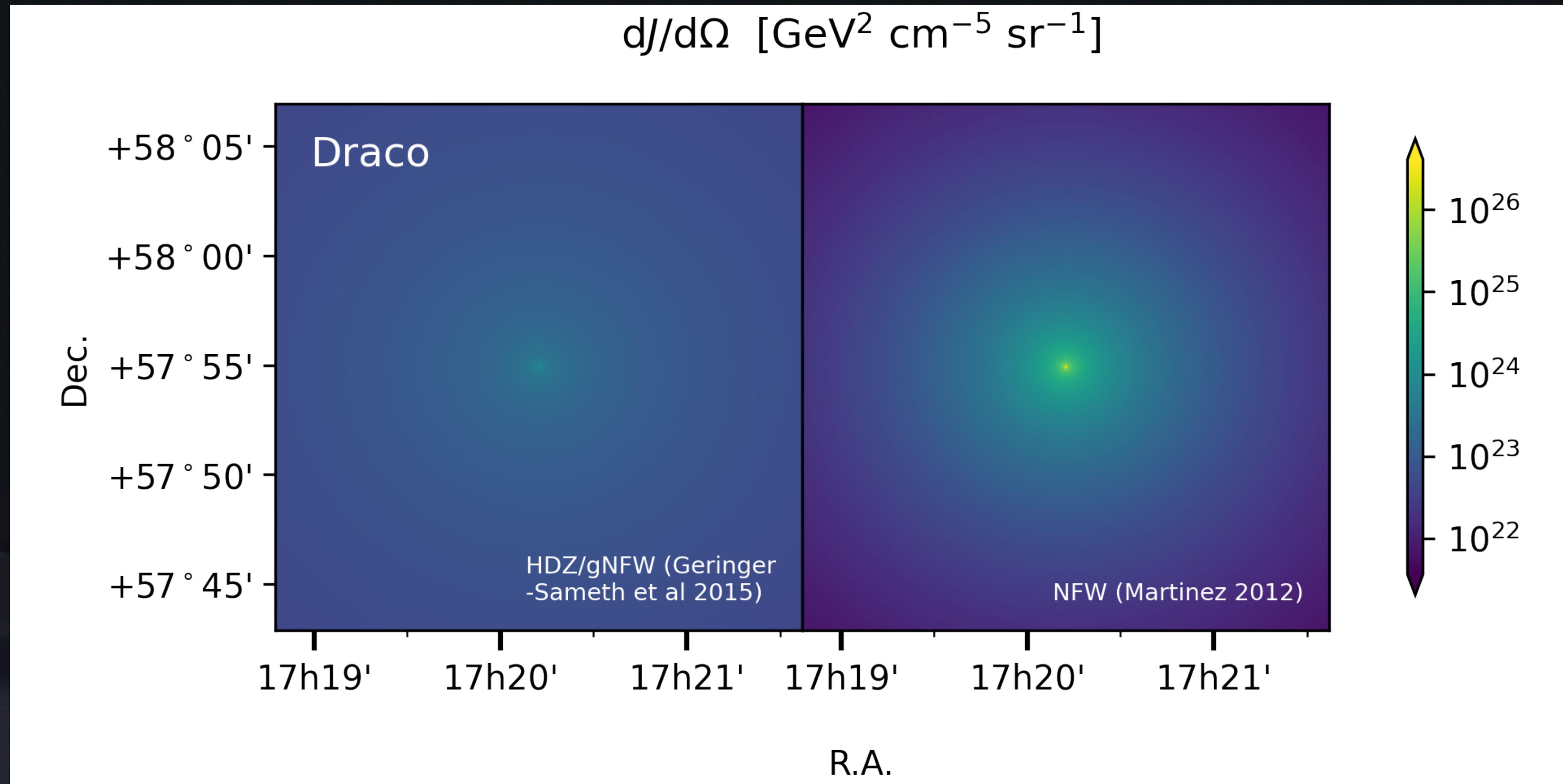
$$\Phi_\gamma = \frac{1}{8\pi m_\chi^2} \times J \times \frac{d\sigma\nu}{dE_\gamma}$$

$$J = \int d\Omega \int_{l.o.s.} ds \rho_{\text{DM}}^2$$

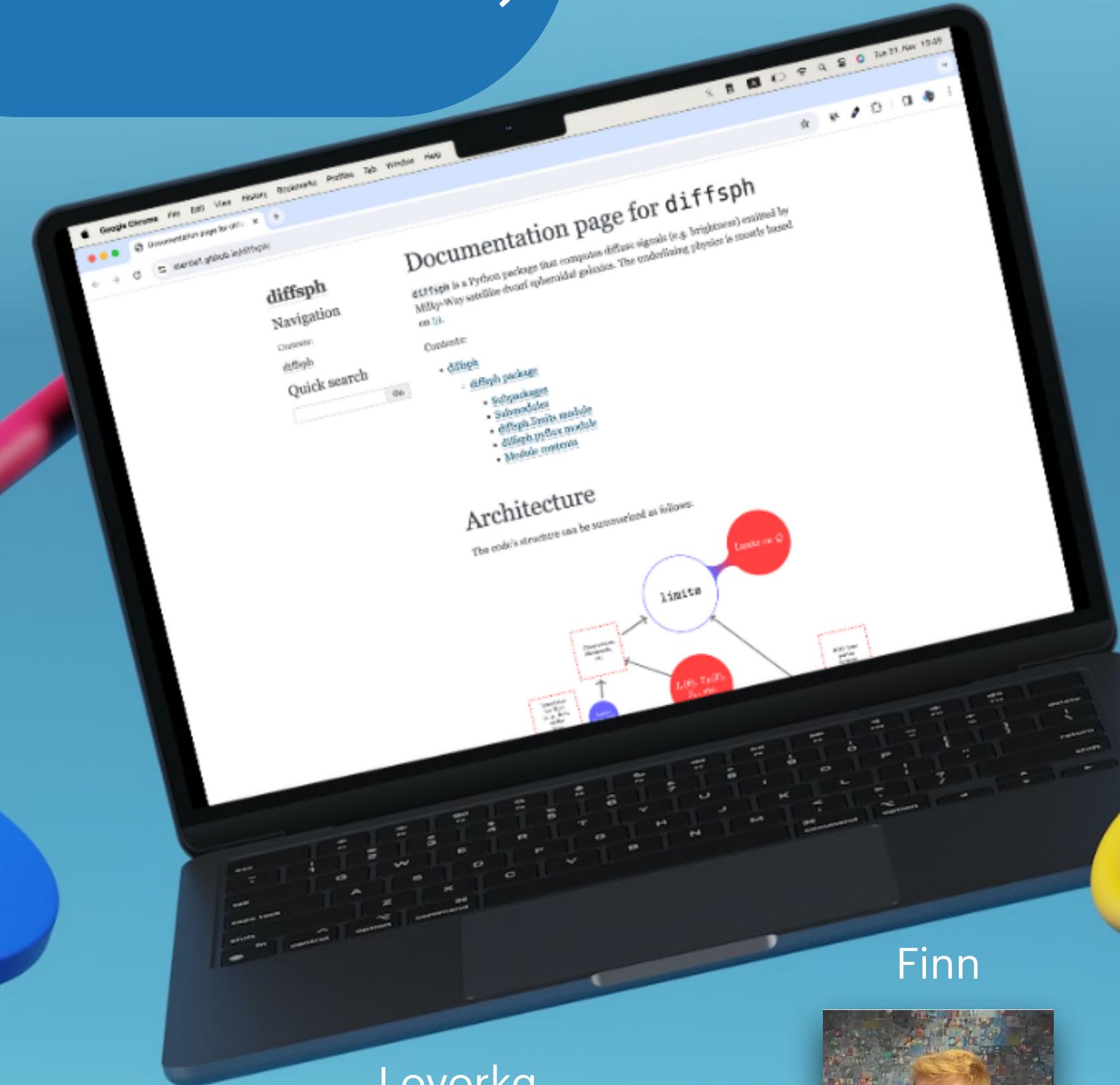


Example (Ad break)

Draco dwarf galaxy with diffSph [2401.05255]



diffSph (2401.05255)



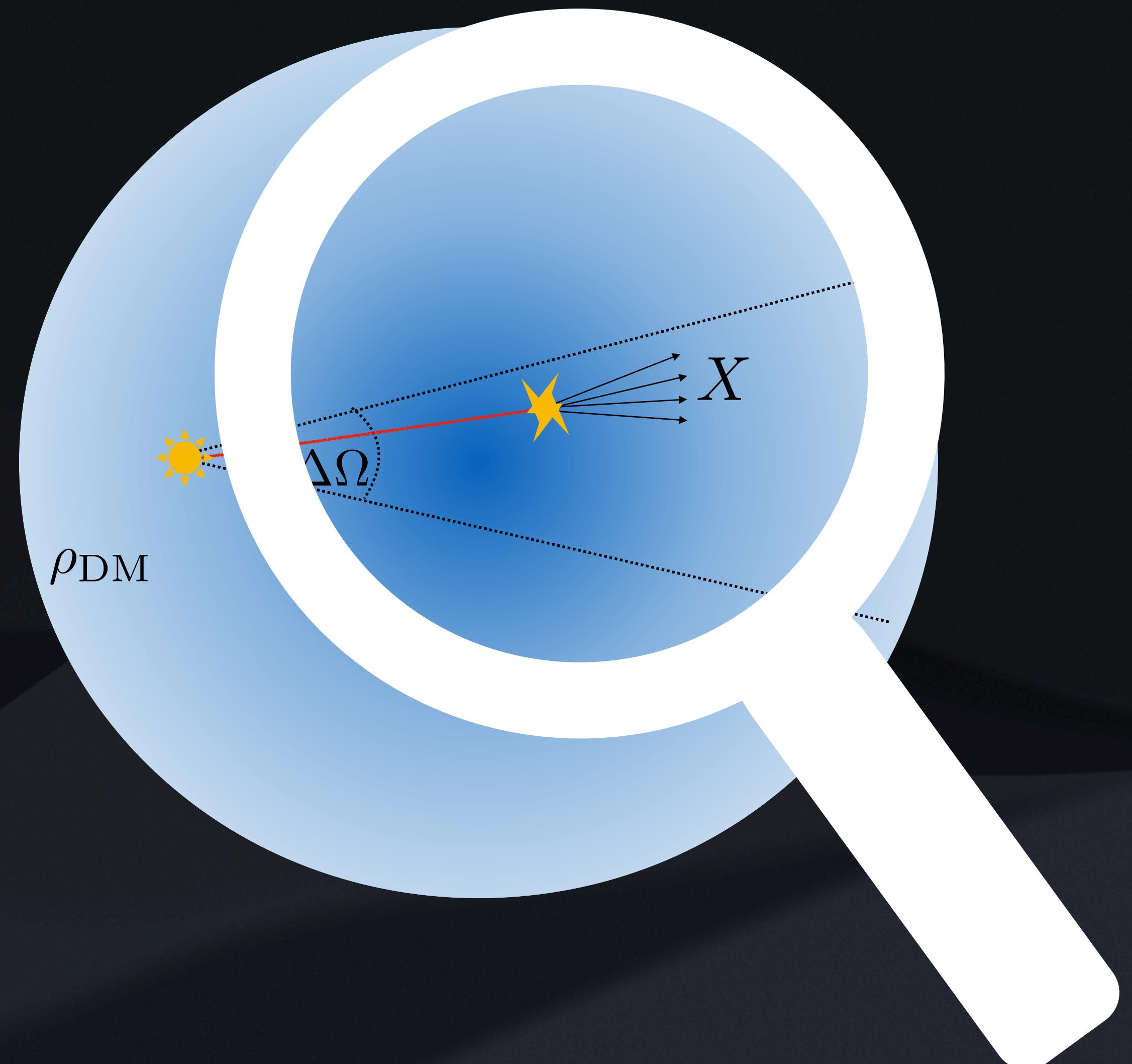
Finn

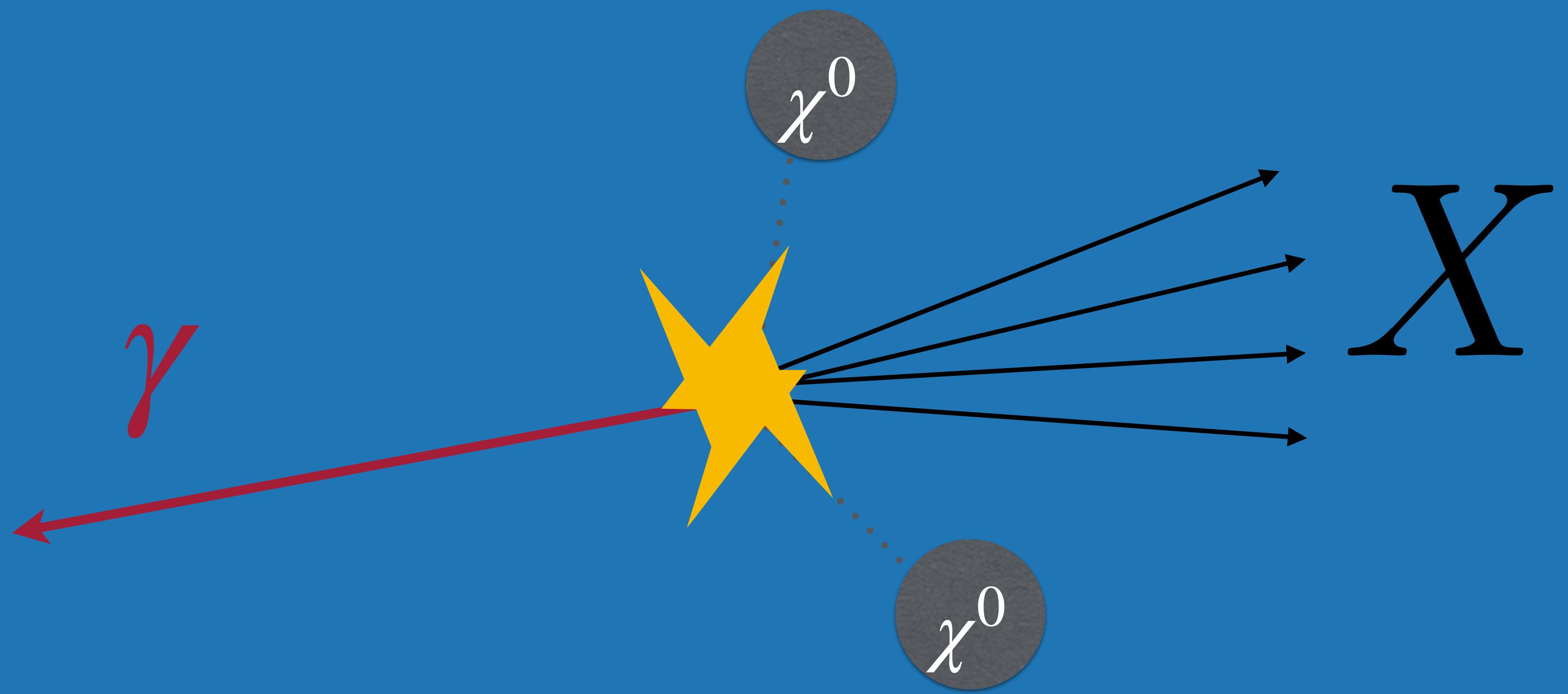


Lovorka

Gamma-ray flux formula

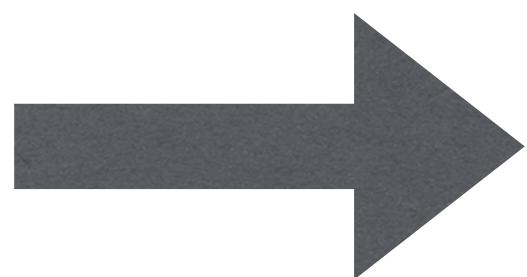
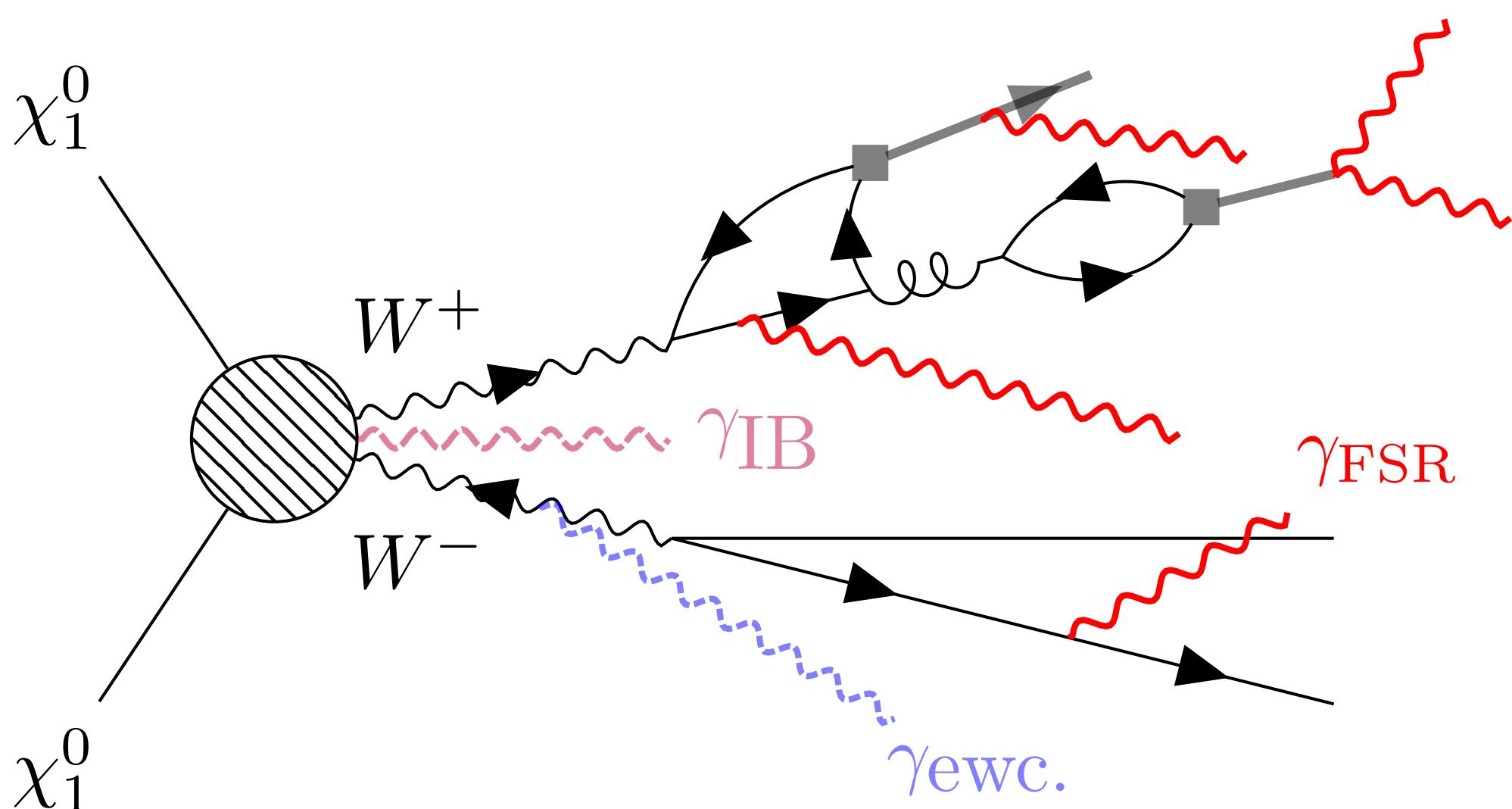
$$\Phi_\gamma = \frac{1}{8\pi m_\chi^2} \times J \times \frac{d\sigma\nu}{dE_\gamma}$$





The problem:

Obtain the annihilation spectrum



$$\frac{d\sigma\nu}{dE_\gamma}$$

TBA

The Basics of Annihilation

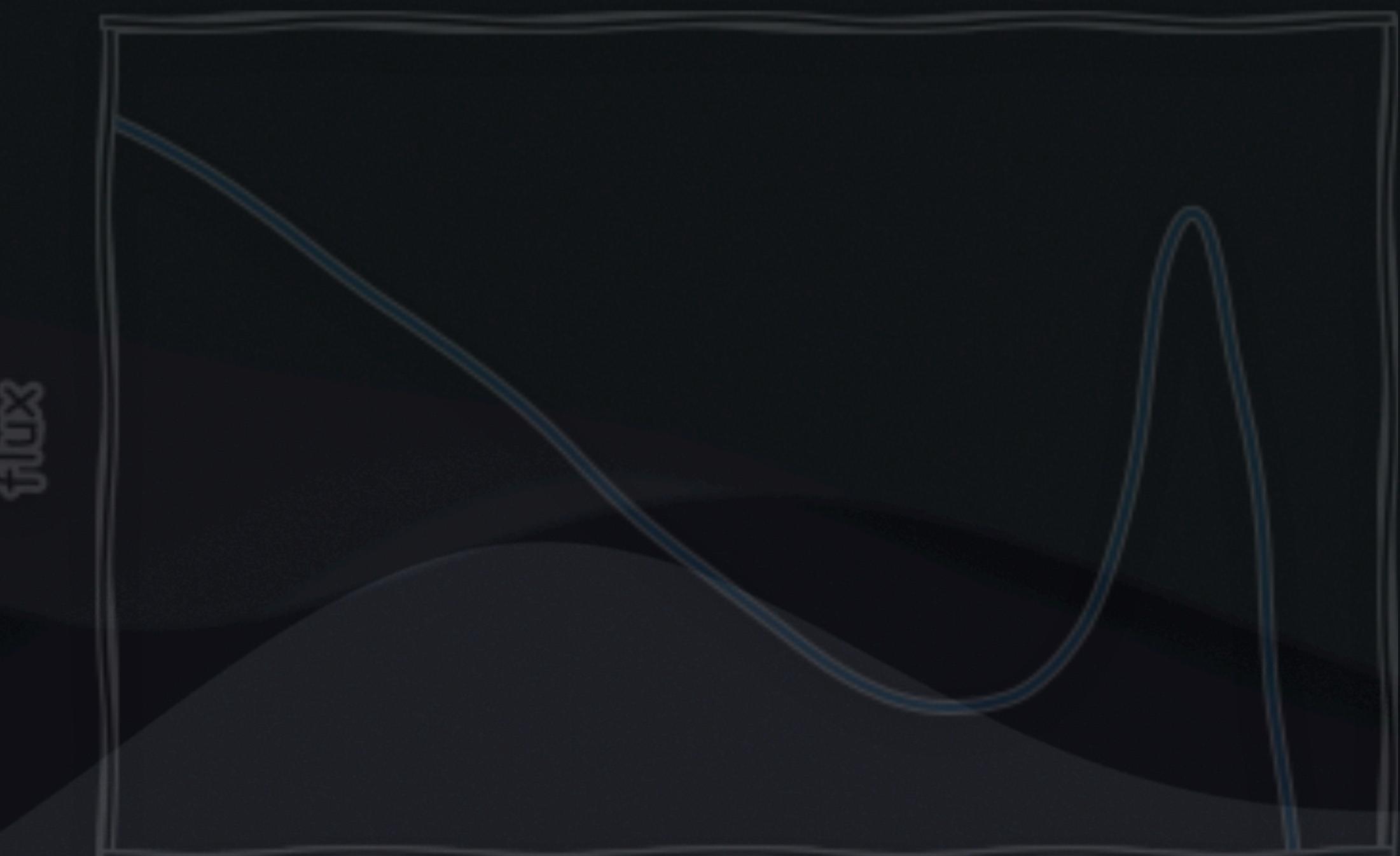


Fixed-order $2 \rightarrow 2$ (tree)
+ Parton Shower

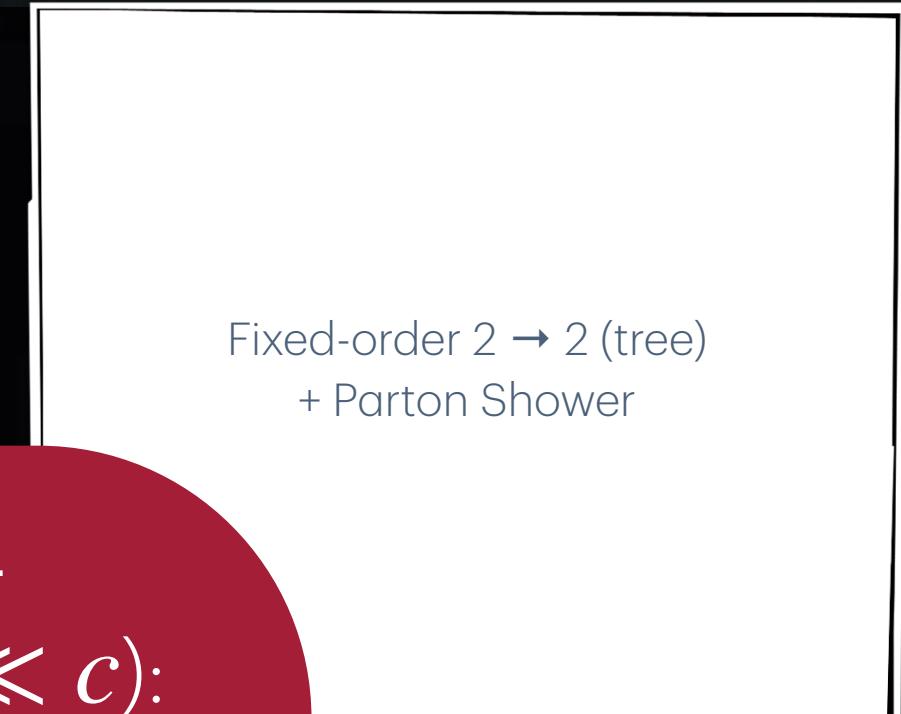
Fixed-order $2 \rightarrow 3$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 2$
+ Parton Shower +
Sommerfeld factor

Fixed-order $2 \rightarrow 3$
+ Parton Shower +
Sommerfeld factor



2000-2010s

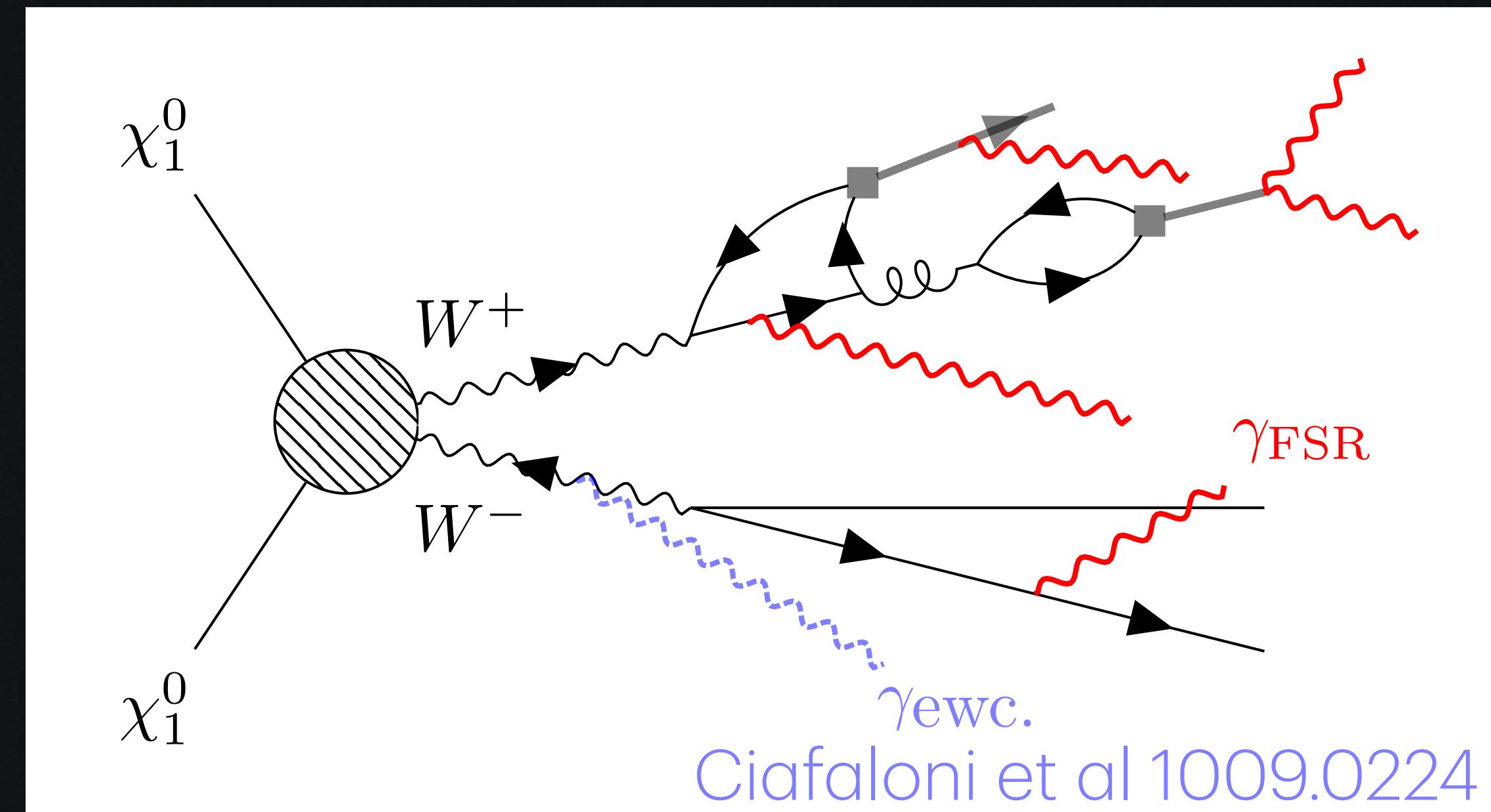


Helicity suppression

If $m_{DM} \gg m_b, m_\tau$ (and $v \ll c$):

$$\langle\sigma v\rangle_{b\bar{b}} \propto \frac{m_b^2}{m_{DM}^2} \rightarrow 0$$

$$\frac{d\sigma v}{dE_\gamma} = (\cancel{\langle\sigma v\rangle})_{\bar{b}b} \frac{dN_{\bar{b}b}^{\text{MC}}}{dE_\gamma} + (\cancel{\langle\sigma v\rangle})_{\tau^+\tau^-} \frac{dN_{\tau^+\tau^-}^{\text{MC}}}{dE_\gamma} + \dots$$



TBA

The Basics of Annihilation



Fixed-order $2 \rightarrow 2$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 3$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 2$
+ Parton Shower +
Sommerfeld factor

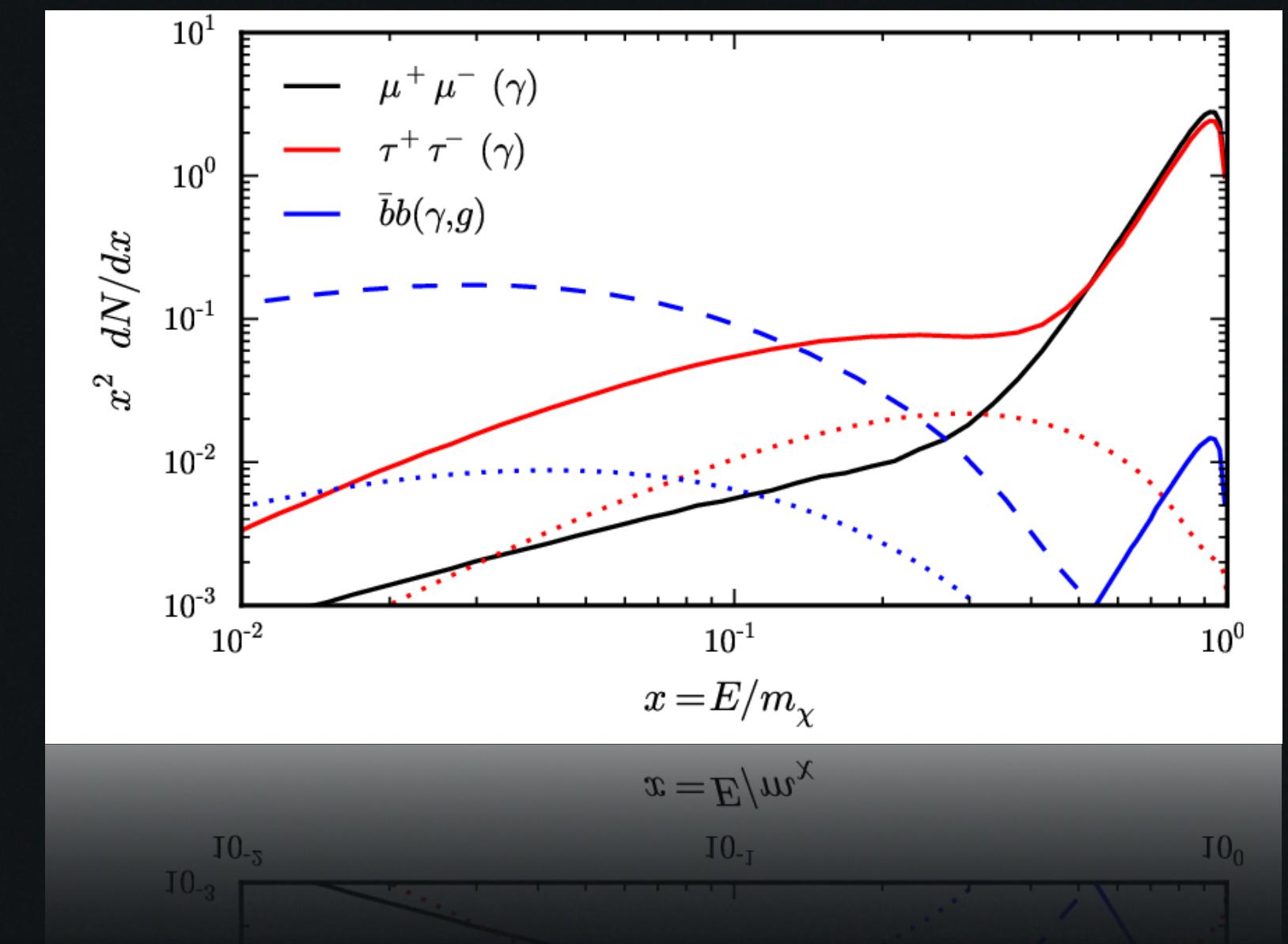
Fixed-order $2 \rightarrow 3$
+ Parton Shower +
Sommerfeld factor



2000-2010s



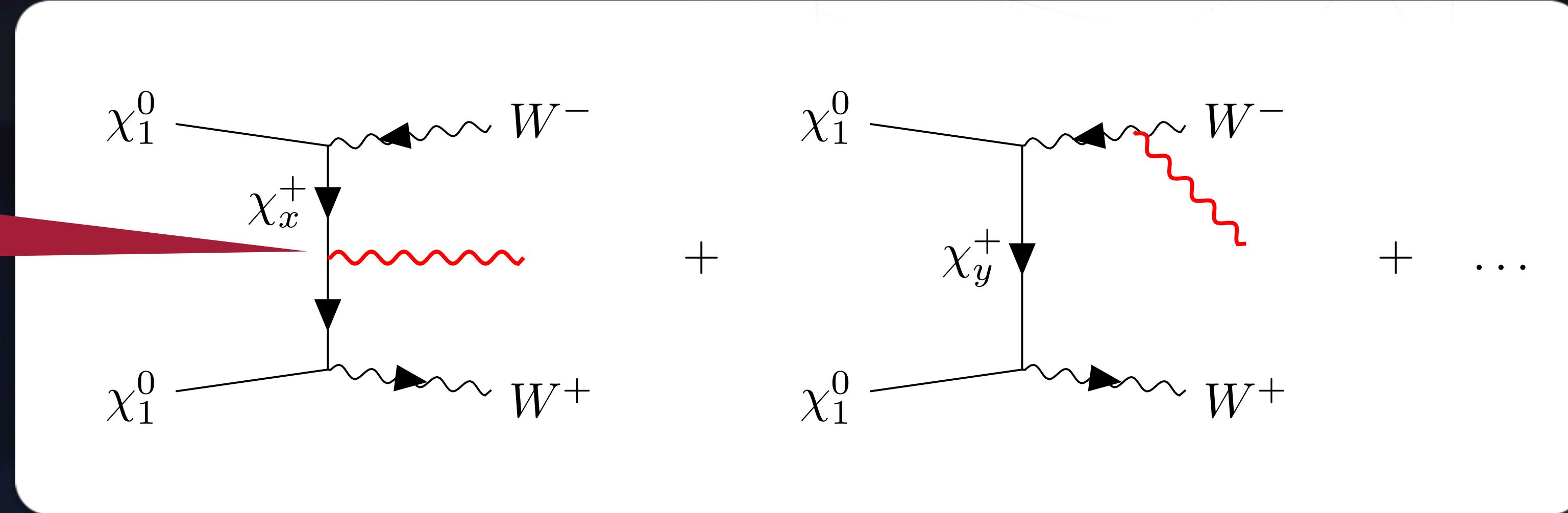
Fixed-order $2 \rightarrow 3$ (tree)
+ Parton Shower



Internal bremsstrahlung

Lift helicity suppression

- $2 \rightarrow 2$ process $\langle \sigma v \rangle_{b\bar{b}} \propto \frac{m_b^2}{m_{DM}^2} \rightarrow 0$
- $2 \rightarrow 3$ process $\langle \sigma v \rangle_{b\bar{b}\gamma} \neq 0$



Bringmann et al 0710.3169

TBA

The Basics of Annihilation

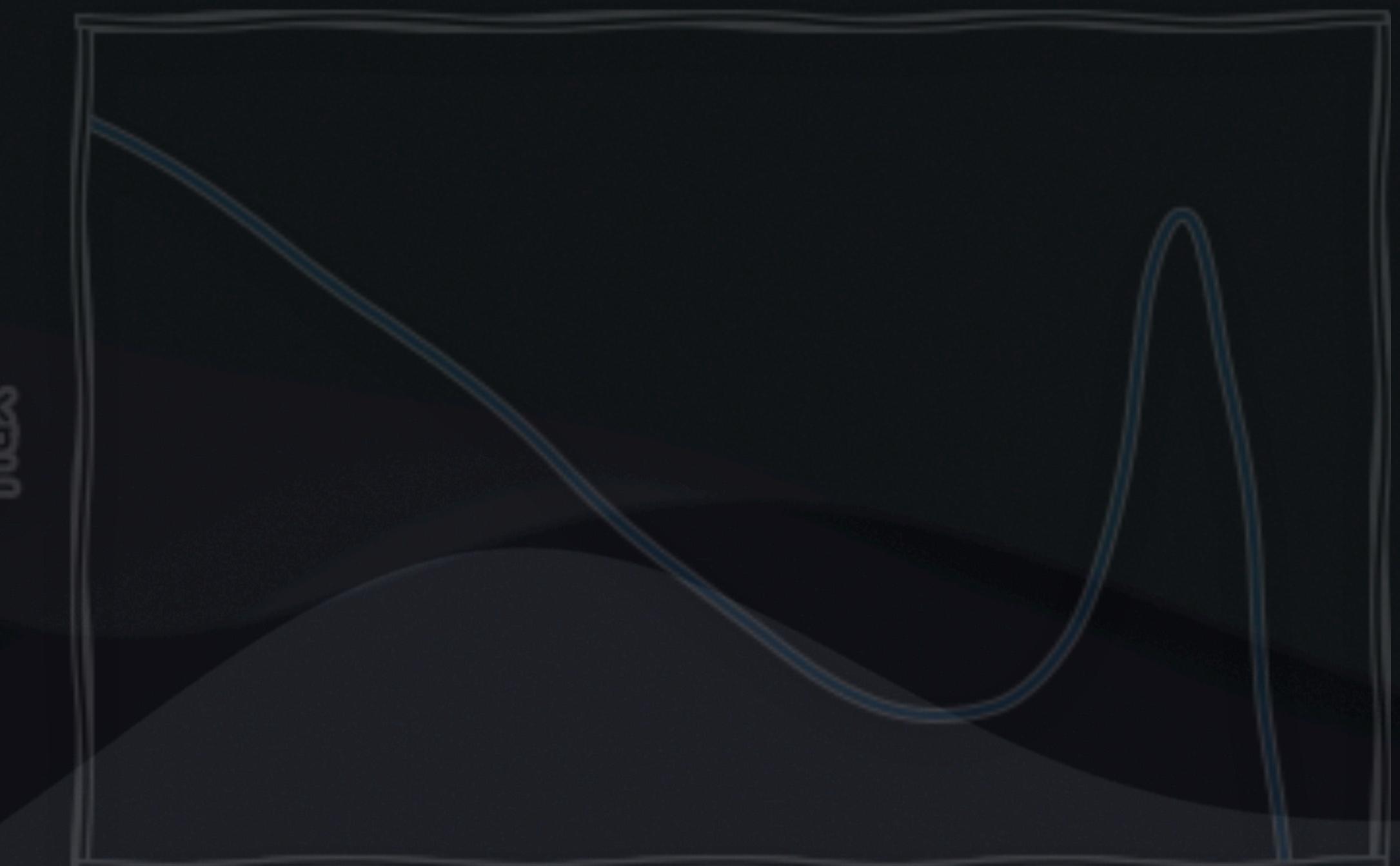


Fixed-order $2 \rightarrow 2$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 3$ (tree)
+ Parton Shower

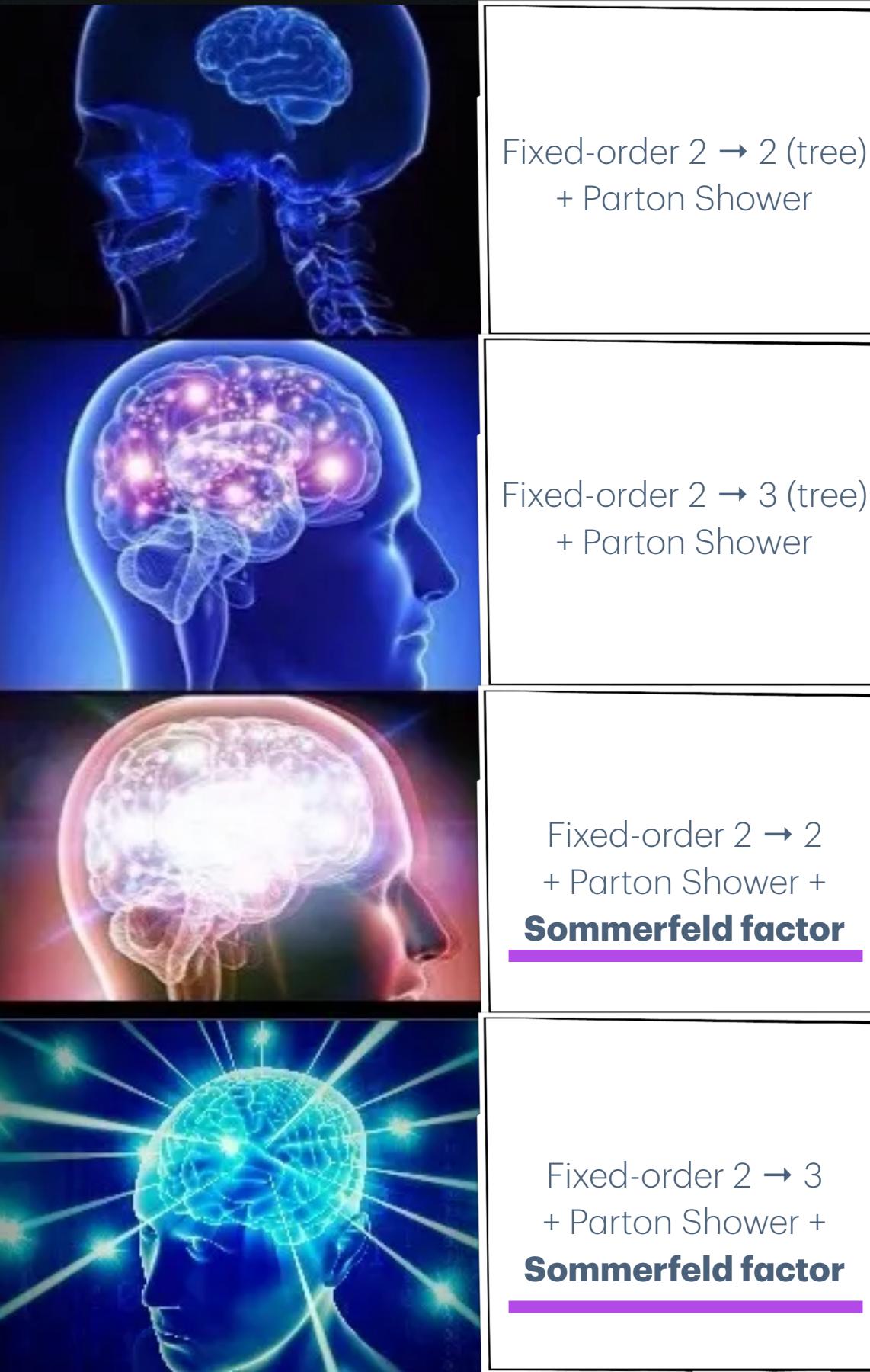
Fixed-order $2 \rightarrow 2$
+ Parton Shower +
Sommerfeld factor

Fixed-order $2 \rightarrow 3$
+ Parton Shower +
Sommerfeld factor



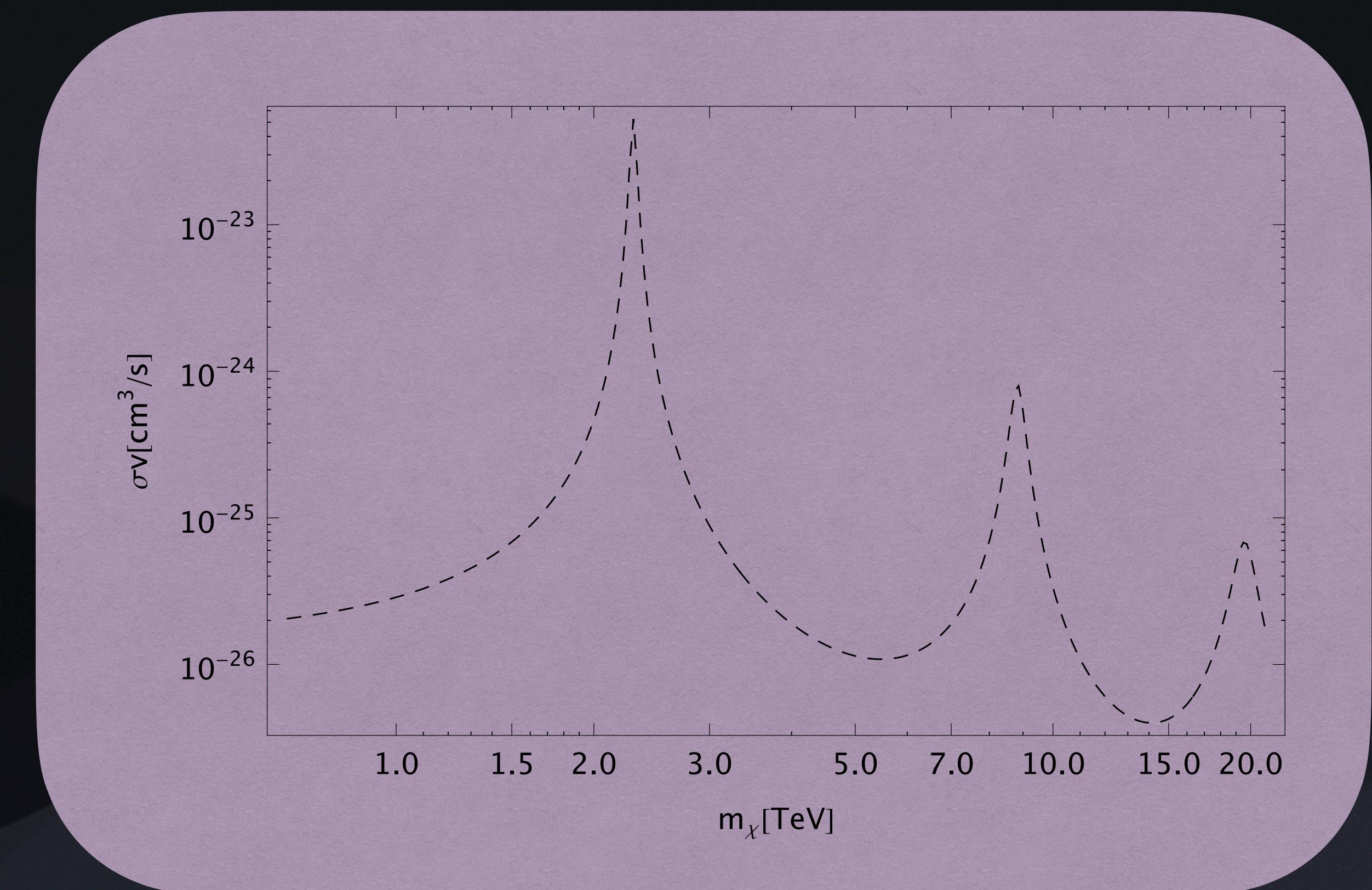
TBA

The Basics of Annihilation

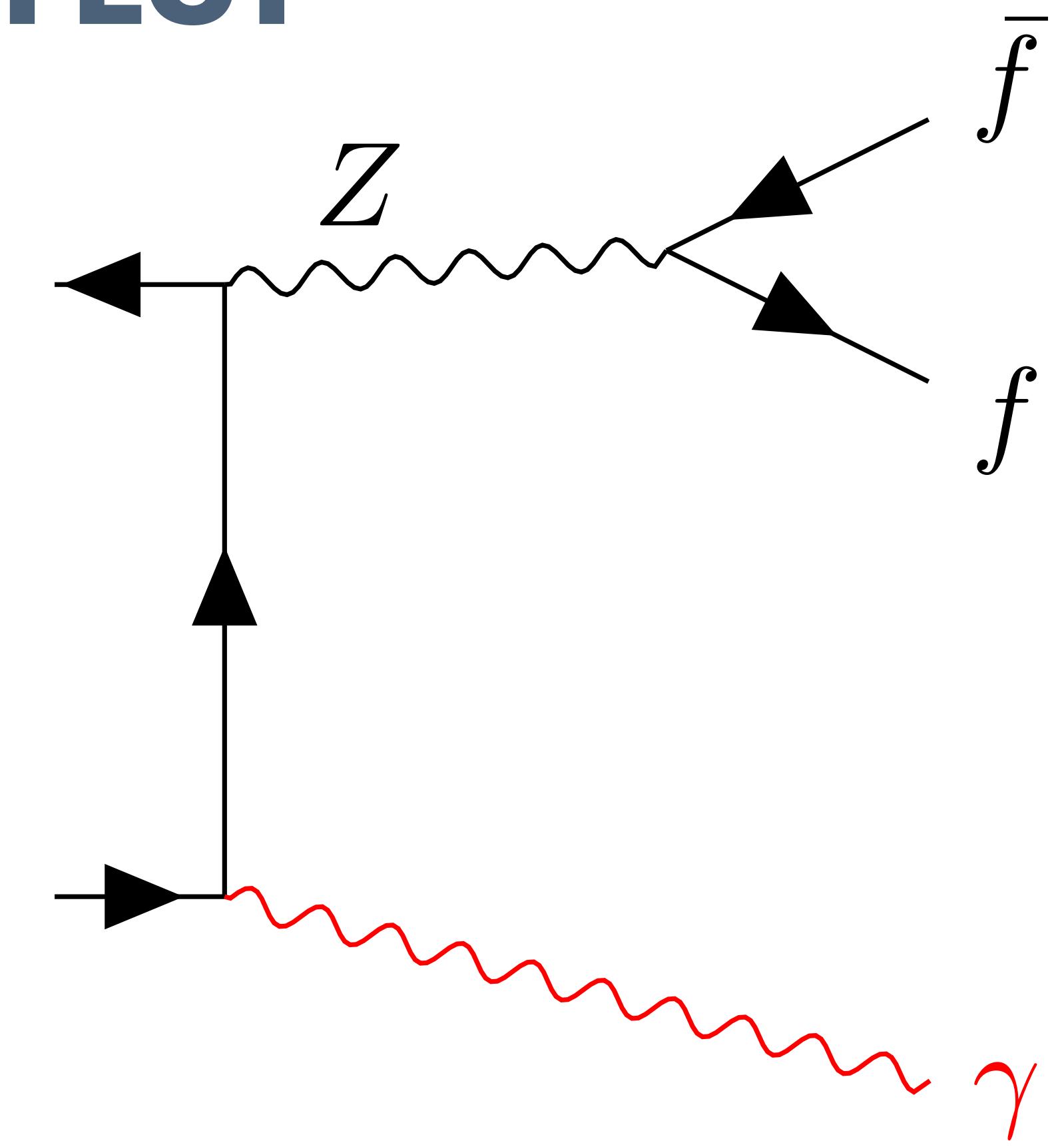
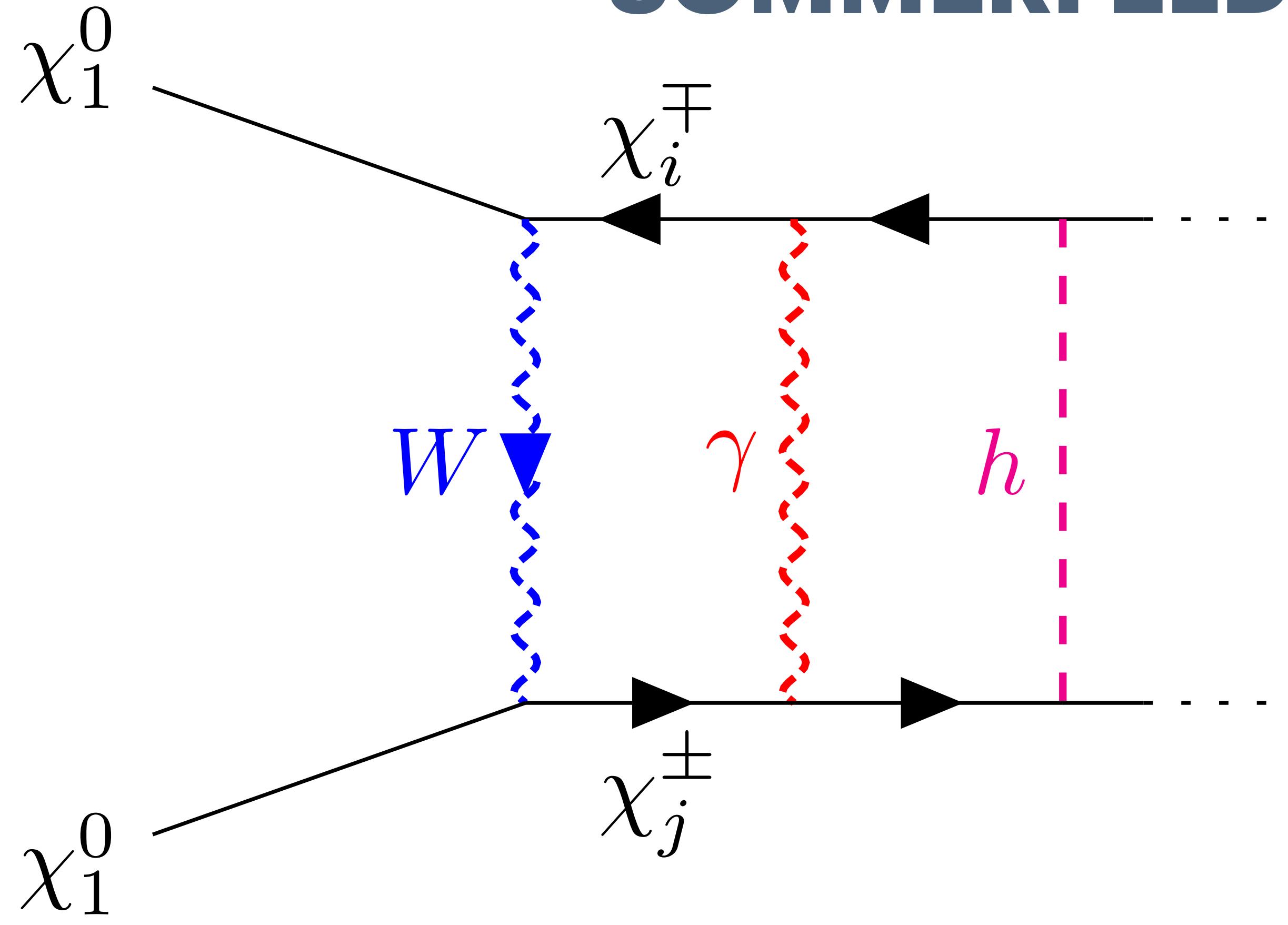


Hisano, Matsumoto, Nojiri [hep-ph/0307216](#)

Arkani-Hamed, Slatyer, Weiner [0810.0713](#)



SOMMERFELD EFFECT



Resummation

Breakdown of perturbative expansion when $m_\chi \gg m_W$

$$\frac{d\sigma v}{dE_\gamma} \sim \# \alpha_{ew}^3 + \alpha_{ew}^4 \left(\text{Silhouette of Gorilla A} + \dots \right) + \alpha_{ew}^5 \left(\text{Silhouette of Gorilla B} + \dots \right) + \dots$$

Tree (LO) 1-loop (NLO) 2-loop (NNLO)

The diagram illustrates the breakdown of a perturbative expansion for the differential cross-section $d\sigma v / dE_\gamma$. It shows three terms: a tree-level term (a single gorilla), a one-loop term (two gorillas), and a two-loop term (three gorillas). Red annotations are present in the one-loop and two-loop terms: a red '#' is placed near the top of the first gorilla's back in the one-loop term; a red 'X' is placed near the top of the second gorilla's back in the two-loop term; and a red minus sign '-' is placed between the two gorillas in the two-loop term.

Resummation

Goal: factor out the “gorillas”

$$\frac{d\sigma_v}{dE_\gamma} = f(\alpha_{ew} \times \text{gorilla}) \times (\# \alpha_{ew}^3 + \mathcal{O}(\alpha_{ew}^4))$$

“safe” to use perturbation theory

TBA

The Basics of Annihilation



Fixed-order $2 \rightarrow 2$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 3$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 2$
+ Parton Shower +
Sommerfeld factor

Fixed-order $2 \rightarrow 3$
+ Parton Shower +
Sommerfeld factor



2020s (before our paper came out)



Fixed-order $2 \rightarrow 2$
+ Parton Shower +
Sommerfeld factor

- Incomplete (missing shower for e.g. $\chi^0 \chi^0 \rightarrow H^\pm W^\mp$)
- Helicity-suppressed cross sections still suppressed
- ...

Fixed-order $2 \rightarrow 3$
+ Parton Shower +
Sommerfeld factor

- Only extrapolations from our endpoint factorization formulas available and for pure wino/higgsino

Outline

Motivation

Indirect detection

Sommerfeld factor

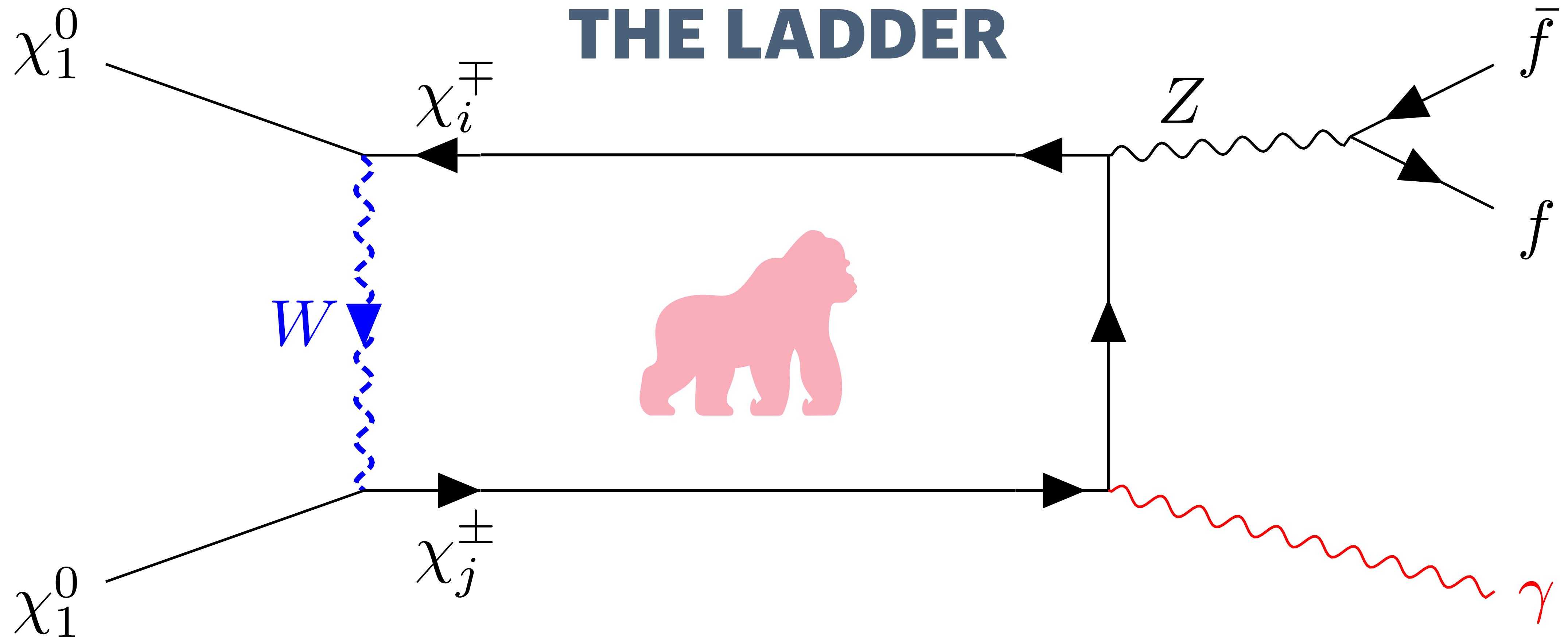
Endpoint resummations

Plots

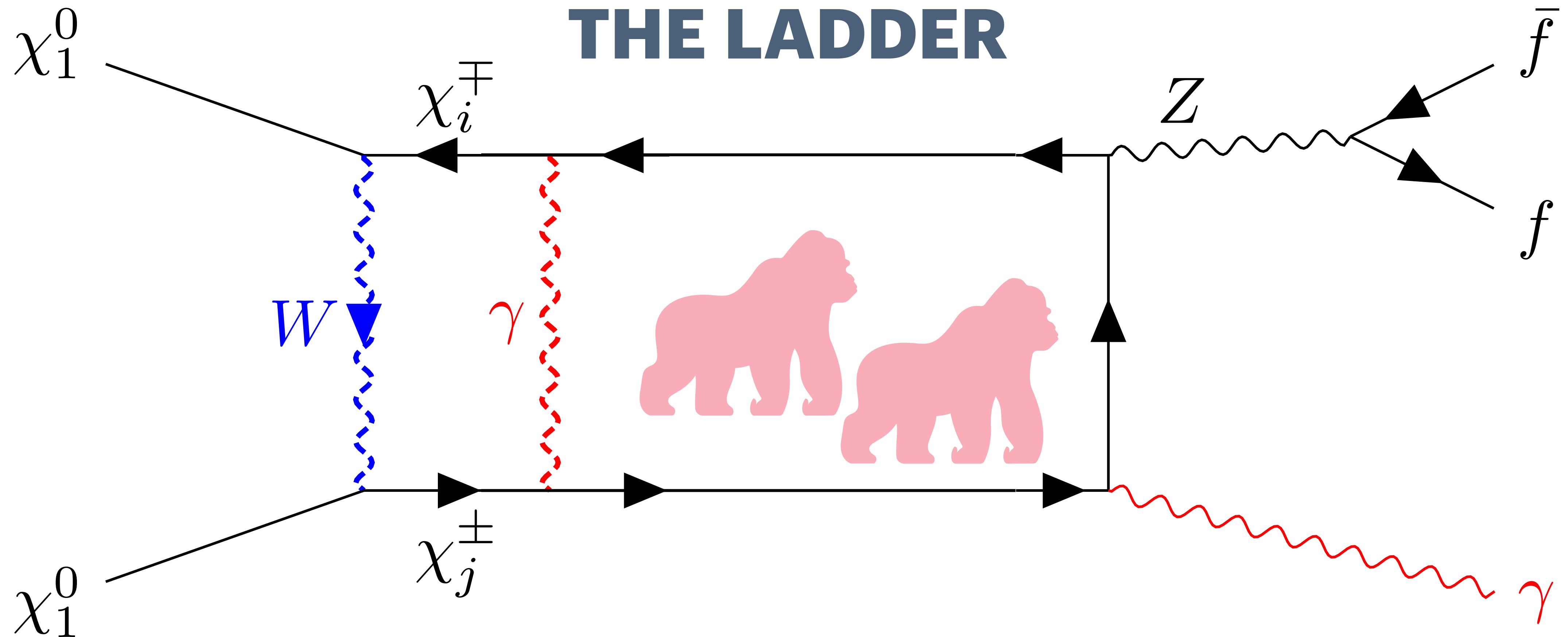
Conclusions

Sommerfeld factor

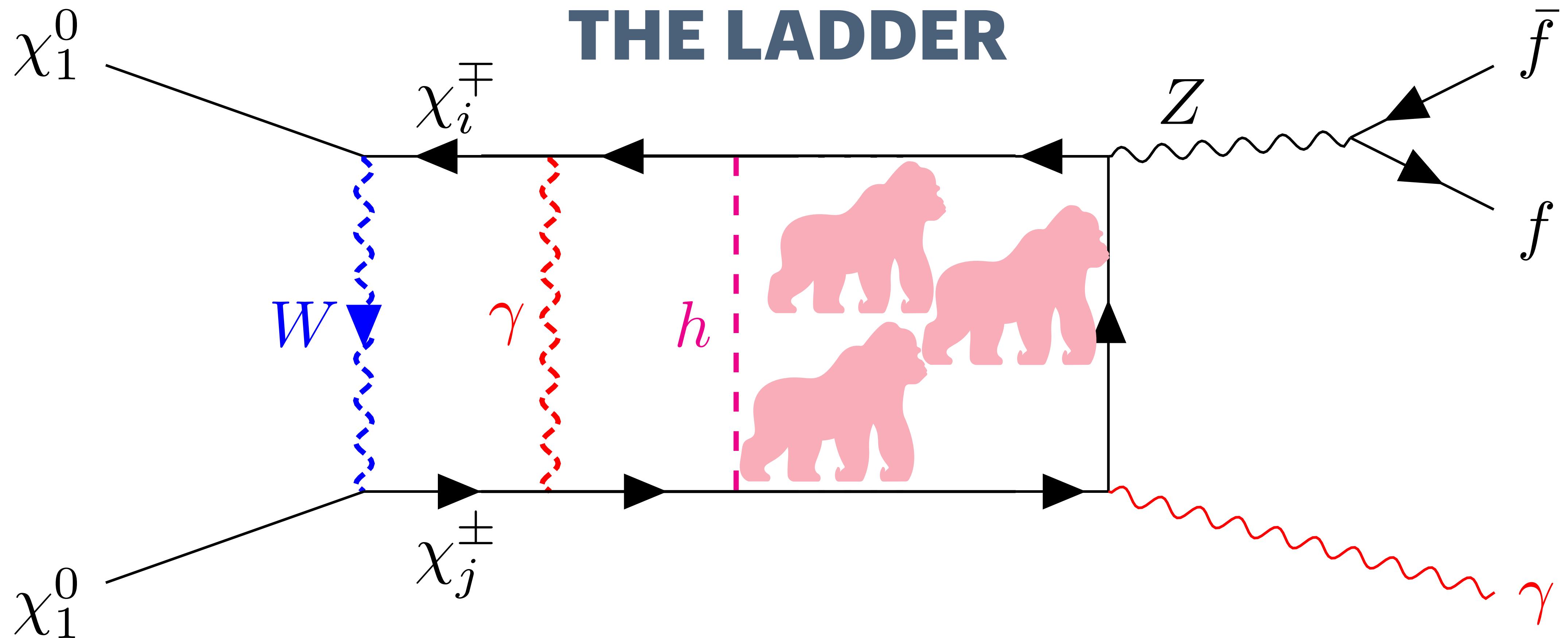
THE LADDER



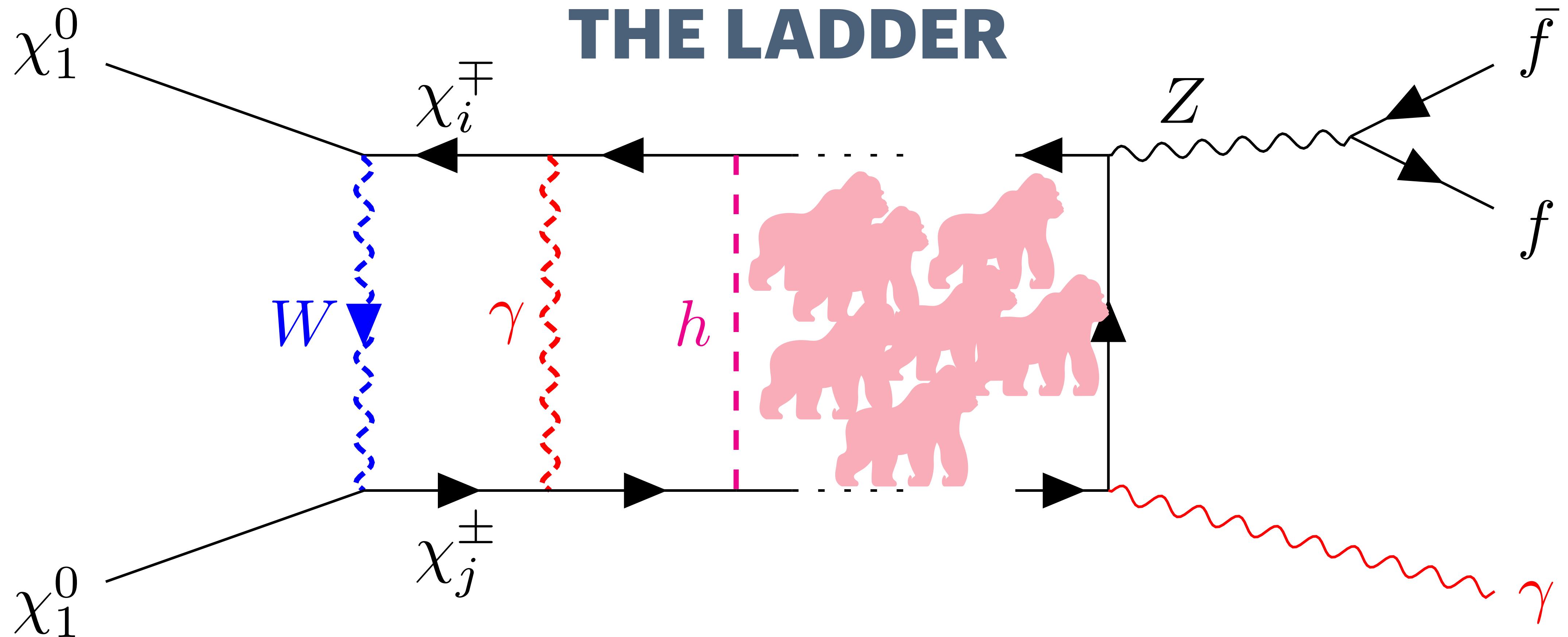
THE LADDER



THE LADDER



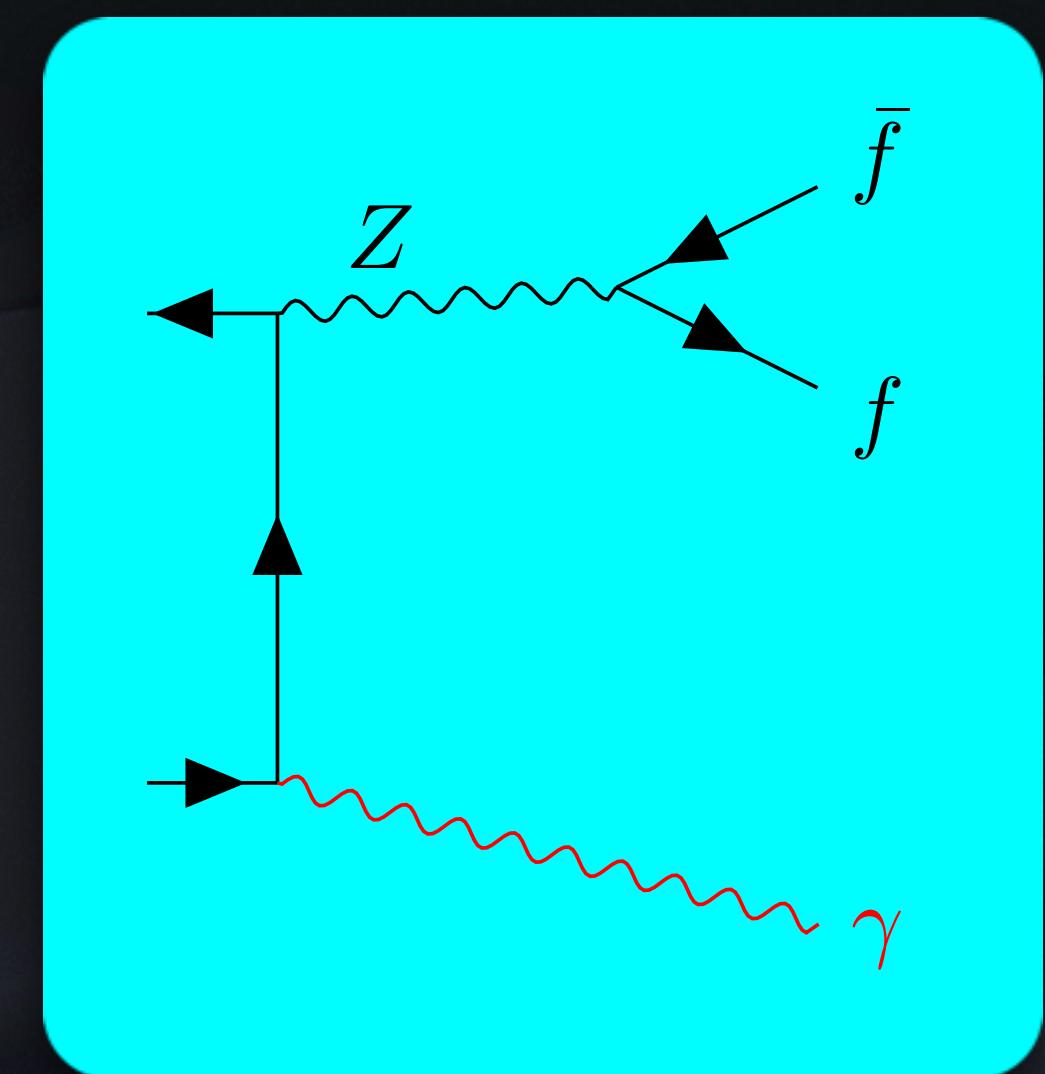
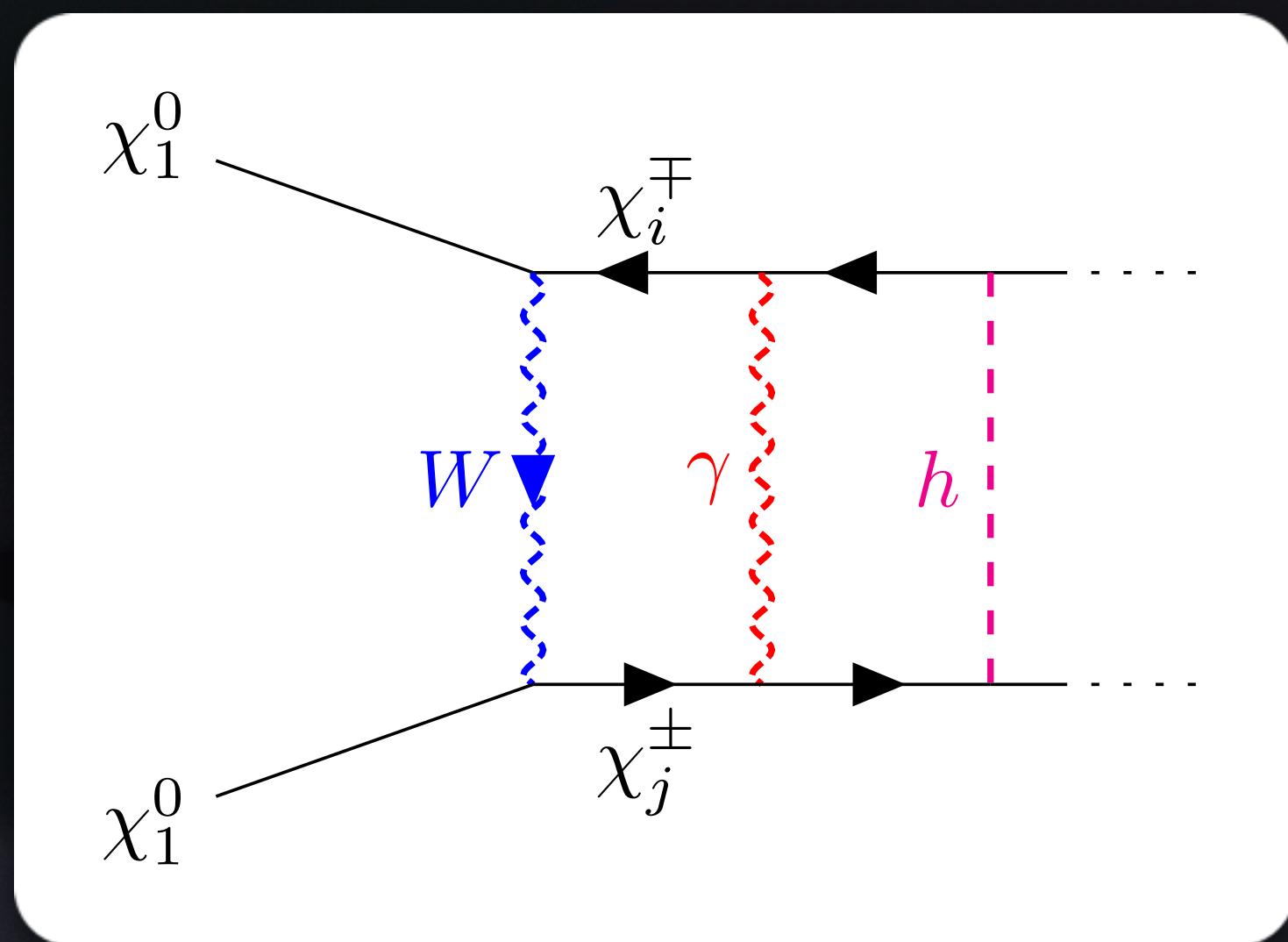
THE LADDER



Resummation

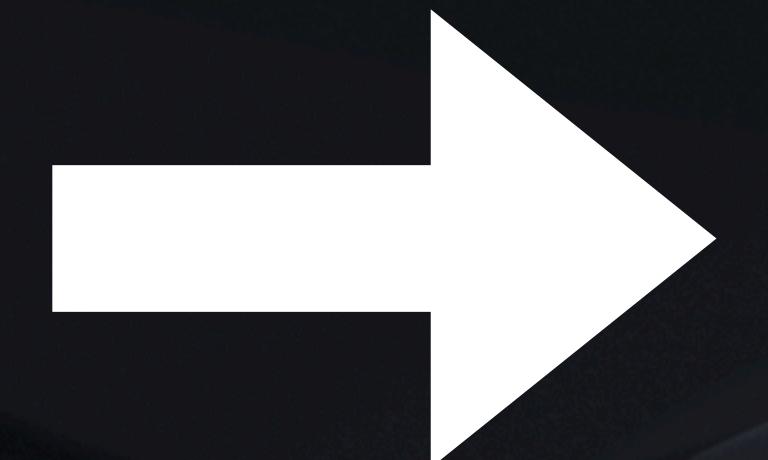
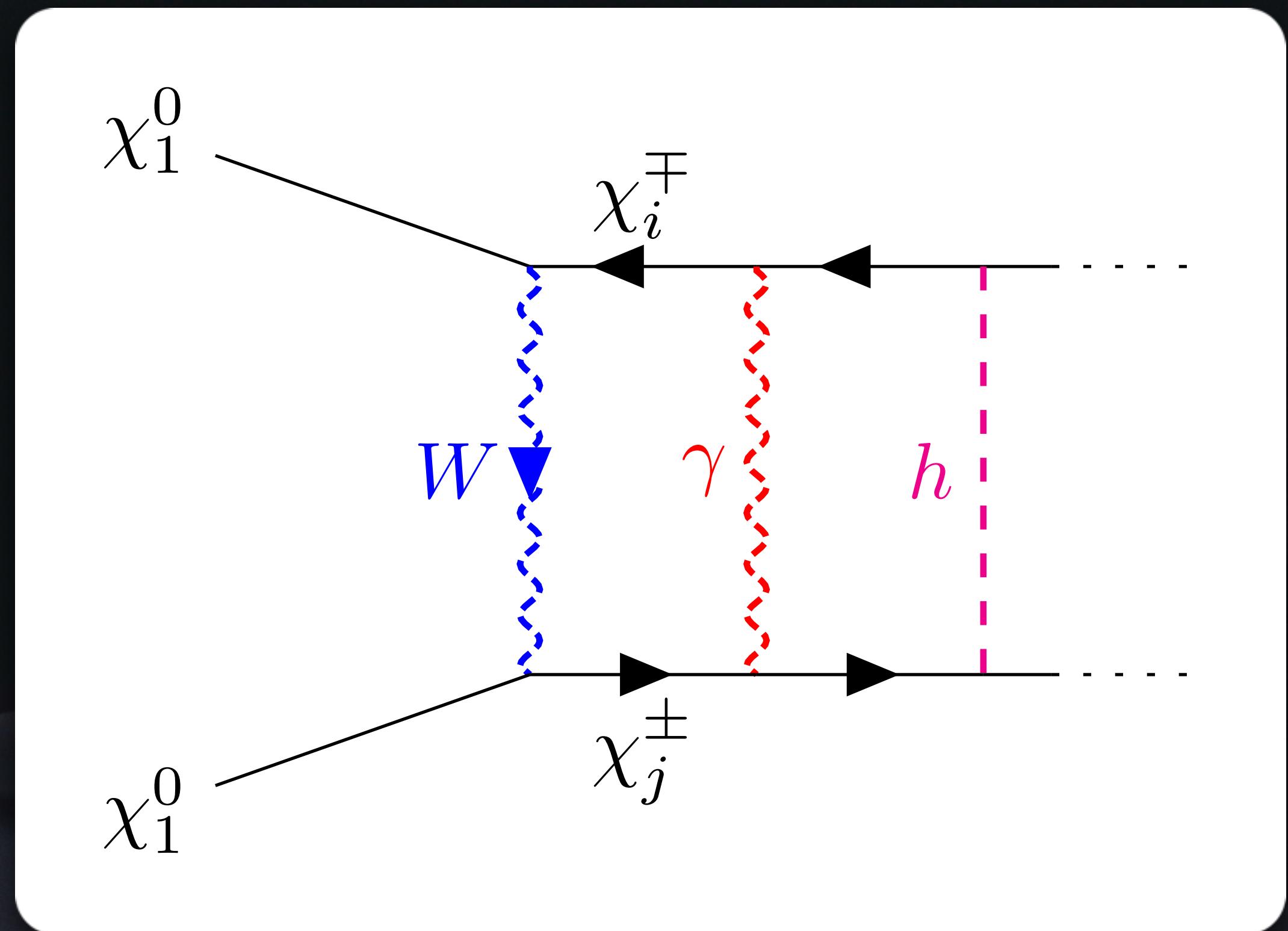
Goal: factor out the “gorillas”

$$\frac{d\sigma v}{dE_\gamma} = f(\alpha_{ew} \times \text{gorilla}) \times (\# \alpha_{ew}^3 + \mathcal{O}(\alpha_{ew}^4))$$



Sommerfeld factor

Reframe the question: QFT  Quantum Mechanics



**Non-relativistic Yukawa
potential-like interactions!!!!**



Sommerfeld factor

Putting all things together

$$\frac{d\sigma\nu}{dE_\gamma} = 2 \sum_{I,J} S_{IJ} \left[\frac{d(\tilde{\sigma}\nu)}{dE_\gamma} \right]_{IJ}$$

Quantum mechanics
(ladder)

Particle pairs

QFT perturbation theory

Go to the blackboard I

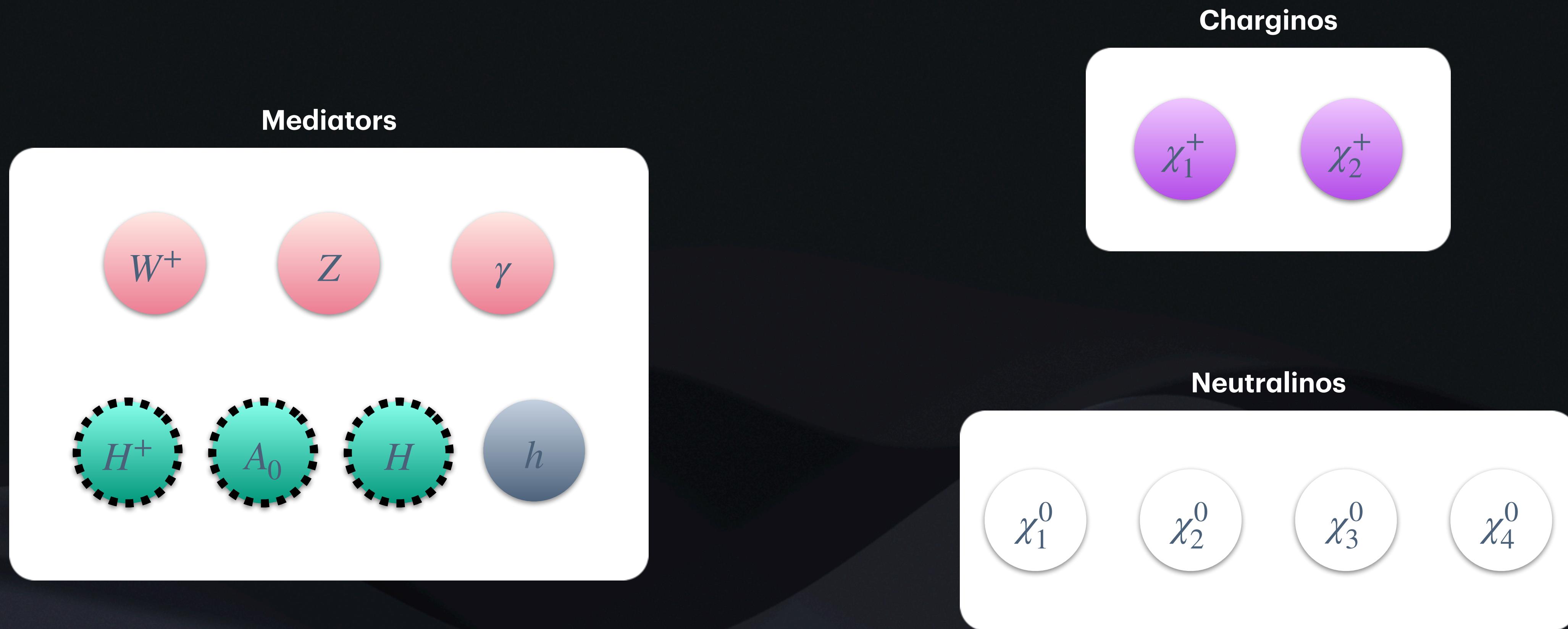
The MSSM in a nutshell

Particles and more particles



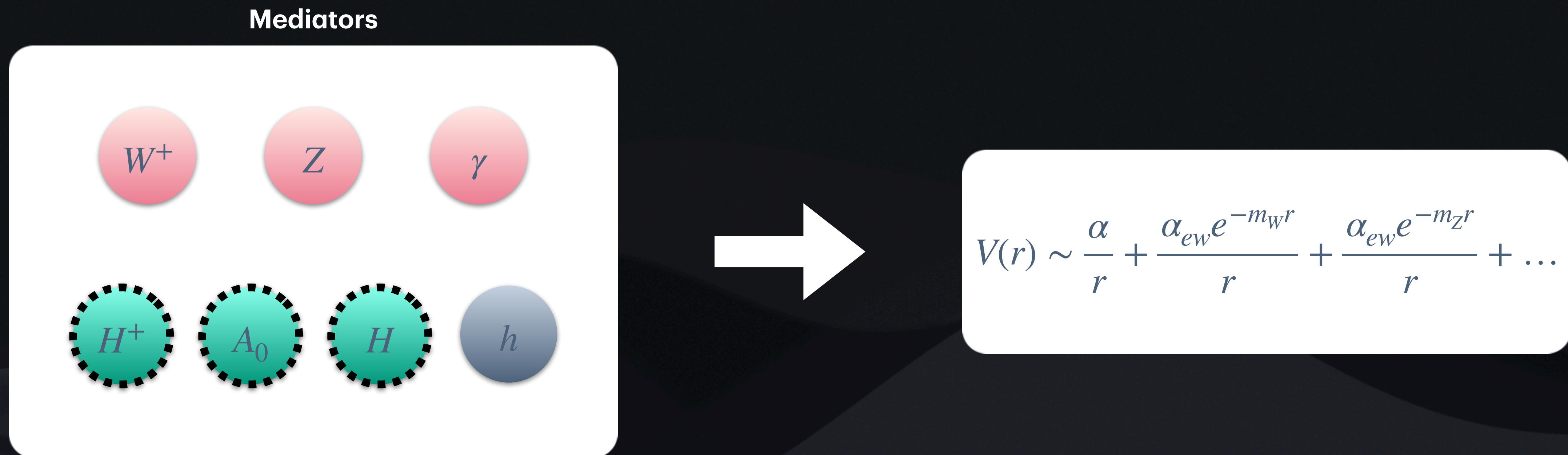
The MSSM in a nutshell

Neutralino/chargino sector



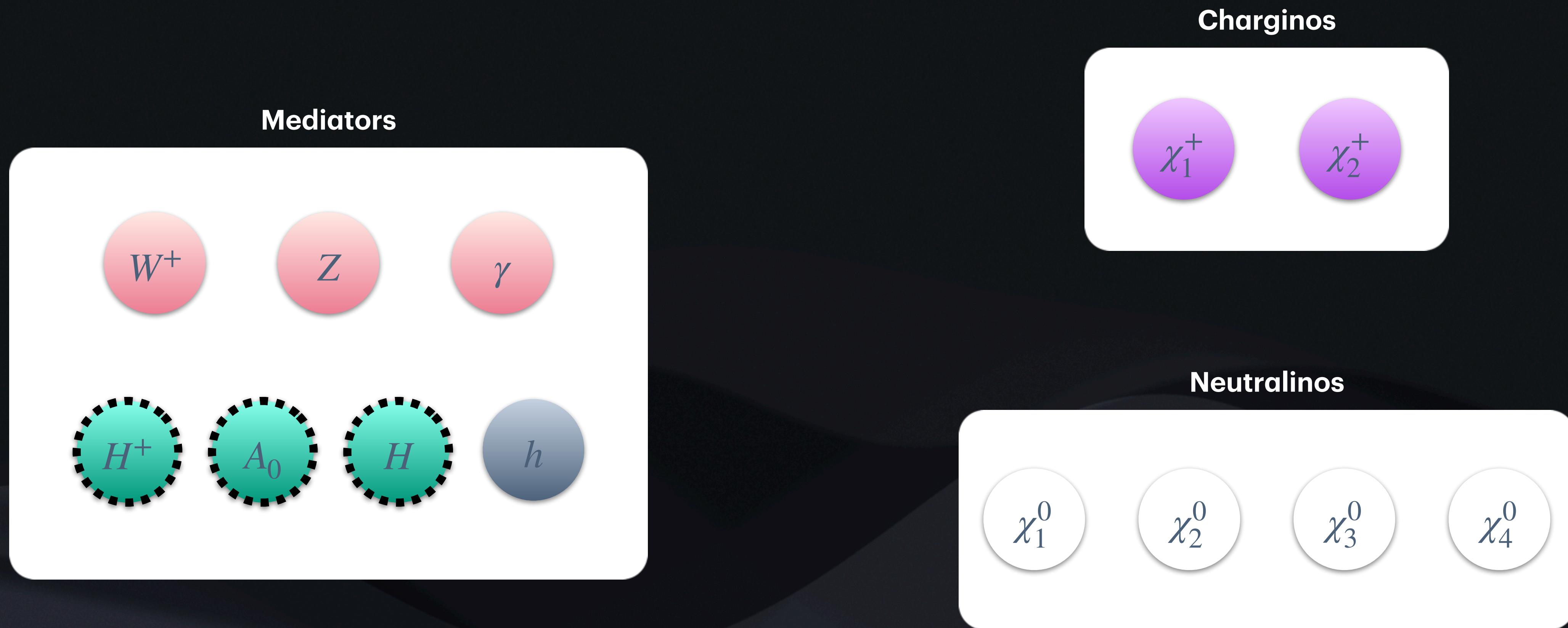
Sommerfeld effect in the MSSM

Dark Matter NR potential



Sommerfeld effect in the MSSM

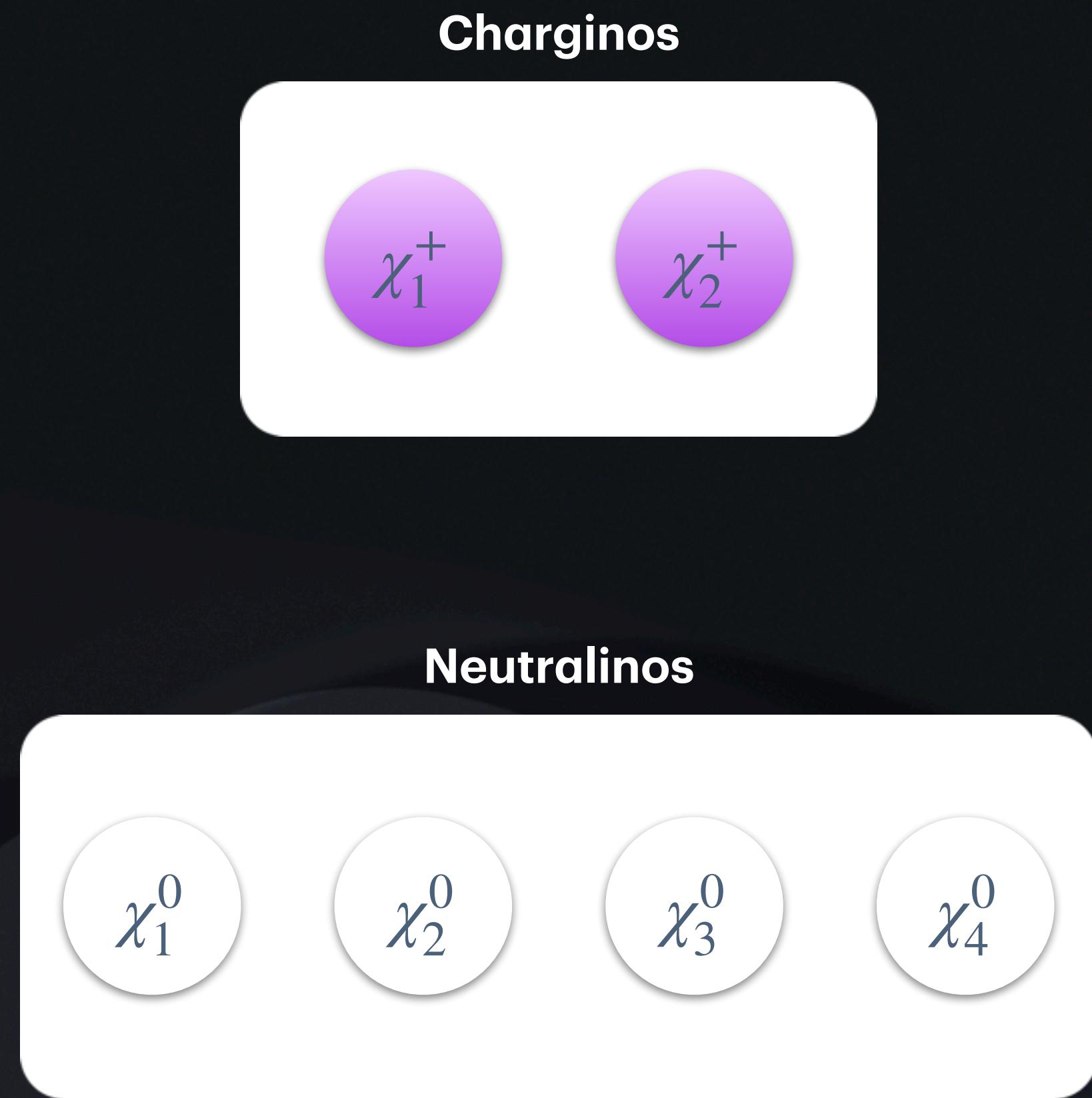
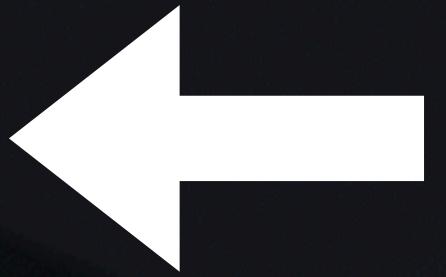
Dark Matter NR potential



Sommerfeld effect in the MSSM

Dark Matter NR potential

$$V(r) \rightarrow V_{IJ}(r)$$



Sommerfeld effect in the MSSM

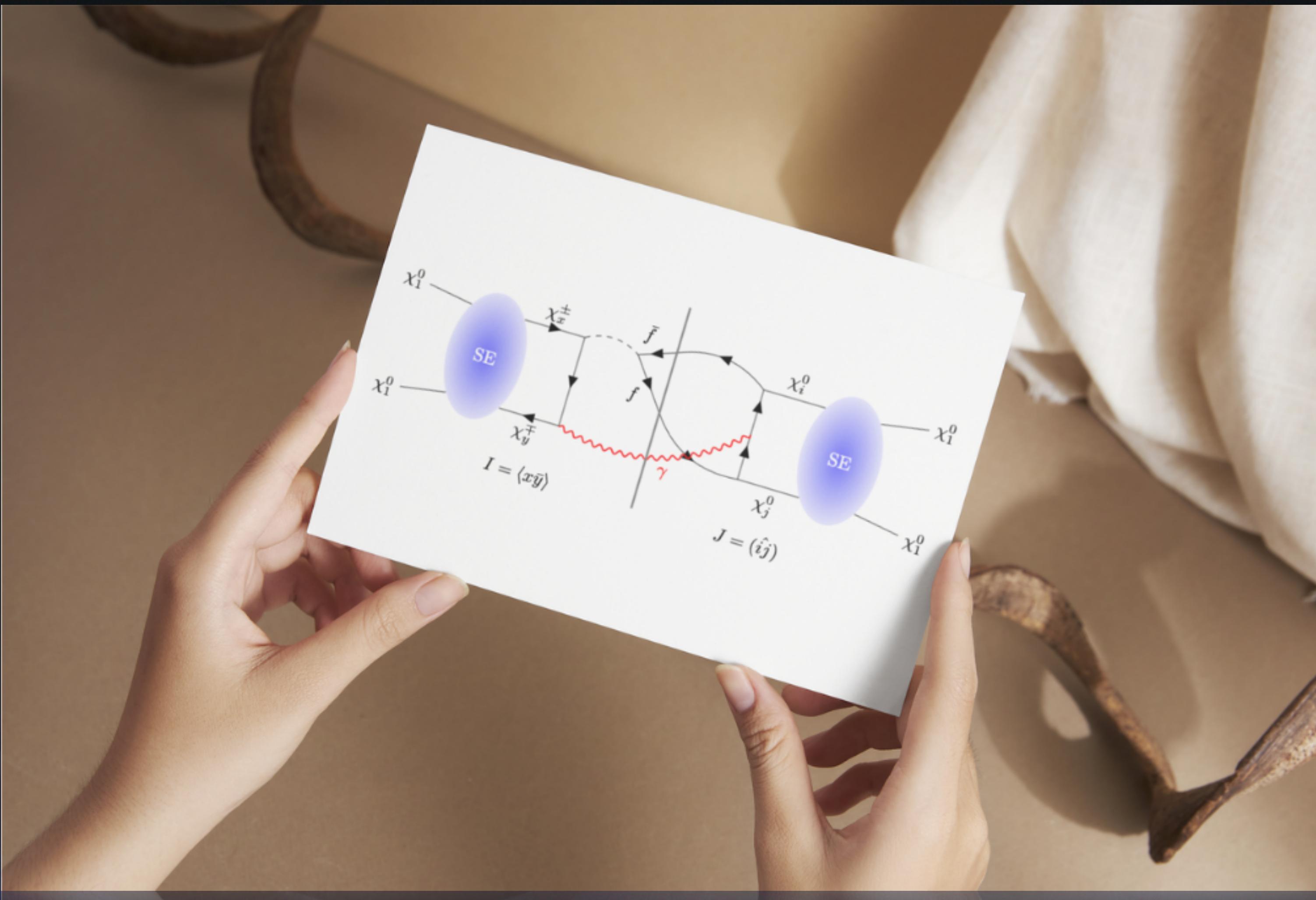
Dark Matter NR potential

Sommerfeld effect in the MSSM

14×14 matrix

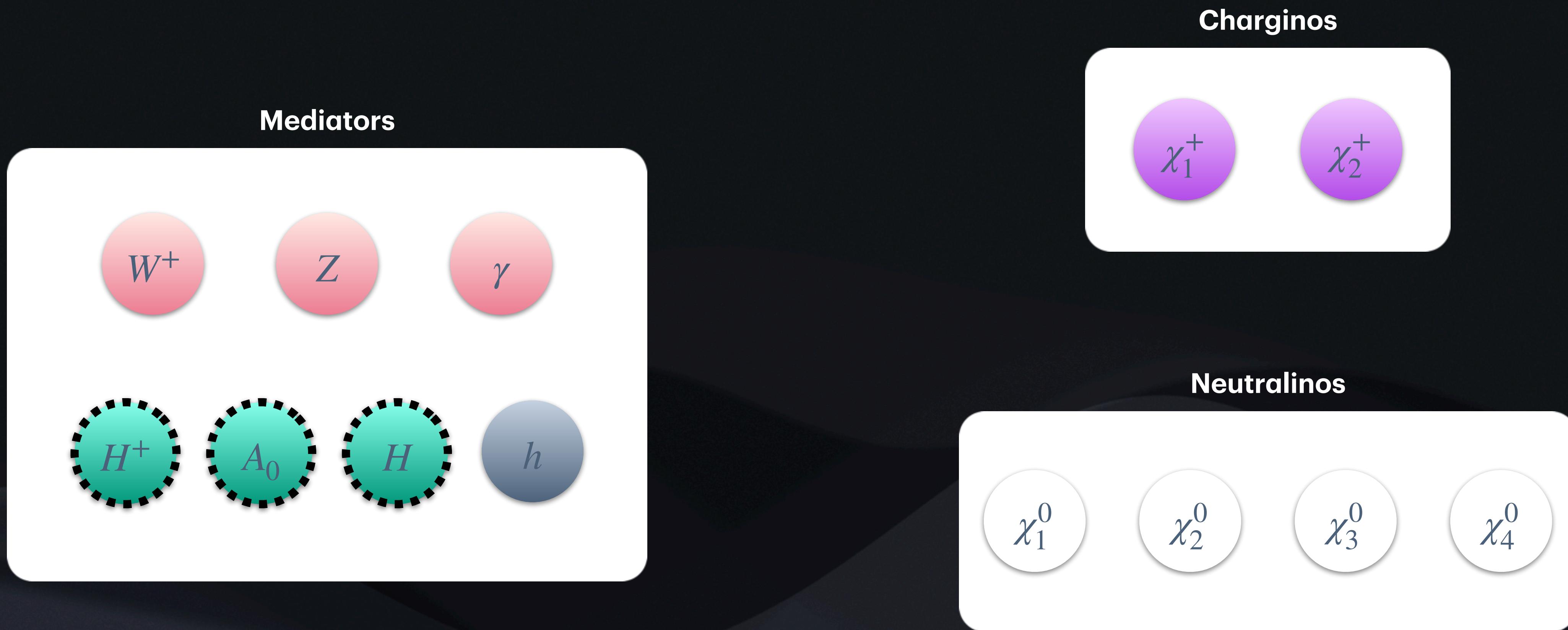
$$\frac{d\sigma\nu}{dE_\gamma} = 2 \sum_{I,J} S_{IJ} \left[\frac{d(\tilde{\sigma}\nu)}{dE_\gamma} \right]_{IJ}$$

105 independent terms



The MSSM in a nutshell

Neutralino/chargino sector

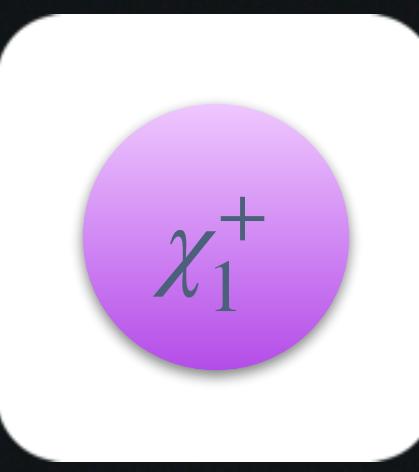


Pure wino

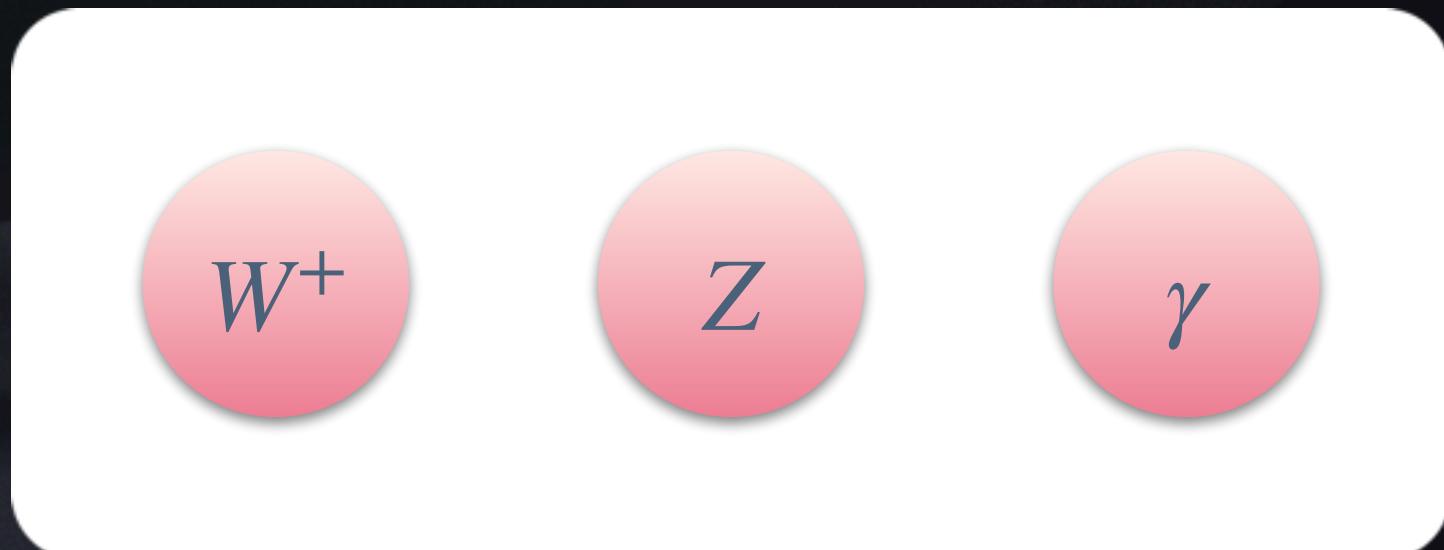
Neutralinos



Charginos



Mediators

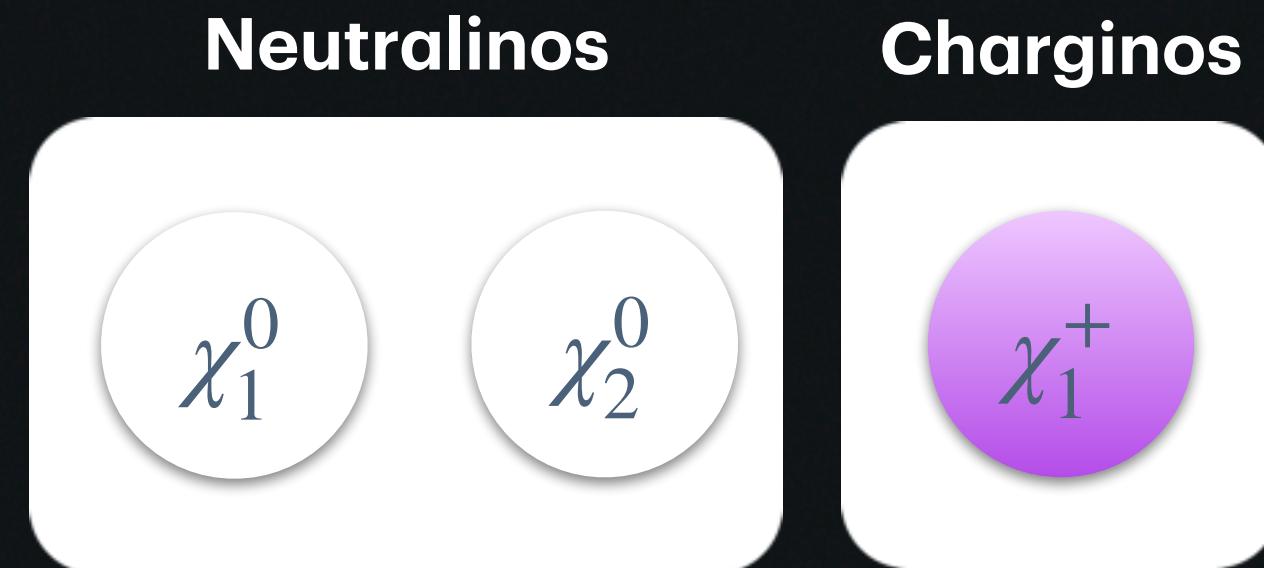


$$V(r) = \begin{pmatrix} 0 & -\sqrt{2}\alpha_2 \frac{e^{-m_W r}}{r} \\ -\sqrt{2}\alpha_2 \frac{e^{-m_W r}}{r} & -\frac{\alpha}{r} - \alpha_2 c_W^2 \frac{e^{-m_Z r}}{r} \end{pmatrix}$$

Further properties:

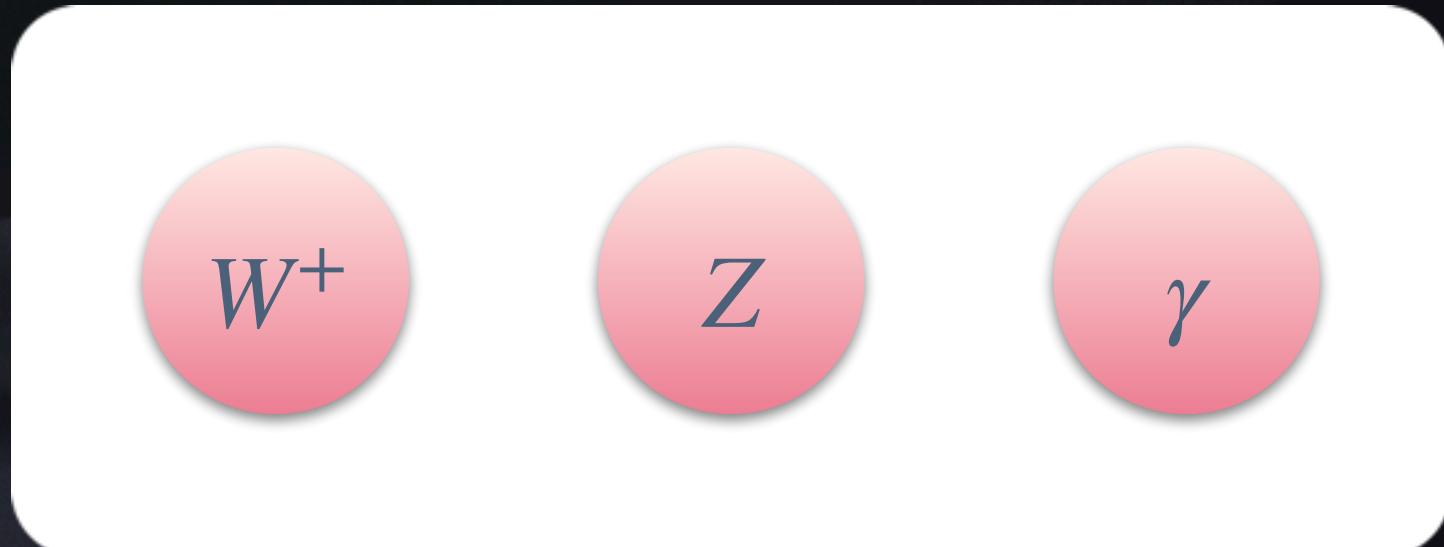
- $m_{\chi_1^0}^{\text{wimp}} \simeq 3 \text{ TeV}$
- $\frac{m_{\chi_1^+} - m_{\chi_1^0}}{m_{\chi_1^0}} \simeq 5.5 \times 10^{-5}$
- No couplings to quarks or gluons

Pure higgsino (see Linda's & Graham's slides)



$$V(r) = \begin{pmatrix} 0 & -\frac{\alpha_2}{4c_W^2} \frac{e^{-m_Z r}}{r} & -\frac{\alpha_2}{2\sqrt{2}} \frac{e^{-m_W r}}{r} \\ -\frac{\alpha_2}{4c_W^2} \frac{e^{-m_Z r}}{r} & 0 & -\frac{\alpha_2}{2\sqrt{2}} \frac{e^{-m_W r}}{r} \\ -\frac{\alpha_2}{2\sqrt{2}} \frac{e^{-m_W r}}{r} & -\frac{\alpha_2}{2\sqrt{2}} \frac{e^{-m_W r}}{r} & \frac{\alpha}{r} - \frac{\alpha_2(s_W^2 - c_W^2)^2}{4c_W^2} \frac{e^{-m_Z r}}{r} \end{pmatrix}$$

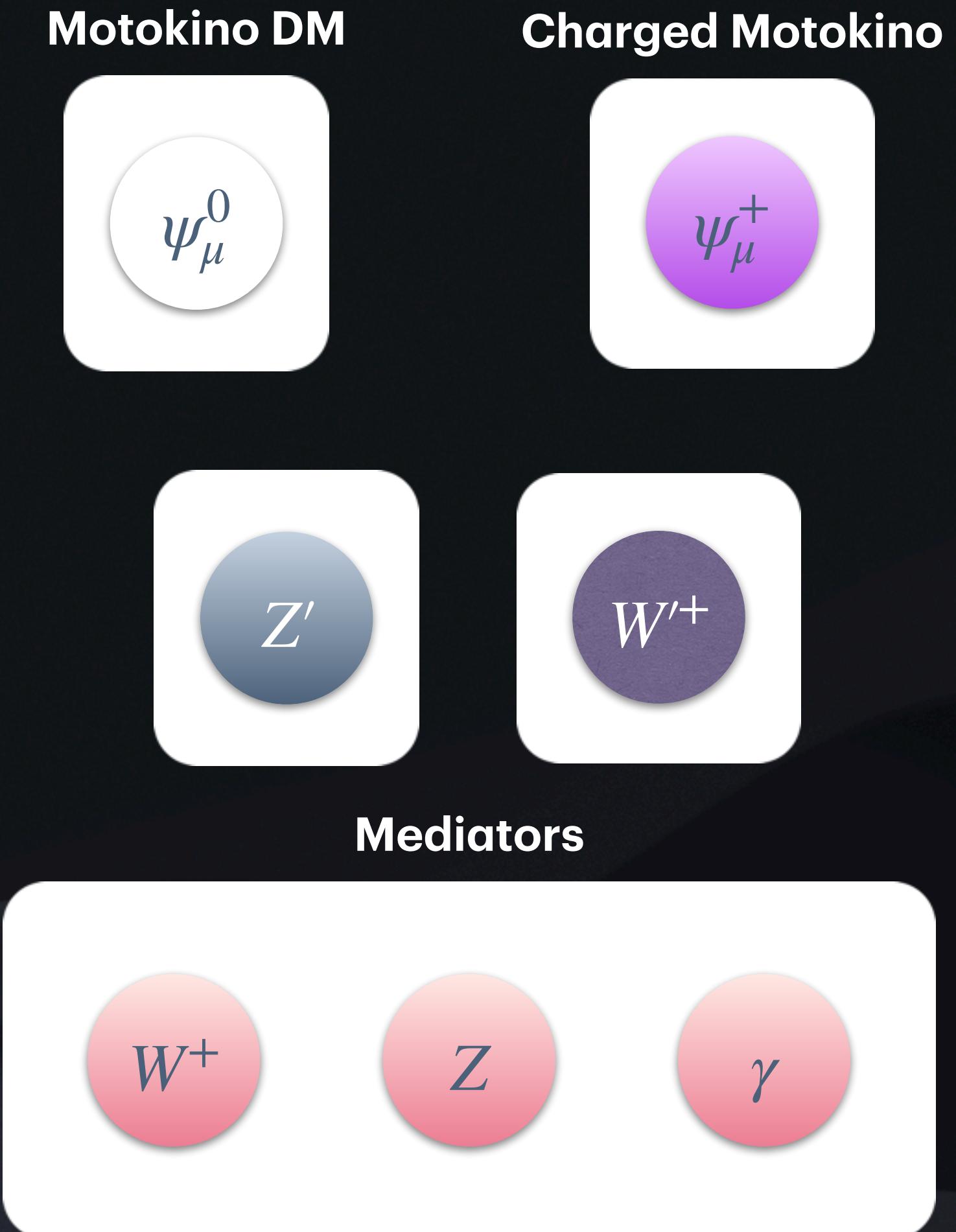
Mediators



Further properties:

- $m_{\chi_1^0}^{\text{wimp}} \simeq 1 \text{ TeV}$
- $\frac{m_{\chi_1^+} - m_{\chi_1^0}}{m_{\chi_1^0}} \simeq 3.5 \times 10^{-4}$
- No couplings to quarks or gluons
- Small admixture with wino/bino required in order to avoid direct-detection constraints

Motokino



$$V^{\text{even}}(r) = \begin{pmatrix} 0 & -\sqrt{2}\alpha_2 \frac{e^{-m_W r}}{r} \\ -\sqrt{2}\alpha_2 \frac{e^{-m_W r}}{r} & -\frac{\alpha}{r} - \alpha_2 c_W^2 \frac{e^{-m_Z r}}{r} \end{pmatrix}$$

Further properties:

- $m_{\psi^0}^{\text{wimp}}$ depend on th. parameters
- No couplings to quarks or gluons
- Identical potential as for wino DM

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Endpoint resummations

TBA

The Basics of Annihilation

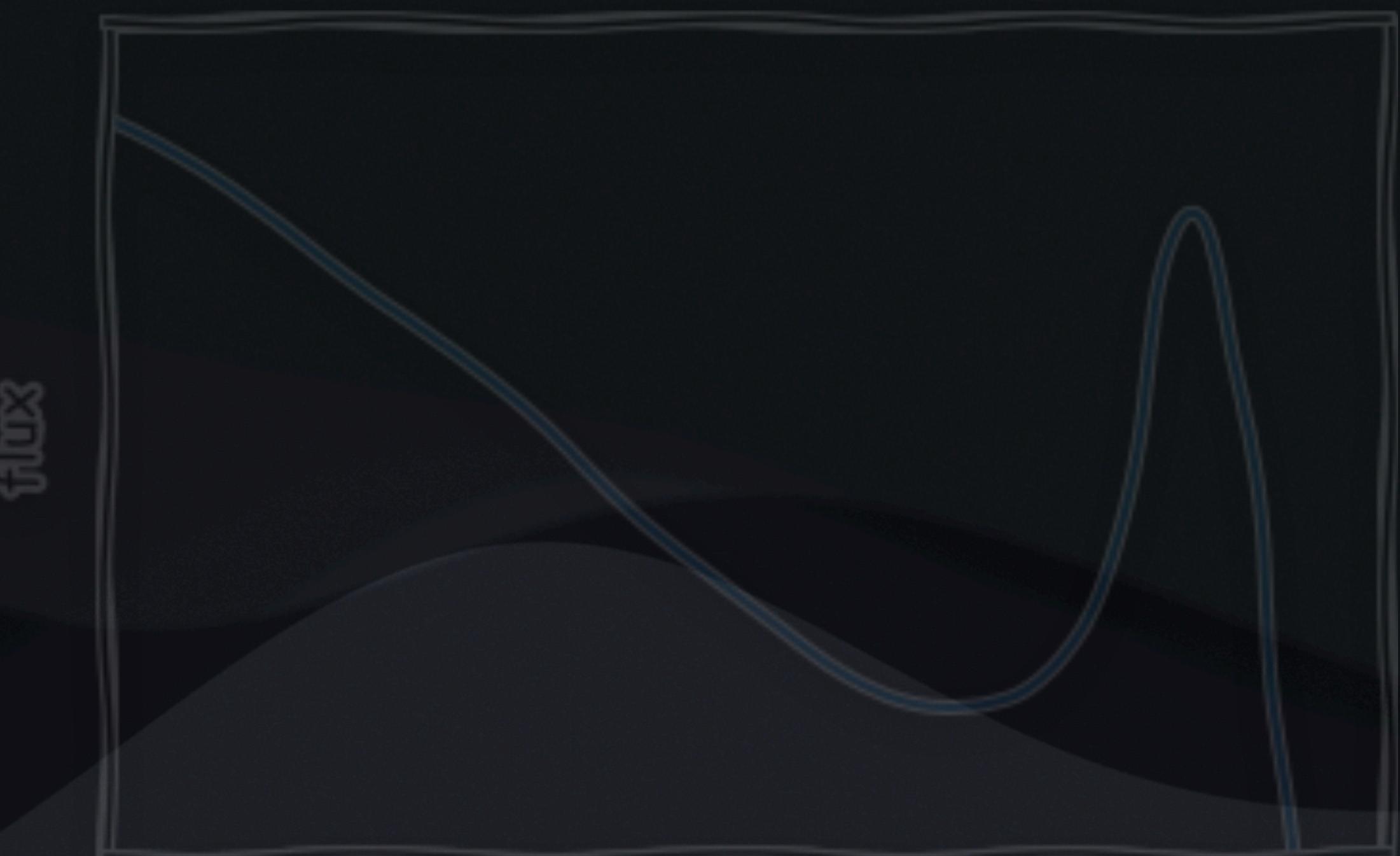


Fixed-order $2 \rightarrow 2$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 3$ (tree)
+ Parton Shower

Fixed-order $2 \rightarrow 2$
+ Parton Shower +
Sommerfeld factor

Fixed-order $2 \rightarrow 3$
+ Parton Shower +
Sommerfeld factor



TBA

The Basics of Annihilation



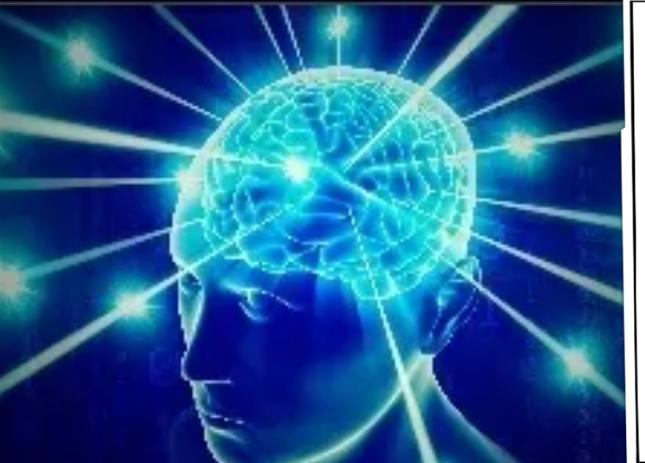
Fixed-order $2 \rightarrow 2$ (tree)
+ Parton Shower



Fixed-order $2 \rightarrow 3$ (tree)
+ Parton Shower



Fixed-order $2 \rightarrow 2$
+ Parton Shower +
Sommerfeld factor



Fixed-order $2 \rightarrow 3$
+ Parton Shower +
Sommerfeld factor



All-order $2 \rightarrow N$ **next-to-leading (prime) Sudakov logs**
+ Parton Shower +
Sommerfeld factor

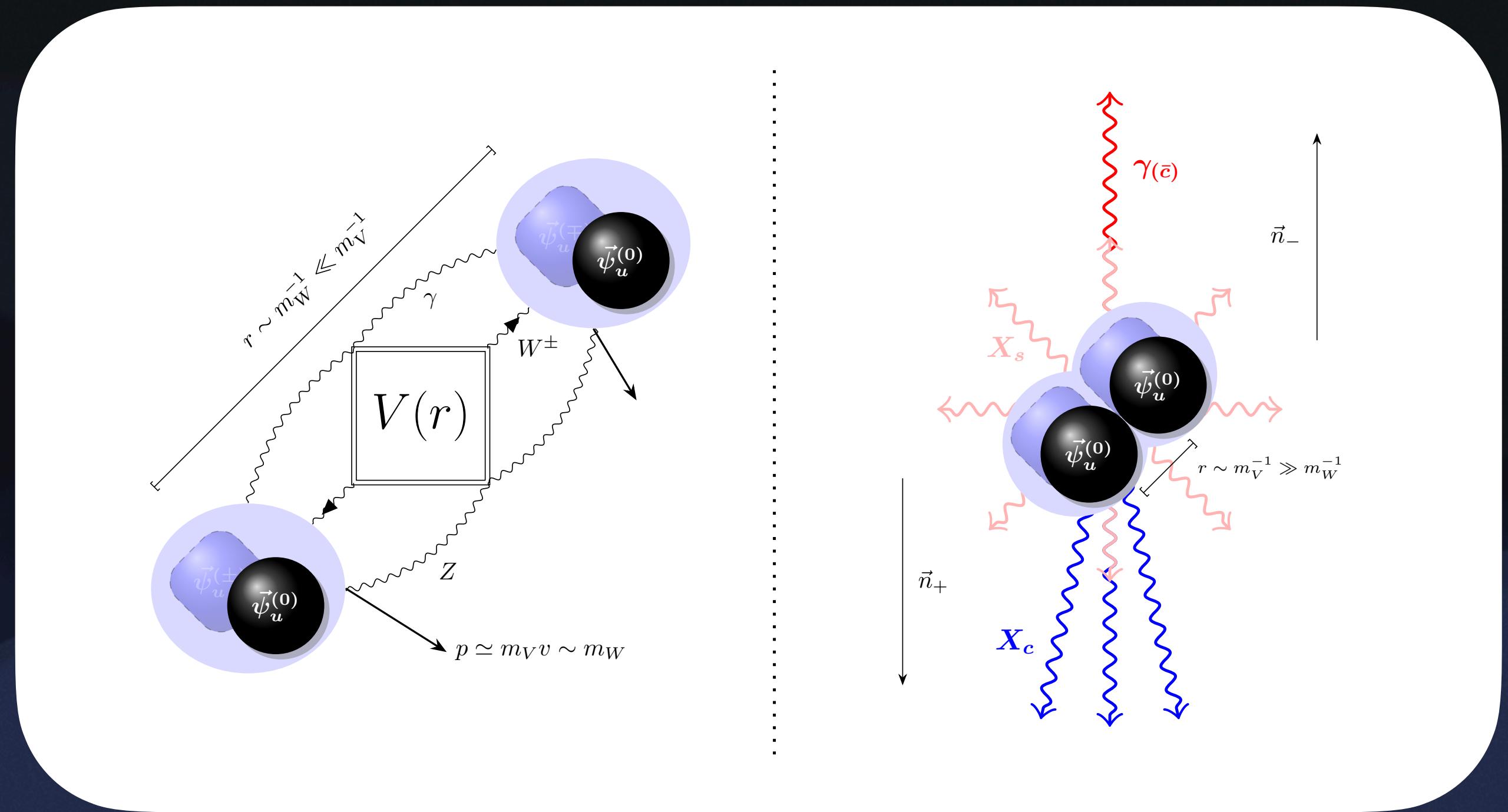


All-order $2 \rightarrow N$ **next-to-leading
(prime) Sudakov logs**
+ Parton Shower +
Sommerfeld factor

Baumgart, Cohen et al – [1712.07656](#)

Beneke, Broggio, Hasner, MV, Urban – [1903.08702](#)

Fujiwara, Vollmann – [2411.XXXX](#)



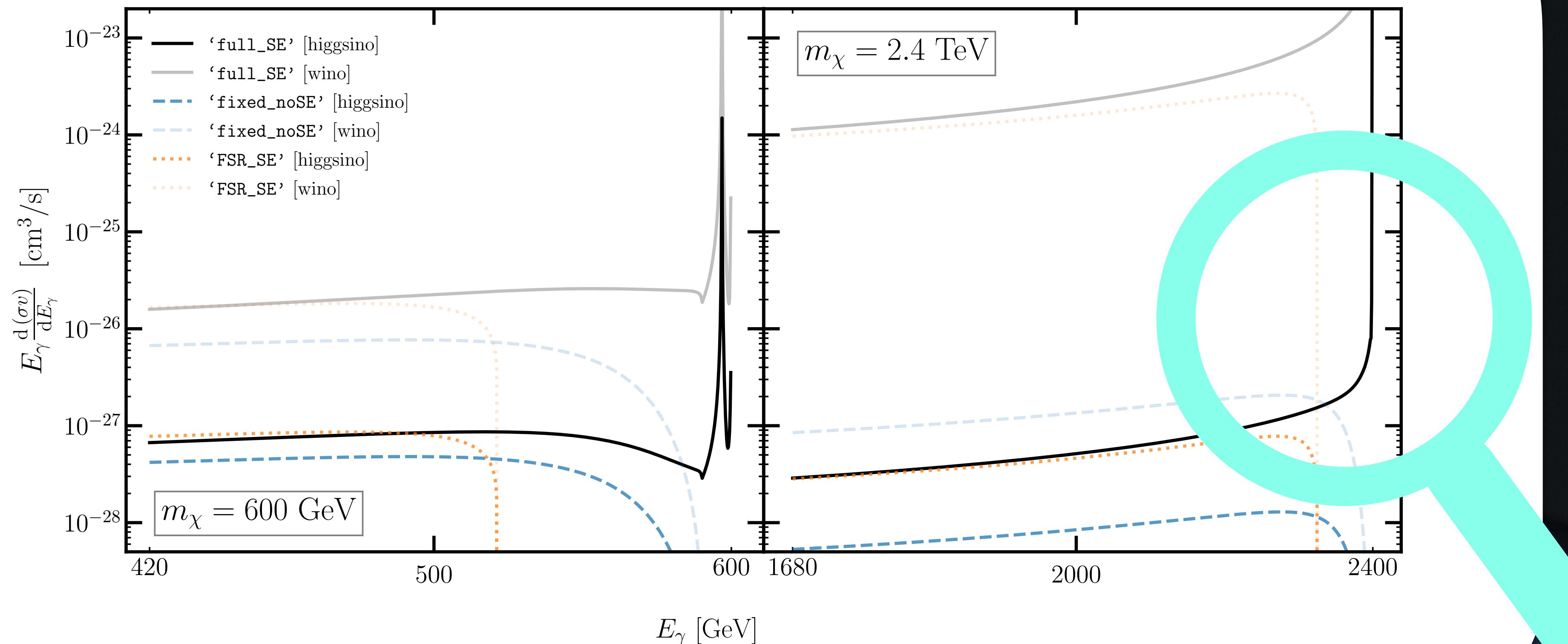
Resummation

Goal: factor out the “gorillas”

$$\frac{d\sigma_v}{dE_\gamma} = f(\alpha_{ew} \times \text{gorilla}) \times (\# \alpha_{ew}^3 + \mathcal{O}(\alpha_{ew}^4))$$

“safe” to use perturbation theory

?



Sudakov double-log resummation

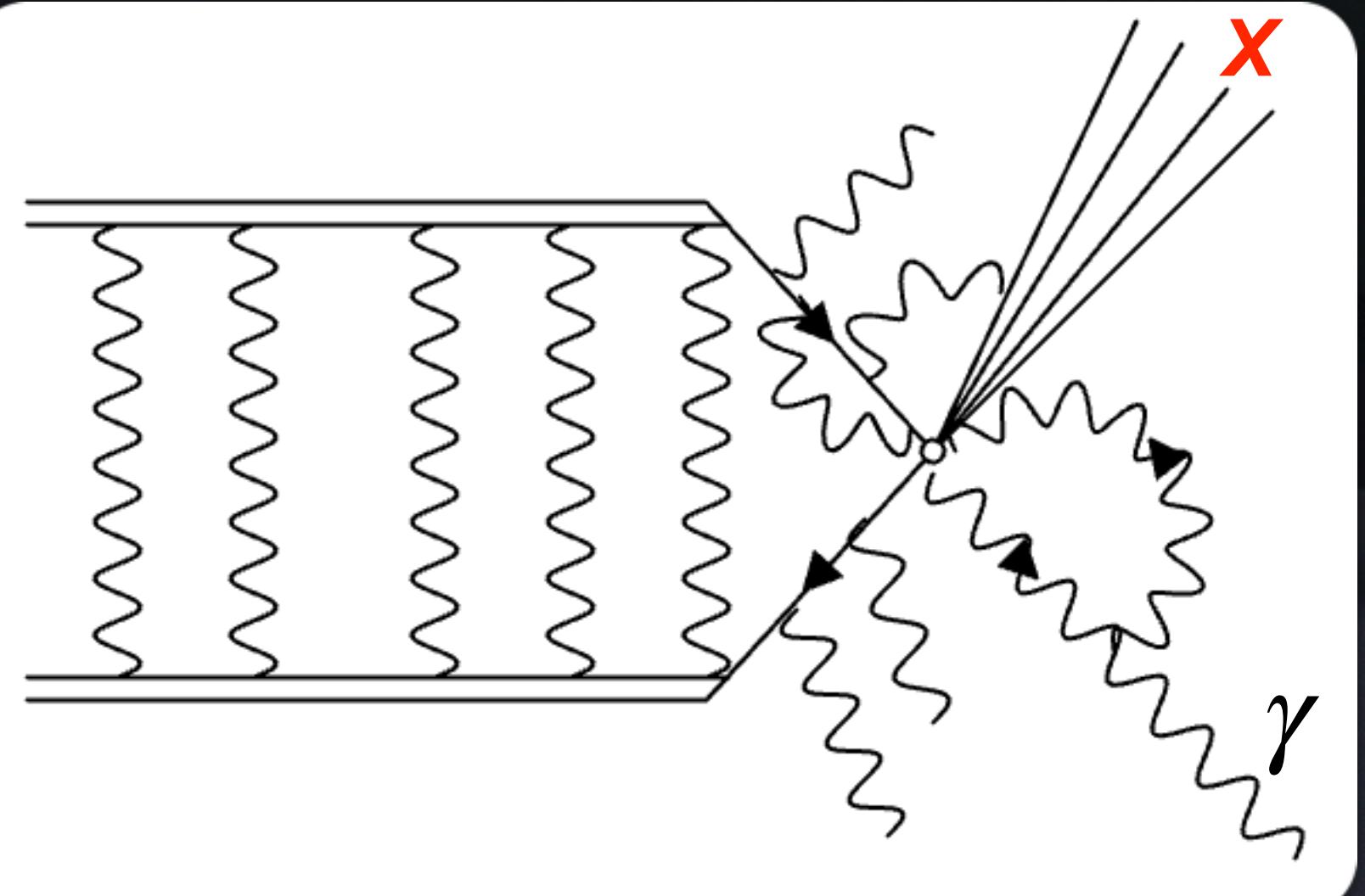
Soft-collinear effective field theory approach

$$\left[\frac{d(\tilde{\sigma}v)^S}{dE_\gamma} \right]_{IJ} = \frac{1}{(\sqrt{2})^{n_{id}}} \frac{1}{9} \frac{2}{\pi m_V} \sum_{\mathcal{J}, \mathcal{J}=0,2} \sum_{i,j} \kappa_{ij}^S H_{\mathcal{J}, \mathcal{J}}^{S;i,j} Z_\gamma^{33} \times \int d\omega J(4m_V(m_V - E - \omega/2)) W_{\mathcal{J}, \mathcal{J};IJ}(\omega)$$

Six Motokino-annihilation operators

$$\begin{aligned} \mathcal{O}_{\mathcal{J}}^{(0)} &= \tilde{\Upsilon}_{u\alpha}^A \eta_u^{\alpha\beta} \tilde{\Upsilon}_{u\beta}^B T_{\mathcal{J}}^{ABCD} \tilde{\mathcal{A}}_{\perp c,\mu}^C(sn_+) \eta_{\perp}^{\mu\nu} \tilde{\mathcal{A}}_{\perp \bar{c},\nu}^D(tn_-) , \\ \mathcal{O}_{\mathcal{J}}^{(2)_1} &= \tilde{\Upsilon}_{u\alpha}^A \tilde{\Upsilon}_{u\beta}^B T_{\mathcal{J}}^{ABCD} (\eta_{\perp}^{\alpha\mu} \eta_{\perp}^{\beta\nu} + \eta_{\perp}^{\alpha\nu} \eta_{\perp}^{\beta\mu}) \tilde{\mathcal{A}}_{\perp c,\mu}^C(sn_+) \tilde{\mathcal{A}}_{\perp \bar{c},\nu}^D(tn_-) - \frac{2}{3} \mathcal{O}_{\mathcal{J}}^{(0)} \\ \mathcal{O}_{\mathcal{J}}^{(2)_2} &= \tilde{\Upsilon}_{u\alpha}^A (n_+ - n_-)^\alpha \tilde{\Upsilon}_{u\beta}^B (n_+ - n_-)^\beta T_{\mathcal{J}}^{ABCD} \tilde{\mathcal{A}}_{\perp c,\mu}^C(sn_+) \eta_{\perp}^{\mu\nu} \tilde{\mathcal{A}}_{\perp \bar{c},\nu}^D(tn_-) \\ &\quad + \frac{4}{3} \mathcal{O}_{\mathcal{J}}^{(0)} ; \\ &\quad + \frac{3}{4} Q_{(0)}^{\mathcal{J}} : \end{aligned}$$

Endpoint $\rightarrow m_X^2 \ll 4m_\chi^2$



Go to the blackboard II

Outline

Motivation

Indirect detection

Sommerfeld factor

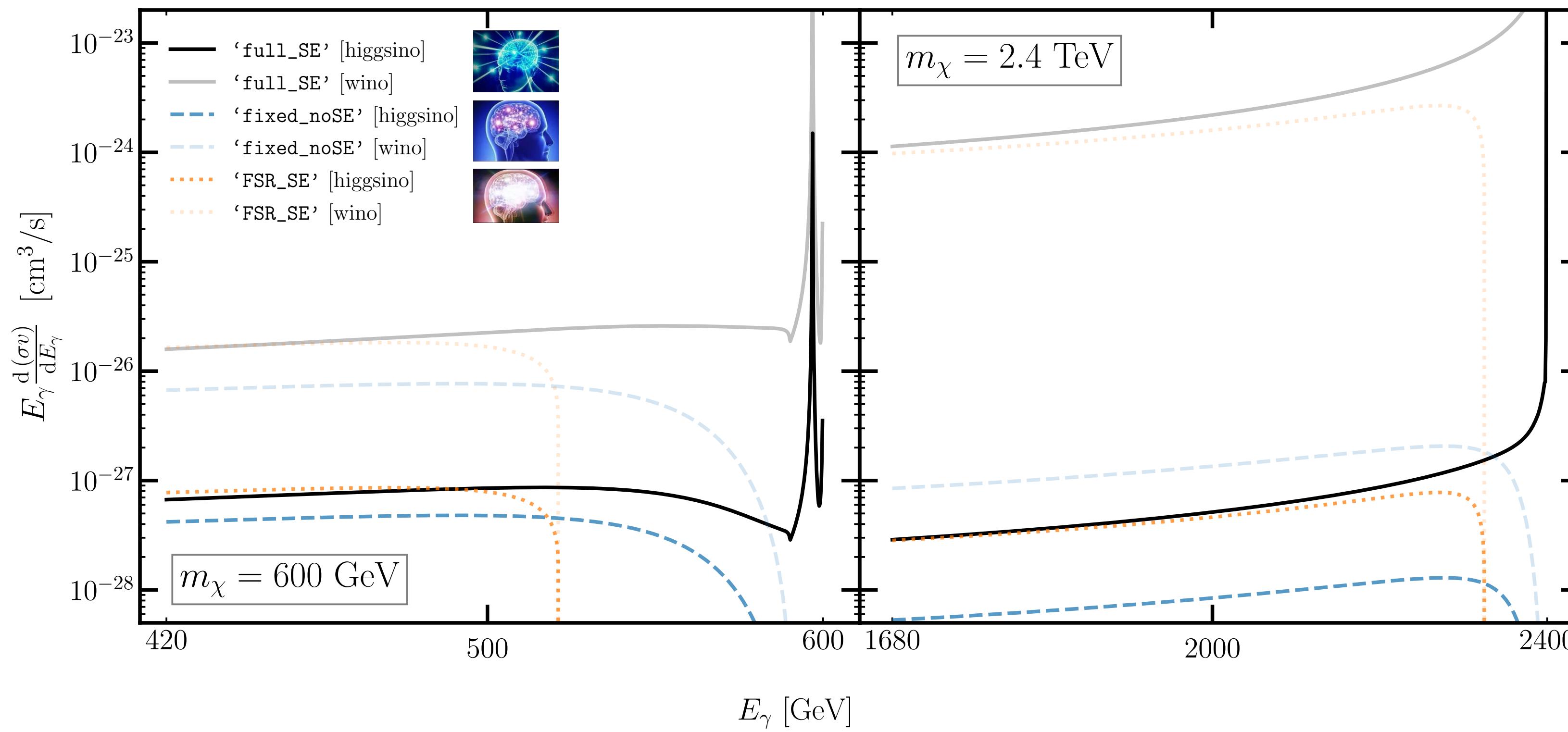
Endpoint resummations

Plots

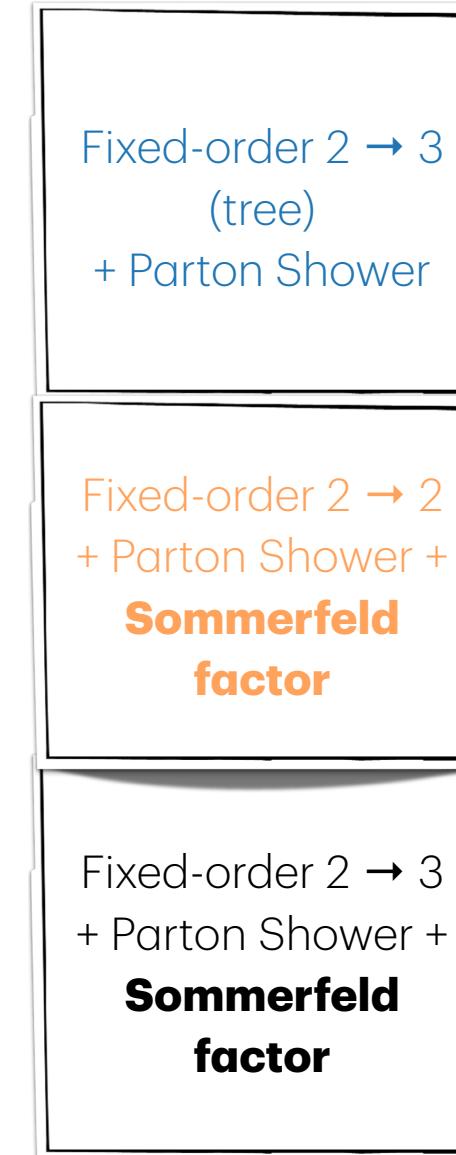
Conclusions

Plots

Without further due...



Meme legend



— ‘full_SE’ [higgsino]

— — ‘full_SE’ [wino]

— — — ‘fixed_noSE’ [higgsino]

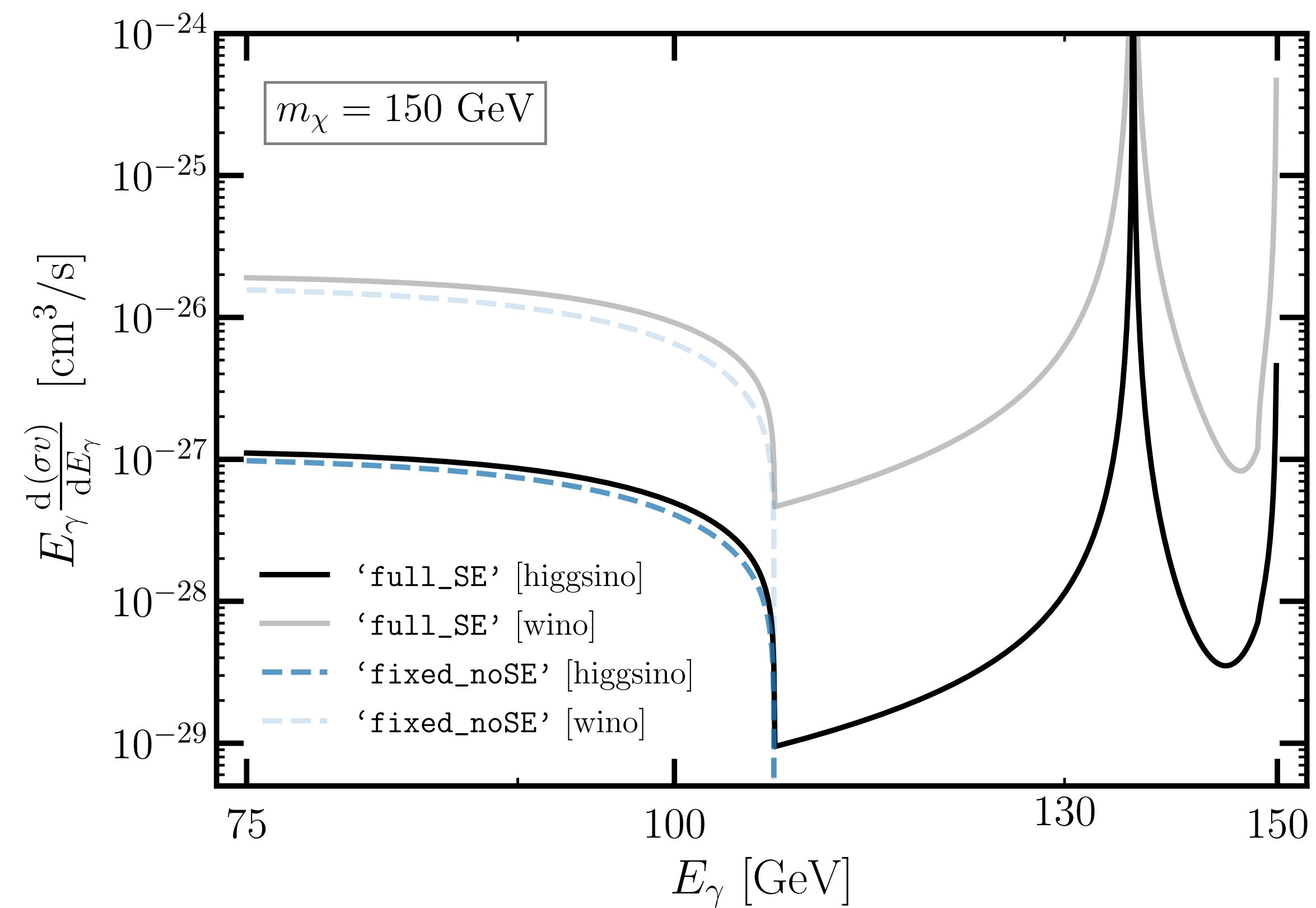
— — — ‘fixed_noSE’ [wino]

····· ‘FSR_SE’ [higgsino]

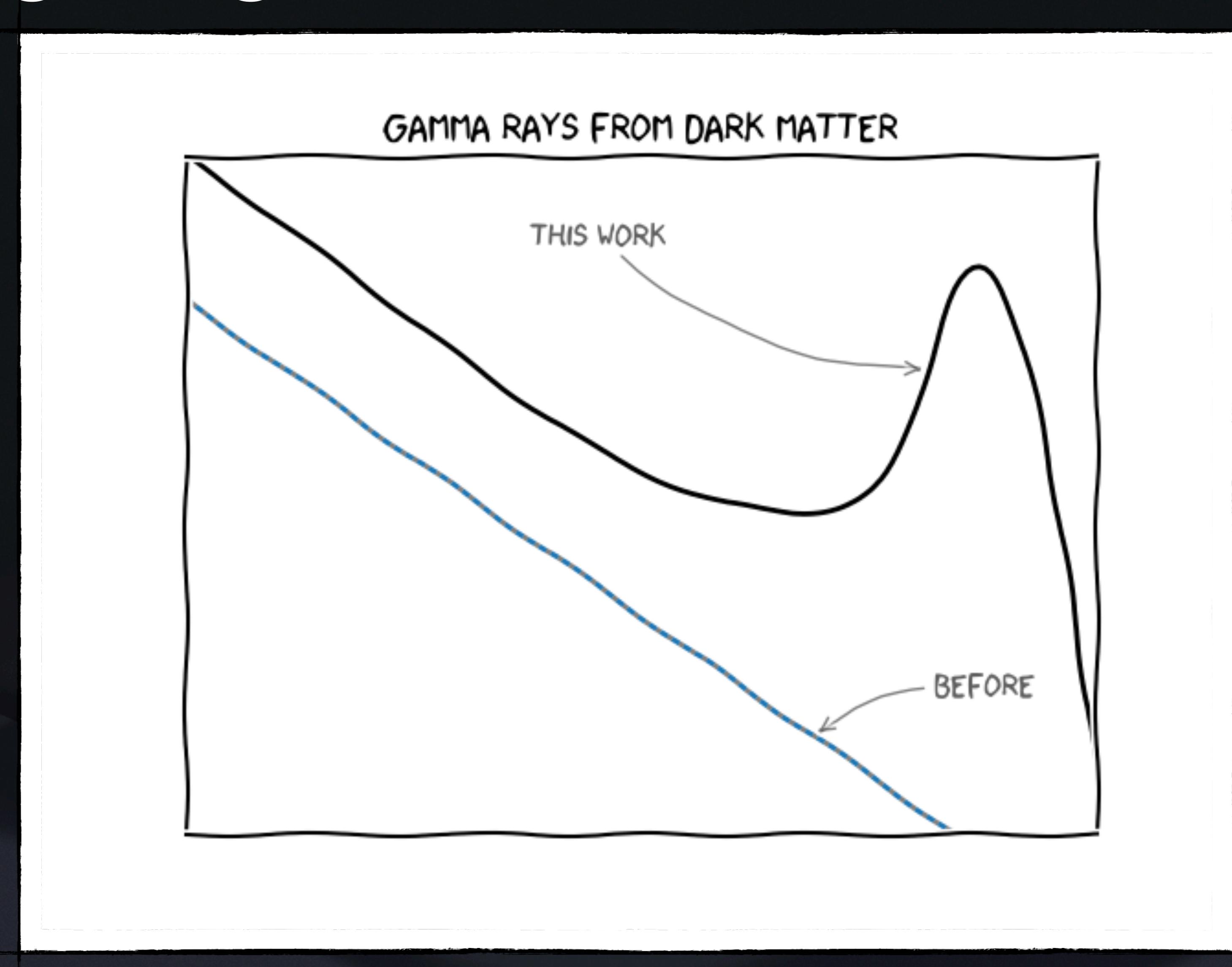
····· ‘FSR_SE’ [wino]

Lighter winos/higgsinos

Kinematic $W^+W^-\gamma$ threshold “lifted” by the Sommerfeld effect

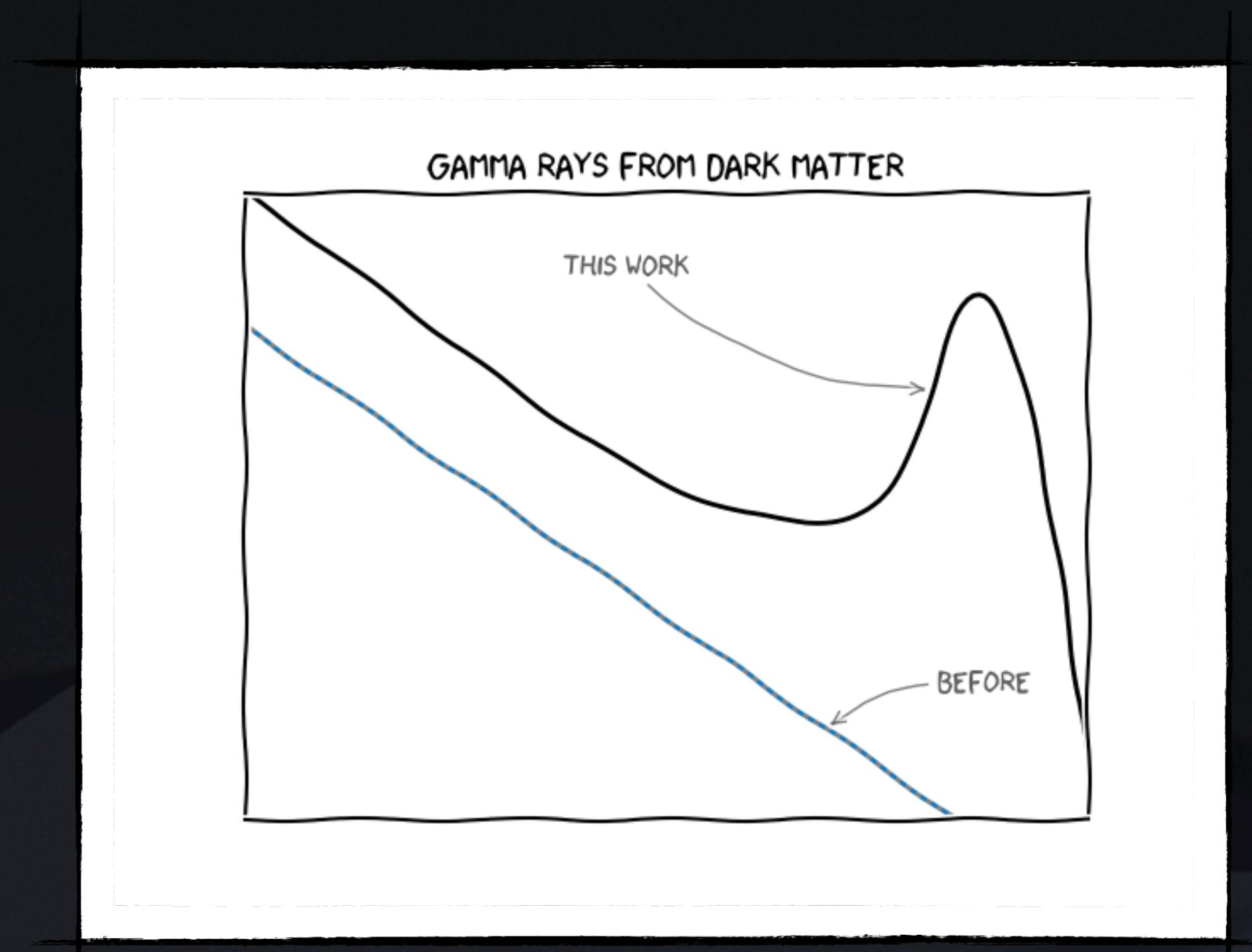
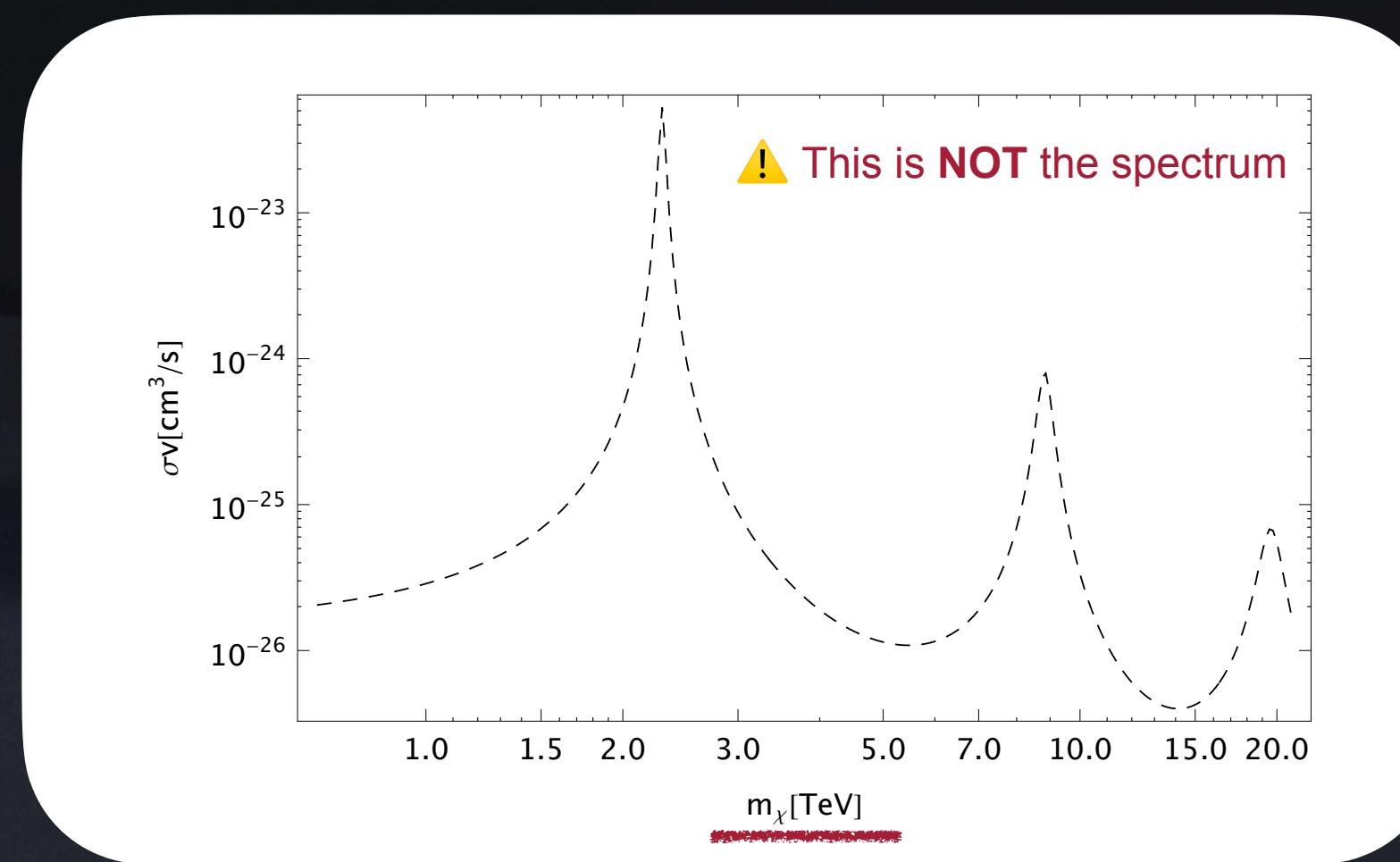
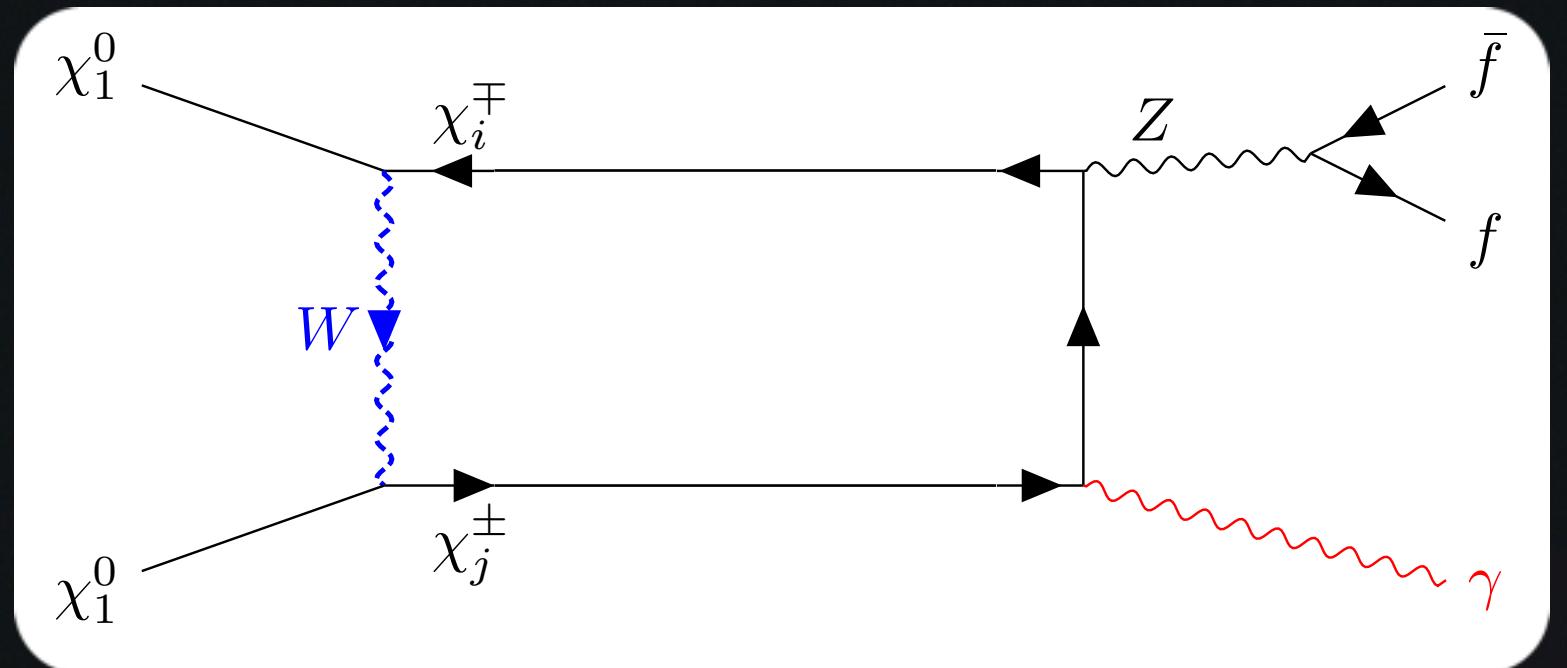


What's going on ?



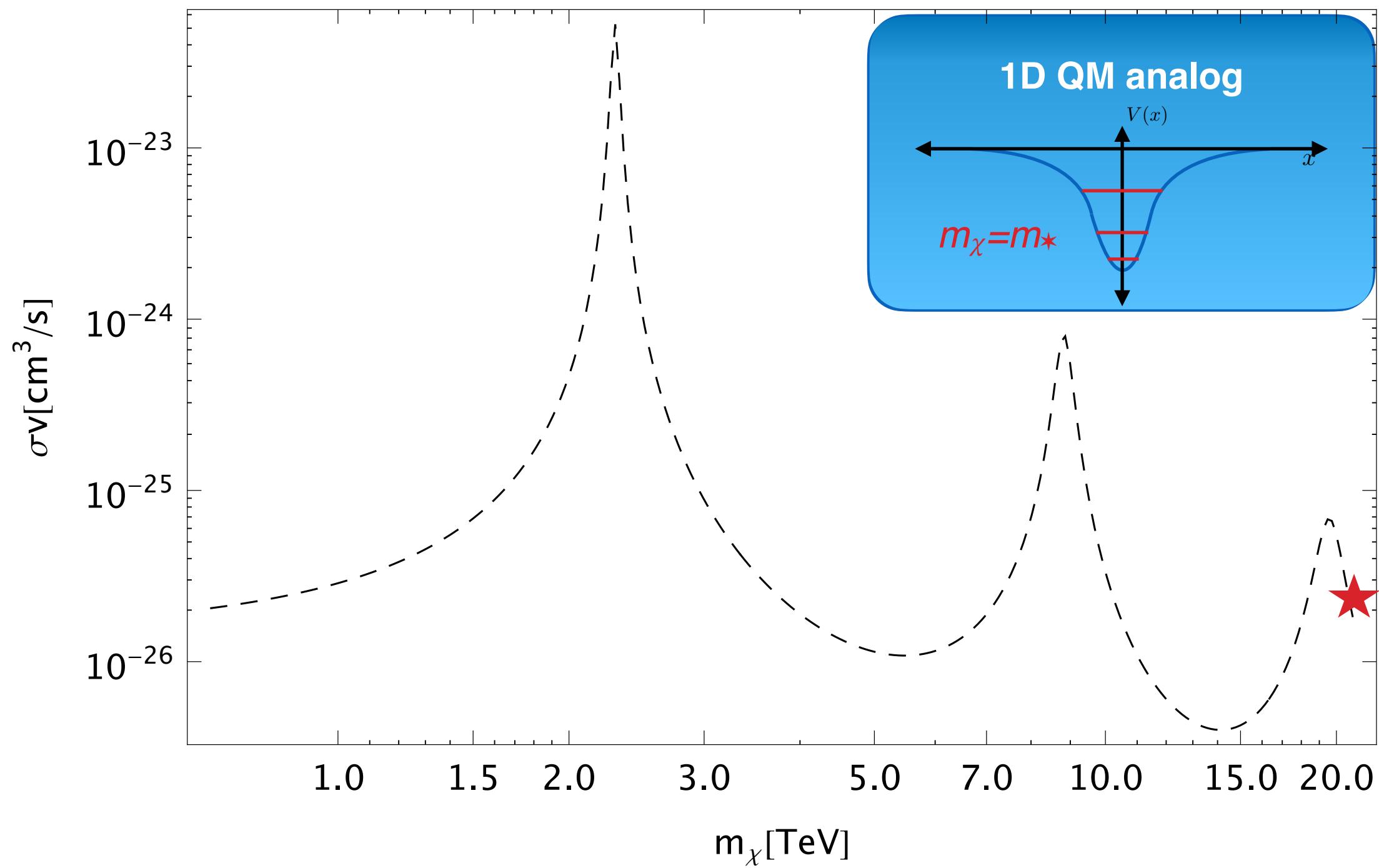
What's going on ?

Charginos are electrically charged / Sommerfeld resonances



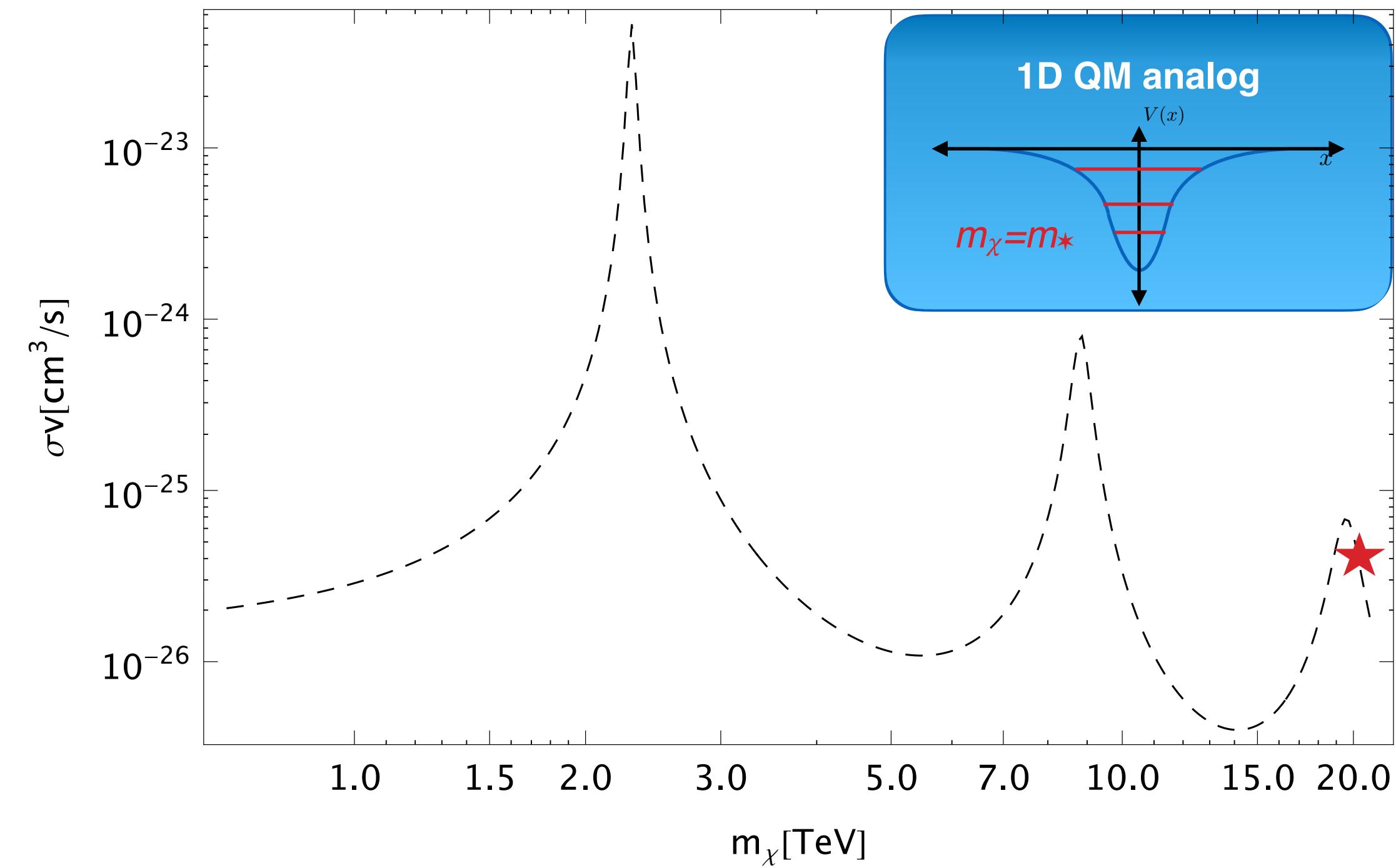
What's going on ?

Sommerfeld resonances



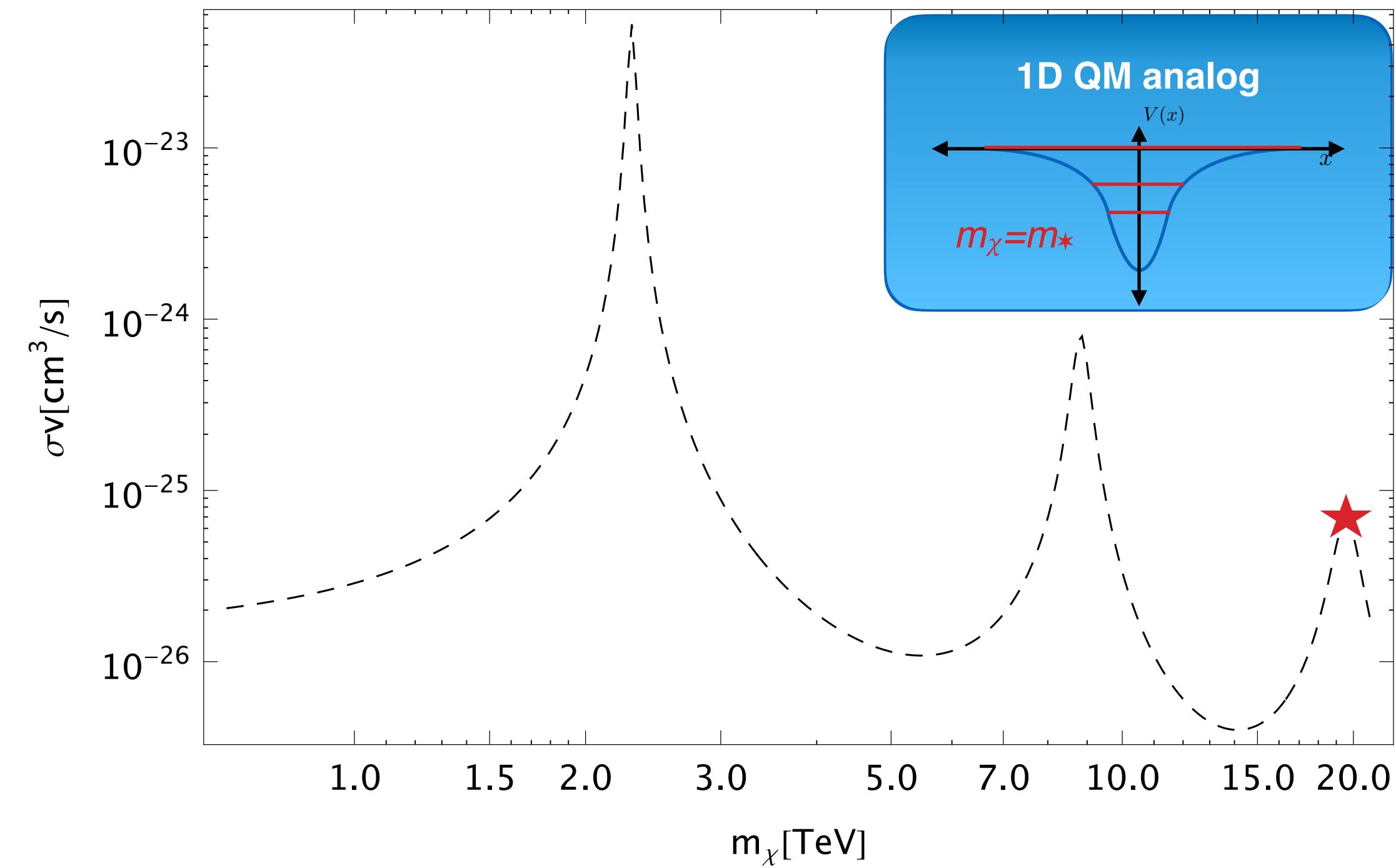
What's going on ?

Sommerfeld resonances



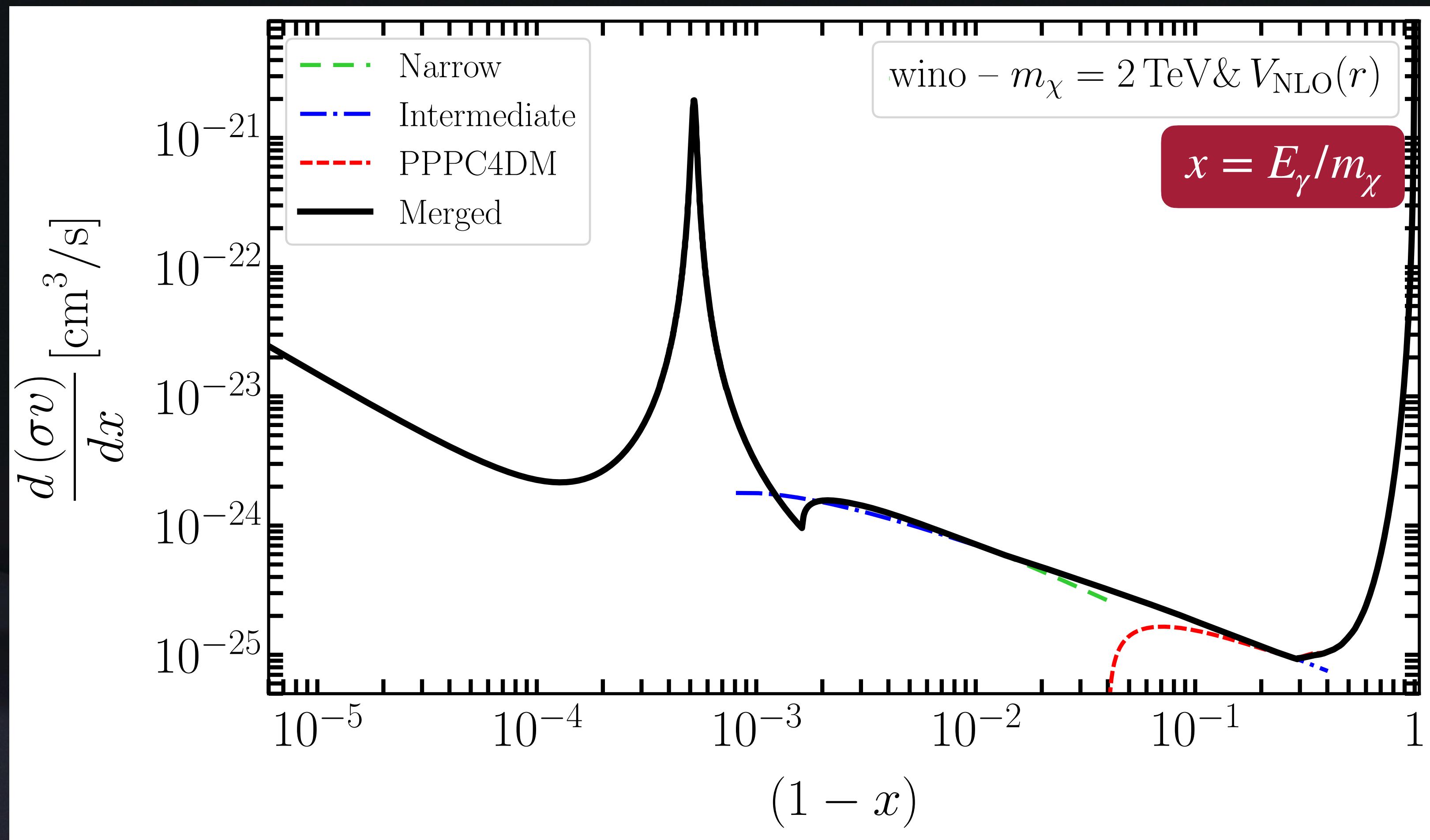
What's going on ?

Sommerfeld resonances

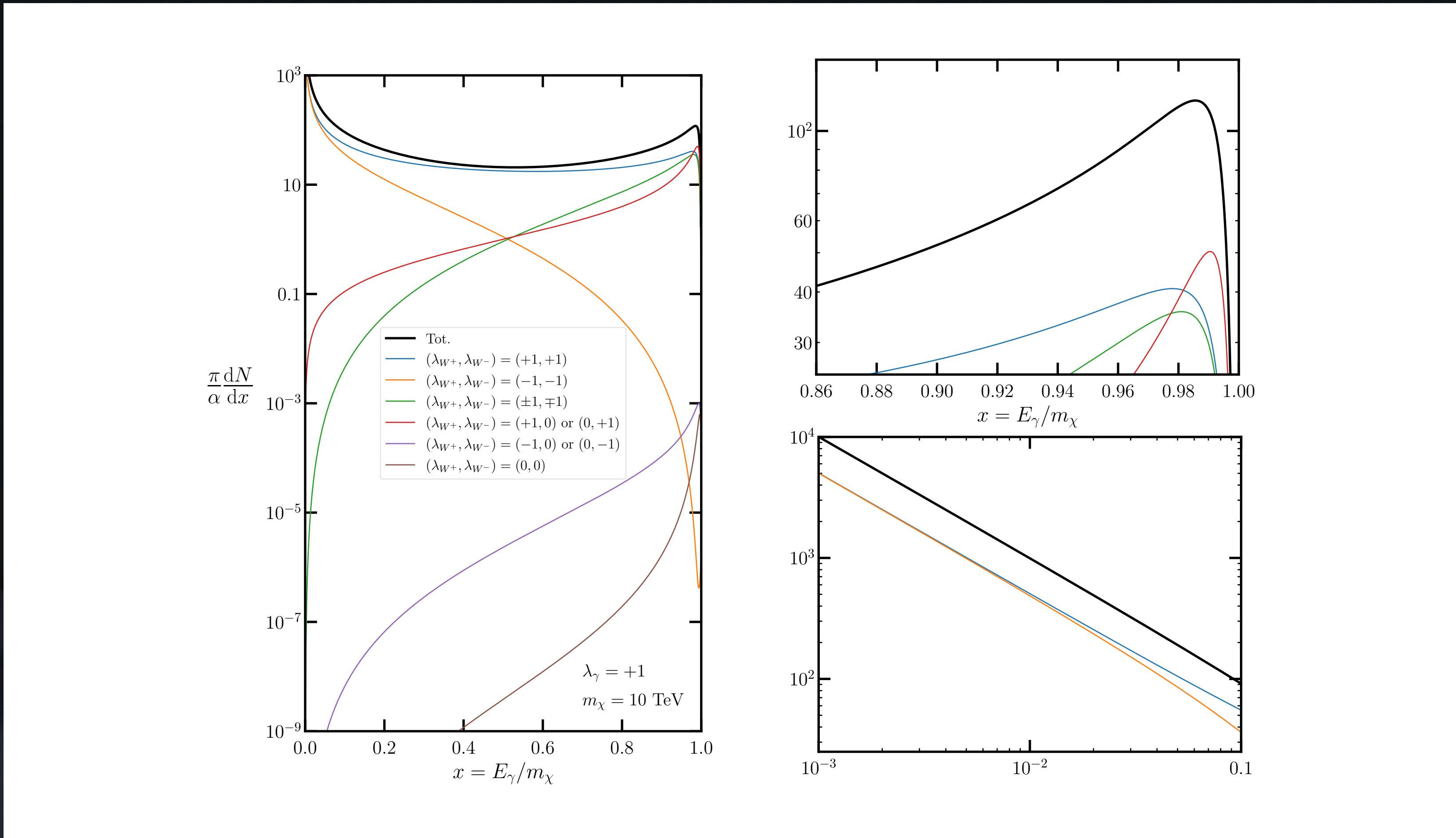


DM γ Spec

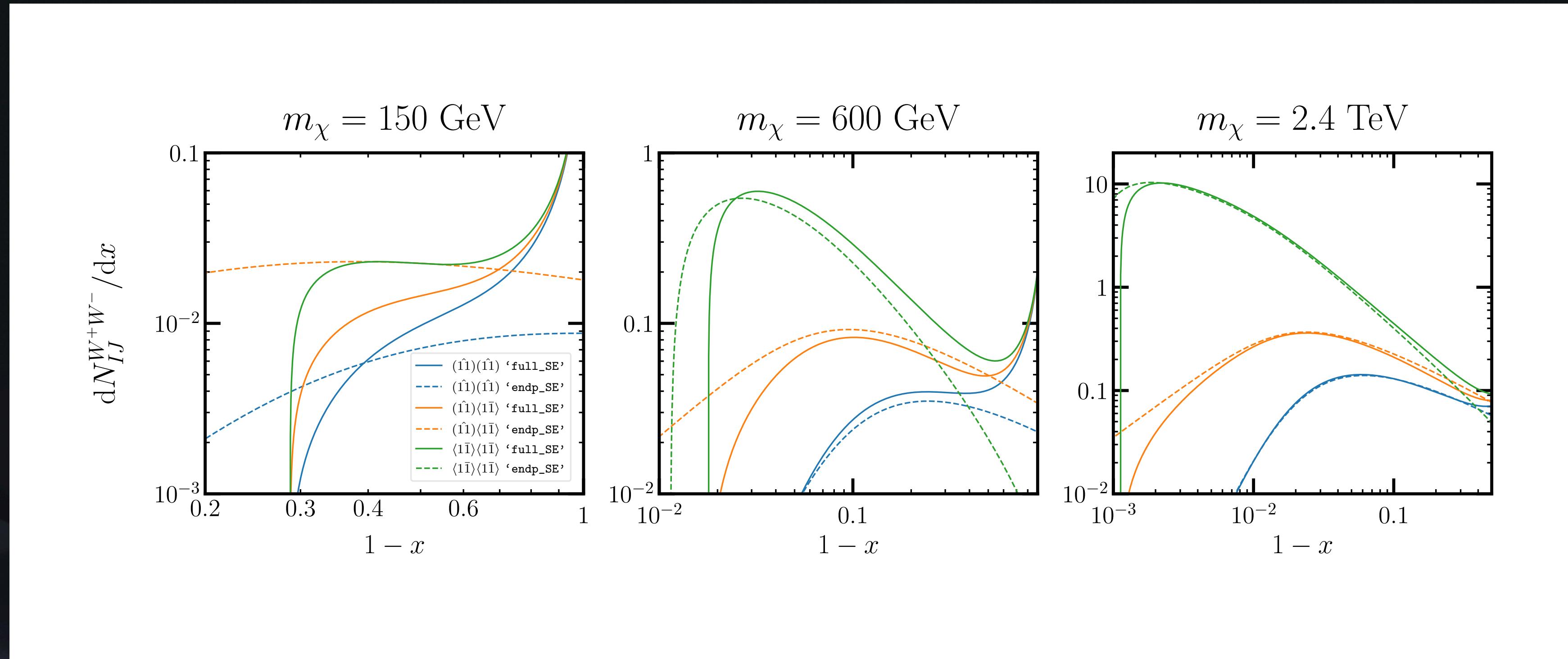
<https://dmyspec.hepforge.org/>



More results to brag about



More results to brag about



Outline

Motivation

Indirect detection

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Endpoint resummations

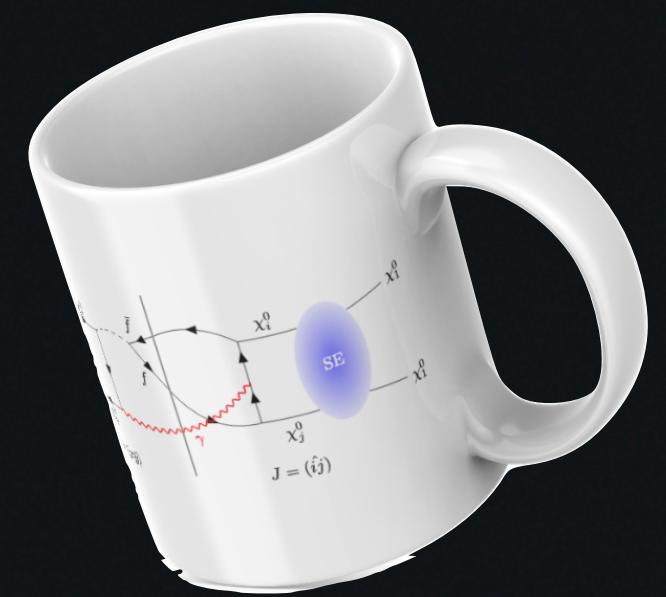
Plots

Conclusions

Conclusions

Conclusions

- TeV-scale WIMP Dark Matter is still “a thing” in 2024
- Cherenkov telescopes are excellent instruments to search for DM
- Beautiful and complex phenomenology for indirect detection
 - Sommerfeld effect, internal bremsstrahlung, resonances, spectral lines, radiative electroweak effects, ...
- Armed with sophisticated EFT tools tackled these obstacles
 - ready for implementation in future observations
- Further explorations in the remaining 2020s **crucial**





MITP 2024

THE DM LANDSCAPE