

Thermal DM with LTR  
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Intro

• Typical BE

$$\frac{dn}{dt} + 3Hn = -\langle \sigma v \rangle (n^2 - n_{eq}^2)$$

$\xrightarrow{\text{Lindke}}$   $\xrightarrow{\text{collision}}$   
 constant  $\xrightarrow{\text{particle physics}}$

• Standard cosmology:

- Entropy SM conserved  $\rightarrow T \sim \frac{1}{a}$
  - Hubble dominated by SM radiation  $\rightarrow H \sim T^2 \sim \frac{1}{a^2}$
  - Very high inflationary reheating
- } guaranteed  $T < T_{BBN}$

dangerous extrapolation



• Reheating

\*  $H \sim \sqrt{P_{inflation}}$

$P_{inflation} \propto a^{-3(1+w)}$

w: effective EoS during reheating

$H(t) \propto a^{-\frac{3}{2}(1+w)}$

- w
- 0: MD
  - $\frac{1}{3}$ : RD
  - +1: Kinetic
  - $\frac{n-2}{n+2}$ :  $V(\phi) \propto \phi^n$

$w \leq \frac{1}{3}$  :  $\phi$  has to decay  
 annihilate  
 $\Rightarrow$  Injection of entropy

\* SM temperature during reheating

$$T(x) = T_{rh} \left( \frac{a_{rh}}{a} \right)^\alpha$$

$$\left\{ \begin{array}{l} a_{rh} = a(T_{rh}) \\ \rho_r(T_{rh}) = \rho_f(T_{rh}) \end{array} \right.$$

• Examples:

\* MD:  $W=0$      $\alpha = 3/4$     ~~constant~~

- inflation in  $V \sim \phi^2$  with constant  $\Gamma$
- EMD (non-relativistic)
- PBH.

Plot

→ Inflation decay  $\Gamma = \Gamma(a)$  or  $\Gamma = \Gamma(\phi)$

$W=0$      $\alpha \neq 3/4$

→  $\phi \rightarrow$  relativistic particle  $\rightarrow$  SM

$W=0$      $\alpha = 0$

\* RD:  $W = 1/3$     mass becomes field dependent

$$\left. \begin{array}{l} \alpha = 1/4 \\ \alpha = 3/4 \end{array} \right\} \begin{array}{l} \text{decay scalars} \\ \text{decay fermions} \end{array}$$

\*  $V(\phi) \sim \phi^m$

$W = \frac{n-2}{n+2}$

$\alpha = \frac{3}{2} \frac{1}{2+n}$     decay scalars

$\alpha = \frac{3}{2} \frac{n-1}{n+2}$     decay fermions

$\alpha = \frac{3}{2n+4}$

annihilation heavy <sup>into</sup> brms

$\alpha = \frac{3(7-2n)}{2n+4}$

annihilation light brms  
mediator

$$\alpha = 1$$

annihilation heavy fermions  
light

$$\alpha = \frac{3(S-n)}{2n+4}$$

\* Kinetic a faster-than radiation  $w > \frac{1}{3}$   $\alpha = -1$

→ If  $\alpha > 0$  :  $T_{\text{max}} \gg T_{\text{rh}}$

$\alpha = 0$  :  $T$  constant during reheating

**PLOT**

\* Hubble:  $H^2 = \frac{\rho_r + \rho_\phi}{3M_p^2}$

$$H(a) = H_{\text{rh}} \times \begin{cases} \left(\frac{a_{\text{rh}}}{a}\right)^{\frac{3(1+w)}{2}} & a \leq a_{\text{rh}} \\ \left(\frac{a_{\text{rh}}}{a}\right)^2 & a_{\text{rh}} \leq a \end{cases}$$

$$H(T) = H_{\text{rh}} \times \begin{cases} \left(\frac{T}{T_{\text{rh}}}\right)^{\frac{3(1+w)}{2}} & T \geq T_{\text{rh}} \\ \left(\frac{T}{T_{\text{rh}}}\right)^2 & T_{\text{rh}} \geq T \end{cases}$$

3 free parameters:  $T_{\text{rh}}, w, \alpha$  ← Effective approach to reheating

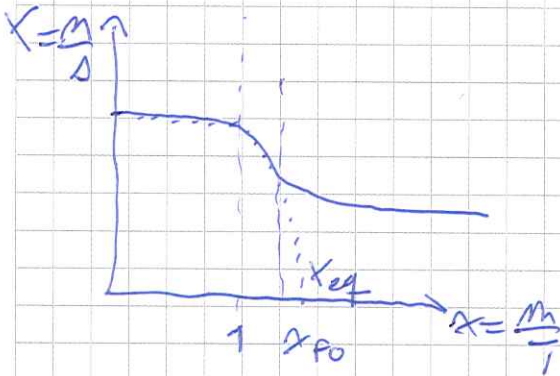
we only assume power-law scalings for  $H$  &  $T$  after decoupling

**PLOT**

production

Thermal DM AFTER reheating

\* WIMPs & SIMPs ( $T_h \gg T_{FO} \gg T_k$ )

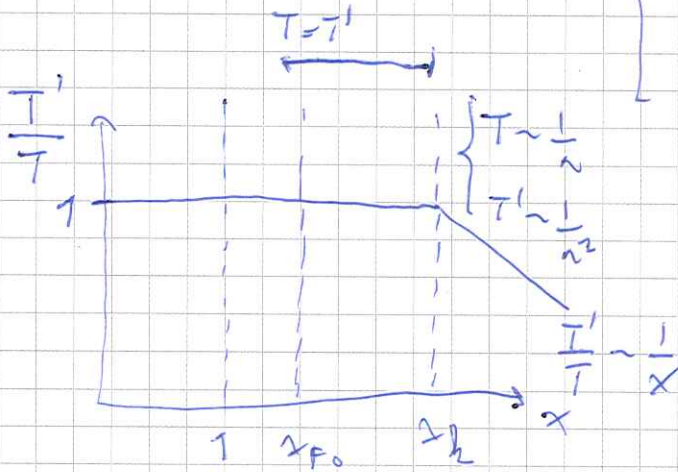


$x_{FO} \sim 20$  chemical FO  
 (number-changing interactions)

WIMP: annihilation, semi-annihilation, coannihilation

SIMP:  $3 \rightarrow 2$ ,  $4 \rightarrow 2$

$0n \rightarrow 0n$ ,  $0n \rightarrow 2n$



$x_h \sim 10^2 - 10^4$  WIMP  
 $x_h \geq x_{FO}$  SIMP

kinetic FO  
 (elastic interactions)

$0n \rightarrow 0n$ ,  $0n \rightarrow 2n$ ,  $2n \rightarrow 0n$

→ Instantaneous decoupling

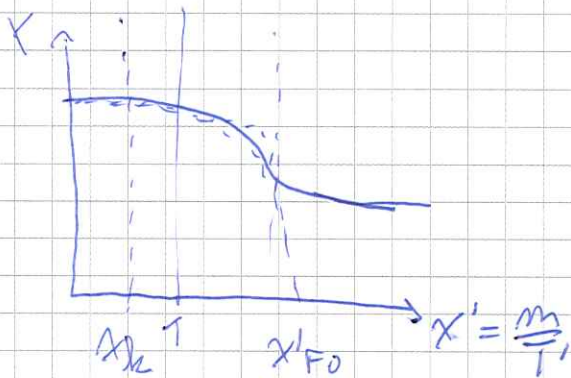
$$x_0 \approx x_{FO} = \frac{M_{eq}(T_{FO})}{s(T_{FO})}$$

$$\left\{ \begin{aligned} M_{eq} &\sim (mT)^{3/2} e^{-\frac{m}{T}} \\ s &\sim T^3 \end{aligned} \right.$$

$$x_0 \approx x_{FO}^{3/2} e^{-x_{FO}}$$

\* ELDERs & cannibals (T<sub>h</sub>) T<sub>h</sub> ≫ T<sub>F0</sub>

$$T_{F0} \neq T'_{F0}$$



$$Y_0 \approx X_{F0} = \frac{M_{eq}(T'_{F0})}{\Delta(T_{F0})}$$

S and S' independently conserved:  $\Delta(T_{F0}) = \Delta(T_h) \left(\frac{a_k}{\Delta_{F0}}\right)^3$

$$= \Delta(T_h) \frac{S'(T_h)}{S'(T'_{F0})}$$

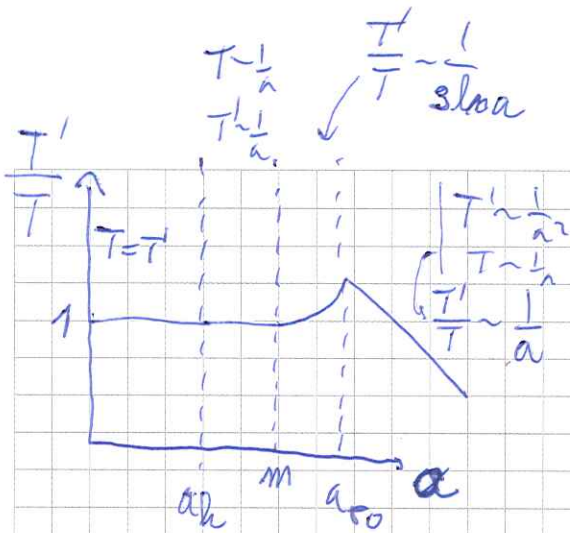
$$Y_0 \approx \frac{M_{eq}(T'_{F0})}{\Delta(T_h)} \frac{S'(T_h)}{S'(T'_{F0})}$$

relativistic or non-Rel.

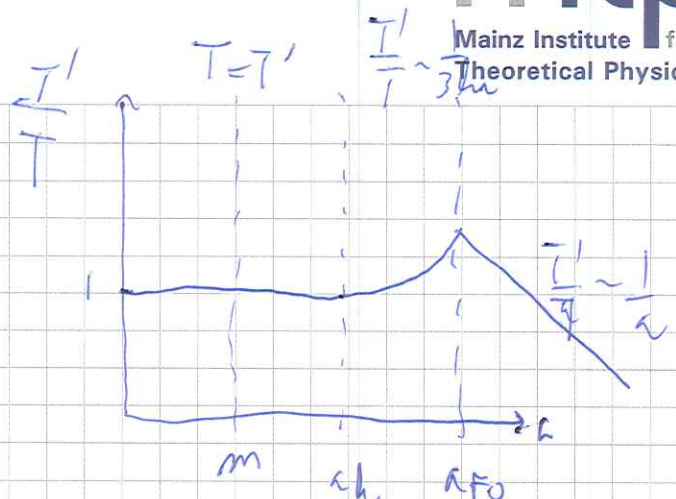
Entropy density in the dark sector:  $S'(T') \sim m^3 K_3\left(\frac{m}{T'}\right)$

$$Y_0 \approx \frac{\Delta_h^3}{x'_{F0}} \frac{K_2(x'_{F0}) K_3(\Delta_h)}{K_3(x'_{F0})}$$

2 limits:  $Y_0 \approx \begin{cases} \frac{1}{x'_{F0}} & x_h \ll 1 \quad \text{cannibal} \\ \frac{x_h^{5/2} e^{-x_h}}{\Delta'_{F0}} & x_h \gg 1 \quad \text{ELDER} \end{cases}$



~~Controlled~~ Controlled  
 maximal "increase"



~~Controlled~~ ELDER  
 controlled "increase"

PLOTS ① T and ② X ③  $x_h$  vs  $x_{F0}$

3 Free parameters:  $m, x_{F0}, x_h$

Thermal DM production DURING reheating

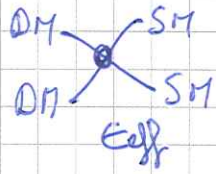
$\phi$  injects entropy to the SM  $\rightarrow$  DM gets diluted

$$\frac{S(T)}{S(T_h)} = \frac{s(T)}{s(T_h)} \left(\frac{a}{a_h}\right)^3 \sim \left(\frac{T_h}{T}\right)^{\frac{3(1-2)}{\alpha}}$$

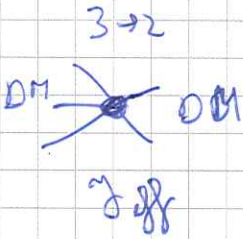
PLOTS

$\Rightarrow$  not really depend independent of  $\underline{\underline{w}}$  !

• Tony Model - contact interactions



$$\langle \sigma v \rangle_d \sim \langle \sigma v \rangle_{2 \rightarrow 2} \sim \left[ \frac{E_{\text{eff}}}{m} \frac{K_1(x)}{K_2(x)} \right]^2 \sim \begin{cases} \frac{E_{\text{eff}}^2}{4T^2} & x \ll 1 \\ \frac{E_{\text{eff}}^2}{m^2} & x \gg 1 \end{cases}$$



$$\langle \sigma v^2 \rangle_{3 \rightarrow 2} \sim \frac{\gamma_{\text{eff}}^3}{m^5}$$

$$\Gamma_d = M_{\text{eq}}^{\text{SM}} \langle \sigma v \rangle_d$$

$$\Gamma_{2 \rightarrow 2} = M_{\text{eq}} \langle \sigma v \rangle_{2 \rightarrow 2}$$

$$\Gamma_{3 \rightarrow 2} = n_{\text{eq}}^2 \langle \sigma v^2 \rangle_{3 \rightarrow 2}$$

