

Inflationary Schwinger Dark Matter Production

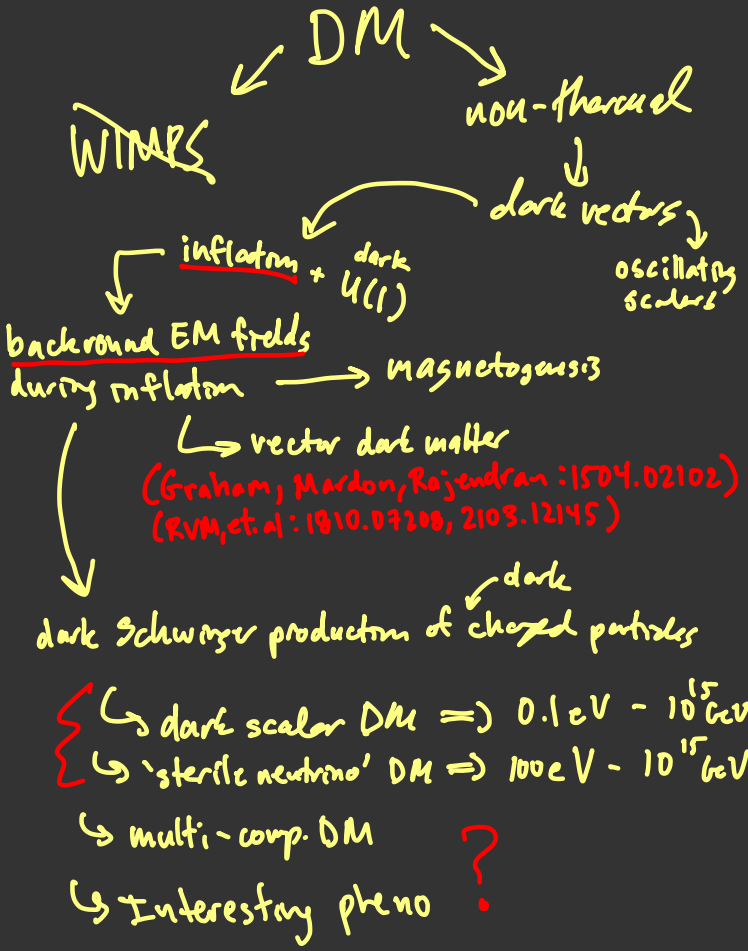
Roberto Vega-Morales
University of Granada

DMLAND Workshop

Mainz Germany, September 4, 2024

In collaboration with: Mar Bastero-Gil
Paulo Ferraz, Jose Santiago, Lorenzo Ubaldi

Motivation / Punchline



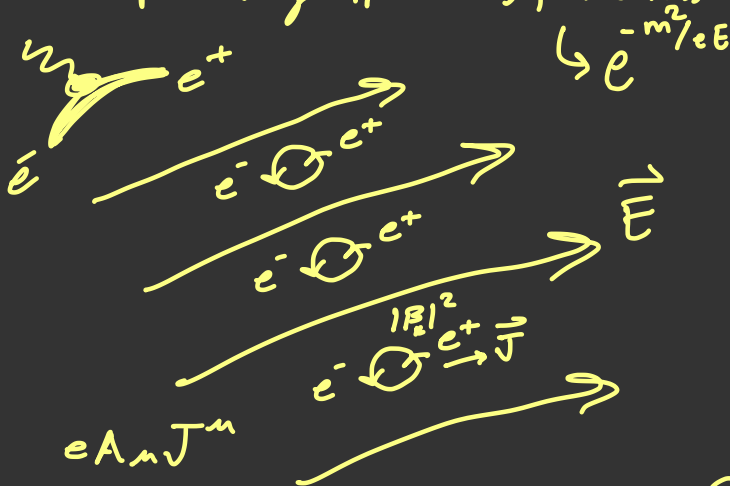
Inflationary Schwinger Dark Matter Production

- Schwinger effect long been known (Schwinger 1951)

↳ pairs of (charged) particles and anti-particles ($e^+ e^-$) created by background E-field


↳ Requires strong E-field in Minkowski

↳ Exponentially suppressed by particle mass



- Background gauge field leads to time-dependent dispersion relation

$$\ddot{g}_k + \omega_k^2 g_k = 0$$



 $\Downarrow \omega_k^2 < 0$

particle production

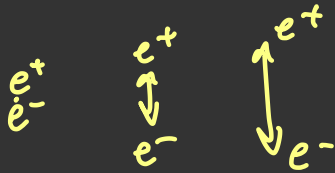
- During inflation the exp. of space leads to time dependence \Rightarrow Particle production

\hookrightarrow Requires breaking of conformal sym.

$$m \neq 0$$

or coupling to gravity / inflaton (2)

- In Schwinger + Inflation both expansion and the E-field contribute to induced charged current



- Leads to novel effects in charged PP
- can generate an induced current even with weak E-fields
- Can have PP even in $m \rightarrow 0$ (conformal) limit in fermion or scalar sector
- Allows for light DM which is typically difficult from inflation (3)

- Schwinger effect utilized in inflationary magnetogenesis scenarios

↳ primordial magnetic fields

- Also potential applications to reheating

Can inflationary Schwinger be used to generate DM relic abundance?

Schwinger Effect in de-Sitter Space

- Starting point is 'QED' action in dS space

$$S = \int d^4x \sqrt{-g} \left[-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \mathcal{L}_{\text{charged}}(A, \chi) \right] g_0 A_\mu J_x^\mu$$

χ can be scalar or fermion 'dark electrons'

$$D_\mu = \partial_\mu + i g_0 A_\mu + \Gamma_\mu$$

- Also need $\mathcal{L}_{\text{source}}(A, \psi)$ to provide BG E-field

- The induced current is given by

$$\frac{\delta \mathcal{L}}{\delta A_\mu} = \partial_\nu F^{\mu\nu} = J_x^\mu$$

- Induced current will be in direction of E -field which we take to be z -direction
- One can directly compute $\langle J_z \rangle$
- Current and E -field related by

$$\langle J_z \rangle = \sigma \langle E_z \rangle = \sigma E$$

Conductivity can be computed in terms of ratios

constant background E -field

$$\lambda \equiv \frac{g_0 E}{H^2}, \bar{m}_x \equiv \frac{m_x}{H}$$

- $\langle J_z \rangle$ is formally divergent \Rightarrow Renorm.
(some discrepancies in literature)
 \hookrightarrow negative currents?

- Need a 'source' to maintain constant background E -field (due to expansion)
 \hookrightarrow various mechanisms possible

- With σ and E can compute energy density of Schwinger produced X during inflation

- From conservation of energy-momentum $\nabla_\mu T^{\mu\nu} = 0$

$$\dot{\rho}_X + n H \rho_X = \langle E \cdot J \rangle = \sigma E^2$$

$n=3$
 $n=4$

↑
 Hubble damping

- Can integrate energy evolution equation to obtain

$$\rho_X \simeq \frac{\bar{\sigma}}{n} E^2 = \frac{2}{n} \bar{\sigma} \rho_E \quad (\bar{\sigma} \equiv \frac{\sigma}{H})$$

- To avoid backreaction on E -field requires $\rho_X \ll \rho_E \Rightarrow \bar{\sigma} < 1$

- We see the energy density in X is propto ρ_E , but suppressed by $\bar{\sigma}$

- Energy density in dark sector dominated by E-field during inflation
- Suppression by $\bar{\sigma}$ can be compensated for by redshifting effects after inflation
- For large enough M_X , X can make up part or even all of the DM ($M_X \gg M_A$)
 $(M_A \gtrsim 10^{-10} \text{ eV})$
- To compute relic abundance, need to know how ρ_X and ρ_E evolve
- We assume ρ_X and ρ_E evolve independently \Rightarrow no thermalization
 \Downarrow
Upper bound on dark gauge coupling

Cosmic Evolution and Relic Abundance

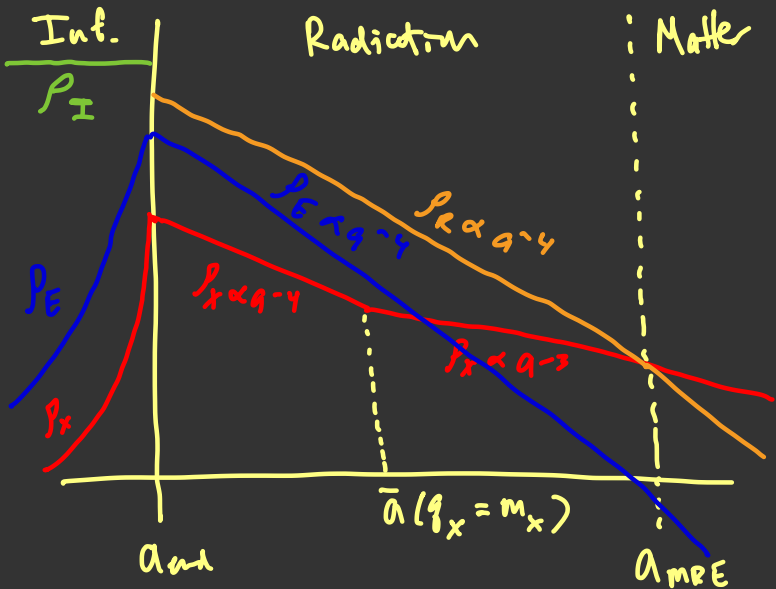
- At the end of inflation we have
- $\rho_{\text{I}}^{\text{end}} = 3H_{\text{end}}^2 M_{\text{pl}}^2$ • $\rho_{\text{E}}^{\text{end}} = \frac{1}{2} E_{\text{end}}^2$
- $\rho_{\text{R}} = E_{\text{R}}^4 \rho_{\text{I}}^{\text{end}}$ • $\rho_{\text{X}}^{\text{end}} = \frac{\bar{\sigma}}{n} E_{\text{end}}^2$
 $= \frac{\pi^2}{30} g_*(T_{\text{RH}}) \underline{T_{\text{RH}}^4}$
- Need to ensure the hierarchy

$$\rho_{\text{I}}^{\text{end}} \gg \rho_{\text{R}} \gg \rho_{\text{E}}^{\text{end}} \gg \rho_{\text{X}}^{\text{end}}$$

↓ ↓ ↓

$$E_{\text{R}} < 1 \quad H_{\text{max}} \quad \bar{\sigma} < 1$$

- No BR on inflation $\rho_{\text{E}}^{\text{end}} < \rho_{\text{I}}^{\text{end}}$
 $\Rightarrow E / M_{\text{end}}^2 < 10^9$
- The A and X energy density spectra which is peaked at small scales
↳ can avoid isocurvature constraints



- The energy density today m_x
- $m_x < \rho_{rad}$

$$\rho_x(T_0) = \left(\frac{T_0}{T_{RH}}\right)^3 \left(\frac{m_x}{\rho_{rad}}\right) \frac{\sigma}{4} E^2$$

$m_x > \rho_{rad}$

$$\rho_x(T_0) = \left(\frac{T_0}{T_{RH}}\right)^3 \frac{\sigma}{3} E^2$$

- Can now examine contours of $\frac{\Omega_x}{\Omega_{\text{com}}} = \frac{\rho_x(T_0)}{\rho_{\text{com}}} = 1$

- See results slides

Summary

- $M_x < M_{\text{end}} : 0.1 \text{ eV} \lesssim m_x \lesssim 10^{12} \text{ GeV}$
 $M_x > M_{\text{end}} : 100 \text{ GeV} \lesssim m_x \lesssim 10^{15} \text{ GeV}$

Ongoing Work

- Computing current in de-Sitter
- Backreaction and reheating effects

Future Work

- Exploring phenomenology
- Multi-component feebly interacting DM

Generating Dark E-fields During Inflation

- Need a 'source' to maintain constant E-field during inflation due to dilution from expansion

↳ Various possibilities from VDM production mechanisms

couplings to metric
 $M_a^2 g_{\mu\nu} A^\mu A^\nu, R A^\mu A_\mu$

coupling to inflaton
 $\alpha F\tilde{F}, \alpha\tilde{F}\tilde{F}$

- VDM production mechanism must generate background E-field (\vec{A}) with coherence length

$$\lambda \sim H^{-1}$$



- Recent mechanism based on $\alpha F\tilde{F}$ can generate necessary E-field during inflation \Rightarrow works for $m_A = 0$

(1810.07208, 2103.12145 : MBG, JS, LU, RM)

- Utilizes the source lagrangian

$$\mathcal{L}_{\text{source}} = (\partial_\mu \psi)^2 - V(\psi) + \frac{\alpha}{f} \psi F \tilde{F}$$

$$\tilde{F} = \epsilon^{\mu\nu\alpha\beta} F_{\alpha\beta}$$

- This leads to EOM for transverse modes

$$\ddot{A}_\pm + H(1+\sigma)\dot{A}_\pm + \left(\frac{k^2}{a^2} + \frac{\alpha}{f}\dot{\psi}\frac{k}{a}\right)A_\pm = 0$$

where we have use $E = \frac{1}{a}\dot{A}_\pm$

- In limit of small conductivity (weak field) we can ignore this correction and find (in conformal time)

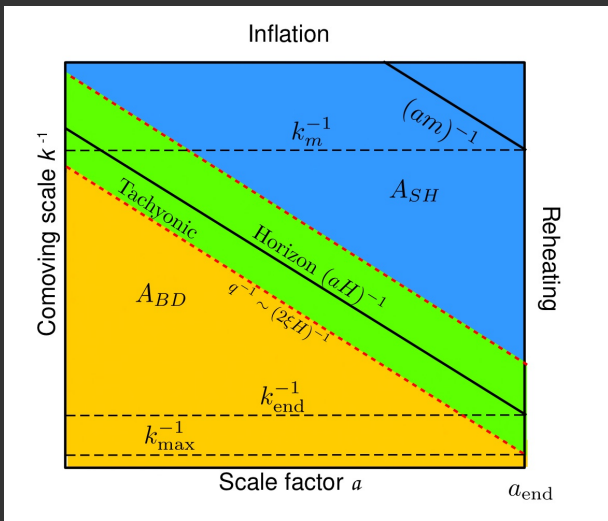
$$\frac{d^2 A_\pm}{d\tau^2} + \omega^2(k, \tau) = 0$$

$$\omega^2(k, \tau) = k^2 \pm \frac{2k\mathcal{E}}{\tau}$$

$$\mathcal{E} = \frac{\alpha\dot{\psi}}{24f}$$

- When $\omega^2 < 0$ there is a tachyonic inst.

- This occurs when $\frac{\dot{u}}{u} = \dot{g} \sim H$
- Leads to exponential production of dark vector modes \Rightarrow high occupation number
- Coherence length $l \sim H^{-1}$
- Classical background E-field within each Hubble patch



Phenomenology of DM sector

- Dark sector may or may not couple to SM

dark vector A

dark scalar ϕ

dark fermion N
 \hookrightarrow "steril neutrino"

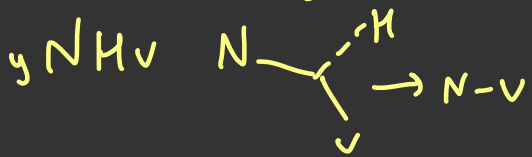
"kinetic mixing portal"



"Higgs portal"



"neutrino portal"



- Working on exploring different portals
- Even if no DM-SM coupling there is a lot of potential phenomenology
 - ↳ structure
 - ↳ superradiance
 - ↳ lensing
 - ↳ gravitational waves
- Studying σ in detail for fermion and scalar χ
- Constructing explicit U(1) models during inflation