



# Overcoming Thresholds for Dark Matter Search

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DM Landscape @ MITP  
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李政道研究所

Tsung-Dao Lee Institute

## 1) Thresholds in DM search

## 2) Fermionic DM Absorption & Nucleon Consumption

**SFG**, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [[JHEP 05 \(2022\) 191](#)]

PandaX + **SFG**, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [[Phys.Rev.Lett. 129 \(2022\) 16, 161804](#)]

**SFG**, Kai Ma, Xiao-Dong Ma, Jie Sheng [[JHEP 11 \(2023\) 190](#)]

Kai Ma, **SFG**, Lin-Yun He, Ning Zhou [[arXiv:2405.16878](#)]

**SFG**, Oleg Titov [[arXiv:2405.05728](#)]

**SFG**, Xiao-Dong Ma [[arXiv:2406.00445](#)]

## 3) Reactivate Forbidden DM @ Supermassive Black Hole

Yu Cheng, **SFG**, Xiao-Gang He, Jie Sheng [[Phys. Lett. B 847 \(2023\) 138294](#)]

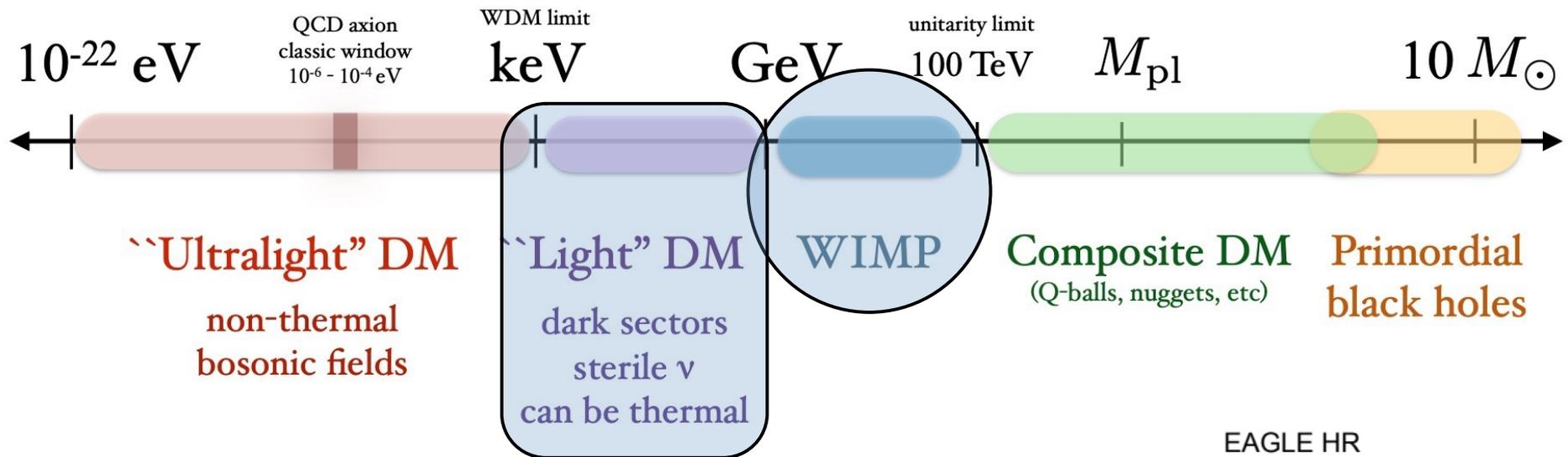
Yu Cheng, **SFG**, Jie Sheng, Tsutomu Yanagida [[PhysRevD.107.123013](#)]

Yu Cheng, **SFG**, Jie Sheng, Tsutomu Yanagida [[arXiv:2309.12043](#)]

## 4) MeV Solar DM

**SFG**, Jie Sheng, Chen Xia, Chuan-Yang Xing [[arXiv:2408.12448](#)]

# Mass Span of DM

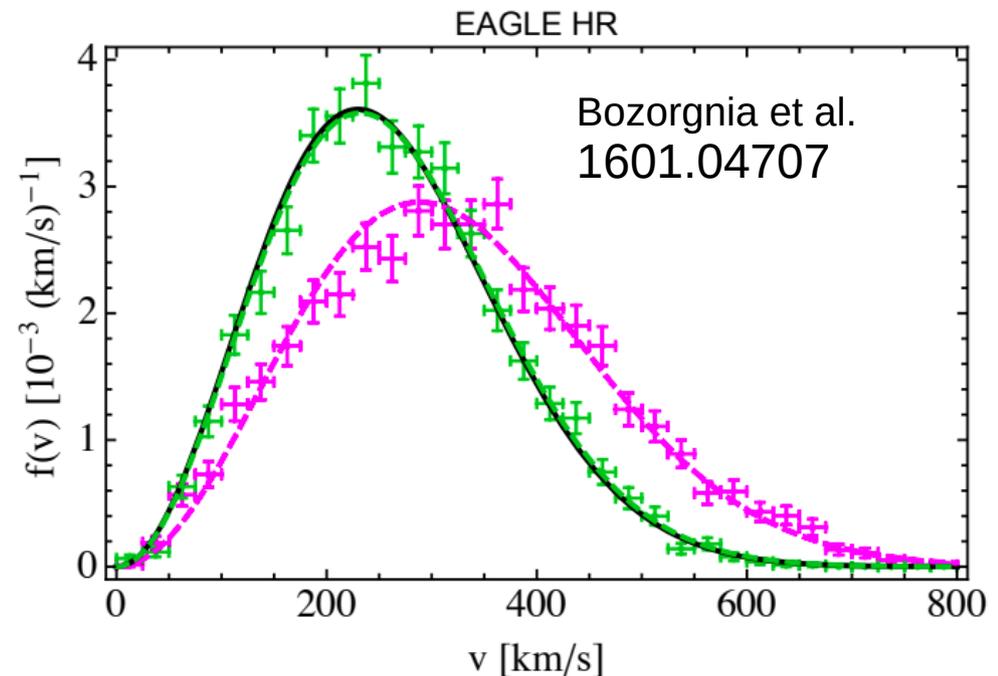


?

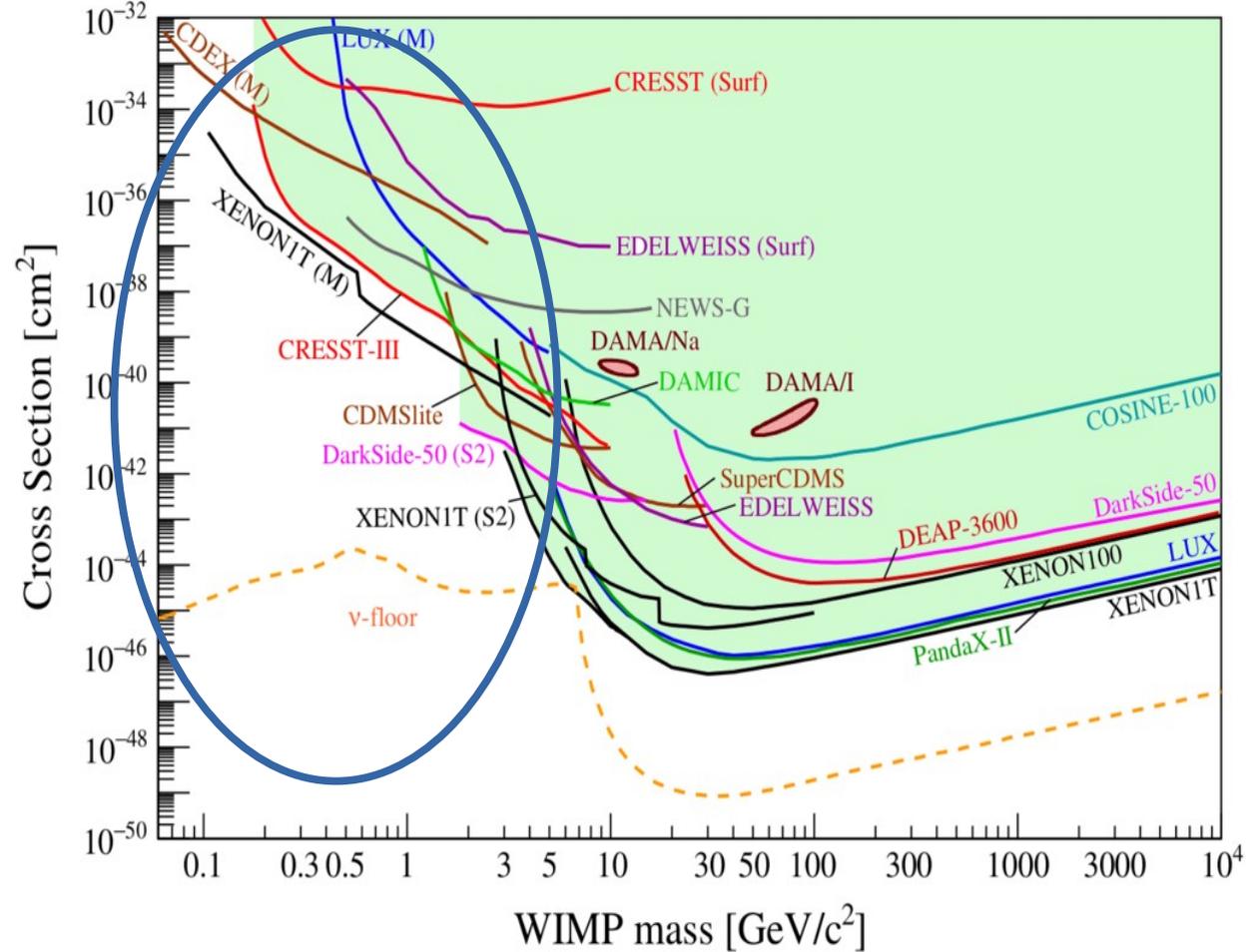
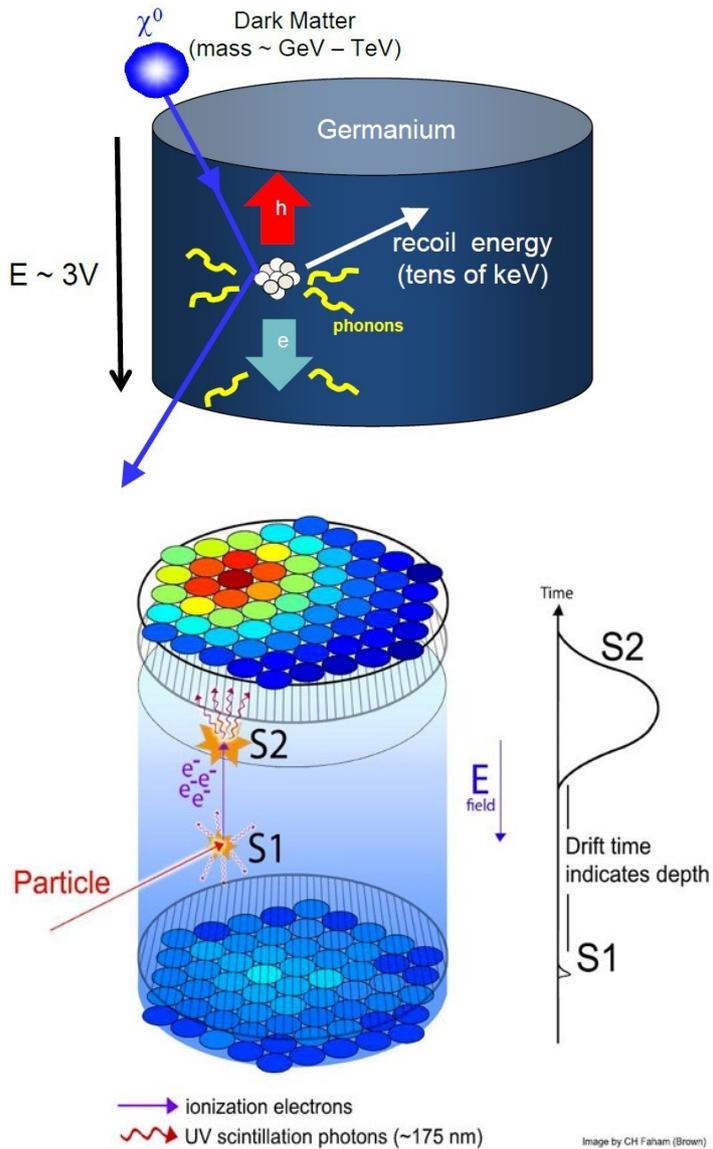
$$E_r \approx \frac{4m_\chi m_N}{(m_\chi + m_N)^2} T_\chi$$

$$\approx \frac{4m_\chi}{m_N} T_\chi \quad T_\chi = \frac{1}{2} m_\chi v_\chi^2$$

$$E_r \propto m_\chi^2$$



# Direct Detection



APPEC Committee Report [2104.07634]

## 1) Lowering the threshold

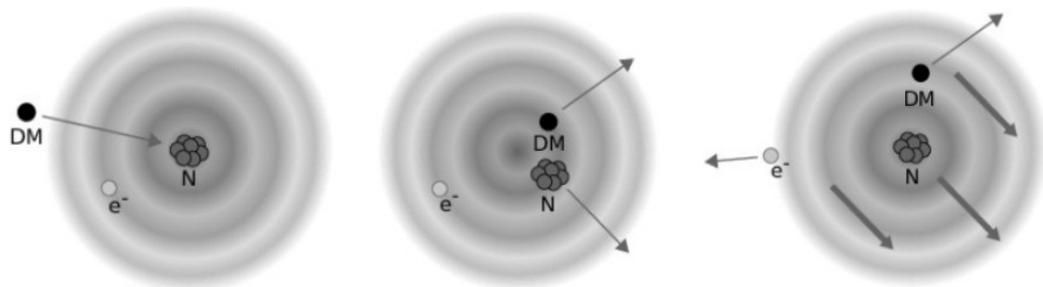
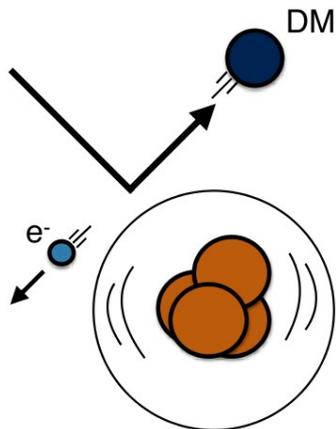
**Bolometer** [1904.00498, Ann.Rev.Nucl.Part.Sci. 67 (2017) 161-181]

**Bremsstrahlung** [Kouvaris & Pradler, PRL 118, 031803 (2017)]



**Migdal effect** [Ibe et al [1707.07258]]

## 2) Electron target



## 3) Boost detector?

## 4) Boost DM

SFG, Jianglai Liu, Qiang Yuan, Ning Zhou, [[Phys.Rev.Lett. 126 \(2021\) 9, 091804](#)]

PandaX-II + SFG & Qiang Yuan [[Phys.Rev.Lett. 128 \(2022\) 17, 171801](#)]

- 1) Thresholds in DM search
- 2) Cosmic ray boosted DM & diurnal modulation
- 3) **Fermionic DM Absorption**
- 4) Reactivate Forbidden DM @ Supermassive Black Hole

# Fermionic Absorption

**SFG**, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [[JHEP 05 \(2022\) 191](#)]

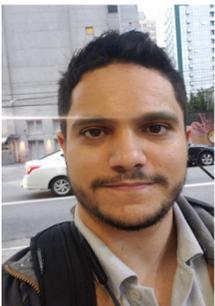
PandaX + **SFG**, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [[Phys.Rev.Lett. 129 \(2022\) 16, 161804](#)]

**SFG**, Pedro Pasquini, Jie Sheng [[JHEP 05 \(2022\) 088](#)]

**SFG**, Kai Ma, Xiao-Dong Ma, Jie Sheng [[JHEP 11 \(2023\) 190](#)]

Kai Ma, **SFG**, Lin-Yun He, Ning Zhou [[arXiv:2405.16878](#)]

**SFG**, Oleg Titov [[arXiv:2405.05728](#), to appear in PRD]



Pedro Pasquini

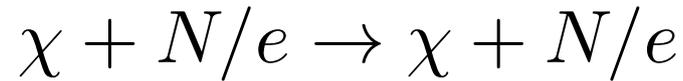


Jie Sheng



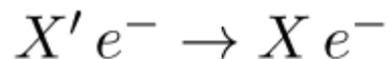
Oleg Titov

- Elastic Scattering



$$E_r \approx \frac{4m_\chi m_N}{(m_\chi + m_N)^2} T_\chi$$

- Exothermic

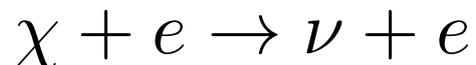


$$E_R \simeq \Delta m \left( 1 - \frac{v_{\text{DM}}}{v_e} \cos \theta_e \right)$$

He, Wang & Zheng [JCAP21, 2007.04963]

Aboubrahim, Althueser, Klasen, Nath & Weinheimer [2207.08621]

- Fermionic absorption



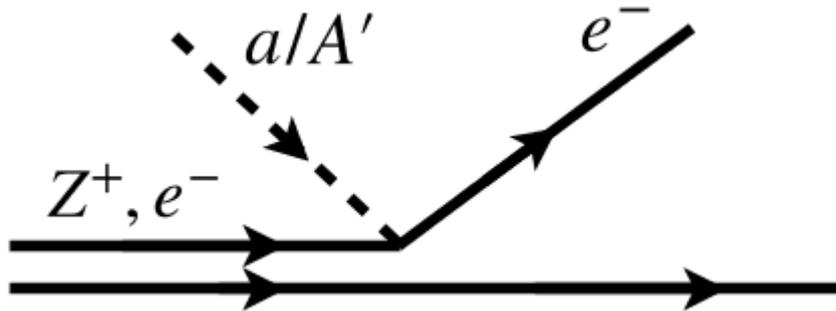
$$E_r \approx \frac{m_\chi^2}{2m_e}$$

See also Dror, Elor & McGehee  
[1905.12635, 1908.10861]

Li, Liao & Zhang [2201.11905]

Dror, Elor, McGehee & Yu [2011.01940]

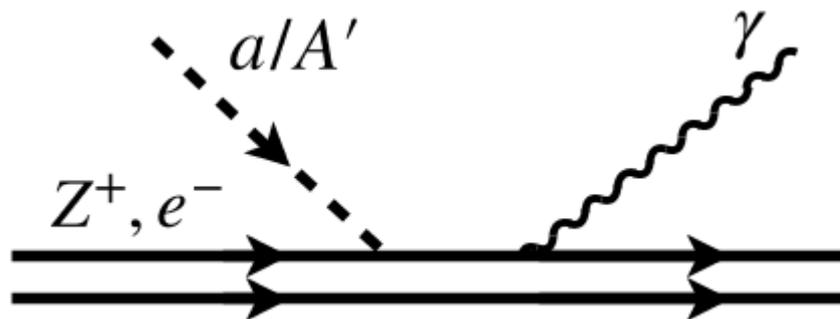
- Dark absorption



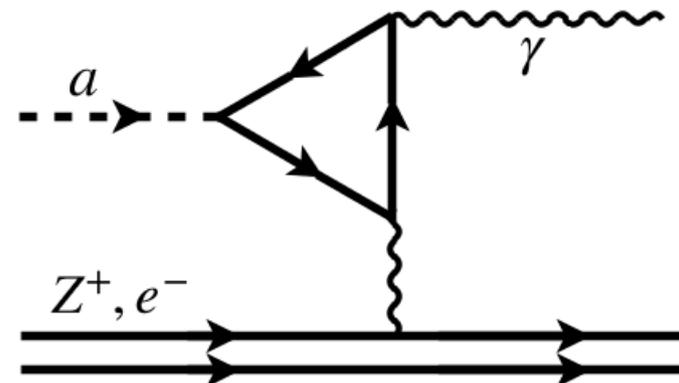
Pospelov, Ritz & Voloshin [0807.3279]

An, Pospelov, Pradler & Ritz [1412.8378]

- Dark Compton



- Inverse Primakoff



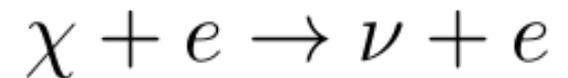
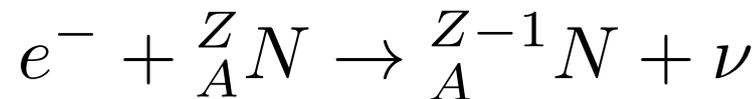
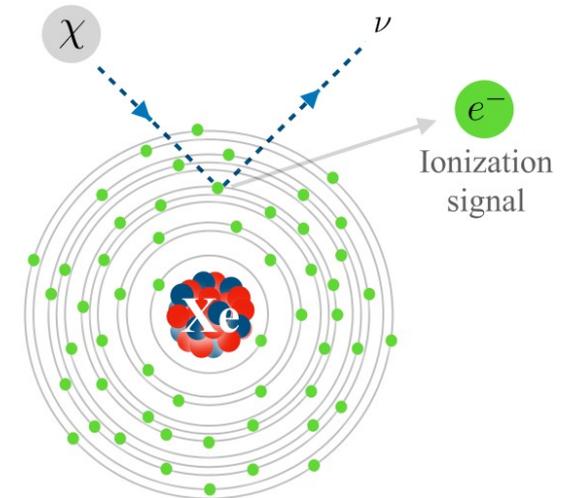
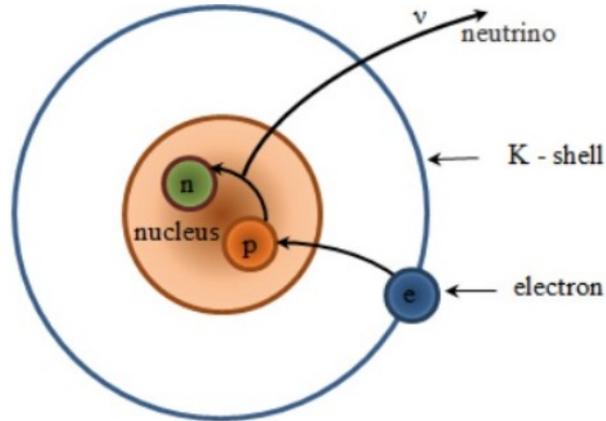
Hochberg, Krosigk, Kuflik & Yu [2109.08168]

# K-Shell Electron Capture for $\nu$



王淦昌

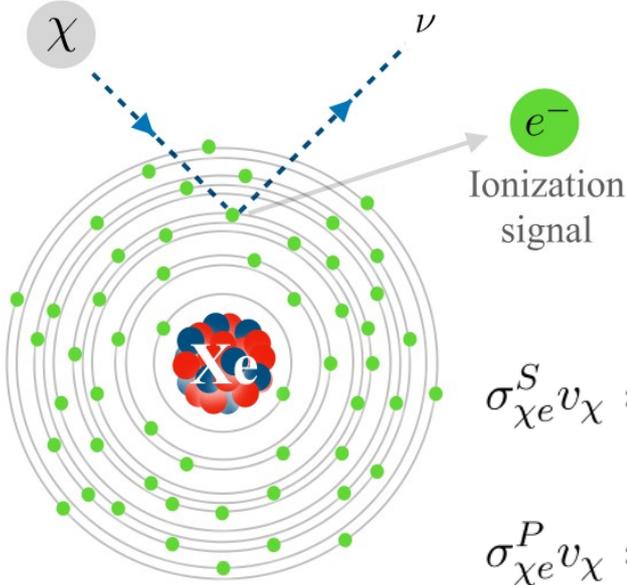
1942 – Kan Chang Wang proposed using K-shell electron capture for detecting neutrino



Kan Chang Wang, *A Suggestion on the Detection of the Neutrino*, Phys. Rev., 61, 97 (1942)

Kan Chang Wang, *Proposed Methods of Detecting the Neutrino*, Phys. Rev., 71, 645-646 (1947)

# Effective Operators



$$\begin{aligned} \mathcal{O}_{ev\chi}^S &\equiv (\bar{e}e)(\bar{\nu}_L\chi_R), \\ \mathcal{O}_{ev\chi}^P &\equiv (\bar{e}i\gamma_5e)(\bar{\nu}_L\chi_R), \\ \mathcal{O}_{ev\chi}^V &\equiv (\bar{e}\gamma_\mu e)(\bar{\nu}_L\gamma^\mu\chi_L), \\ \mathcal{O}_{ev\chi}^A &\equiv (\bar{e}\gamma_\mu\gamma_5e)(\bar{\nu}_L\gamma^\mu\chi_L), \\ \mathcal{O}_{ev\chi}^T &\equiv (\bar{e}\sigma_{\mu\nu}e)(\bar{\nu}_L\sigma^{\mu\nu}\chi_R), \end{aligned}$$

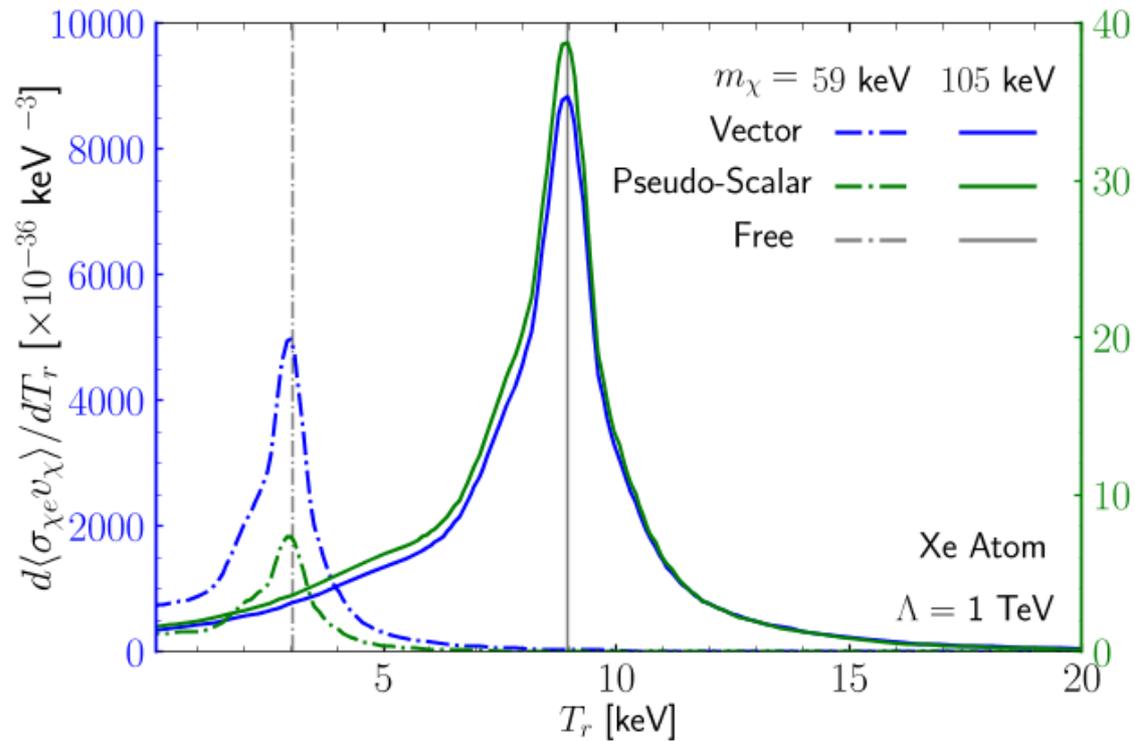
$$\begin{aligned} \sigma_{\chi e}^S v_\chi &\approx \frac{1}{\Lambda^4} \frac{m_\chi^2(2m_e + m_\chi)^4}{64\pi(m_e + m_\chi)^4}, \\ \sigma_{\chi e}^P v_\chi &\approx \frac{1}{\Lambda^4} \frac{m_\chi^4(2m_e + m_\chi)^2}{64\pi(m_e + m_\chi)^4}, \\ \sigma_{\chi e}^V v_\chi &\approx \frac{1}{\Lambda^4} \frac{m_\chi^2(2m_e + m_\chi)^2(2m_e^2 + 4m_em_\chi + 3m_\chi^2)}{32\pi(m_e + m_\chi)^4}, \\ \sigma_{\chi e}^A v_\chi &\approx \frac{1}{\Lambda^4} \frac{m_\chi^2(2m_e + m_\chi)^2(6m_e^2 + 8m_em_\chi + 3m_\chi^2)}{32\pi(m_e + m_\chi)^4}, \\ \sigma_{\chi e}^T v_\chi &\approx \frac{1}{\Lambda^4} \frac{m_\chi^2(2m_e + m_\chi)^2(6m_e^2 + 10m_em_\chi + 5m_\chi^2)}{8\pi(m_e + m_\chi)^4}. \end{aligned}$$

Dror, Elor, McGehee & Yu [2011.01940]

SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]

# Kinematics of Absorption DM

$$\chi + e \rightarrow \nu + e \quad E_r = \frac{m_\chi^2}{2(m_e + m_\chi)}$$

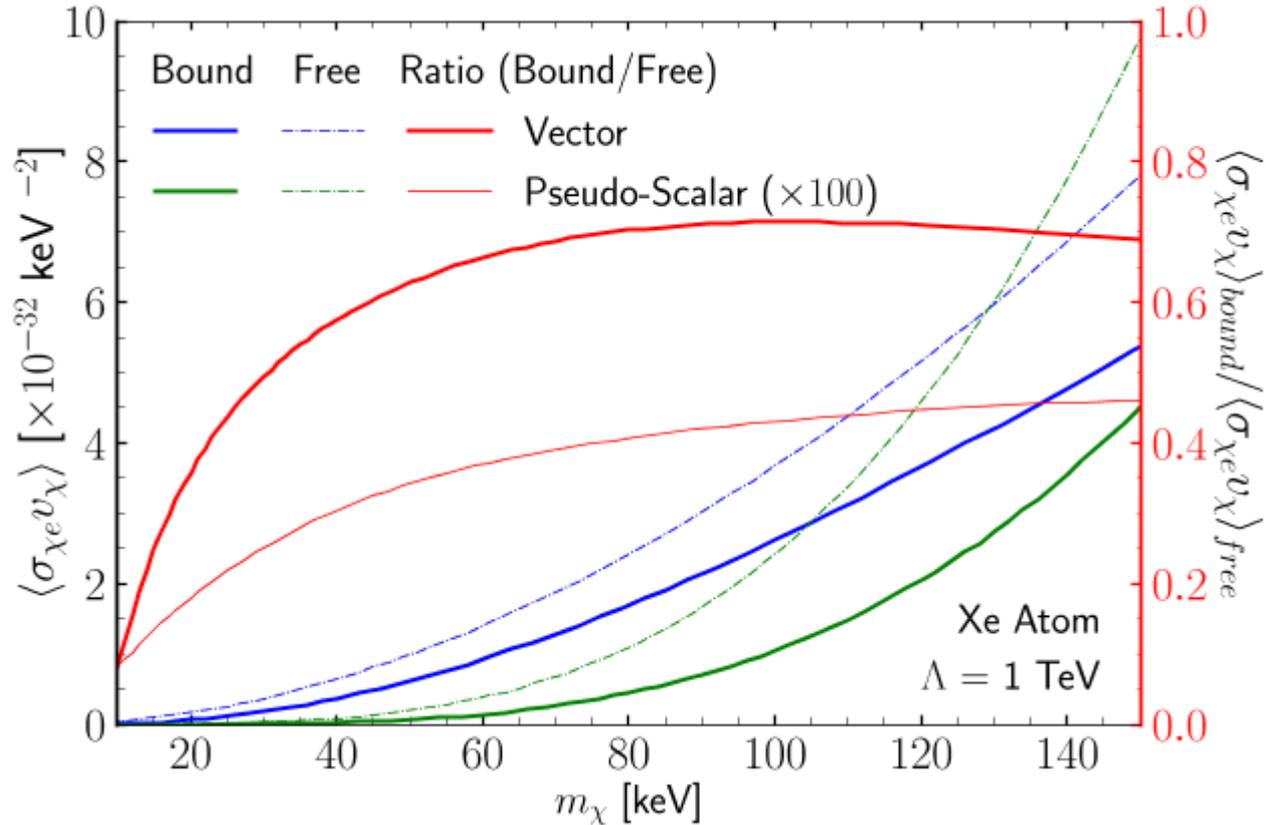


$$\frac{d\langle\sigma_{\chi e\nu_\chi}\rangle}{dT_r} = \sum_{nl} (4l + 2) \frac{1}{T_r} \frac{m_\chi - \Delta E_{nl}}{16\pi m_e^2 m_\chi} |\mathcal{M}|^2(\mathbf{q}) K_{nl}(T_r, |\mathbf{q}|),$$

SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]

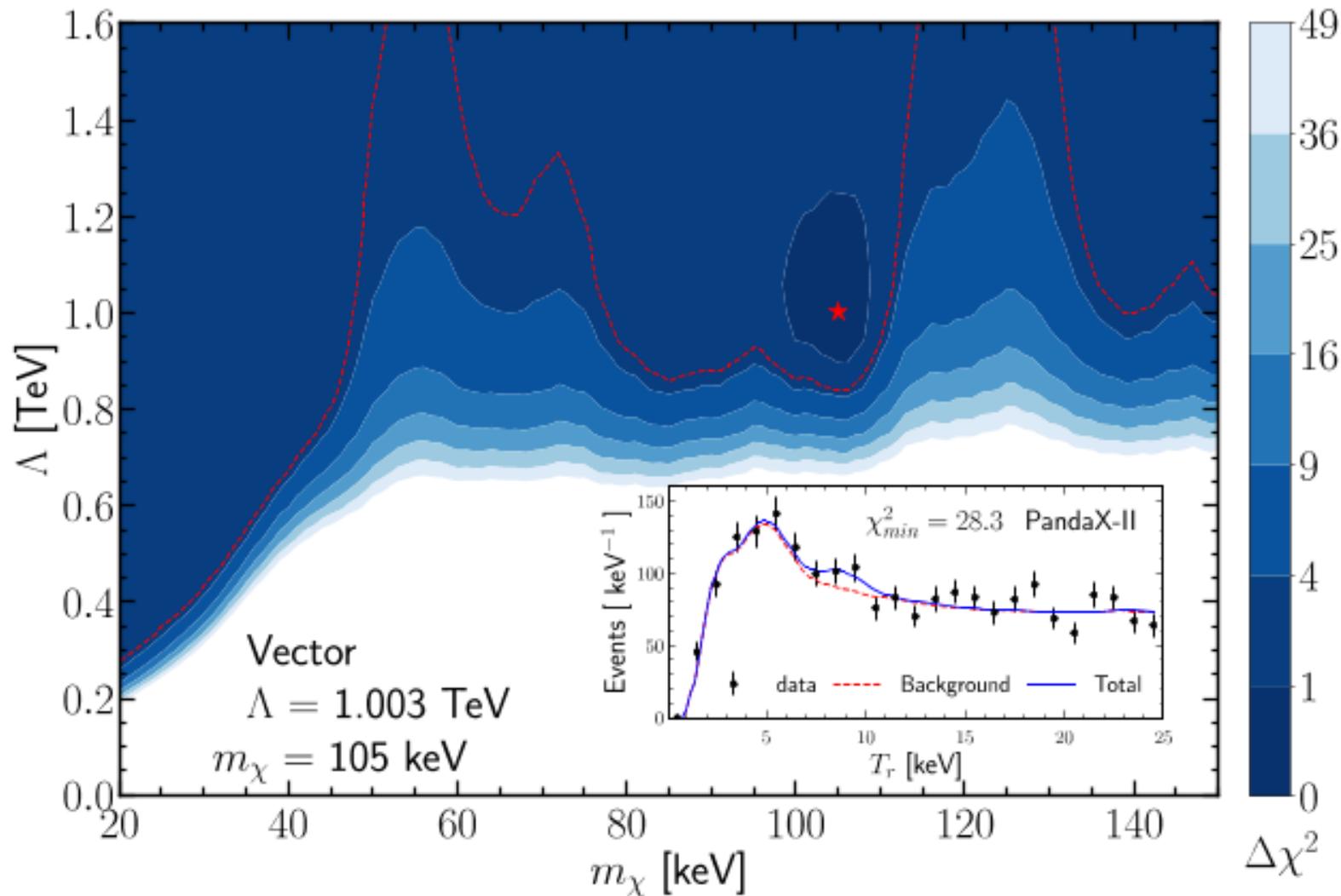
SFG, Pedro Pasquini, Jie Sheng [JHEP 05 (2022) 088]

$$\chi + e \rightarrow \nu + e$$



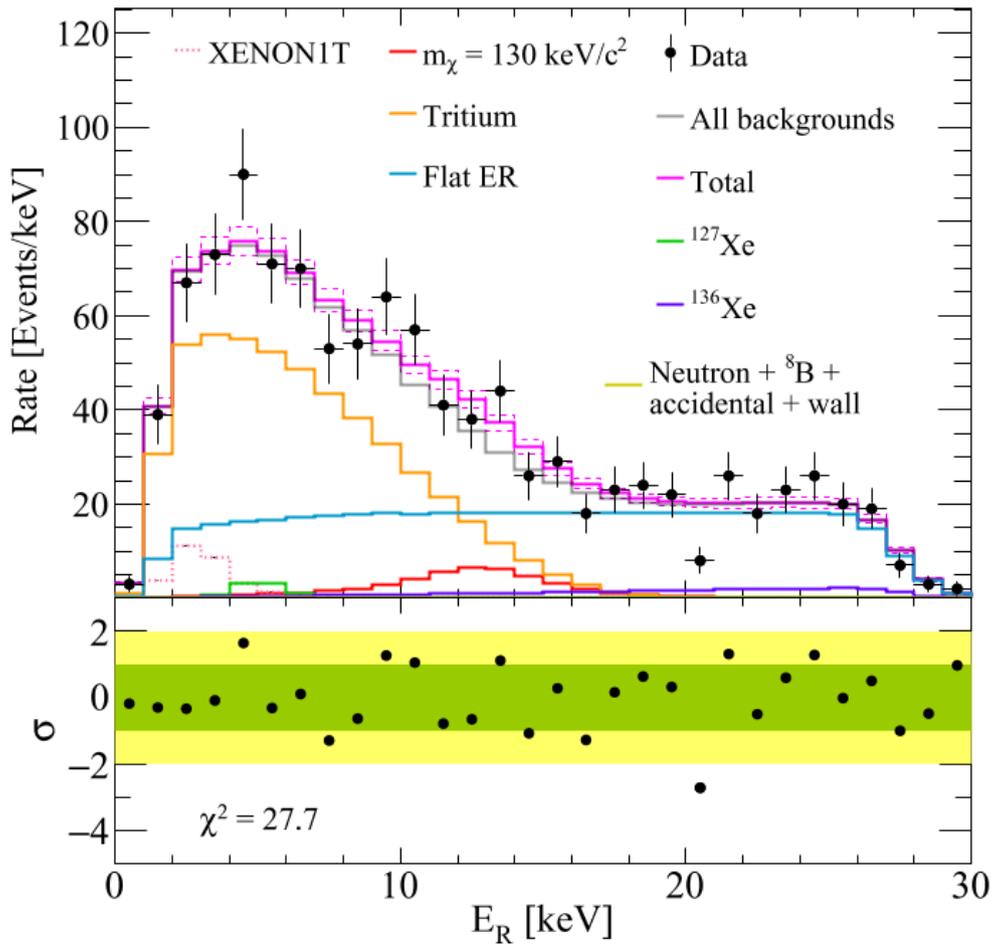
$$\frac{d\langle \sigma_{\chi e \nu \chi} \rangle}{dT_r} = \sum_{nl} (4l + 2) \frac{1}{T_r} \frac{m_\chi - \Delta E_{nl}}{16\pi m_e^2 m_\chi} |\mathcal{M}|^2(\mathbf{q}) K_{nl}(T_r, |\mathbf{q}|),$$

SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]



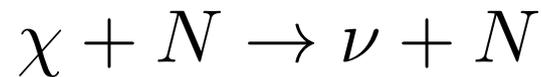
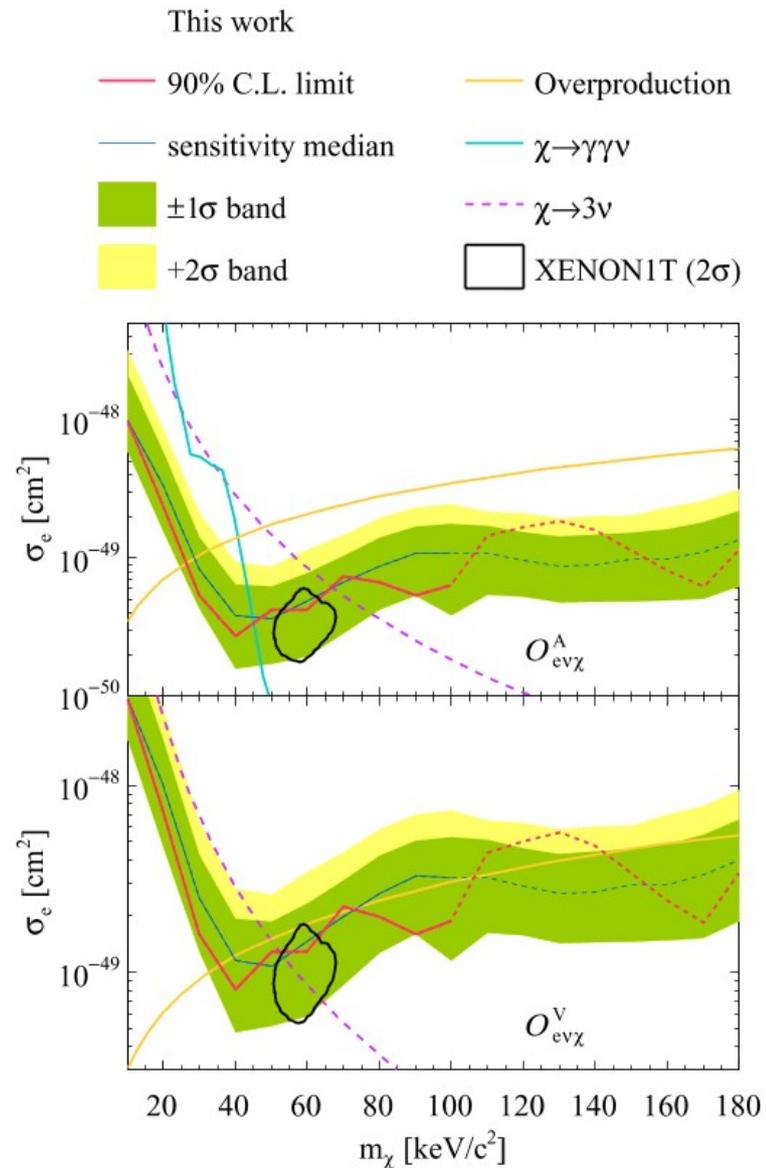
SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]

# PandaX-4T Results



PandaX + **SFG**, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng,  
*Phys.Rev.Lett.* 129 (2022) 16, 161804 [2206.02339]

See also PandaX, *Phys.Rev.Lett.* 129 (2022) 16, 161803 [2205.15771]  
CDEX, *Phys.Rev.Lett.* 129 (2022) 22, 221802 [2209.00861]

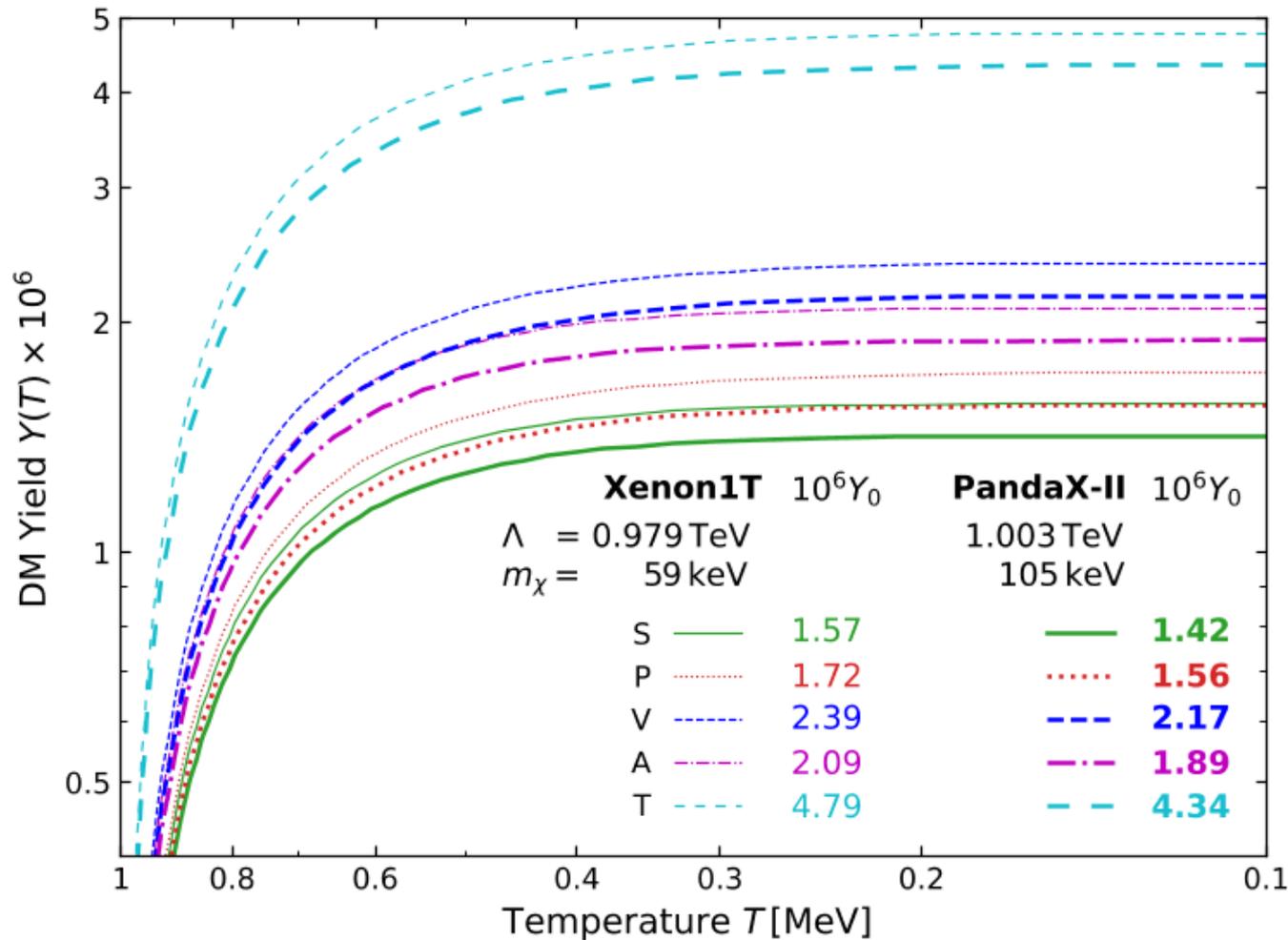


# Overproduction (Freeze-In)

$$e^-e^+ \rightarrow \nu\bar{\chi}/\bar{\nu}\chi$$

$$e^\pm\nu \rightarrow e^\pm\chi$$

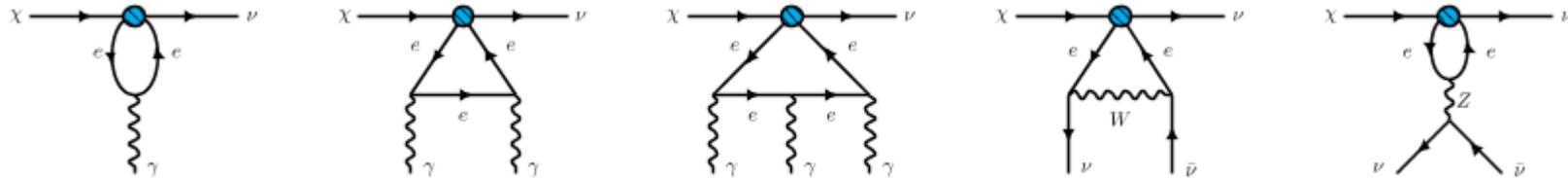
$$e^\pm\bar{\nu} \rightarrow e^\pm\bar{\chi}$$



SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [[JHEP 05 \(2022\) 191](#)]

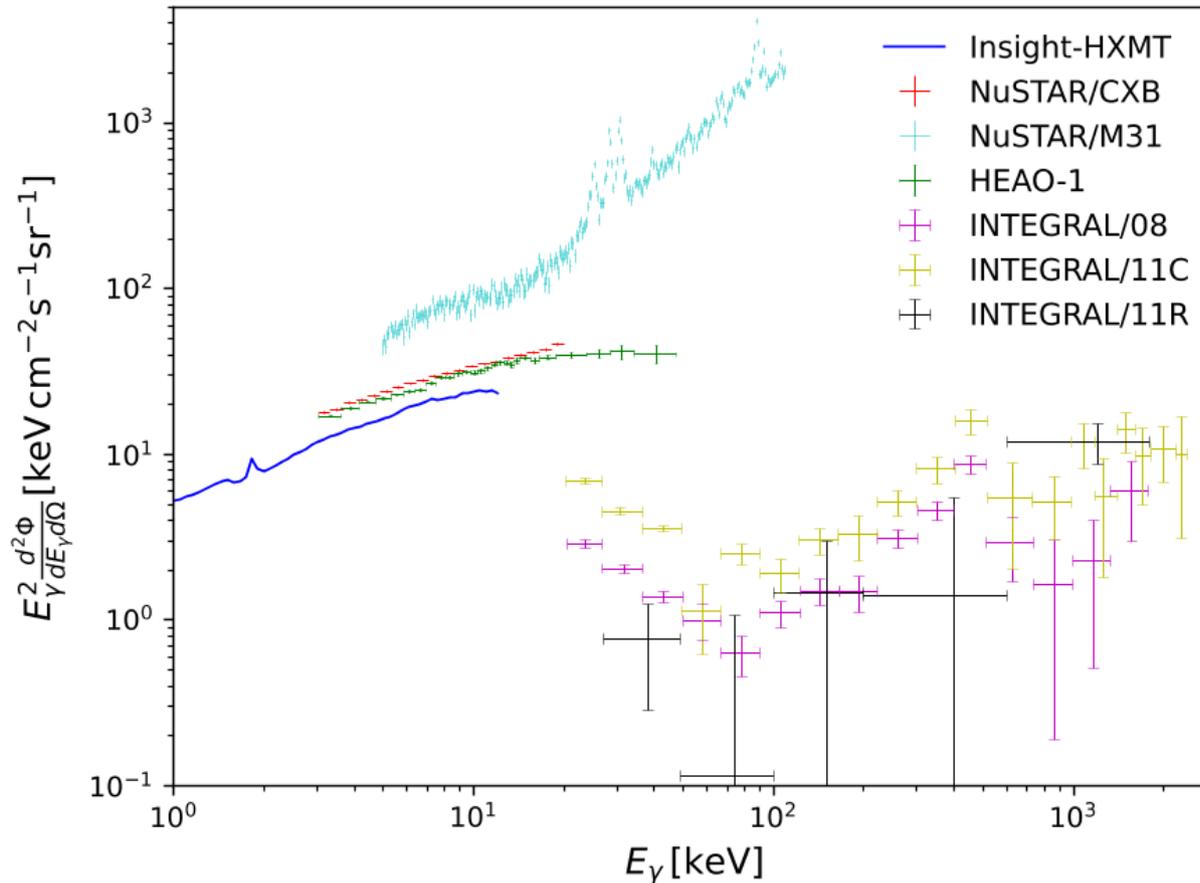
SFG, Kai Ma, Xiao-Dong Ma, Jie Sheng [[JHEP 11 \(2023\) 190](#)]

# Visible & Invisible Decays



Operator	Process			
	$\chi \rightarrow \nu \gamma$	$\chi \rightarrow \nu \gamma \gamma$	$\chi \rightarrow \nu \gamma \gamma \gamma$	$\chi \rightarrow 3\nu$
S: $\mathcal{O}_{e\nu\chi}^S$	×	✓	×	×
P: $\mathcal{O}_{e\nu\chi}^P$	×	✓	×	×
V: $\mathcal{O}_{e\nu\chi}^V$	×	×	✓	✓
A: $\mathcal{O}_{e\nu\chi}^A$	×	✓	×	✓
T: $\mathcal{O}_{e\nu\chi}^T$	✓	×	×!	×!

SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]



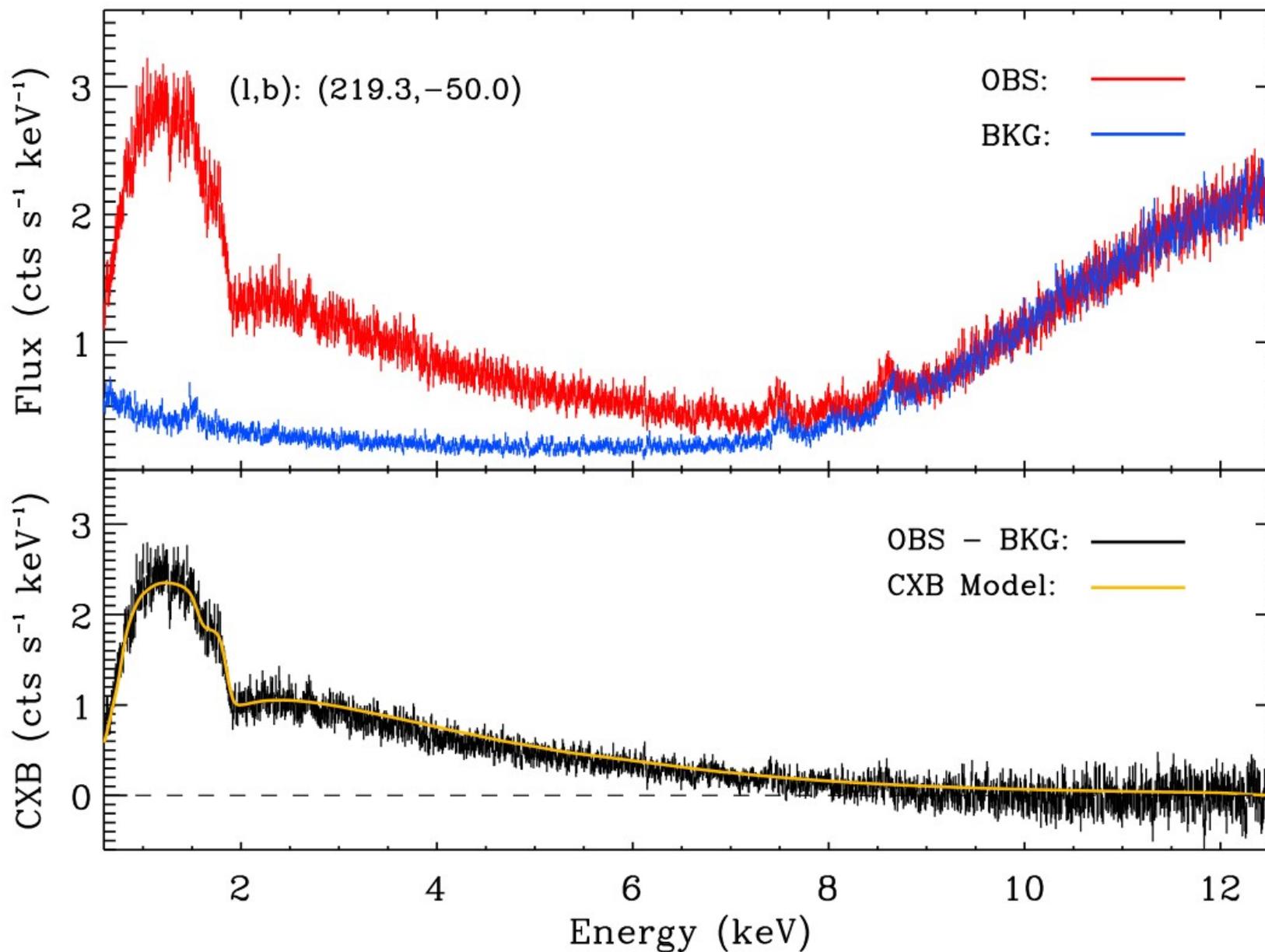
- Galactic

$$\frac{d^2\Phi_\gamma}{dE_\gamma d\Omega} = \frac{1}{4\pi} \frac{d\Gamma_\chi}{dE_\gamma} \int_{\text{l.o.s.}}^{s_{\text{max}}} \frac{\rho_\chi(r)}{m_\chi} ds$$

- Extra-Galactic

$$\frac{d^2\Phi_r^{\text{EG}}}{dE_\gamma d\Omega} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_\chi H_0 \sqrt{\Omega_m}} \int_0^\infty \frac{d\Gamma_\chi}{dE_\gamma(z)} \frac{dz}{\sqrt{\kappa + (1+z)^3}}$$

SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]

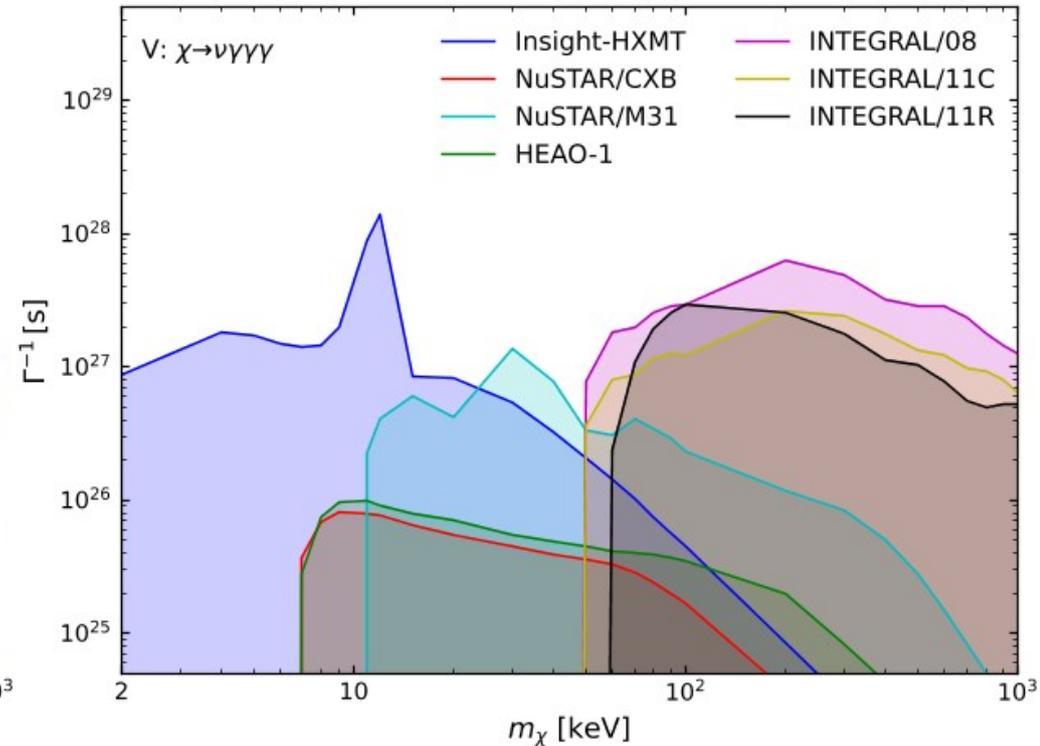
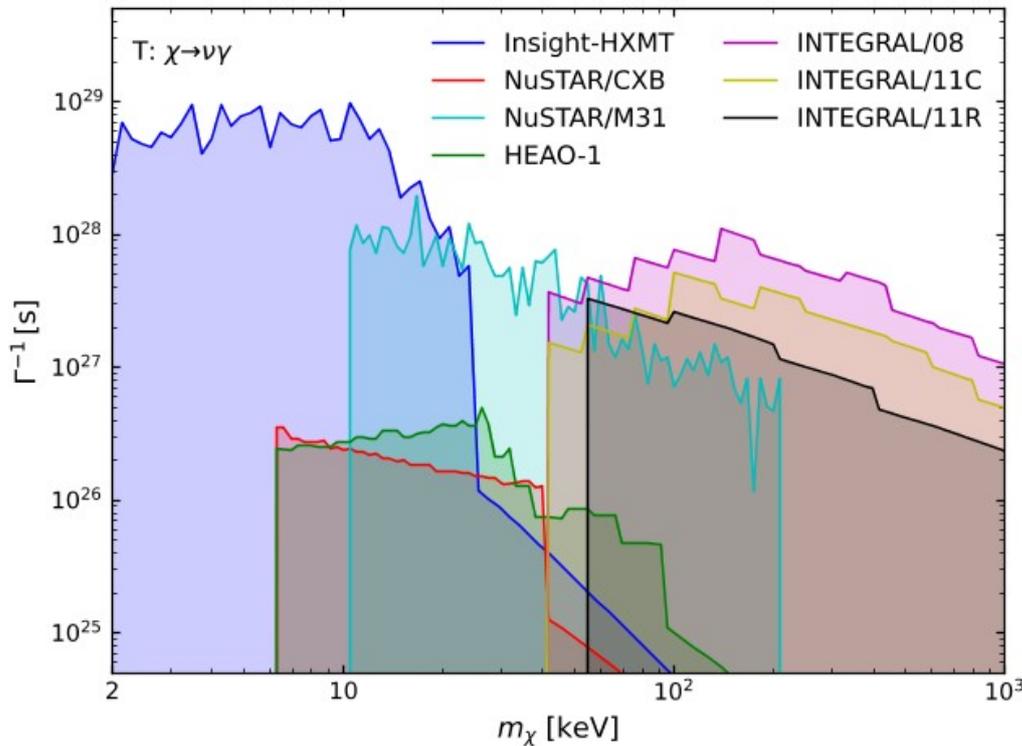


Jin-Yuan Liao et al [JHEAp 27 (2020) 24-32, arXiv:2004.01432]

$$N_i^{\text{th}} \leq N_i^{\text{obs}} \equiv A_{\text{eff}} T_{\text{obs}} \Delta\Omega \left( \frac{d^2\Phi_\gamma}{dE_\gamma d\Omega} \right)_{\text{exp@95\%}}^i \Delta E_i$$

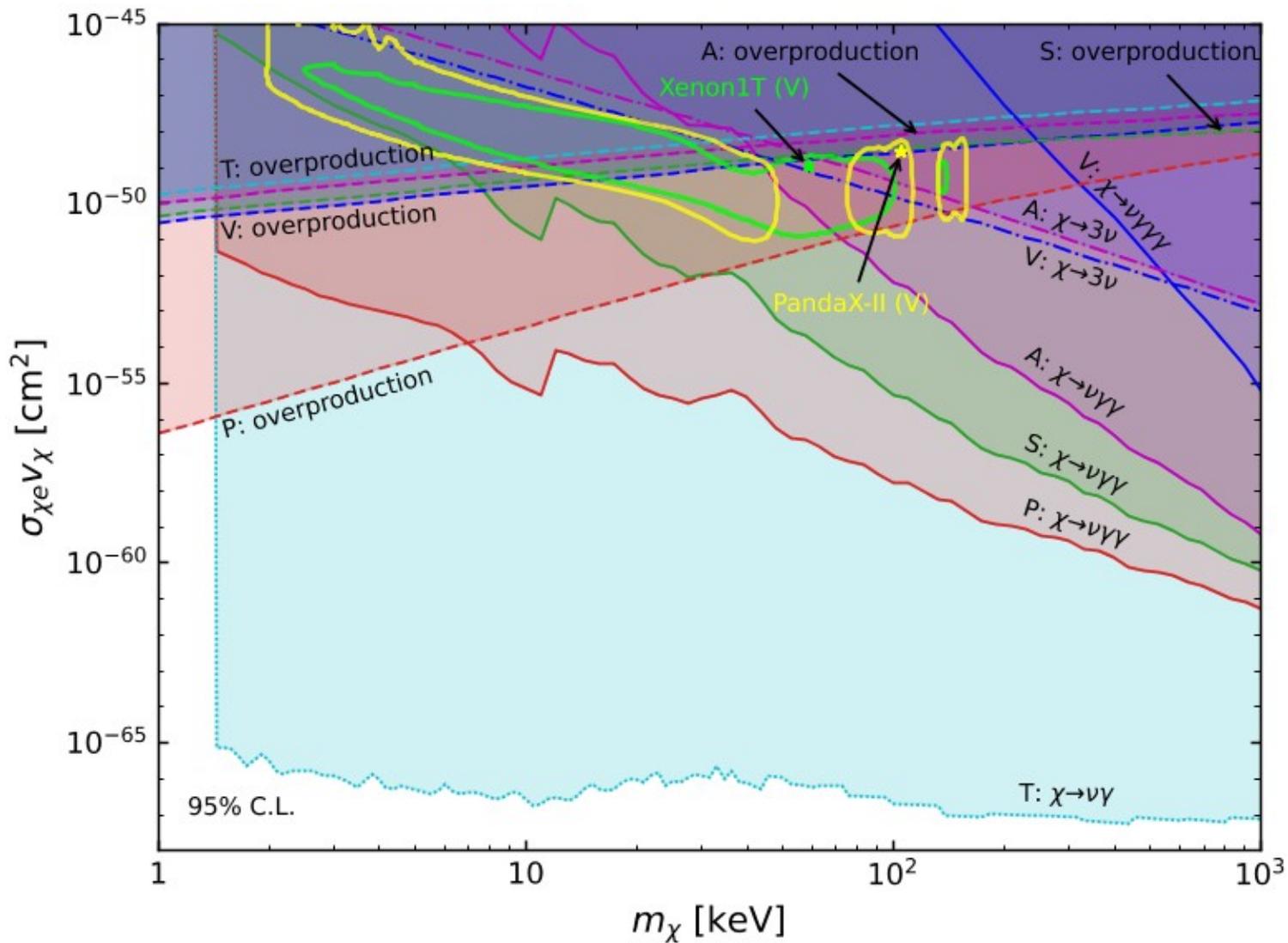
- Mono-energetic  $\gamma$

- Continuous Spectrum



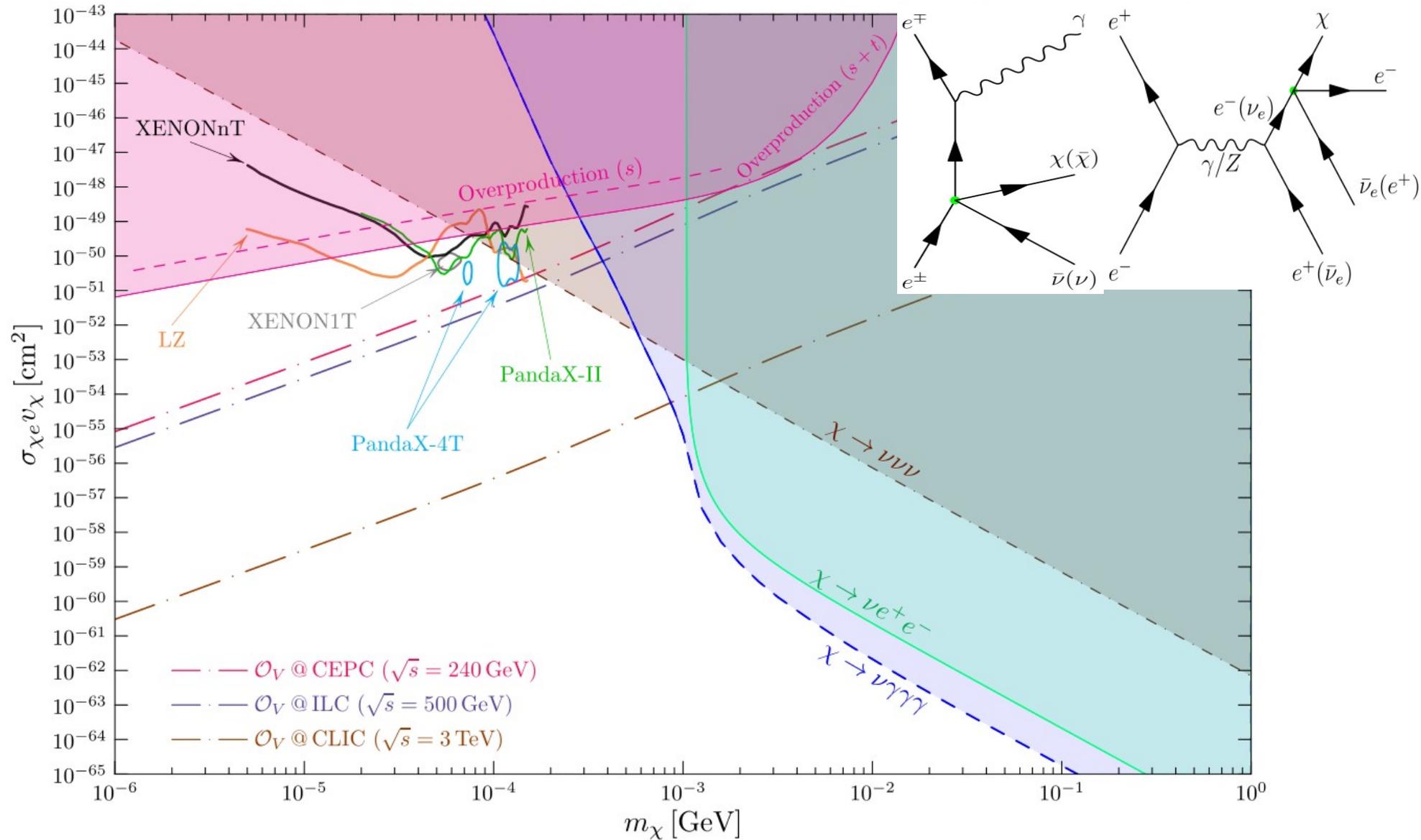
SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [[JHEP 05 \(2022\) 191](#)]

# Constraints from Astro & Cosmo



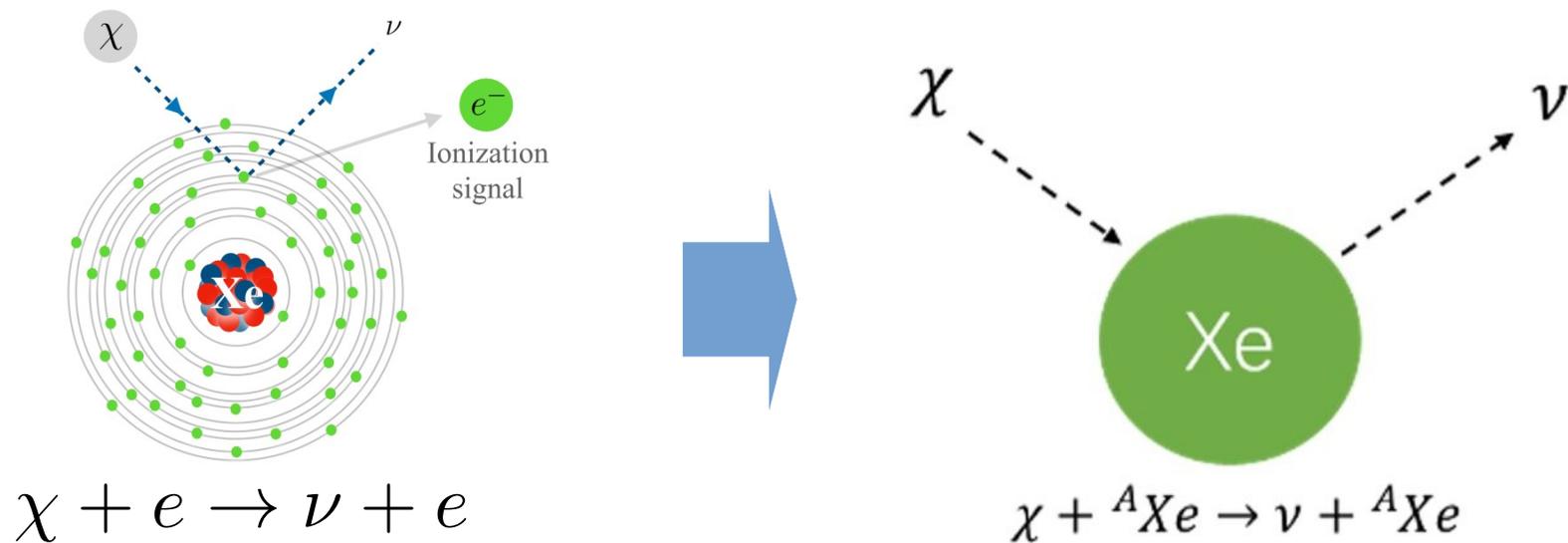
SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]

# PandaX-4T & Lepton Colliders



SFG, Kai Ma, Xiao-Dong Ma, Jie Sheng [JHEP 11 (2023) 190]

# Absorption @ Nuclear Targets



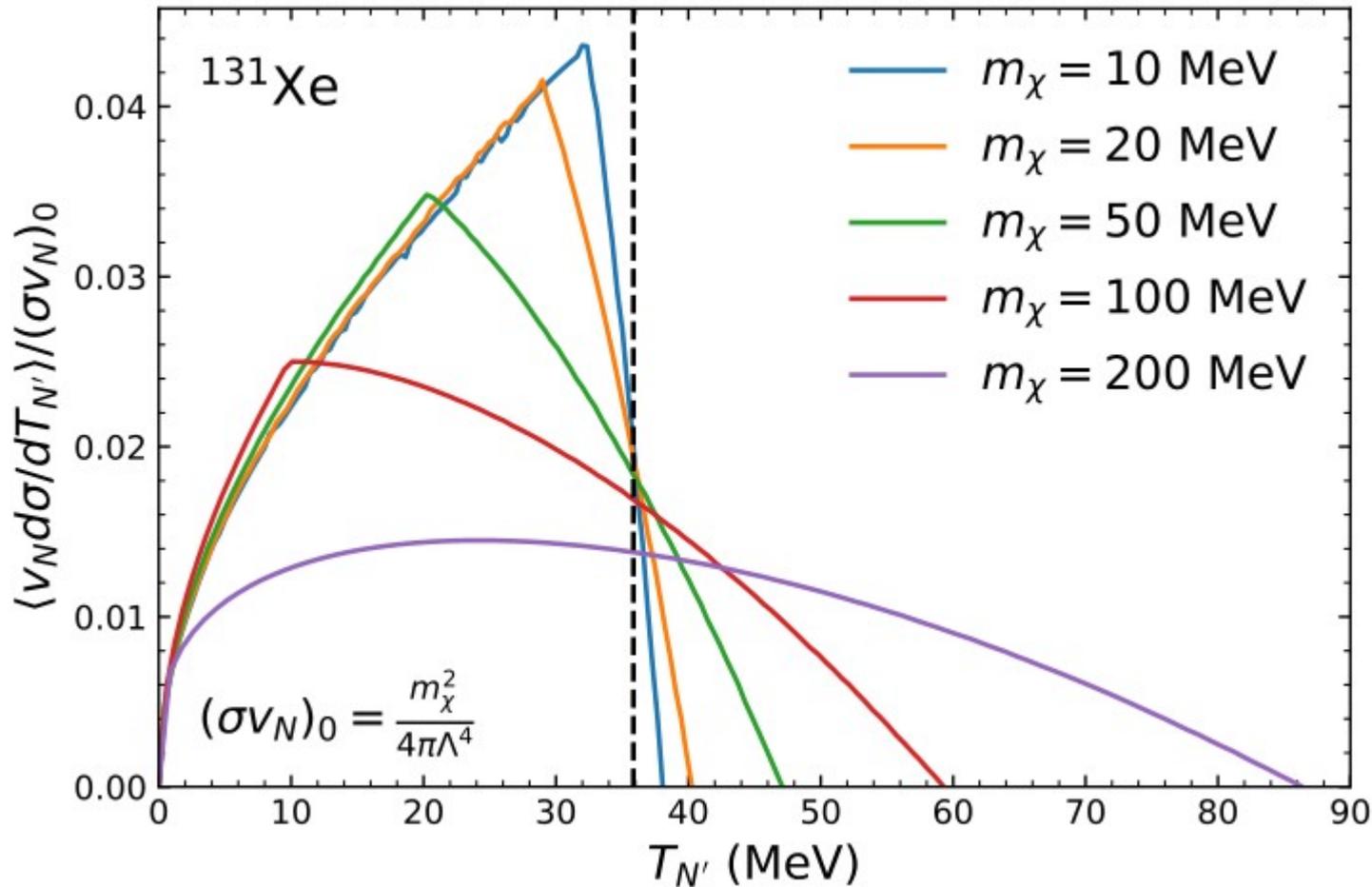
$$T_r = \frac{m_\chi^2}{2(m_A + m_\chi)} \gtrsim 1 \text{ keV} \quad \longrightarrow \quad m_\chi \gtrsim \mathcal{O}(10) \text{ MeV}$$

See also Dror, Elor & McGehee  
[1905.12635, 1908.10861]

SFG, Oleg Titov [[arXiv:2405.05728](https://arxiv.org/abs/2405.05728), to appear in PRD]

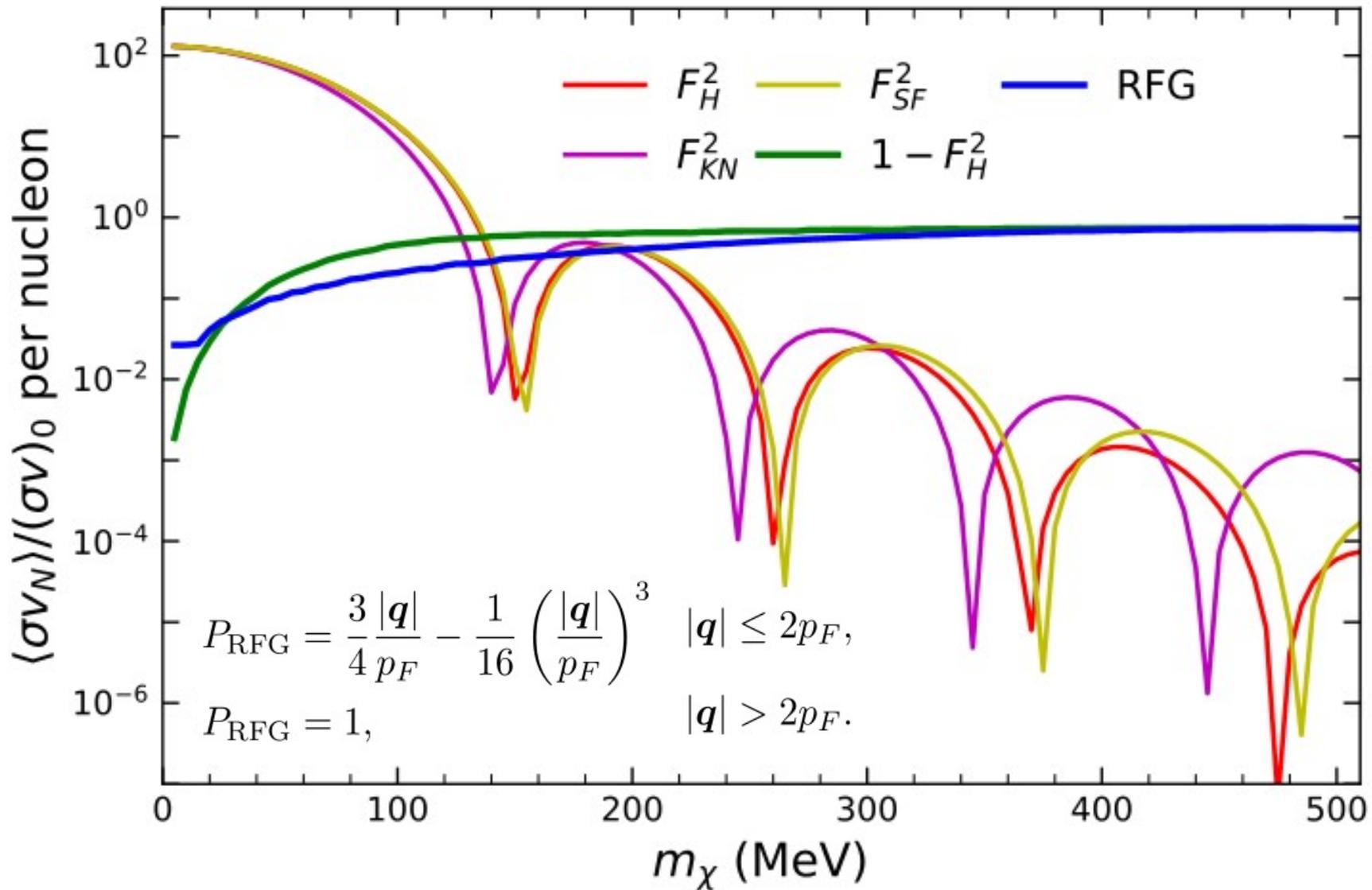
# Fermi Motion & Pauli Blocking

$$|\mathbf{q}| = |\mathbf{p}_\nu| = m_\chi - \frac{m_\chi^2}{2(m_N + m_\chi)} \approx m_\chi \quad T_{N'}^\pm \approx \frac{m_\chi \pm |\mathbf{p}_N|)^2}{2m_N}$$



SFG, Oleg Titov [arXiv:2405.05728, to appear in PRD]

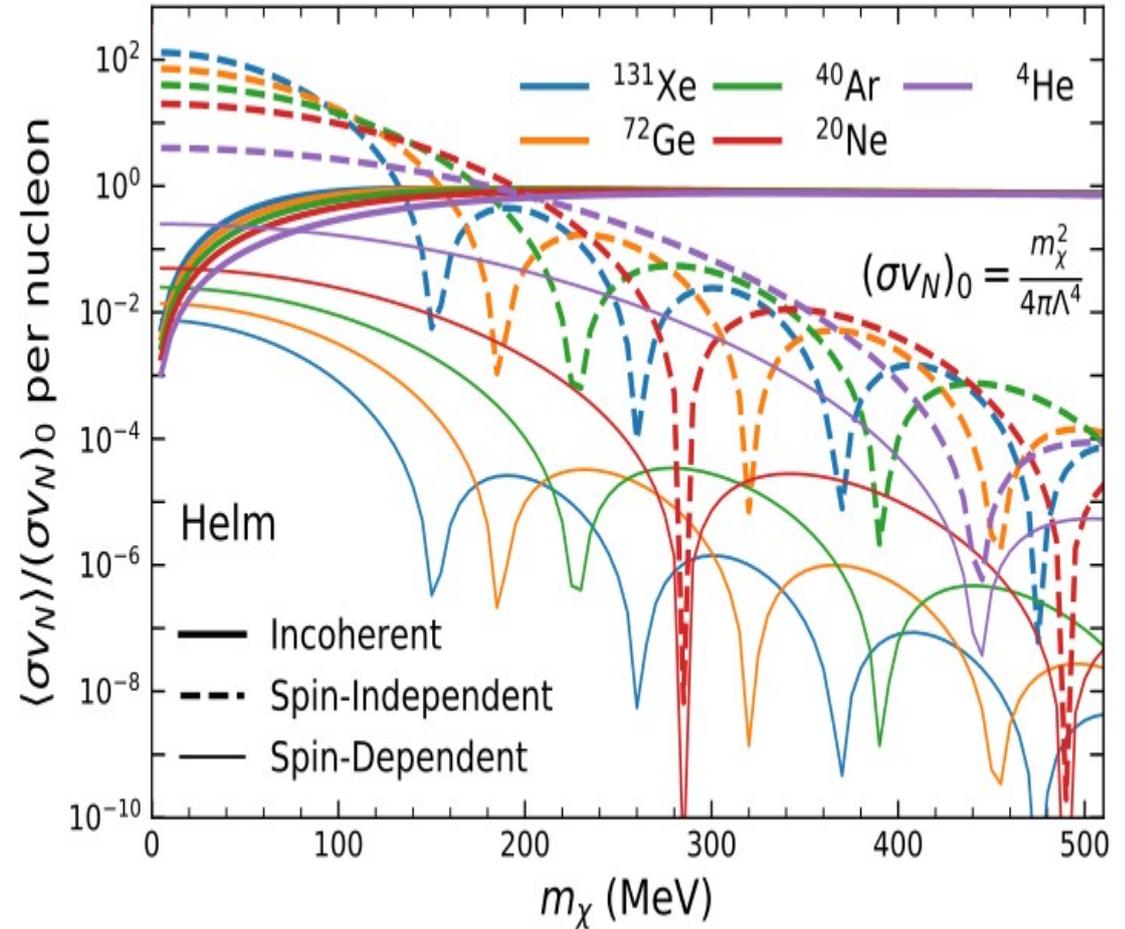
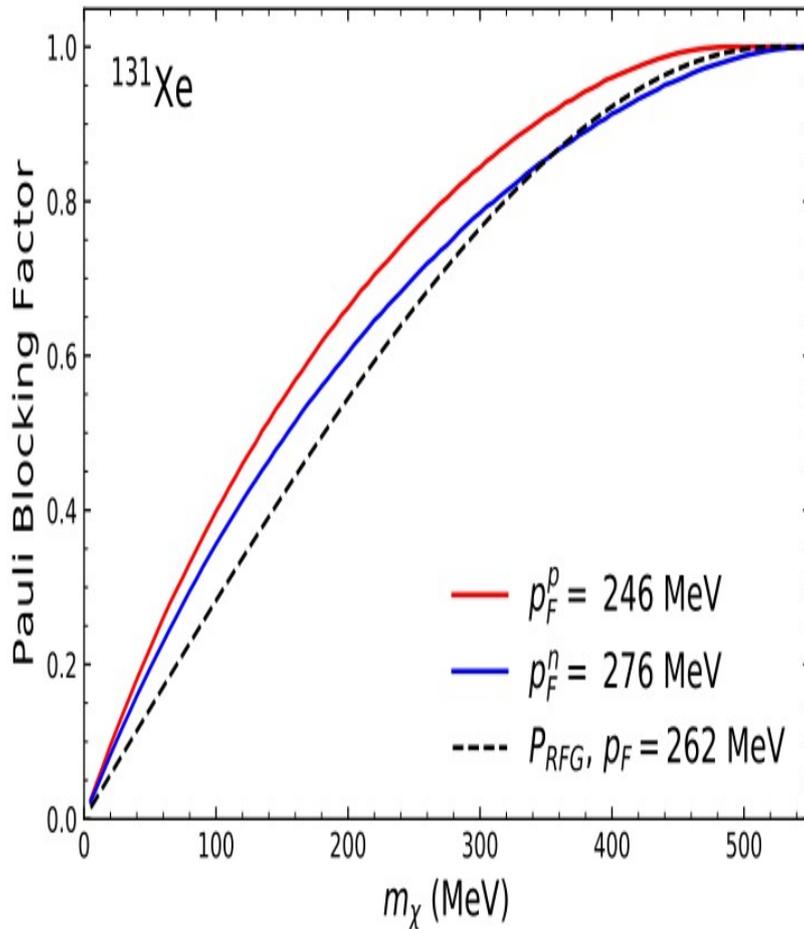
# Incoherent vs Coherent



SFG, Oleg Titov [arXiv:2405.05728, to appear in PRD]

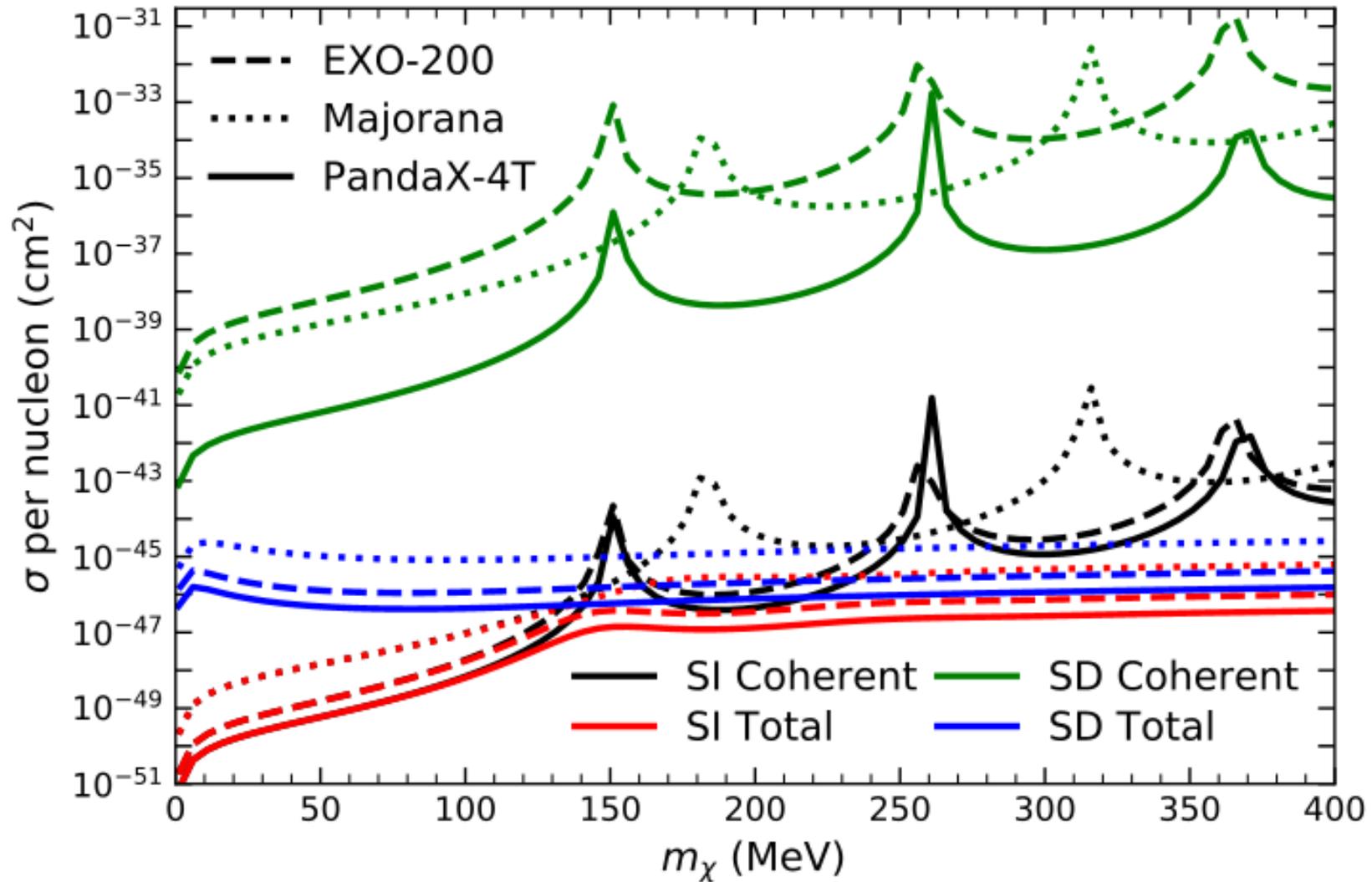
# Incoherent vs Coherent

$$|\mathbf{q}| = |\mathbf{p}_\nu| = m_\chi - \frac{m_\chi^2}{2(m_N + m_\chi)} \approx m_\chi$$



SFG, Oleg Titov [[arXiv:2405.05728](https://arxiv.org/abs/2405.05728), to appear in PRD]

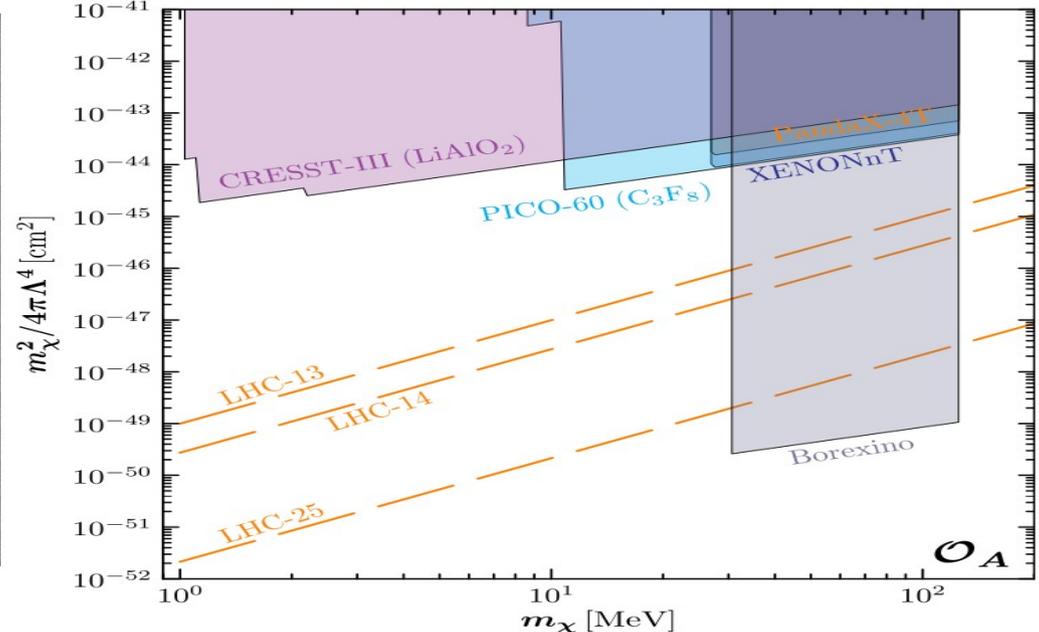
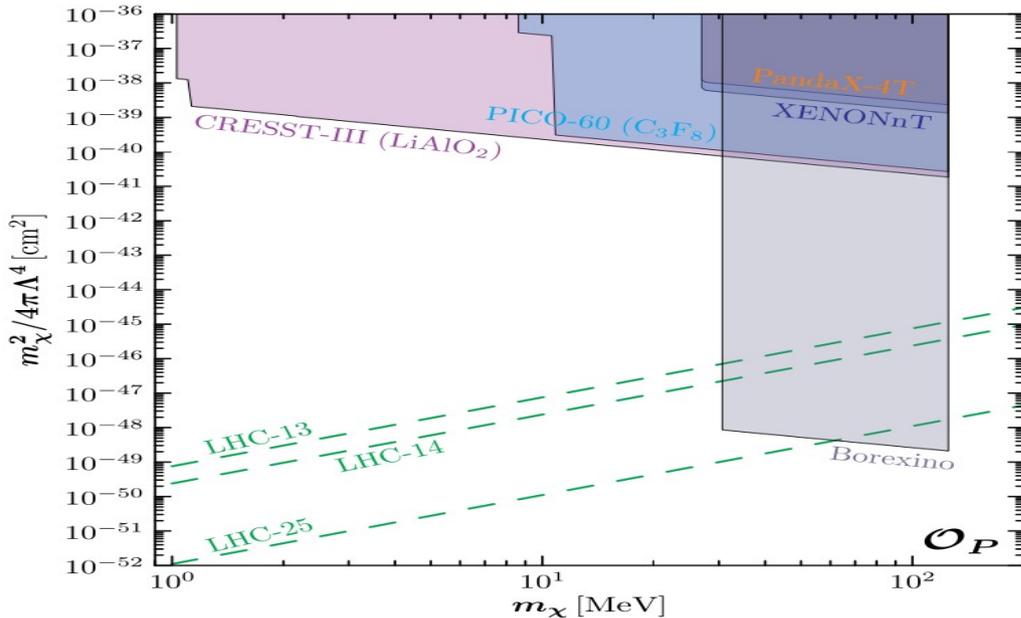
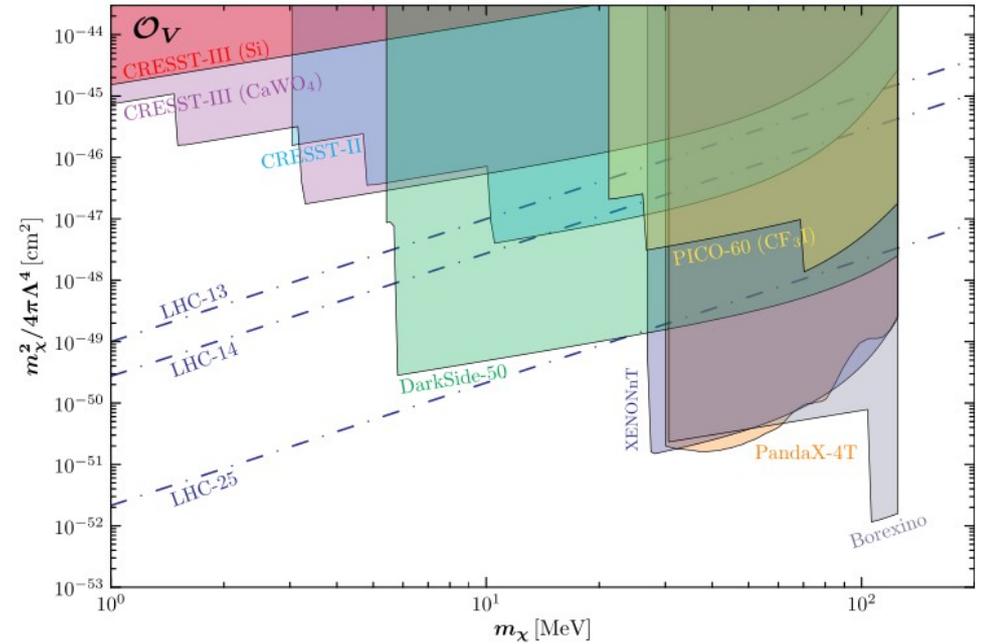
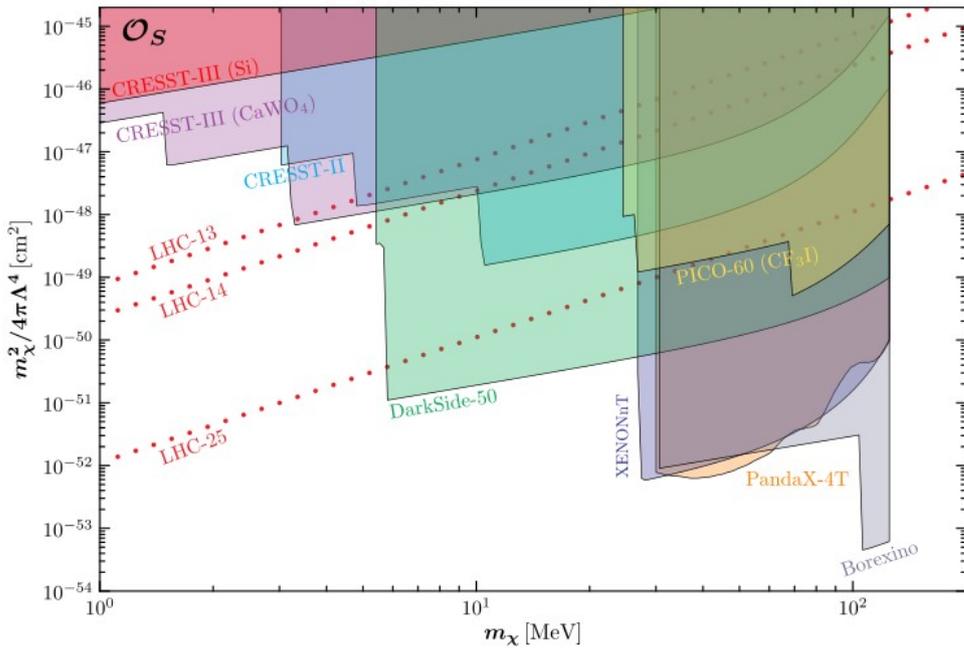
# Incoherent Domination @ large $m$



SFG, Oleg Titov [[arXiv:2405.05728](https://arxiv.org/abs/2405.05728), to appear in PRD]

# Nuclear Absorption @ Collider

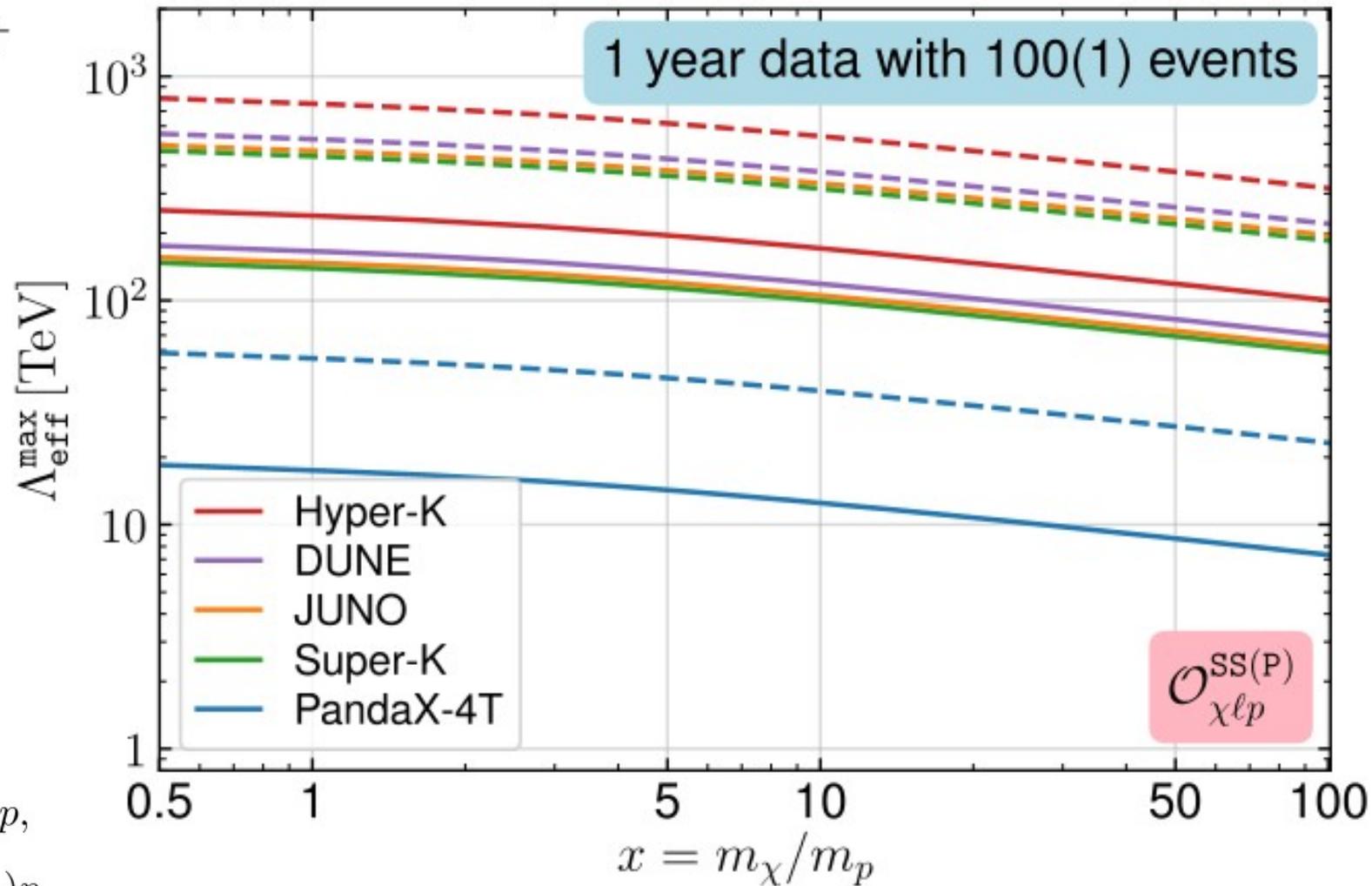
Kai Ma, SFG, Lin-Yun He, Ning Zhou [arXiv:2405.16878]



# Nucleon Consumption

$$\chi + p \rightarrow \chi + \ell^+$$

$$\chi + n \rightarrow \chi + \nu$$



$$O_{\chi lp}^{\text{SS(P)}} = \bar{\chi} \chi \bar{\ell}^c (i\gamma_5) p,$$

$$O_{\chi lp}^{\text{PS(P)}} = \bar{\chi} i\gamma_5 \chi \bar{\ell}^c (i\gamma_5) p,$$

$$O_{\chi lp}^{\text{VV(A)}} = \bar{\chi} \gamma_{\mu} \chi \bar{\ell}^c \gamma^{\mu} (\gamma_5) p,$$

$$O_{\chi lp}^{\text{AV(A)}} = \bar{\chi} \gamma_{\mu} \gamma_5 \chi \bar{\ell}^c \gamma^{\mu} (\gamma_5) p,$$

$$O_{\chi lp}^{\text{TT}(\tilde{T})} = \bar{\chi} \sigma_{\mu\nu} \chi \bar{\ell}^c \sigma^{\mu\nu} (i\gamma_5) p,$$

SFG, Xiao-Dong Ma [arXiv:2406.00445]

- 1) Thresholds in DM search
- 2) Cosmic ray boosted DM & diurnal modulation
- 3) Fermionic DM Absorption
- 4) **Reactivate Forbidden DM @ Supermassive Black Hole**

# Forbidden DM

Yu Cheng, **SFG**, Xiao-Gang He, Jie Sheng [[Phys. Lett. B 847 \(2023\) 138294](#), [arXiv:2211.05643](#)]

Yu Cheng, **SFG**, Jie Sheng, Tsutomu Yanagida [[PhysRevD.107.123013](#); [arXiv:2309.12043](#)]



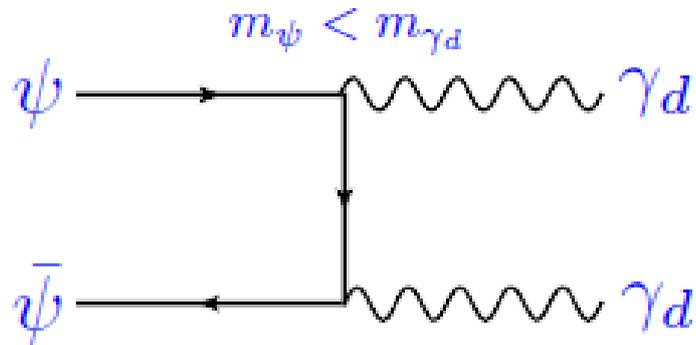
Yu Cheng



Jie Sheng

# Forbidden DM

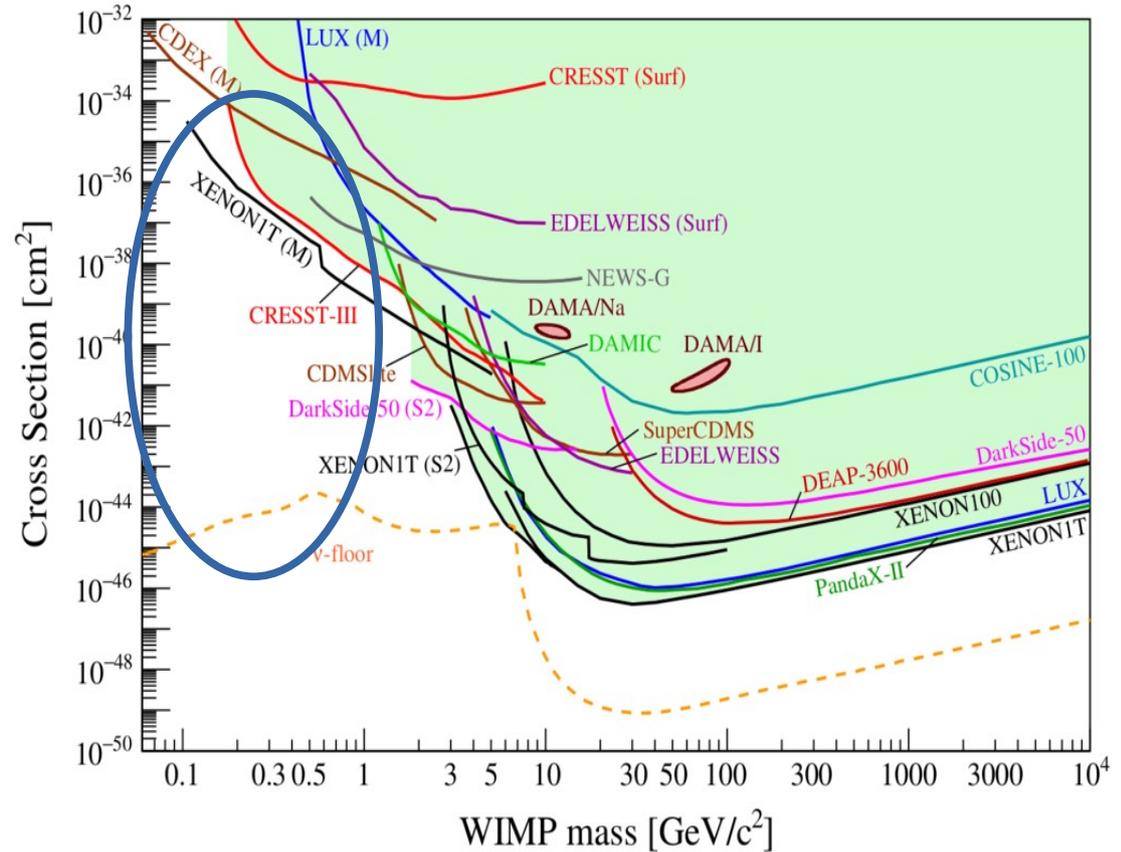
1) forbidden annihilations:



$$\langle \sigma_{\gamma_d \gamma_d} v \rangle \sim \alpha_d^2 / m_{\gamma_d}^2$$



$$\langle \sigma_{\chi \bar{\chi}} v \rangle = \frac{(n_{\gamma_d}^{eq})^2}{(n_\psi^{eq})^2} \langle \sigma_{\gamma_d \gamma_d} v \rangle \approx 8\pi f_\Delta \frac{\alpha_d^2}{m_\psi^2} e^{-2\Delta x}$$



D'Agnolo & Ruderman, Phys.Rev.Lett. 115 (2015) 6, 061301 [1505.07107]

## Three exceptions in the calculation of relic abundances

Kim Griest

*Center for Particle Astrophysics and Astronomy Department, University of California, Berkeley, California 94720*

David Seckel

*Bartol Research Institute, University of Delaware, Newark, Delaware 19716*

(Received 15 November 1990)

The calculation of relic abundances of elementary particles by following their annihilation and freeze-out in the early Universe has become an important and standard tool in discussing particle dark-matter candidates. We find three situations, all occurring in the literature, in which the standard methods of calculating relic abundances fail. The **(first)** situation occurs when another particle lies near in mass to the relic particle and shares a quantum number with it. An example is a light squark with neutralino dark matter. The additional particle must be included in the reaction network, since its annihilation can control the relic abundance. The **(second)** situation occurs when the relic particle lies near a mass threshold. Previously, annihilation into particles heavier than the relic particle was considered kinematically forbidden, but we show that if the mass difference is  $\sim 5-15\%$ , these “forbidden” channels can dominate the cross section and determine the relic abundance. The **(third)** situation occurs when the annihilation takes place near a pole in the cross section. Proper treatment of the thermal averaging and the annihilation after freeze-out shows that the dip in relic abundance caused by a pole is not nearly as sharp or deep as previously thought.

Coannihilation

Forbidden DM

Breit-Wigner

# How to test Forbidden DM?

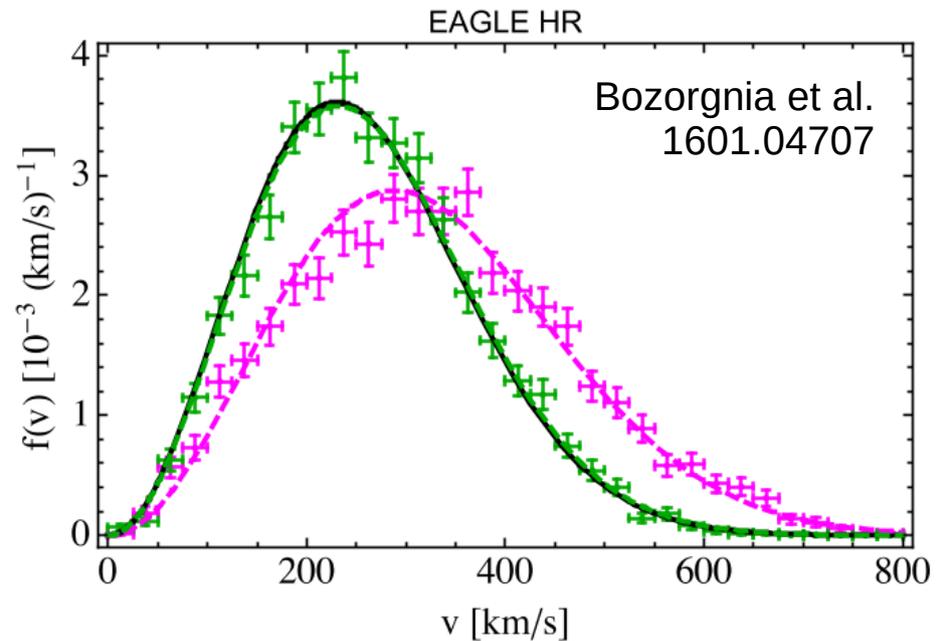
- Boltzmann Distribution

$$f_0(v) = \frac{\rho_\chi}{m_\chi} \frac{4}{\sqrt{\pi}} \frac{v^2}{(\sqrt{2/3}v_d)^3} e^{-\frac{v^2}{(\sqrt{2/3}v_d)^2}}$$

$$\Delta_{F\chi} \equiv \frac{m_F - m_\chi}{m_\chi} \sim 1\%$$



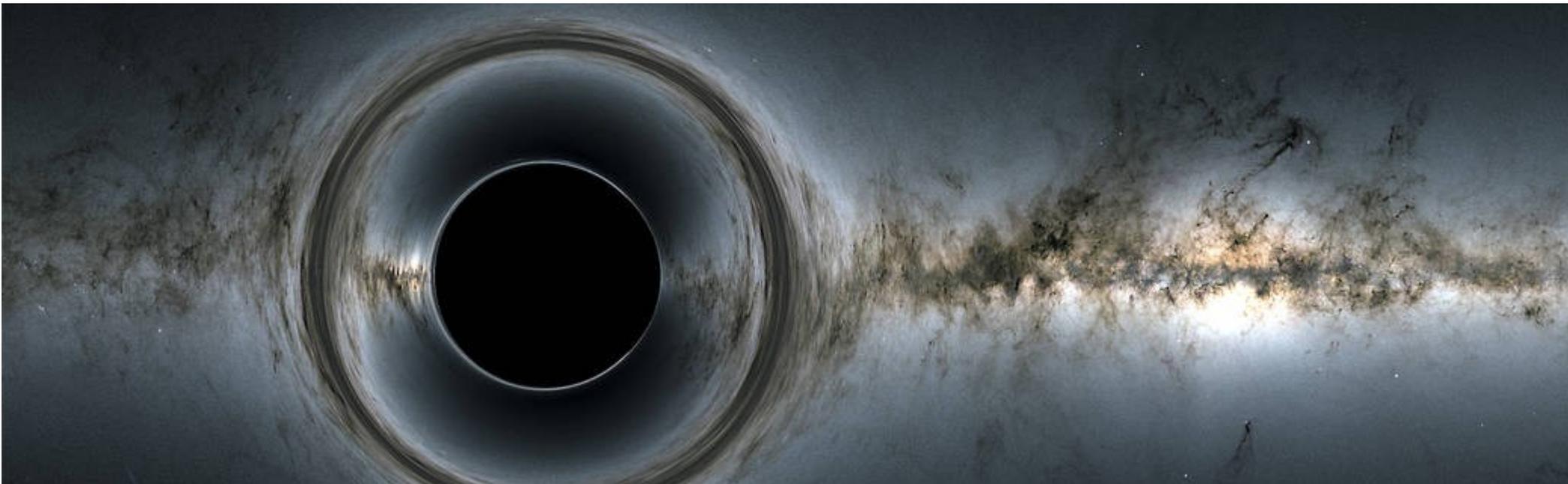
$$v_d \sim 10\%$$



- Juttner Distribution

$$f_J(\mathbf{p}) = \frac{1}{4\pi T m_\chi^2 K_2(x)} e^{-\frac{\sqrt{|\mathbf{p}|^2 + m_\chi^2}}{T}} \quad x \equiv m_\chi/T$$

# Velocity Scaling @ Black Holes



$$v^2 \sim \frac{GM}{r} \quad \Rightarrow \quad v(r) \sim \frac{1}{\sqrt{r}}$$
$$\rho \propto r^?$$

## ➤ NFW Profile

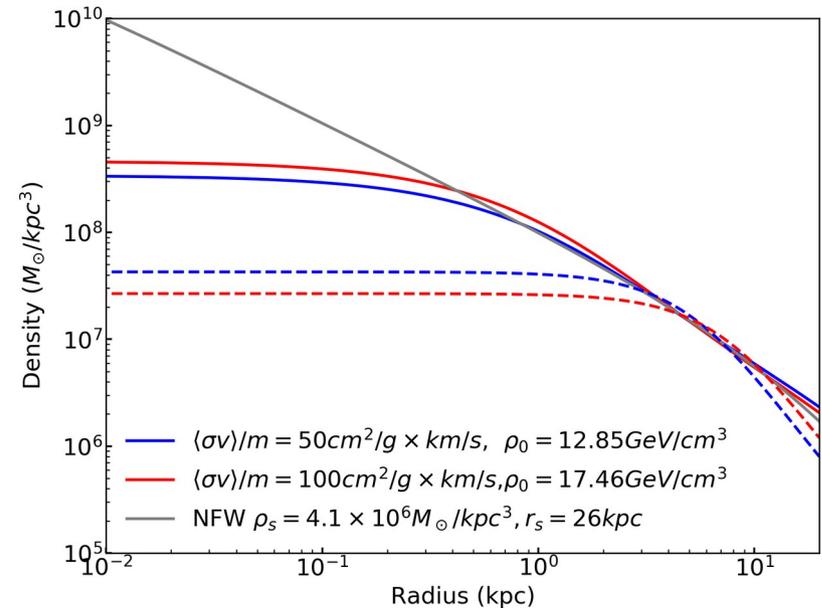
$$\rho(r) = \begin{cases} \rho_{\text{iso}}(r), & r < r_1 \\ \rho_{\text{NFW}}(r), & r > r_1 \end{cases}$$

$$\rho_{\text{NFW}}(r) = \frac{\rho_s}{(r/r_s)(1+r/r_s)^2}$$

## ➤ Jeans equation

$$\frac{d}{dr} \left( r^2 \frac{d \ln \rho}{dr} \right) = - \frac{4\pi G r^2 (\rho + \rho_b)}{\sigma_0^2}$$

1. Isothermal gas: constant velocity dispersion  $\sigma_0$
2. Mass contained inside core the same
3. Continuity



Kaplinghat, Keeley, Linden & H. B. Yu, Phys. Rev. Lett. 113, 021302 (2014) [arXiv:1311.6524]  
Kaplinghat, Tulin & H. B. Yu, Phys. Rev. Lett. 116, no.4, 041302 (2016) [arXiv:1508.03339]

# Spike: Conductive Fluid

➤ Inner region more influenced by BH potential

$$\rho(r) = \begin{cases} \rho_{\text{NFW}}(r), & r > r_1, \\ \rho_{\text{iso}}(r), & r_0 < r < r_1 \\ \rho_{\text{spike}}(r), & r < r_0. \end{cases}$$

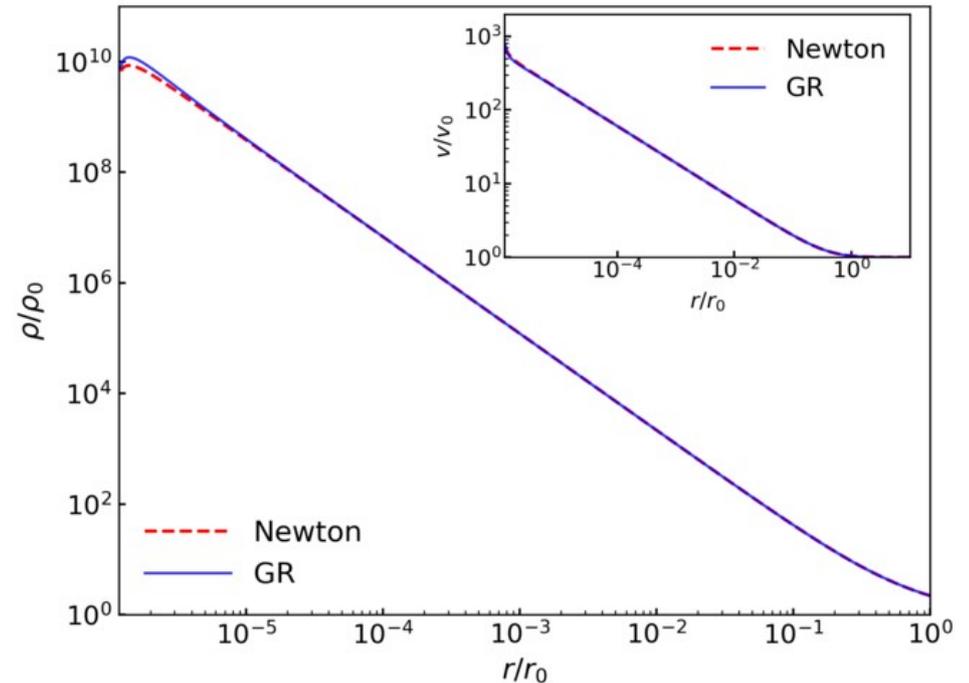
$$r_0 \equiv GM/v_0^2$$

➤ Conductive fluid

$$\frac{dv}{dr} = \frac{D}{v^{2-a} \rho^2 r^4}$$

$$\frac{d\rho}{dr} = -\frac{\rho}{v^2 r^2} - \frac{2D}{v^{3-a} \rho r^4}$$

$$\rho = 0, \quad r = r_{\text{in}} \equiv 4GM$$



$$v_d \propto r^{-1/2}$$

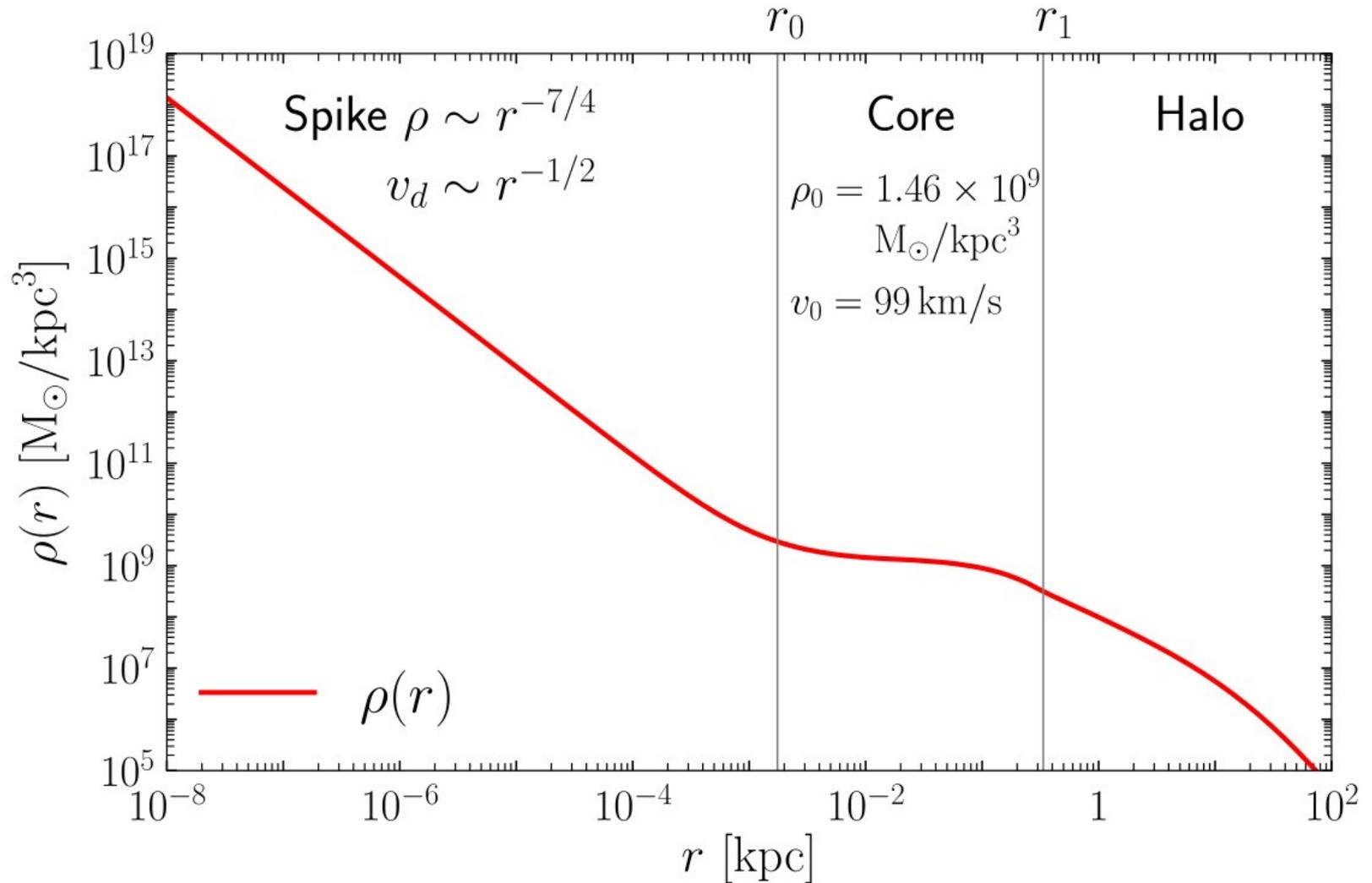
$$\rho \propto r^{-(3+a)/4}$$

$$\sigma = \sigma_0 (v_d/v_0)^a$$

Shapiro & Paschalidis, Phys. Rev. D 89, no.2, 023506 (2014) [arXiv:1402.0005]

# DM Profiles

$$\rho(r) = \begin{cases} \rho_{\text{NFW}}(r), & r > r_1, \\ \rho_{\text{iso}}(r), & r_0 < r < r_1 \\ \rho_{\text{spike}}(r), & r < r_0. \end{cases}$$



$$\mathcal{L}_{\text{DM}} = g_\chi \bar{\chi} \gamma^\mu \chi \phi_\mu + g_{F\chi} \bar{F} \gamma^\mu F \phi_\mu$$

## ➤ Forbidden channel

$$\chi \bar{\chi} \rightarrow F \bar{F} \quad m_\chi \lesssim m_F$$

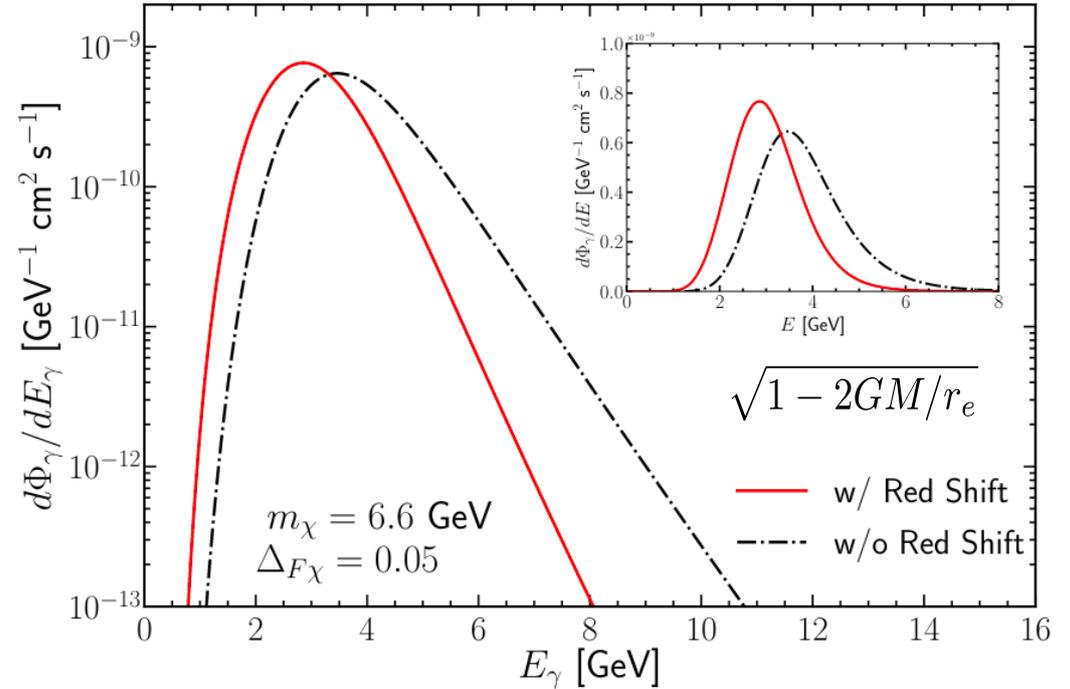
$$\frac{d\sigma}{d\Omega} \approx \sqrt{\frac{s - 4m_F^2}{s - 4m_\chi^2}} \frac{4m_F^2 + 4m_\chi^2 + s}{64\pi^2 (s - m_\phi^2)^2}$$

Isotropic in the C.O.M. frame

## ➤ Signal channel

$$F \rightarrow \nu \gamma$$

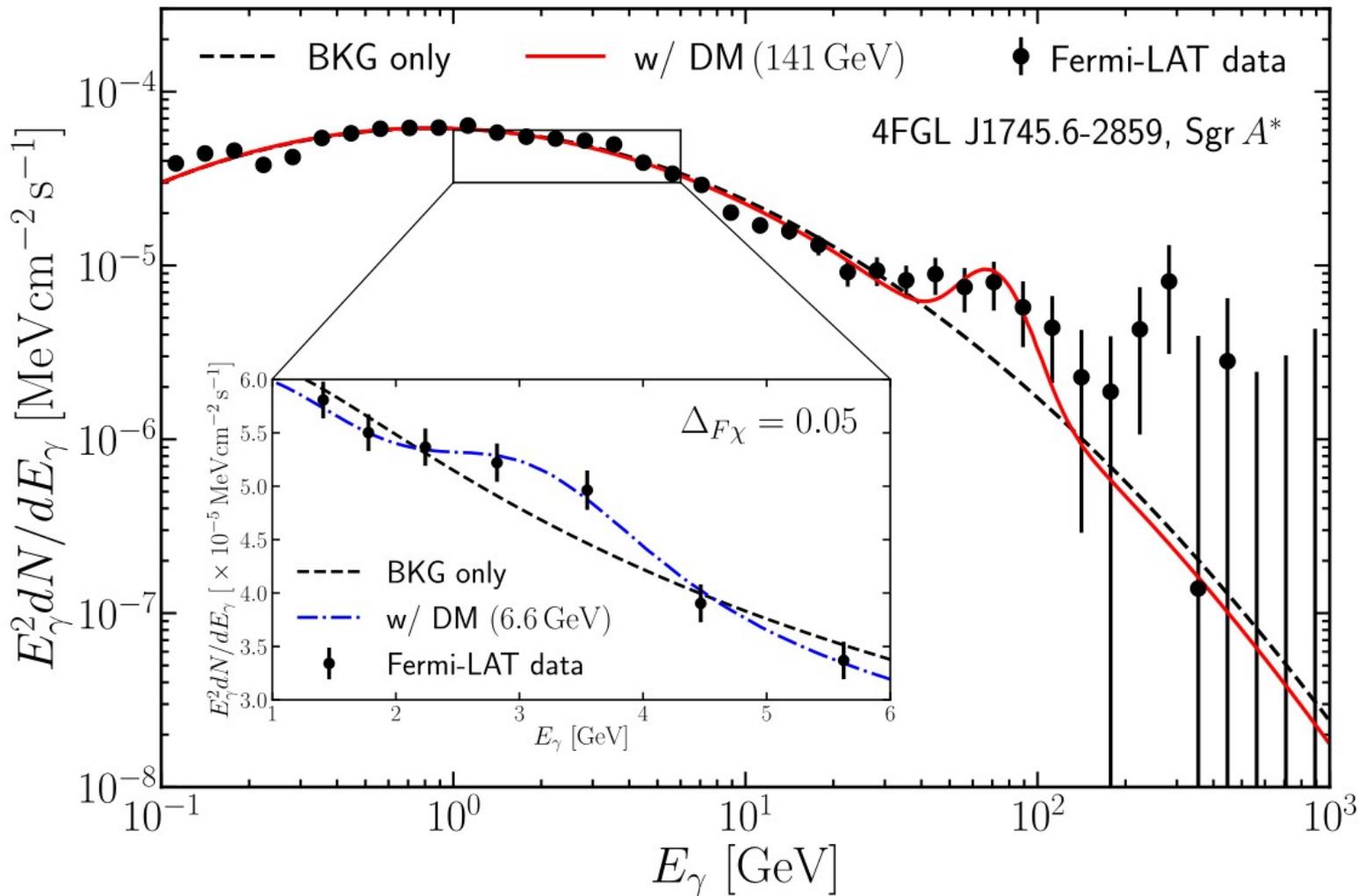
Box-shaped spectrum in C.O.M.



$$\frac{dF_\gamma}{dE_\gamma}(r) = \int_0^1 dV_r dV_c \mathcal{P}_r(V_r, V_c) \sigma_{V_r} \frac{dN_\gamma}{dE_\gamma}(V_r, V_c)$$

$$\mathcal{P}_r(V_r, V_c) \equiv \frac{x^2}{K_2^2(x)} \frac{\gamma_r^3 (\gamma_r^2 - 1) V_c^2}{(1 - V_c^2)^2} e^{-x \sqrt{(2+2\gamma_r)/(1-V_c^2)}}$$

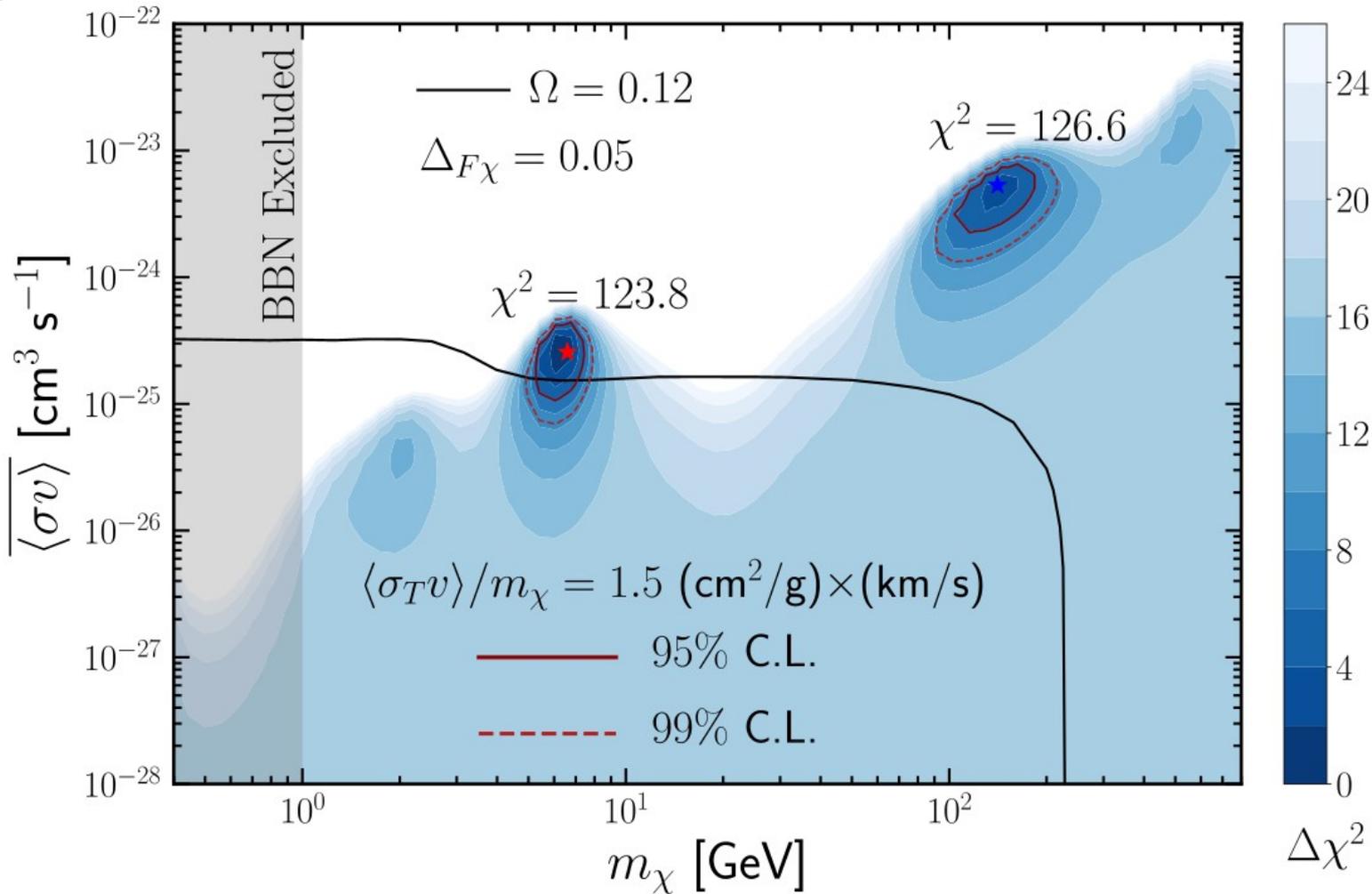
# Fitting the Fermi-LAT data



Aug/4, 2008  
~  
Oct/26, 2022

**Background Model:**  $\frac{dN}{dE} = N_0 \left( \frac{E}{E_0} \right)^{-\alpha - \beta \log(E/E_0)}$

# Sensitivity Plots



● **Bkg-only hypothesis**

$$\chi_{\text{BKG}}^2 = 140.8$$

● **1<sup>st</sup> Peak @ 6.6GeV**

● **2<sup>nd</sup> Peak @ 141GeV**

$$\langle\sigma v\rangle = 1.99 \times 10^{-22} \text{ cm}^3 \text{ s}^{-1}$$

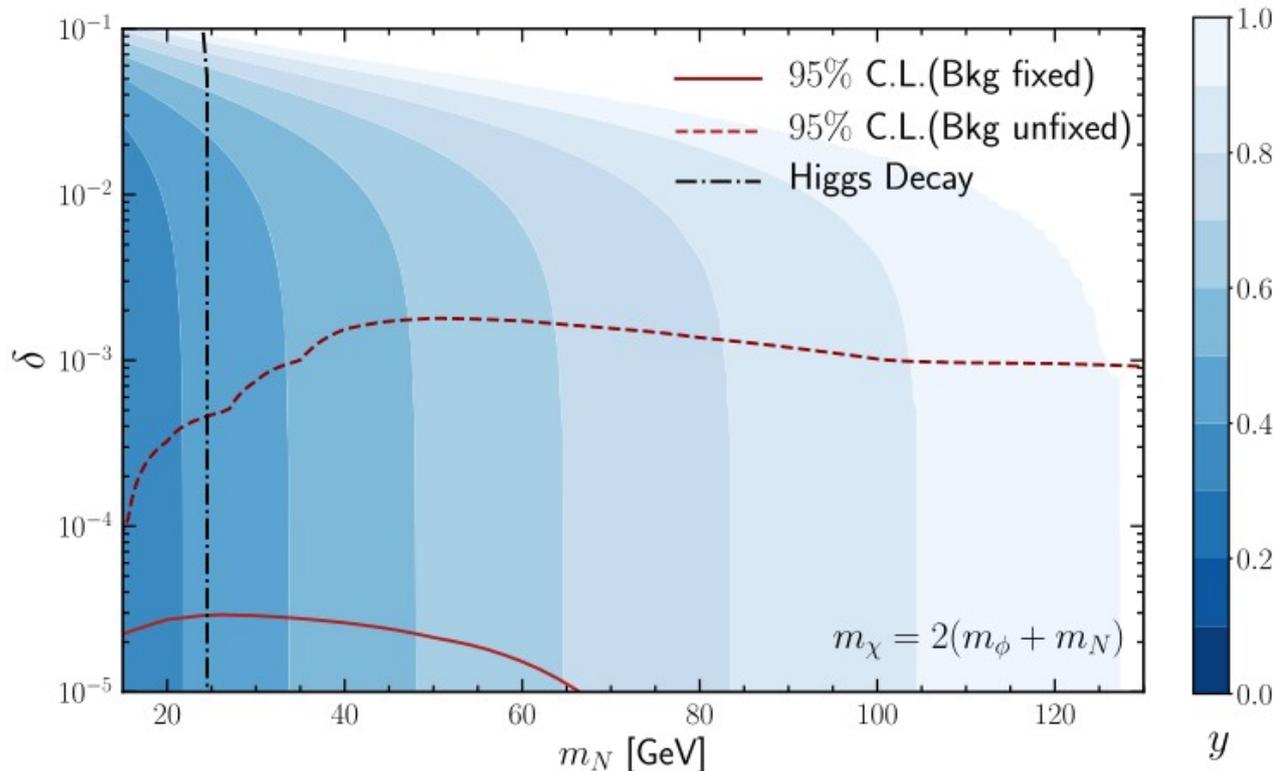
$$\overline{\langle\sigma v\rangle} = 4.12 \times 10^{-21} \text{ cm}^3 \text{ s}^{-1}$$

# Forbidden with Right-Handed $\nu$

$$\mathcal{L}_{\text{int}} = (y\phi N\chi + h.c.) + \lambda m_\phi \phi H^\dagger H$$

DM freezes out through the forbidden channel  $N + N \leftrightarrow \phi + \phi$

$$\langle \sigma_{NN\nu} \rangle \equiv \left( \frac{n_\phi^{\text{eq}}}{n_N^{\text{eq}}} \right)^2 \langle \sigma_{\phi\phi\nu} \rangle = \frac{(1 + \delta)^3}{4} e^{-2x\delta} \langle \sigma_{\phi\phi\nu} \rangle$$

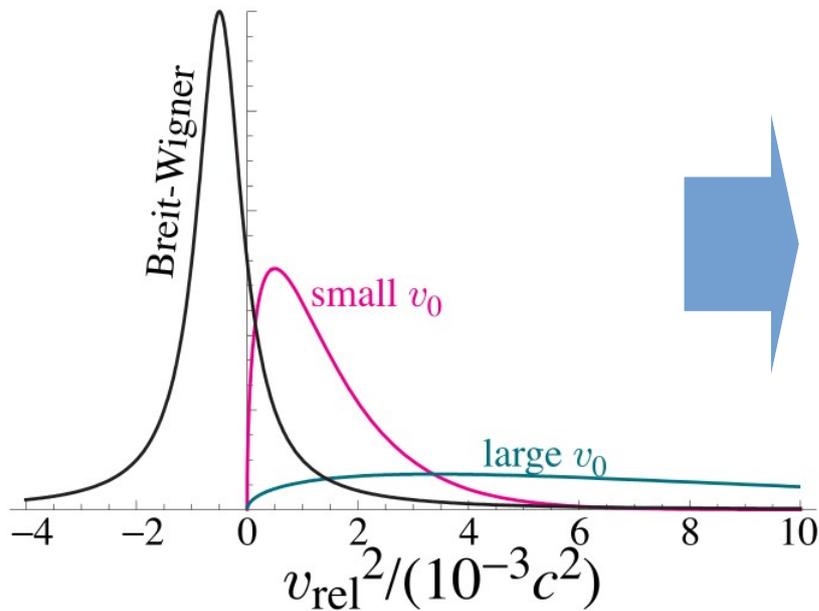


$$\delta \equiv \frac{m_\phi - m_N}{m_N}$$

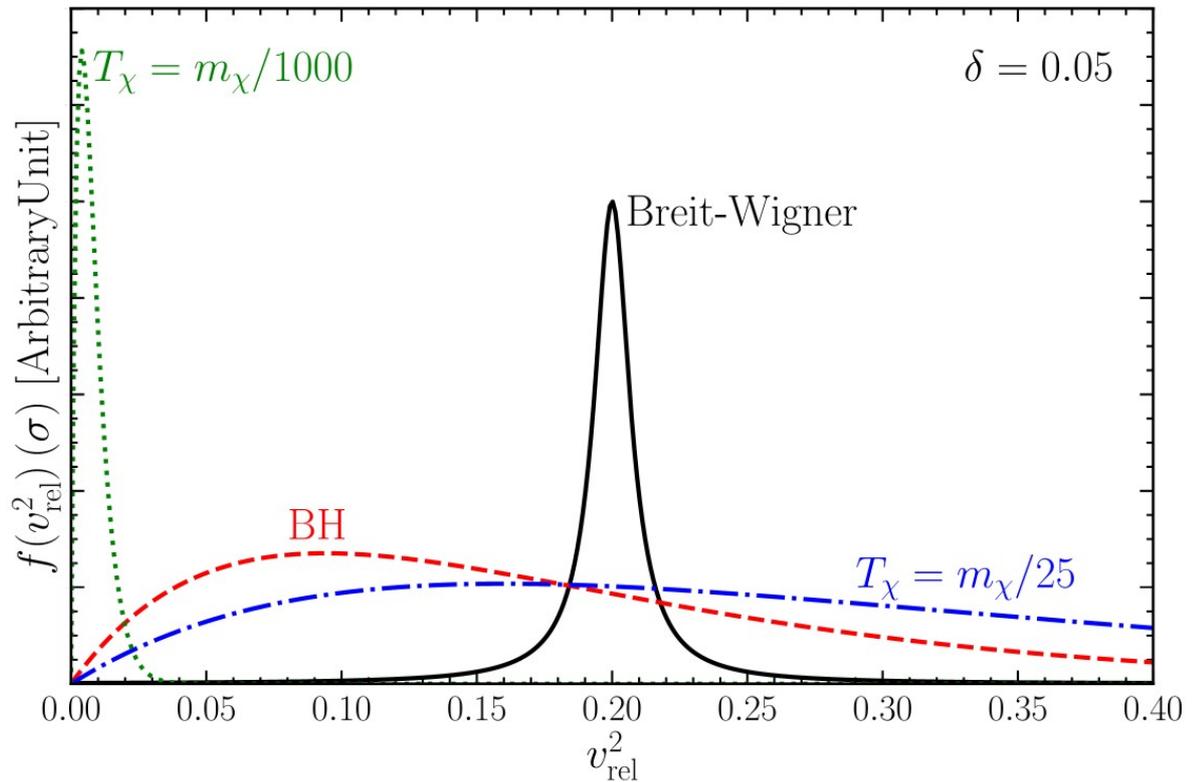
Yu Cheng, **SFG**, Jie Sheng, Tsutomu Yanagida  
[PhysRevD.107.123013]

# BW Resonance with Heavy Mediator

$$\sigma = \frac{16\pi\beta_f}{s\bar{\beta}_i\bar{\beta}_f\beta_i} \frac{\gamma^2}{\left(\frac{\vec{v}_{\text{rel}}^2/4 - \delta}{1+\delta}\right)^2 + \gamma^2} B_i B_f \quad \mathcal{L}_{\text{int}} = (y\phi N_\chi^T \epsilon N_\chi + h.c.) + \lambda m_\phi \phi H^\dagger H$$



Ibe, Murayama, Yanagida [Phys.Rev.D 79 (2009) 095009]



Yu Cheng, **SFG**, Jie Sheng, Tsutomu Yanagida [[PhysRevD.107.123013](#); [arXiv:2309.12043](#)]

- 1) Thresholds in DM search
- 2) Cosmic ray boosted DM & diurnal modulation
- 3) Fermionic DM Absorption
- 4) **Reactivate Forbidden DM @ Supermassive Black Hole**
- 5) **Energetic MeV Solar DM from pp-chain**

# MeV Solar DM

SFG, Jie Sheng, Chen Xia, Chuan-Yang Xing [[arXiv:2408.12448](https://arxiv.org/abs/2408.12448)]



Jie Sheng

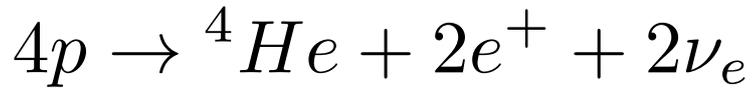


Chen Xia



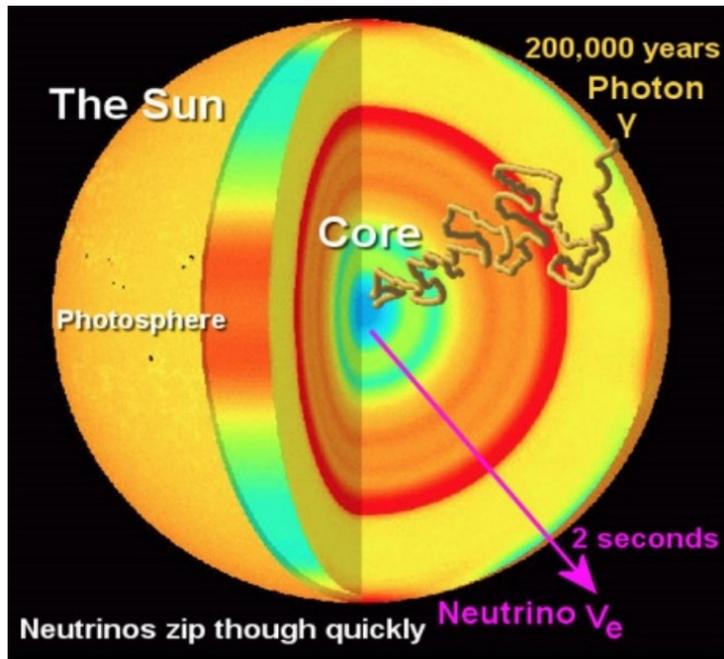
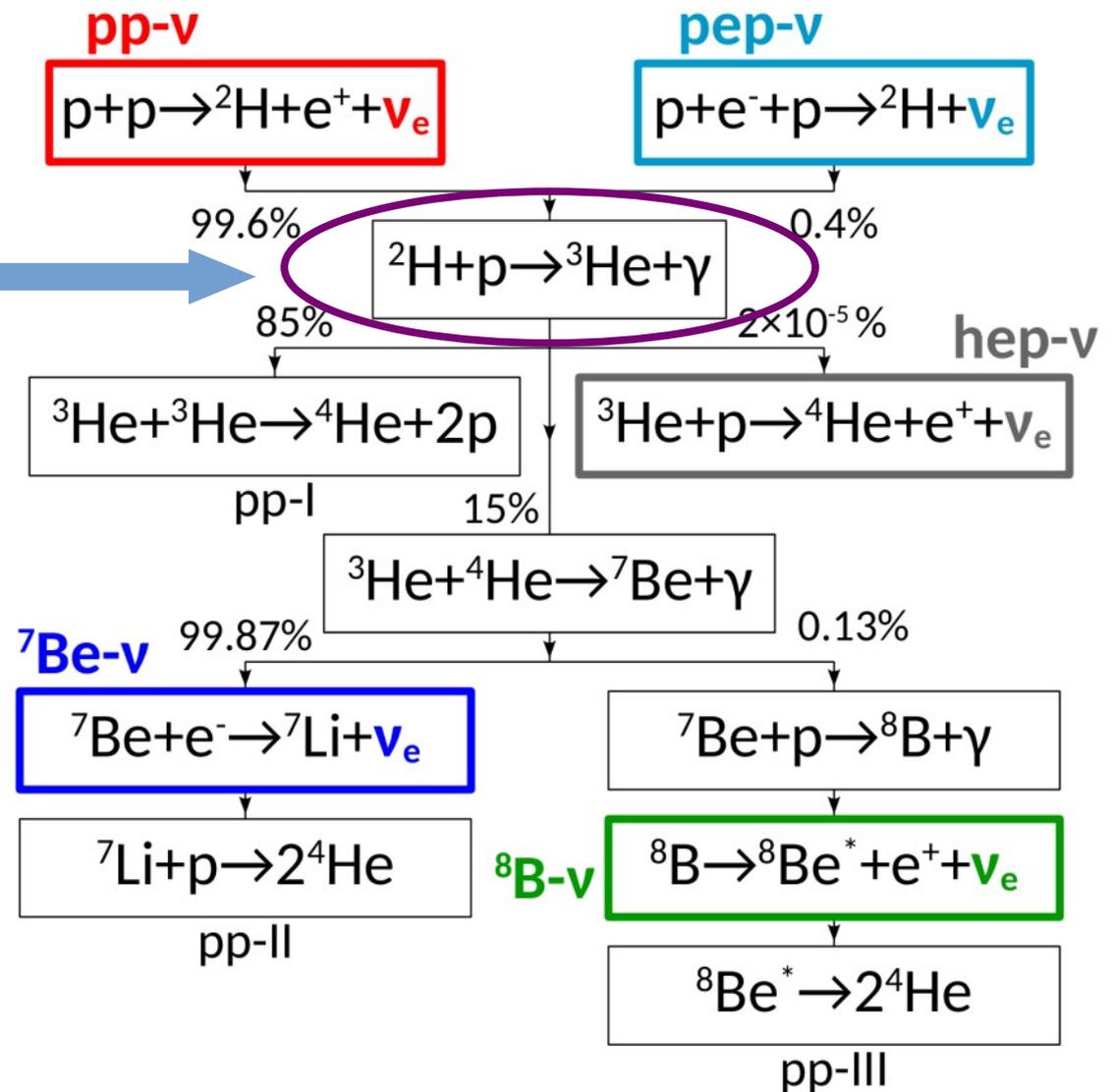
Chuan-Yang Xing

# MeV Energy in Sun

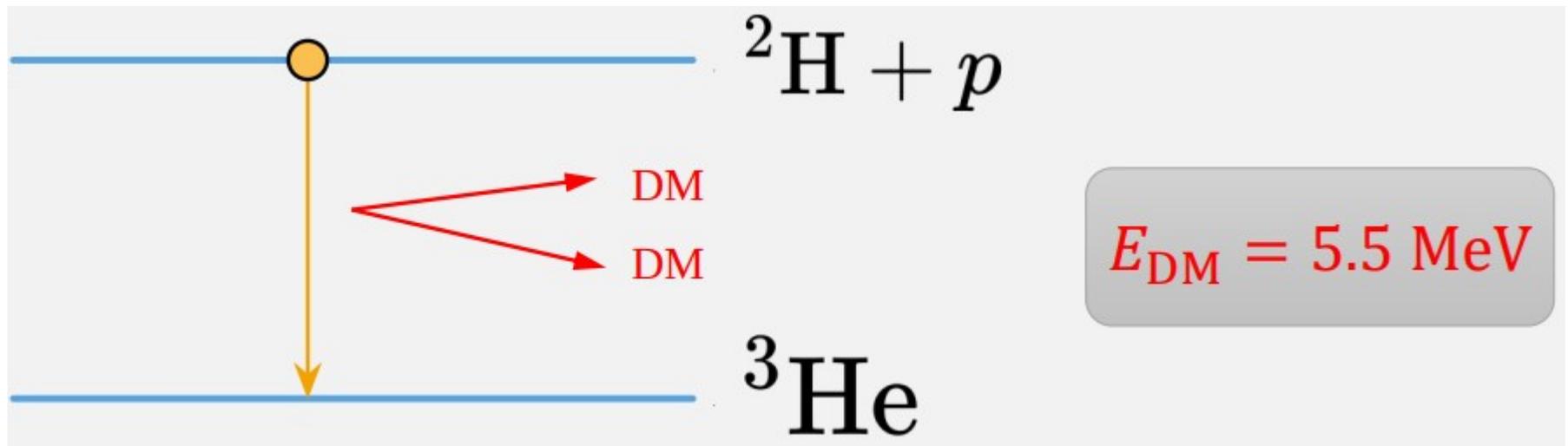
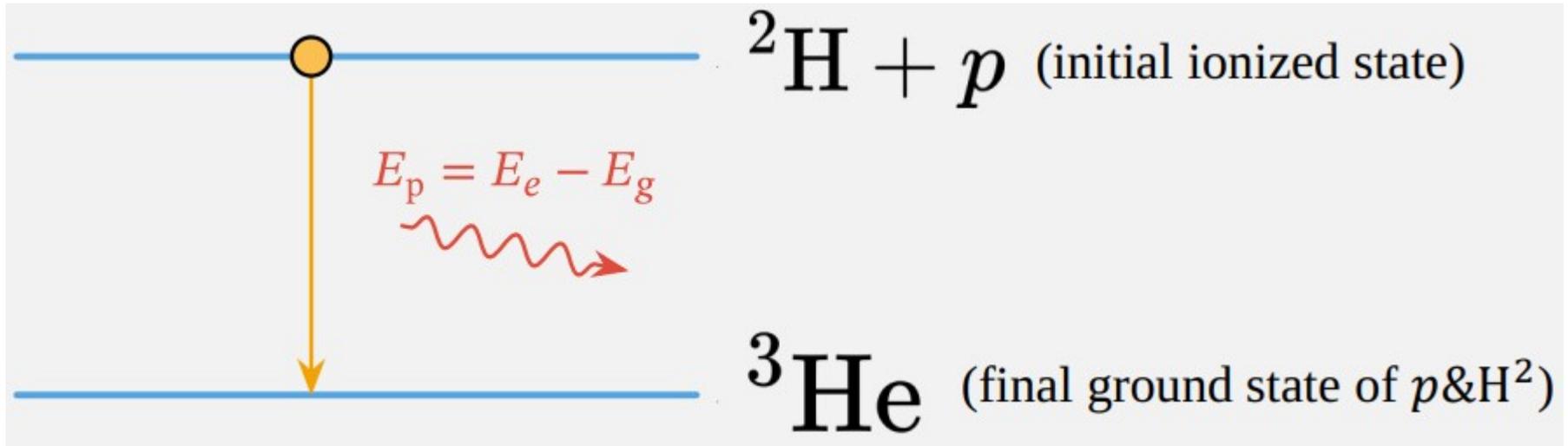
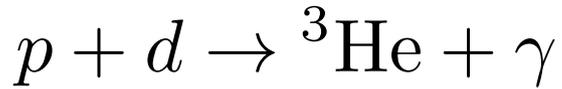


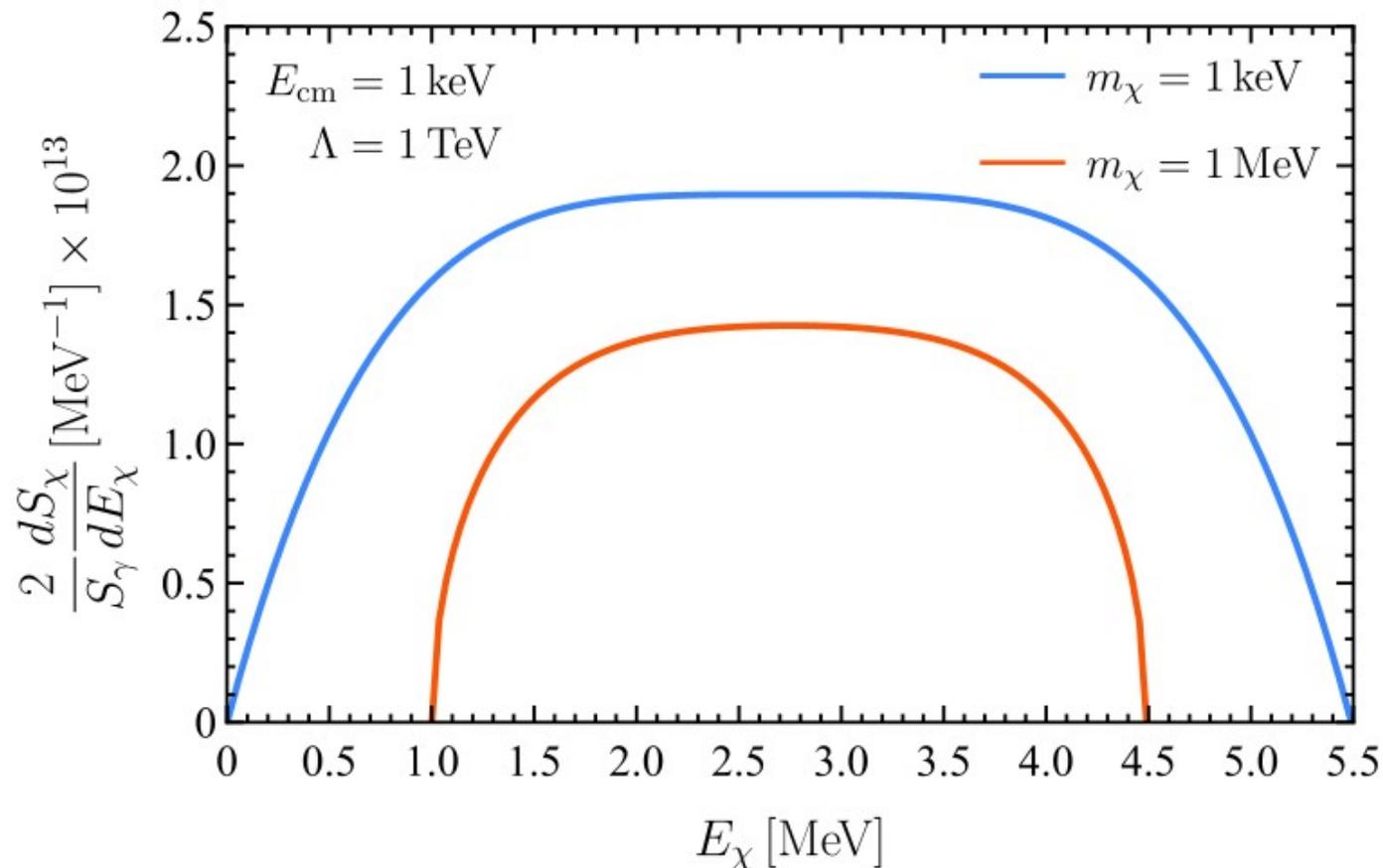
## pp chain

5.5MeV



# MeV DM Production in Sun



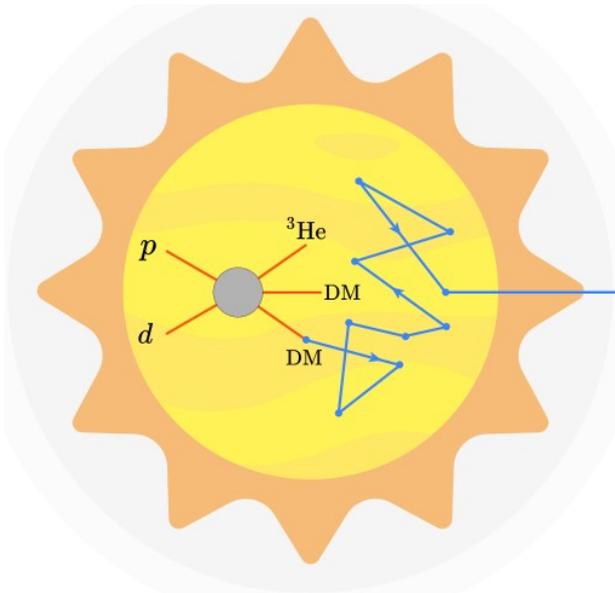


$$V_0 + V_S(\mathbf{L} \cdot \mathbf{S}) + V_C$$

$$\frac{d^3 N_\chi}{dt dE_\chi dV_\odot} \simeq \frac{2}{S_\gamma} \frac{dS_\chi}{dE_\chi} \frac{d^2 N_\gamma}{dt dV_\odot}$$

$$S(E_{\text{cm}}) \equiv \sigma(E_{\text{cm}}) E_{\text{cm}} e^{2\pi\eta}$$

# Analytical Attenuation in Sun

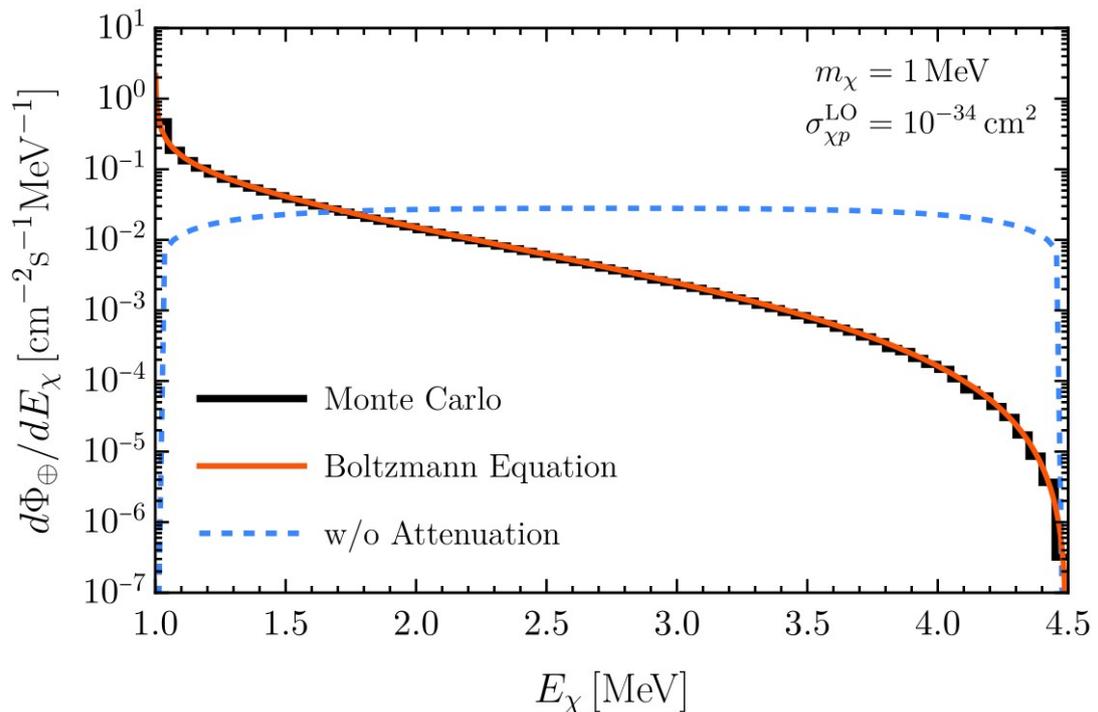
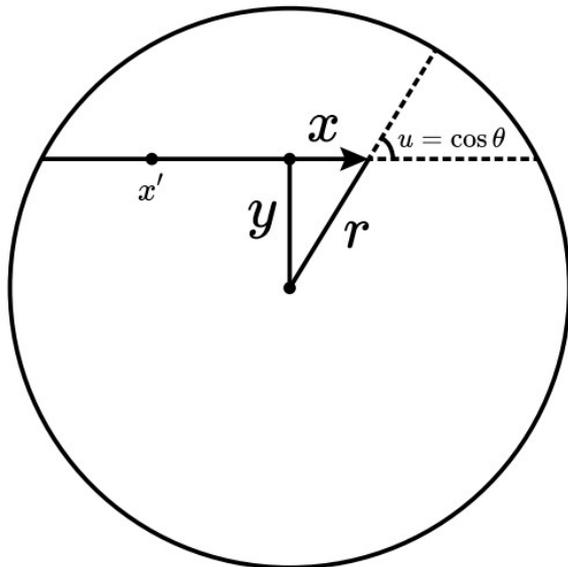


$$\hat{\mathbf{L}}[f_\chi] = |\mathbf{p}_\chi| \left( u \frac{\partial f_\chi}{\partial r} + \frac{1 - u^2}{r} \frac{\partial f_\chi}{\partial u} \right)$$

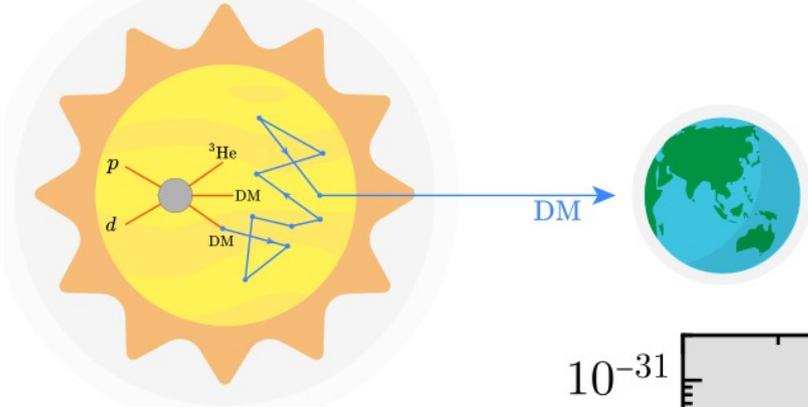


$$x \equiv ru \quad y \equiv r\sqrt{1 - u^2}$$

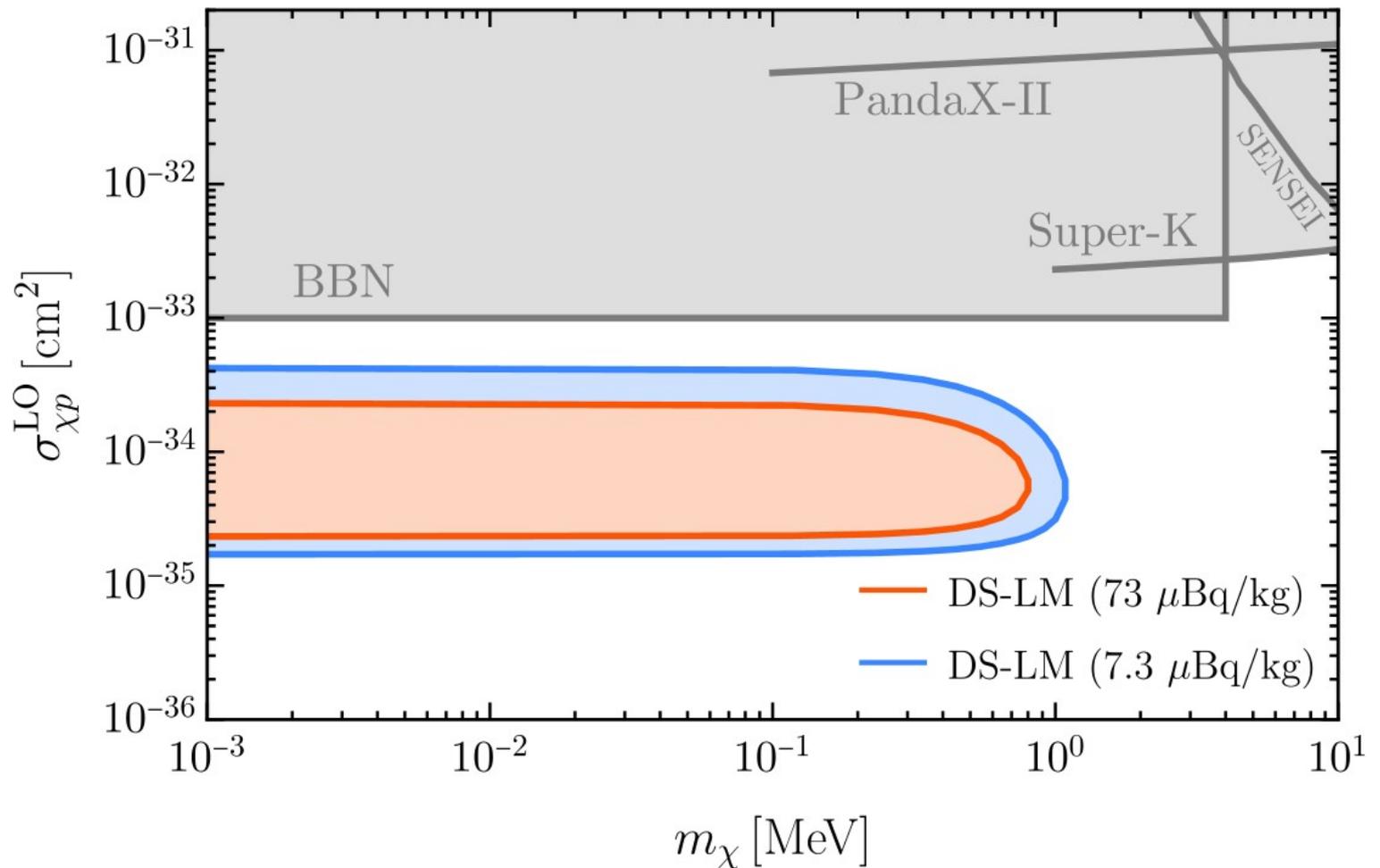
$$\hat{\mathbf{L}}[f_\chi] = |\mathbf{p}_\chi| \partial_x f_\chi$$



# Detection at DarkSide-Low



Production Attenuation



Threshold exists in not just direct detection but also indirect detection!

0) Cosmic Ray Boosted DM & Diurnal Modulation

1) **Fermionic DM Absorption & Nucleon Consumption**

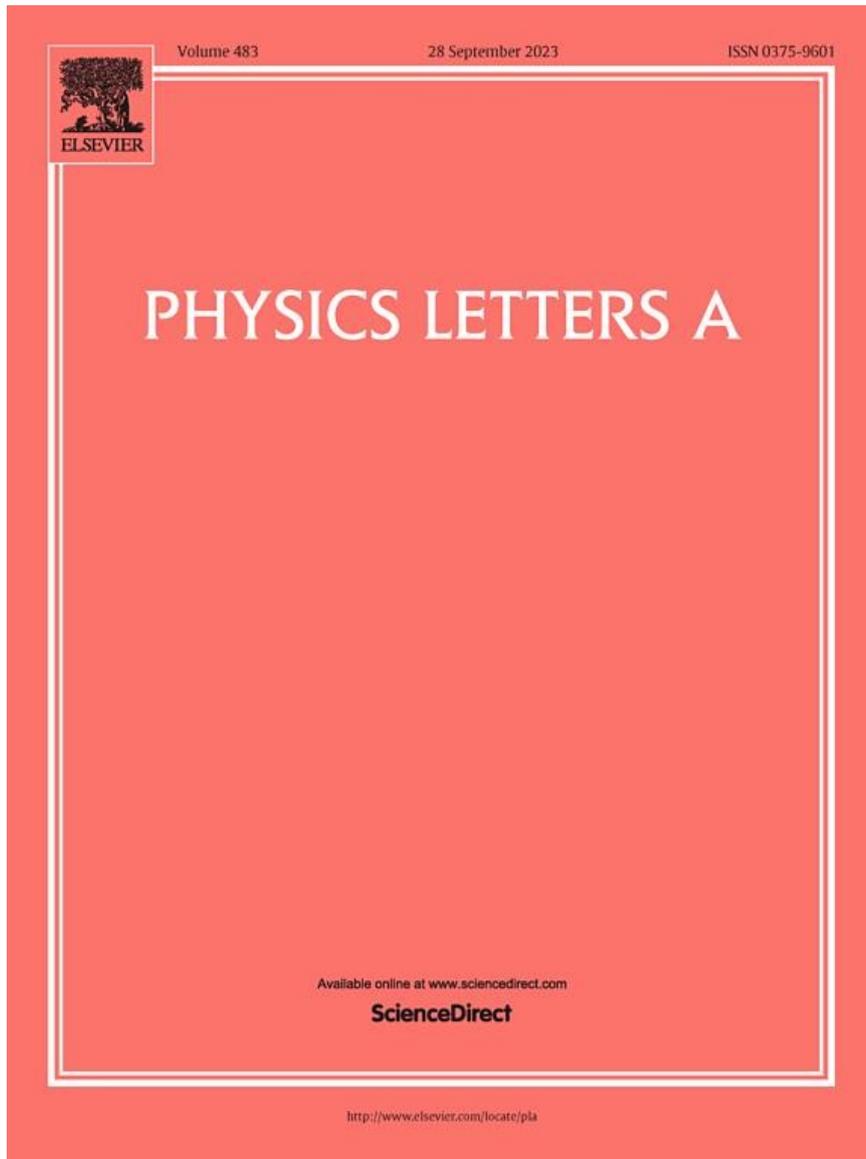
- Efficient energy release
- Peak shape by atomic effects

2) **Forbidden DM @ SMBH**

- How to uniquely test forbidden DM?
- Point source around supermassive BH

3) **MeV Solar DM**

- 5.5MeV energy release from pp-chain
- Analytic attenuation



## Aims & Scope

- Nonlinear science,
- Statistical physics,
- Mathematical and computational physics,
- AMO and physics of complex systems,
- Plasma and fluid physics,
- Optical physics,
- General and cross-disciplinary physics,
- Biological physics and nanoscience,
- Astrophysics, Particle physics and Cosmology.



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Tsung-Dao Lee Institute

# Thank You

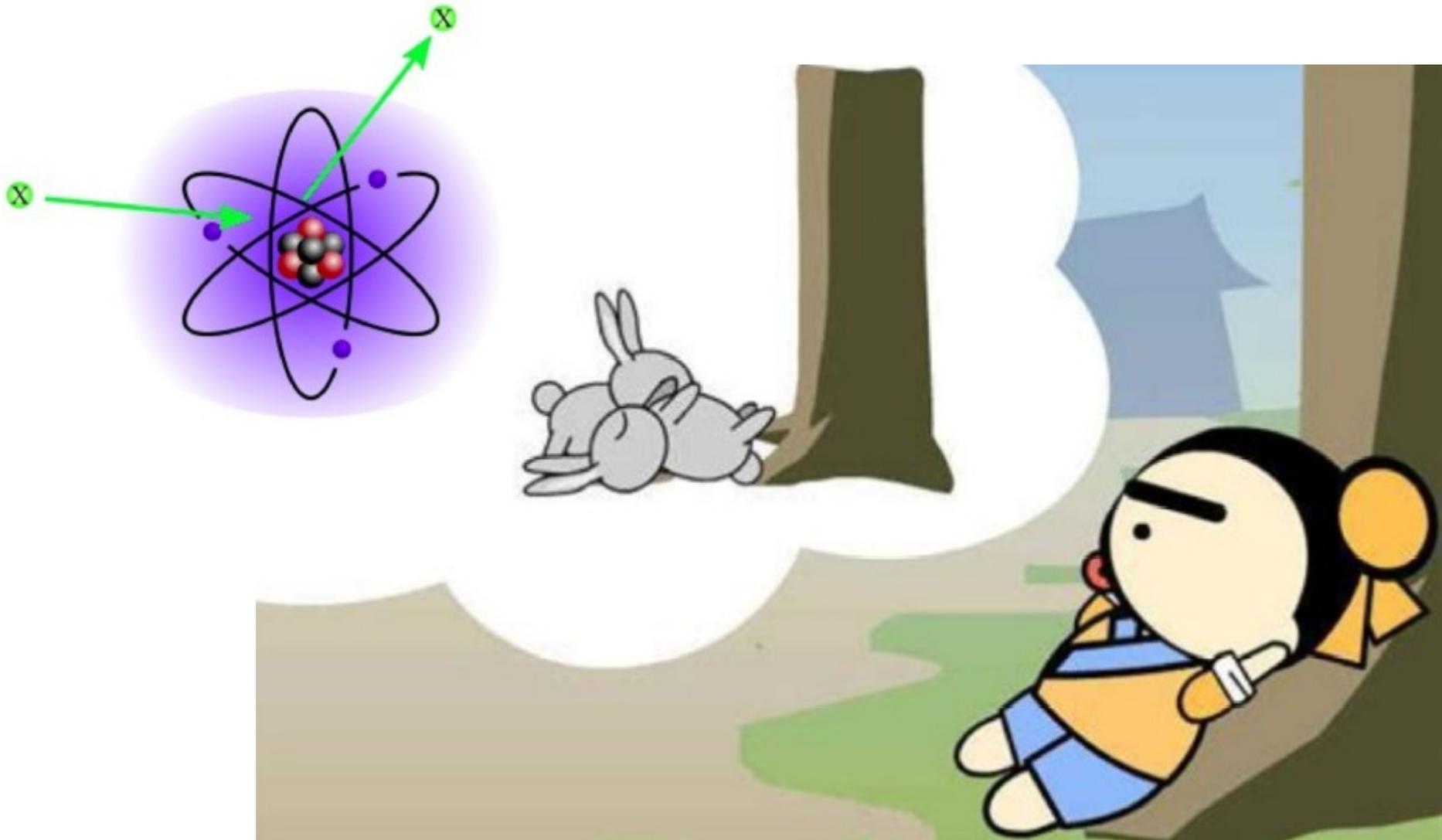
- 1) Thresholds in DM search
- 2) **Cosmic ray boosted DM & diurnal modulation**
- 3) Fermionic DM Absorption
- 4) Reactivate Forbidden DM @ Supermassive Black Hole

# CR Boosted DM

SFG, Jianglai Liu, Qiang Yuan, Ning Zhou [[Phys.Rev.Lett. 126 \(2021\) 9, 091804](#)]

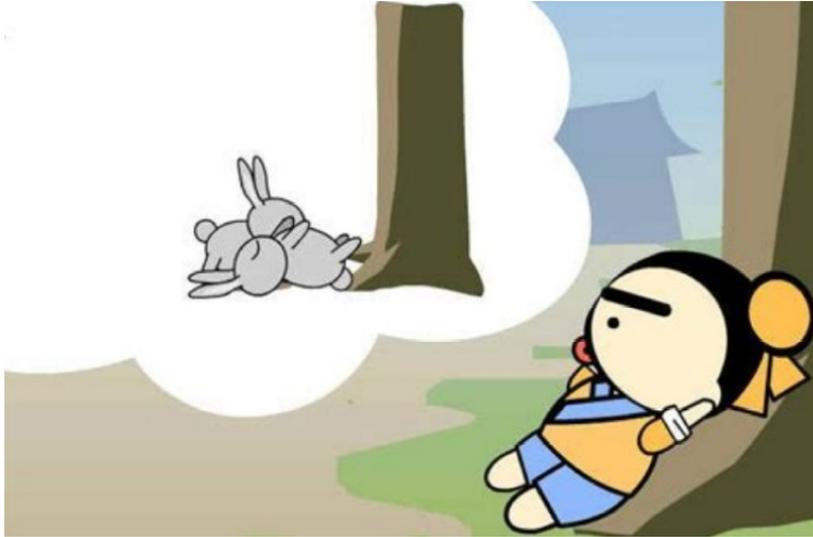
PandaX-II + SFG, Qiang Yuan [[Phys.Rev.Lett. 128 \(2022\) 17, 171801](#)]

# DM Detection vs 守株待兔



If the sleepy rabbit is too small, the tree cannot feel it!

# Beating Rabbits



**Stick has the same material as tree!**

**So long as direct detection is possible.**

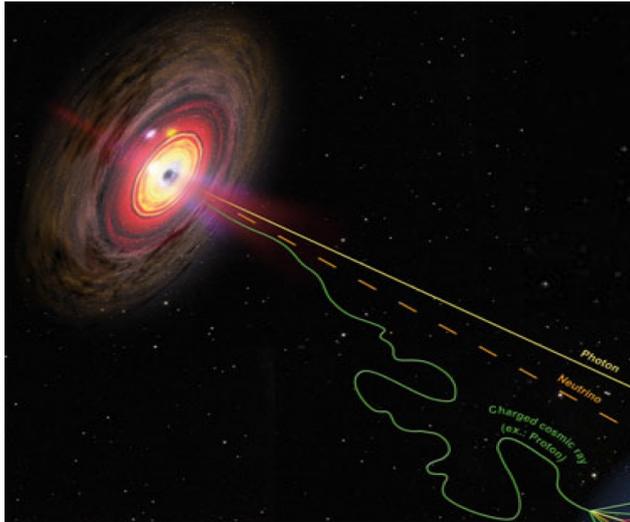


## Cosmic Ray Boosted DM

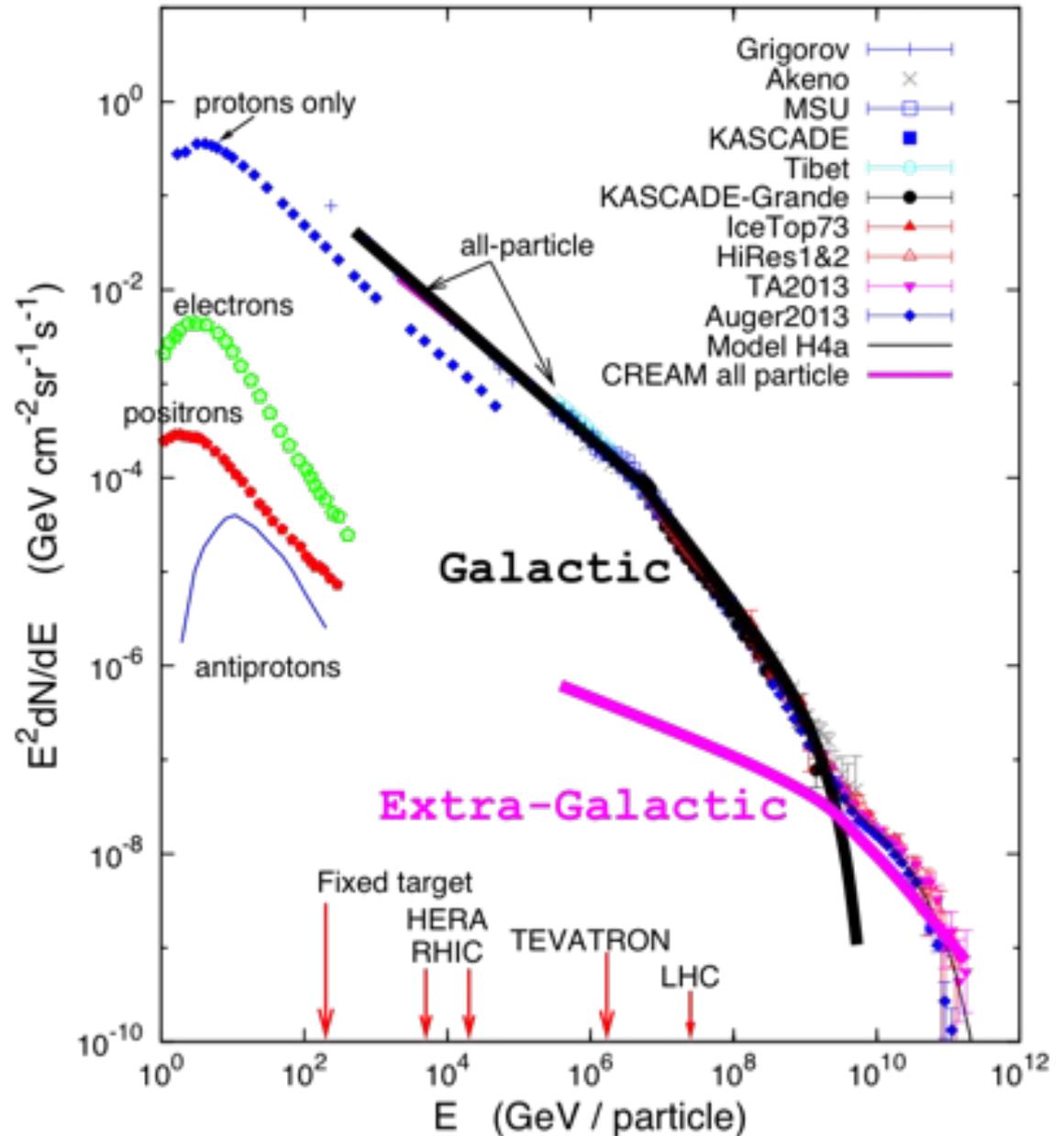
Cappiello, Ng & Beacom, PRD 2019 [arXiv:1810.07705]  
Bringmann & Pospelov, PRL 2019 [arXiv:1810.10543]  
Ema, Sala & Sato, PRL 2019 [arXiv:1811.00520]

# Cosmic Rays

Energies and rates of the cosmic-ray particles



- 1) Cosmic rays are almost everywhere in the galaxy
- 2) Cosmic rays are very energetic
- 3) Cosmic rays have several components (proton, electron, ...)

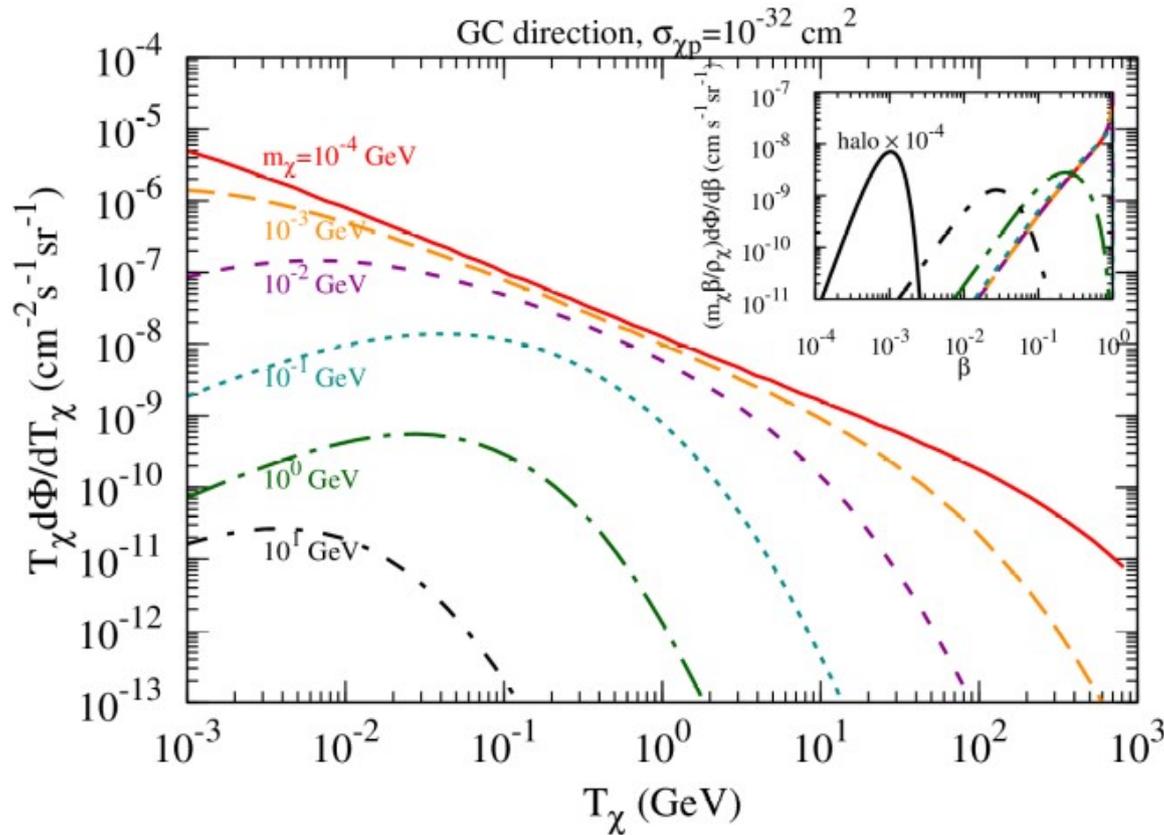


# Cosmic Ray Boosted DM

$$\zeta_{\chi}(\mathbf{r}, T_{\chi}) = \frac{\rho_{\chi}(|\mathbf{r}|)}{m_{\chi}} \sum_{i=p, \text{He}} \int_{T_i^{\min}}^{\infty} dT_i \frac{n_{\text{CR},i}(\mathbf{r}, T_i)}{T_{\chi}^{\max}(T_i)} \times v_i \sigma_{\chi i} G_i^2(Q^2),$$

with constant  $\sigma_{\chi p}$  & form factor  $G_i(Q^2) \equiv 1/(1 + Q^2/\Lambda_i^2)^2$ .

$$\frac{d\Phi}{dT_{\chi}}(\hat{\mathbf{n}}, T_{\chi}) = \frac{1}{4\pi} \int \zeta_{\chi}(\mathbf{r}, T_{\chi}) dl$$



1. Heavier DM  $\rightarrow$  smaller  $n_{\chi}$

$$n_{\chi} = \frac{\rho_{\chi}}{m_{\chi}} \propto \frac{1}{m_{\chi}}$$

2. lighter DM  $\rightarrow$  softer spectrum

$$0 \leq T_{\chi} \leq \frac{T_i(T_i + 2m_i)}{T_i + m_{\mu}^i}$$

$$\approx \frac{T_i^2}{T_i + m_i^2/2m_{\chi}}$$

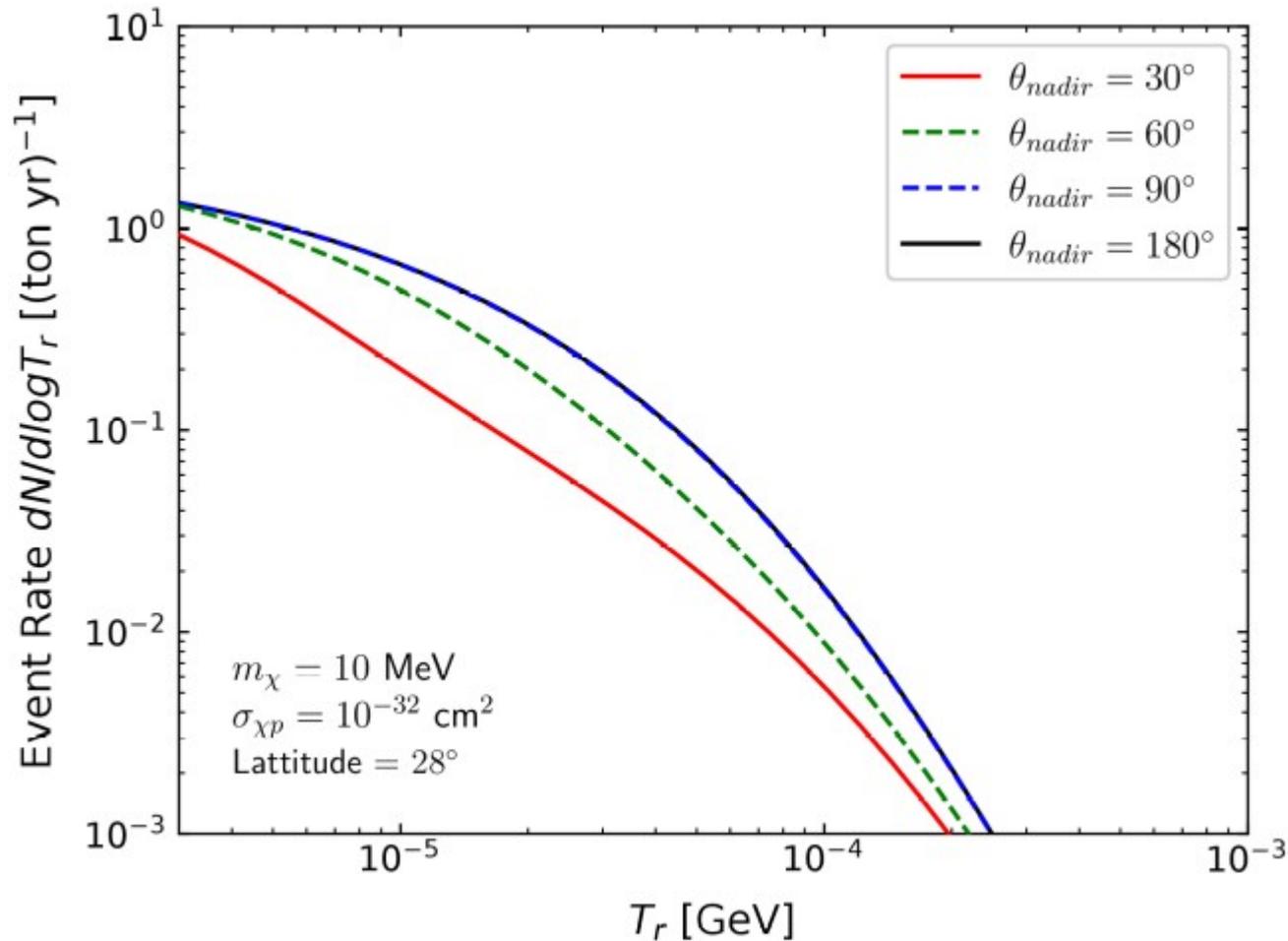
for  $m_{\chi} \ll m_i$

SFG, Jianglai Liu, Ning Zhou, Qiang Yuan, PRL 2021 [arXiv:2005.09480]

See also Bringmann & Pospelov, PRL 2019 [arXiv:1810.10543]

$$\sigma_{\chi A} = \sigma_{\chi p} A^2 \left[ \frac{m_A(m_\chi + m_p)}{m_p(m_\chi + m_A)} \right]^2$$

$$G_i(Q^2) \equiv 1/(1 + Q^2/\Lambda_i^2)^2$$

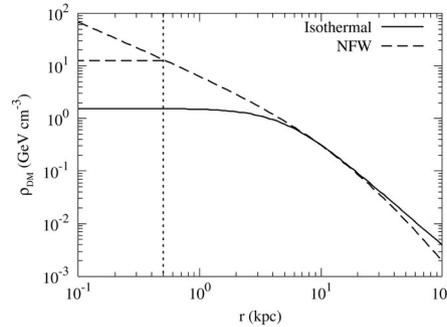


The nuclear recoil from CRDM is much more energetic!

SFG, Jianglai Liu, Ning Zhou, Qiang Yuan, PRL 2021 [arXiv:2005.09480]

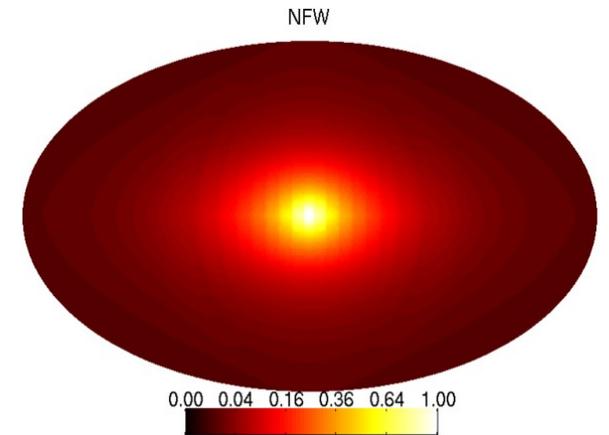
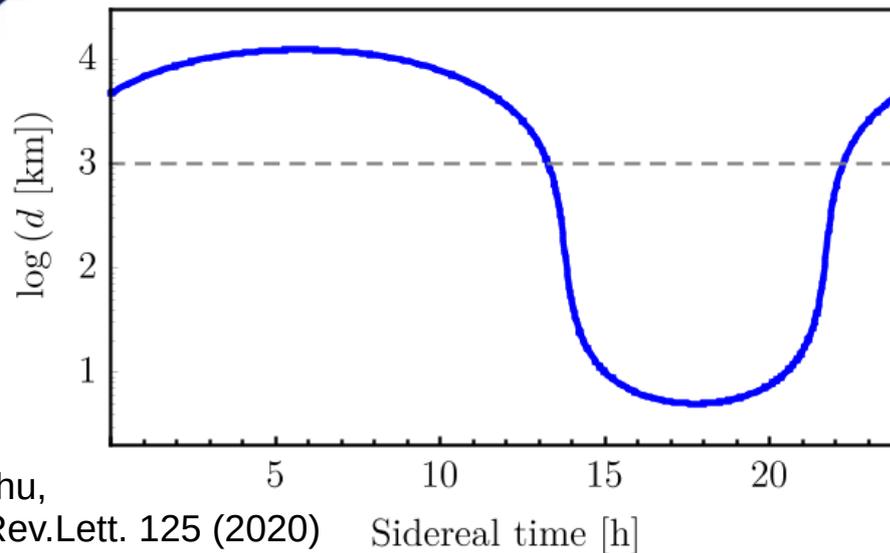
# Anisotropies → Diurnal Effect

SFG, Jianglai Liu, Ning Zhou, Qiang Yuan,  
PRL 2021 [arXiv:2005.09480]



DM

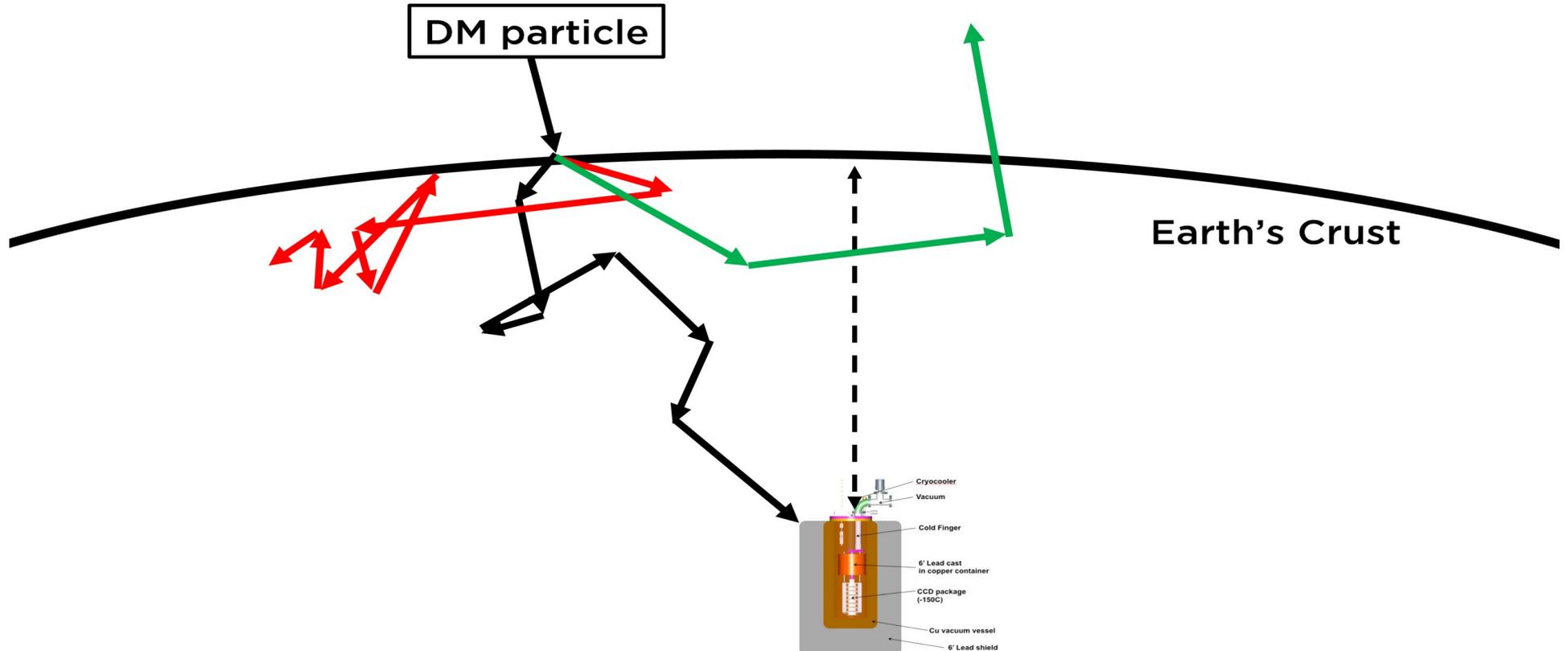
Cosmic Rays



Fornal, Sandick, Shu,  
Su & Zhao, Phys.Rev.Lett. 125 (2020)  
16, 161804 [2006.11264]

Sidereal time [h]

See also Xia, Xu & Zhou [2206.11454]



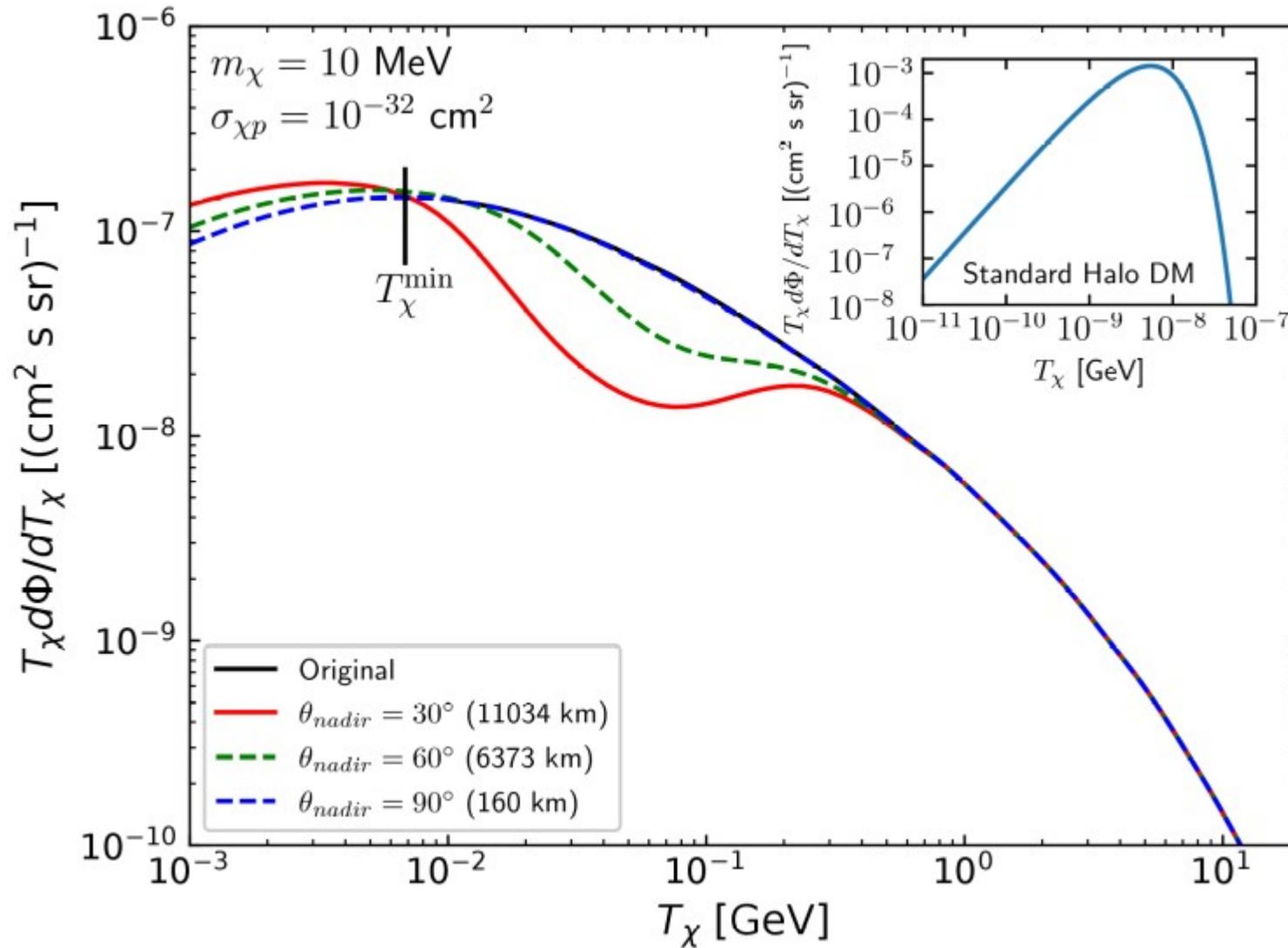
$$\frac{\partial}{\partial l} \frac{d\Phi(l, T_\chi)}{d \ln T_\chi} = \frac{\rho_N(l)}{m_N} \sigma_{\chi N} G_N^2(Q^2) \left[ -\frac{d\Phi(l, T_\chi)}{d \ln T_\chi} + \int \frac{d\Phi(l, T'_\chi)}{d \ln T'_\chi} \frac{T_\chi (T'_\chi + m_\mu^N)}{T'_\chi (T'_\chi + 2m_\chi)} d \ln T'_\chi \right]$$

Transfer to lower energy

Come from higher energy

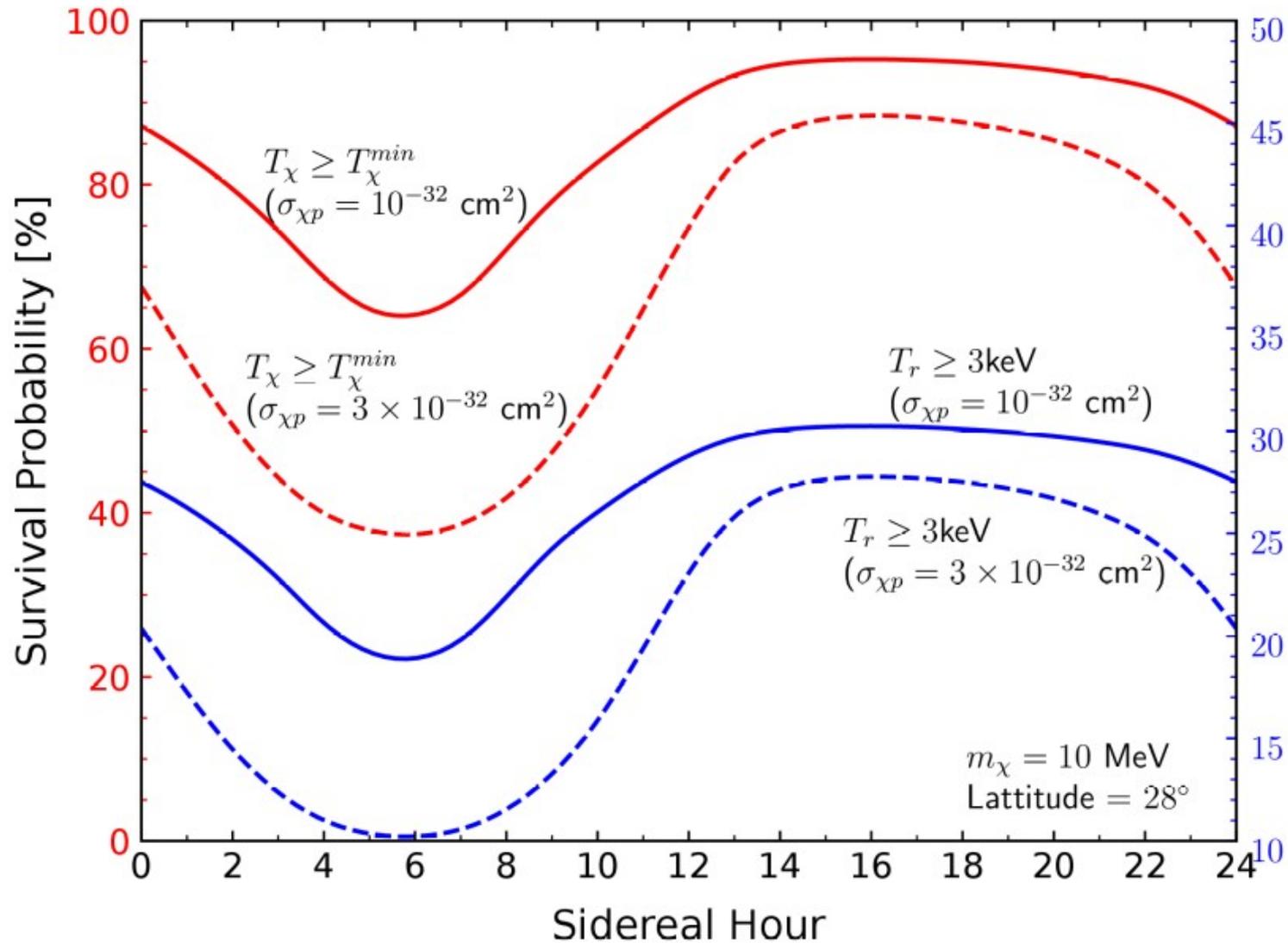
Xia, Xu & Zhou, JCAP 02 (2022) 02, 028 [2111.05559]

# Attenuated Flux

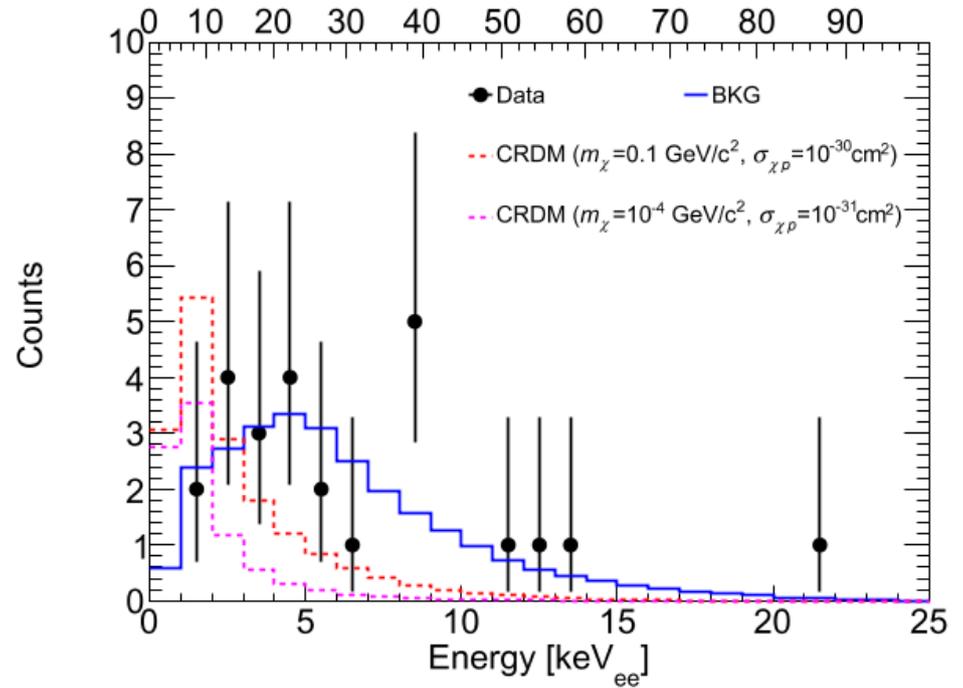
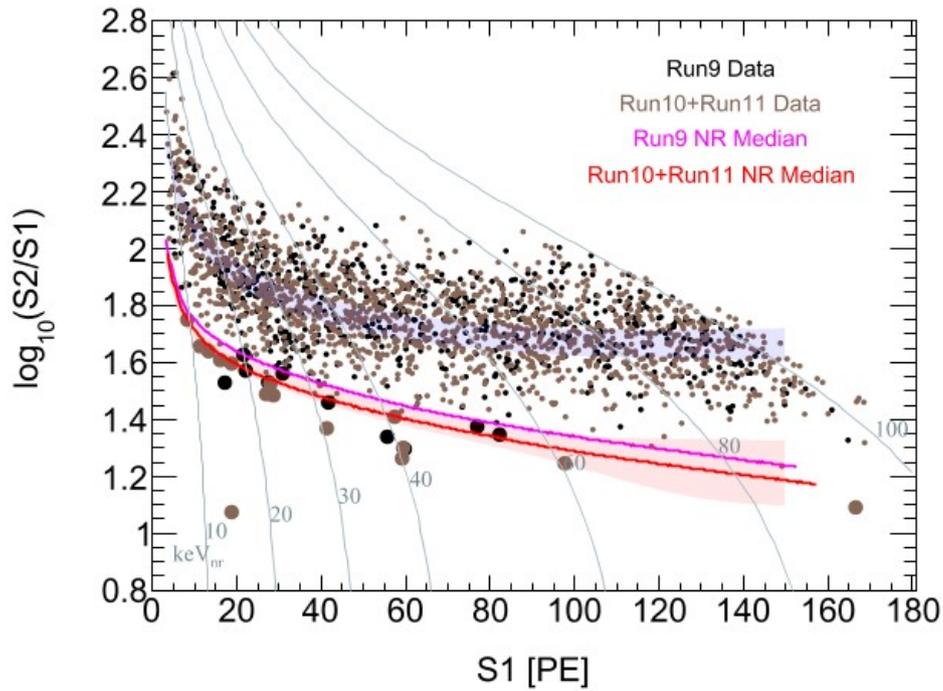


SFG, Jianglei Liu, Ning Zhou, Qiang Yuan, PRL 2021 [arXiv:2005.09480]

# Diurnal Effect

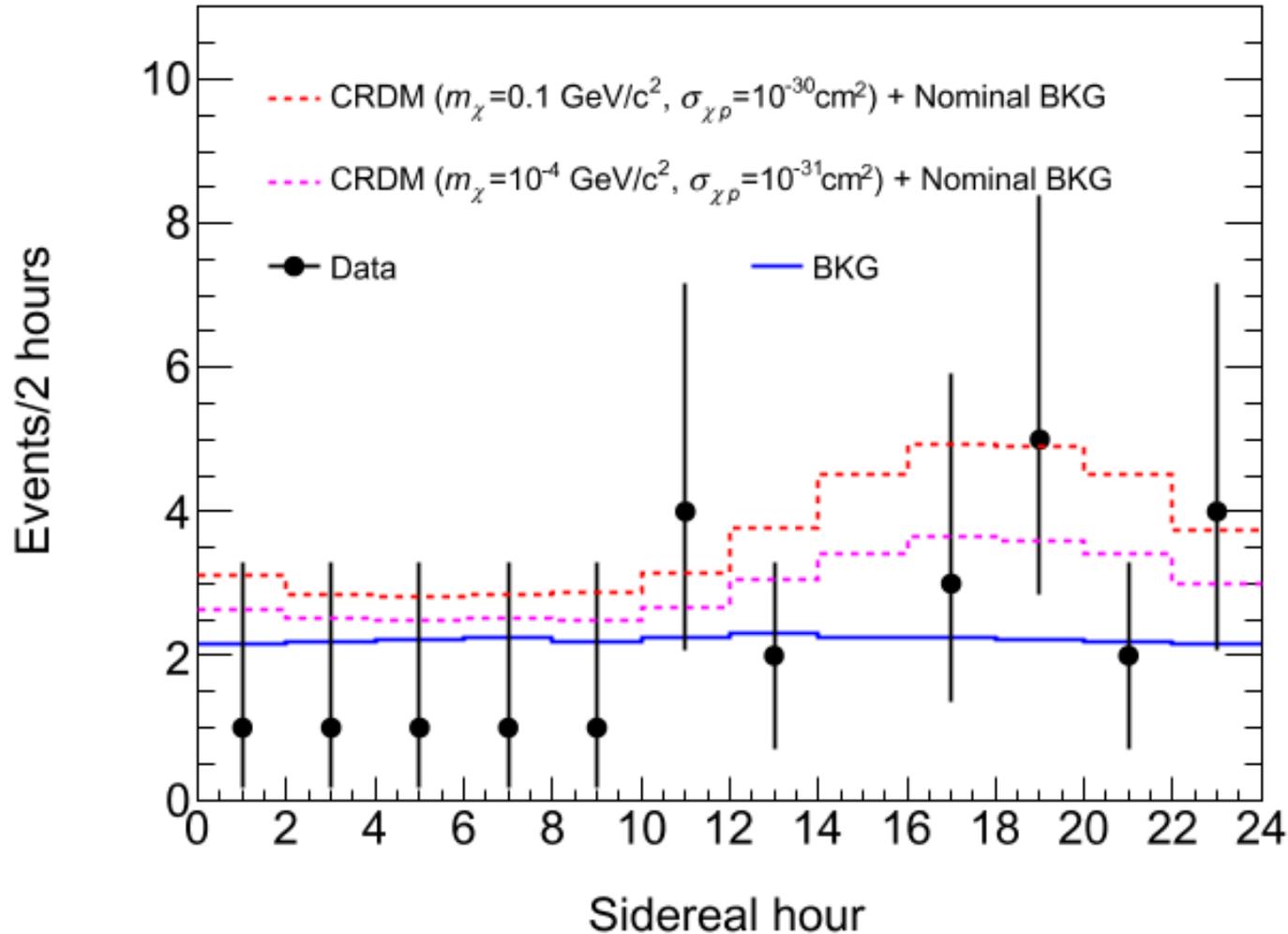


SFG, Jianglei Liu, Ning Zhou, Qiang Yuan, PRL 2021 [arXiv:2005.09480]



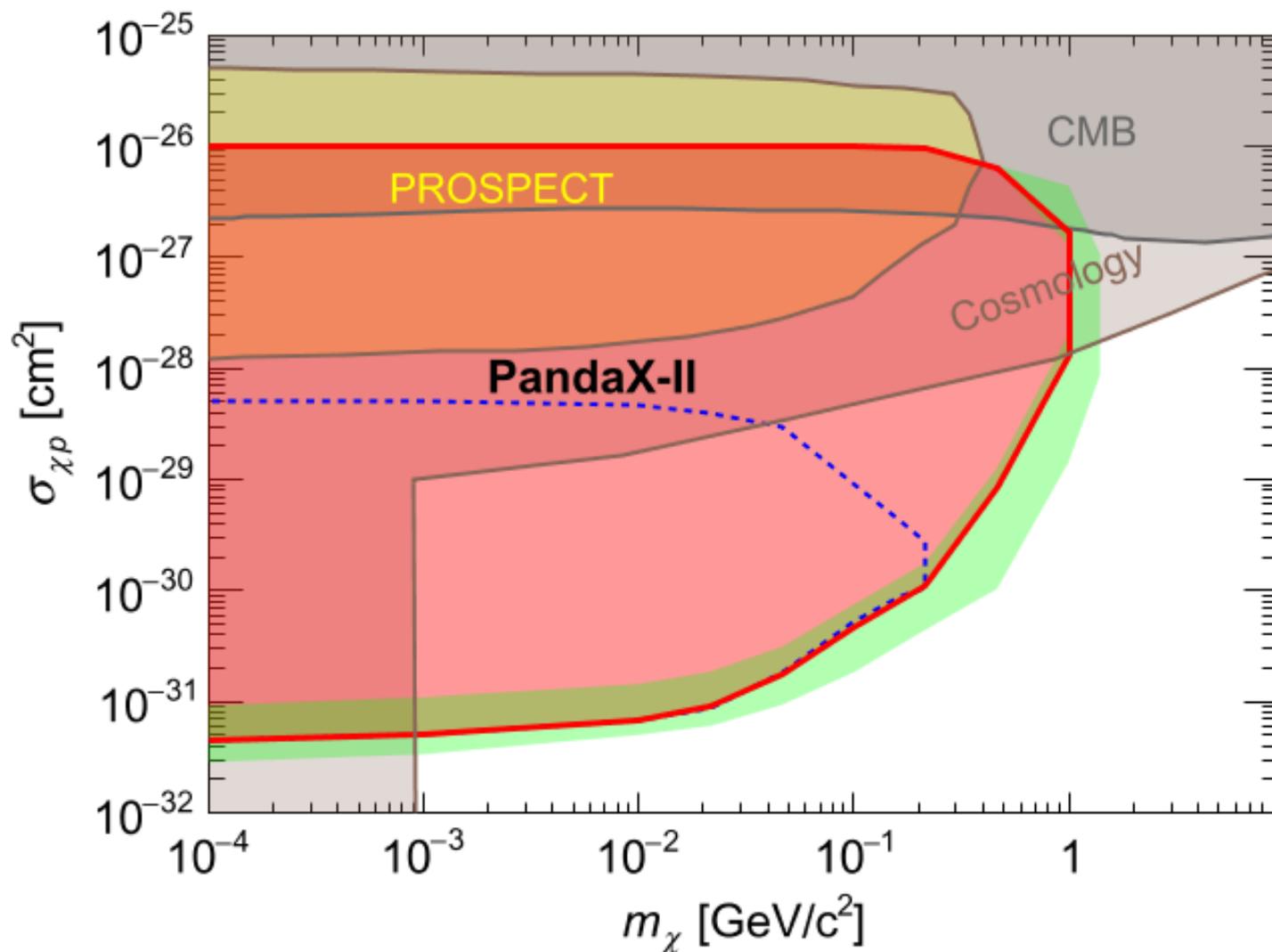
PandaX-II + **SFG** & Qiang Yuan, Phys.Rev.Lett. 128 (2022) 17, 171801 [2112.08957]

# Diurnal Modulation @ PandaX



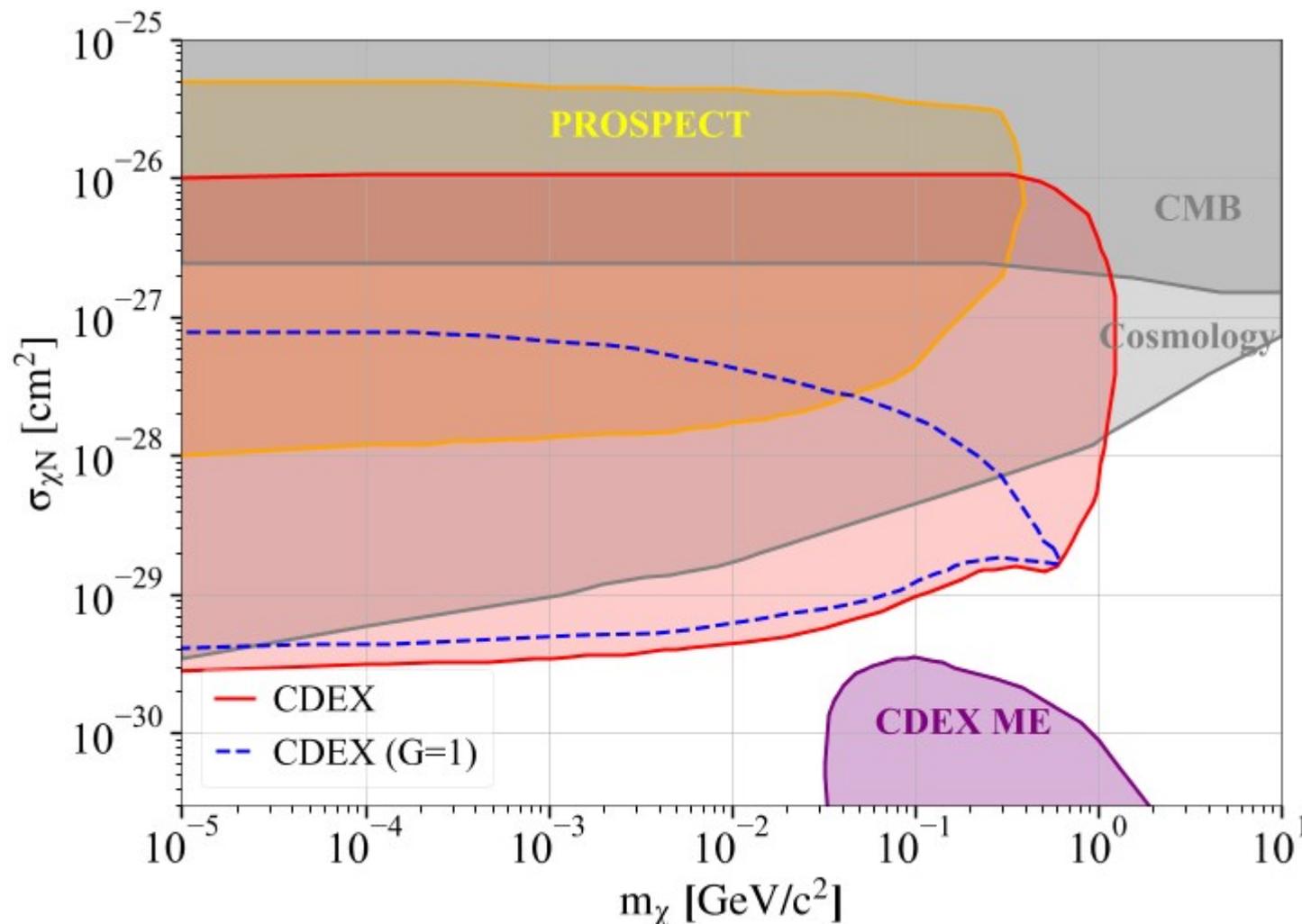
PandaX-II + **SFG** & Qiang Yuan, Phys.Rev.Lett. 128 (2022) 17, 171801 [2112.08957]

# Sensitivity @ PandaX



PandaX-II + **SFG** & Qiang Yuan, Phys.Rev.Lett. 128 (2022) 17, 171801 [2112.08957]

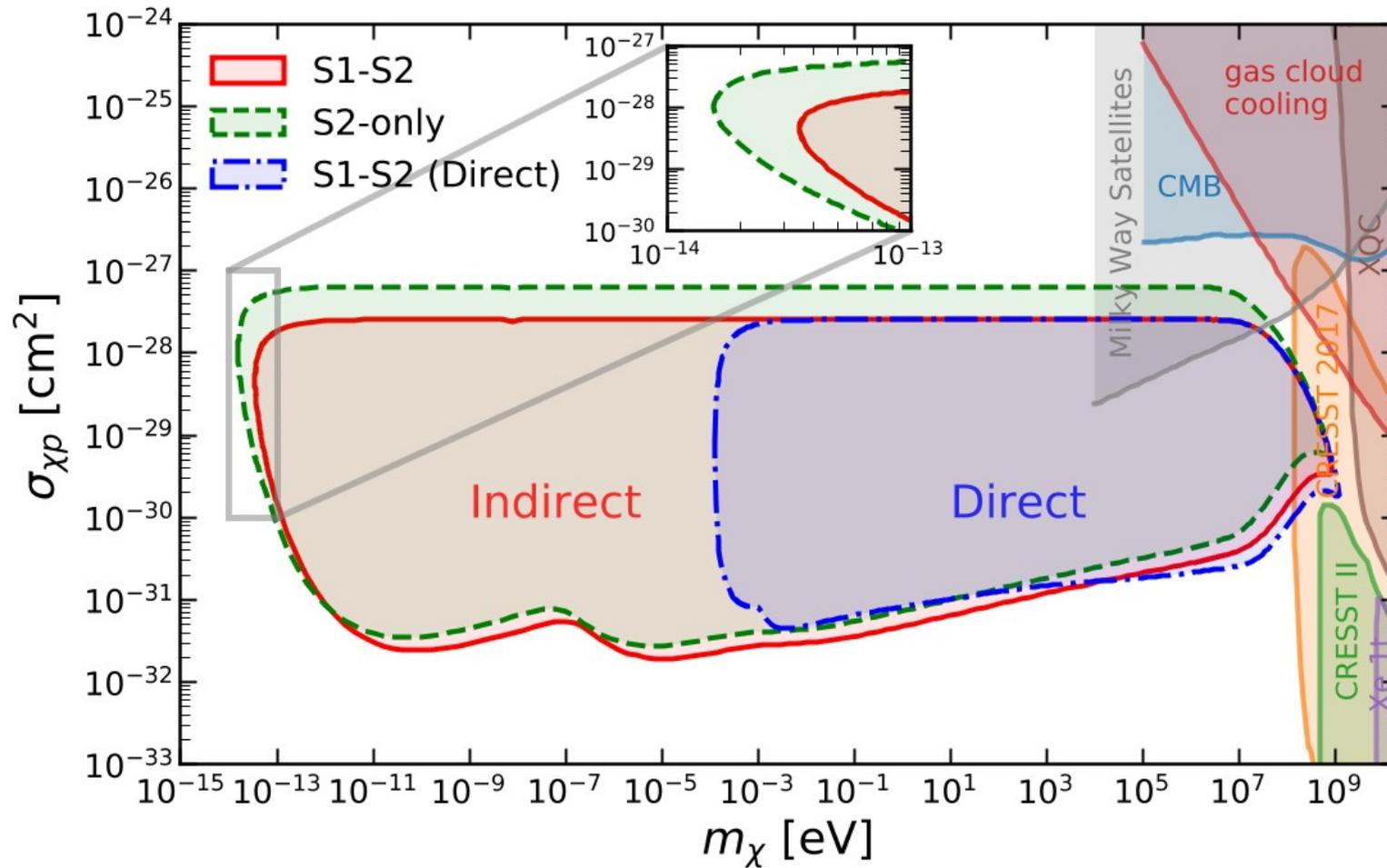
# Sensitivity @ CDEX



Zhang, Lei & Tang, Chin.Phys.C 46  
(2022) 8, 085103 [2008.07116]

CDEX, Phys.Rev.D 106 (2022) 5,  
052008 [2201.01704]

# Kinematic Limits

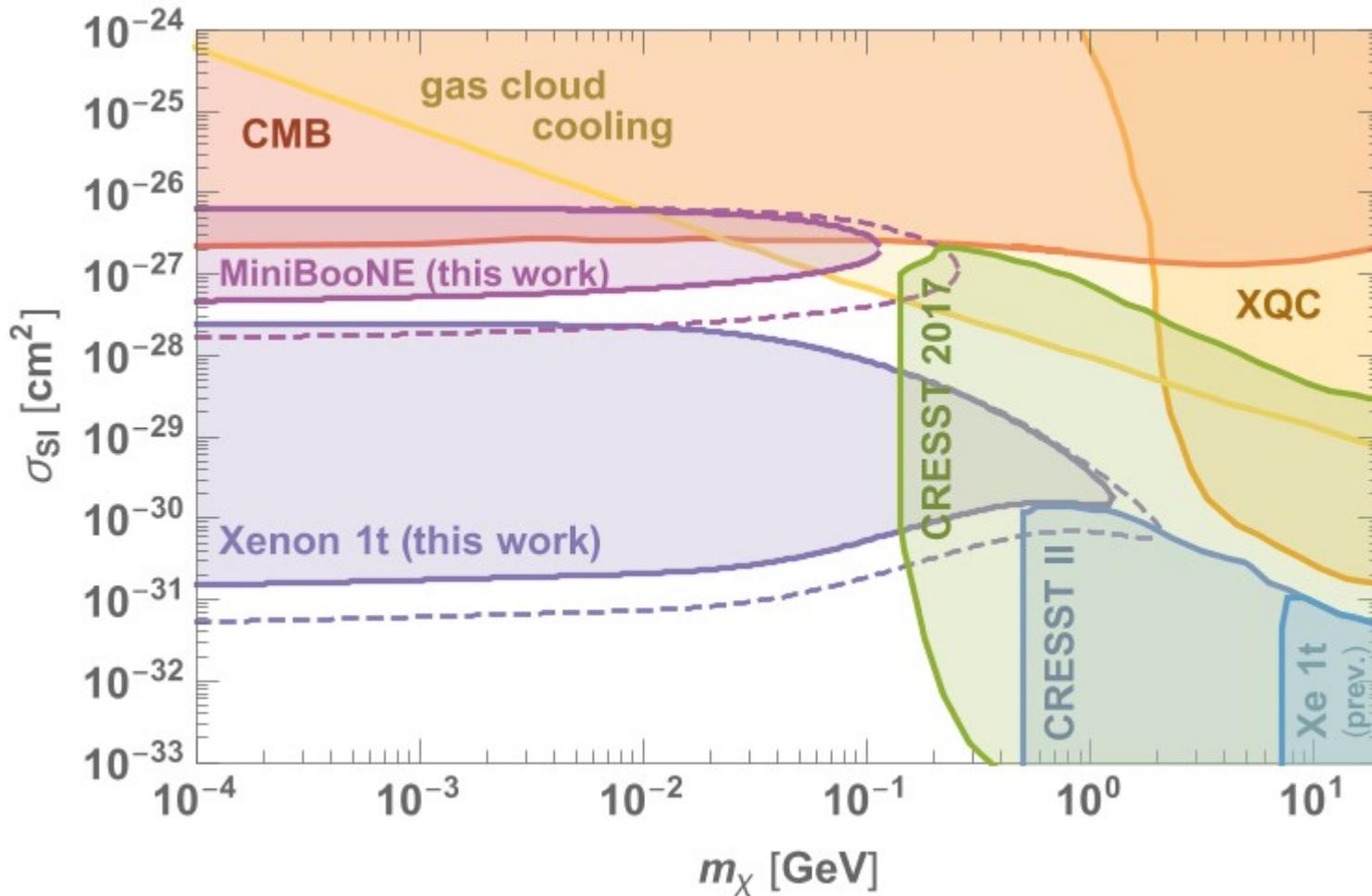


$$0 \leq T_\chi \leq \frac{T_i(T_i + 2m_i)}{T_i + m_\mu^i} \approx \frac{T_i^2}{T_i + m_i^2/2m_\chi} \quad \text{for } m_\chi \ll m_i$$

Xia, Xu & Zhou, Nucl.Phys.B 969 (2021) 115470 [2009.00353]

# Backup Slides

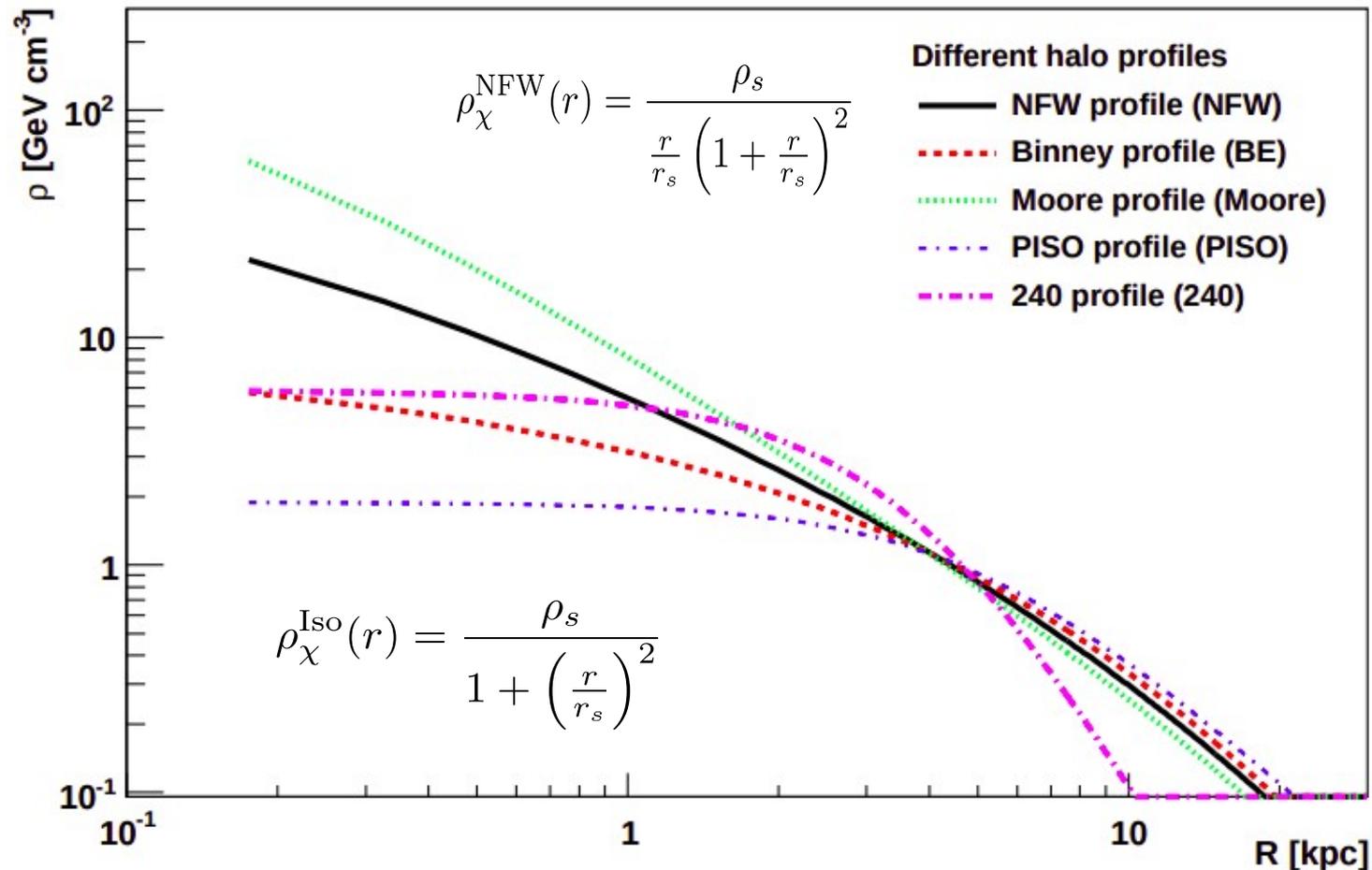
# Projected Sensitivities



Bringmann & Pospelov, PRL 2019 [arXiv:1810.10543]

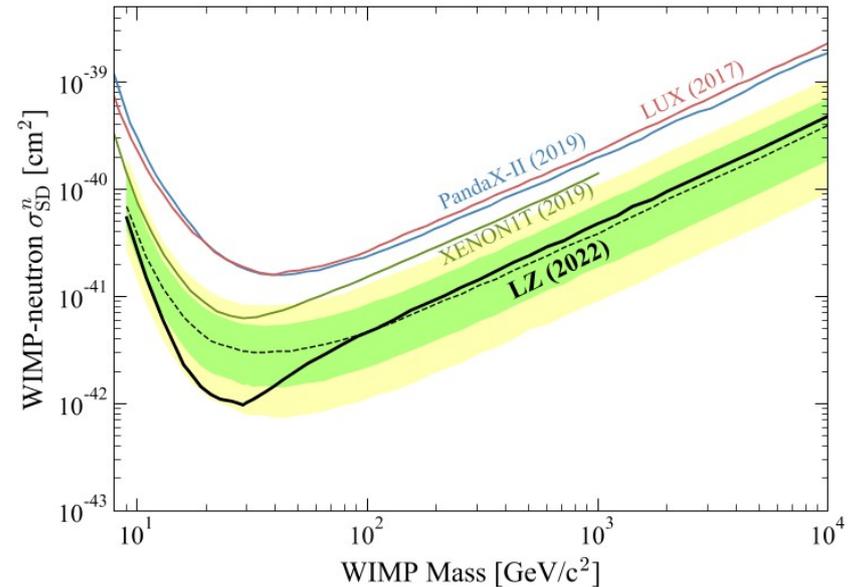
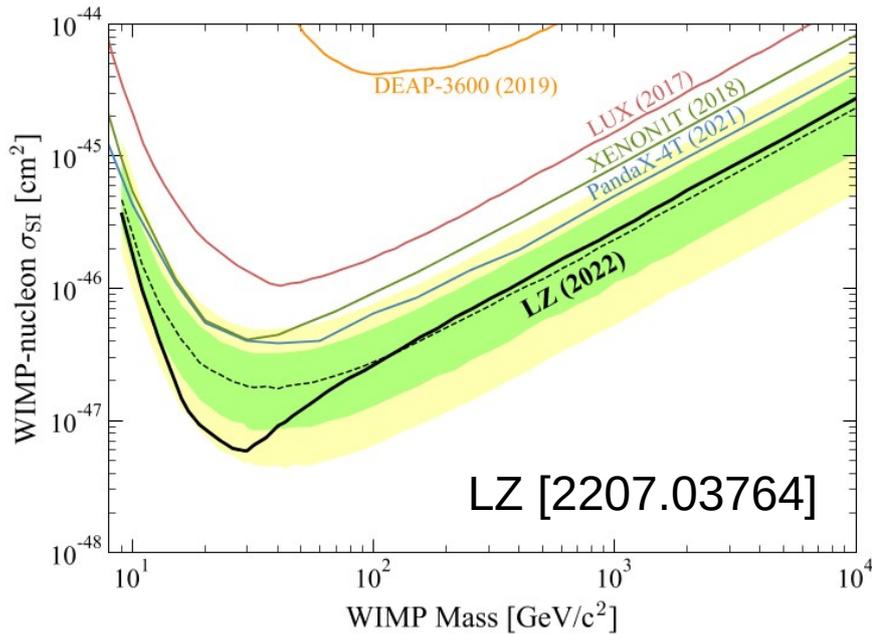
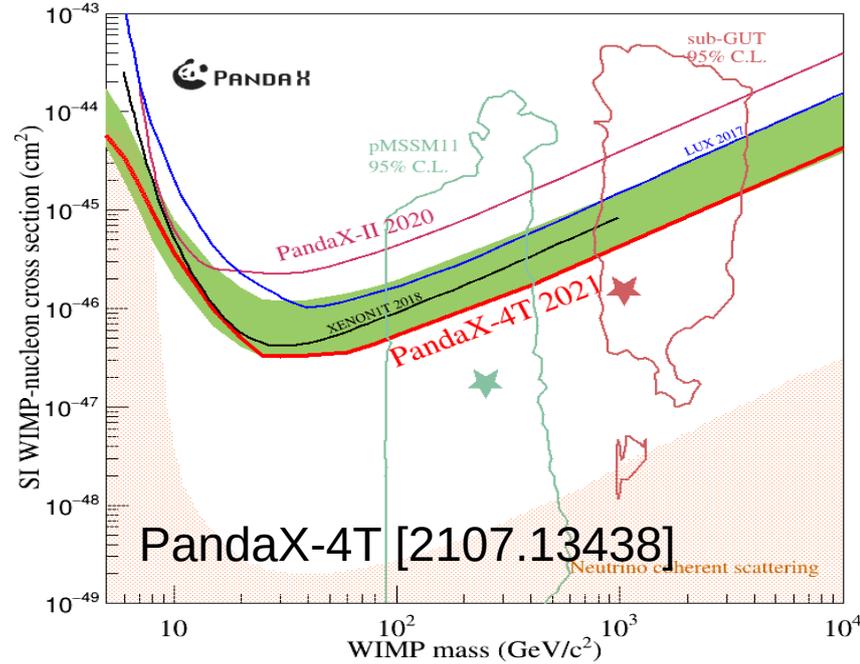
# Galaxy DM Density Profile

Weber & Boer [0910.4272]

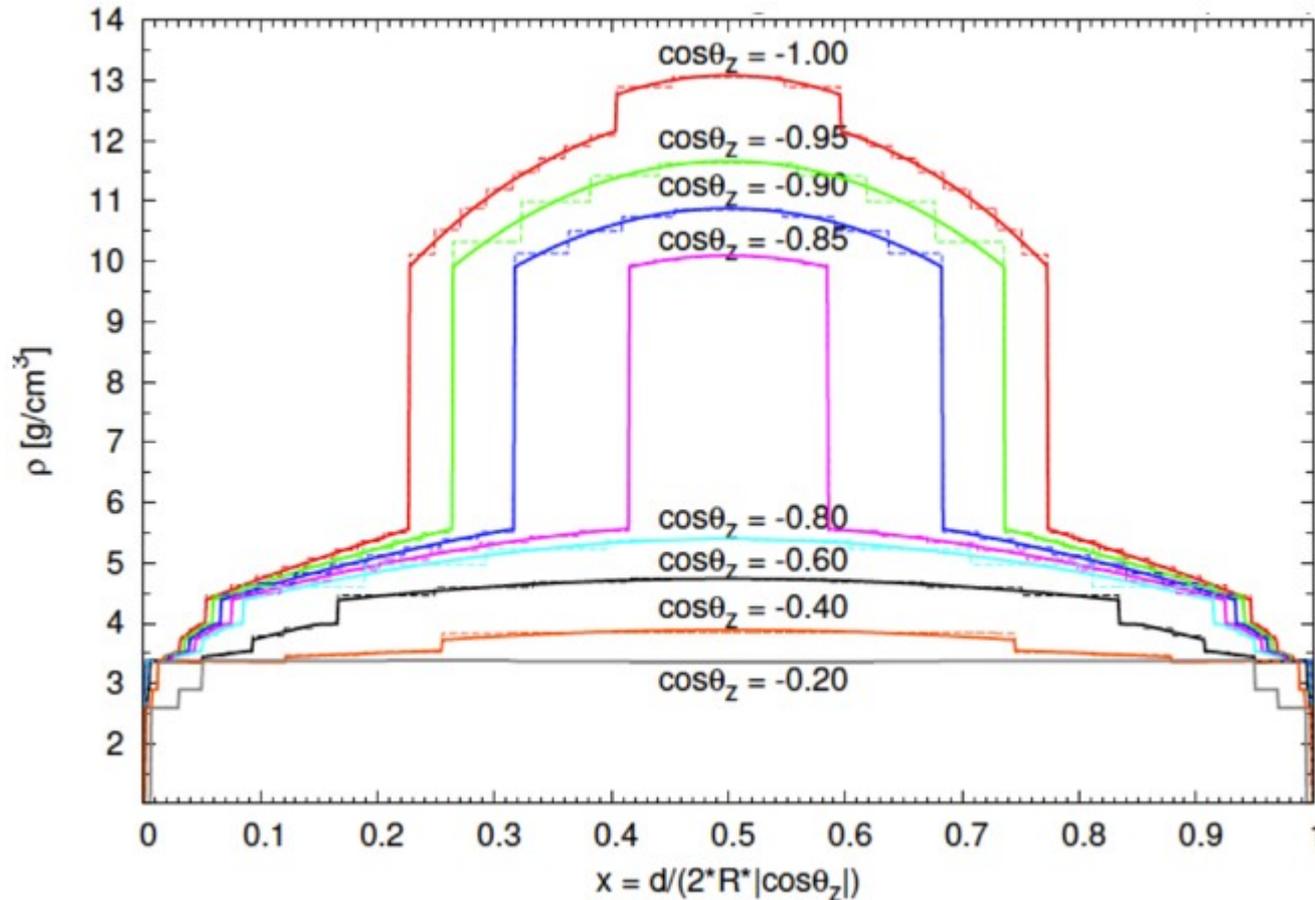


DM particles are almost at rest.

# Current Sensitivity



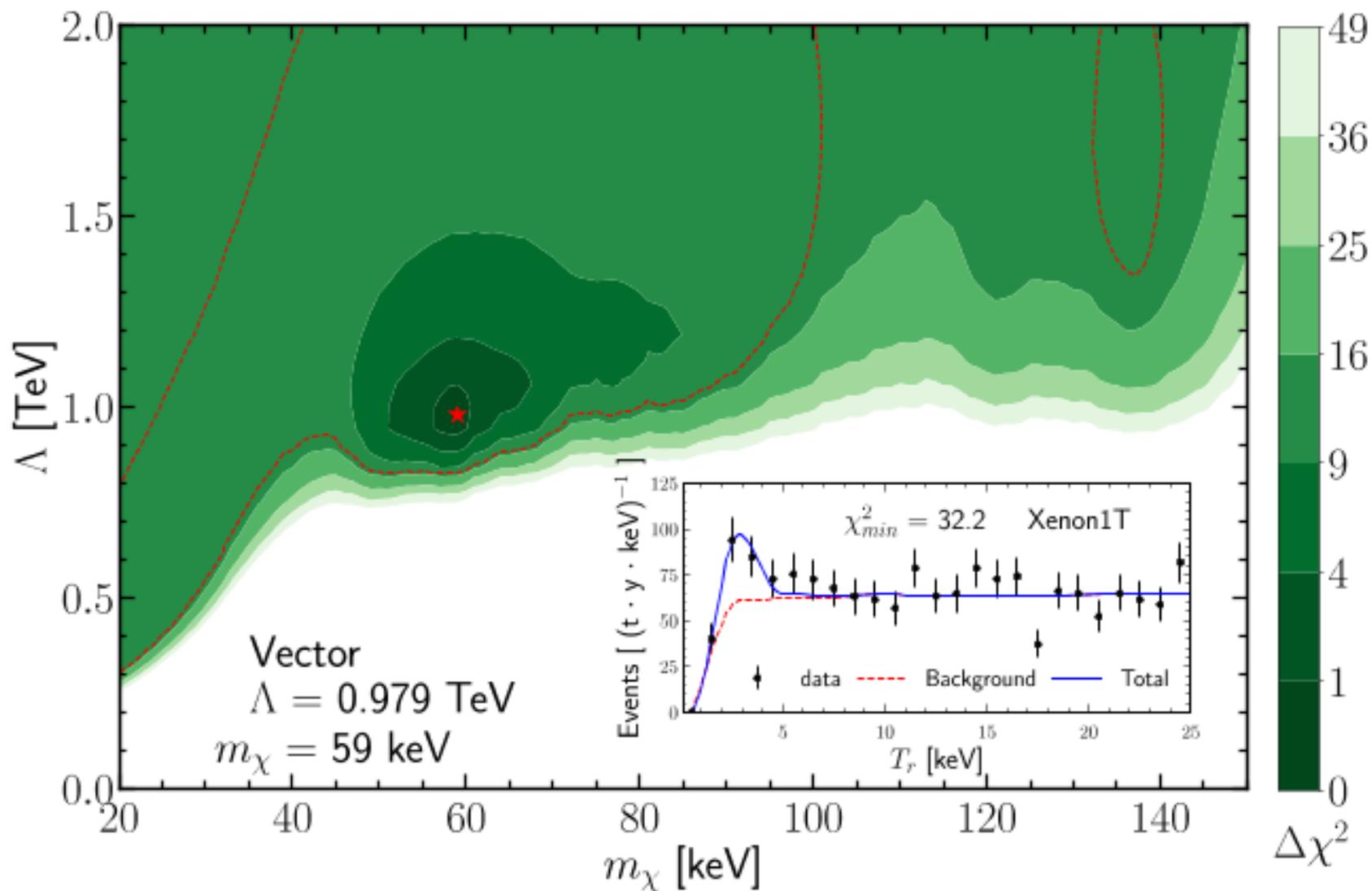
# Earth Density Profile



SFG, Hagiwara, Rott [1309.3176]

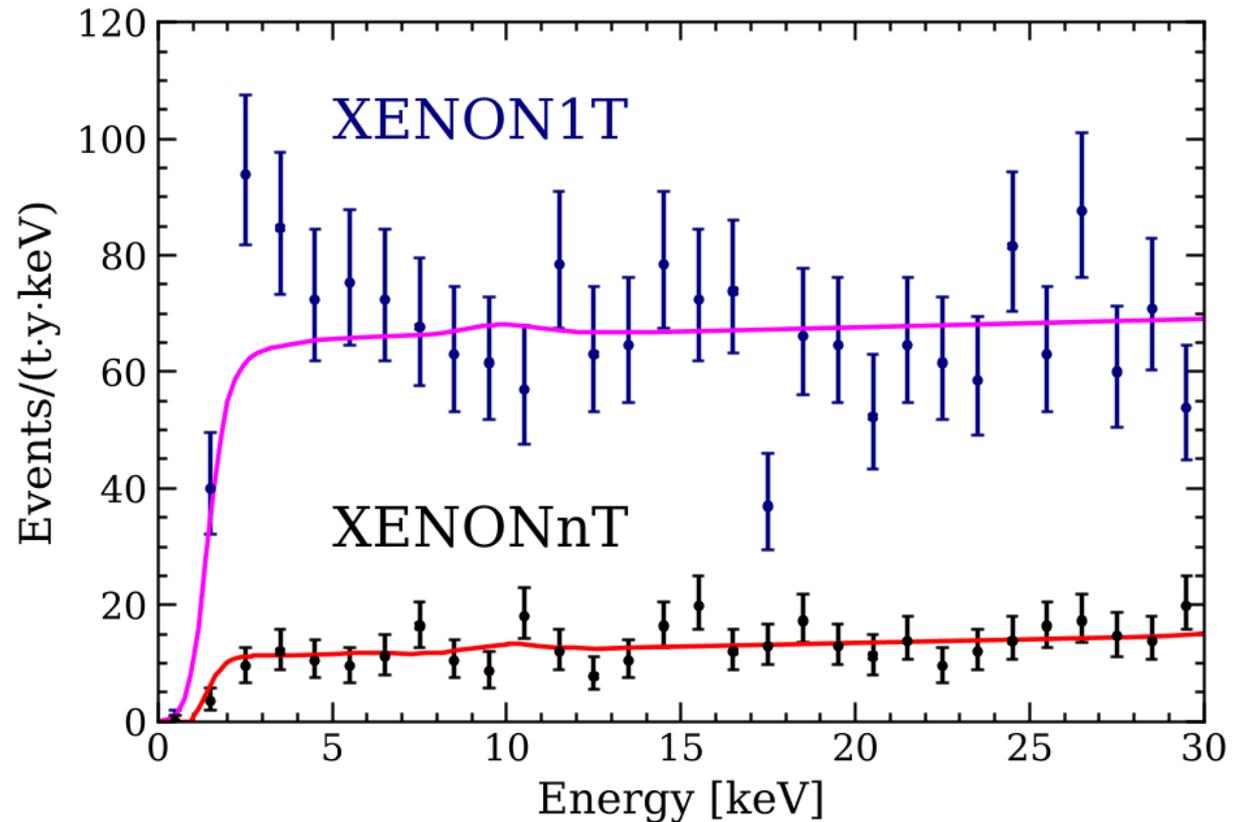
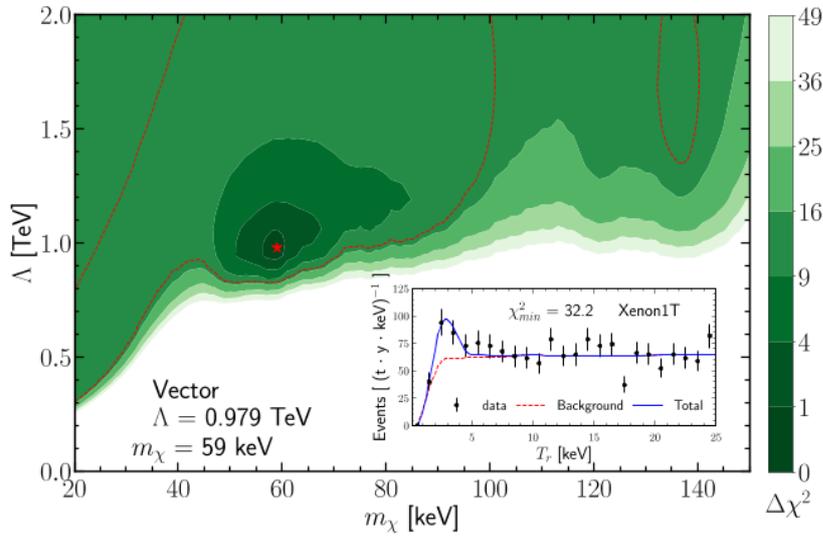
The earth matter can be approximated by 2-step profile.

# Xenon1T excess



SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]

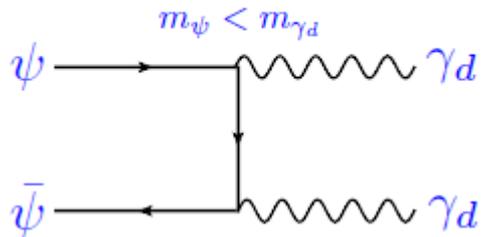
# Xenon1T excess



Knut Dundas Morå @ IDM2022  
XENONnT, Phys.Rev.Lett. 129 (2022) 16, 161805 [2207.11330]

# Forbidden DM

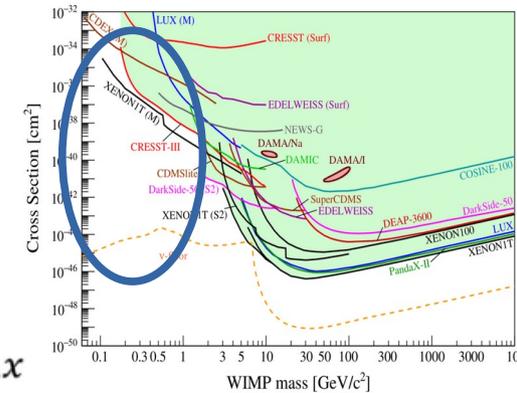
1) forbidden annihilations:



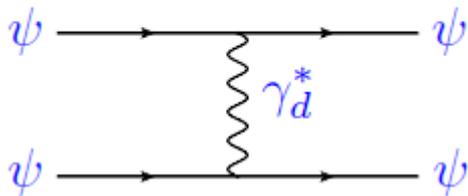
$$\langle \sigma_{\gamma_d \gamma_d} v \rangle \sim \alpha_d^2 / m_{\gamma_d}^2$$



$$\langle \sigma_{\chi \bar{\chi}} v \rangle = \frac{(n_{\gamma_d}^{eq})^2}{(n_{\psi}^{eq})^2} \langle \sigma_{\gamma_d \gamma_d} v \rangle \approx 8\pi f_{\Delta} \frac{\alpha_d^2}{m_{\psi}^2} e^{-2\Delta x}$$

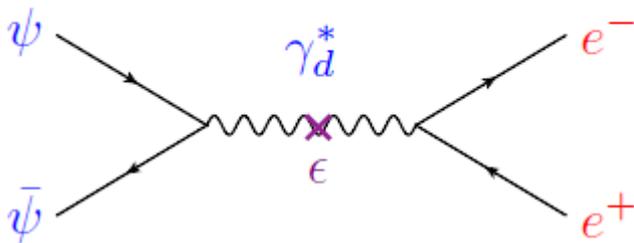


2) self-interactions:

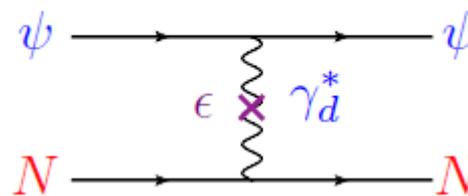


- Naturally includes large self-interactions.
- An exponentially larger cross section than the forbidden annihilation rate.

3) indirect detection:



4) direct detection:



- Signal suppressed by

D'Agnolo & Ruderman, Phys.Rev.Lett. 115 (2015) 6, 061301 [1505.07107]