

Minimal decaying dark matter: from cosmological tensions to neutrino constraints

Lea Fuß, Mathias Garny, Alejandro Ibarra
based on arXiv:2403.15543

The Dark Matter Landscape: From Feeble to Strong Interactions (MITP)

August 30, 2024



- I. Cosmological model
- II. Theoretical model building
- III. New phenomenology
 - Neutrinos
 - Production via freeze-in
 - Low-mass signatures
- IV. Outlook and Summary

- I. Cosmological model
- II. Theoretical model building
- III. New phenomenology
 - Neutrinos
 - Production via freeze-in
 - Low-mass signatures
- IV. Outlook and Summary

Cosmological tensions: A hint for something new?



S_8 tension:

persistent tension of $2 - 3\sigma$
between early and late universe
measurements in the clustering on
small scales

Cosmological tensions: A hint for something new?



S_8 tension:

persistent tension of $2 - 3\sigma$
between early and late universe
measurements in the clustering on
small scales

$$S_8 = \sigma_8 \sqrt{\Omega_m / 0.3}$$

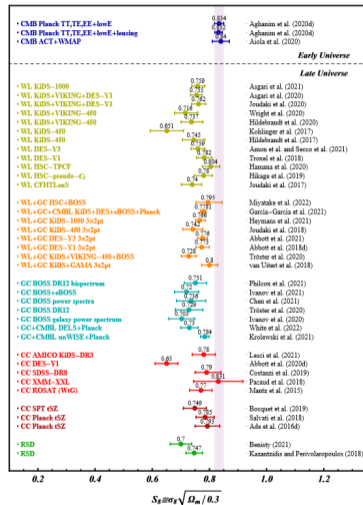
$$\sigma_R^2 = \int_0^\infty \frac{dk}{k} \Delta_m(k) \tilde{W}_R(k)^2$$

Cosmological tensions: A hint for something new?



S_8 tension:

persistent tension of $2 - 3\sigma$
between early and late universe
measurements in the clustering on
small scales



[Abdalla et. al., arXiv:2203.06142]

Decaying Cold Dark Matter

DM model that generates suppression on small scales

→ **Decaying Cold Dark Matter (DCDM)**

DCDM → WDM + DR

Decaying Cold Dark Matter

DM model that generates suppression on small scales

→ **Decaying Cold Dark Matter (DCDM)**

DCDM → WDM + DR

2 parameters: **lifetime** τ , **mass splitting** $\epsilon = \frac{1}{2} \left(1 - \frac{m^2}{M^2}\right)$

Decaying Cold Dark Matter

DM model that generates suppression on small scales

→ **Decaying Cold Dark Matter (DCDM)**

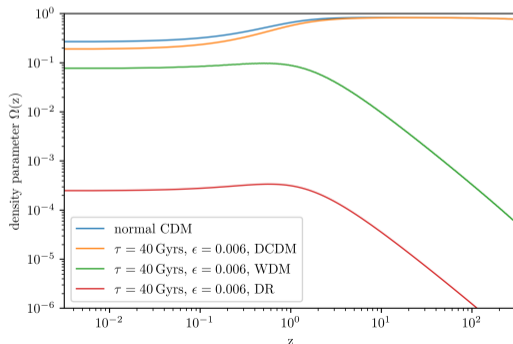
DCDM → WDM + DR

2 parameters: **lifetime** τ , **mass splitting** $\epsilon = \frac{1}{2} \left(1 - \frac{m^2}{M^2}\right)$

$$\dot{\bar{\rho}}_{\text{dcdm}} = -3\mathcal{H}\bar{\rho}_{\text{dcdm}} - a\Gamma\bar{\rho}_{\text{dcdm}}$$

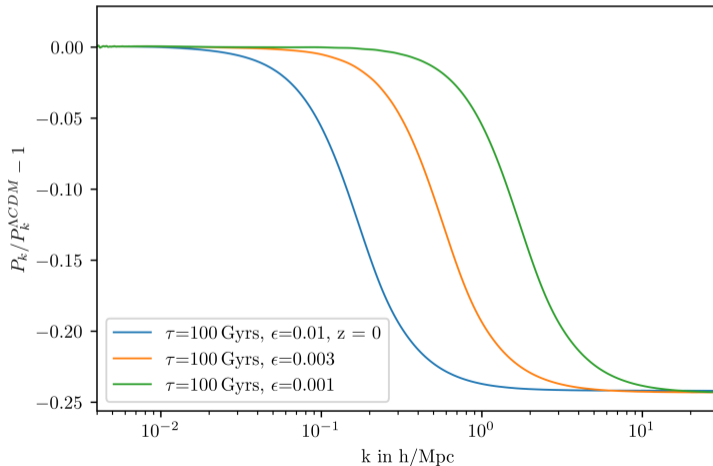
$$\dot{\bar{\rho}}_{\text{wdm}} = -3(1 + \omega)\mathcal{H}\bar{\rho}_{\text{wdm}} \\ + (1 - \epsilon)a\Gamma\bar{\rho}_{\text{dcdm}}$$

$$\dot{\bar{\rho}}_{\text{dr}} = -4\mathcal{H}\bar{\rho}_{\text{dr}} + \epsilon a\Gamma\bar{\rho}_{\text{dcdm}}$$



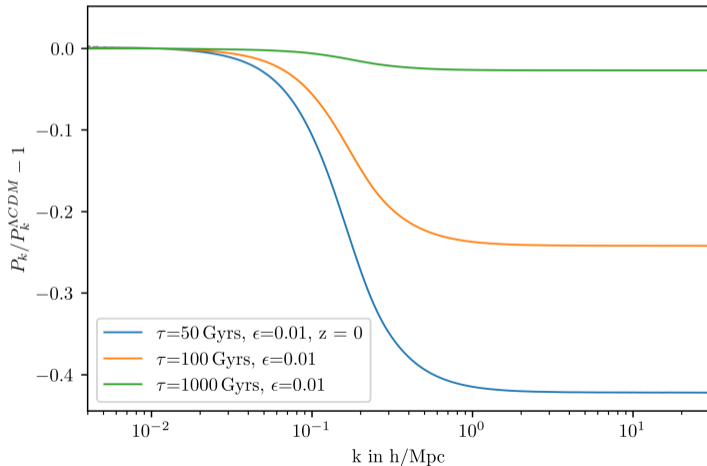
Suppression through decay

- ▶ Compute power spectrum with modified CLASS code for DCDM from [Abellan, Murgia, Poulin, arXiv:2102.12498]



Suppression through decay

- ▶ Compute power spectrum with modified CLASS code for DCDM from [Abellan, Murgia, Poulin, arXiv:2102.12498]



Lyman- α forest

[LF, Garny, arXiv:2210.06117]

CMB and BAO

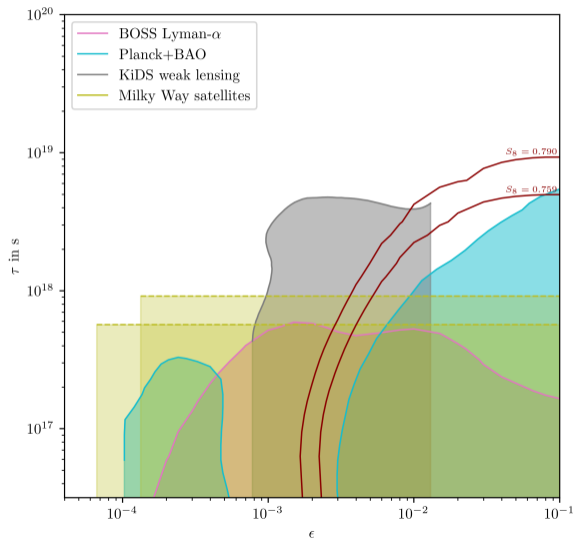
[Simon et al., arXiv:2203.07440]

Weak lensing shear data

[Bucko et al., arXiv:2307.03222]

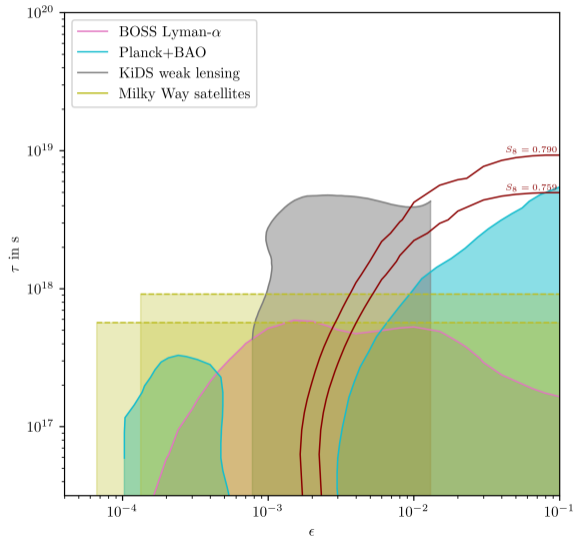
DM halo evolution

[DES Collab., arXiv:2201.11740]



singles out parameter space of interest to address S_8 tension:

- ▶ $\tau \sim 10^{18} \text{ s} \sim 100 \text{ Gyrs}$
- ▶ $\epsilon \sim 10^{-2}$



- I. Cosmological model
- II. Theoretical model building**
- III. New phenomenology
 - Neutrinos
 - Production via freeze-in
 - Low-mass signatures
- IV. Outlook and Summary

A theoretical framework

Question: *How can such a model be realized theoretically?*

A theoretical framework

Question: *How can such a model be realized theoretically?*

Idea: “DR” only has to couple sufficiently weakly to the SM particles to be considered dark

Can DM decay instead into neutrinos?

Question: *How can such a model be realized theoretically?*

Idea: “DR” only has to couple sufficiently weakly to the SM particles to be considered dark

Can DM decay instead into neutrinos?

Minimal approach: as few ingredients as possible

- ▶ 2 new fermionic particles N_1 and N_2 as DM
- ▶ SM neutrinos as “DR”
- ▶ described by effective interaction

What we want:
(for S_8)

decay into neutrinos
with $\tau \sim 10^{18}$ s

The challenge

What we want:
(for S_8)

decay into neutrinos
with $\tau \sim 10^{18}$ s

What we need:
(indirect detection constraints)

decay into $e^+/e^-/\gamma$
with $\tau \gtrsim 10^{26} - 10^{30}$ s

The challenge

What we want:
(for S_8)

decay into neutrinos
with $\tau \sim 10^{18} \text{ s}$

What we need:
(indirect detection constraints)

decay into $e^+/e^-/\gamma$
with $\tau \gtrsim 10^{26} - 10^{30} \text{ s}$

Challenge!

\Rightarrow coupling to SM visible particles needs to be suppressed around 10 orders of magnitude

easiest operators:

$$\mathcal{L} \sim (\bar{L}N_1)(\bar{N}_2L) + \text{h.c.}$$

$$\mathcal{L} \sim (\bar{L}N_1)(\bar{N}_2^cL) + \text{h.c.}$$

$L\bar{L}$ pair leads to decay into $\nu\bar{\nu}$, e^+e^-

easiest operators:

$$\mathcal{L} \sim (\bar{L}N_1)(\bar{N}_2L) + \text{h.c.}$$

$$\mathcal{L} \sim (\bar{L}N_1)(\bar{N}_2^cL) + \text{h.c.}$$

⇒ operators need to be avoided!

$L\bar{L}$ pair leads to decay into $\nu\bar{\nu}$, e^+e^-

New symmetries

easiest operators:

$$\mathcal{L} \sim (\bar{L}N_1)(\bar{N}_2L) + \text{h.c.}$$

$$\mathcal{L} \sim (\bar{L}N_1)(\bar{N}_2^cL) + \text{h.c.}$$

$L\bar{L}$ pair leads to decay into $\nu\bar{\nu}$, e^+e^-

⇒ operators need to be avoided!

impose 2 U(1) symmetries:

L	N
$N_2 \rightarrow e^{i\alpha} N_2$	$N_2 \rightarrow e^{i\alpha} N_2$
$N_1 \rightarrow e^{i\alpha} N_1$	$N_1 \rightarrow e^{-i\alpha} N_1$

New symmetries

easiest operators:

$$\begin{aligned}\mathcal{L} &\sim (\bar{L}N_1)(\bar{N}_2L) + \text{h.c.} \\ \mathcal{L} &\sim (\bar{L}N_1)(\bar{N}_2^cL) + \text{h.c.}\end{aligned}$$

$L\bar{L}$ pair leads to decay into $\nu\bar{\nu}$, e^+e^-

⇒ operators need to be avoided!

impose 2 U(1) symmetries:

L	N
$N_2 \rightarrow e^{i\alpha} N_2$	$N_2 \rightarrow e^{i\alpha} N_2$
$N_1 \rightarrow e^{i\alpha} N_1$	$N_1 \rightarrow e^{-i\alpha} N_1$

$$\Rightarrow \mathcal{L}_{\text{int}} = \frac{1}{\Lambda^4} (\bar{L}\tilde{H}P_R N_2) (\bar{L}\tilde{H}P_R N_1) + \text{h.c.}$$

$$\text{with } \tilde{H} = \left(\frac{v_{\text{EW}} + h - iG^0}{\sqrt{2}}, -G^- \right)$$

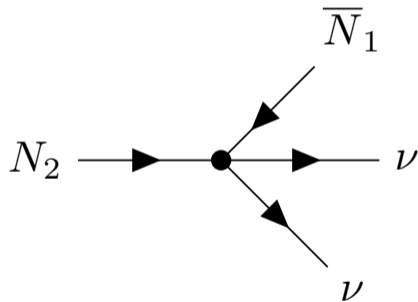
after electroweak symmetry breaking:

$$\mathcal{L}_{\text{eff}} = \frac{v_{\text{EW}}^2}{2\Lambda^4} \bar{\nu} P_R N_2 \bar{\nu} P_R N_1 + \text{h.c.}$$

Recognizing DCDM

after electroweak symmetry breaking:

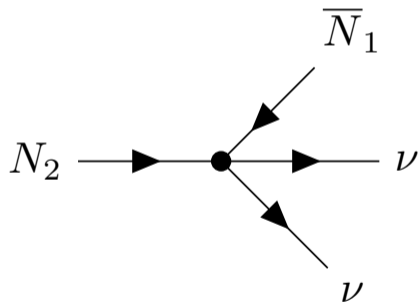
$$\mathcal{L}_{\text{eff}} = \frac{v_{\text{EW}}^2}{2\Lambda^4} \bar{\nu} P_R N_2 \bar{\nu} P_R N_1 + \text{h.c.}$$



Recognizing DCDM

after electroweak symmetry breaking:

$$\mathcal{L}_{\text{eff}} = \frac{v_{\text{EW}}^2}{2\Lambda^4} \bar{\nu} P_R N_2 \bar{\nu} P_R N_1 + \text{h.c.}$$



$$\Gamma_{N_2 \rightarrow N_1 \nu \nu} = \frac{v_{\text{EW}}^4}{1280\pi^3 \Lambda^8} (\epsilon M)^5 = \frac{1}{\tau}$$

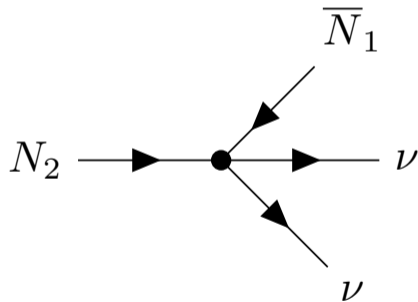
→ Λ only dependent on model parameters ϵ , τ plus the DM mass M :

$$\Lambda = \left(\frac{v_{\text{EW}}^4}{1280\pi^3} \tau (\epsilon M)^5 \right)^{1/8}$$

Recognizing DCDM

after electroweak symmetry breaking:

$$\mathcal{L}_{\text{eff}} = \frac{v_{\text{EW}}^2}{2\Lambda^4} \bar{\nu} P_R N_2 \bar{\nu} P_R N_1 + \text{h.c.}$$



$$\Gamma_{N_2 \rightarrow N_1 \nu \nu} = \frac{v_{\text{EW}}^4}{1280\pi^3 \Lambda^8} (\epsilon M)^5 = \frac{1}{\tau}$$

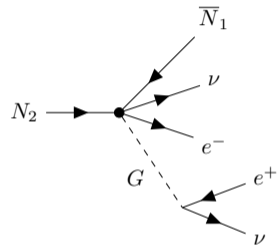
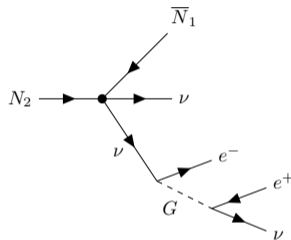
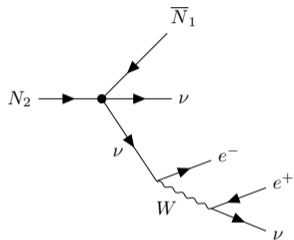
→ Λ only dependent on model parameters ϵ , τ plus the DM mass M :

$$\Lambda = \left(\frac{v_{\text{EW}}^4}{1280\pi^3} \tau (\epsilon M)^5 \right)^{1/8}$$

What we wanted!

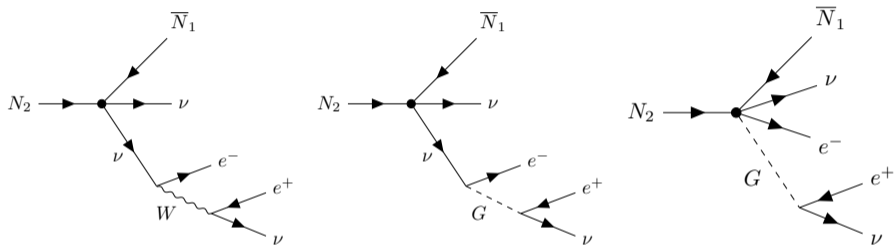
Charged particles?

- ▶ e^+e^- production possible via W and Goldstone boson



Charged particles?

- ▶ e^+e^- production possible via W and Goldstone boson

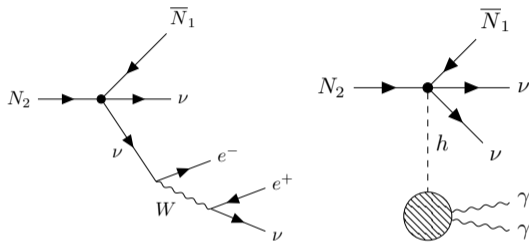


- ▶ diagrams heavily suppressed with branching ratio scaling as

$$\frac{\Gamma_{N_2 \rightarrow \bar{N}_1 \nu \nu e^+ e^-}}{\Gamma_{N_2 \rightarrow \bar{N}_1 \nu \nu}} \propto \frac{(\epsilon M)^4}{v_{EW}^4}$$

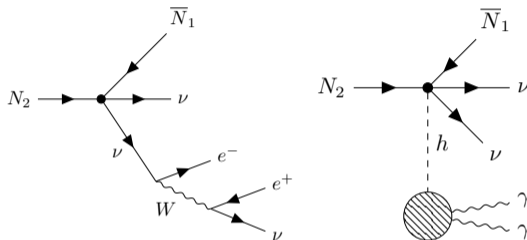
Photons?

- ▶ γ production via previous diagram or Higgs loop



Photons?

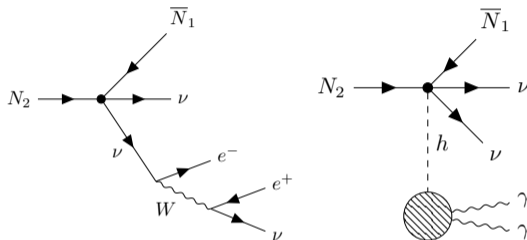
- ▶ γ production via previous diagram or Higgs loop



- ▶ similarly suppressed with branching ratio

$$\frac{\Gamma_{N_2 \rightarrow \bar{N}_1 \nu \nu \gamma \gamma}}{\Gamma_{N_2 \rightarrow \bar{N}_1 \nu \nu}} \propto \frac{(\epsilon M)^8}{m_h^4 v_{EW}^4}$$

- ▶ γ production via previous diagram or Higgs loop



- ▶ similarly suppressed with branching ratio

$$\frac{\Gamma_{N_2 \rightarrow \bar{N}_1 \nu \nu \gamma \gamma}}{\Gamma_{N_2 \rightarrow \bar{N}_1 \nu \nu}} \propto \frac{(\epsilon M)^8}{m_h^4 v_{EW}^4}$$

⇒ **We solved the challenge!**

Adjusting for a 3-body decay

$$\text{DCDM} \rightarrow \text{WDM} + \text{DR} + \text{DR}$$

- ▶ now: momentum distribution

Adjusting for a 3-body decay

$$\text{DCDM} \rightarrow \text{WDM} + \text{DR} + \text{DR}$$

- ▶ now: **momentum distribution**
- ▶ important effect given by perturbations capturing the heating of WDM instead of background evolution

Adjusting for a 3-body decay

$$\text{DCDM} \rightarrow \text{WDM} + \text{DR} + \text{DR}$$

- ▶ now: **momentum distribution**
- ▶ important effect given by perturbations capturing the heating of WDM instead of background evolution
- ▶ scale ϵ (3-body) to a new ϵ' (2-body) that produce the same perturbations

\Rightarrow small re-scaling with $\epsilon'(\epsilon) = \sqrt{\frac{13}{21}}\epsilon + \mathcal{O}(\epsilon^2) \approx 0.79\epsilon$

- I. Cosmological model
- II. Theoretical model building
- III. **New phenomenology**
 - Neutrinos
 - Production via freeze-in
 - Low-mass signatures
- IV. Outlook and Summary

A new neutrino flux

diffuse neutrino flux induced by N_2
decay:

$$\frac{d\Phi_\nu}{dE_\nu} \simeq \frac{1}{4\pi} \frac{1}{\tau M} \frac{1}{3} \frac{dN}{dE_\nu} D(\Omega)$$

D-factor:

$$D(\Omega) = \int d\Omega \int \rho(l) dl,$$

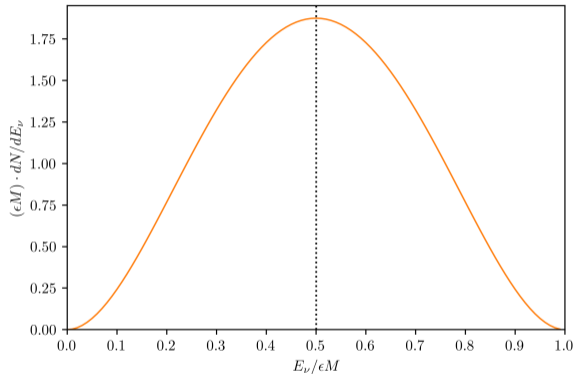
A new neutrino flux

diffuse neutrino flux induced by N_2 decay:

$$\frac{d\Phi_\nu}{dE_\nu} \simeq \frac{1}{4\pi} \frac{1}{\tau M} \frac{1}{3} \frac{dN}{dE_\nu} D(\Omega)$$

D-factor:

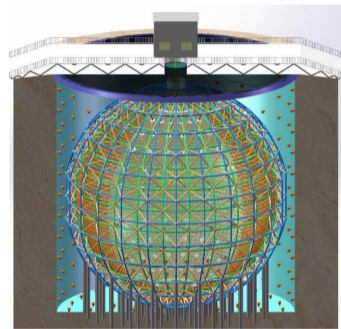
$$D(\Omega) = \int d\Omega \int \rho(l) dl,$$



neutrino spectrum $\frac{dN}{dE_\nu}$ with $\langle E_\nu \rangle = \epsilon M/2$

Neutrino experiments

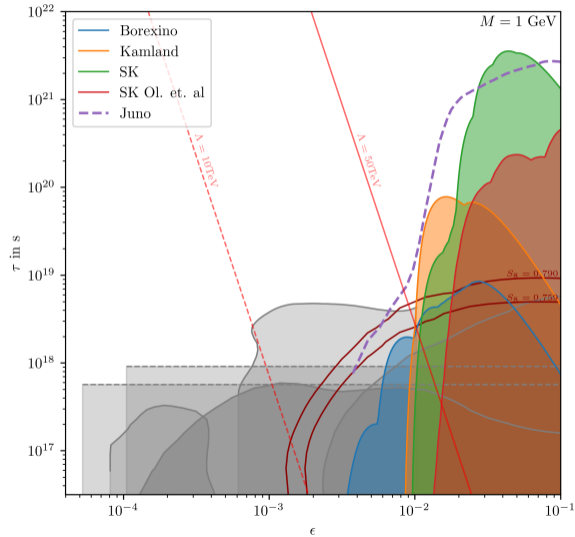
- ▶ **Borexino** (1.8 – 16.8 MeV)
[Borexino Collab., arXiv:1909.02422]
- ▶ **KamLAND** (8.3 – 30.8 MeV)
[KamLAND Collab., arXiv:2108.08527]
- ▶ **Super-Kamiokande** (9.3 – 200 MeV)
[SK Collab., arXiv:2109.11174;
Olivares-Del Campo et al., arXiv:1711.05283]
- ▶ **JUNO** (2.75 – 100 MeV)
[Akita et al., arXiv:2206.06755]



https://www.weltmaschine.de/neuigkeiten/neuigkeiten_archiv/2016/neutrinonos_auf_der_goldwaage_das_juno_experiment/

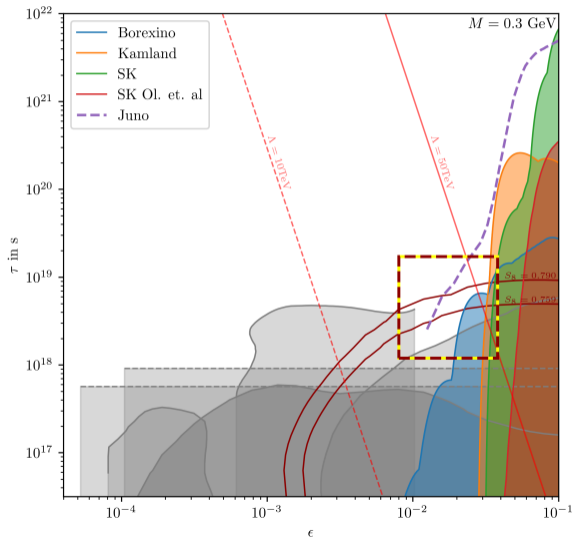
Measurement via inverse- β -decay: $\bar{\nu}_e + p \rightarrow e^+ + n$

$M = 1 \text{ GeV}$



...but opening it again!

$M = 0.3 \text{ GeV}$

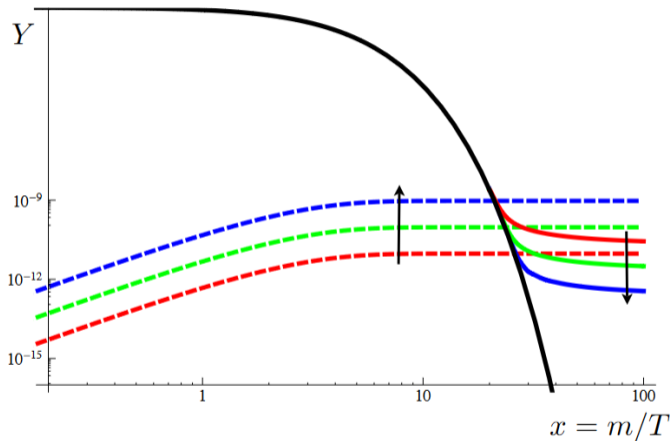


How to produce DM?

Freeze-out

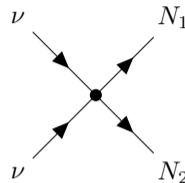
vs.

Freeze-in
for typical scales
 $\Lambda \sim \text{TeV}$



[Hall et al., arXiv:0911.1120]

- ▶ effective interaction can also produce DM after EW symmetry-breaking: $\nu\nu \rightarrow N_1 N_2$, $\bar{\nu}\bar{\nu} \rightarrow \bar{N}_1 \bar{N}_2$

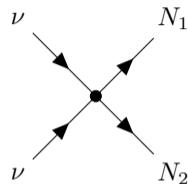


leads to 50% N_1 , 50% N_2

- ▶ effective interaction can also produce DM after EW symmetry-breaking: $\nu\nu \rightarrow N_1 N_2$, $\bar{\nu}\bar{\nu} \rightarrow \bar{N}_1 \bar{N}_2$
- ▶ freeze-in assumption: neglect back-reaction

$$\frac{dY}{dx} \propto x^4 \gamma_{N_1 N_2}$$

$$\gamma_{N_1 N_2} = \frac{v_{EW}^4 M^8}{256\pi^5 \Lambda^8} \frac{1}{x^8} \left(x^6 K_1(x)^2 + 2x^5 K_1(x) K_2(x) + (4 + x^2)x^4 K_2(x)^2 \right)$$



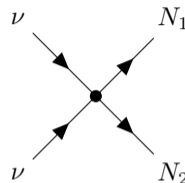
leads to 50% N_1 , 50% N_2

- ▶ effective interaction can also produce DM after EW symmetry-breaking: $\nu\nu \rightarrow N_1 N_2$, $\bar{\nu}\bar{\nu} \rightarrow \bar{N}_1 \bar{N}_2$
- ▶ freeze-in assumption: neglect back-reaction

$$\frac{dY}{dx} \propto x^4 \gamma_{N_1 N_2}$$

$$\gamma_{N_1 N_2} = \frac{v_{EW}^4 M^8}{256\pi^5 \Lambda^8} \frac{1}{x^8} \left(x^6 K_1(x)^2 + 2x^5 K_1(x) K_2(x) + (4 + x^2)x^4 K_2(x)^2 \right)$$

- ▶ depends strongly on temperature!



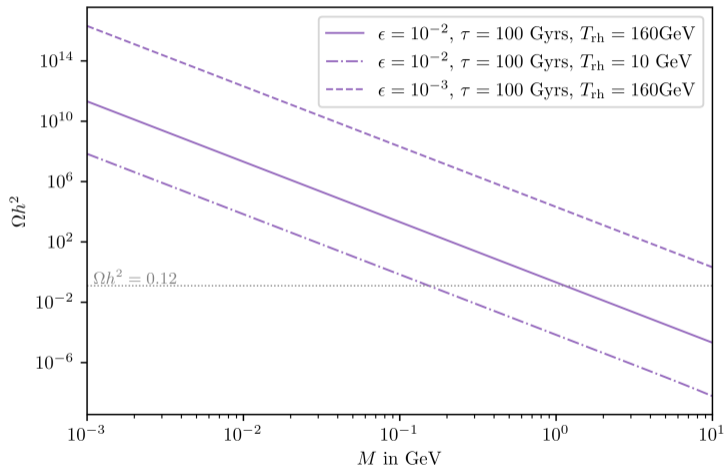
leads to 50% N_1 , 50% N_2

- ▶ restriction to broken phase with $T < 160 \text{ GeV}$

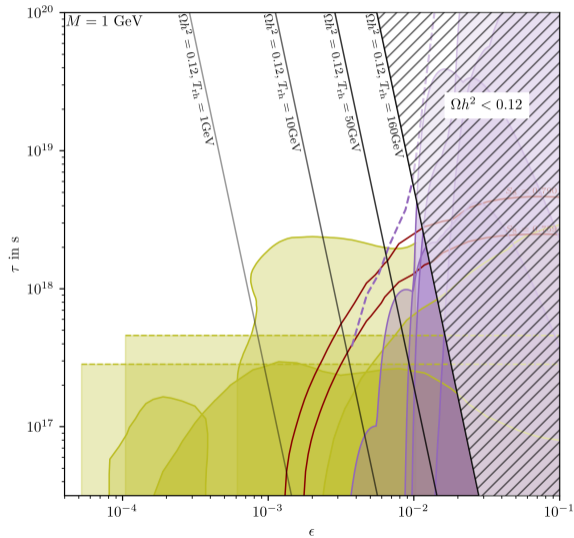
- ▶ restriction to broken phase with $T < 160 \text{ GeV}$
- ▶ vary reheating temperature T_{rh} up to this limit

...after reheating

- ▶ restriction to broken phase with $T < 160 \text{ GeV}$
- ▶ vary reheating temperature T_{rh} up to this limit

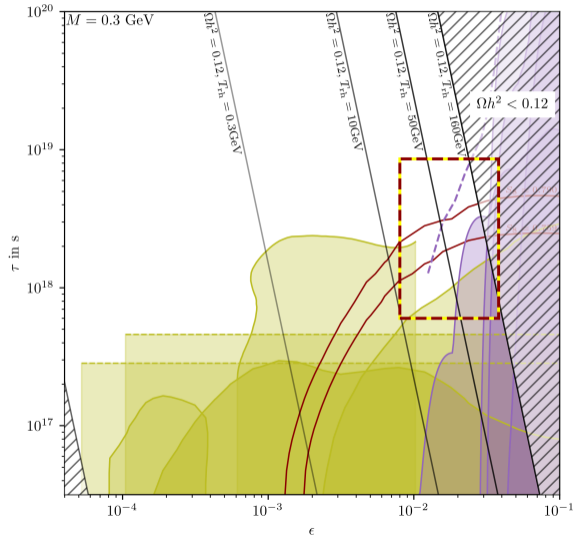


M = 1 GeV

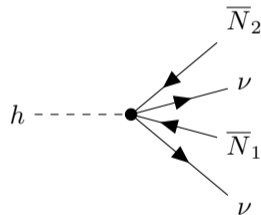


one window still open!

$M = 0.3 \text{ GeV}$

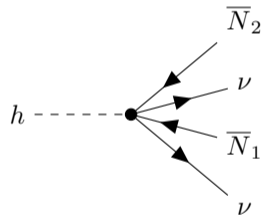


1. Invisible Higgs decay



1. Invisible Higgs decay

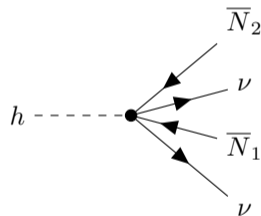
- ▶ $\Gamma_h^{\text{SM}} \simeq 3.2\text{MeV}$ with invisible branching ratio constrained to $< 12\%$



1. Invisible Higgs decay

- ▶ $\Gamma_h^{\text{SM}} \simeq 3.2\text{MeV}$ with invisible branching ratio constrained to $< 12\%$

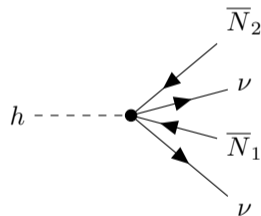
$$\begin{aligned}\Gamma_h^{\text{inv}} &= \frac{1}{4m_h} \frac{v_{\text{EW}}^2}{30\pi^5\Lambda^8} \left(\frac{m_h}{4}\right)^8 \\ &\approx 1.37 \cdot 10^{-20} \text{MeV} \left(\frac{\text{MeV}}{\epsilon M}\right)^5 \left(\frac{100 \text{ Gyrs}}{\tau}\right)\end{aligned}$$



1. Invisible Higgs decay

- ▶ $\Gamma_h^{\text{SM}} \simeq 3.2\text{MeV}$ with invisible branching ratio constrained to $< 12\%$

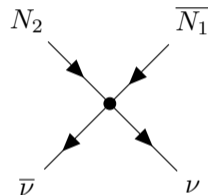
$$\begin{aligned}\Gamma_h^{\text{inv}} &= \frac{1}{4m_h} \frac{v_{\text{EW}}^2}{30\pi^5\Lambda^8} \left(\frac{m_h}{4}\right)^8 \\ &\approx 1.37 \cdot 10^{-20} \text{MeV} \left(\frac{\text{MeV}}{\epsilon M}\right)^5 \left(\frac{100 \text{ Gyrs}}{\tau}\right)\end{aligned}$$



- ▶ small effect only relevant at very small M , ϵ and τ

2. Neutrino-DM scattering

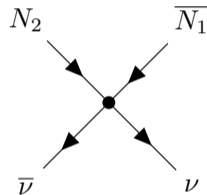
- ▶ typical constraints at **low** neutrino energies from CMB/LSS



$$\sigma_{N_1 \overline{\nu} \rightarrow \overline{N_2} \nu} = \frac{v_{EW}^4}{256\pi\Lambda^8} \frac{(s - M^2)^2}{s}$$

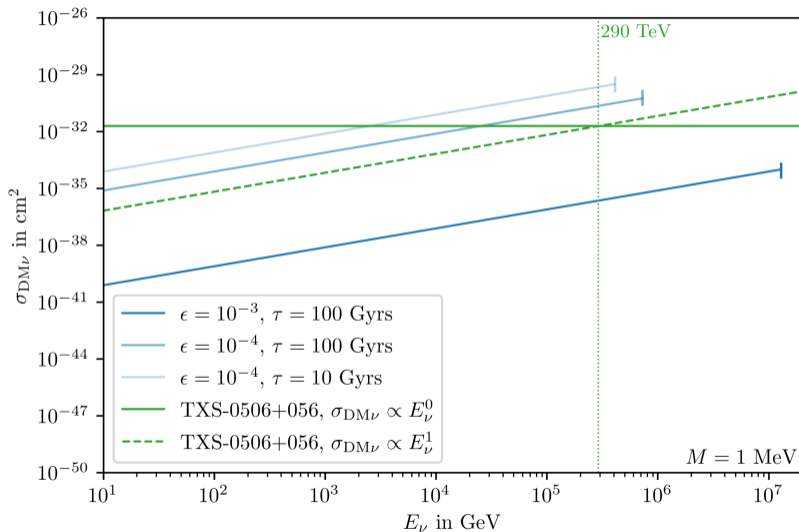
2. Neutrino-DM scattering

- ▶ typical constraints at **low** neutrino energies from CMB/LSS
- ▶ cross section boosted at **high** energies
- ▶ limits from blazar TXS-0506+056 with $E_\nu \sim 290$ TeV measured by IceCube [Ferrer, Herrera, Ibarra, arXiv:2209.06339]

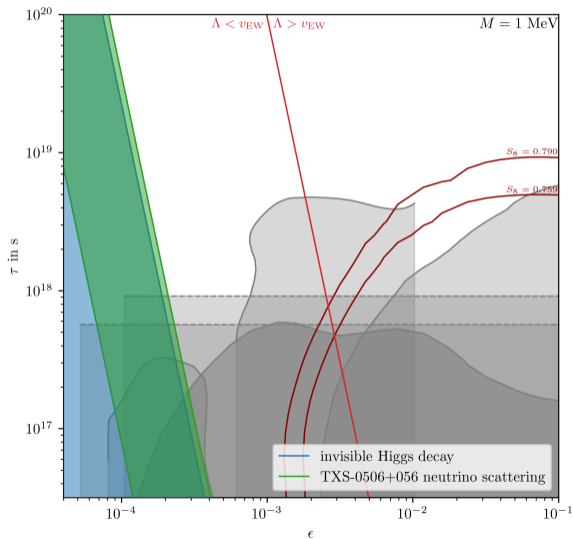


$$\sigma_{N_1 \bar{\nu} \rightarrow \bar{N}_2 \nu} = \frac{v_{EW}^4}{256\pi\Lambda^8} \frac{(s - M^2)^2}{s}$$

Help from high energy sources

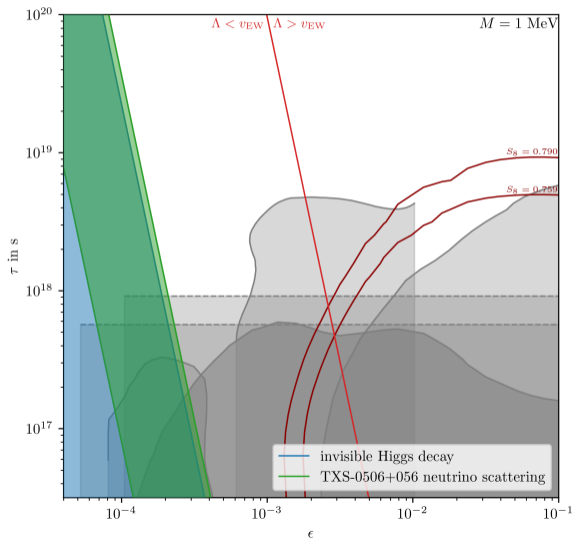


- ▶ **neutrino-DM scattering:**
constraints only shown
for $\sqrt{s} < \Lambda$
 - ▶ **invisible Higgs decay:**
constraints in $\Lambda < v_{EW}$
regime
- ⇒ limit of EFT description



Low Mass Constraints

- ▶ **neutrino-DM scattering:**
constraints only shown
for $\sqrt{s} < \Lambda$
 - ▶ **invisible Higgs decay:**
constraints in $\Lambda < v_{EW}$
regime
- ⇒ limit of EFT description



- I. Cosmological model
- II. Theoretical model building
- III. New phenomenology
 - Neutrinos
 - Production via freeze-in
 - Low-mass signatures
- IV. Outlook and Summary

What comes next?

One step further: going to a **UV complete theory**

What comes next?

One step further: going to a **UV complete theory**

1. New and/or improved phenomenology? (low mass regime, reheating temperature, collider)

What comes next?

One step further: going to a **UV complete theory**

1. New and/or improved phenomenology? (low mass regime, reheating temperature, collider)
2. Connection to neutrino masses via e.g. Seesaw mechanism and heavy neutral leptons that carry lepton number?

What comes next?

One step further: going to a **UV complete theory**

1. New and/or improved phenomenology? (low mass regime, reheating temperature, collider)
2. Connection to neutrino masses via e.g. Seesaw mechanism and heavy neutral leptons that carry lepton number?
3. Natural explanation for the mass splitting between N_1 and N_2 ?

- ▶ Found minimal and effective realization of decaying DM that opens up new phenomenology
- ▶ Complementary constraints from cosmology, neutrino experiments, and freeze-in production
- ▶ Window in parameter space where all constraints and lower S_8 are satisfied for $M \lesssim 1 \text{ GeV}$
- ▶ Possible future testability: JUNO, Euclid (?)

- ▶ Found minimal and effective realization of decaying DM that opens up new phenomenology
- ▶ Complementary constraints from cosmology, neutrino experiments, and freeze-in production
- ▶ Window in parameter space where all constraints and lower S_8 are satisfied for $M \lesssim 1 \text{ GeV}$
- ▶ Possible future testability: JUNO, Euclid (?)

Thank you for your attention!