## *GRAVITATIONAL WAVES FROM DARK CONFINEMENT WITH HOLOGRAPHY*

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## MOTIVATION & MAIN POINTS



- Gravitational Waves from FOPT in Strongly coupled QFTs ("resembling known theories") is likely not visible in the future due to lack of supercooling. (In contrast to previous attempts in RS using AdS/CFT)
- The AdS/CFT Correspondence for FOPTs in strongly coupled QFTs described by HP PT provide an bound on the amount of supercooling.

#### OUTLINE



Figure from Baggioli 1908.02667

#### Interlude to AdS/CFT

Our Holographic model of interest for SU(N) pYM.

Gravitational waves from an FOPT

#### **Outlook & Future interests**

## ADS/CFT CONJECTURE! (9711200)

#### In this work we wish to study SU(N) pYM

$\mathcal{N}$ = 4, SUSY SU(N <sub>c</sub> ) YM = IIB String Theory on AdS <sub>5</sub> ×	$S_5$
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C(Q)FT in d-Dimensions	String/Gravity in d+1-Dimensions
$Z_{Q\mathcal{FT}}$	$Z_{String/Gravity}$
$\mathcal{W}_{\mathcal{QFT}}$	$-S_{OS}$
Temperature T	Hawking Temperature BH
Entropy S	Hawking Entropy
Large N Limit	Classical GR in $AdS_5 \times S_5$



Two immediate questions emerge:

- SU(N) pYM Theory is not a CFT?
- Furthermore it is also not Supersymmetric.

## NON CONFORMALITY/NO SUSY?



## *PROS/CONS* ??

#### TOP DOWN HOLOGRAPHY

- Dictionary known precisely, one knows exactly what one talks about.
- Dual theories stemming from string theory somewhat confident in no pathologies
- Limited set of theories with known explicit duals (SUSY, CFT)
- Enormous large amounts of fields and content in each theory to account for

#### BOTTOM UP HOLOGRAPHY

- Large amount of flexibility as one imposes the sought dynamics one wish to study
- In certain cases the models can be relatively simple yet possess rich dynamics
  - This approach has been useful for understanding various dynamics and features in Hydrodynamics, QCD and Condensed Matter

In principle there could be unwanted/hidden pathologies in the constructed theory

# INPROVED HOLOGRAPHIC

0707.1324, 0707.1349, 0812.0792, 0903.2859, 0906.1980

> This model is not provided from a string theory. Hence we possess a large N gauge theory at strong coupling with certain features.

## IHQCD VACUUM THEORY

 $ds^2 = b^2(r)(dr^2 - dt^2 + dx^m dx_m)$ ,  $\Phi(r)$ ,

$$\mathcal{S}_{5} = -M_{p}^{3}N_{c}^{2}\int d^{5}x\sqrt{g}\left(R - \frac{4}{3}(\partial\Phi)^{2} + V(\Phi)\right) + 2M_{p}^{3}\int_{\partial\mathcal{M}} d^{4}x\sqrt{h}\mathcal{K} + CT$$

AdS<sub>5</sub> Einstein-dilaton gravity

4D Strongly coupled QFT

Radial 5-D coordinate r

• RG scale

• Scalar field  $\lambda = e^{\phi}$ 

• T'hooft coupling  $\lambda_{\rm t} = N_c g_{YM}^2$ 

## IHQCD VACUUM IN THE UV REGIME

• Vacuum equations of motion

• 
$$6\frac{\dot{b^2}}{b^2} + 3\frac{\ddot{b}}{b^2} = b^2 V$$
,  $6\frac{\dot{b^2}}{b^2} - 3\frac{\ddot{b}}{b^2} = \frac{4}{3}\dot{\Phi^2}$ ,

- 4-D energy scale given by  $E = E_0 b(r)$ .
- Furthermore imposing the identification  $\lambda = \kappa \lambda_t$
- Defining the holographic beta function as
- $\beta(\lambda) = \frac{d\lambda}{d\log E} = \lambda \frac{\dot{\Phi}}{\dot{A}}$ ,  $A(r) = \log b(r)$ ,
- Introducing scalar variables  $X(\lambda) = \frac{\beta(\lambda)}{3\lambda}$  . ("Gluon Condensate")

$$\lambda \frac{dX}{d\lambda} = -\frac{4}{3} (1 - X^2) \left( 1 + \frac{3}{8X} \lambda \frac{d \log V}{d\lambda} \right) ,$$

RG scale 
$$r = const.$$
 Mink<sub>4</sub>  
 $r = 0$  Horizon IR  
Figure from Mateos 0709.1523

MAIN BOTTOM UP ASPECT!

$$V = \frac{12}{l^2} \left( 1 + V_0 \lambda + V_1 \lambda^{\frac{4}{3}} \left( \log \left[ 1 + V_2 \lambda^{\frac{4}{3}} + V_3 \lambda^2 \right]^{\frac{1}{2}} \right) \right)$$

- Fix available parameters  $V_0, V_1, V_2, V_3$
- V<sub>0</sub>, V<sub>2</sub> fixed to reproduce YM β function at 2 loops.
- V<sub>1</sub>, V<sub>3</sub> fitted against SU(3) lattice data in the IR.

Here we display the direct comparison using the set of potential parameters provided by original authors. Data from



SOLUTIONS AT  $T \neq 0$ 

AdS Schwarzchild BH solution Deconfined Phase

$$ds^{2} = b^{2}(r) \left( \frac{dr^{2}}{f(r)} - f(r)dt^{2} + dx^{m}dx_{m} \right)$$

$$\Phi = \Phi(r), \qquad r \in (0, r_h), \qquad f(r_h) = 0$$

Thermal Graviton Gas Solution Confined Phase

$$ds^{2} = b_{0}^{2}(r)(dr^{2} - dt^{2} + dx^{m}dx_{m})$$
  

$$\Phi = \Phi_{0}(r), \qquad r \in (0, \infty)$$





#### THERMODYNAMICS

- Temperature, Time periodicity  $\tau \rightarrow \tau + \frac{1}{\tau}$ .
- $S = \frac{Area}{4G_5} =$
- $4\pi M_p^3 N_c^2 V_3 b(r_h)^3$

• 
$$\mathcal{F} = \frac{\beta}{V_3} \left( \mathcal{S}_{\mathcal{dC}} - \mathcal{S}_{\mathcal{CN}} \right)$$

$$T_h \equiv \frac{|\dot{f}(r_h)|}{4\pi} = T$$

#### CENTRAL IDEA/WHY USING ADS/CFT!!

 $\mathcal{Z}_{\phi}[\phi_0]_{Gravity} = \mathcal{Z}_{\mathcal{O}}[\phi_{\mathcal{O}}]_{CFT}$ 



Figure from 2308.02159.



*EFFECTIVE ACTION FOR TUNNELING I* 

- Interpolate between BBH and SBH
- Violate the condition  $T_h \neq T$
- 1. BH not in thermal eq.
- 2. Conical singularity
- Regularize with spherical cap



$$V_{\text{eff}}(\lambda_h, T) = \mathcal{F}(\lambda_h) - 4\pi M_p^3 N_c^2 b(\lambda_h)^3 \left(1 - \frac{T_h}{T}\right)$$



## GRAVITATIONAL WAVES FROM FOPT

- Release of vacuum energy  $\Delta v$ ,
- Drives bubble expansion and fluid motion which sources Gravitational waves







### GRAVITATIONAL WAVES FROM FOPT II

#### How to compute GW signal?

- Hydrodynamical lattice simulations of a scalar field coupled to the plasma.
- Imposing Local Thermal Equilibrium simplifies things

#### PT is characterized by few parameters:

- PT strength (Latent Heat)  $\alpha \sim \frac{\Delta^2 vac}{2}$
- Bubble wall velocity  $v_w$   $\Omega_{rad}$
- $\bullet$  Bubble nucleation rate  $\beta_*$
- PT temperature T\*



#### TATIONAL WAVES FI $\begin{aligned} f &= f^3, & f < f_p \\ f &= f^{-4}, & f > f_p \end{aligned}$ $F_{GW,0} = (3.57 \pm 0.05) \cdot 10^{-5} \left(\frac{100}{g_*}\right)^{\frac{1}{3}}$ Numerical Prefactor $\sim 10^{-2}$ $\bullet \quad \frac{d\Omega_{GW,0}}{d\ln f} = 0.687F_{GW,0}K^{\frac{3}{2}}\left(H(T_*)R(T_*)\right)^2\overline{\Omega_{GW}}C\left(\frac{f}{f_{n,0}}\right)$ $K = \frac{\kappa(\alpha)\alpha}{1+\alpha}$ $f_{p,0} \simeq 26 \left(\frac{1}{R_* H(T_*)}\right) \left(\frac{z_p}{10}\right) \left(\frac{T_*}{100 \text{ GeV}}\right) \left(\frac{g_*}{100}\right)^{\frac{1}{6}} \mu \text{Hz}$ $\kappa(\alpha) = \frac{\alpha}{0.73 + 0.083\sqrt{\alpha} + \alpha}$ Source duration time of acoustic production $\frac{\beta_*}{H_*}^{-2}$

One Major Drawback here: The simulation at which these semi-analyic templates have only been calculated at weak coupling/ weak transition strengths. Hence spectral shapes/dependencies may change at strong coupling

*EFFECTIVE ACTION FOR TUNNELING II* 

- Kinetic term normalization:  $c \frac{N_c^2}{16\pi^2} (\vec{\nabla} \lambda_h)^2$
- Effective action for  $\mathcal{O}(3)$  tunneling configurations (Thermal fluctuation!)

• 
$$S_{\mathcal{B}} = \frac{4\pi}{T} \int dr r^2 \left[ c \frac{N_c^2}{16\pi^2} (\partial_r \lambda_h(r))^2 + V_{\text{eff}}(\lambda_h(r), T) \right]$$

• Bubble Nucleation Rate

• 
$$\Gamma \approx T^4 \left(\frac{S_B}{2\pi}\right)^{3/2} e^{-S_B}$$
, Nucleation  $\Gamma \approx H^4$ .

- $T_n \approx T_p \approx 0.99T_c$
- PT stength (energy released)

• 
$$\alpha = \frac{4}{3} \frac{\Delta \theta}{\Delta w} = \frac{1}{3} \frac{\Delta \rho - 3\Delta p}{\Delta w} \sim 0.34$$

- Inverse PT rate
- $\frac{\beta}{H} = T\left(\frac{dS_{\mathcal{B}}}{dT}\right) \sim 10^5$



#### *GW SPECTRA SU(3) YANG-MILLS*

- Bubble wall velocity realistically?
- Kinetic term coefficient c=1 ??
- These curves should be thought of as "naive" estimates.
- Other methods have been attempted/applied like PLM, Matrix Models, and Thin Wall estimations.
- Nothing beyond the Bounce Solution has been attempted here so far for these purposes to my knowledge

#### KINETIC TERM COEFFICIENT Preliminary $S_{\mathcal{B}} = \frac{4\pi}{T} \int dr \ r^2 \left[ c \frac{N_c^2}{16\pi^2} (\partial_r \lambda_h(r))^2 + V_{\text{eff}}(\lambda_h(r), T) \right]$



We want to evaluate the action on a bubble configuration.

$$S \supset \int d^5 x \dots (\vec{\nabla} \lambda_h)^2 + \dots (\vec{\nabla} \lambda_h)^4 + \dots$$

Is c  $\sim$  1 a reasonable approximation??







- The PT is faster hence further GW suppression
- A Thin wall approximation might be valid for these transitions

Dual Quantity: Wave function renormalization factor and its evolution w.r.t the t'hooft coupling.



Kinetic term pole at 0.3 seems to be a numerical issue. Matching at a lower cut off  $\lambda_0$  in value removes this feature.



## CONCLUSIONS & OUTLOOK

- AdS/CFT provides a upper bound on supercooling from its geometrical construction of the theories of concern.
- We should not expect to observe GWs from strongly coupled FOPT in theories with limited supercooling.
- Still much to understand in regards to frameworks methods and their various applicability.

- Exploring options for calculating bubble wall velocities
- Inclusion of flavor/axion contributions.
- Phase separated phenomena dynamical instabilities.
- Hydrodynamical analysis of plasma and its properties.
- Black Hole perturbation theory for fluctuating determinant
- What about considering Top Down examples?