GRAVITATIONAL WAVES FROM DARK CONFINEMENT WITH HOLOGRAPHY

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MOTIVATION & MAIN POINTS

- Gravitational Waves from FOPT in Strongly coupled QFTs ("resembling known theories") is likely not visible in the future due to lack of supercooling. (In contrast to previous attempts in RS using AdS/CFT)
- The AdS/CFT Correspondence for FOPTs in strongly coupled QFTs described by HP PT provide an bound on the amount of supercooling.

what strongly coupled you want to compute QFT **DUAL** what weakly coupled you know classical gravity to compute

Figure from Baggioli 1908.02667

OUTLINE Interlude to AdS/CFT

Our Holographic model of interest for SU(N) pYM.

> Gravitational waves from an FOPT

Outlook & Future interests

ADS/CFT CONJECTURE! (9711200)

In this work we wish to study SU(N) pYM

Two immediate questions emerge:

- SU(N) pYM Theory is not a CFT?
- Furthermore it is also not Supersymmetric.

NON CONFORMALITY/ NO SUSY?

PROS/CONS ??

TOP DOWN HOLOGRAPHY

- Dictionary known precisely, one knows exactly what one talks about.
- Dual theories stemming from string theory somewhat confident in no pathologies
- Limited set of theories with known explicit duals (SUSY, CFT)
- Enormous large amounts of fields and content in each theory to account for

BOTTOM UP HOLOGRAPHY

- Large amount of flexibility as one imposes the sought dynamics one wish to study
- In certain cases the models can be relatively simple yet possess rich dynamics
	- This approach has been useful for understanding various dynamics and features in Hydrodynamics, QCD and Condensed Matter

In principle there could be unwanted/hidden pathologies in the constructed theory

IMPROVED HOLOGRAPHIC

QCD<sup>0707.1324, 0707.1349, 0812.0792,
0903.2859, 0906.1980</sup> 0903.2859, 0906.1980

> This model is not provided from a string theory. Hence we possess a large N gauge theory at strong coupling with certain features.

IHQCD VACUUM THEORY

 $ds^2 = b^2(r)(dr^2 - dt^2 + dx^m dx_m), \qquad \Phi(r),$

$$
\mathcal{S}_5 = -M_p^3 N_c^2 \int d^5 x \sqrt{g} \left(R - \frac{4}{3} (\partial \Phi)^2 + V(\Phi) \right) + 2M_p^3 \int_{\partial \mathcal{M}} d^4 x \sqrt{h} \mathcal{K} + C T
$$

• AdS_5 Einstein-dilaton gravity

• 4D Strongly coupled QFT

• Radial 5-D coordinate r

• RG scale

• Scalar field $\lambda = e^{\Phi}$

• T'hooft coupling $\lambda_t = N_c g_{YM}^2$

IHQCD VACUUM IN THE UV REGIME

• Vacuum equations of motion

•
$$
6\frac{b^2}{b^2} + 3\frac{b}{b^2} = b^2V
$$
, $6\frac{b^2}{b^2} - 3\frac{b}{b^2} = \frac{4}{3}\dot{\phi}^2$,

- 4-D energy scale given by $E = E_0 b(r)$.
- Furthermore imposing the identification $\lambda = \kappa \lambda_t$
- Defining the holographic beta function as
- $\beta(\lambda) = \frac{d\lambda}{d\log \lambda}$ $rac{d\lambda}{d \log E} = \lambda \frac{\dot{\Phi}}{\dot{A}}$ $\frac{\Phi}{\dot{A}}$, $A(r) = \log b(r)$,
- Introducing scalar variables $X(\lambda) = \frac{\beta(\lambda)}{3\lambda}$. ("Gluon" Condensate")

$$
\lambda \frac{dX}{d\lambda} = -\frac{4}{3} (1-X^2) \left(1 + \frac{3}{8X} \lambda \frac{d \log V}{d\lambda} \right) ,
$$

Figure from Mateos 0709.1523

MAIN BOTTOM UP ASPECT!

$$
V = \frac{12}{1^2} \left(1 + V_0 \lambda + V_1 \lambda^{\frac{4}{3}} \left(\log \left[1 + V_2 \lambda^{\frac{4}{3}} + V_3 \lambda^2 \right]^{\frac{1}{2}} \right) \right)
$$

- Fix available parameters V_0 , V_1 , V_2 , V_3
- V_0 , V_2 fixed to reproduce YM β function at 2 loops.
- V_1 , V_3 fitted against SU(3) lattice data in the IR.

Here we display the direct comparison using the set of potential parameters provided by original authors. Data from

 $0.\ell$

 0.5

 0.4

0.3

 0.2

 0.1

 Ω

 1.2

 1.4

 16

 18

 2.0

 2.2

 2.4

 $SOLUTIONSATT \neq 0$

AdS Schwarzchild BH solution Deconfined Phase

$$
ds^2 = b^2(r) \left(\frac{dr^2}{f(r)} - f(r)dt^2 + dx^m dx_m \right)
$$

$$
\Phi = \Phi(r), \qquad r \in (0, r_h), \qquad f(r_h) = 0
$$

Thermal Graviton Gas Solution Confined Phase

$$
ds^{2} = b_{0}^{2}(r)(dr^{2} - dt^{2} + dx^{m}dx_{m})
$$

\n
$$
\Phi = \Phi_{0}(r), \qquad r \in (0, \infty)
$$

THERMODYNAMICS

- Temperature, Time periodicity $\tau \to \tau + \frac{1}{\tau}$ $\frac{1}{T}$.
- $S =$ Area $4G_5$ =
- $4\pi M_p^3 N_c^2 V_3 b(r_h)^3$

$$
\cdot \quad \mathcal{F} = \frac{\beta}{V_3} \big(\mathcal{S}_{dc} - \mathcal{S}_{cn.} \big)
$$

•
$$
T_h \equiv \frac{|f(r_h)|}{4\pi} = T
$$

CENTRAL IDEA/WHY USING ADS/CFT!!

 $Z_{\phi}[\phi_0]_{gravity} = Z_{\mathcal{O}}[\phi_0]_{CFT}$

Figure from 2308.02159.

EFFECTIVE ACTION FOR TUNNELING I

- Interpolate between BBH and SBH
- Violate the condition $T_h \neq T$
- 1. BH not in thermal eq.
- 2. Conical singularity
- Regularize with spherical cap

$$
V_{\text{eff}}(\lambda_h, T) = \mathcal{F}(\lambda_h) - 4\pi M_p^3 N_c^2 b(\lambda_h)^3 \left(1 - \frac{T_h}{T}\right)
$$

GRAVITATIONAL WAVES FROM FOPT

- Release of vacuum energy ΔV ,
- Drives bubble expansion and fluid motion which sources Gravitational waves

GRAVITATIONAL WAVES FROM FOPT II

How to compute GW signal?

- Hydrodynamical lattice simulations of a scalar field coupled to the plasma.
- Imposing Local Thermal Equilibrium simplifies things

PT is characterized by few parameters:

- PT strength (Latent Heat) $\alpha \sim$ Ω_{vac}
- Bubble wall velocity V_W $\Omega_{\rm rad}$
- Bubble nucleation rate β_*
- PT temperature T_*

GRATIONAL WAVES FI $f = f^3$, $f < f_p$ 1 **Numerical** 100 3 $f = f^{-4}$, $f > f_p$ $F_{GW,0} = (3.57 \pm 0.05) \cdot 10^{-5}$ Prefactor $\sim 10^{-2}$ g_{*} • $\frac{d\Omega_{GW,0}}{d \ln f} = 0.687 F_{GW,0} K^{\frac{3}{2}} (H(T_*)R(T_*)^2) \frac{d\Omega_{GW} C}{d\Omega_{GW}} \left(\frac{f}{f_{no}}\right)$ κ (α) α 1 $K=$ 1 z_p T_* g_{*} 6 $1 + \alpha$ $f_{p,0} \simeq 26$ μHz $R_*H(T_*$ 10 100 GeV 100 α $\kappa(\alpha) =$ $0.73 + 0.083\sqrt{\alpha} + \alpha$ Source duration time of acoustic production β_* −2 H_*

One Major Drawback here: The simulation at which these semi-analyic templates have only been calculated at weak coupling/ weak transition strengths. Hence spectral shapes/dependencies may change at strong coupling *EFFECTIVE ACTION FOR TUNNELING II*

- Kinetic term normalization: $c \frac{N_c^2}{16\pi}$ $\frac{N_C}{16\pi^2}(\nabla\lambda_h$ 2
- Effective action for $O(3)$ tunneling configurations (Thermal fluctuation!)

•
$$
S_B = \frac{4\pi}{T} \int dr r^2 \left[c \frac{N_c^2}{16\pi^2} \left(\partial_r \lambda_h(r) \right)^2 + V_{\text{eff}}(\lambda_h(r), T) \right]
$$

• Bubble Nucleation Rate

•
$$
\Gamma \approx T^4 \left(\frac{S_B}{2\pi}\right)^{3/2} e^{-S_B}
$$
, Nucleation $\Gamma \approx H^4$.

- $T_n \approx T_p \approx 0.99 T_c$
- PT stength (energy released)

•
$$
\alpha = \frac{4}{3} \frac{\Delta \theta}{\Delta w} = \frac{1}{3} \frac{\Delta \rho - 3\Delta p}{\Delta w} \sim 0.34
$$

• Inverse PT rate

$$
\bullet \ \ \frac{\beta}{H} = T\left(\frac{d\mathcal{S_B}}{dT}\right) \sim 10^5
$$

GW SPECTRA SU(3) YANG-MILLS

- Bubble wall velocity realistically?
- Kinetic term coefficient c=1 ??
- These curves should be thought of as "naive" estimates.
- Other methods have been attempted/applied like PLM, Matrix Models, and Thin Wall estimations.
- Nothing beyond the Bounce Solution has been attempted here so far for these purposes to my knowledge

KINETIC TERM COEFFICIENT $\mathcal{S_{B}}=% \begin{bmatrix} \omega_{\mathrm{d}}\mathcal{S}_{\mathrm{d}} & \frac{\omega_{\mathrm{d}}}{2} & \frac{\$ 4π \overline{T} $\int dr r^2 \Big| c$ N_c^2 $16π²$ $\partial_r \lambda_h(r)$ 2 + $V_{\text{eff}}(\lambda_h(r), T)$ **Preliminary**

We want to evaluate the action on a bubble configuration.

$$
S \supset \int d^5x \ldots (\vec{\nabla} \lambda_h)^2 + \ldots (\vec{\nabla} \lambda_h)^4 + \ldots
$$

Is c ∼ 1 a reasonable approximation??

 $\mathcal{S_{B}}=% \begin{bmatrix} \omega_{\mathrm{d}}\sqrt{2\pi\epsilon_{\mathrm{d}}}\ \frac{\omega_{\mathrm{d}}}{2} & \omega_{\mathrm{d}}\sqrt{2\pi\epsilon_{\mathrm{$

 4π

 \overline{T}

 $\int dr r^2 \Big| c(\lambda_h$

Kinetic term pole at 0.3 seems to be a numerical issue. Matching at a lower cut off λ_0 in value removes this feature.

The kinetic term is small. We shall expect 2 features.

 $\partial_r \lambda_h(r)$

2

+ $V_{\text{eff}}(\lambda_h(r), T)$

The PT is faster hence further GW suppression

 N_c^2

 $16π²$

• A Thin wall approximation might be valid for these transitions

Dual Quantity: Wave function renormalization factor and its evolution w.r.t the t'hooft coupling.

CONCLUSIONS & OUTLOOK

- AdS/CFT provides a upper bound on supercooling from its geometrical construction of the theories of concern.
- We should not expect to observe GWs from strongly coupled FOPT in theories with limited supercooling.
- Still much to understand in regards to frameworks methods and their various applicability.
- Exploring options for calculating bubble wall velocities
- Inclusion of flavor/axion contributions.
- Phase separated phenomena dynamical instabilities.
- Hydrodynamical analysis of plasma and its properties.
- Black Hole perturbation theory for fluctuating determinant
- What about considering Top Down examples?