

Domain wall networks and their cosmological signatures

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Effective Theories
for non-perturbative
Physics
5-17 August 2024
14-8-24



Preliminary evidence : Römische Theater



Overview

- Left-right symmetric model
 - A minimal esthetic extension of SM
- Generic domain walls in left-right symmetric models
 - Leptogenesis
 - Gravitational waves
- Supersymmetric version and $SO(10)$ embedding
 - strategies for DW removal
 - Relating CP violation to EDM
- Possibility of a two peak Gravitational Wave signal

1 Genesis of baryogenesis

Just a little asymmetry

Annus Mirabilis of Cosmology

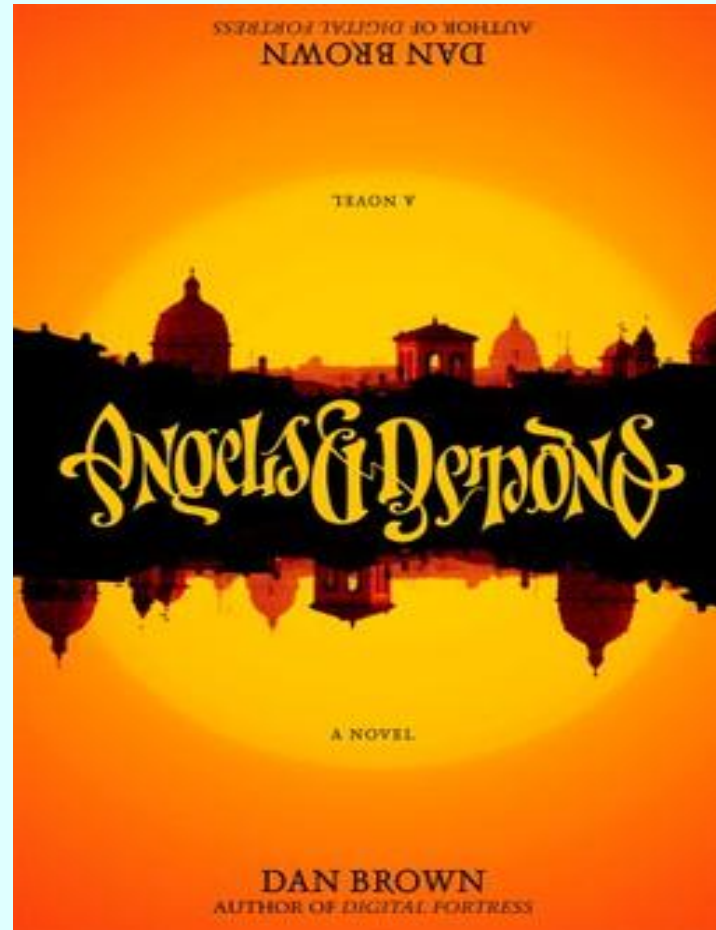
Two discoveries of 1964

→ CP violation in $K^0 - \bar{K}^0$

→ CMB !!!

Weinberg comment in Brandeis Lectures 1964.

Sakharov model elucidating the criteria 1967



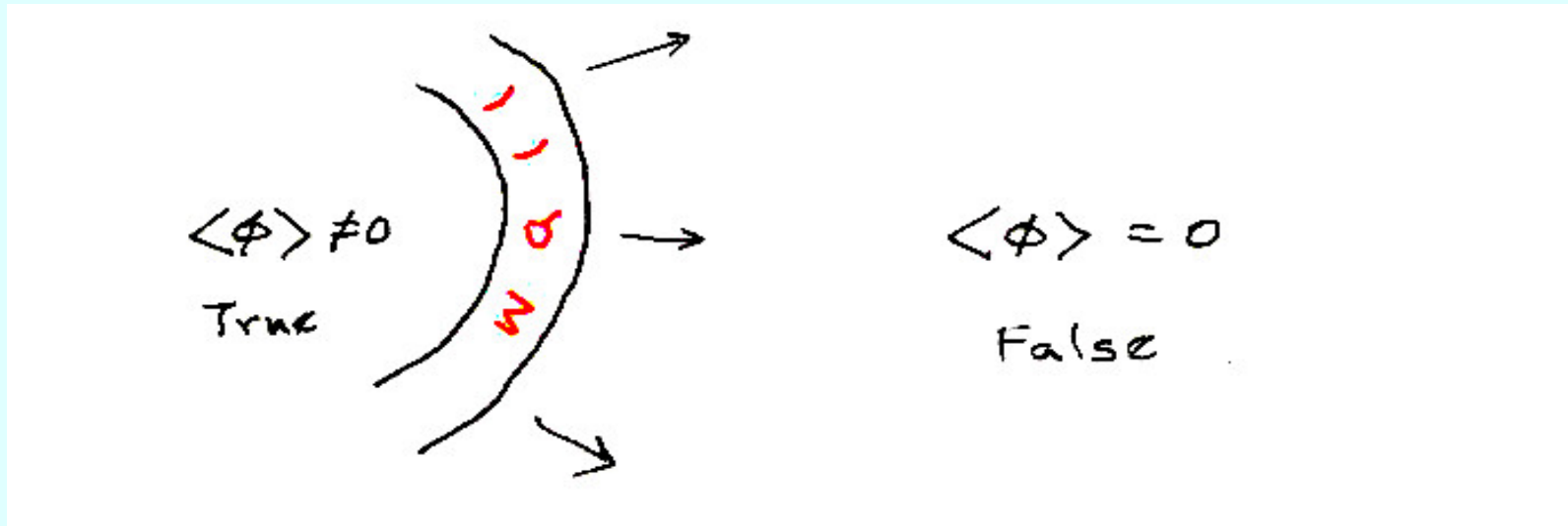
1.1 GUT baryogenesis

- Out of equilibrium decays of GUT scale leptoquarks
- The Particle Physics rates and expansion rate of the Universe compete : **out-of-equilibrium decays**

$$\Gamma_x \cong \alpha_x m_x^2 / T; \quad H \cong g_*^{1/2} T^2 / M_{\text{Pl}}$$

1.2 Electroweak baryogenesis (Low scale)

- Expansion rate H too slow at electroweak scale
 - need another source of out of equilibrium conditions
 - moving phase boundaries of a **First Order Phase Transition (FOPT)**



- But **First order phase transition (FOPT) in SM** requires Higgs mass to be $\lesssim 90\text{GeV}$

2 Baryogenesis from leptogenesis

Replay the Baryon asymmetry recipes for Leptons

- Thermal solution : Out of equilibrium decays of Majorana neutrinos – high scale
- Phase transition : **Moving phase boundaries of Left-Right symmetric model** – can be low scale

2.1 Difficulties of High scale leptogenesis

out-of-equilibrium decays of heavy Majorana neutrinos
[see eg. Buchmüller, Di Bari and Plümacher (2004)]

- Getting Majorana neutrinos to be in equilibrium

$$M_N \gtrsim O(10^9)\text{GeV} \left(\frac{2.5 \times 10^{-3}}{Y_N} \right) \left(\frac{0.05\text{eV}}{m_\nu} \right)$$

- Have sufficiently large CP violation – assuming see-saw mechanism and 3 generations

$$|\varepsilon_{CP}| \lesssim 10^{-7} \left(\frac{M_1}{10^9\text{GeV}} \right) \left(\frac{m_3}{0.05\text{eV}} \right)$$

- Preventing washout of the produced asymmetry by the same Majorana neutrino mediated processes

3 Left-right symmetric model

(Mohapatra and Senjanovic 1970's; predecessor Pati and Salam 1974)

$$\begin{array}{l}
 \\
 \\
 \begin{bmatrix} \nu_L \\ e_L^- \end{bmatrix} \\
 \begin{bmatrix} \nu_R \\ e_R^- \end{bmatrix}
 \end{array}
 \begin{array}{ll}
 T_L^3 & T_R^3 \\
 +\frac{1}{2} & 0 \\
 -\frac{1}{2} & 0 \\
 0 & +\frac{1}{2} \\
 0 & -\frac{1}{2}
 \end{array}
 \begin{array}{ll}
 \frac{1}{2}X & Q \\
 -\frac{1}{2} & 0 \\
 -\frac{1}{2} & -1 \\
 -\frac{1}{2} & 0 \\
 -\frac{1}{2} & -1
 \end{array}$$

$$\begin{array}{l}
 \\
 \\
 \begin{bmatrix} u_L \\ d_L \end{bmatrix} \\
 \begin{bmatrix} u_R \\ d_R \end{bmatrix}
 \end{array}
 \begin{array}{ll}
 T_L^3 & T_R^3 \\
 +\frac{1}{2} & 0 \\
 -\frac{1}{2} & 0 \\
 0 & +\frac{1}{2} \\
 0 & -\frac{1}{2}
 \end{array}
 \begin{array}{ll}
 \frac{1}{2}X & Q \\
 +\frac{1}{6} & +\frac{2}{3} \\
 +\frac{1}{6} & -\frac{1}{3} \\
 +\frac{1}{6} & +\frac{2}{3} \\
 +\frac{1}{6} & -\frac{1}{3}
 \end{array}$$

3.0.1 Higgs sector – suitable for neutrino see-saw

$$\Phi = (1, 2, 2, 0)$$
$$\Delta_L = (1, 3, 1, 2), \quad \Delta_R = (1, 1, 3, 2)$$

In the notation

$$\Delta_L \equiv \Delta_L^i T_L^i = \begin{pmatrix} \Delta^+ & \Delta^{++} \\ \Delta^0 & \Delta^- \end{pmatrix}$$

Choice of vev

$$\langle \Delta_L \rangle = \begin{pmatrix} 0 & 0 \\ l & 0 \end{pmatrix}, \quad \langle \Delta_R \rangle = \begin{pmatrix} 0 & 0 \\ r & 0 \end{pmatrix},$$

$$\Phi = \begin{pmatrix} K & 0 \\ 0 & k' \end{pmatrix}$$

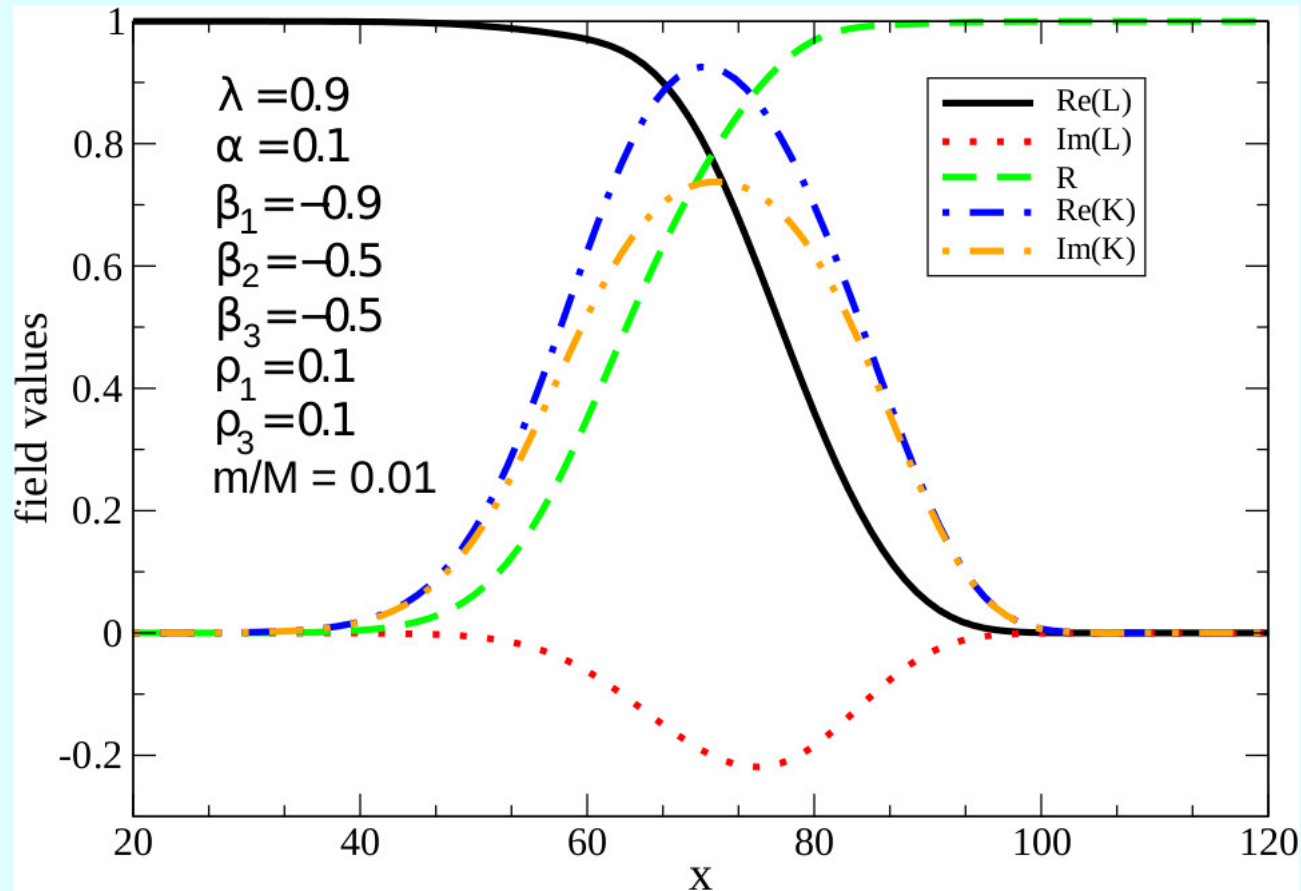
- Introduced new species ν_R --> as a partner to e_R^-
- New gauge symmetry $SU(2)_R$
- Need a new hypercharge X --> turns out to be exactly $B - L$
- In praise of $B - L$...
 - the only conserved charge of SM which is not gauged!
 - Hereby it gains the status of being gauged
- Emerges naturally in $SO(10)$ unification

4 L -genesis in left-right symmetric world

- Assume flip symmetry $SU(2)_L \leftrightarrow SU(2)_R$ (possible in $SO(10)$)
- $B - L$ automatically a local symmetry – ensures we start with a clean slate $B - L = 0$ at the Big Bang
- Two kinds of vacua – $SU(2)_L$ breaking or $SU(2)_R$ breaking
 - one desirable, the other accidental
- Big Bang universe has horizons
 - patchwork of both kinds of domains
- So we have
 - L number violation for Majorana neutrinos
 - Out of equilibrium wall motion (bring us to SM)
 - CP violation – transient values in the core of the DW

Moving phase boundaries at $SU(2)_R$ breaking

[Sarkar, Abhishek and UAY]



How can we verify this? **The parable of the cow and the grass.**
Electroweak baryogenesis models rely on

- a cosmological phase transition
- Movement of bubble walls
- CP violation within the width of the wall

The CP phase is transient : both time and space dependent.

Further,

- Thermal leptogenesis with high scale has difficulties
- The EW Bgenesis scenarios can be adapted to BSM for leptogenesis

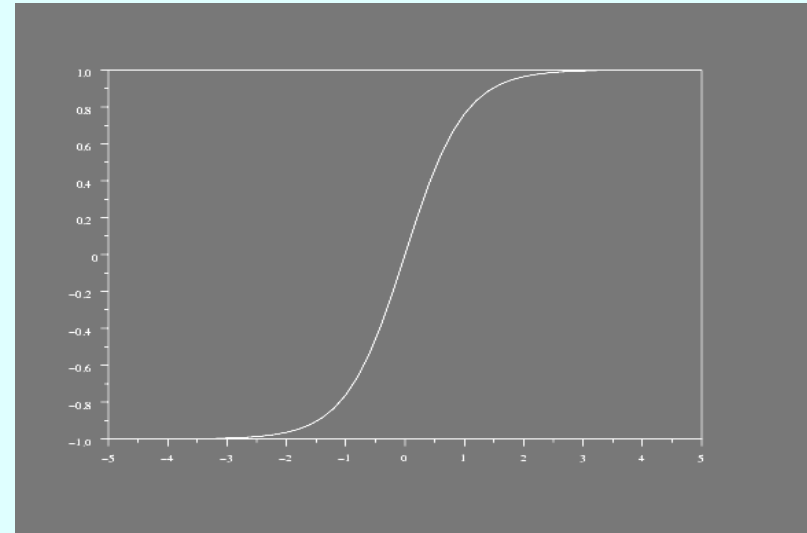
How to relate transient CP violation to low energy physics.

Key issues :

- Bubble walls are solitonic solutions
 - Space dependent CP phase is also solitonic

Machine errors in end values can produce a completely different curve.

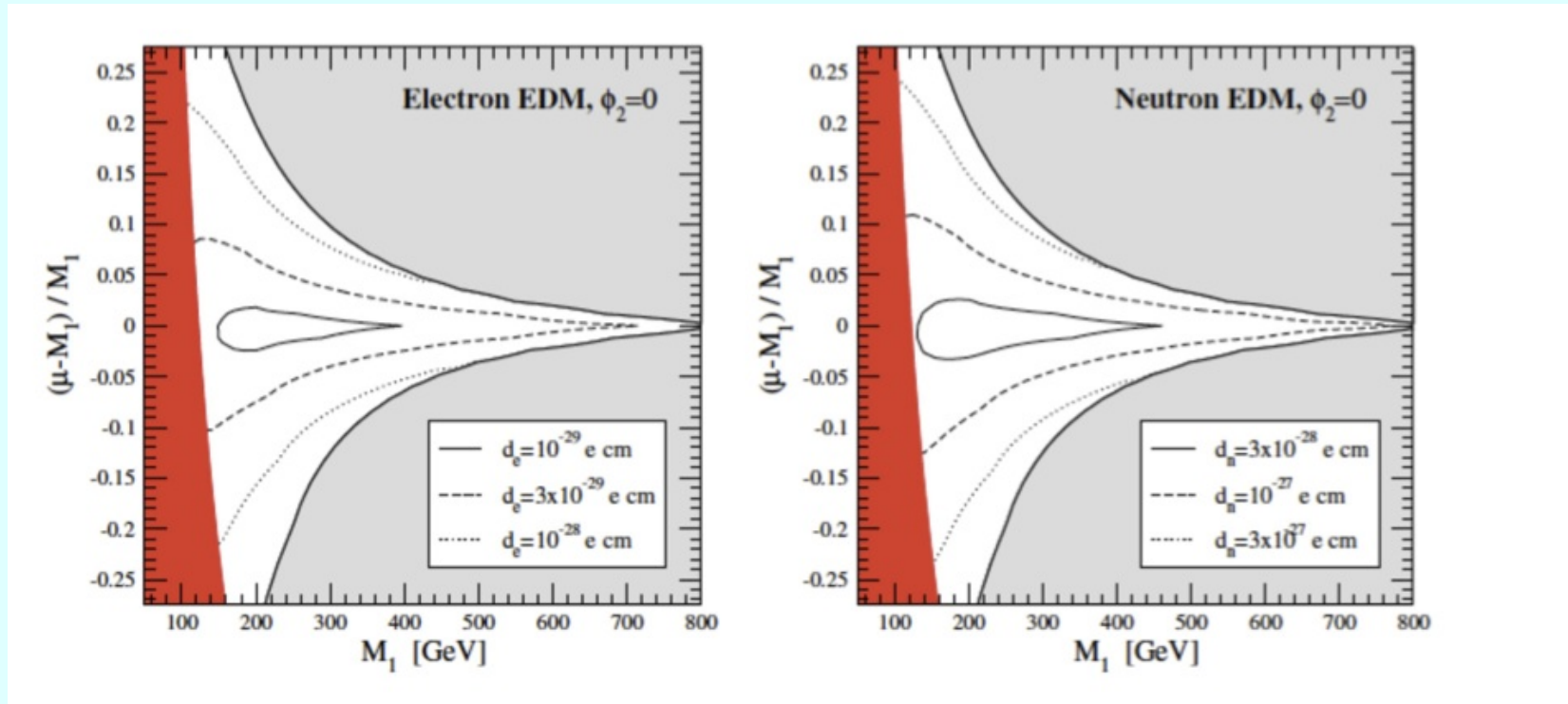
- Difficulty relating values at a **finite boundary** to interior values.



- Bubble walls occur at finite temperature. Need to relate finite temperature parameter values to observable zero temperature values.

4.1 Relating CP violation to EDM

In MSSM++ models,



Morrissey and Ramsey-Musolf (2012)

$$d_f \cong \sin \delta_{CP} \left(\frac{m_f}{\text{MeV}} \right) \left(\frac{1 \text{ TeV}}{M} \right)^2 \times 10^{-26} \text{ e cm}$$

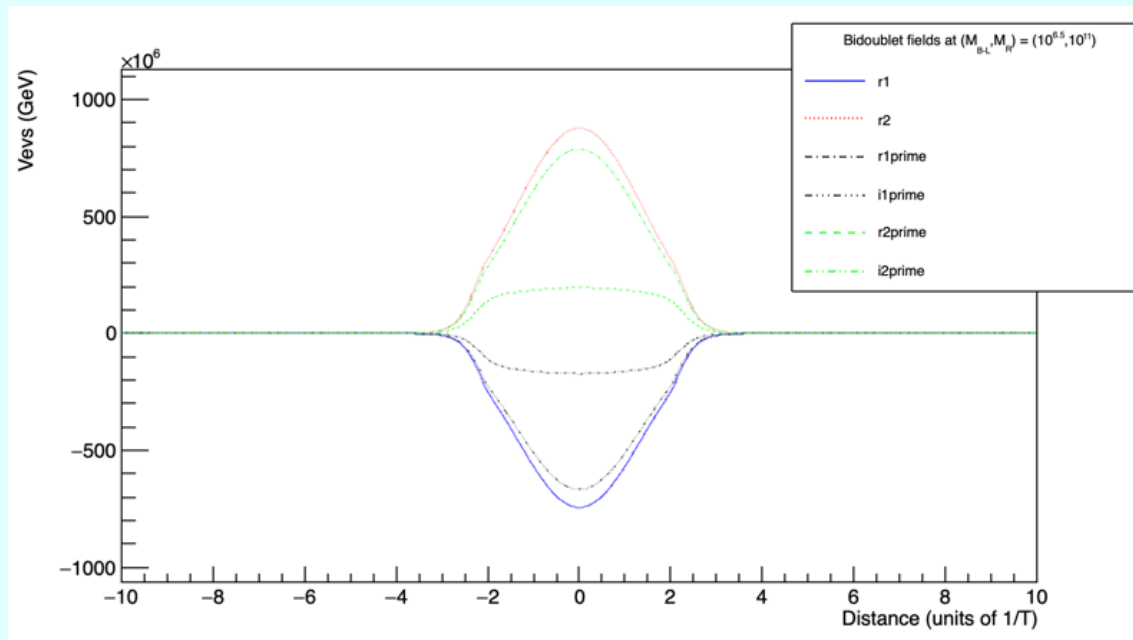
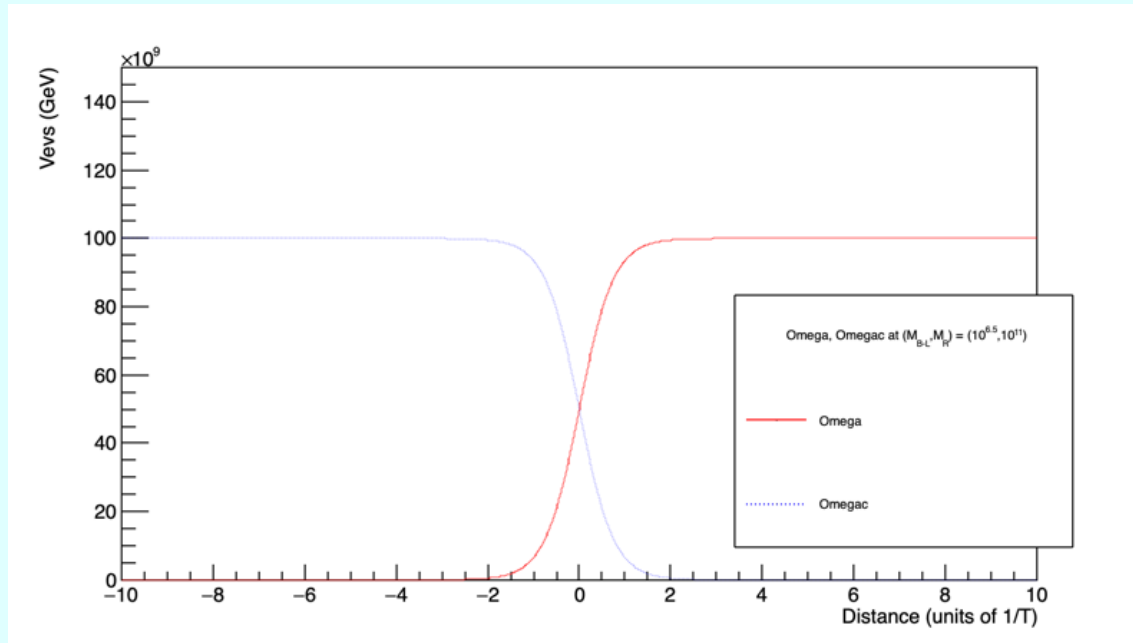
5 A renormalisable SUSY LR model

[Benakli, Aulakh, Senjanovic (1997)]

- Higgs content : superfields Bidoublet Φ , Triplets $\Delta_L, \Delta_R, \Delta_L^c, \Delta_R^c$ with $B - L = \pm 2$, and new Triplets Ω_L, Ω_R with $B - L = 0$.
 - Renormalisable model
- Two stage gauge symm breaking : $M_R \sim SU(2)_R \rightarrow U(1)_R$
and $M_{B-L} \sim U(1)_R \otimes U(1)_{B-L} \rightarrow U(1)_Y$
 - Avoid new mass scale by imposing an R symmetry $\rightarrow M_{B-L}^2 \approx M_{EW} M_R$

$$\begin{aligned}
 W = & m_\Delta (\text{Tr} \Delta \bar{\Delta} + \text{Tr} \Delta_c \bar{\Delta}_c) + \frac{m_\Omega}{2} (\text{Tr} \Omega^2 + \text{Tr} \Omega_c^2) + \mu_{ij} \text{Tr} T_2 \Phi_i^T T_2 \Phi_j \\
 & + a (\text{Tr} \Delta \Omega \bar{\Delta} + \text{Tr} \Delta_c \Omega_c \bar{\Delta}_c) \\
 & + \alpha_{ij} (\text{Tr} \Omega \Phi_i T_2 \Phi_j^T T_2 + \text{Tr} \Omega_c \Phi_i^T T_2 \Phi_j T_2)
 \end{aligned}$$

Piyali Banerjee
and UAY JHEP
(2021)



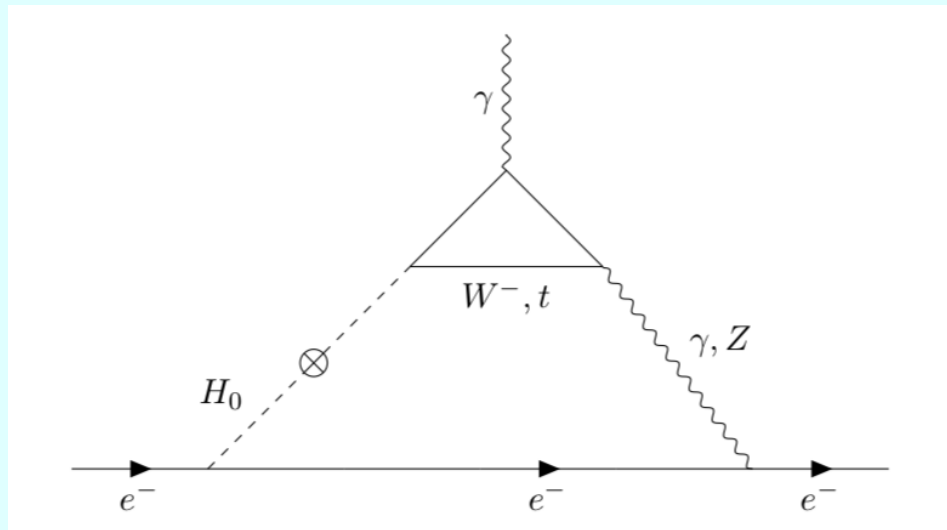
5.1 Corroborating with EDM

- Assume all scalar vevs entering the DW are taken to be corrected by temperature correction $O(g^2 T^2)$
- In a simple bidouble Higgs model, a 1-loop formula for EDM is

$$\frac{d_e}{e} \sim \frac{\alpha m_e}{4\pi M_h^2} \sin \delta$$

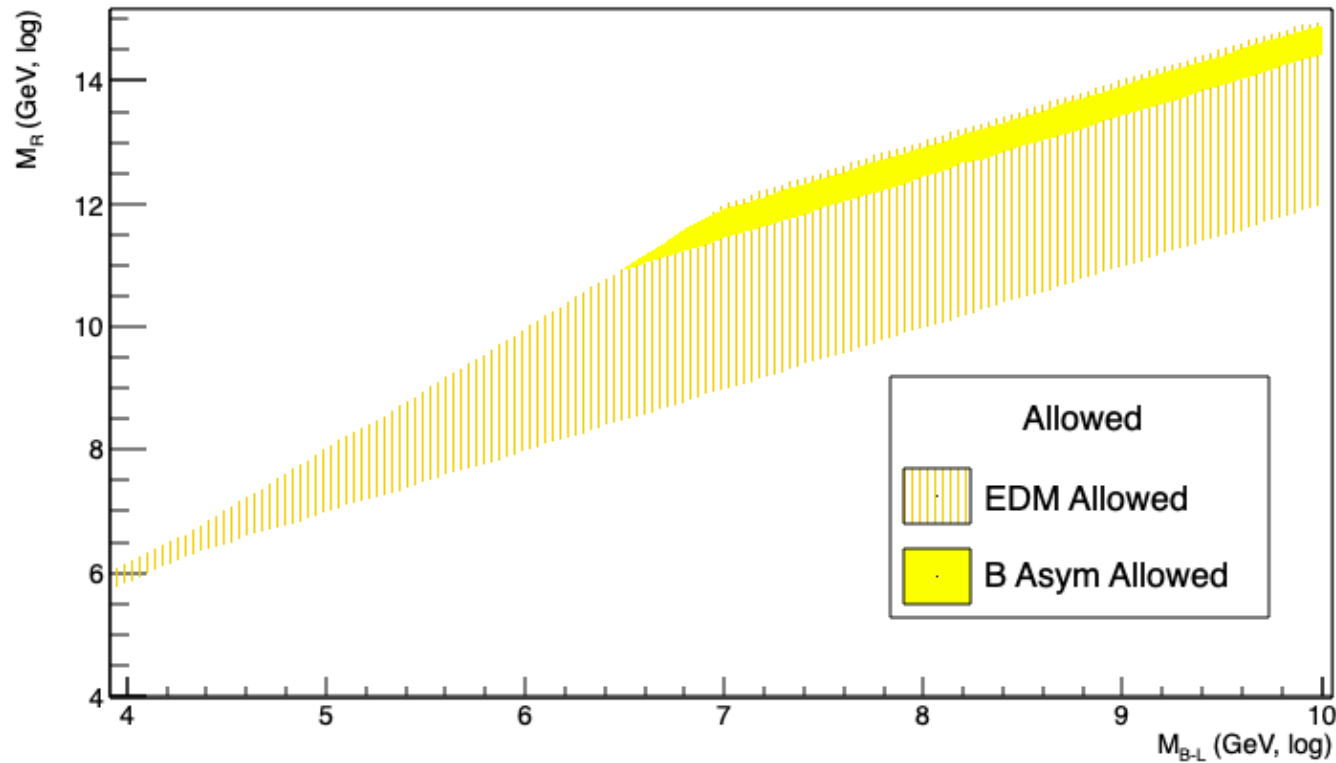
- For large values of M_{B-L} and M_R two loop effects arising from the neutral scalars dominate the one loop

$$(d_e / e)|_{\text{two loop}} = \frac{G_F m_e \alpha \sin \delta}{\pi^3 \sqrt{2}} (f_{W, HY\gamma} (M_W^2 / M_h^2) + f_{W, HZ\gamma} (M_W^2 / M_h^2) + f_{t, HY\gamma} (M_t^2 / M_h^2) + f_{t, HZ\gamma} (M_t^2 / M_h^2)).$$



and four other such diagrams

Thus we obtain constraints on the mass scales M_R and M_{B-L} .



- Interesting lesson : the R -symmetry compatible formula of Benaqli, Aulakh and Senjanovic is in tension.
- Can be repaired by including the $m_\Omega \Omega \Omega$ term in superpotential

Left – Right models a tale of two phase transitions

6 Phase transition gravitational waves

mostly Caprini, Durrer, Servant, Binétruy, Hindmarsh ... 2009 – 2019;

Kosowsky and Turner(1993); Huber and Konstandin(2008); Weir (2016)

1. Bubble collisions

$$h^2 \Omega_{\text{env}}(f) = 1.67 \times 10^{-5} \left(\frac{H_*}{\beta} \right)^2 \left(\frac{K_c \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{g_*} \right)^{\frac{1}{3}} \left(\frac{0.11 v_w^3}{0.42 + v_w^2} \right) \frac{3.8 (f / f_{\text{env}})^{2.8}}{1 + 2.8 (f / f_{\text{env}})^{3.8}}$$

where $H_* = H(T_n)$

2. Sound waves

Pressure waves created in the plasma by movement of DW

$$h^2 \Omega_{\text{sw}}(f) = 2.65 \times 10^{-6} \left(\frac{H_*}{\beta} \right) \left(\frac{K_{\text{sw}} \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{g_*} \right)^{\frac{1}{3}} v_w \left(\frac{f}{f_{\text{sw}}} \right)^3 \left(\frac{7}{4 + 3 (f / f_{\text{sw}})^2} \right)^{7/2}$$

with peak frequency

$$f_{\text{sw}} = \frac{1.9 \times 10^{-5}}{v_w} \left(\frac{\beta}{H_*} \right) \left(\frac{T_*}{100 \text{ GeV}} \right) \left(\frac{g_*}{100} \right)^{\frac{1}{6}} \text{ Hz}$$

3. MHD turbulence

$$h^2 \Omega_{\text{turb}}(f) = \left(\frac{H_*}{\beta} \right) \left(\frac{K_{\text{turb}} \alpha}{1 + \alpha} \right)^{\frac{3}{2}} \left(\frac{100}{g_*} \right)^{1/3} v_w \frac{3.35 \times 10^{-4} (f / f_{\text{turb}})^3}{[1 + (f / f_{\text{turb}})]^{\frac{11}{3}} (1 + 8 \pi f / h_*)}$$

6.1 The two phase transitions

Zafri A. Borboruah and UAY ArXiv 2022; PRD2024

Tree level Higgs potential for Δ_R vev r

$$V_0(r) \approx \frac{\rho_1}{4} (r^2 - \eta^2)^2$$

and similar for l .

The effective potential to be used

$$V_{\text{eff}}(r, T) = V_0(r) + V_{\text{CW}}(r) + V_{\text{FT}}(r, T) + V_{\text{D}}(r, T)$$

including Coleman-Weinberg, Finite temperature and daisy digrams.

Finally, with both L and R contributions :

$$V_{\text{eff}}^{\text{total}}(r, l, T) = V_{\text{eff}}(r, T) + V_{\text{eff}}(l, T)$$

Two types of phase transitions

I. Kibble mechanism - "Causal horizon limited SOPT"

Characterised by Ginzberg temperature

$$\xi_G^3 \Delta_c V_{eff} = T_G$$

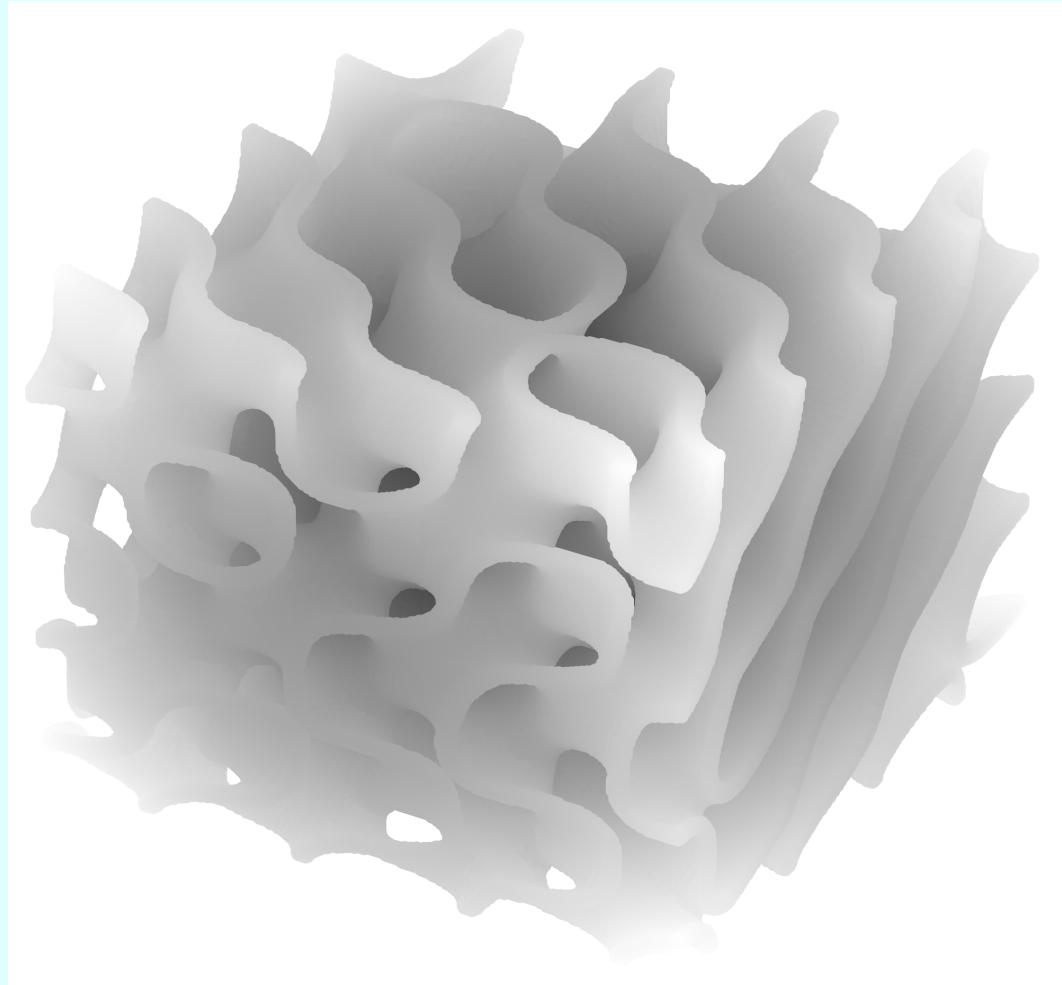
and its length scale ρ_1 a quartic coupling, η a vev

$$\xi_G \approx \frac{1}{2 \rho_1 \eta}$$

Instead of diverging indefinitely, the putative SOPT has a scale

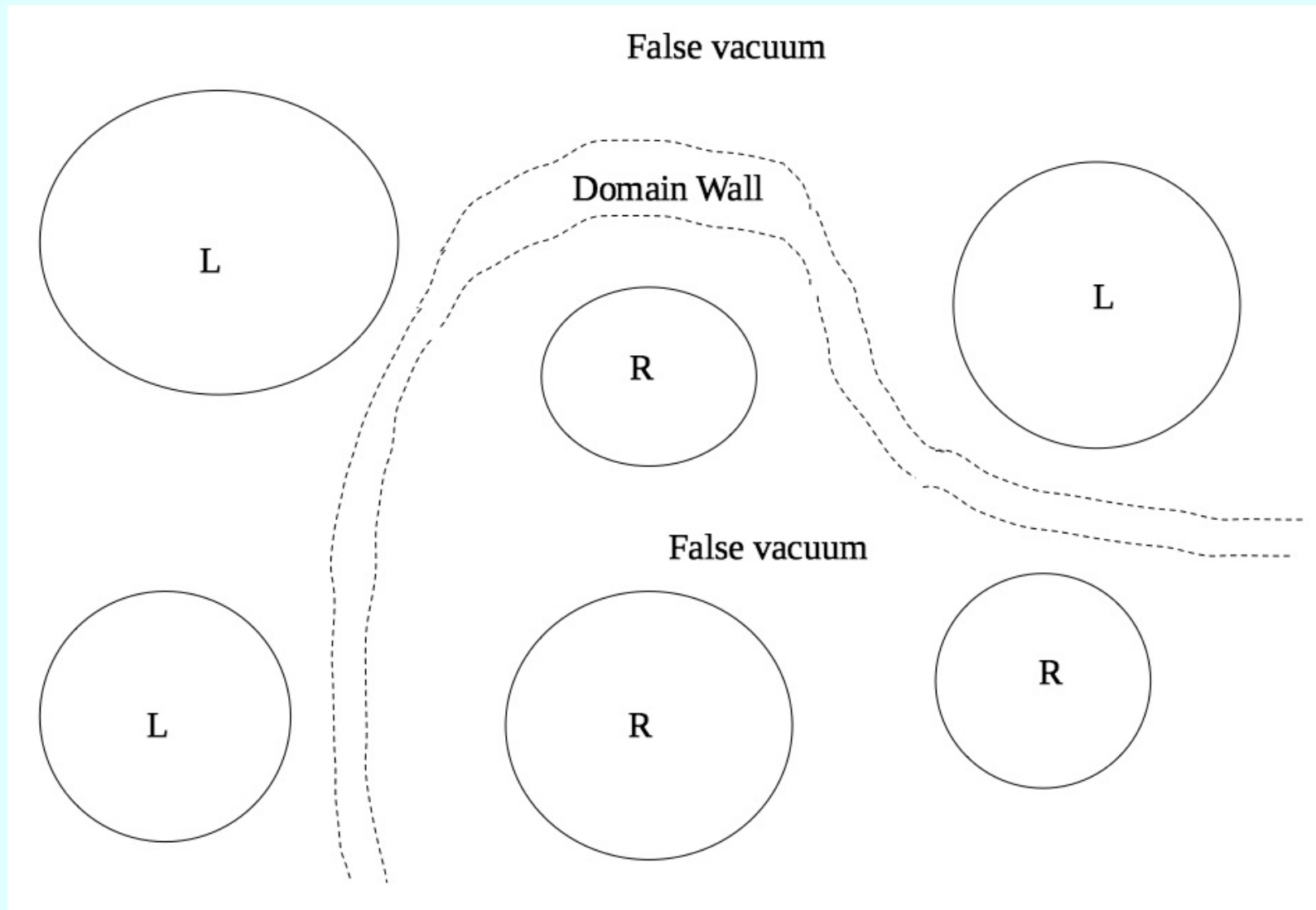
$$\xi_{causal} = \left(\frac{M_{pl}}{\sqrt{\mathcal{N}} m_r^2 T_c^2} \right)^{1/3}$$

We treat these walls by direct simulation "crumbling walls" by introducing a bias term to break L-R



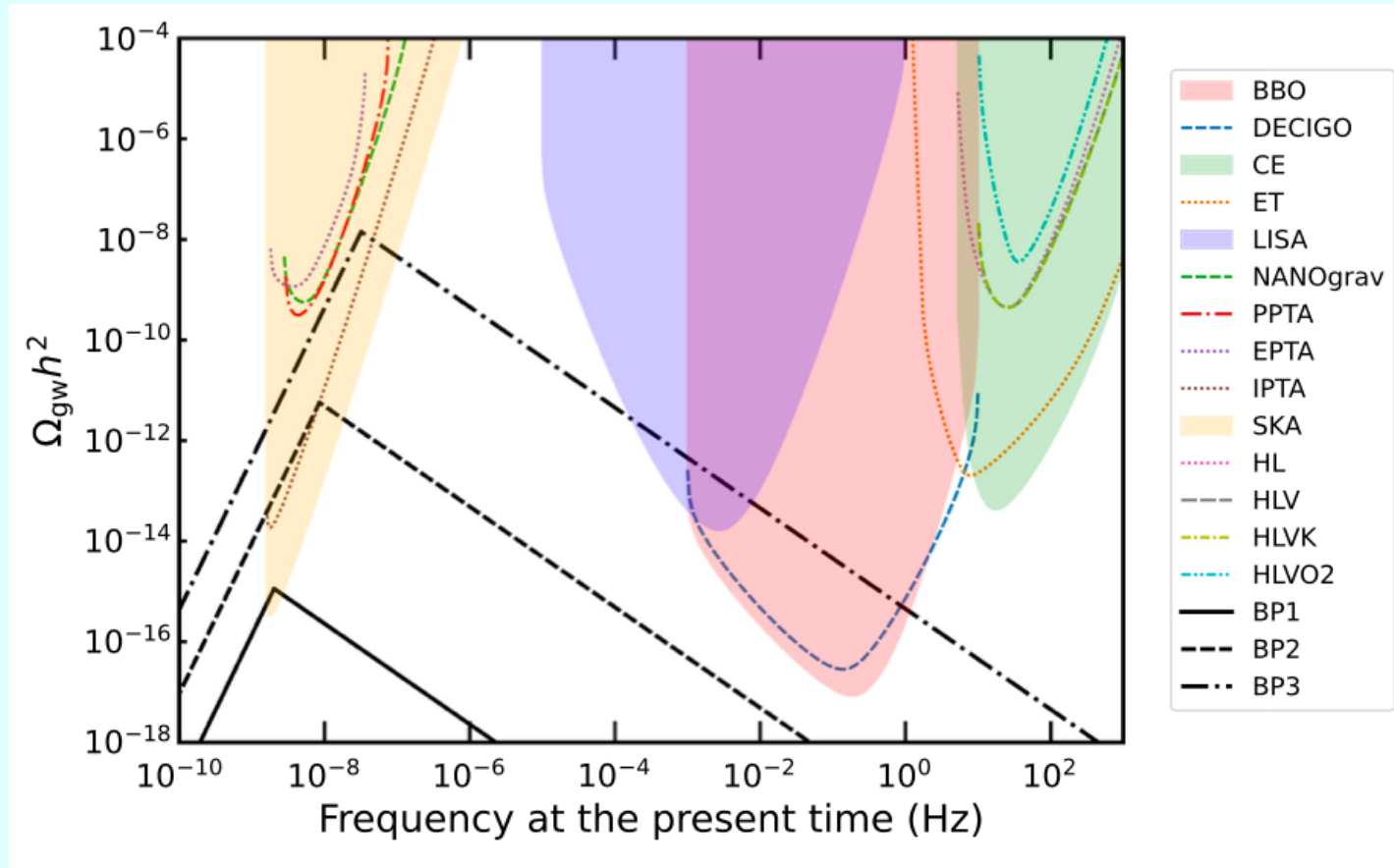
II. Degenerate field FOPT

- Z_2 of D -parity for each L and R sectors effectively gives a Z_4 with two distinct fields Δ_L and Δ_R
- The left-like and the right-like phases percolate individually
- Where the two percolated regions meet, there is a domain wall, until the whole Universe is filled with a frustrated network of domain walls
- This needs to be treated as
 - standard FOPT for bubbles
 - followed by the result of crumbling walls separating percolated regions



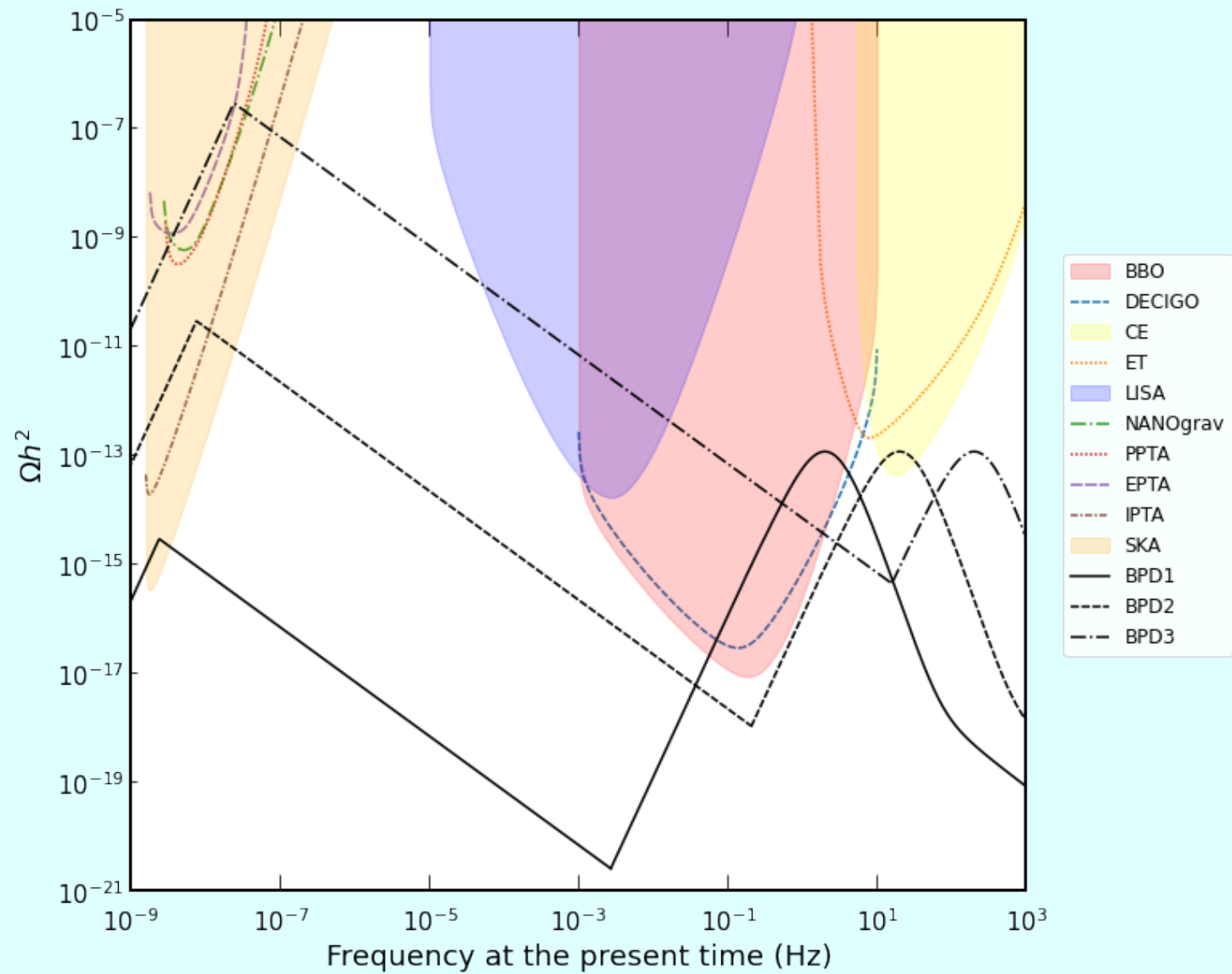
6.2 L-R two phase transition results

The Kibble or “lightcone limited SOPT” case

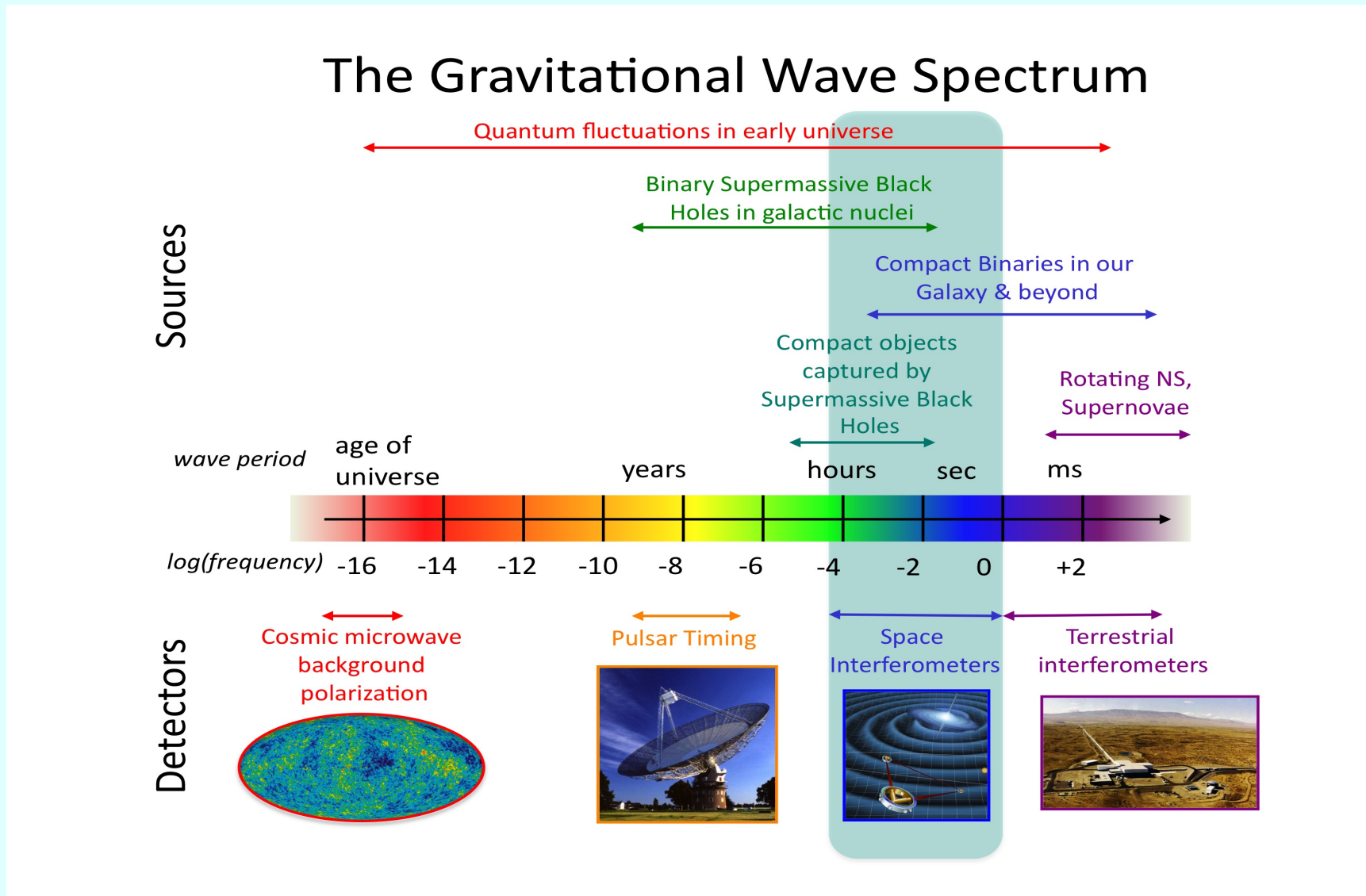


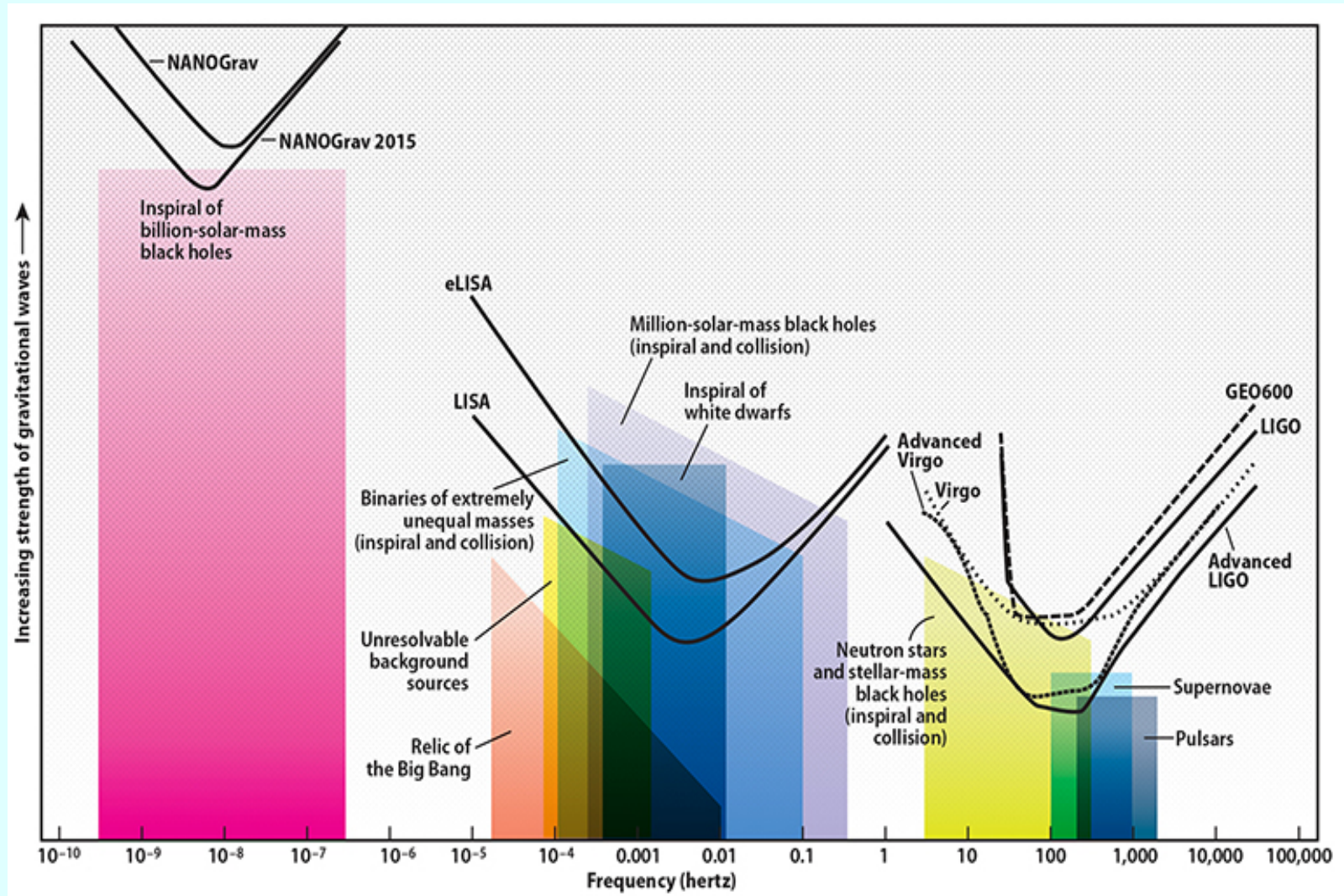
Benchmark points 1,2,3 $M_R = 10^4, 10^5, 10^6$ GeV

The Degenerate field FOPT case - two peaked spectrum



6.3 GW the observational prospects





6.4 Key experiments beyond LISA range

Already functioning

- Pulsar timing array - radio telescopes
 - PPTA, IPTA, EPTA → NANOGravuGMRT (2017) correlated data of 300–500 MHz and 1260–1460 MHz
- GAIA - orbiting sky scanning optical telescope till 2025
 - Astrometry** - μ arc-sec accuracy for about 2 billion objects
 - GWs from individually resolvable supermassive black hole binaries

7 Conclusion

- GUT Bgenesis and EW Bgenesis are unrealistic possibilities
- Thermal leptogenesis requires fine tuning
- Low (TeV to PeV) scale leptogenesis viable through phase transition Domain Walls
- Presented the case of L-R models transient CP phase relation to its zero temperature value $M\ddot{u}$
 - to be verified through Electron EDM
- Two possibilities for the L-R case; 2-peak signal for FOPT.

Danke schön!

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