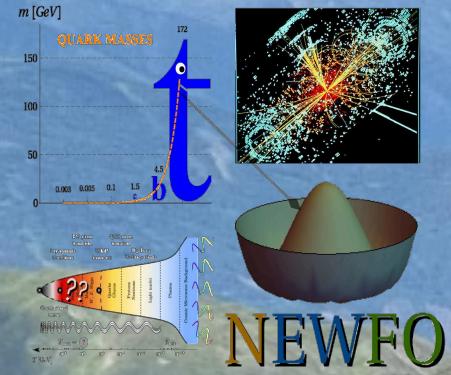


Cartography of gauge-chiral dynamics

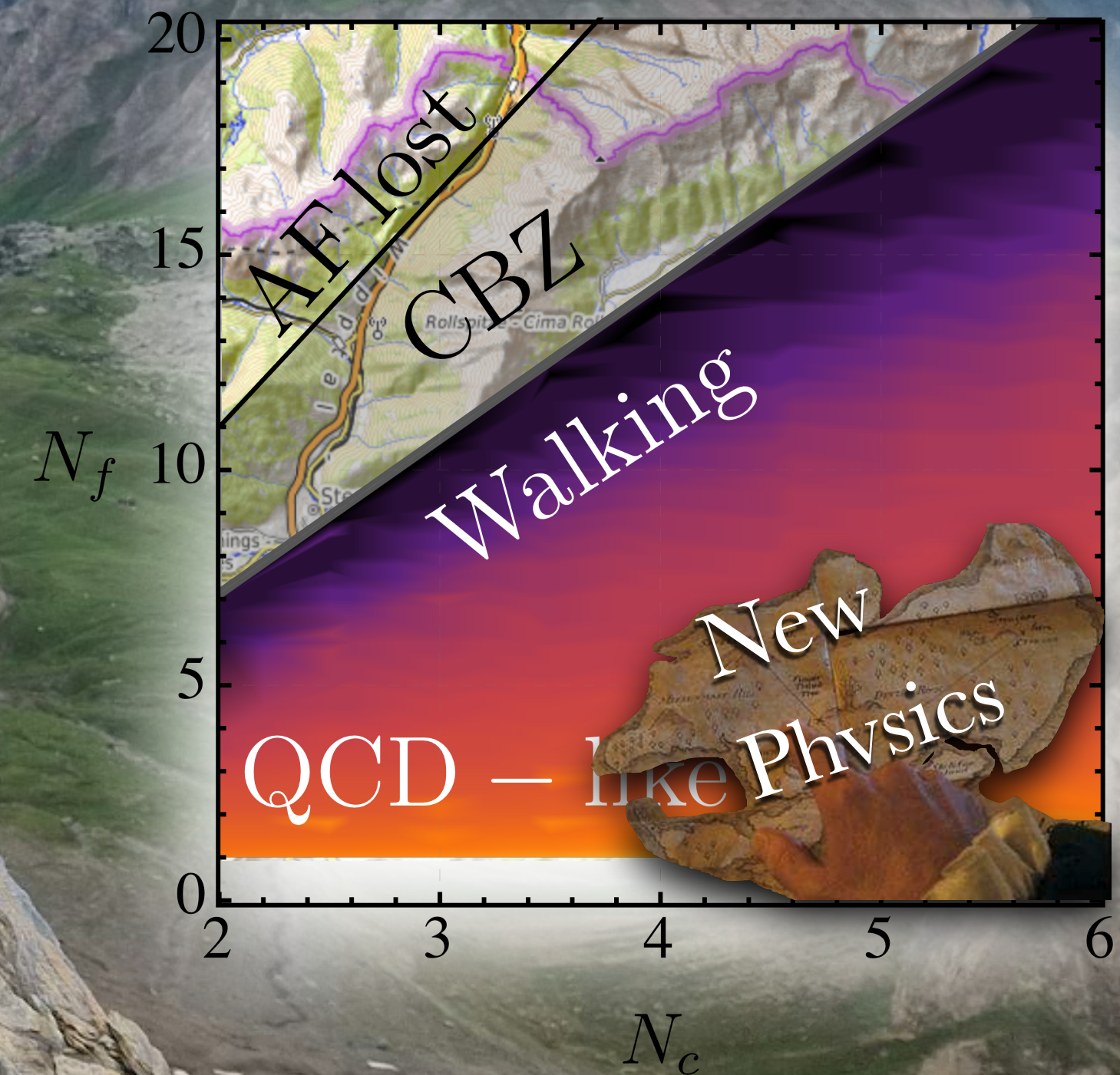
1. confinement



2. $d\chi$ SB

3. conformality

Work in collaboration with
Florian Goertz and Jan M. Pawłowski



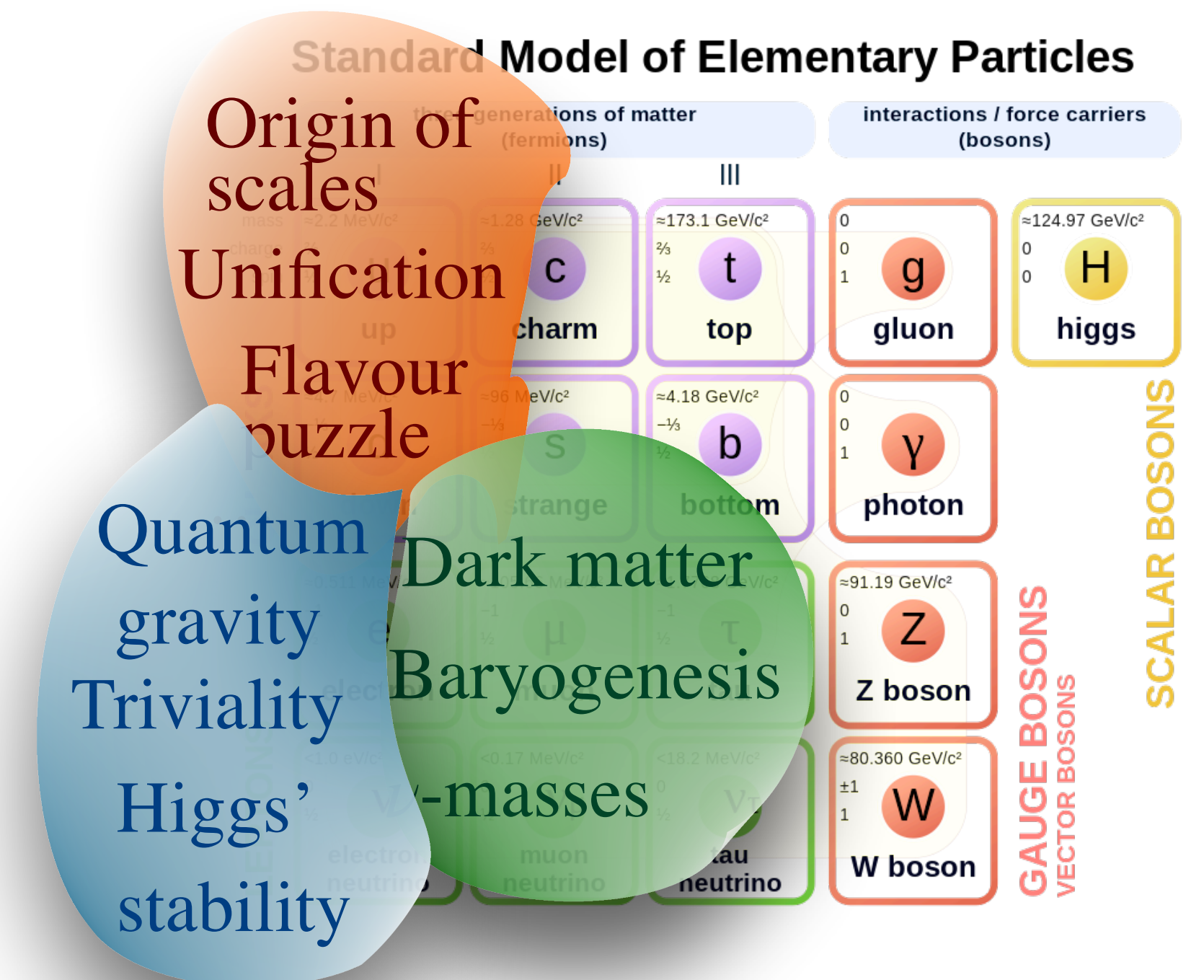
Álvaro Pastor Gutiérrez

MITP 5.08.2024

Gauge-chiral theories

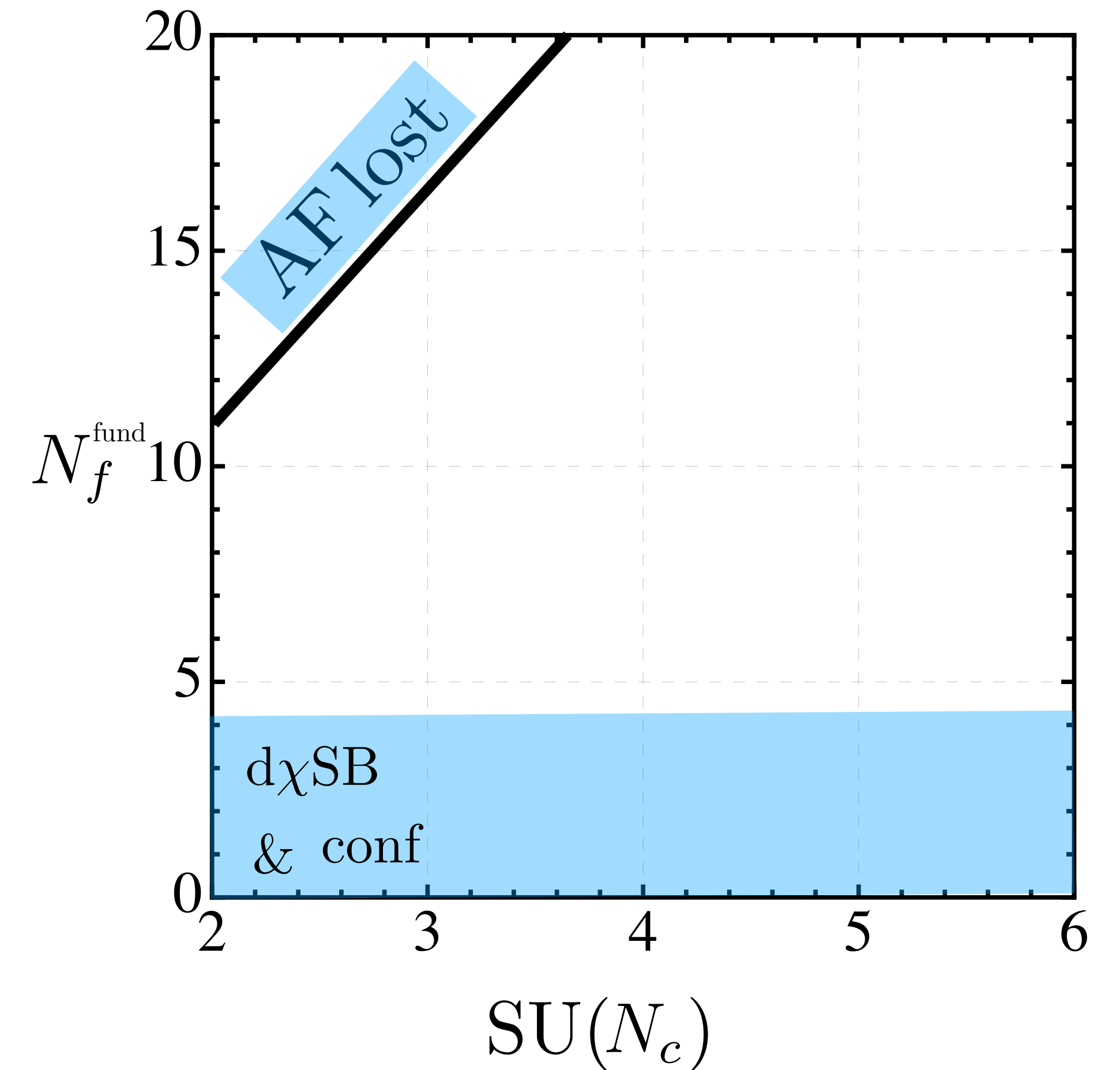
$$S = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} + \bar{\psi} \not{D} \psi$$

- ◆ Importance in the **natural world**
- ◆ Non-trivial **dynamics** and **phenomena**:
 - Confinement
 - Dynamical chiral symmetry breaking (dχSB)
 - Near-conformality (walking)
- ◆ Can be the **solution** to our **problems** and **puzzles**
 - Dark strong sectors
 - Composite Higgs and Technicolour
 - Cosmological phase transitions and gravitational wave signatures



What we know so far from first-principles?

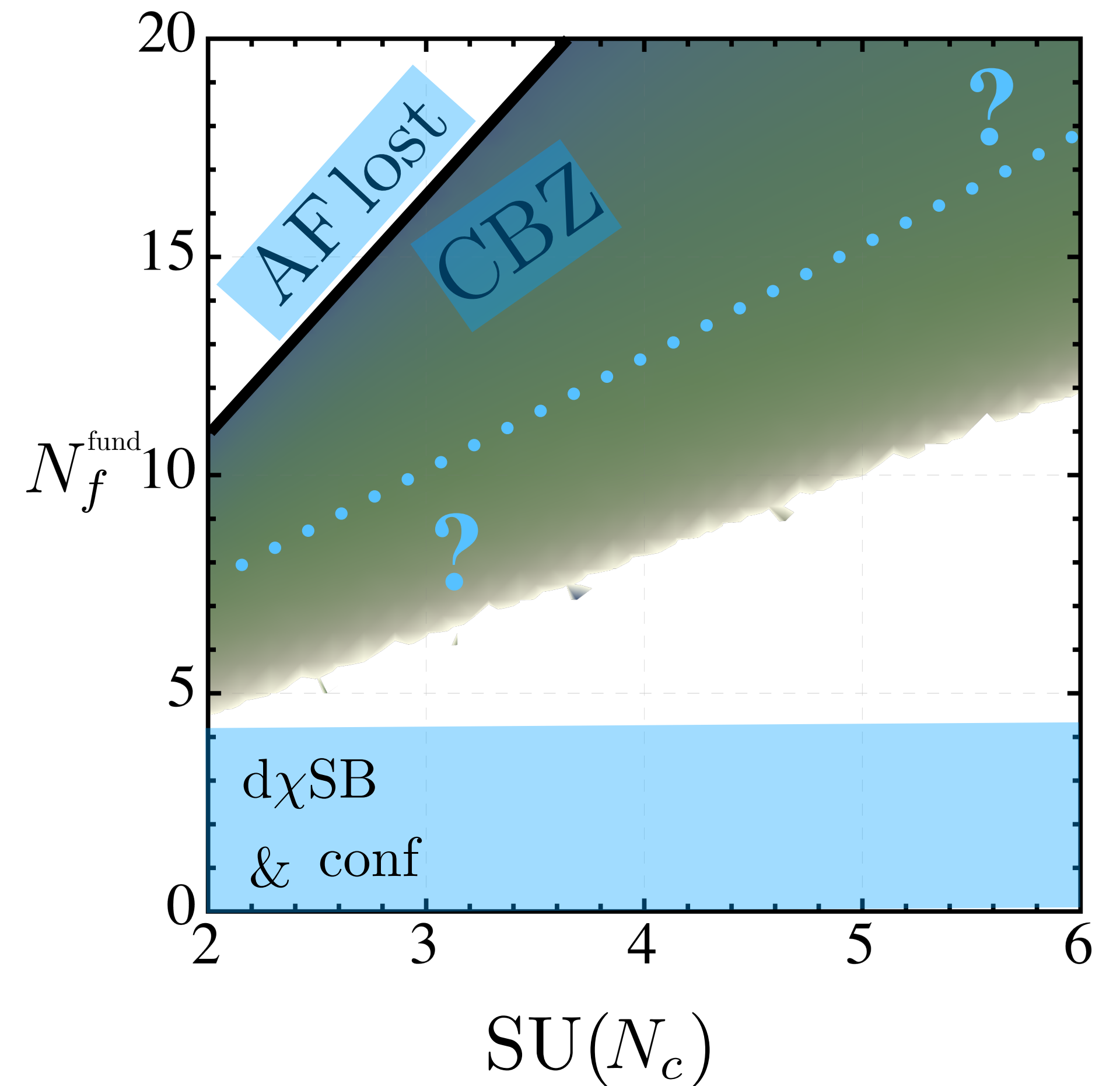
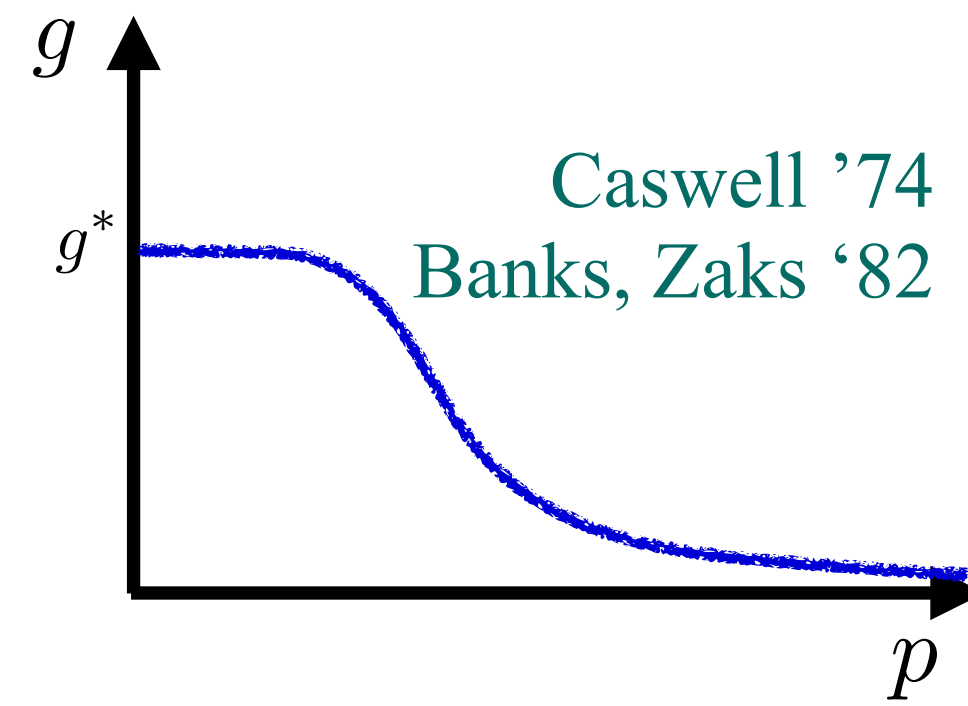
- ◆ Where **asymptotic freedom** is lost
- ◆ **Few flavours**: QCD-like theories



What we know so far from first-principles?

- ◆ Where **asymptotic freedom** is lost
- ◆ **Few flavours**: QCD-like theories
- ◆ **CBZ IR fixed-point** (perturbative)

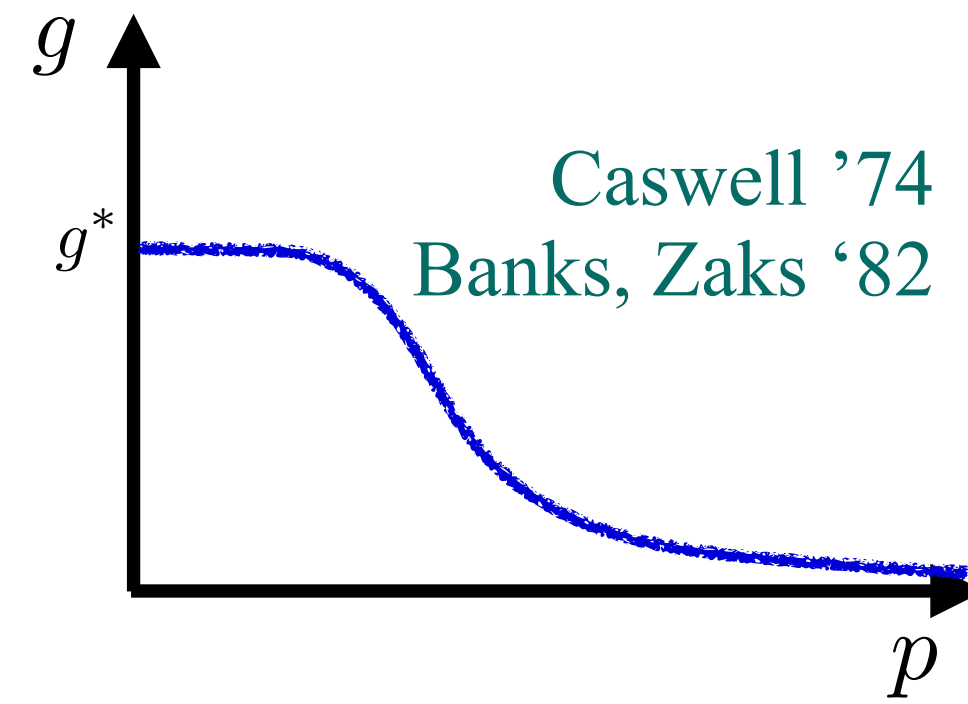
$$\beta_g = -\frac{g^3}{(4\pi)^2} \left(\frac{11}{3} C_A - \frac{4}{3} T_F N_f \right) - \frac{g^5}{(4\pi)^4} \left(\frac{34}{3} C_A - 4 C_F T_F N_f - \frac{20}{3} C_A T_F N_f \right) + \dots$$



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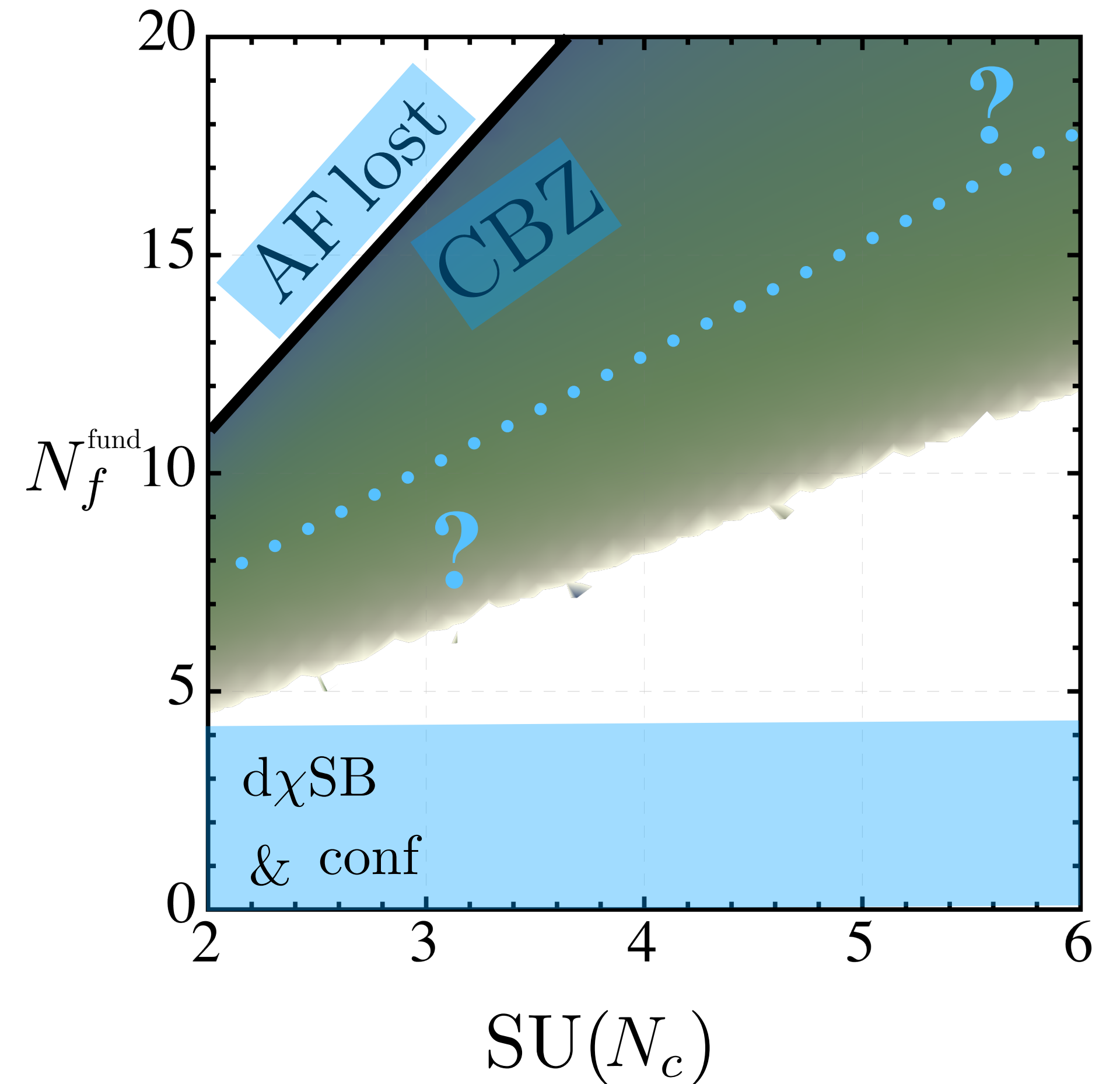
- ◆ Where **asymptotic freedom** is lost
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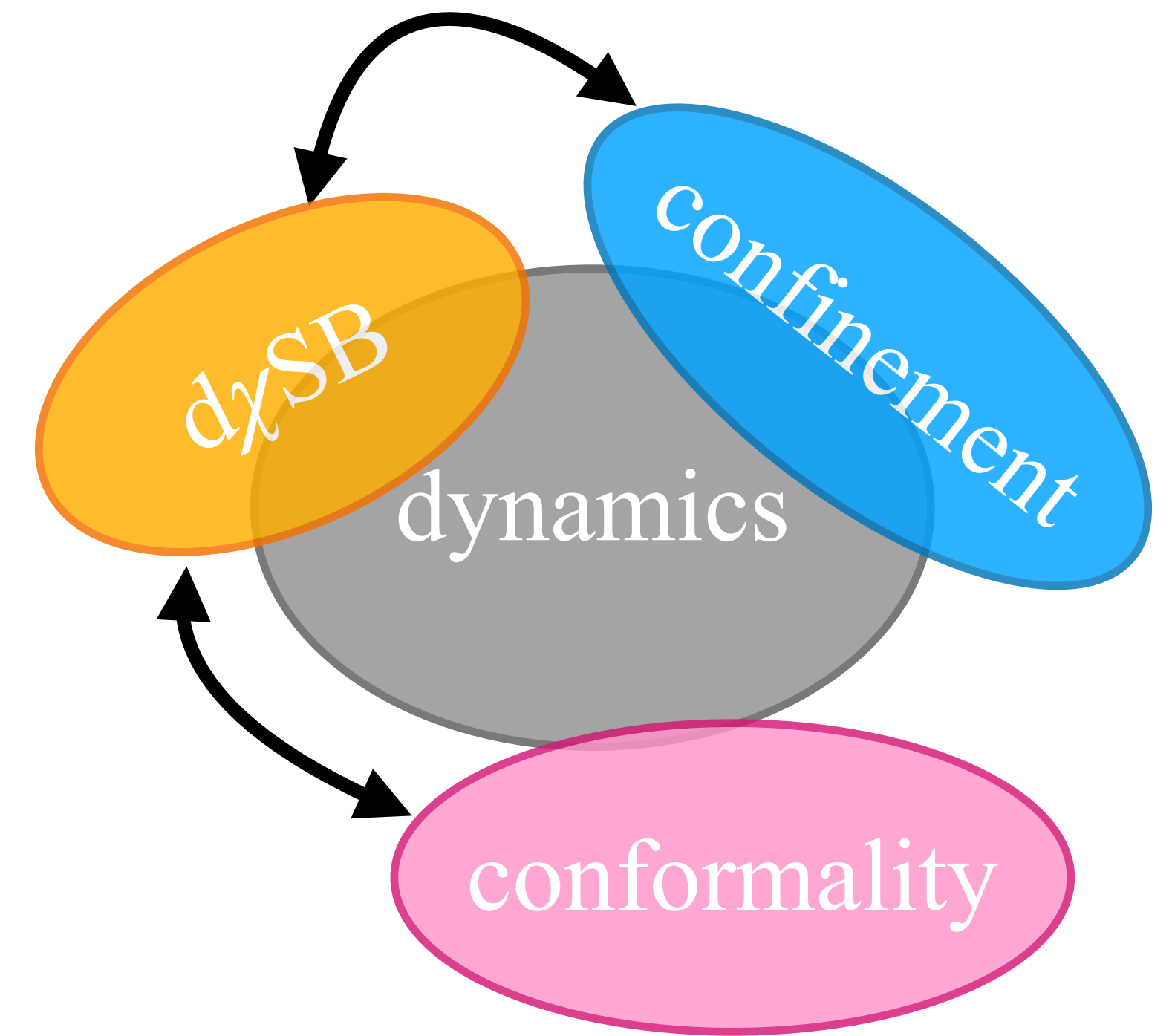
- ◆ Lattice deficiencies:

- Individual theory investigations
- The chiral limit is difficult: singular fermion determinant



Towards the uncharted regime

1. Chart the **landscape** of gauge-chiral theories
2. Understand **dynamics** in the many-flavour limit from **first principles**
3. Obtain **fundamental parameters** and the relation between scales
4. Determine the **boundary** between phases: conformal & non-conformal
5. Provide **UV input** to low-energy effective theories



- ◆ Functional methods in the continuum limit: the **functional Renormalisation Group**
 - No difficulties in the exact chiral limit

$$\Gamma_k[\phi]$$

Functional Renormalisation Group

- ◆ Effective **average** action: $\Gamma_k[\phi]$
 - Average action of fields over a k^{-d} volume
 - Kadanoff's block-spinning idea in the continuum limit

◆ Pedagogical introduction to fRG: [Gies \[0611146\]](#)

◆ Review of achievements: [Dupuis, Canet, Eichhorn, Metzner, Pawłowski, Tissier, Wschebor \[2006.04853\]](#)

Functional Renormalisation Group

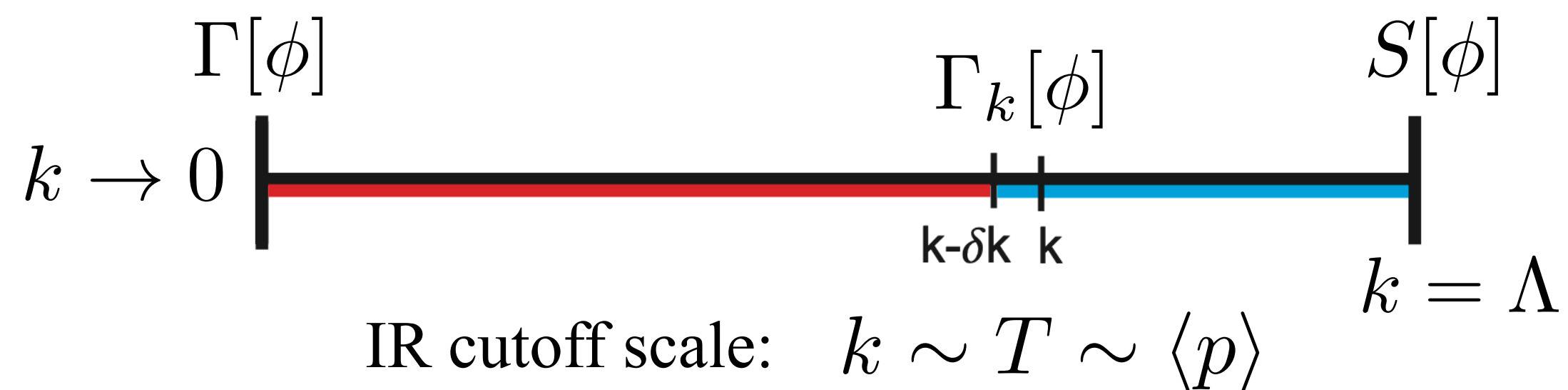
◆ Effective **average** action: $\Gamma_k[\phi]$

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$$\Delta S_k[\phi] = \int_p \phi(p) R_k \phi(-p)$$

$$\Gamma_k[\phi] = \int_x J(x)\phi(x) - \mathcal{W}_k[J] - \Delta S_k[\phi]$$

Wetterich '89



◆ Pedagogical introduction to fRG: [Gies \[0611146\]](#)

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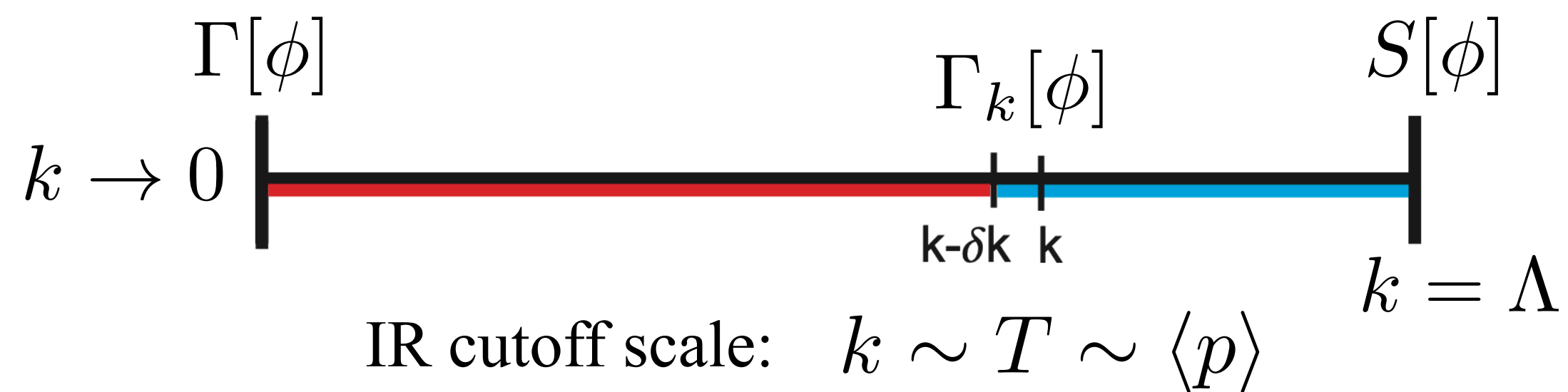
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Wetterich '89



Flow equation:

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \left[\frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k \right] = \frac{1}{2} \text{Diagram}$$

The diagram is a circle with a cross inside, representing a one-loop trace.

$$\partial_t \equiv k \partial_k$$

Wetterich '93

- ◆ One loop exact
- ◆ **Non-perturbative**
- ◆ **Mass-dependent**
- ◆ **Analytic regulators**
- ◆ **Versatile**
- ◆ Systematic expansion schemes
- ◆ UV-IR finite
- ◆ Diagrammatic
- ◆ Real-time formulation

◆ Pedagogical introduction to fRG: [Gies \[0611146\]](#)

◆ Review of achievements: [Dupuis, Canet, Eichhorn, Metzner, Pawłowski, Tissier, Wschebor \[2006.04853\]](#)

Deriving correlation functions

$$\Gamma_k[\phi] = \sum_n^{N_{\text{tr}}} \int_p \Gamma_{k, \phi_{i_1} \dots \phi_{i_n}}^{(n)}(p_1, \dots, p_n) \phi_{i_n}(p_n) \dots \phi_{i_1}(p_1)$$

$$\Gamma_{k, \phi_{i_1} \dots \phi_{i_n}}^{(n)}(p_1, \dots, p_n) = \frac{\delta}{\delta \phi_{i_1}(p_1)} \dots \frac{\delta}{\delta \phi_{i_n}(p_n)} \Gamma_k[\phi]$$

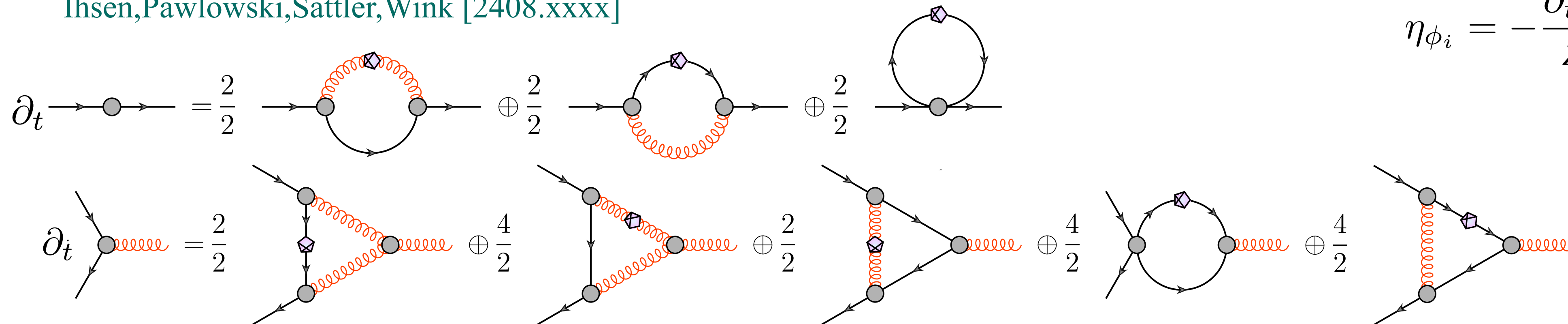
Example: quark-gluon vertex and gauge coupling

$$\partial_t \Gamma_k^{(\bar{\psi}\psi A)} = \partial_t \left(Z_A^{1/2} Z_\psi g_{\bar{\psi}\psi A} \cdot \mathcal{T}_\mu \right) \longrightarrow \partial_t g_{\bar{\psi}\psi A} = \frac{\text{Tr} \left[\mathcal{T}_\mu \partial_t \Gamma_k^{(\bar{\psi}\psi A)} \right]}{\text{Tr} \left[\mathcal{T}_\mu^2 \right] Z_A^{1/2} Z_\psi} + \left(\frac{1}{2} \eta_A + \eta_\psi \right) g_{\bar{\psi}\psi A}$$

with anomalous dimensions:

$$\eta_{\phi_i} = -\frac{\partial_t Z_{\phi_i}}{Z_{\phi_i}} = -\frac{\partial_{p^2} \partial_t \Gamma_k^{(\phi_i \phi_i)}}{Z_{\phi_i}}$$

Ihsen, Pawłowski, Sattler, Wink [2408.xxxx]



Colour confinement: gluon mass gap

♦ **Absence of colour asymptotic states**

♦ **Massive spectrum** of bound states (glueballs)

♦ Existence of a gluon **mass gap**:

- Wilson area law $Z_c(p) \propto (p^2)^{\kappa_{\text{ghost}}}$ $\kappa_{\text{ghost}} = 0.579 \pm 0.005$
- Kugo-Ojima conditions $Z_A(p) \propto (p^2)^{-2\kappa_{\text{gluon}}}$ $\kappa_{\text{gluon}} = 0.573 \pm 0.002$
- Confinement-deconfinement phase transition

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◆ **fRG bootstrap approach to confinement:**

- **Mass gap** generated by **quantum fluctuations**

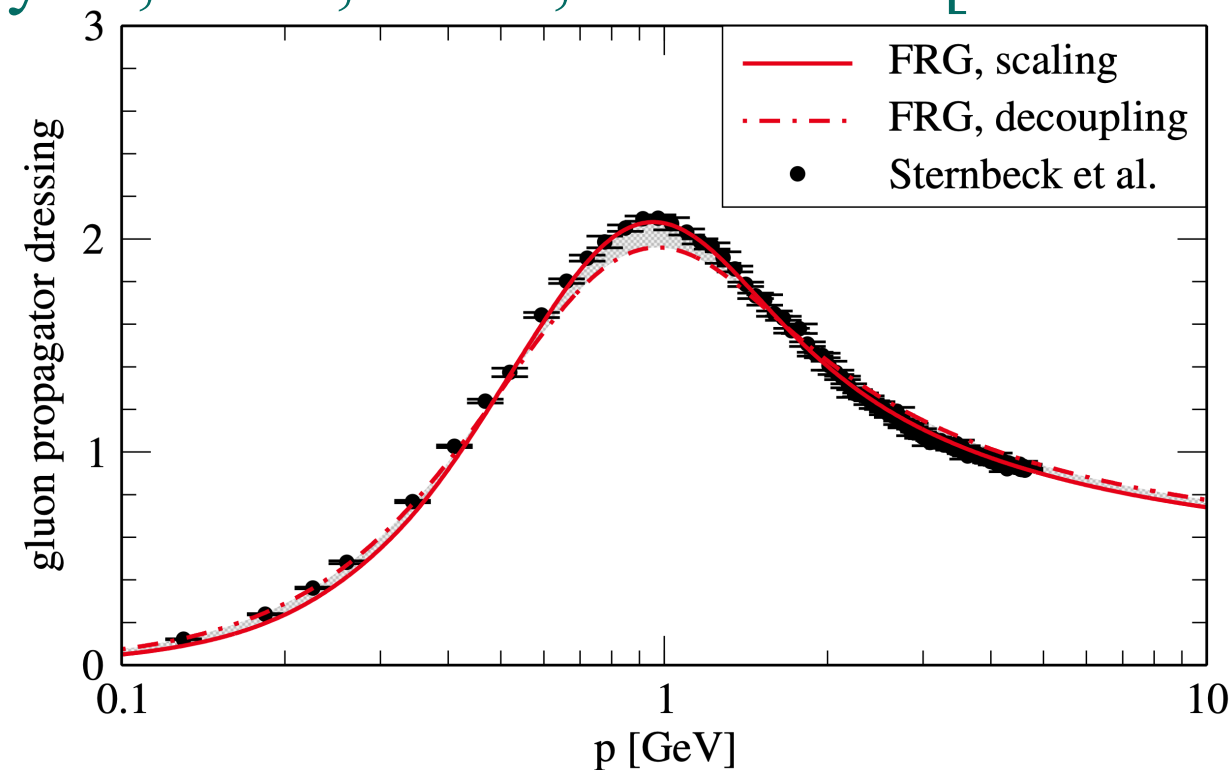
$$\Gamma_k^{(AA)}(p^2) = Z_{A,k}(p) (p^2 + m_{A,k}^2) = \hat{Z}_{A,k}(p) p^2$$

- **Uniquely defined** confinement condition
- Confinement from correlation function

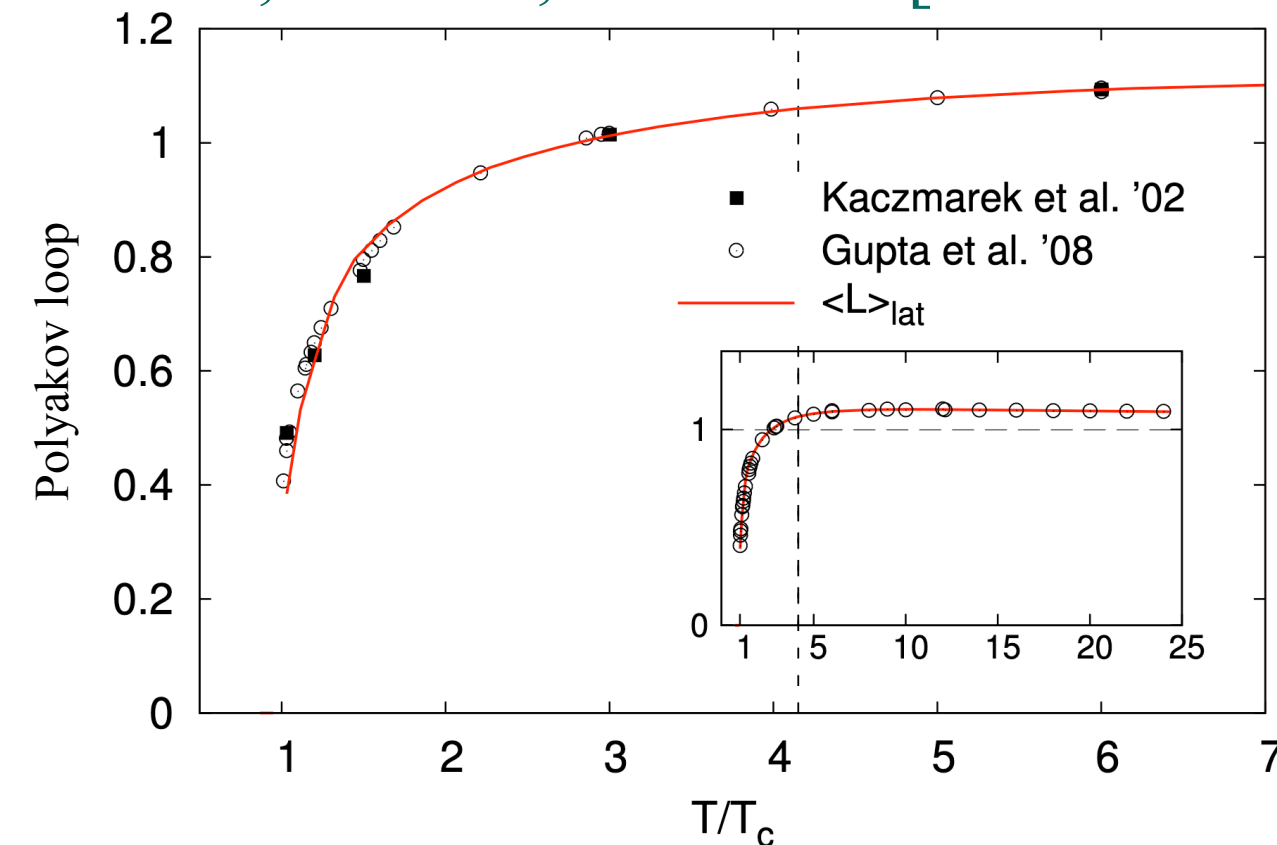
◆ **New: “easy” confinement**

- **Cutoff dependences suffice**
- Semi-analytical
- Facilitate study beyond QCD-limit

Cyrol,Fister,Mitter,Pawlowski [1605.01856]



Herbst,Luecker,Pawlowski[1510.03830]

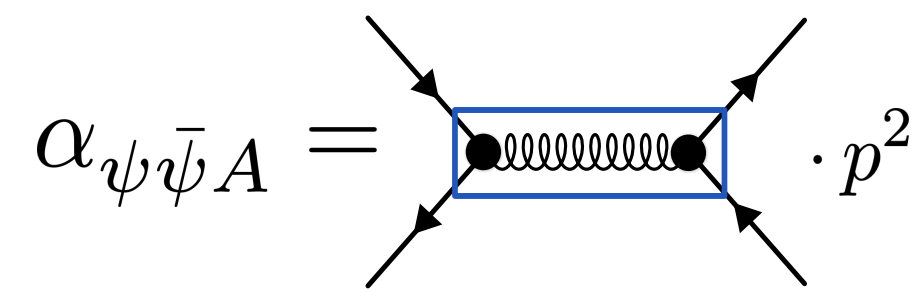


Confinement in correlation functions

◆ Flows computed: $\{g_{A\bar{\psi}\psi}, g_{A\bar{c}c}, g_{A^3}, g_{A^4}, \bar{m}_A, Z_A, Z_c\}$

◆ **Exchange couplings:**

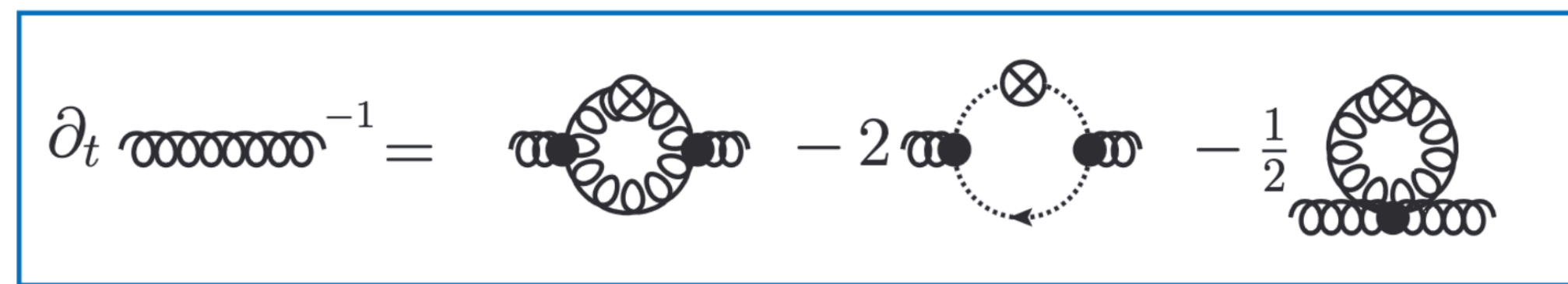
$$\alpha_{\psi\bar{\psi}A} = \frac{\left[\Gamma_k^{(\psi\bar{\psi}A)}\right]_{\mathcal{T}=1}^2}{4\pi Z_A Z_\psi^2} = \frac{g_{\psi\bar{\psi}A}^2}{4\pi(1 + \bar{m}_A^2)}$$



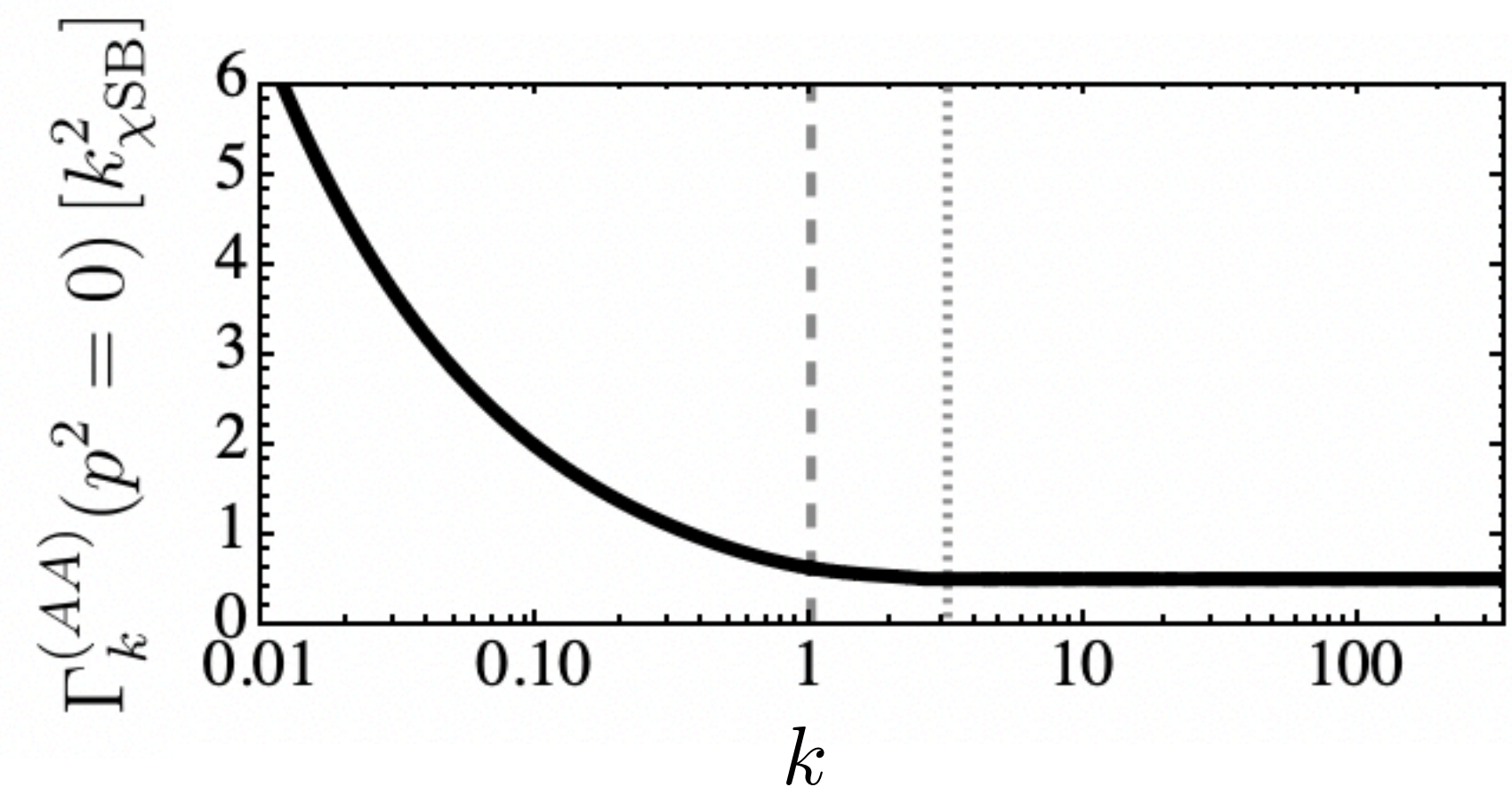
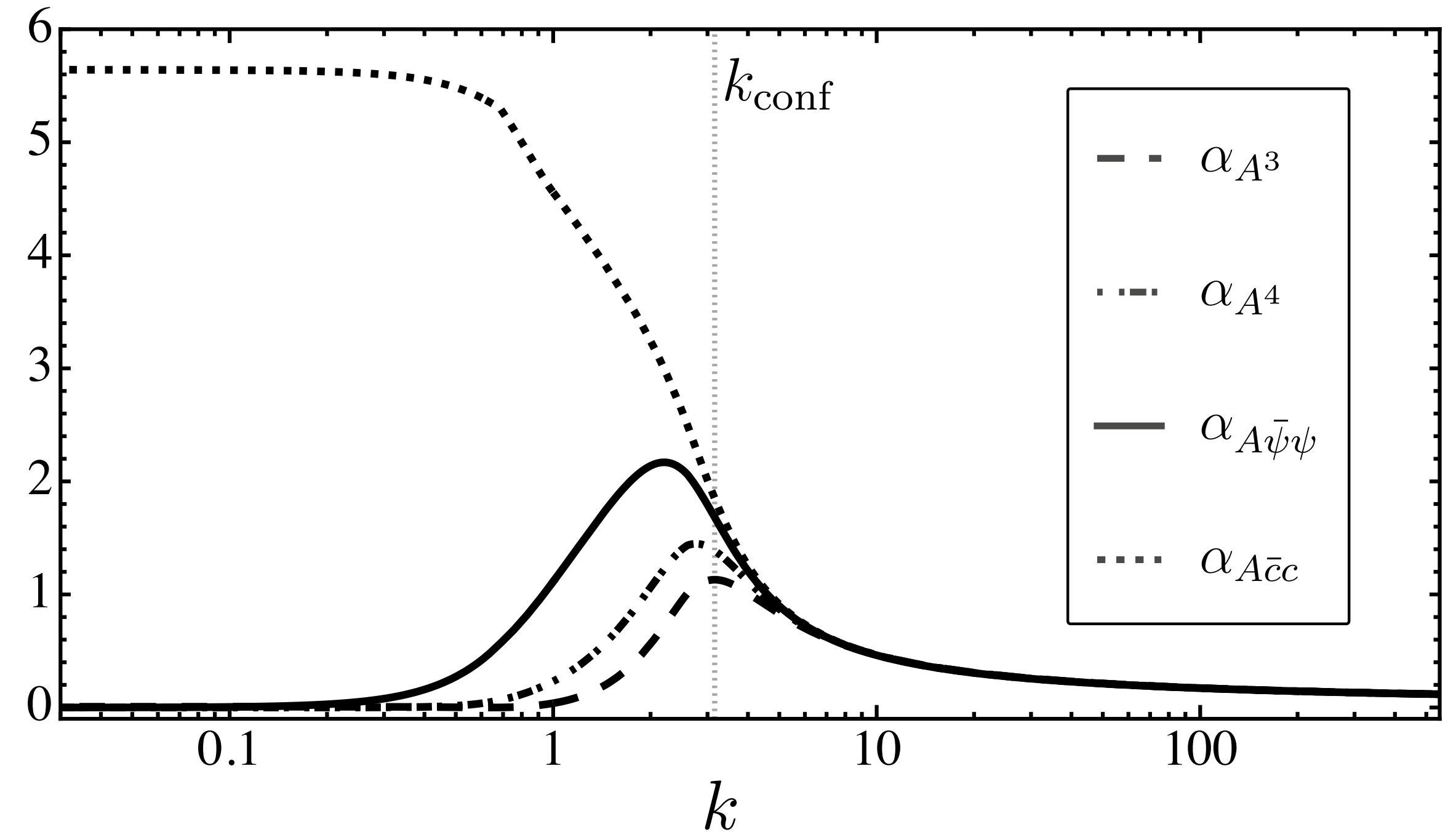
◆ **Decay of correlation functions below the mass gap scale**

$$k_{\text{conf}} \sim m_{A, \text{gap}} \sim T_{\text{conf}} \sim \Lambda_{\text{QCD}}$$

◆ The **interplay of gapped gauge and ghost contributions**

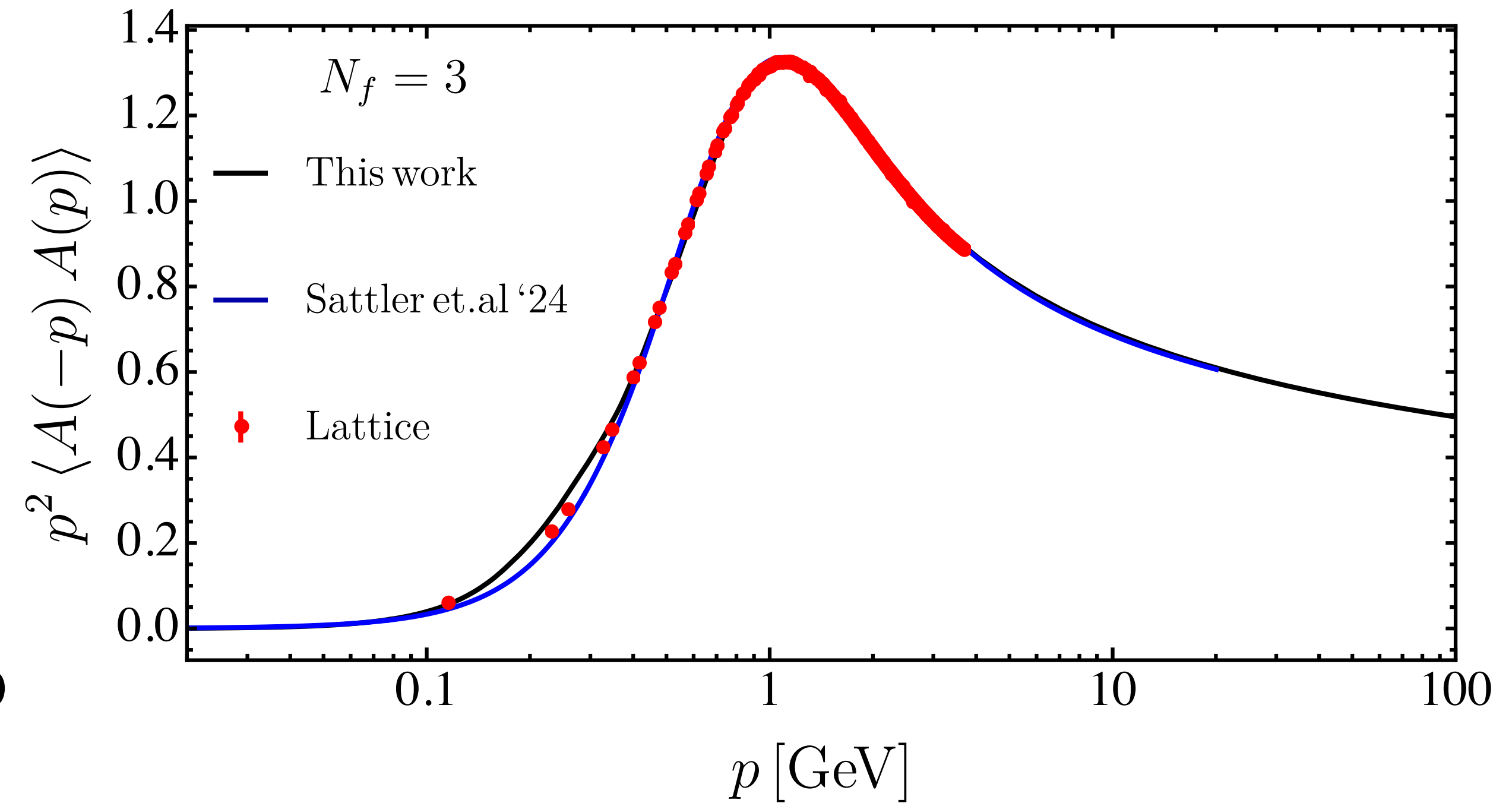
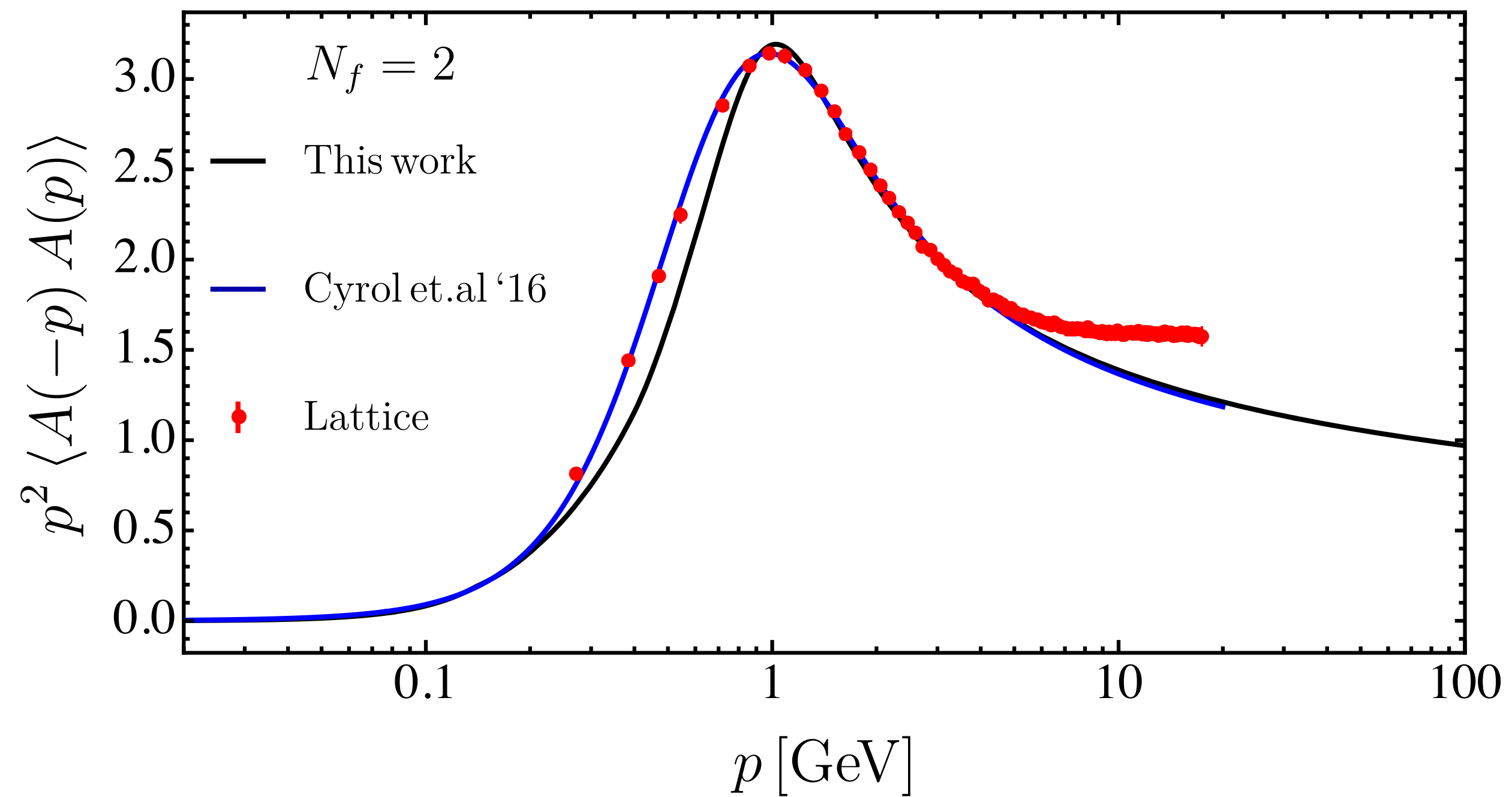


Cyrol, Fister, Mitter, Pawłowski [1605.01856]



“Easy” confinement

Gluon propagator dressing:
$$p^2 \langle A(-p)A(p) \rangle = p^2 \left[\Gamma_k^{(AA)}(p^2 = 0) \right]_{k=cp}^{-1}$$



Dynamical χ SB

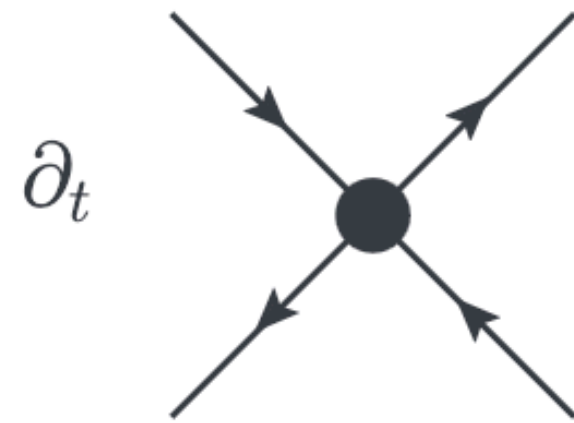
$$S = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} + \bar{\psi} \gamma_\mu D_\mu \psi$$

Dynamical χ SB

$$\Gamma = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} + \bar{\psi} \gamma_\mu D_\mu \psi + \lambda (\bar{\psi} \mathcal{T} \psi)^2 + \kappa (\bar{\psi} \mathcal{T} \psi)^3 + \dots$$

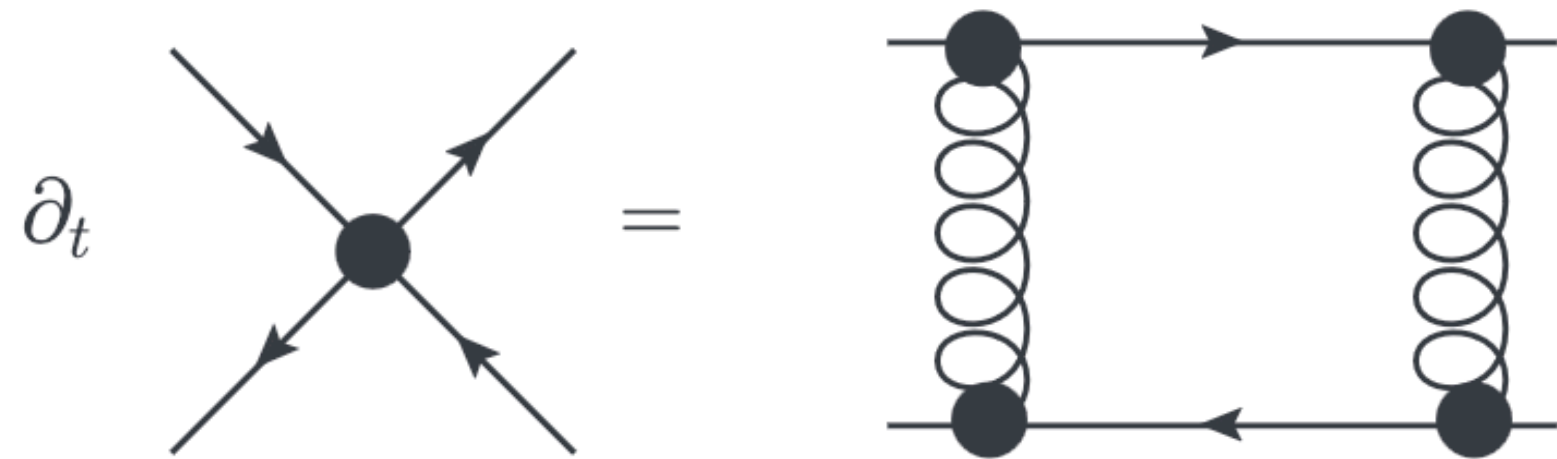
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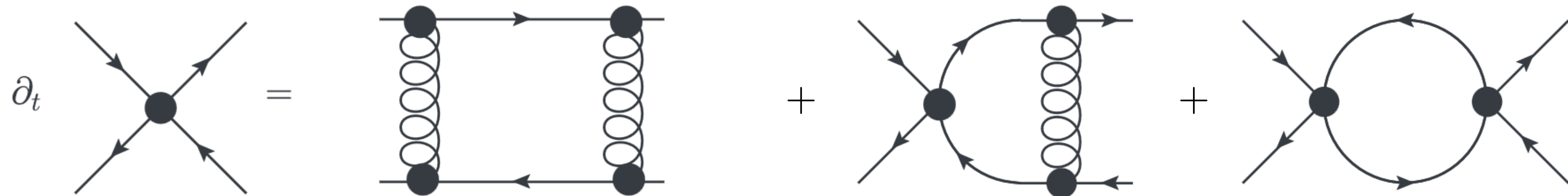
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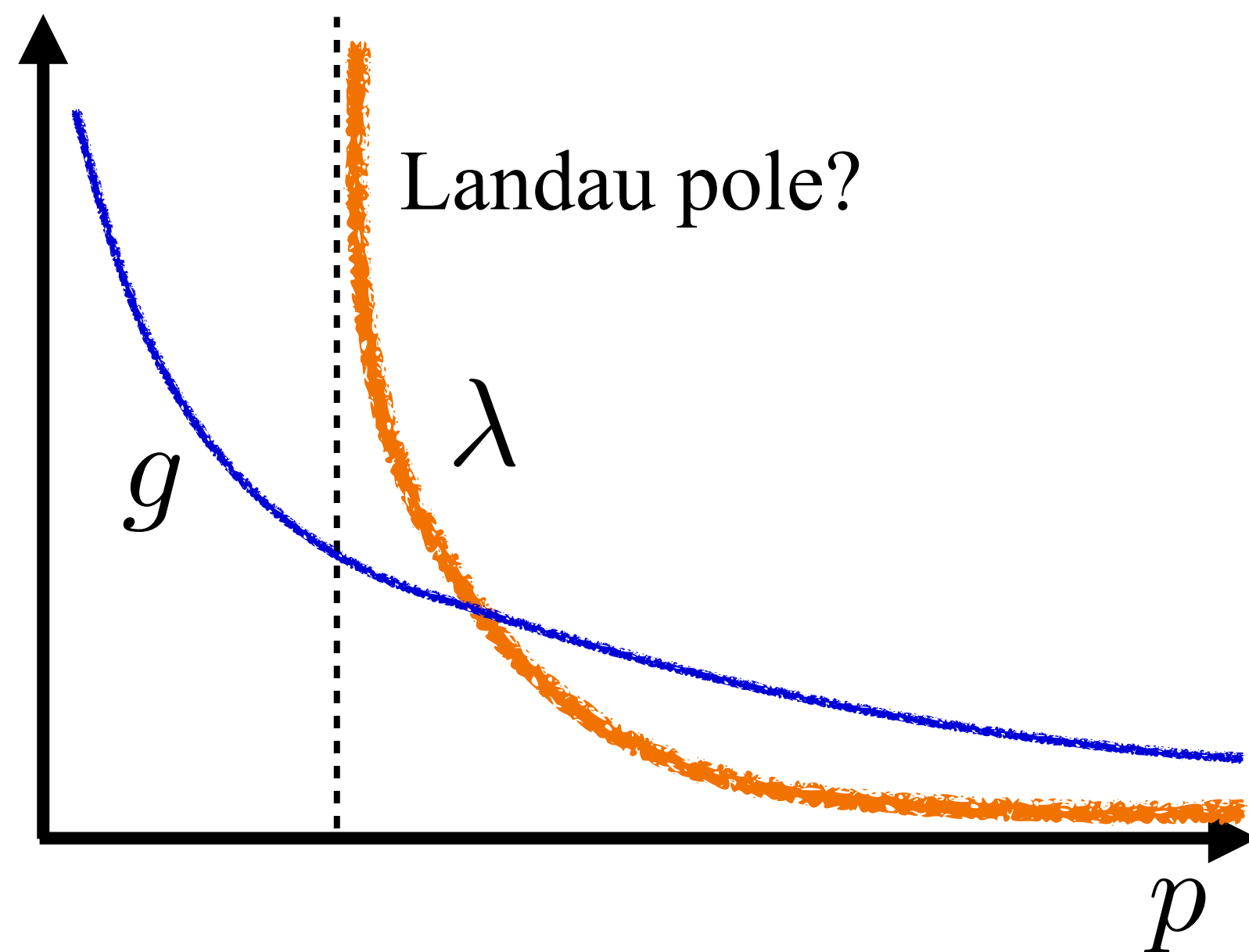
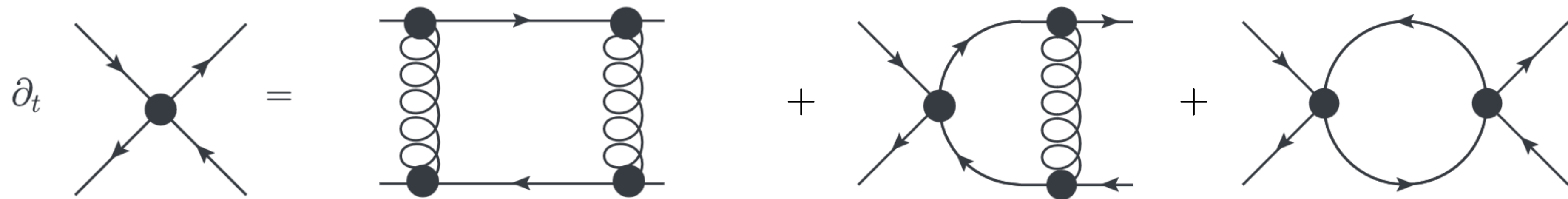
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Dynamical χ SB

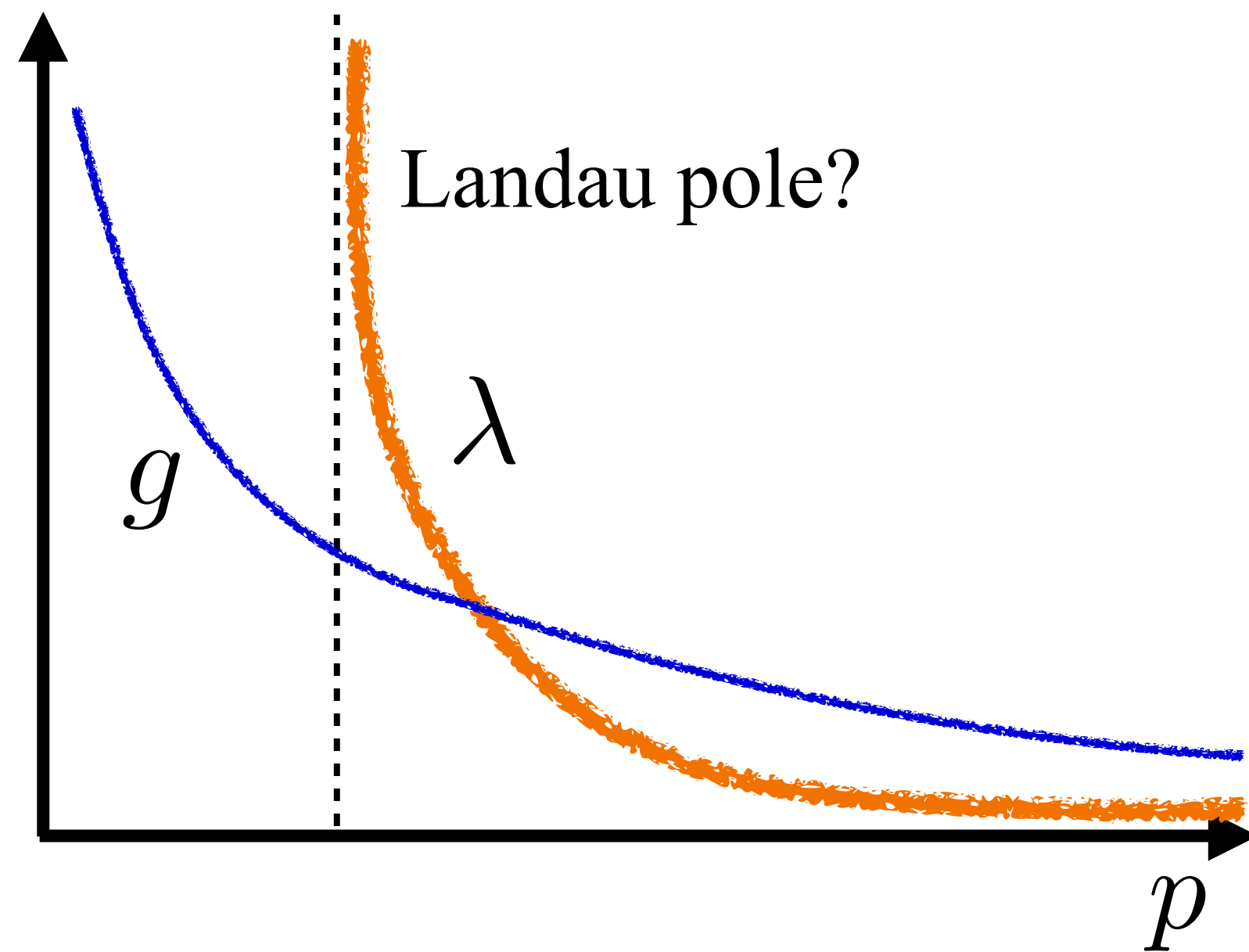
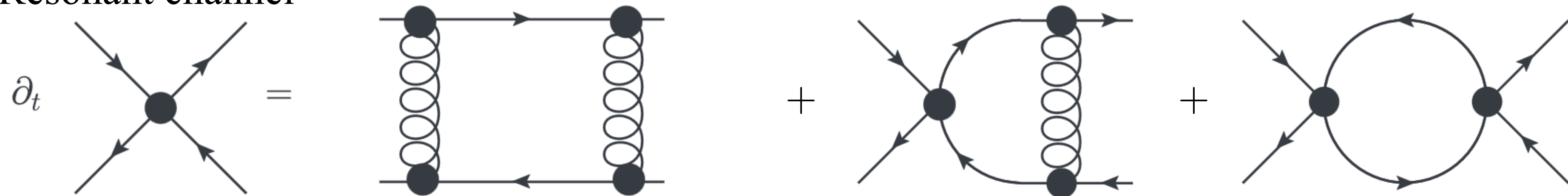
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Resonant channel

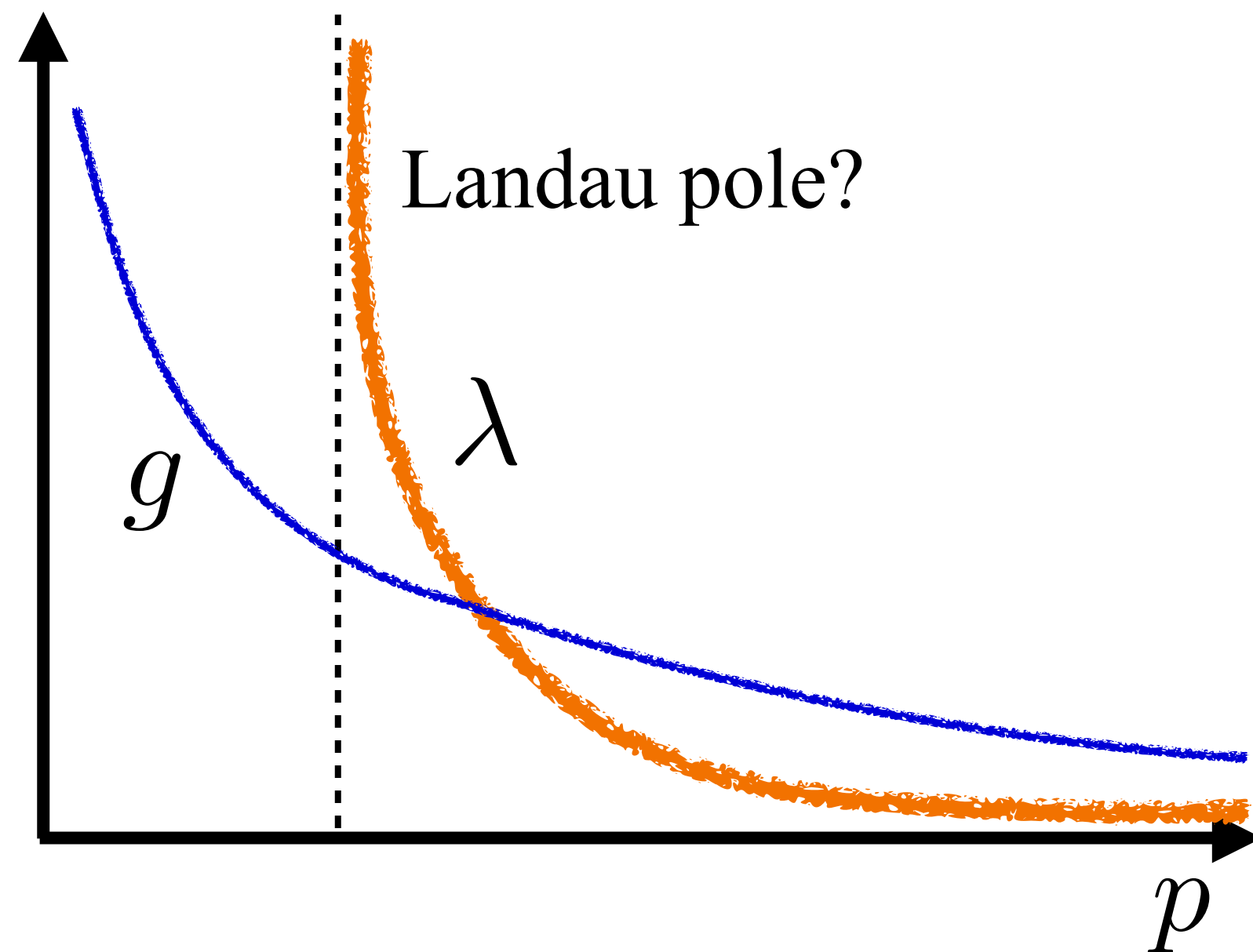
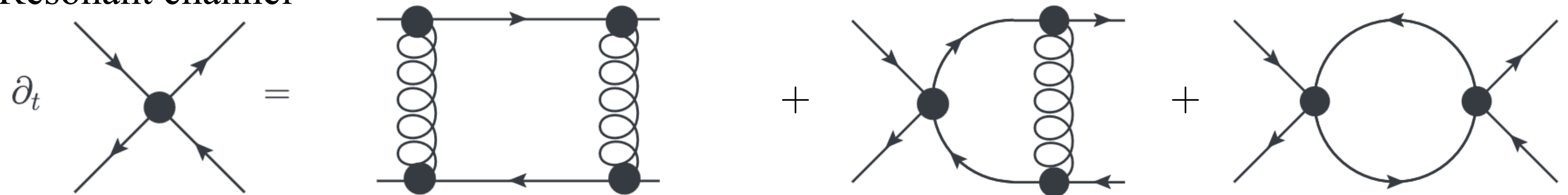


$$\langle (\bar{\psi} \psi)^2 \rangle \sim \langle \phi \phi \rangle^{-1} \longrightarrow \lambda \propto \frac{1}{p^2 + m_\phi^2}$$

Dynamical χ SB

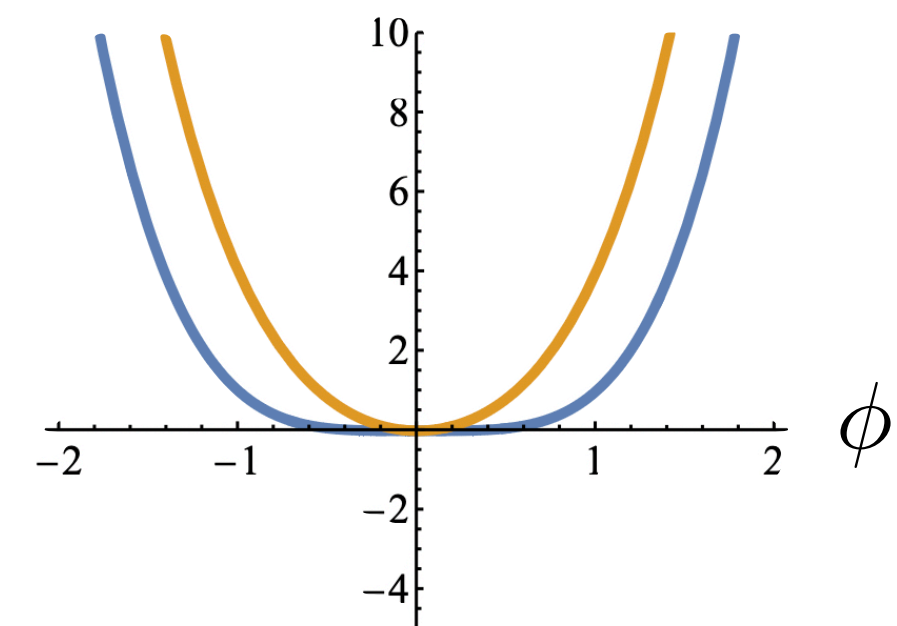
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Resonant channel



$$\langle (\bar{\psi} \psi)^2 \rangle \sim \langle \phi \phi \rangle^{-1} \longrightarrow \lambda \propto \frac{1}{p^2 + m_\phi^2}$$

$$\chi\text{SB} \longrightarrow (\lambda = \infty) \equiv (m_\phi = 0)$$

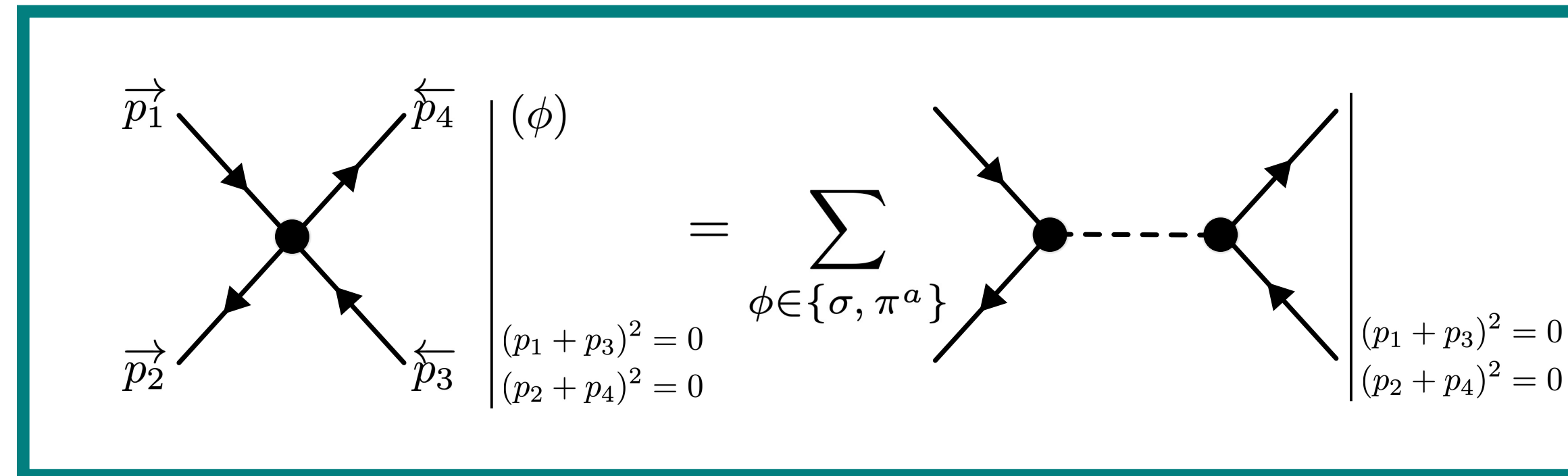


Higher dimensional operators carry information of bound states formation and $\mathbf{d}\chi\text{SB}$

Emergent composites

$$\Gamma = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + (\partial_\mu \bar{c}^a) D_\mu^{ab} c^b + \frac{1}{2\xi} (\partial_\mu A_\mu^a)^2 + \bar{\psi} [(\gamma_\mu D_\mu)] \psi - \lambda \left[(\bar{\psi} T_f^0 \psi)^2 + (\bar{\psi} i\gamma_5 T_f^a \psi)^2 \right]$$

Stratonovich'57 Hubbard'59 Gies, Wetterich '01

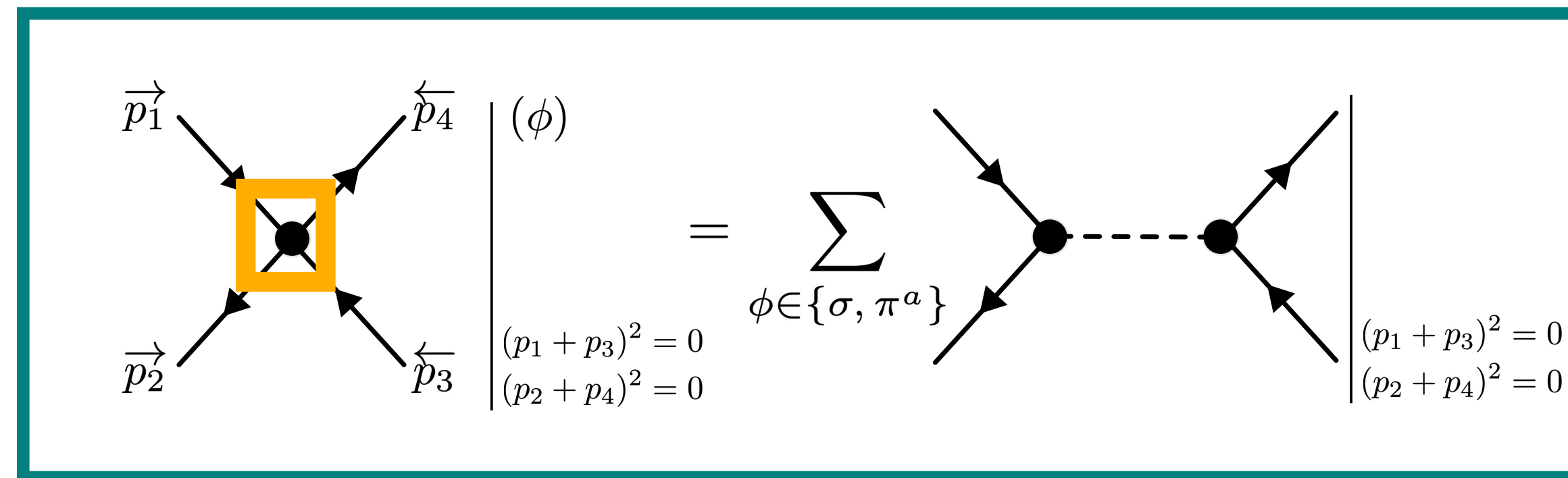


Pawlowski [hep-th/0512261] Fukushima, Pawlowski, Strodthoff [2103.01129]

Emergent composites

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Stratonovich'57 Hubbard'59 Gies, Wetterich '01

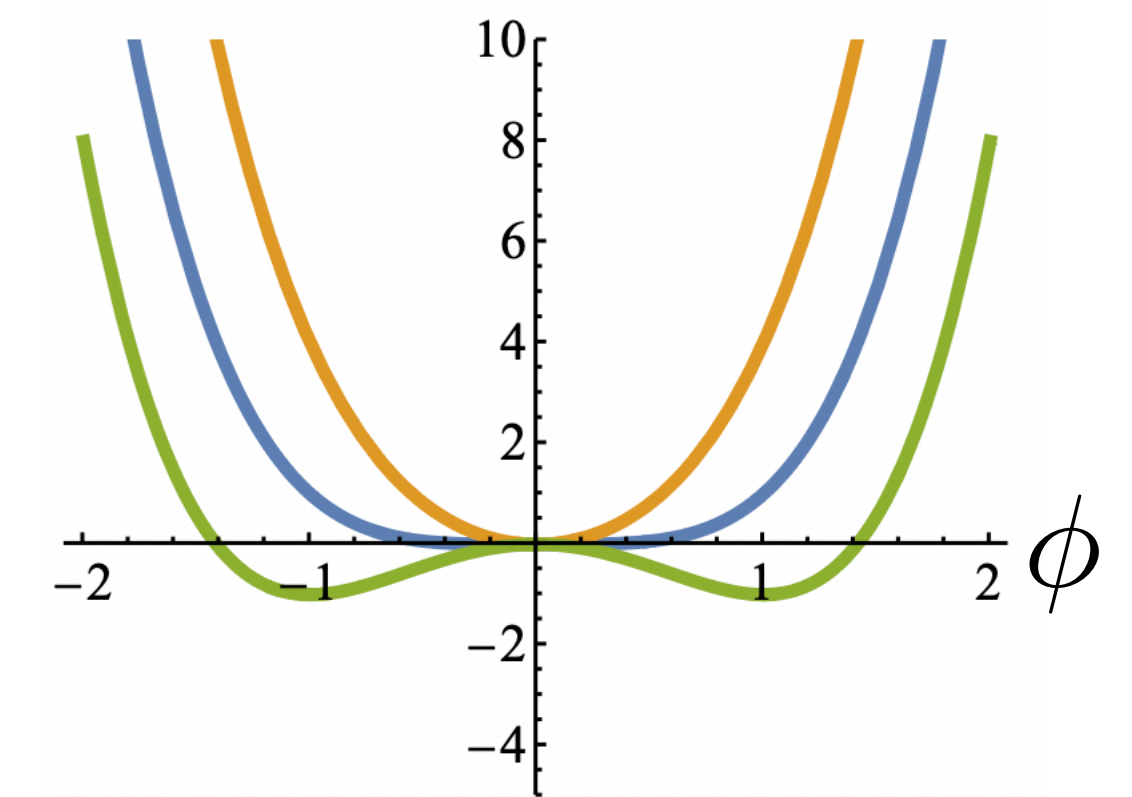
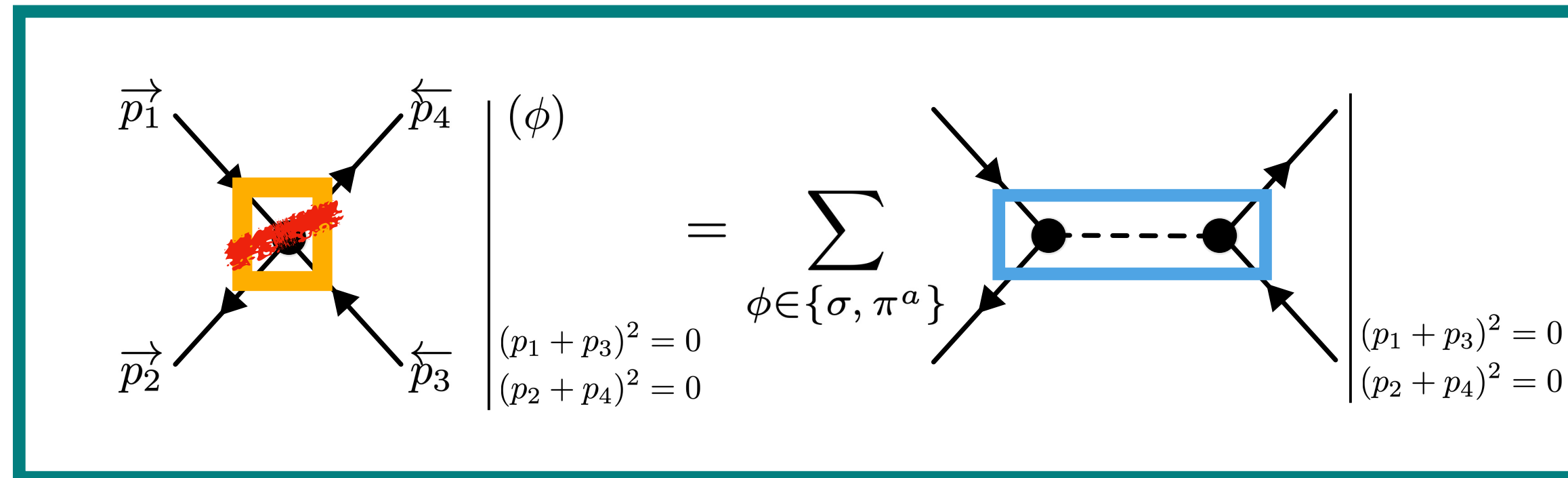


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Stratonovich '57 Hubbard '59 Gies, Wetterich '01



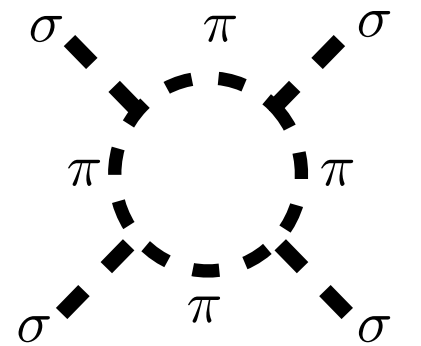
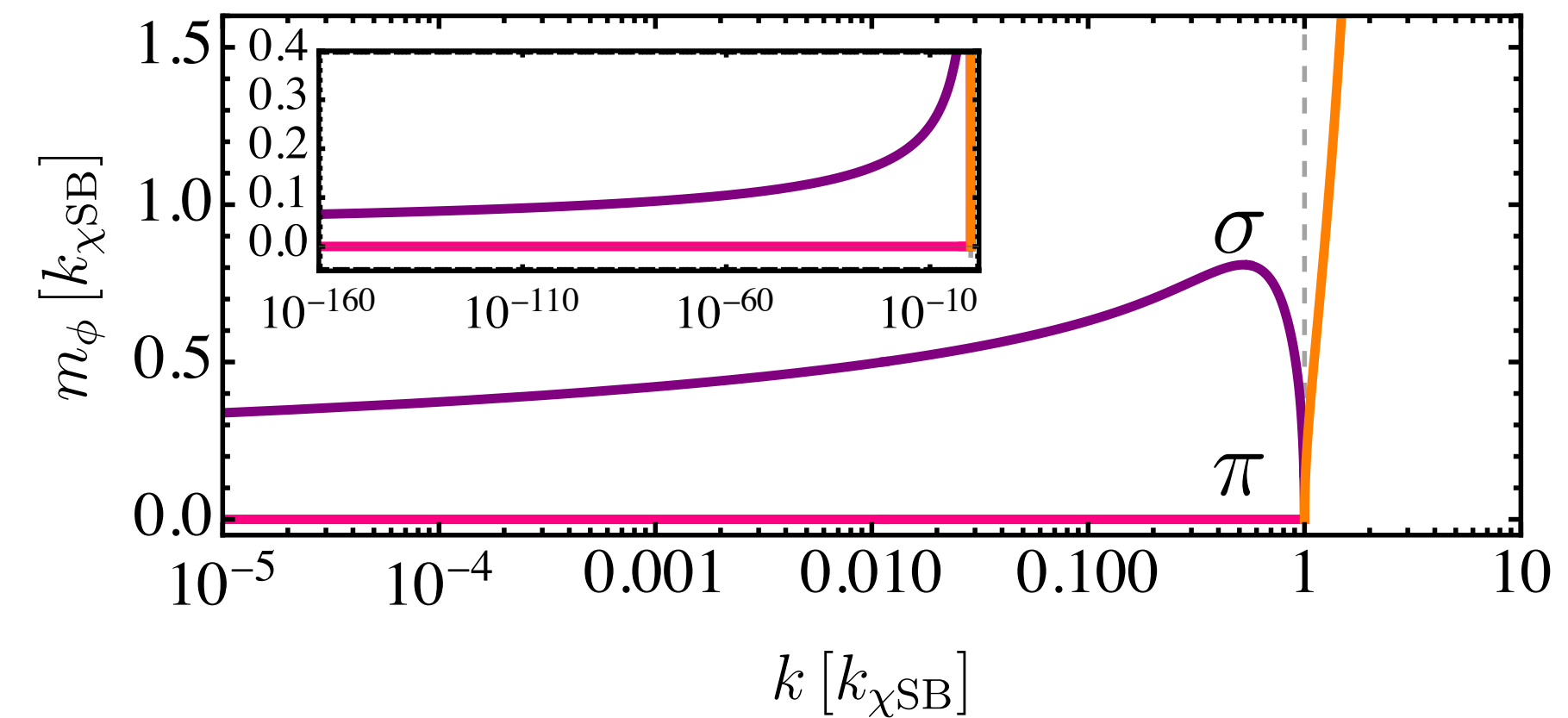
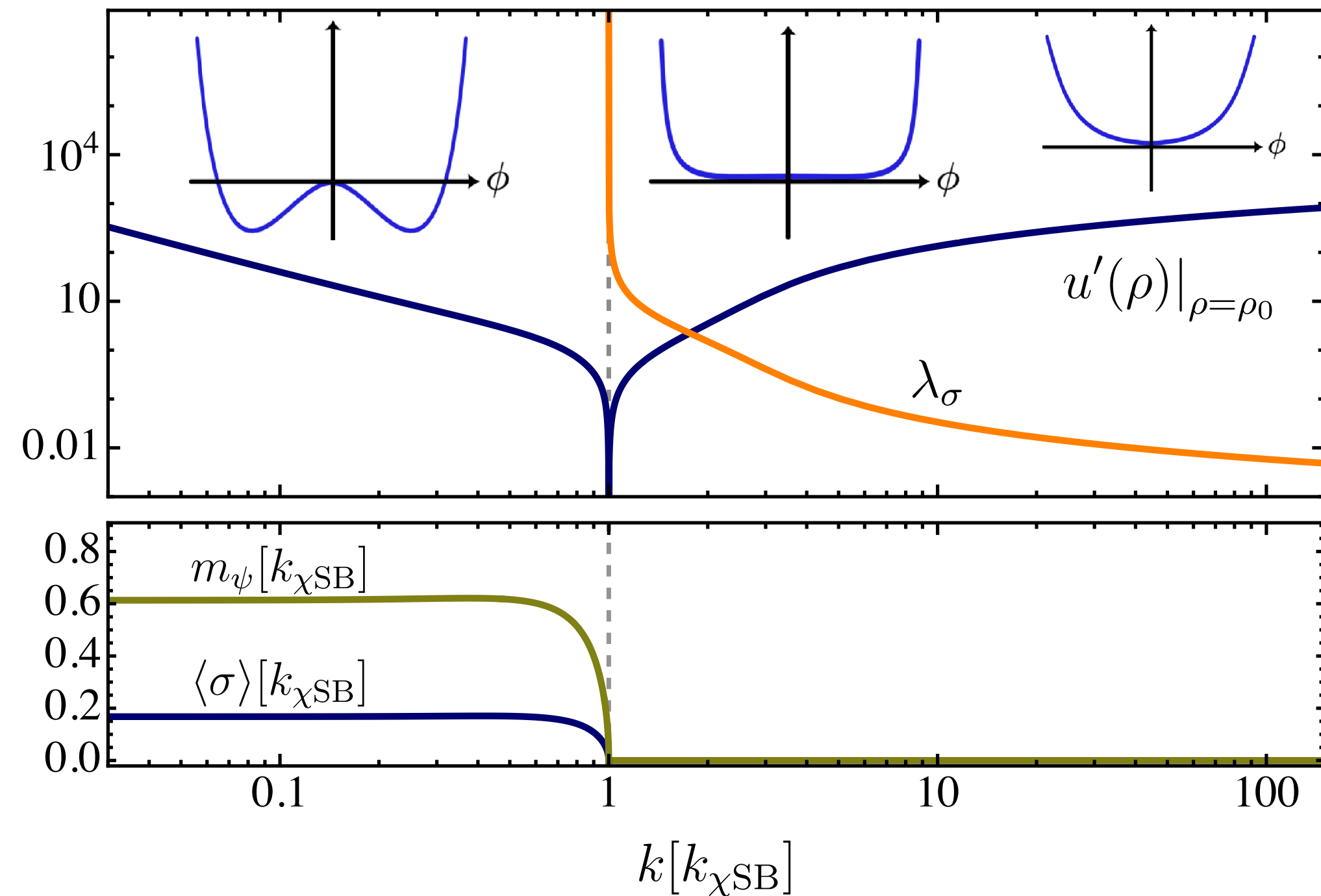
Pawlowski [hep-th/0512261] Fukushima, Pawlowski, Strodtzoff [2103.01129]

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$$+ h \bar{\psi} (T_f^0 \sigma + i\gamma_5 T_f^a \pi^a) \psi + \frac{1}{2} (\partial_\mu \phi)^2 + V(\phi^2)$$

$$V(\phi^2) = \sum_{n=1}^{N_{\max}} \frac{\lambda_n}{n!} \left(\frac{\phi^2}{2} \right)^n$$

Dynamical chiral symmetry breaking



- ◆ Flows computed: $\{h, V(\phi), Z_\psi, Z_\phi, \lambda_i\}$
- ◆ **Continuous** interpolation between chirally **symmetric** and **broken** regimes
- ◆ A **clear** and **precise** way to **diagnose** χ SB

◆ Obtaining **fundamental parameters**

- Constituent fermion masses: m_ψ
- Chiral condensate: $\langle \sigma \rangle \sim f_\pi$
- Composite masses of bosonised channels: m_σ, m_π

- ◆ Account for **higher dimensional fermionic** operators via higher-order scalar potential

Simplified truncation

◆ Purpose:

- **Practical** simplification of the pure gauge sector
- **Connection** with high-order **perturbative computations** towards the **CBZ** limit
- Implement **IR fixed-point** solutions from perturbative approaches

◆ Implementation:

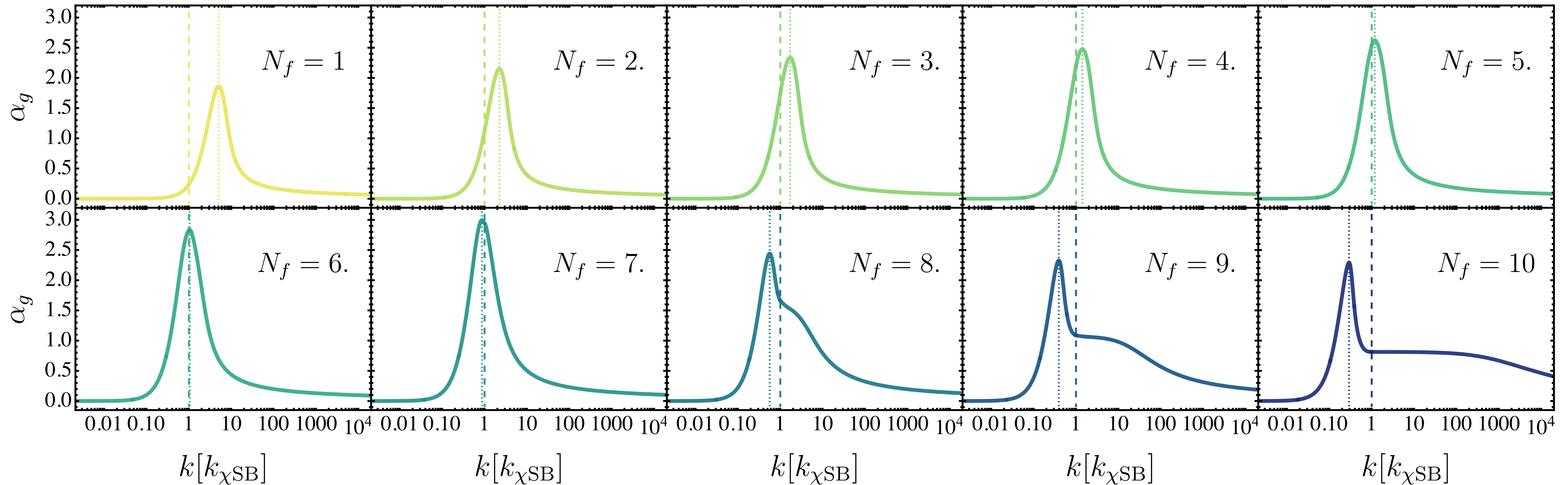
- Employ a **single adapted** gauge coupling **avatar**
- **Bosonised sector in full glory**

Smooth decoupling of dof

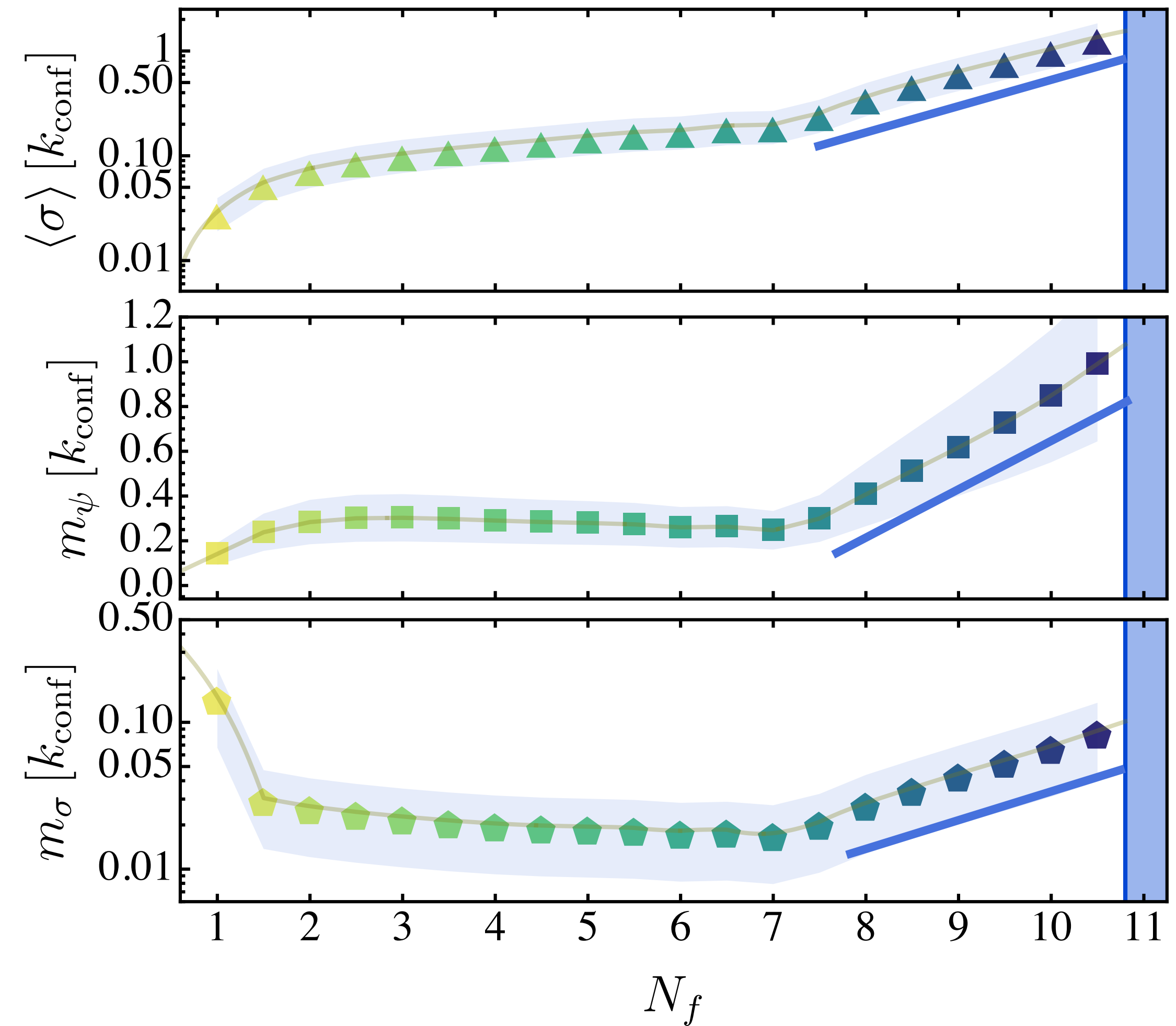
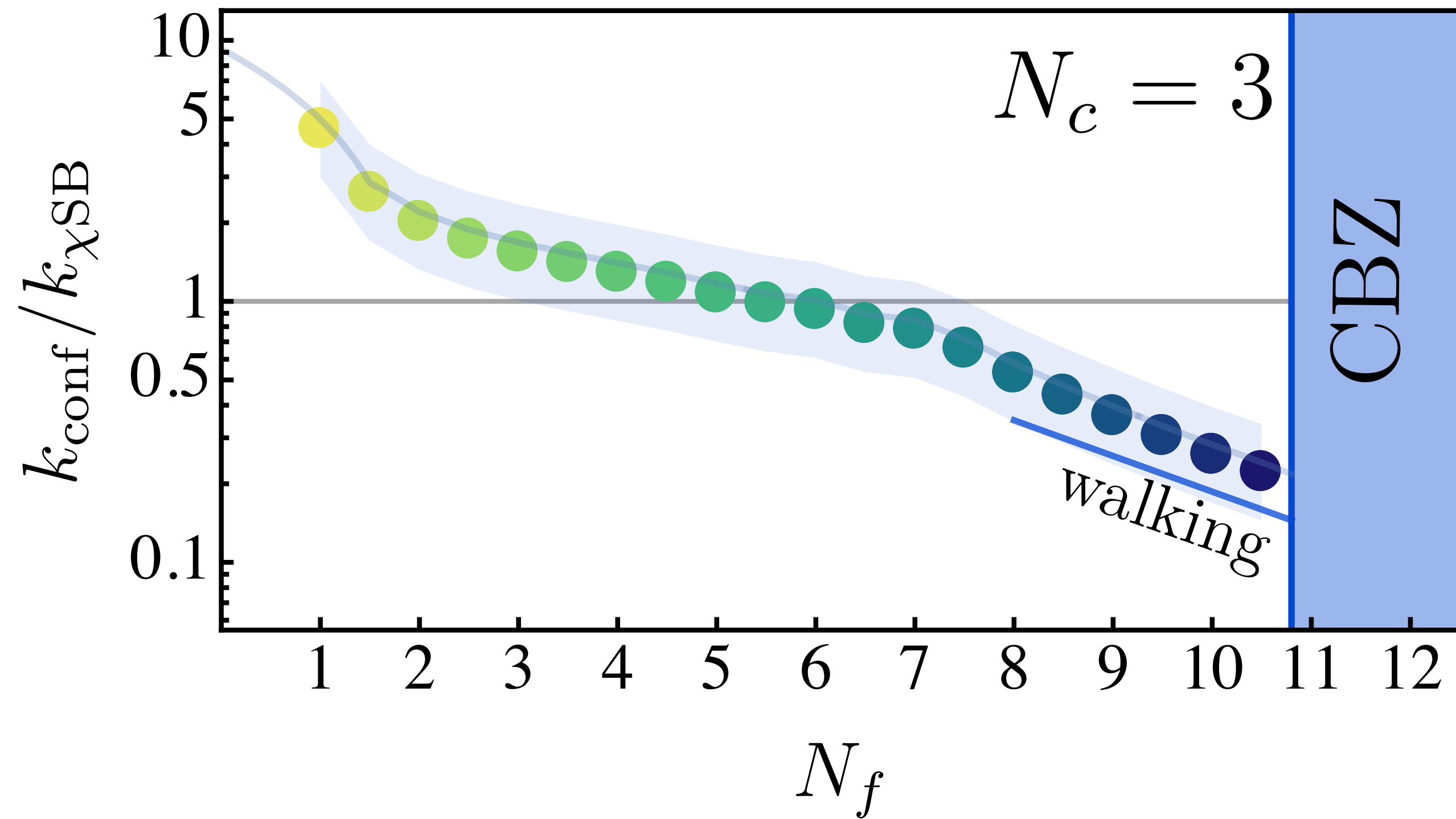
$$\partial_t g = \underbrace{\partial_t g_{A\bar{\psi}\psi}^{(1)}}_{\text{Non-perturbative}} + \frac{\partial_t g_{A\bar{\psi}\psi}^{(>1)}}{\partial_t g_{A\bar{\psi}\psi}^{(>1)} \Big|_{\{\bar{m}_A, \bar{m}_\psi\}=0}} \sum_{n=2}^{N_{\max}} \left\{ \sum_{i=1}^{i_{\max}} \frac{1}{(1+i_{\max})} \frac{1}{(1+\bar{m}_A^2)^{3n+i}} \right\} \left(\underbrace{\beta_A^{(n)}}_{\text{Perturbative}} + \left[\frac{(2+\bar{m}_\psi^2)}{2(1+\bar{m}_\psi^2)^3} \right]^{N_f \text{ in loop}} \underbrace{\beta_\psi^{(n)}} \right)$$

* Fundamental parameters in good agreement with best truncations

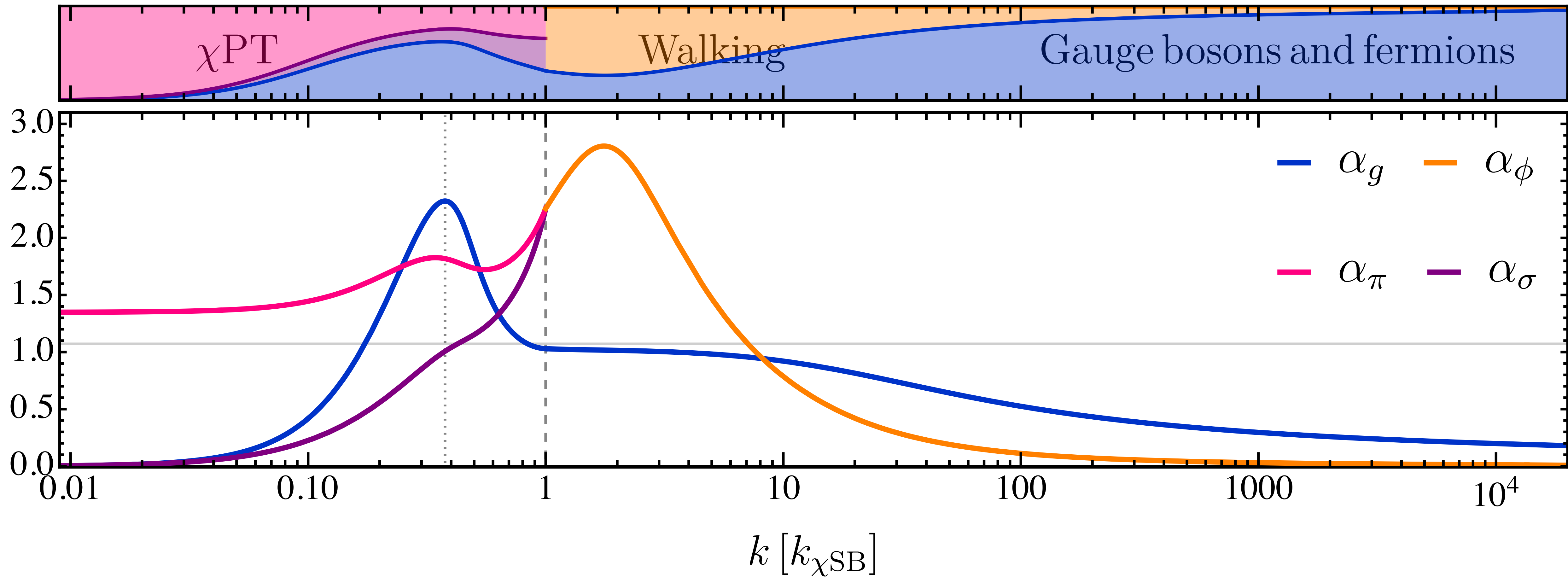
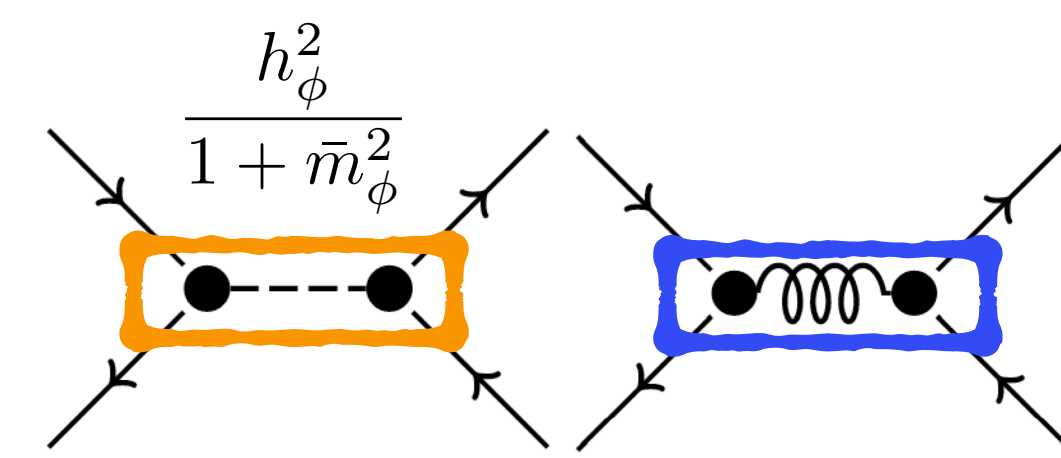
Many flavour gauge dynamics



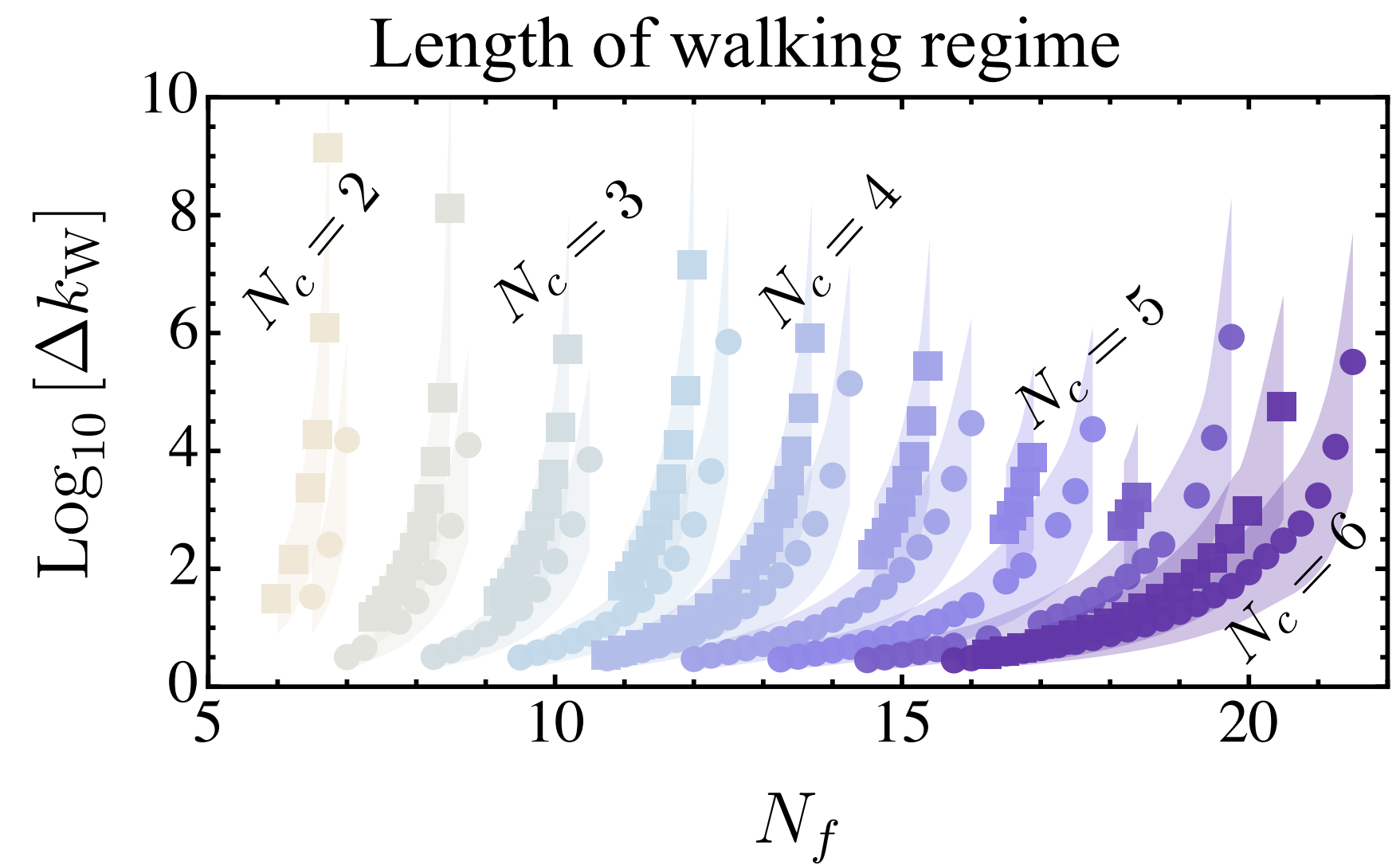
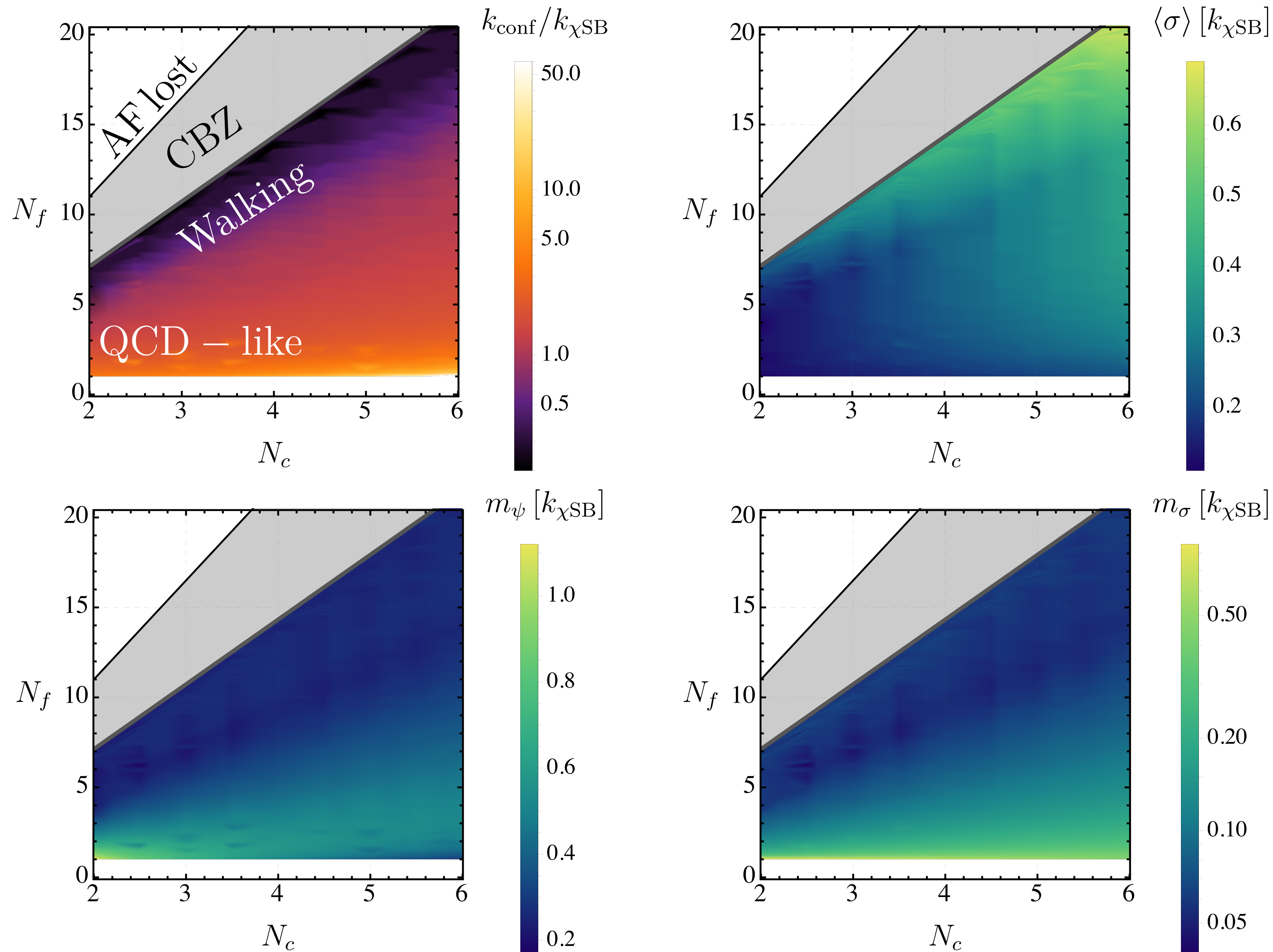
Interplay of scales and fundamental parameters



Walking dynamics and flow over EFTs

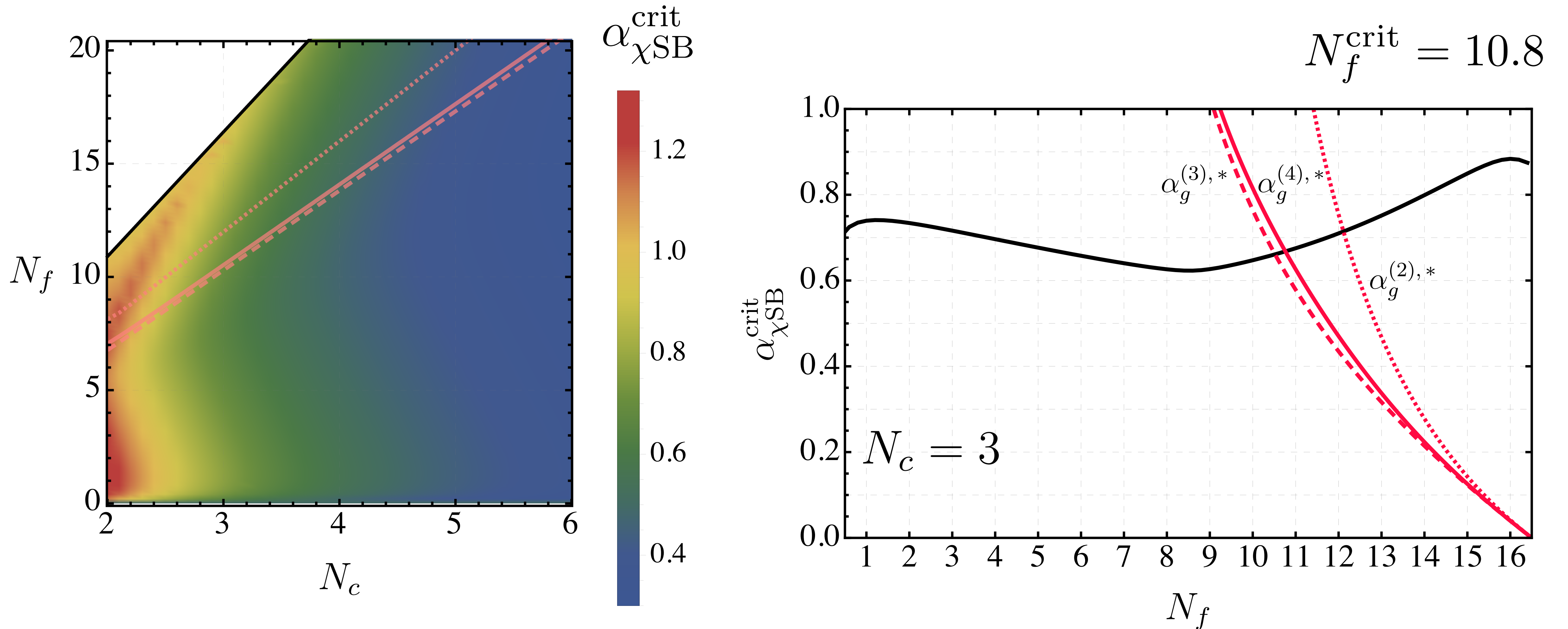


Cartography



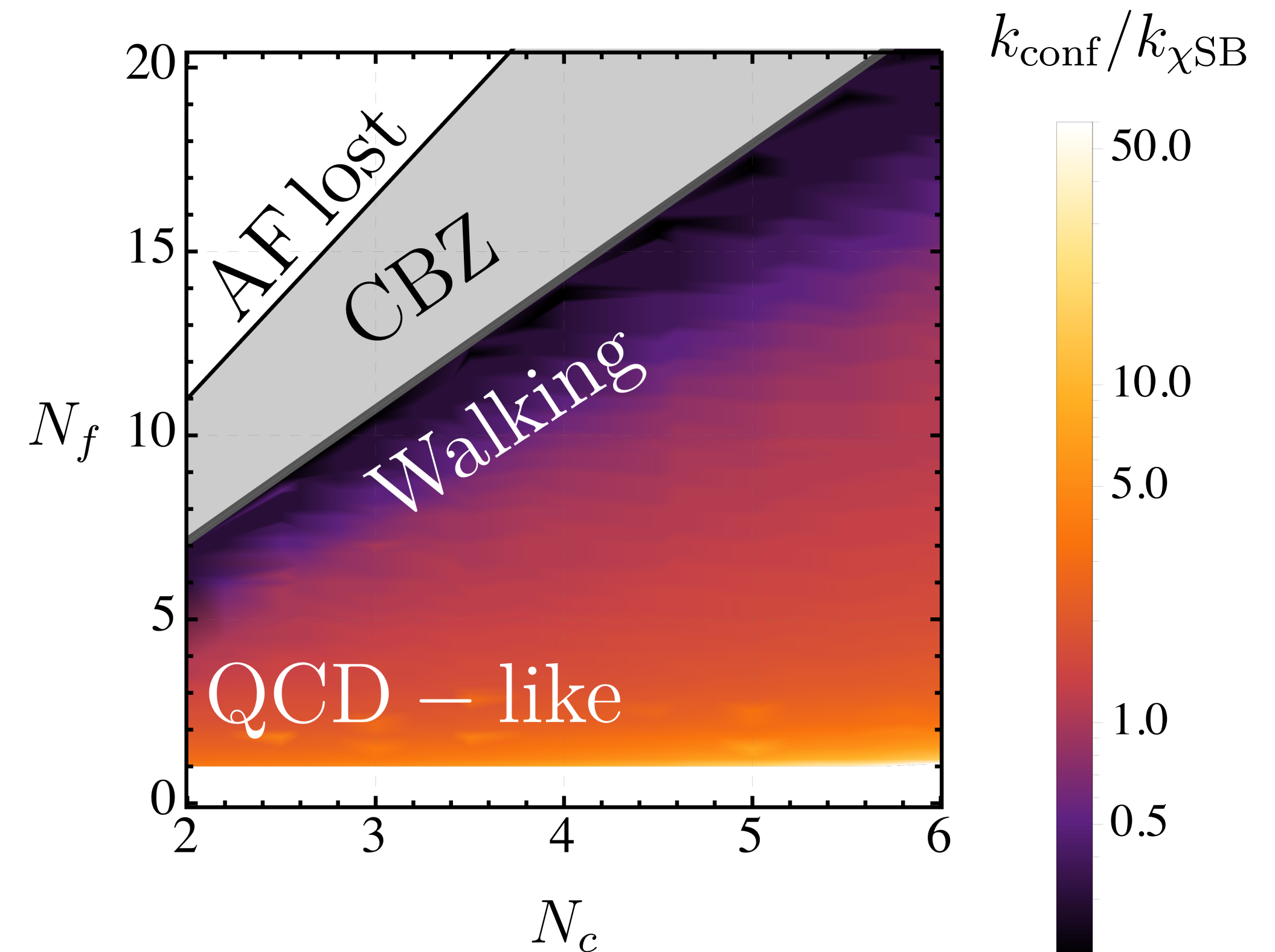
Boundary of the conformal window

$\alpha_{\chi\text{SB}}^{\text{crit}}$: Critical strength of **gauge dynamics** necessary to trigger dynamically χSB



Summary and conclusions

- ◆ The fRG is a **powerful approach** to strong dynamics:
 - Qualitatively new **treatment of confinement**
 - Charted the **many-flavour landscape** from **first principles**
 - Determination of **fundamental parameters**:
 m_ψ , $\langle\sigma\rangle$, m_σ , size of walking regime, ...
 - Provided a new **prediction** for the **lower boundary** of the CBZ window



Outlook

♦ A large number of **applications**:

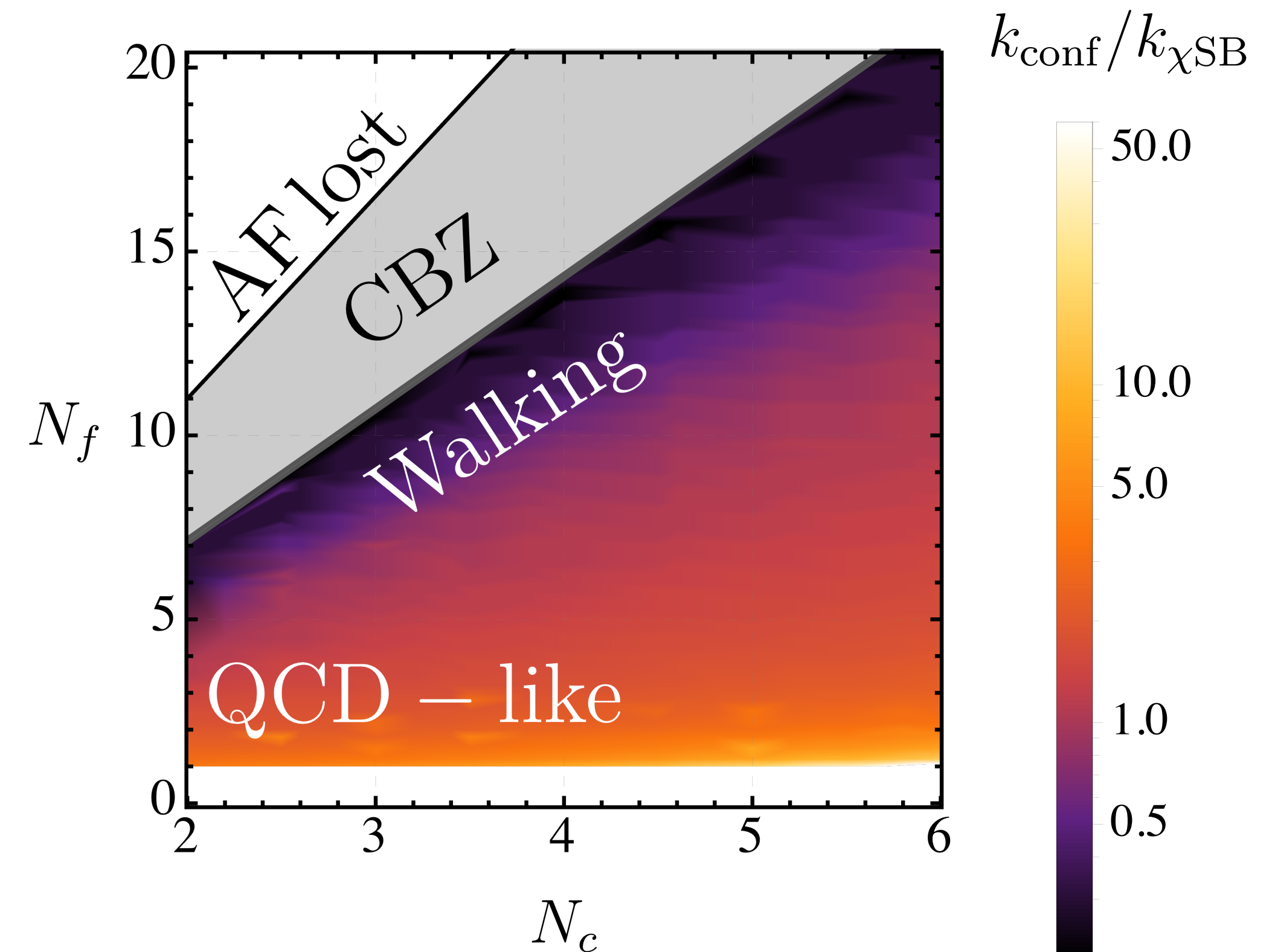
- **Ph: Solving the Dark Matter - Baryon Coincidence:**

Exploring Theories of Dark-Color Unification (with Yi Chung, Florian Goertz and Aika Tada)

- **Th:** Dynamical symmetry breaking in the **Bars-Yankielowicz** model (with Shahram Vahani and Haolin Li)

- **Th-Ph:** Phase diagram of **scalar-QCD** theories (with Maciej Kierkla, Masatoshi Yamada and Jisuke Kubo)

- ...



Outlook

♦ A large number of **applications**:

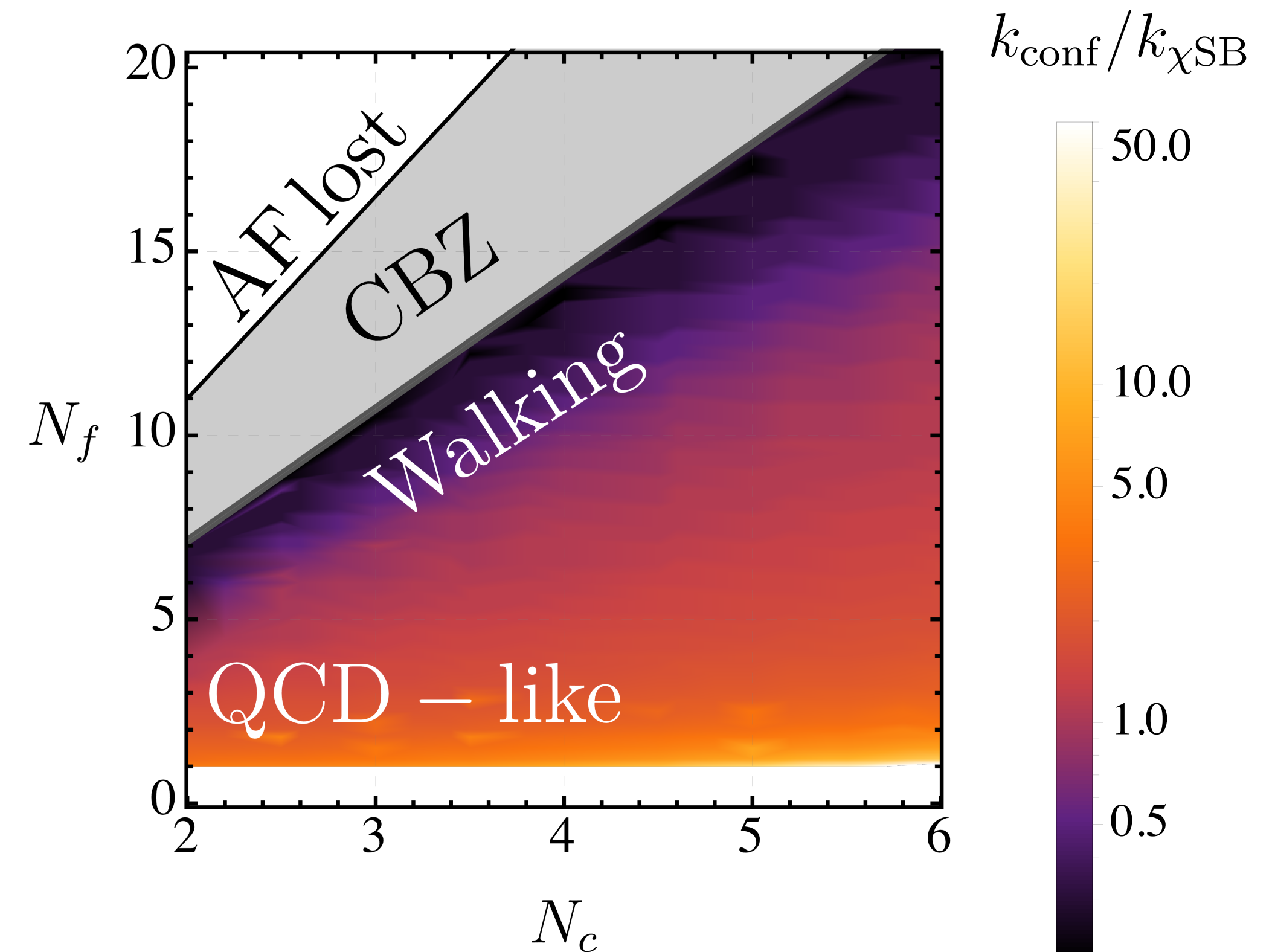
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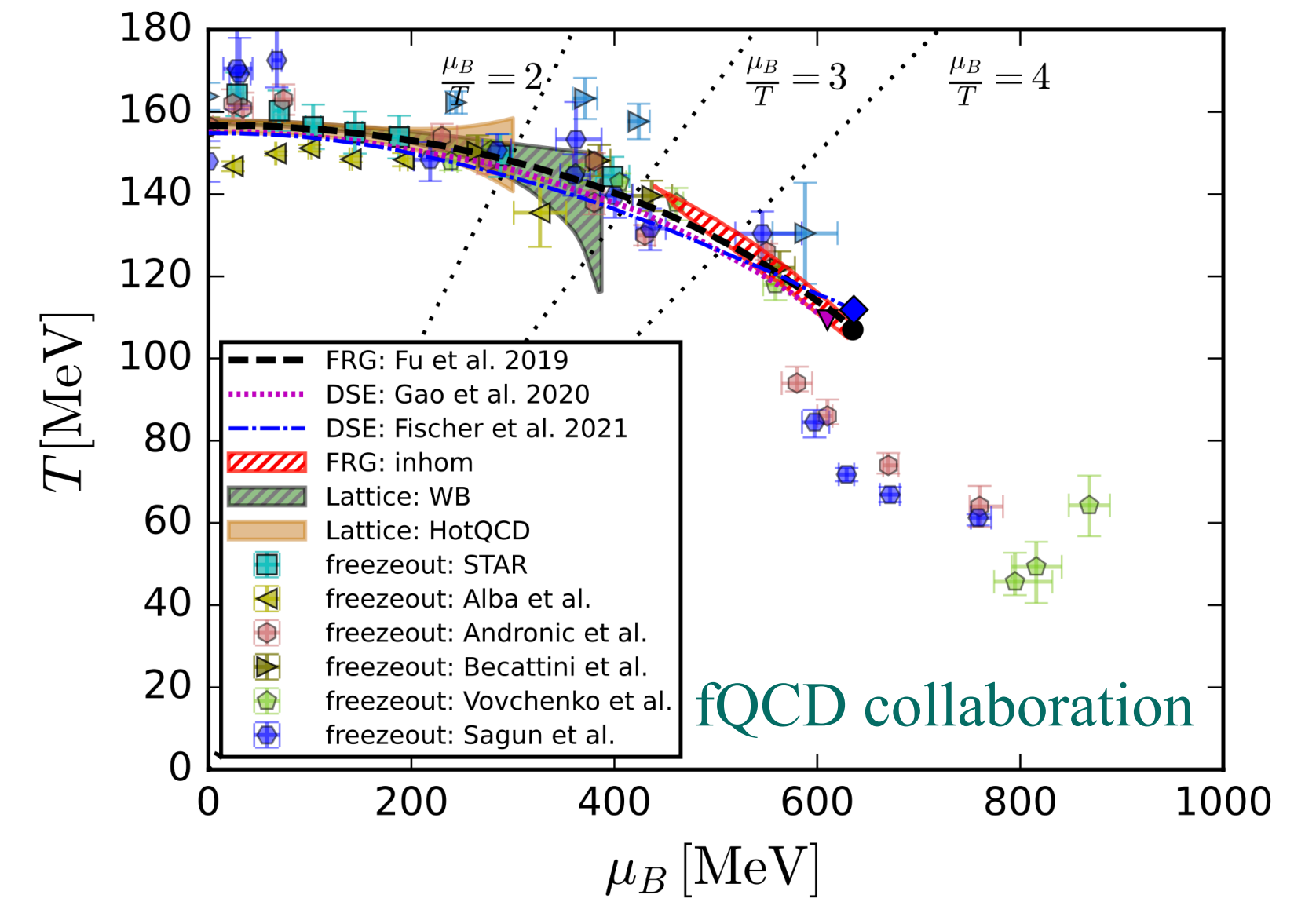
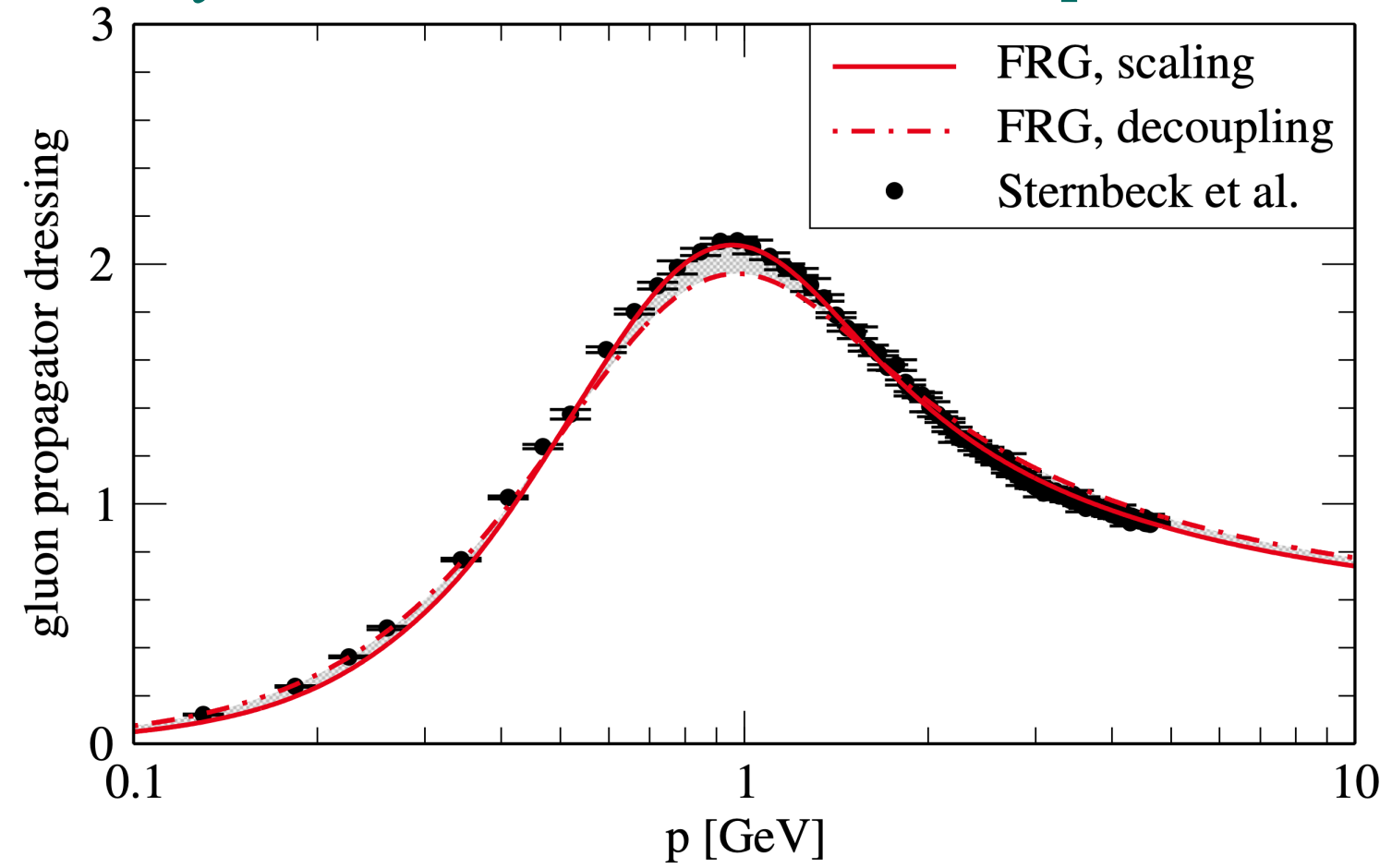


Thank you for your attention!

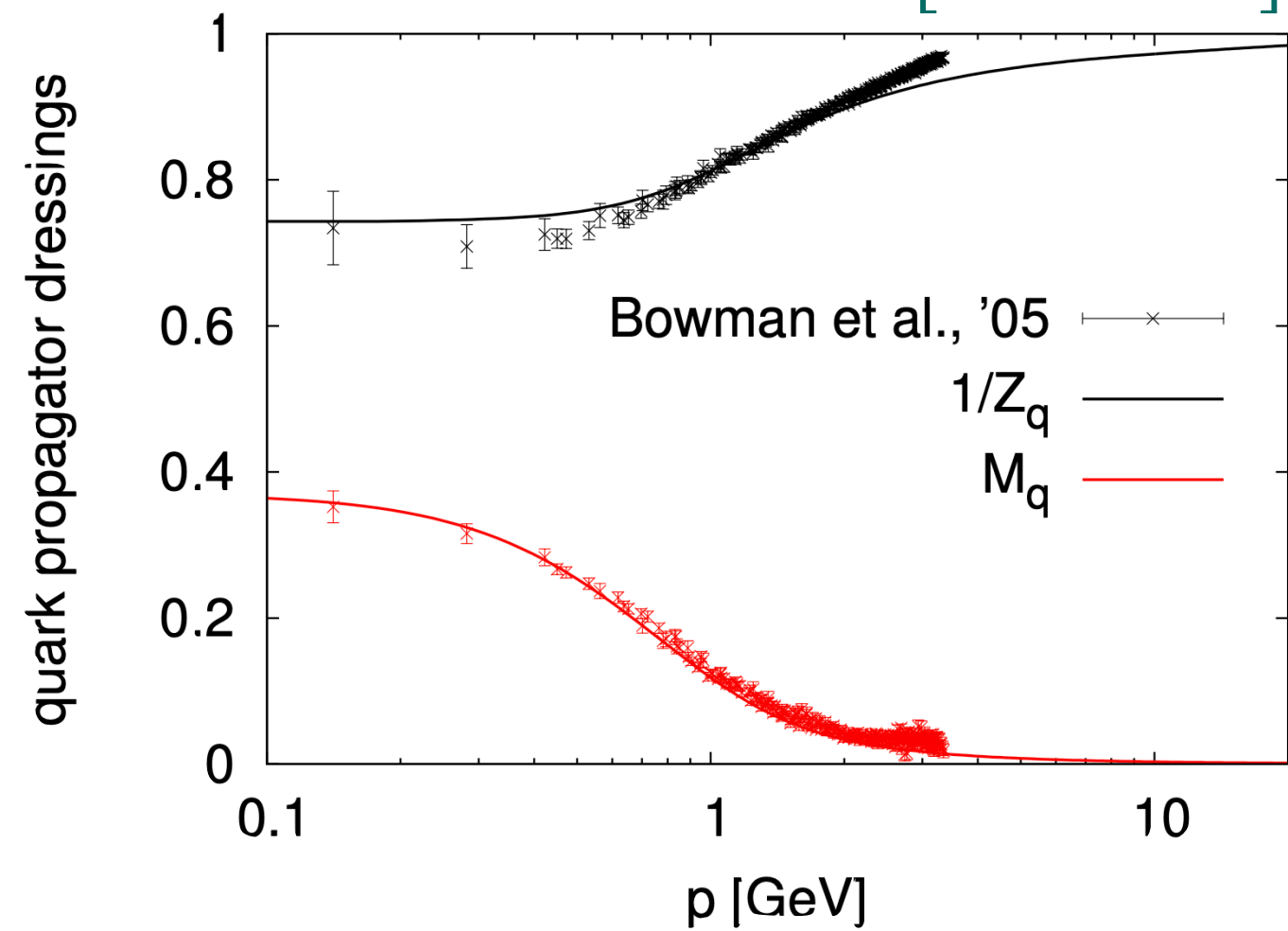
Additional slides

Functional Renormalisation Group in QCD

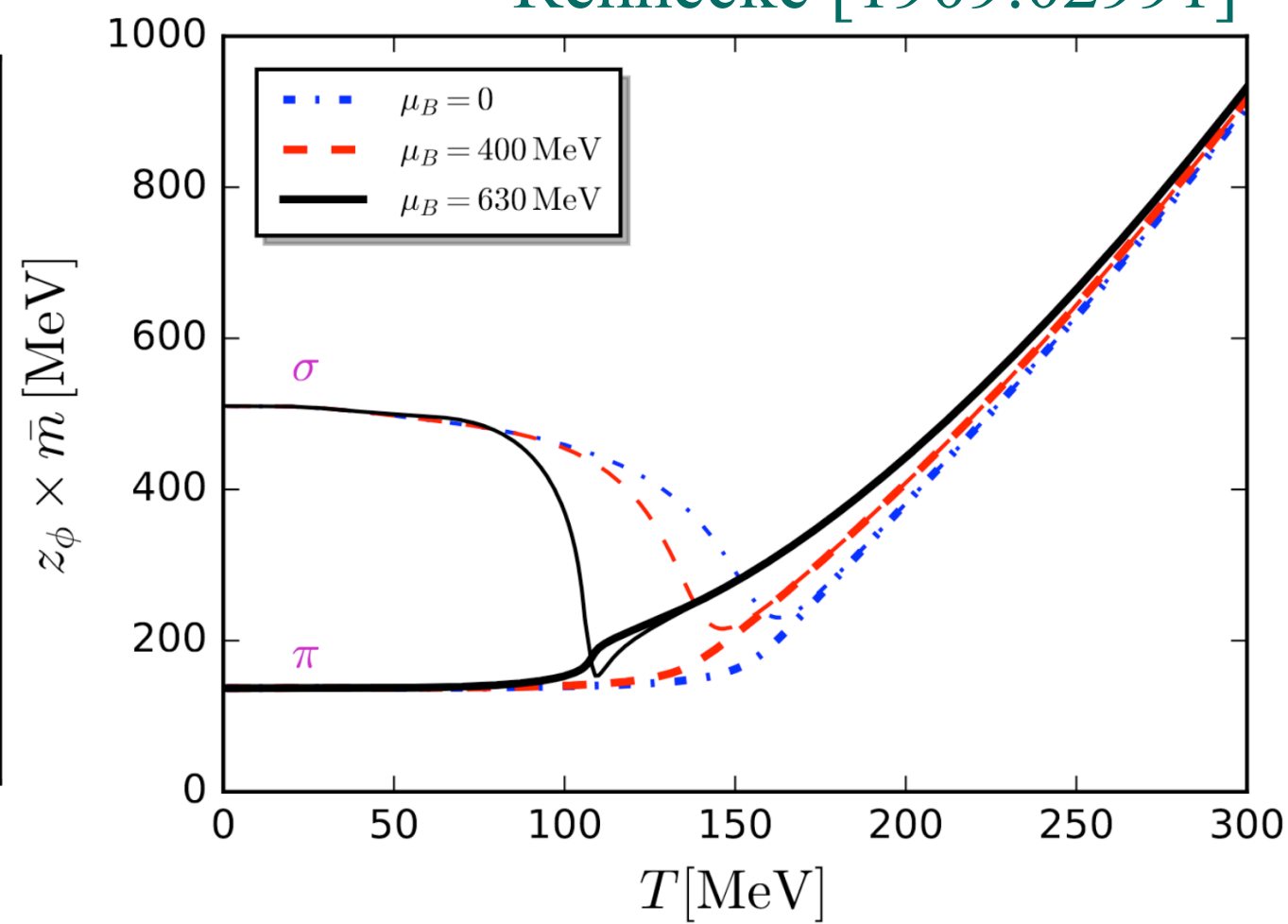
Cyrol, Fister, Mitter, Pawłowski [1605.01856]



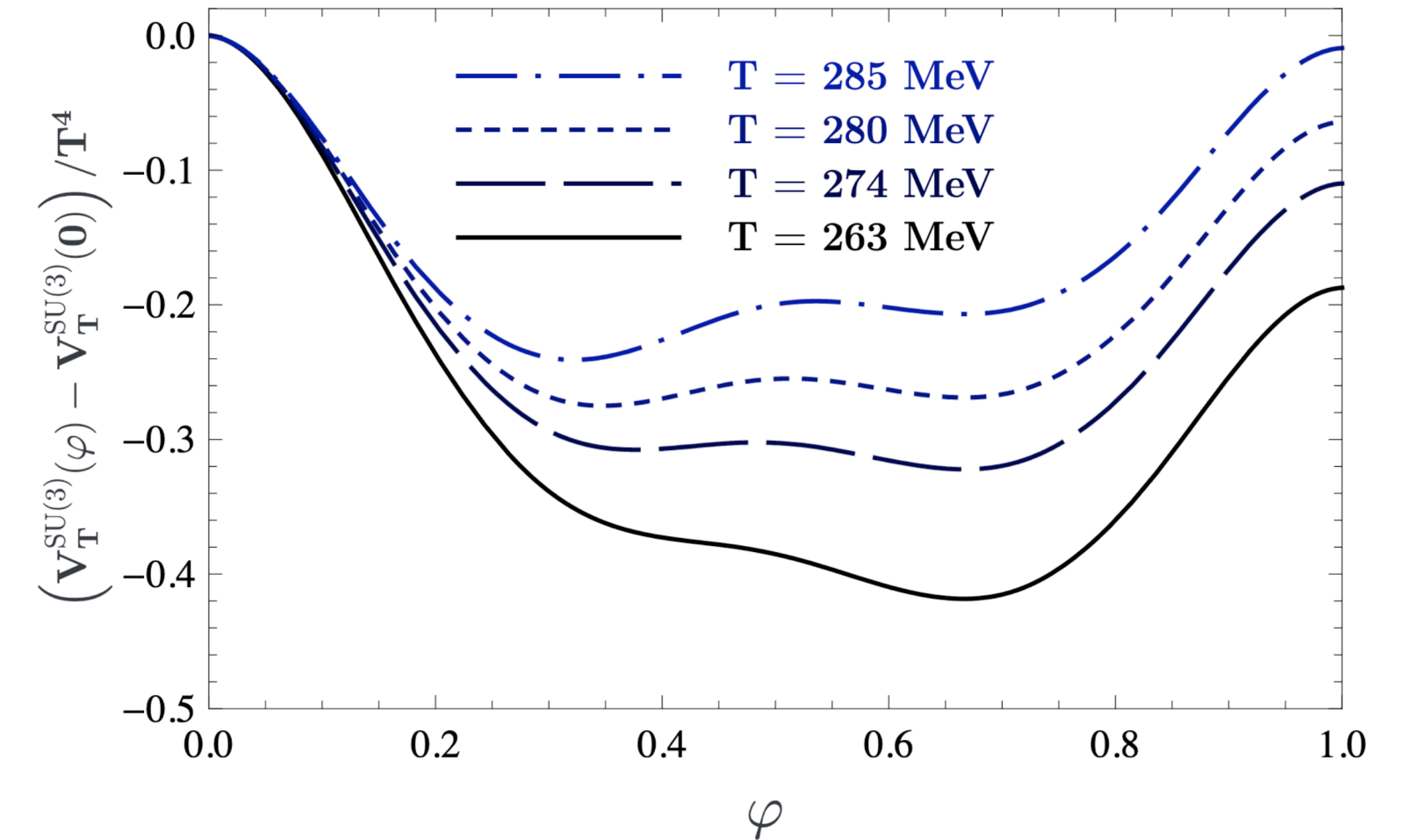
Mitter, Pawłowski, Strothoff [1411.7978]



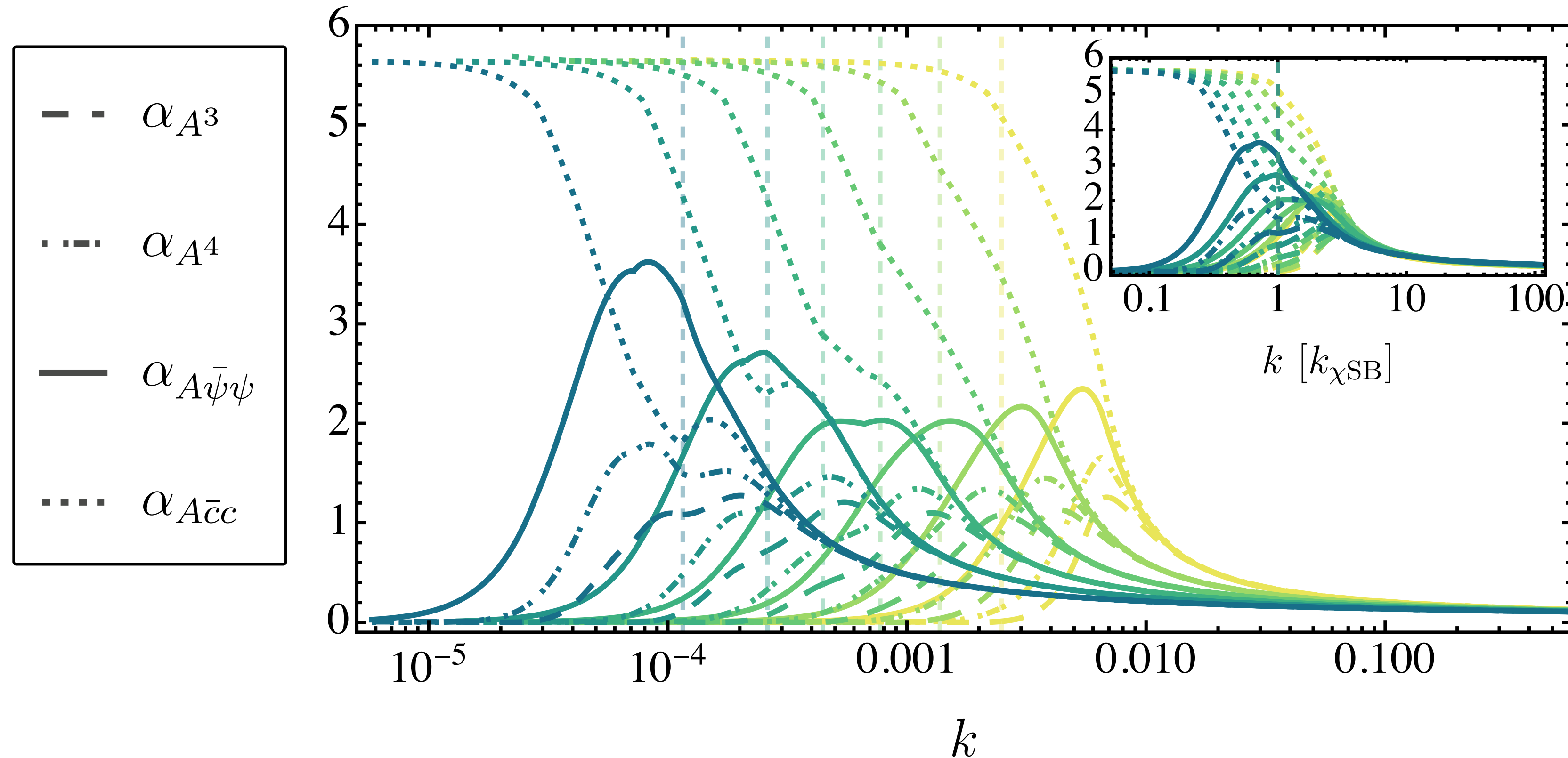
Fu, Pawłowski, Rennecke [1909.02991]



Fister, Pawłowski [1301.4163]



Many flavour full truncation



◆ Observations:

- The peak of the pure-gluon vertex is rather constant and independent of the number of flavours
- The ghost-gauge coupling in the deep IR always reaches the same constant coupling given it is only colour-group dependent
- Increment of gauge-chiral coupling strength. The stronger value of this avatar causes a flattening of the chiral potential and consequently an earlier χ SB.

Yang-Mills phases and confinement

$$\Gamma_k^{(AA)}(p^2) = Z_{A,k}(p) (p^2 + m_{A,k}^2)$$

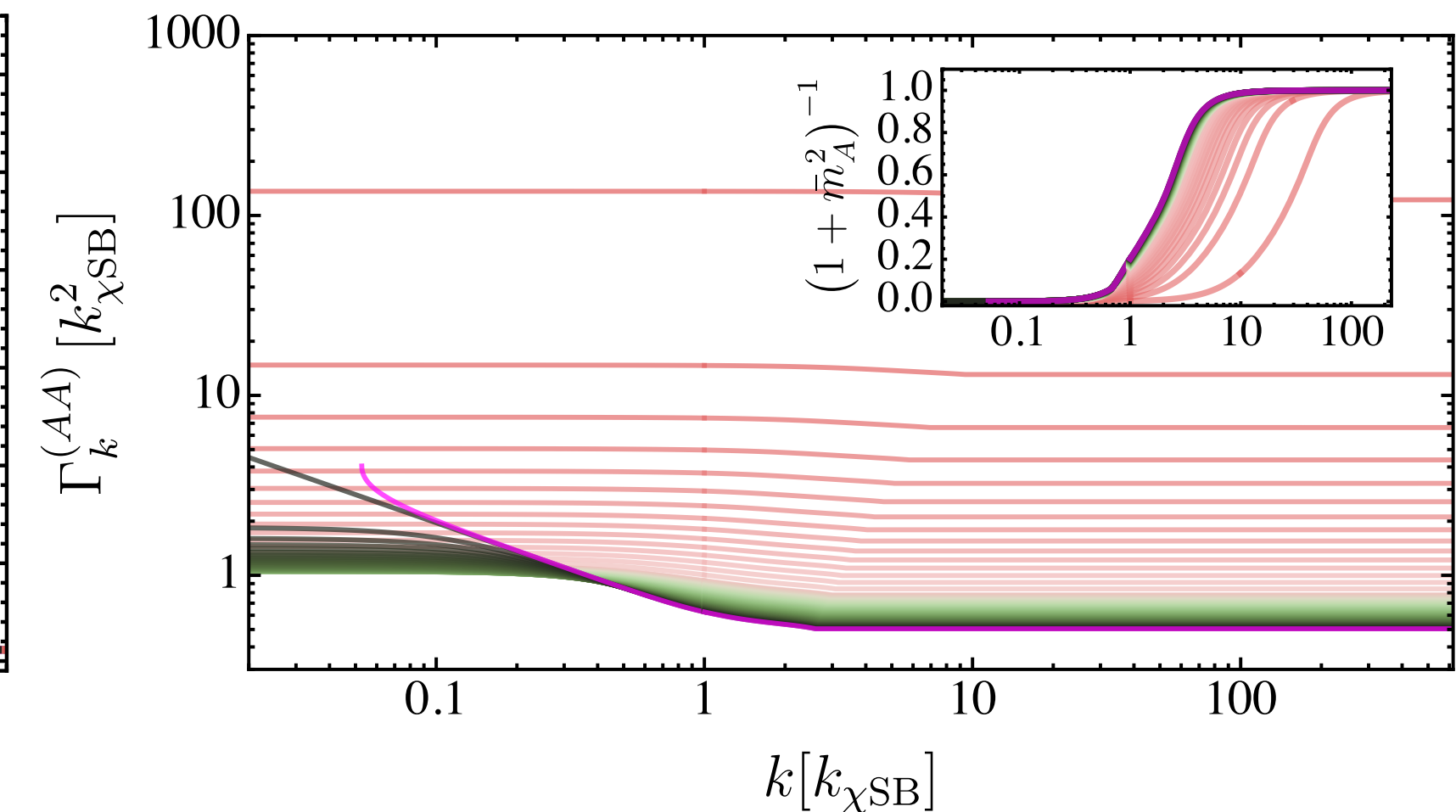
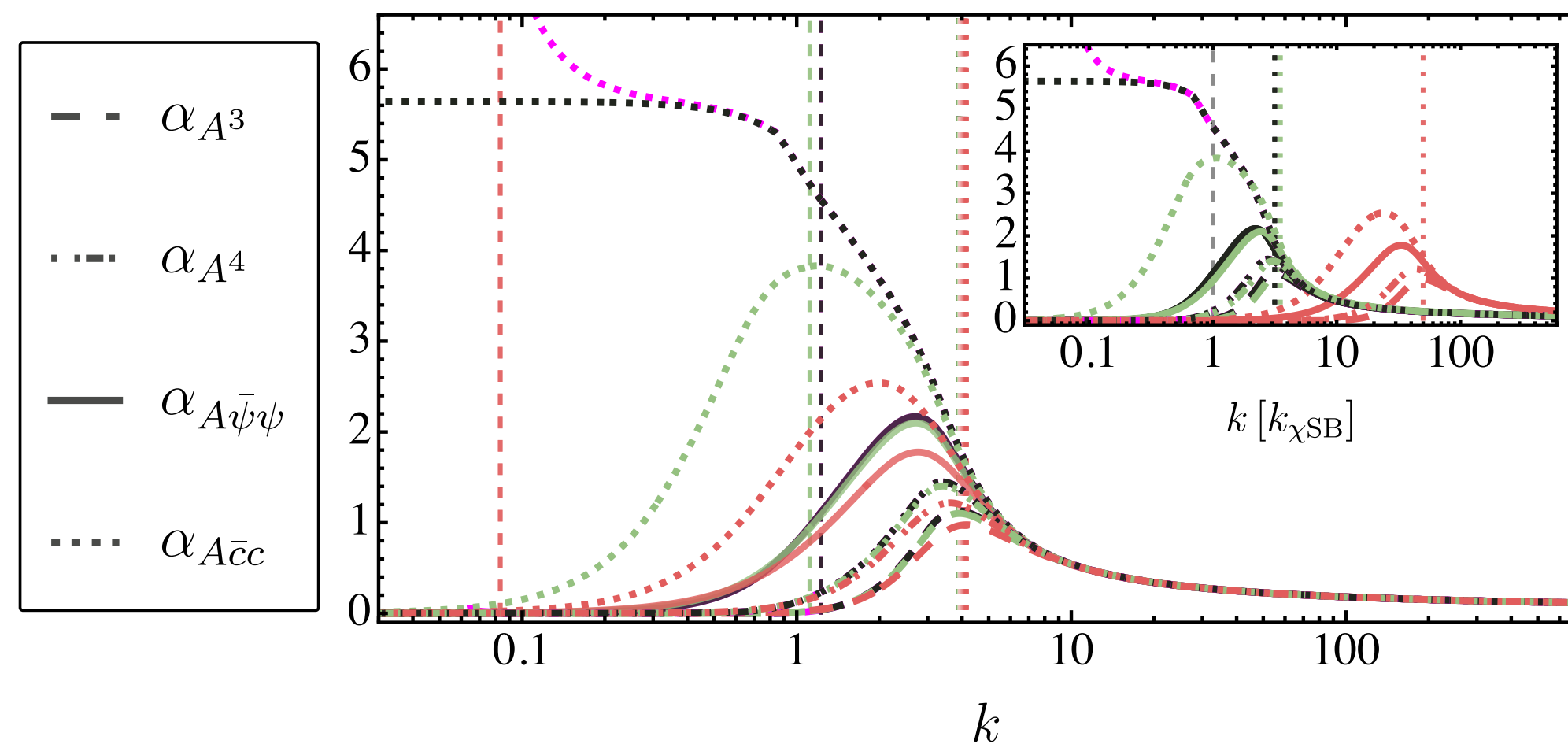
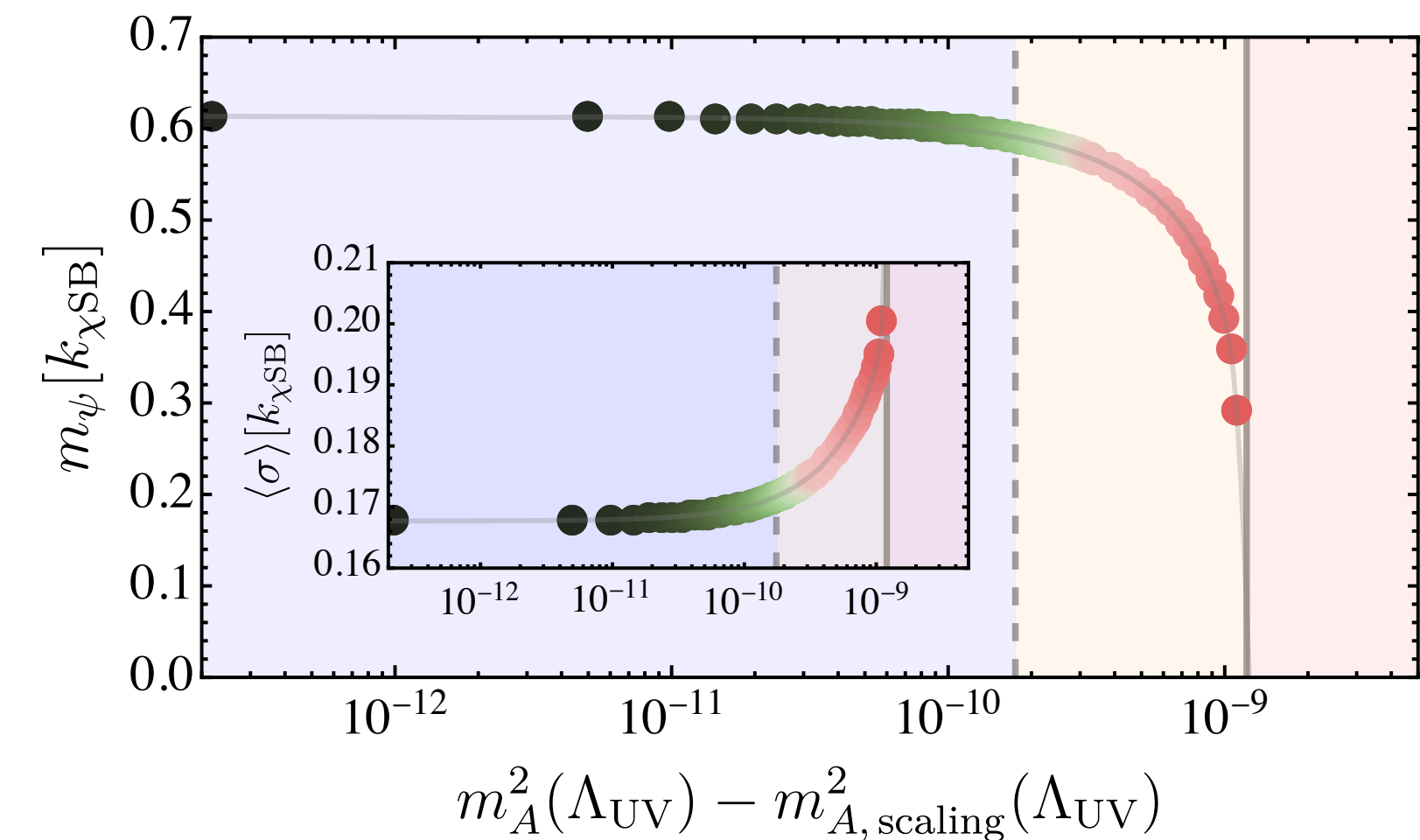
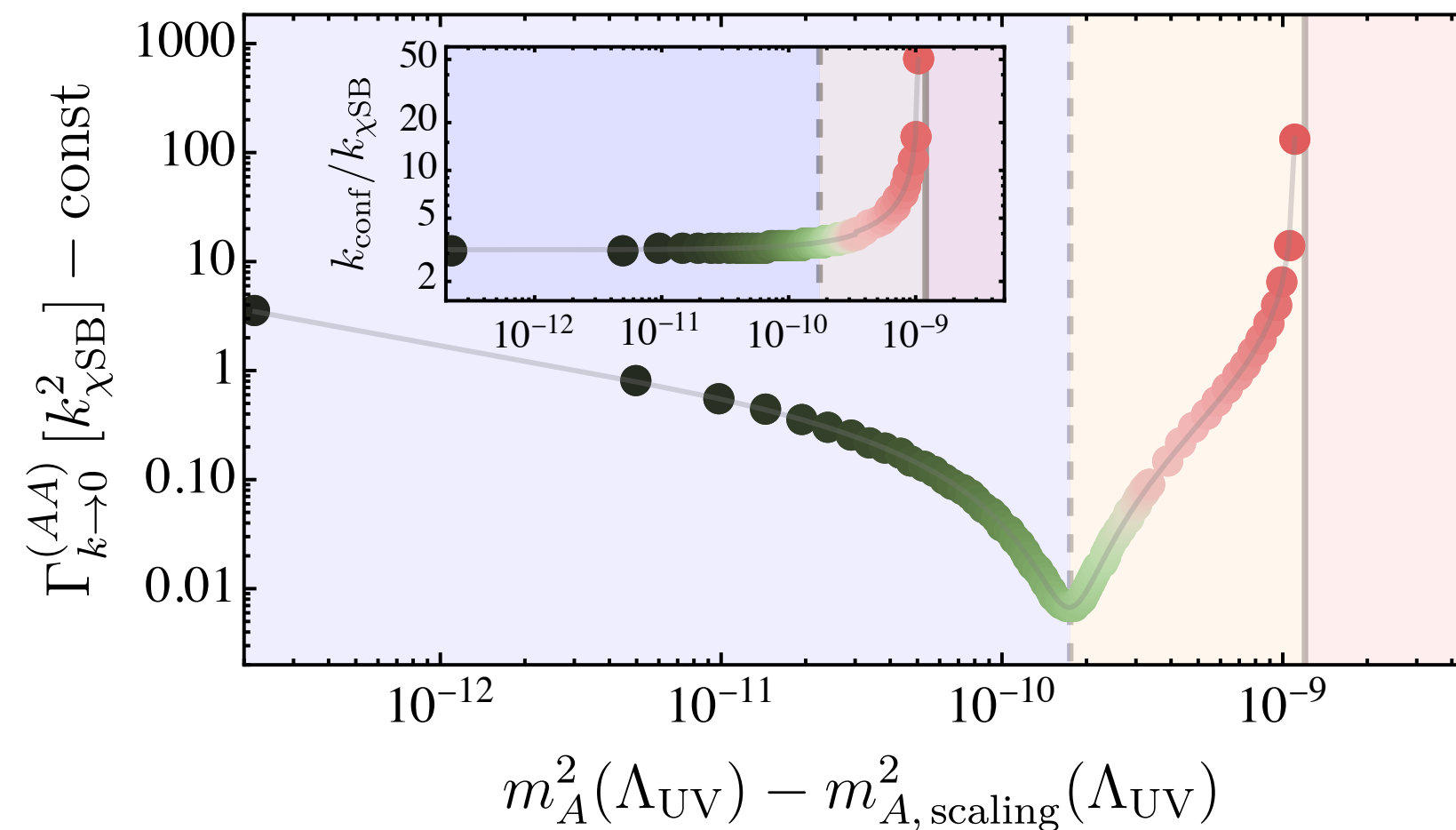
◆ Mass gap as a dialing parameter

◆ Confining solutions:

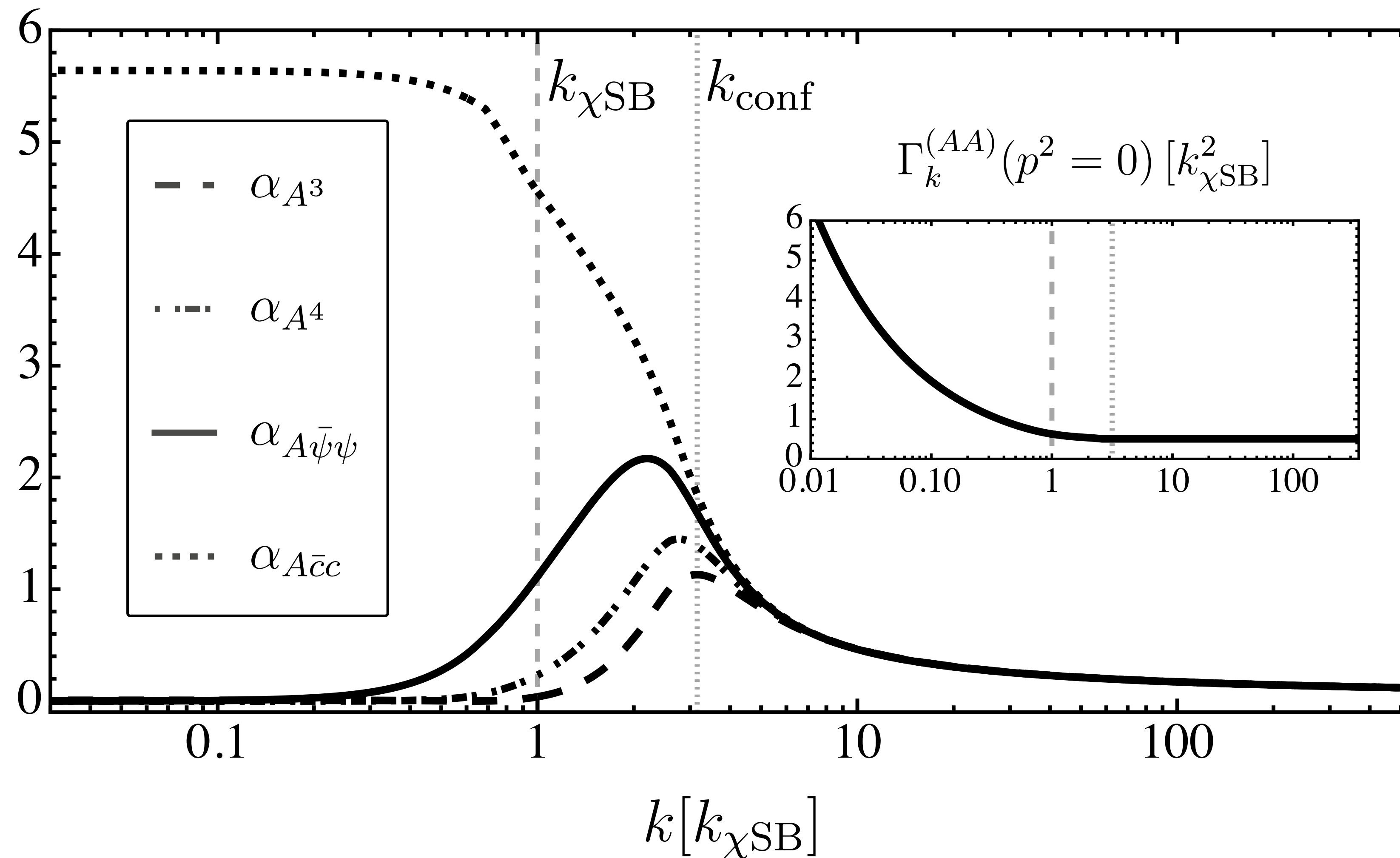
- Scaling
- Decoupling

◆ Non-confining solutions:

- Coulomb phase
- Massive Yang-Mills



Criticality of low N_f



Conformality and scale invariance

$$S = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} + \bar{\psi} \not{D} \psi$$

Conformality and scale invariance

$$S = \int_x \left[\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} \right]_{\mathcal{G}_{\text{TC}}} + \bar{\psi} \not{D} \psi \Big|_{\text{SU}(N_f)_L \times \text{SU}(N_f)_R}$$

Conformality and scale invariance

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$\text{SU}(N_f)_L \times \text{SU}(N_f)_R$

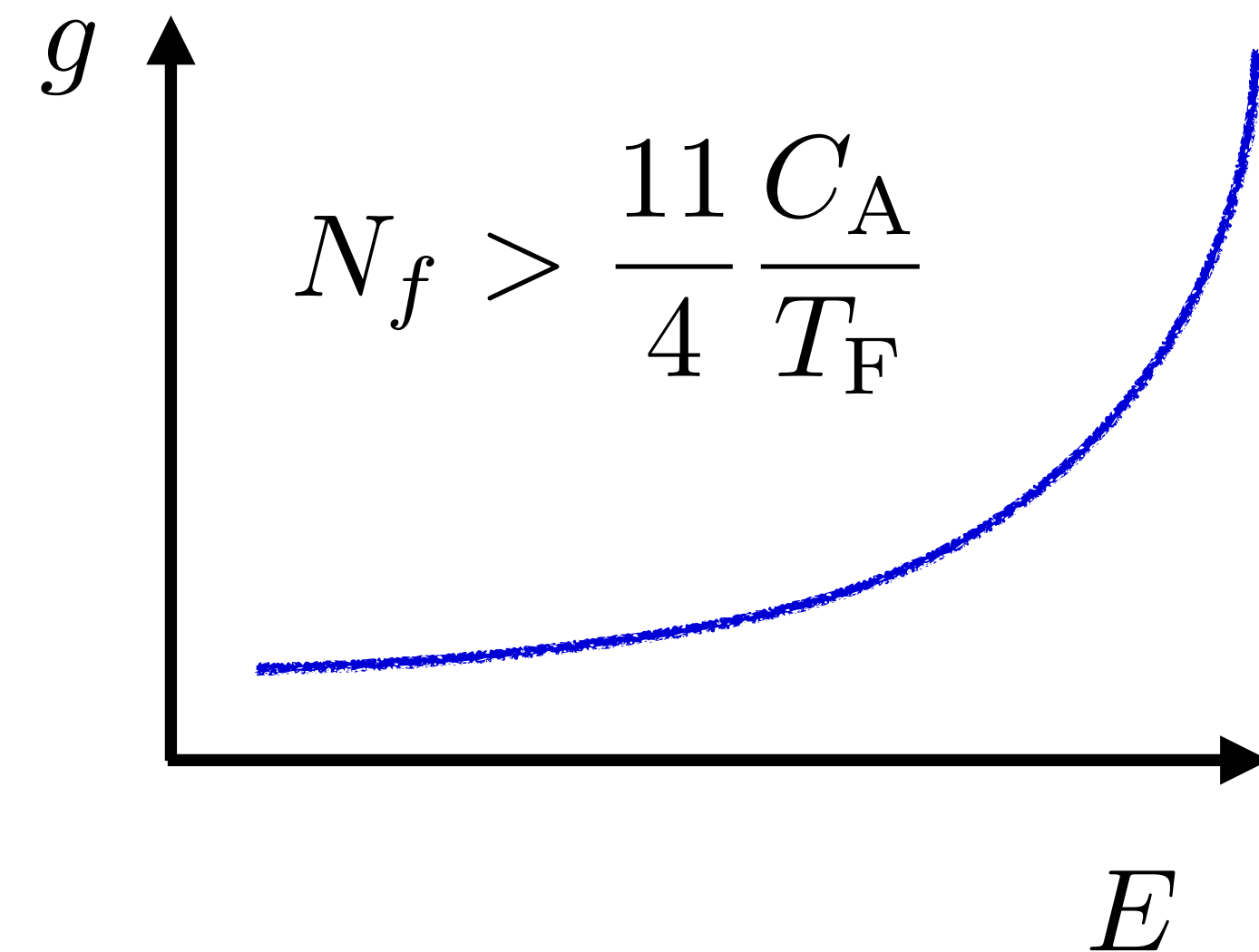
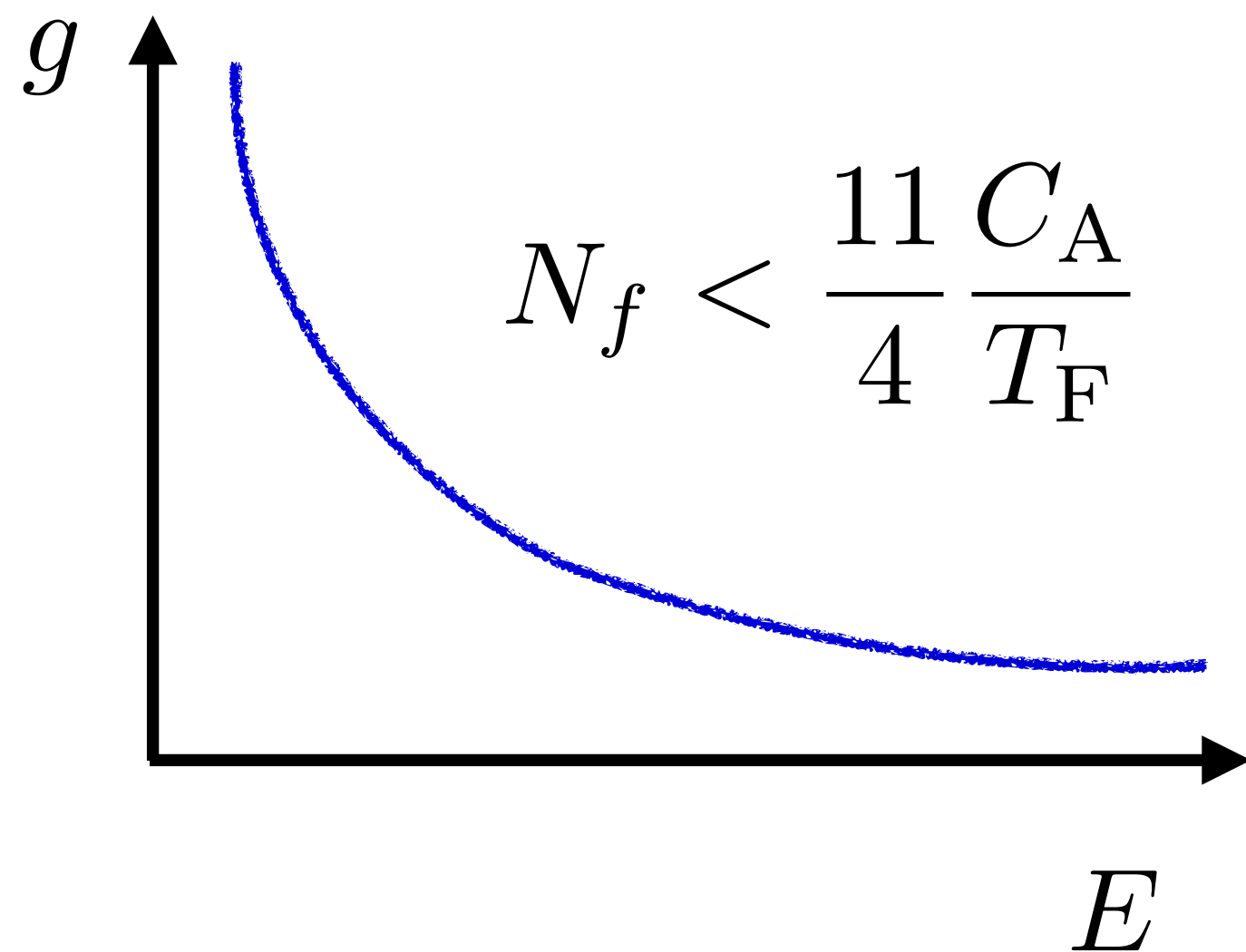
$$\rightarrow \beta_g = -\frac{g^3}{(4\pi)^2} \left(\frac{11}{3} C_A - \frac{4}{3} T_F N_f \right)$$

Conformality and scale invariance

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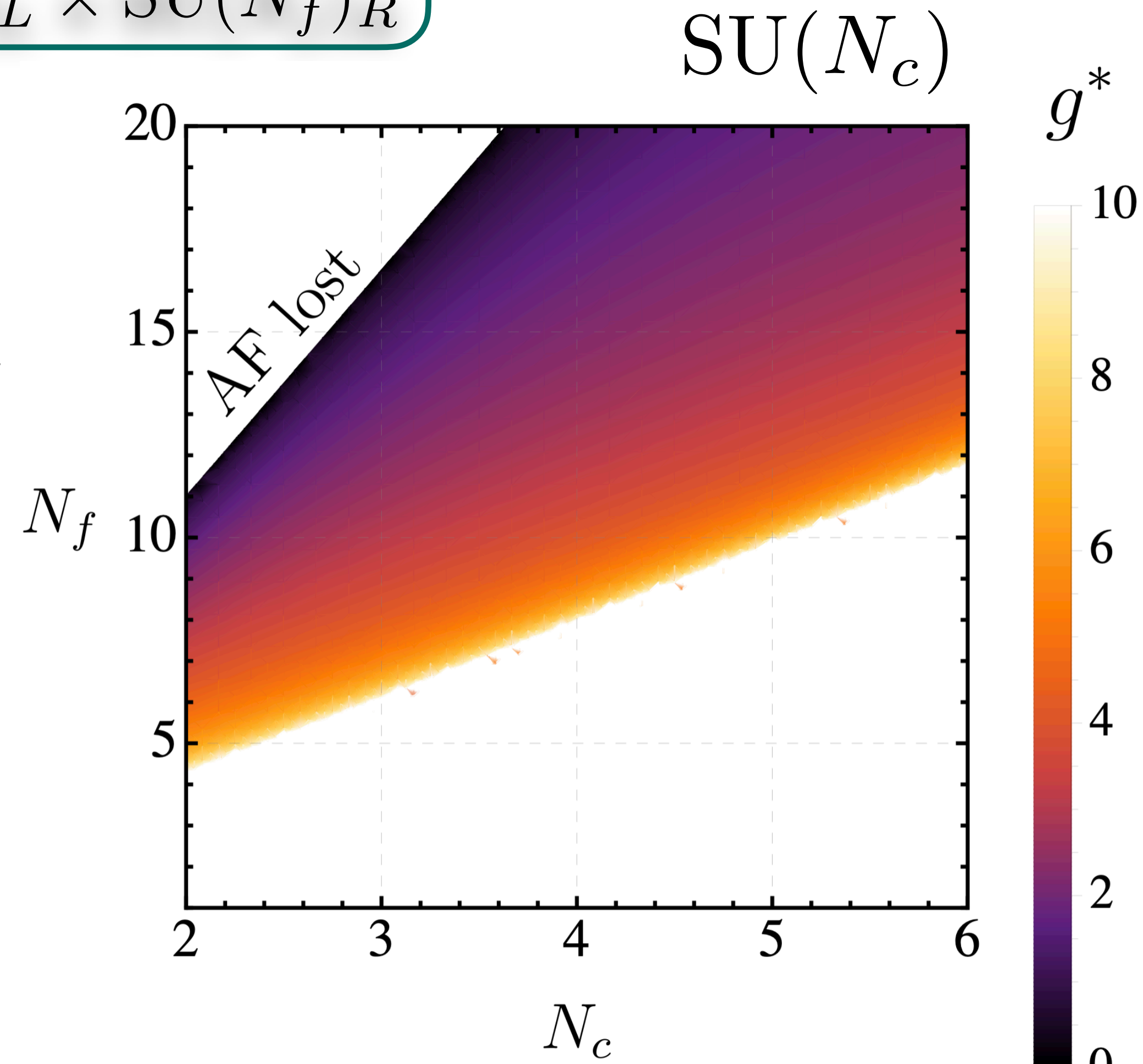
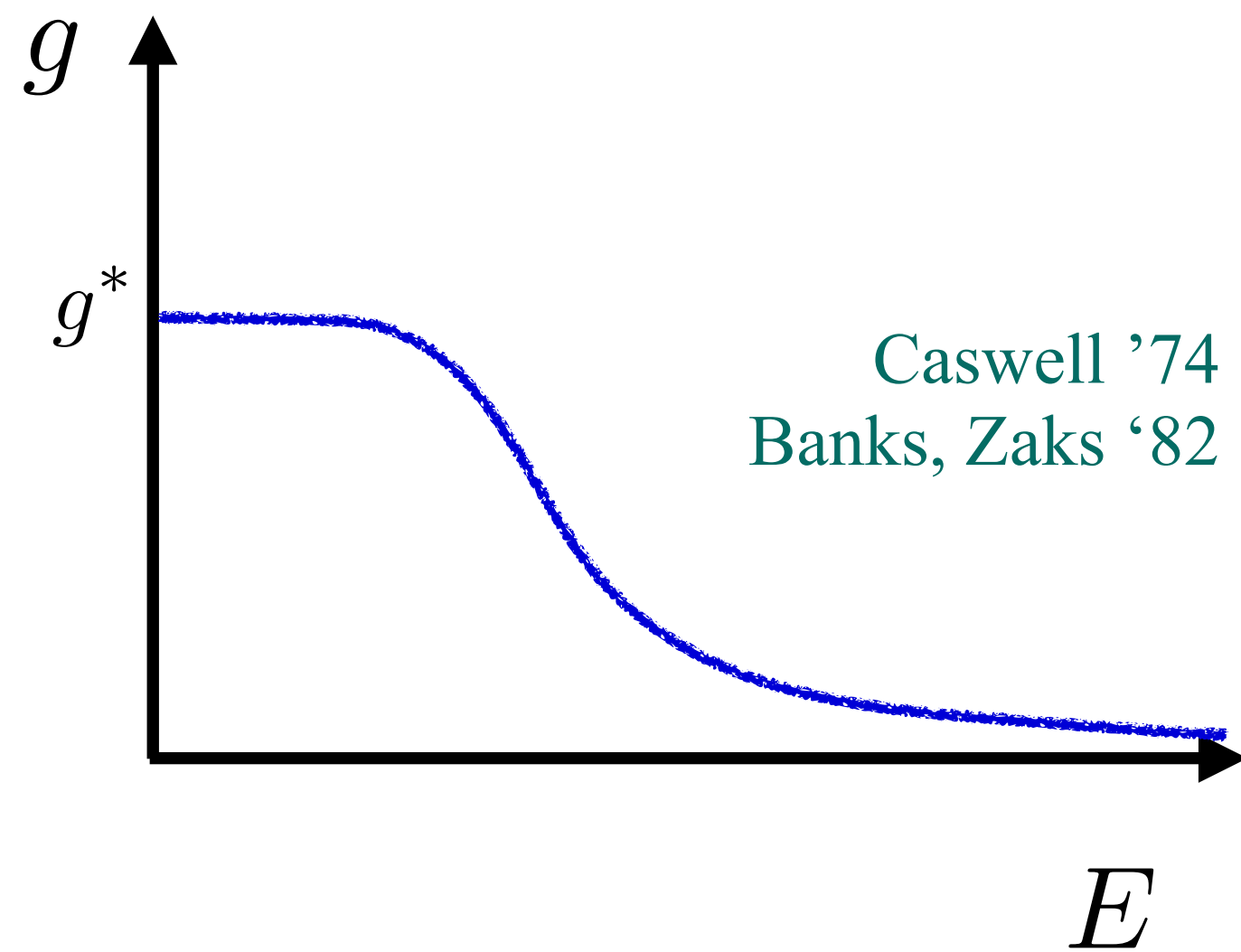
$$\begin{aligned} \rightarrow \beta_g &= -\frac{g^3}{(4\pi)^2} \left(\frac{11}{3} C_A - \frac{4}{3} T_F N_f \right) \\ &\quad - \frac{g^5}{(4\pi)^4} \left(\frac{34}{3} C_A - 4C_F T_F N_f - \frac{20}{3} C_A T_F N_f \right) + \dots \end{aligned}$$

Conformality and scale invariance

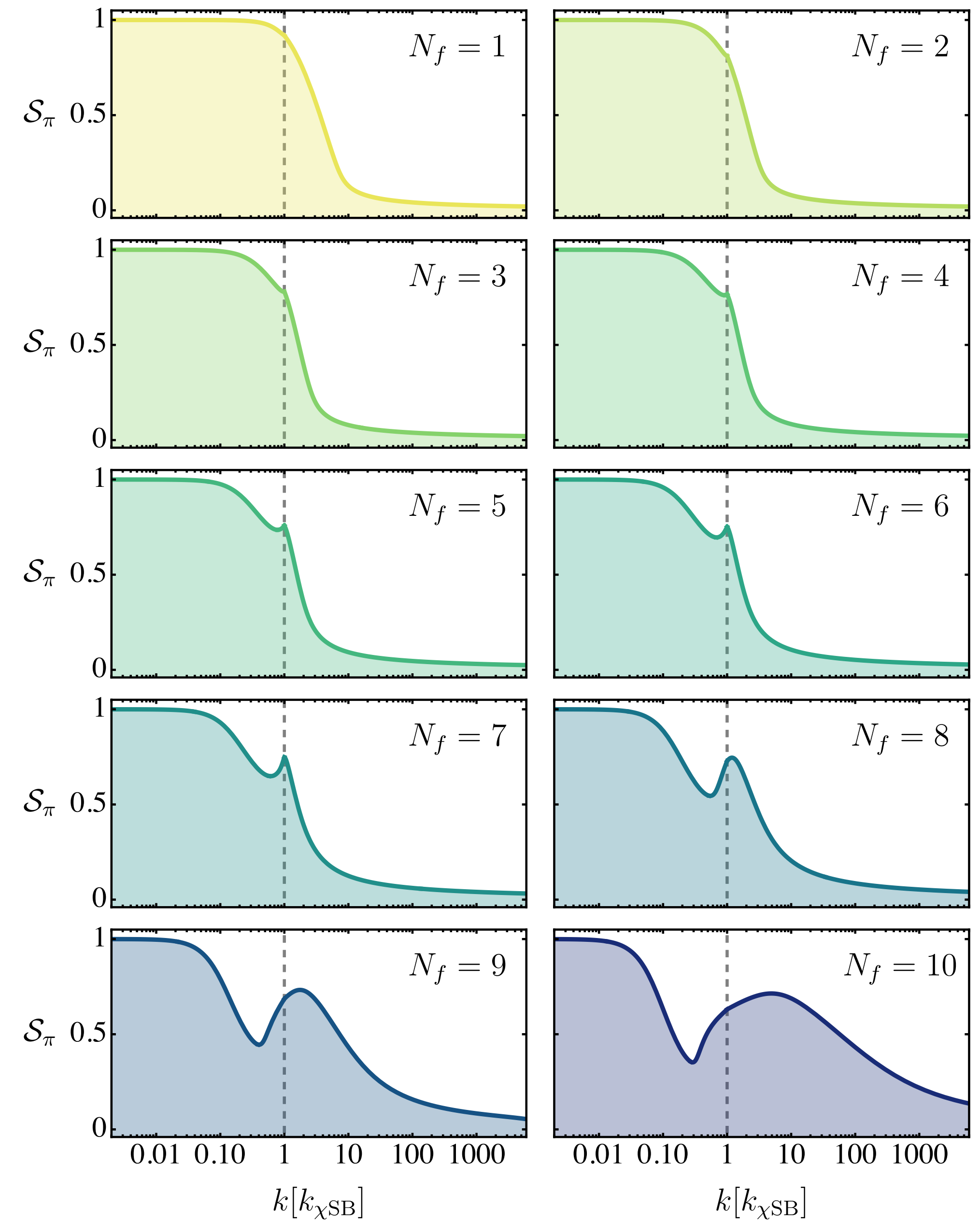
$$S = \int_x \left[\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} \right]_{\mathcal{G}_{\text{TC}}} + \bar{\psi} \not{D} \psi$$

$\text{SU}(N_f)_L \times \text{SU}(N_f)_R$

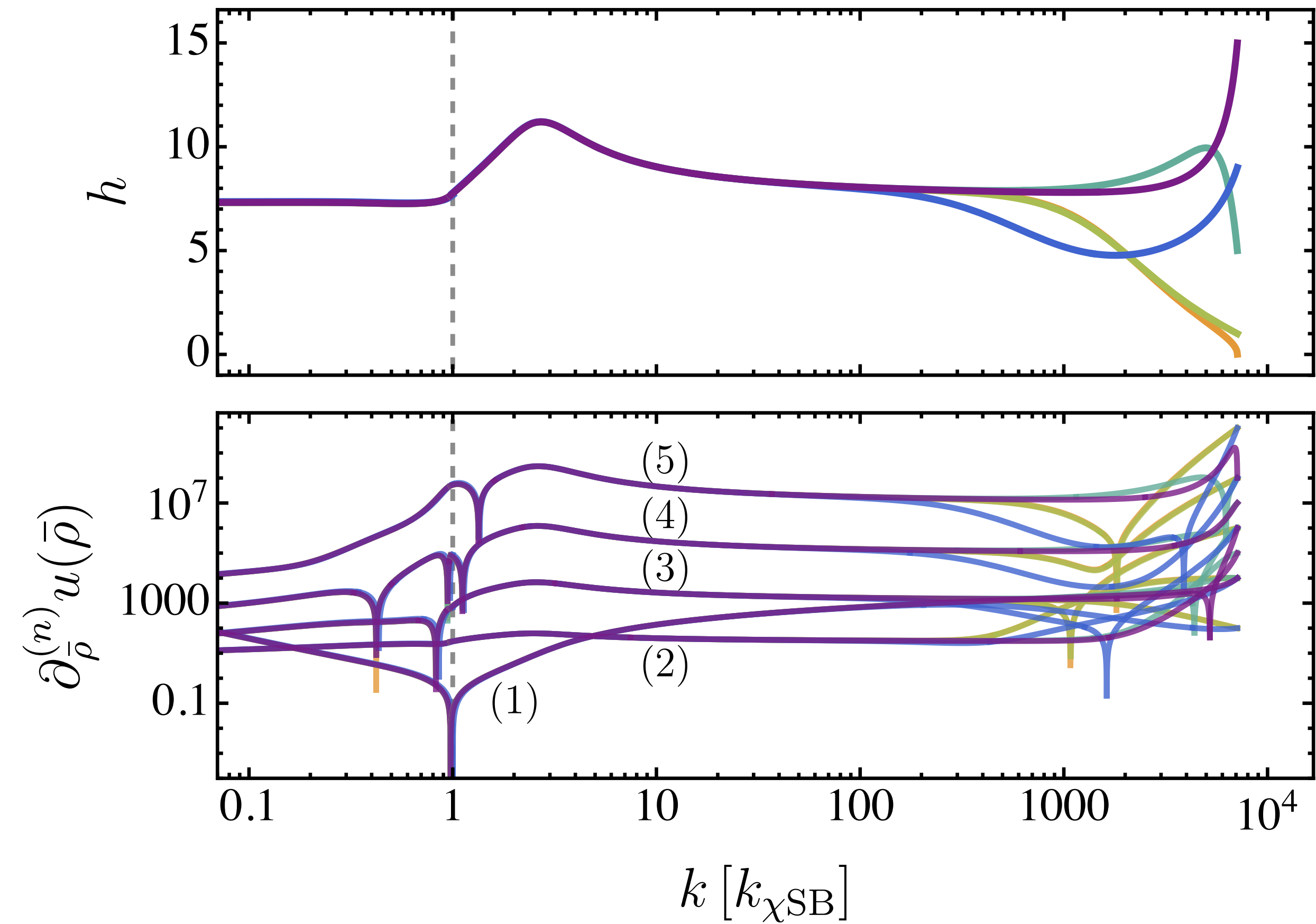
$$\rightarrow \beta_g = -\frac{g^3}{(4\pi)^2} \left(\frac{11}{3} C_A - \frac{4}{3} T_F N_f \right) - \frac{g^5}{(4\pi)^4} \left(\frac{34}{3} C_A - 4 C_F T_F N_f - \frac{20}{3} C_A T_F N_f \right) + \dots$$



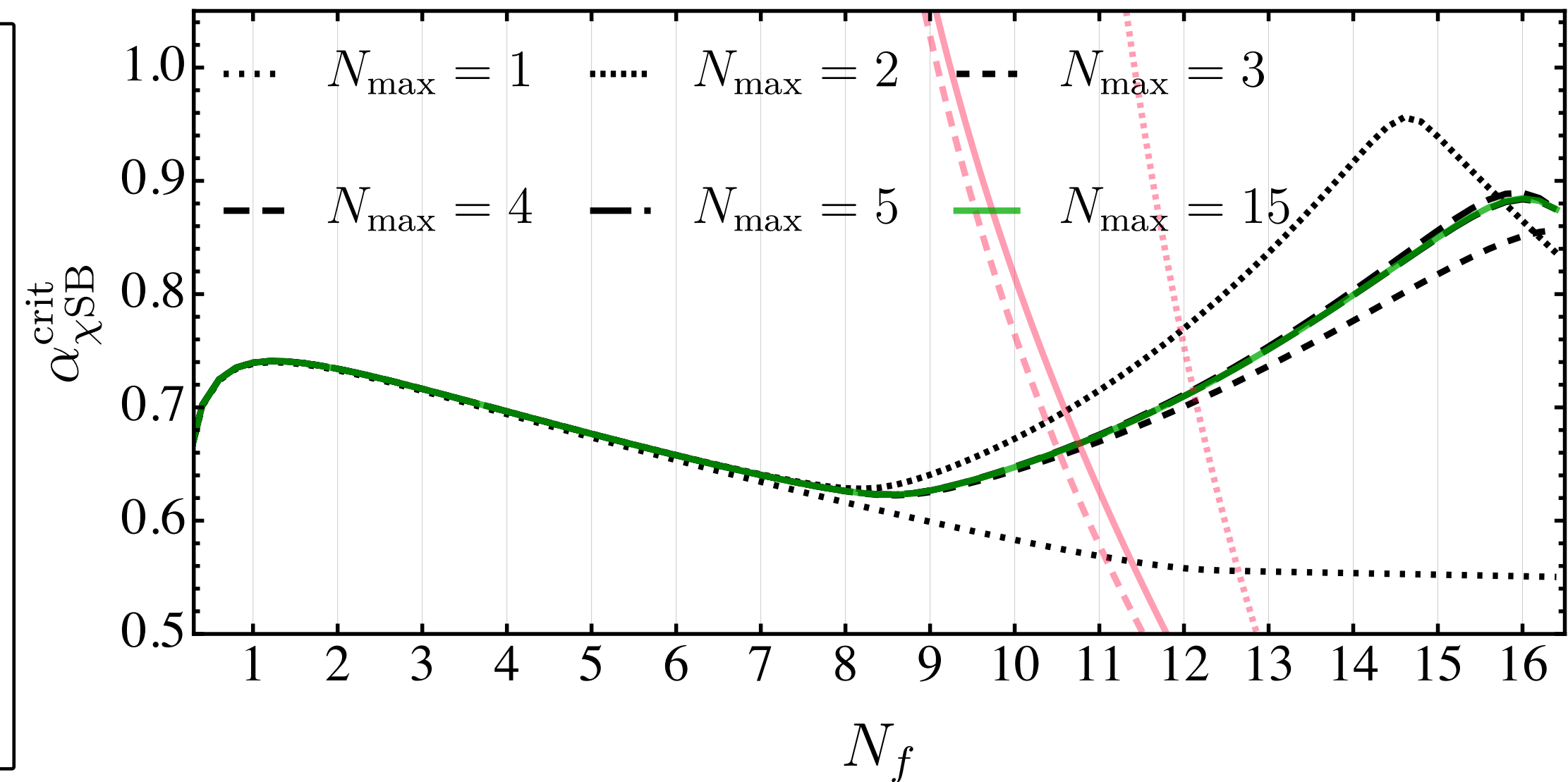
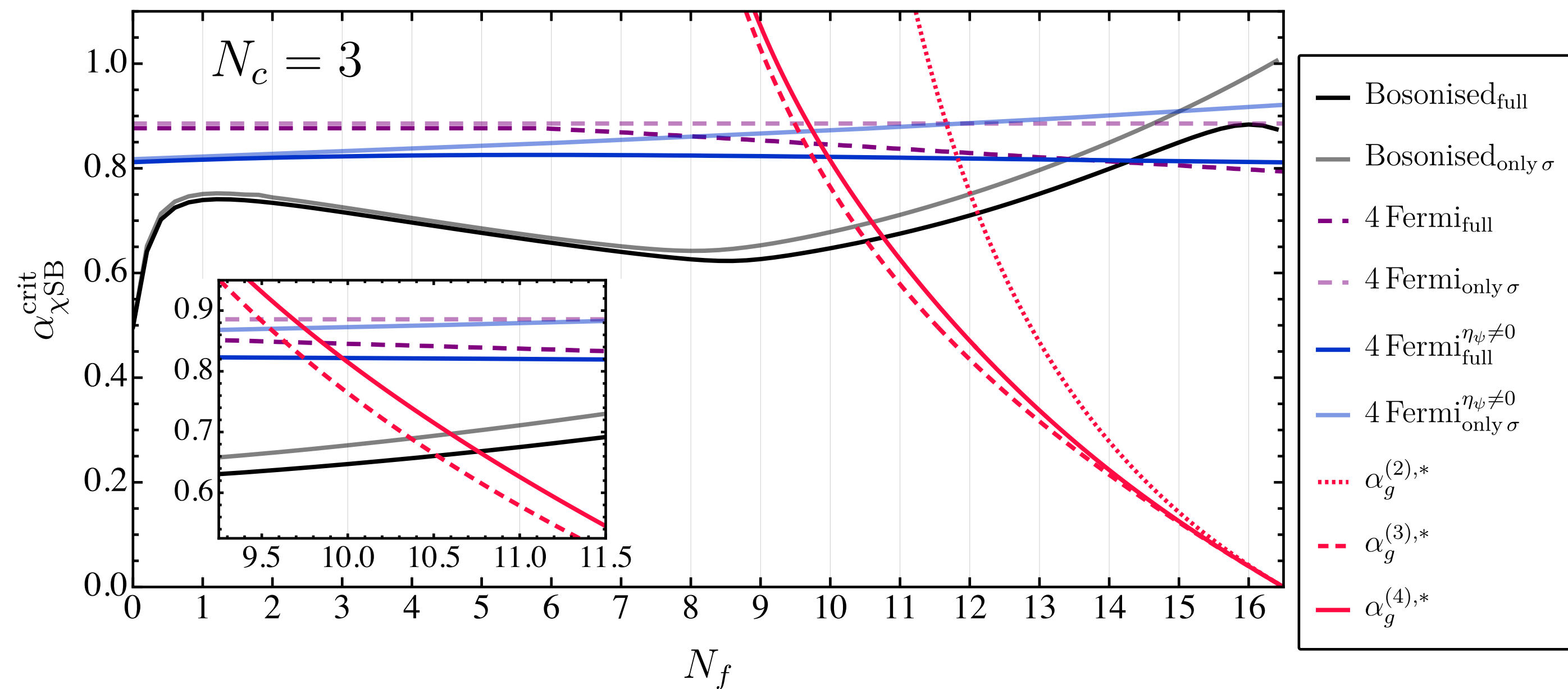
Validity of LEFTS



Dynamical bosonisation



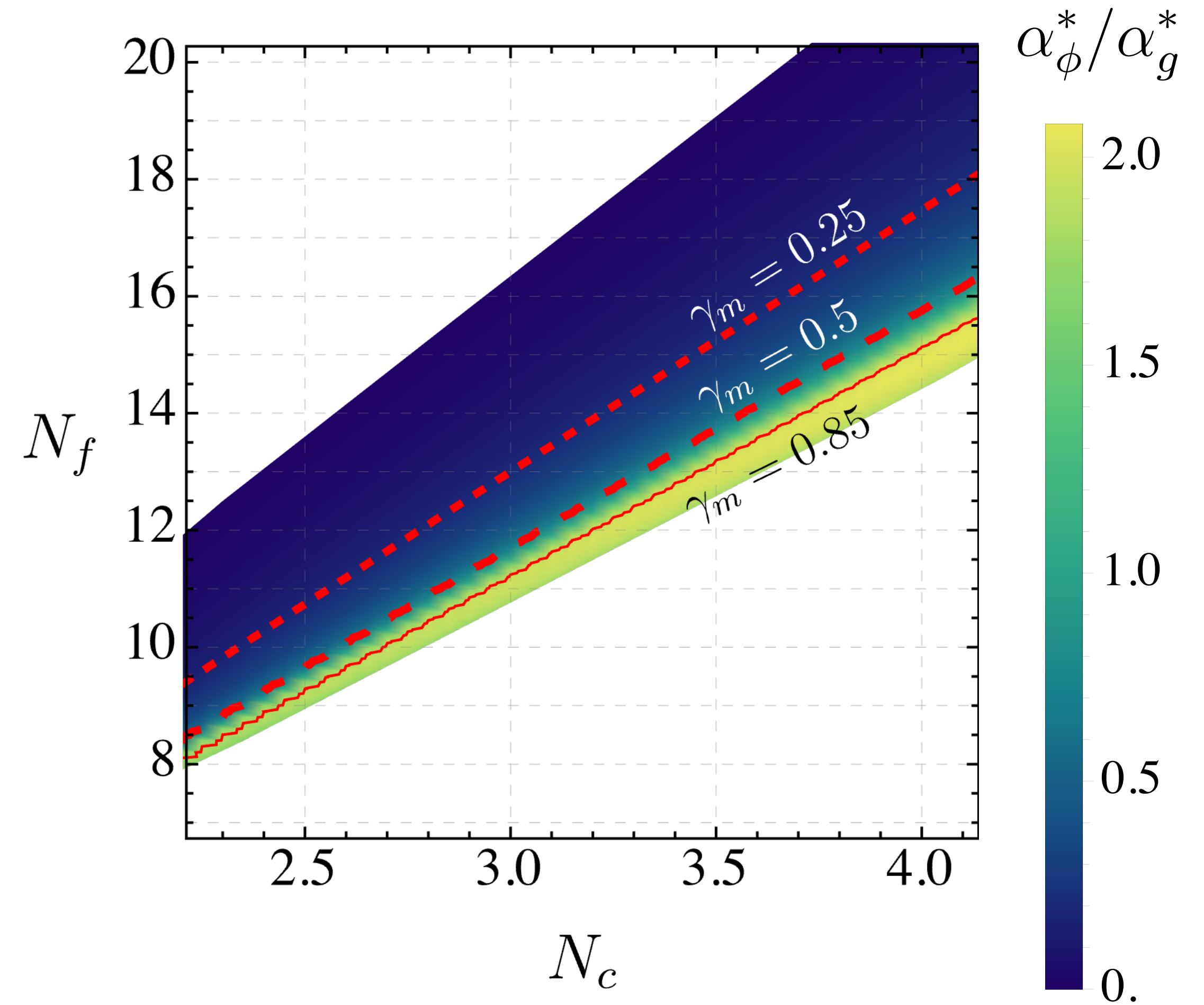
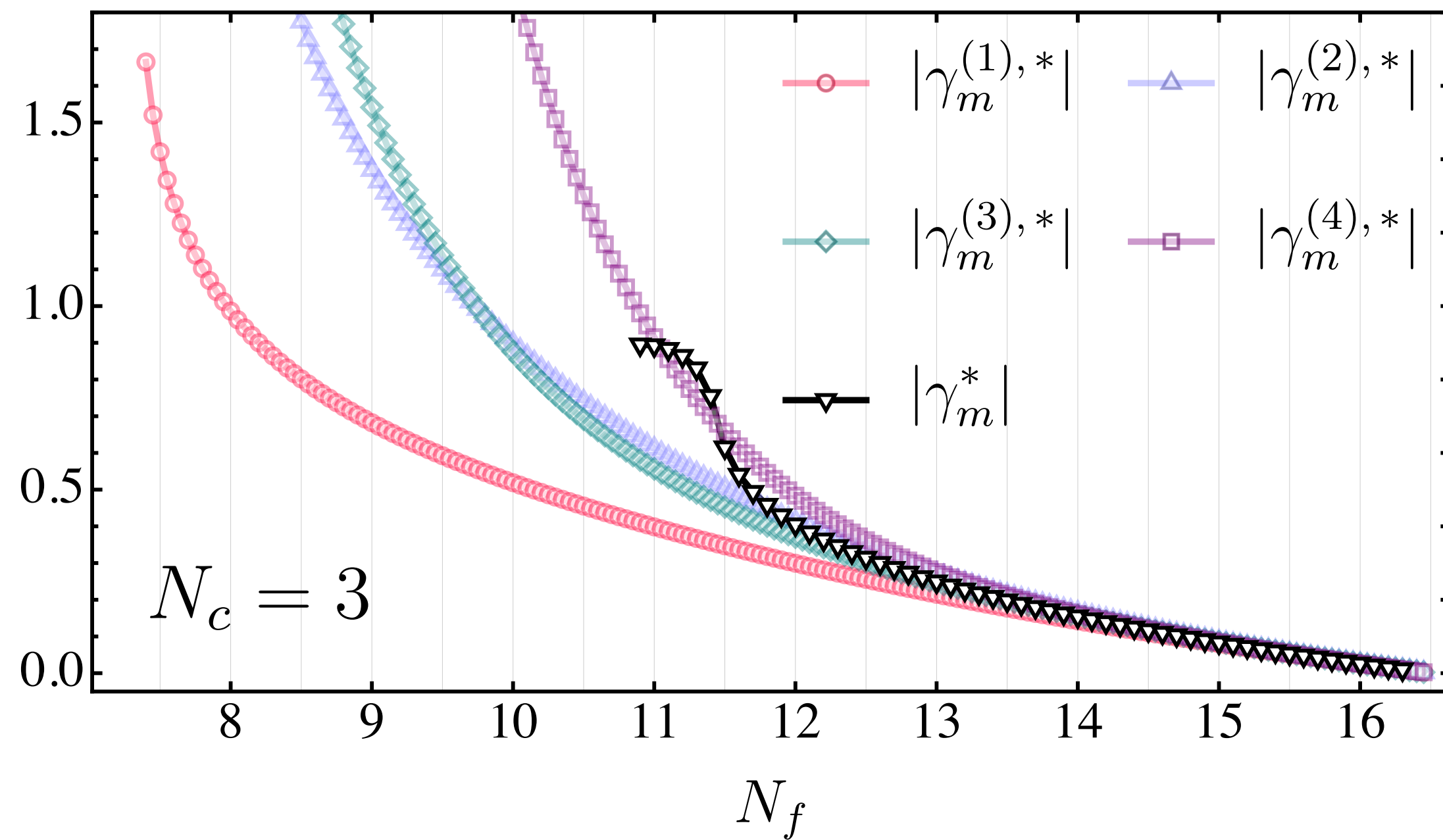
Boundary of the conformal window: systematics



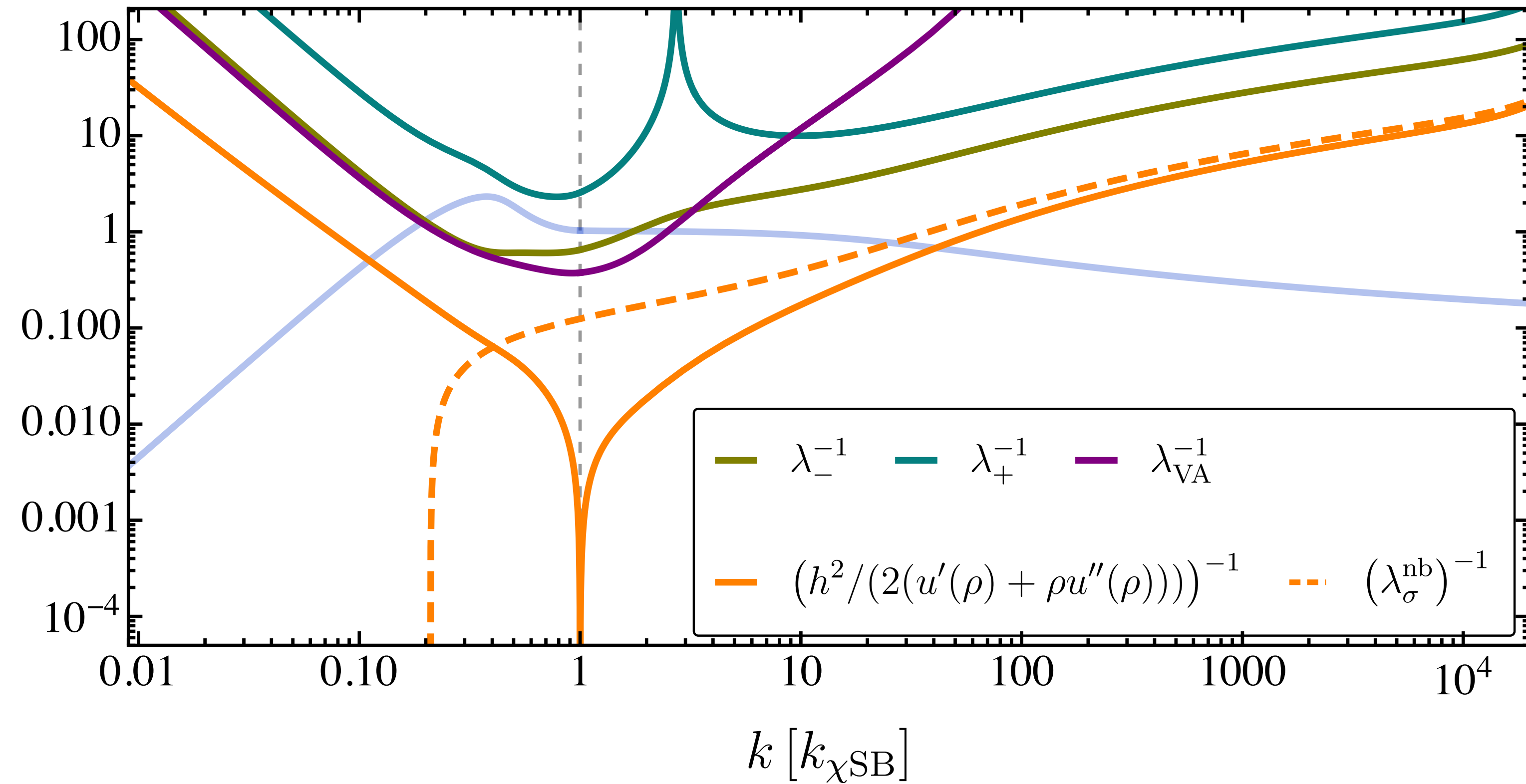
Boundary of the conformal window: γ_m

$$\gamma_m = \frac{\partial_t \bar{m}_\psi}{\bar{m}_\psi} = \left[\frac{\partial_t \left(T_f^0 \Gamma_k^{(\bar{\psi}\psi)}(p^2=0) \right)}{Z_\psi m_\psi} + \eta_\psi \right]$$

$$= \gamma_m^{(1)} + \gamma_m^{(>1)} = \frac{3 \alpha_g C_F}{2\pi} + \dots$$



Different tensor structures



Dynamical hadronisation and flow over EFTs

✓ Well defined implementation with the fRG

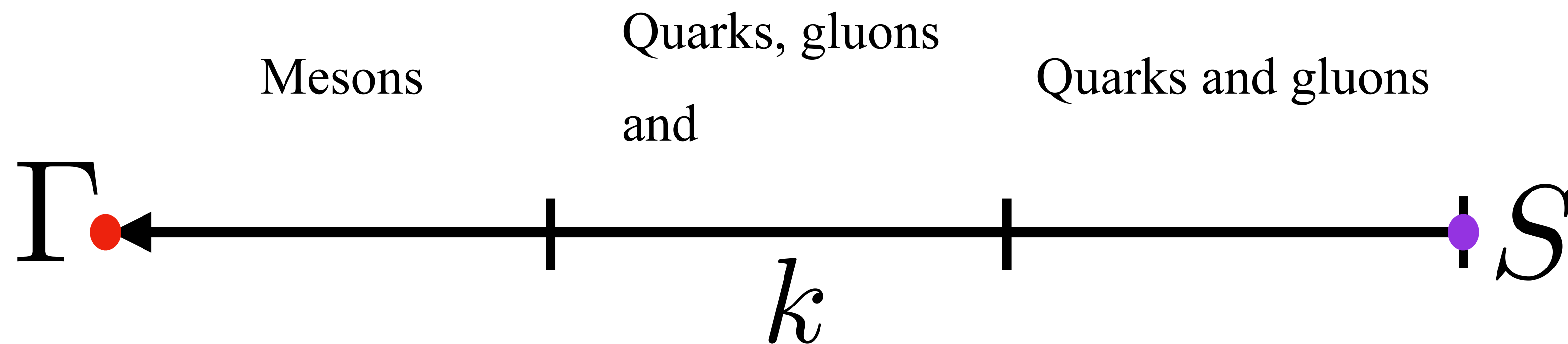
$$\begin{aligned} & \left(\partial_t + \int \dot{\phi}_c \frac{\delta}{\delta \phi_c} \right) \Gamma_k[\phi] \\ &= \frac{1}{2} \text{Tr} \left[\left(\frac{1}{\Gamma_k^{(2)} + R_k} \right)_{ij} \left(\partial_t \delta^{jn} + 2 \frac{\delta \dot{\phi}_n}{\delta \phi_j} \right) R_k^{mi} \right] \end{aligned}$$

Dynamical hadronisation and flow over EFTs

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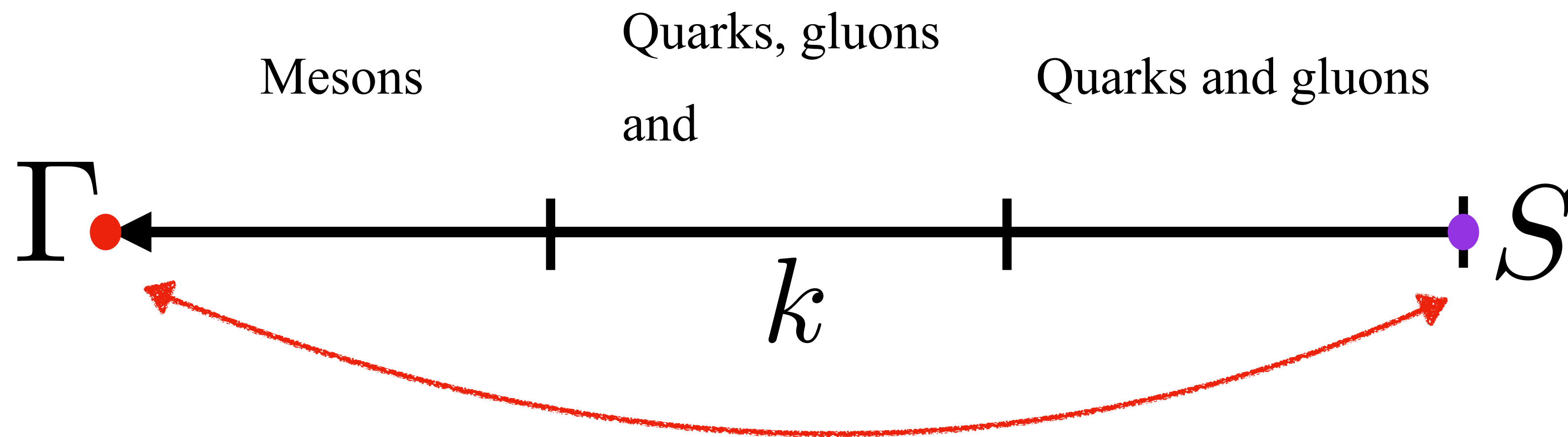


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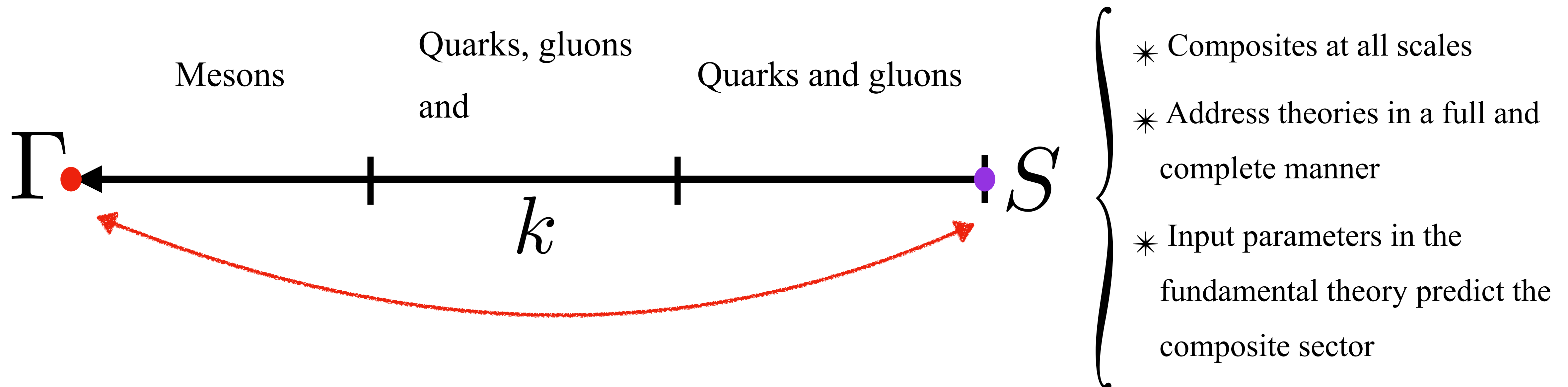


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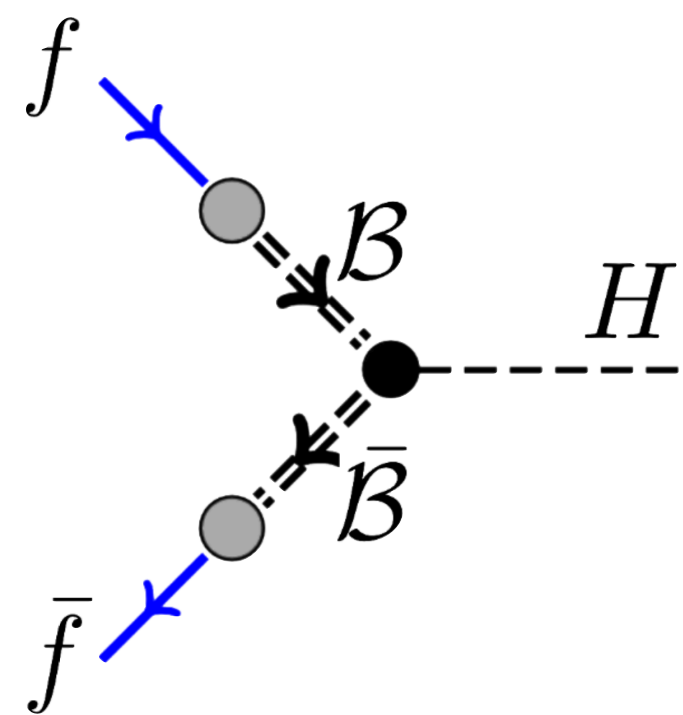


Flavour hierarchies from Fundamental Partial Compositeness

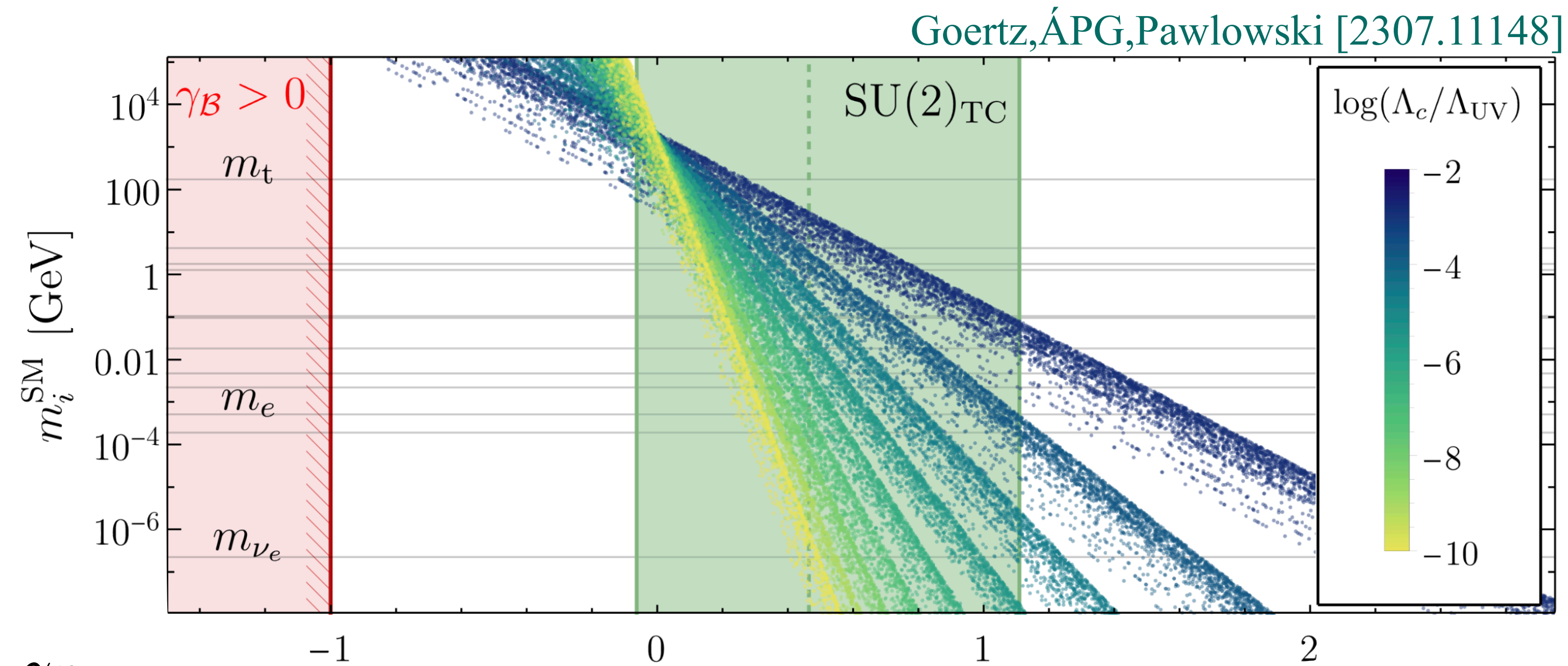
Sannino, Strumia, Tesi, Vigiani [1607.01659]
Cacciapaglia, Gertov, Sannino, Thomsen [1704.07845]

$$\bar{\Gamma}_{\text{mix}} = \int_x \left\{ \lambda_t^L \bar{q}_L \mathcal{B}_R^q + \lambda_t^R \bar{t}_R \mathcal{B}_L^t + \text{h.c.} \right\}$$

$$\mathcal{B} \sim \mathcal{F}\mathcal{S} \quad \underline{\text{MFPC}} \begin{cases} \text{Sp}(24)_{\mathcal{S}} \\ \text{SU}(4)_{\mathcal{F}} \end{cases}$$



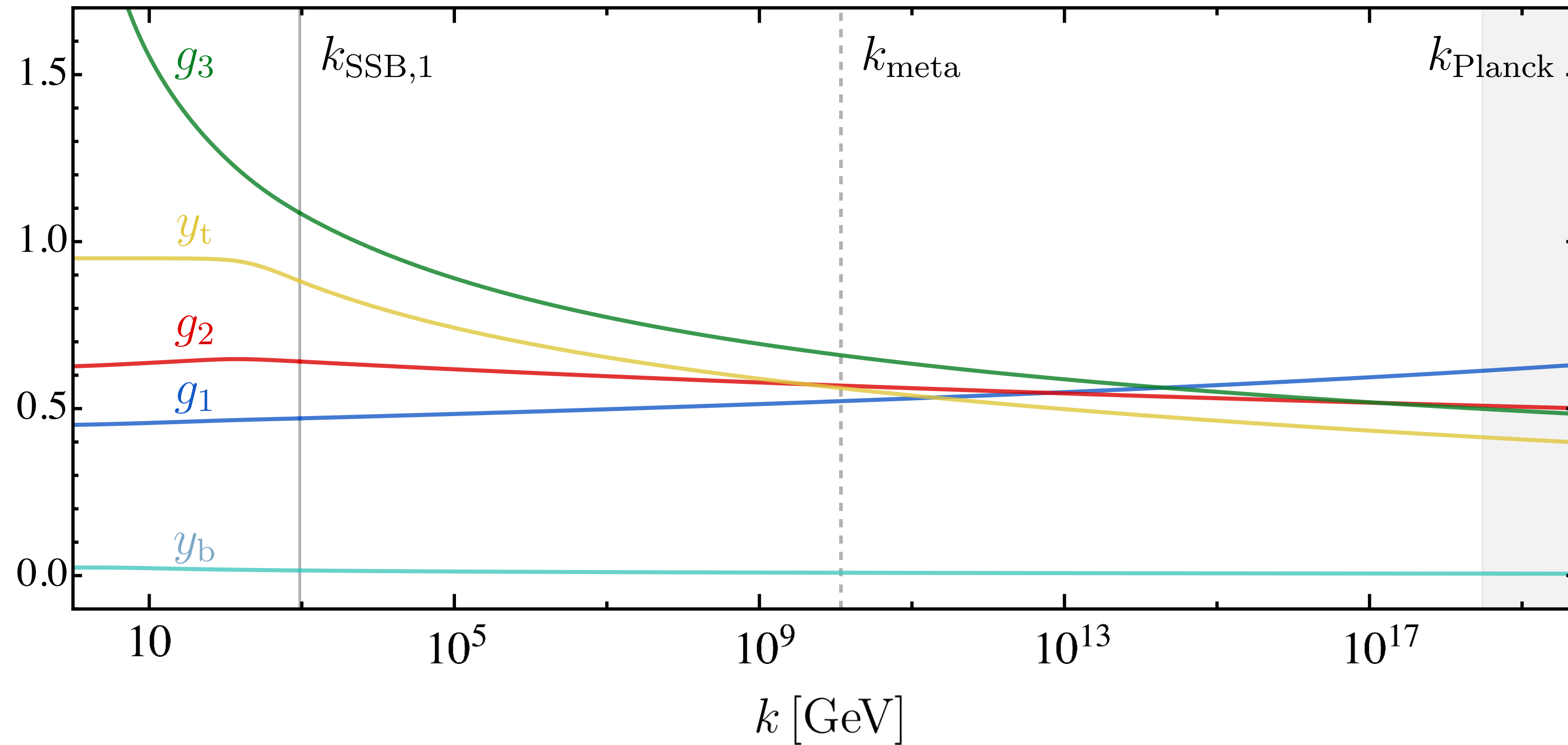
$$m_f \sim \frac{\mathcal{Y}_f v}{\sqrt{2}} \left(\frac{\Lambda_c}{\Lambda_{\text{UV}}} \right)^{-2-\gamma_{\mathcal{B}}}$$



The SM from the fRG

ÁPG,Pawlowski,Reichert [2207.09817]

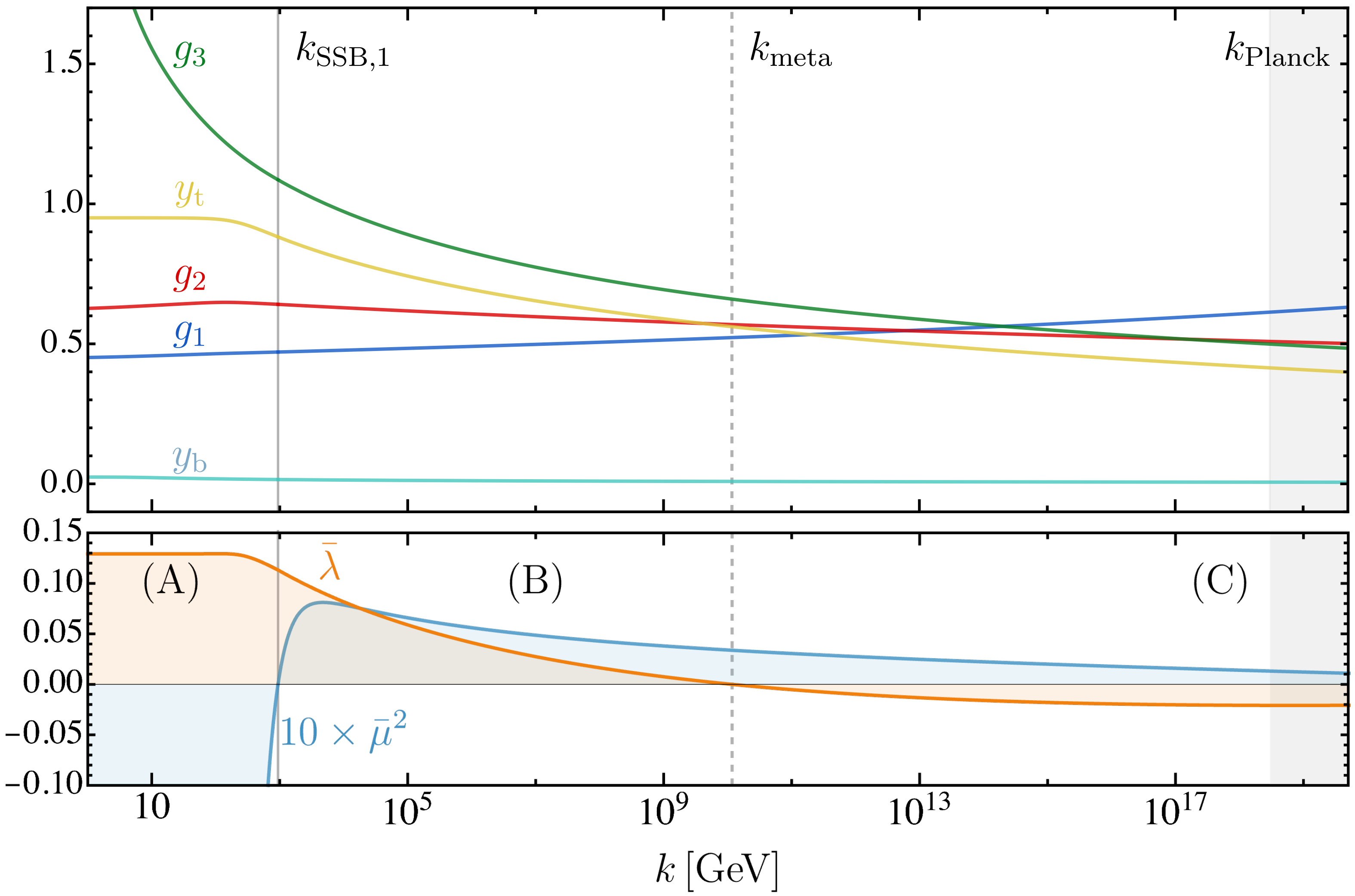
Goertz,ÁPG [2308.13594]



The SM from the fRG

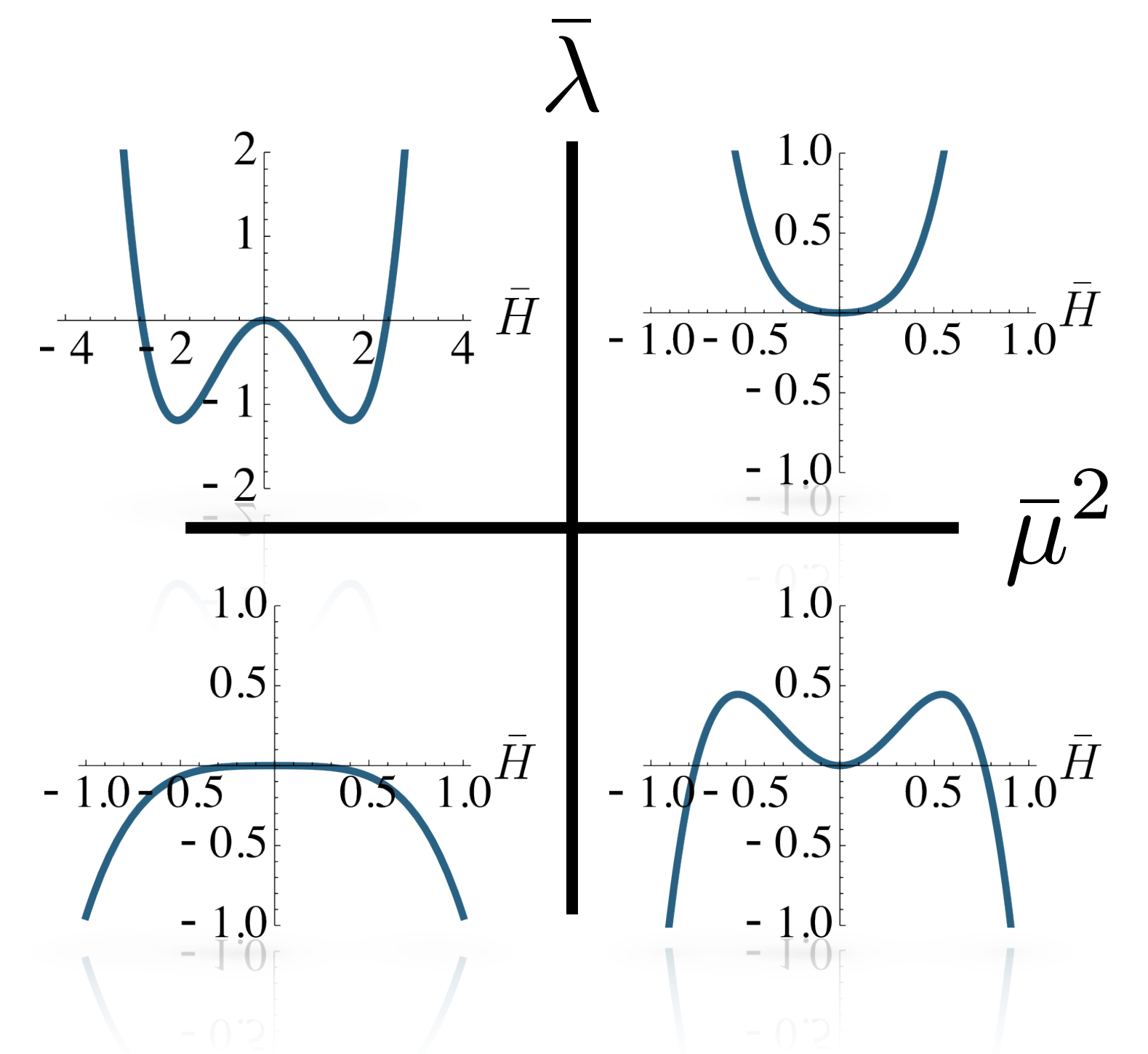
ÁPG, Pawłowski, Reichert [2207.09817]

Goertz, ÁPG [2308.13594]



$$u(\bar{\rho}) = V_{\text{eff}}(\rho) / k^4 = \bar{\mu}^2 \bar{\rho} + \bar{\lambda} \bar{\rho}^2$$

$$\bar{\rho} = Z_{\Phi} \frac{\text{tr} \Phi^{\dagger} \Phi}{k^2} \quad \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \mathcal{G}_1 + i\mathcal{G}_2 \\ H + i\mathcal{G}_3 \end{pmatrix}$$



The SM from the fRG

ÁPG,Pawlowski,Reichert [2207.09817]

Goertz,ÁPG [2308.13594]

ÁPG,Pawlowski,Reichert [2207.09817]

