



Instituto de  
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# Dynamical Cobordism without Supersymmetry

Based on work with R. Blumenhagen, N. Cribiori, C. Kneissl

Andriana Makridou

Strings Breaking SUSY @ MITP

November 22, 2023



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Based on work with R. Blumenhagen, N. Cribiori, C. Kneissl

See also talk by Roberta Angius!

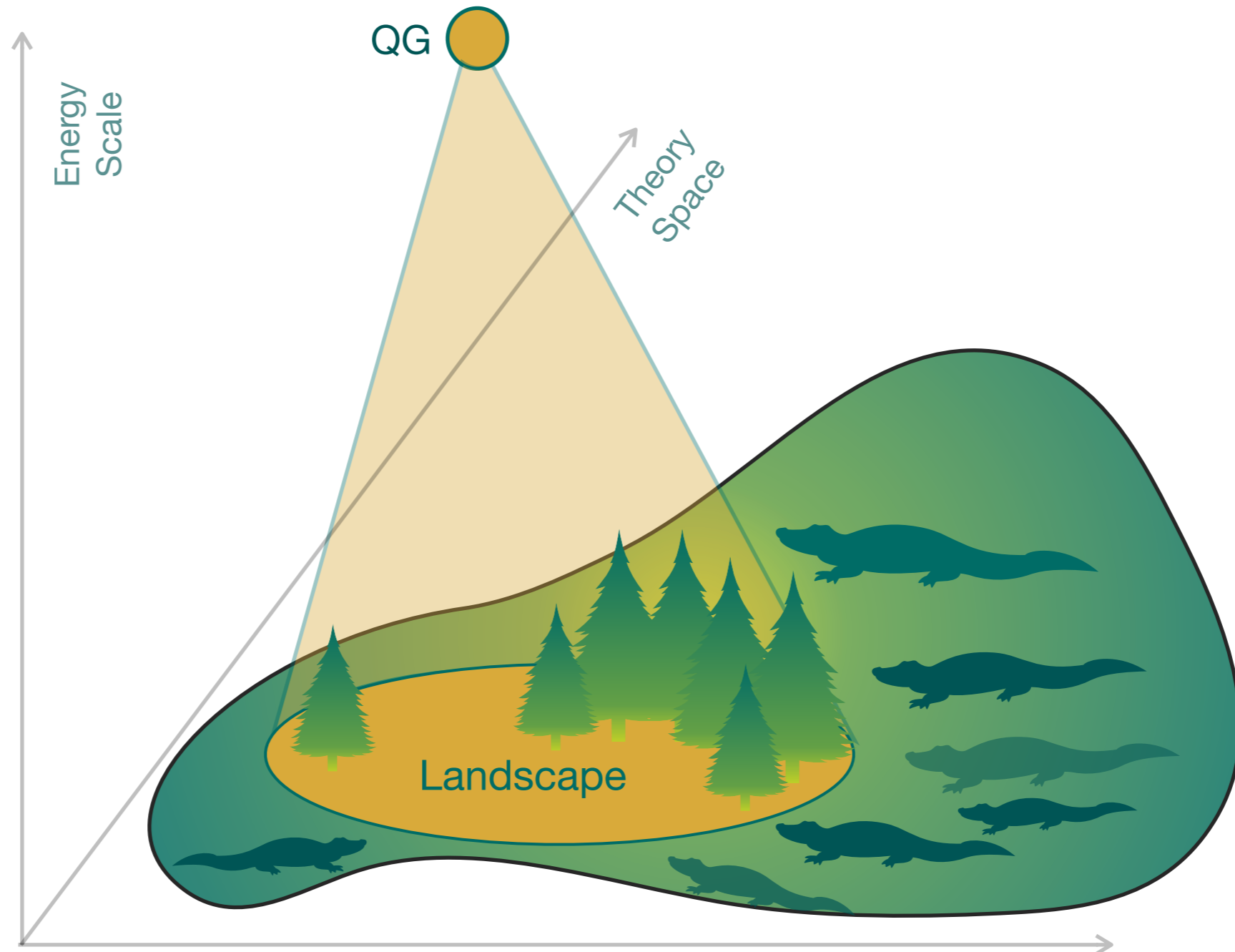
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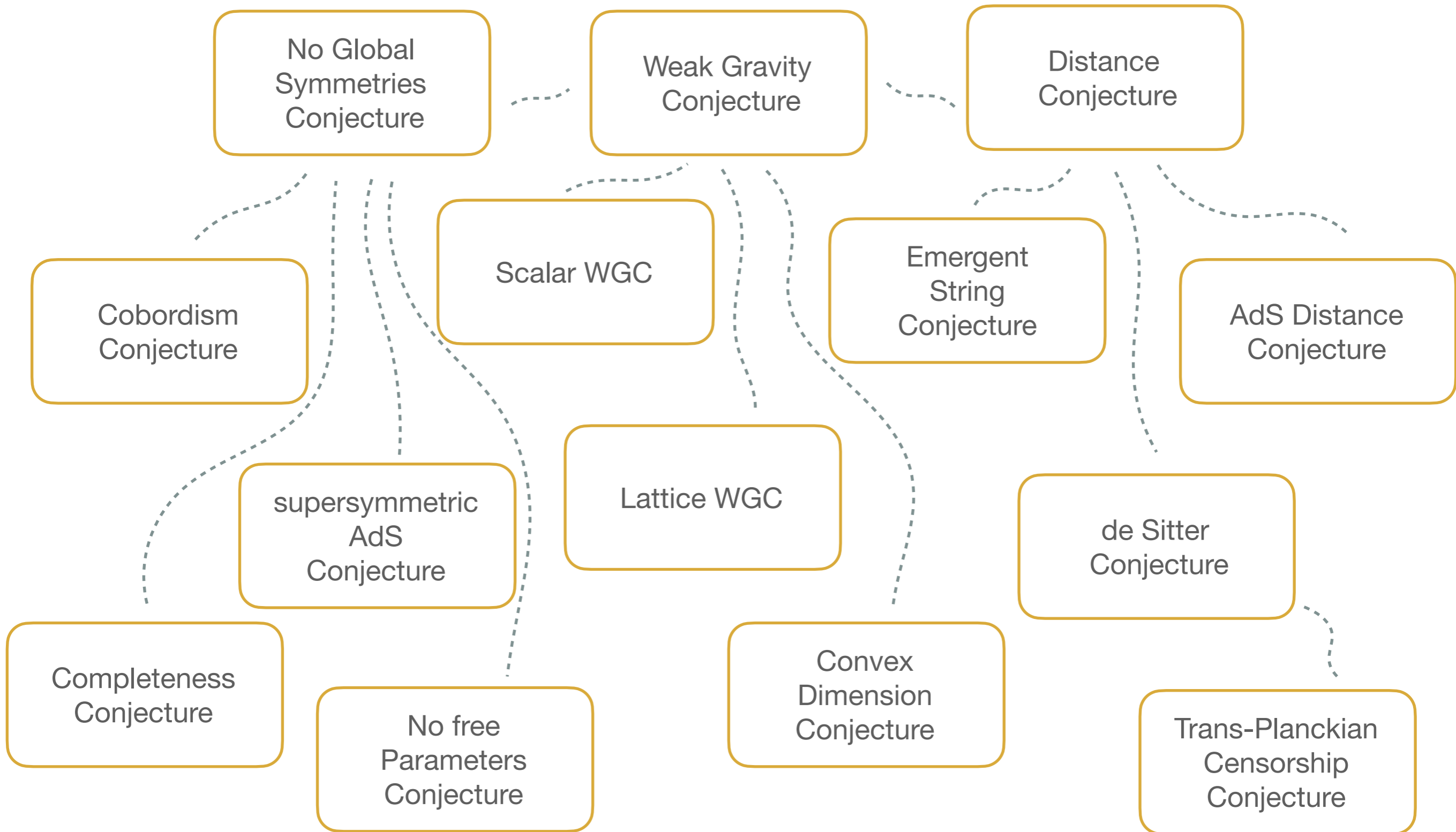
# The Swampland Program

Idea: Not all low-energy EFTs can be UV-completed to Quantum Gravity

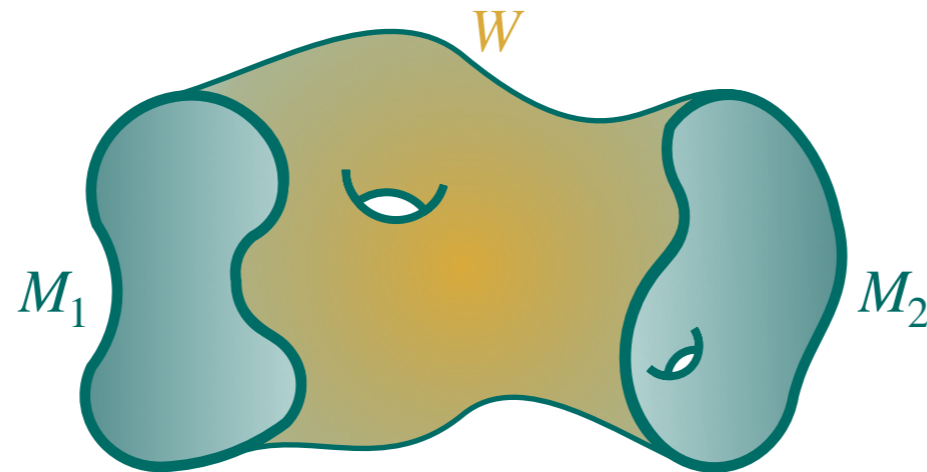


[Recent reviews: Palti '19, Valenzuela et al '21, Vafa et al '22]

# A web of swampland conjectures



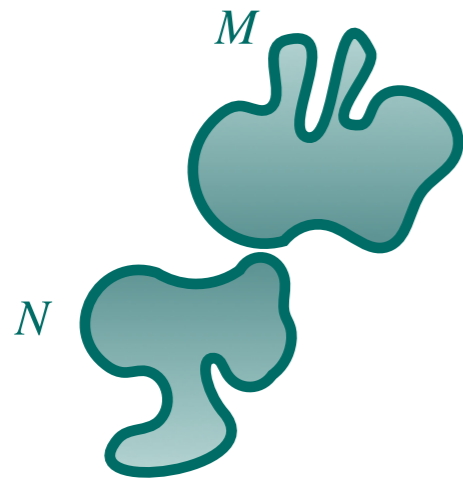
# Cobordism



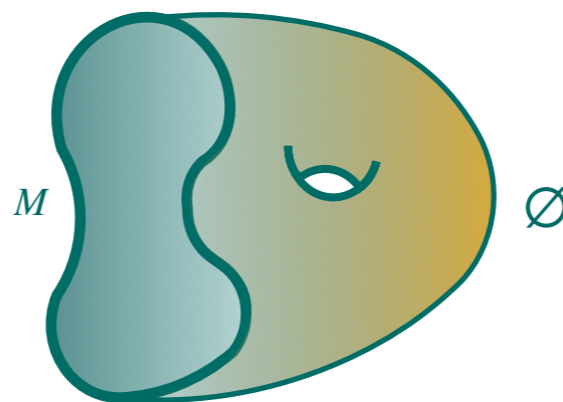
Allowed topology changes  
(Encoded in  $\xi$ )

$$M_1 \sim M_2 \Leftrightarrow \exists W \text{ s.t. } \partial W = M_1 \sqcup M_2$$

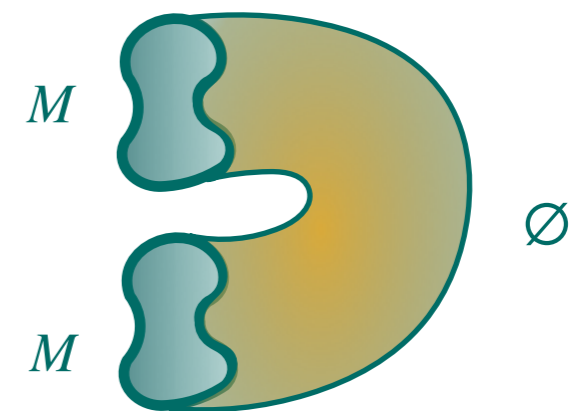
$$\Omega_k^\xi = \{\text{compact, closed, } k\text{-dimensional manifolds}\} / \sim$$



$$[M \sqcup N] = [M] + [N]$$

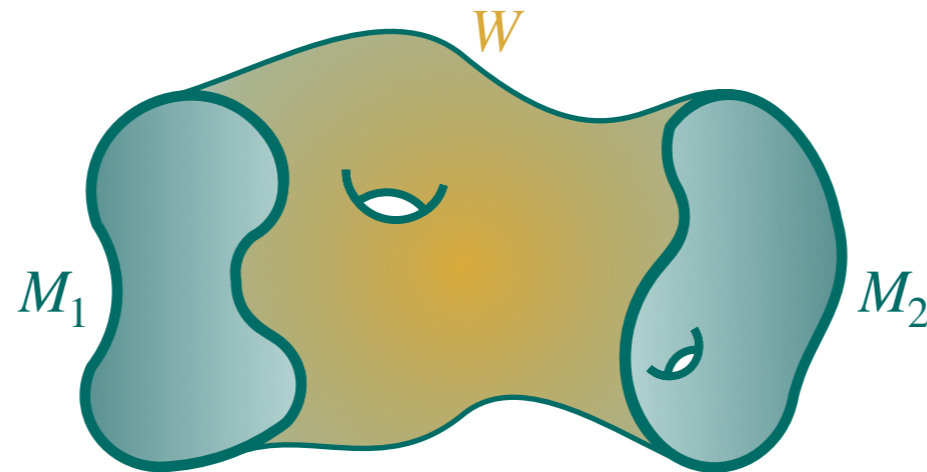


$$[\emptyset] = 0$$



$$[M] + [M] = [M \sqcup M] = 0$$

# Cobordism



Allowed topology changes  
(Encoded in  $\xi$ )

$$M_1 \sim M_2 \Leftrightarrow \exists W \text{ s.t. } \partial W = M_1 \sqcup M_2$$

$$\Omega_k^\xi = \{\text{compact, closed, } k\text{-dimensional manifolds}\} / \sim$$

Compactify  $d$ -dimensional theory on  $M^k$  down to  $D=d-k$  dimensions:



$EFT_1$   $\text{Domain Wall}$   $EFT_2$   
 $(d-k)$  – dimensional  $(d-k-1)$  – dimensional  $(d-k)$  – dimensional

→ Finite-energy transition between EFTs

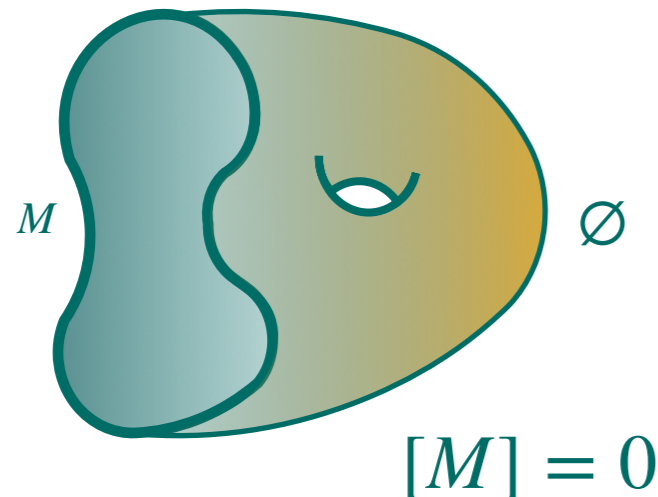
→ Domain walls detected by cobordism can be well-known objects, e.g. D-branes

# Cobordism Conjecture

No global symmetries in quantum gravity  $\rightarrow$  Cobordism Conjecture

[e.g., Banks, Seiberg '11]

[McNamara, Vafa '19]



Cobordism Conjecture:

All Cobordism Classes should be trivial

$$\Omega_k^{QG} = 0$$

Cobordism Group  $\Omega_k^\xi \leftrightarrow$  Cobordism Invariant  $\mu_k$

For trivial cobordism class:  $\mu_k[\emptyset] = 0$

If cobordism class  $[M] \neq 0 \leftrightarrow$  obstruction to decay into “nothing”

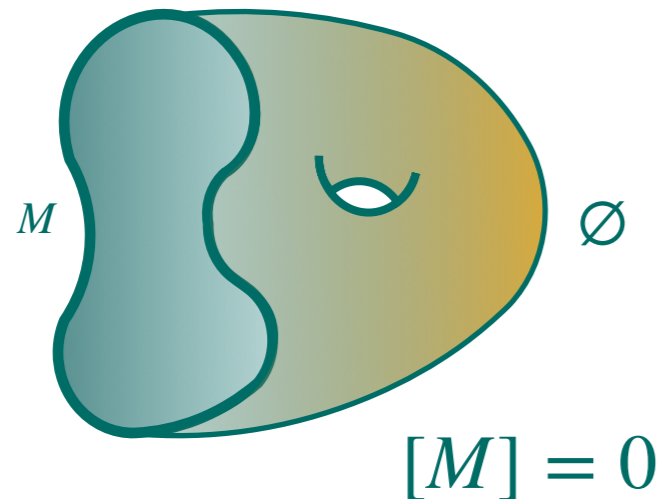
$\Omega_k^\xi \neq 0 \Leftrightarrow (d - k - 1)$ -dim. global symmetry  
with charges labelled by classes  $[M] \in \Omega_k^\xi$

# Cobordism Conjecture

No global symmetries in quantum gravity  $\rightarrow$  Cobordism Conjecture

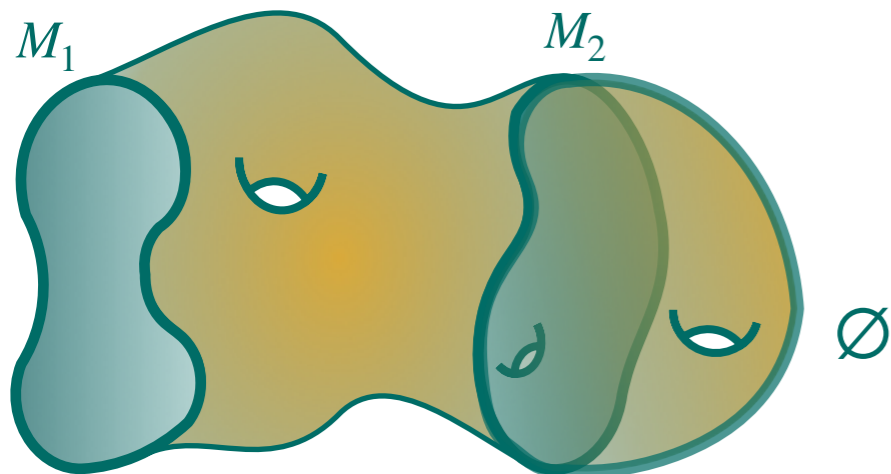
[e.g., Banks, Seiberg '11]

[McNamara, Vafa '19]



Cobordism Conjecture:  
All Cobordism Classes should be trivial

$$\Omega_k^{QG} = 0$$



In  $D=d-k$  dimensions:



*Major Implication: there is a "single" theory of quantum gravity*

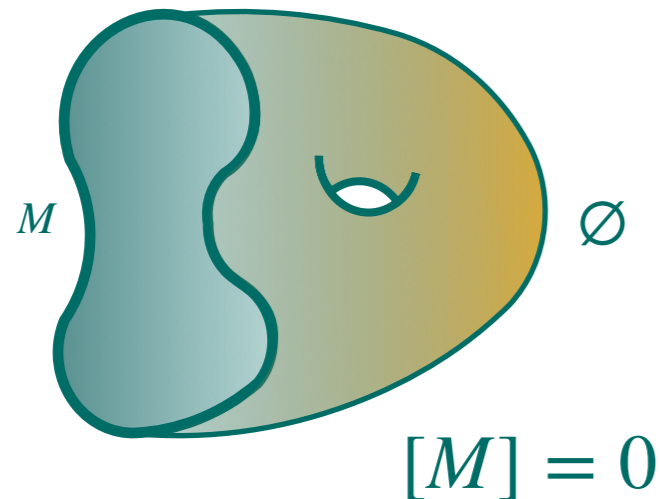


# Cobordism Conjecture

No global symmetries in quantum gravity  $\rightarrow$  Cobordism Conjecture

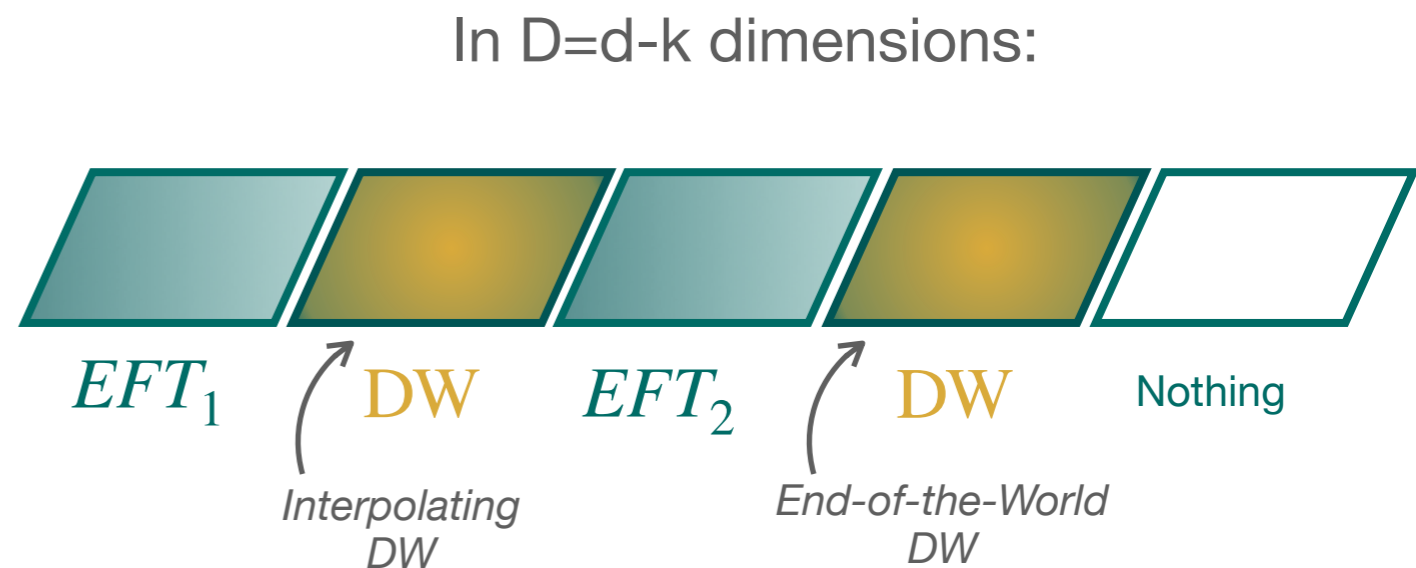
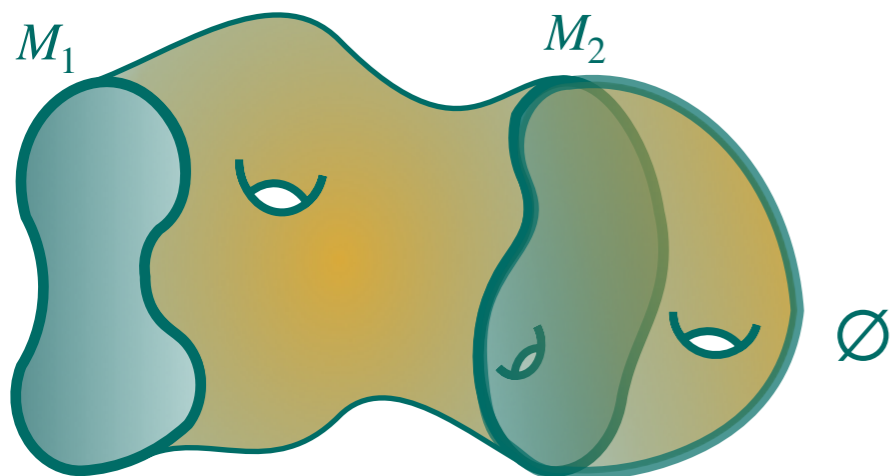
[e.g., Banks, Seiberg '11]

[McNamara, Vafa '19]



Cobordism Conjecture:  
All Cobordism Classes should be trivial

$$\Omega_k^{QG} = 0$$



*Major Implication: there is a "single" theory of quantum gravity*

# Cobordism Conjecture in practice

$$\Omega_k^{\widetilde{QG}} \neq 0$$

Breaking

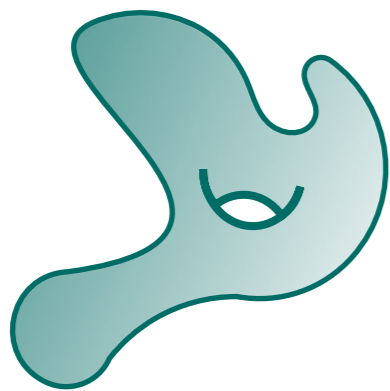
Gauging

$$\Omega_k^{\widetilde{QG}} \rightarrow \Omega_k^{\widetilde{QG}+\text{defects}} = 0$$

$$0 = [M] \in \Omega_k^{\widetilde{QG}} \neq 0$$

$$\Omega_k^{\widetilde{QG}+\text{g.fields}} \rightarrow \Omega_k^{\widetilde{QG}}$$

[Blumenhagen, Cribiori '21],  
[Blumenhagen, Cribiori, Kneissl, AM '22/2]



$$0 \neq [M] \in \Omega_k^{\widetilde{QG}}$$



$$[M] = 0 = \Omega_k^{\widetilde{QG}+\text{defects}}$$



New defects possibly detected!

[McNamara, Vafa '19], [Montero, Vafa '20], [Dierigl, Heckman '21]

[Debray, Dierigl, Heckman, Montero '21]

[Dierigl, Heckman, Montero, Torres '22/23]

[Blumenhagen, Cribiori, Kneissl, AM '22]

[Blumenhagen, Kneissl, Wang '23]

[Kaidi, Ohmori, Tachikawa, Yonekura '23]

[Friedrich, Hebecker, Walcher '23]

# Cobordism and string theory

What are the relevant structures for String Theory?

e.g. Type I  $\leftrightarrow \xi = \text{Spin}$ , Type II  $\leftrightarrow \xi = \text{Spin}^c$

[see also e.g. Andriot, Carqueville, Cribiori '22, Debray '23]

$n$	0	1	2	3	4	5	6	7	8	9	10
$\Omega_n^{\text{Spin}}$	$\mathbb{Z}$	$\mathbb{Z}_2$	$\mathbb{Z}_2$	0	$\mathbb{Z}$	0	0	0	$2\mathbb{Z}$	$2\mathbb{Z}_2$	$3\mathbb{Z}_2$
$\Omega_n^{\text{Spin}^c}$	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$2\mathbb{Z}$	0	$2\mathbb{Z}$	0	$4\mathbb{Z}$	0	$4\mathbb{Z} \oplus \mathbb{Z}_2$

Example:  $\Omega_{p+1}^{\text{Spin}, \text{U}(1)_p} \supset \mathbb{Z}$

$$dF_{p+1} = 0 \Leftrightarrow \int_M F_{p+1} \in \mathbb{Z}$$

[McNamara, Vafa '19, see also Dierigl @ASC '23]

Can kill cobordism charges by adding source in Bianchi identity:  $dF_{p+1} \neq 0$

Physical object: magnetically charged (p-3)-manifold  $\rightarrow$  Dp-branes

*Predicted objects can be totally new! e.g. R7-branes, heterotic branes*

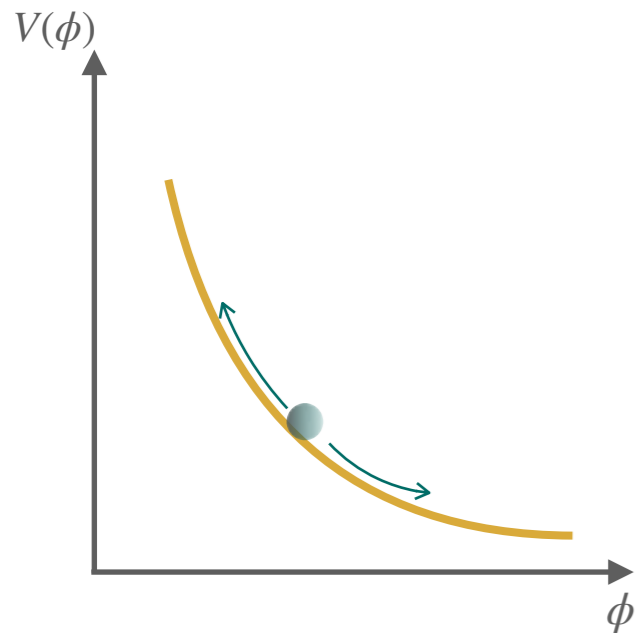
[Dierigl, Heckman, Montero, Torres '22/23] [Kaidi, Ohmori, Tachikawa, Yonekura '23]

**See talk by Justin Kaidi this afternoon!**

*Side remark: Cobordism also relevant in study of anomalies, with recent application to non-SUSY strings!*

[Basile, Debray, Delgado, Montero '23]

# Dynamical Cobordism : Motivation



Dynamical tadpoles (vs RR tadpoles)

Naturally occurring in supersymmetry-breaking potentials

Indicate lack of maximally-symmetric vacuum

[Sugimoto '99]  
 [Antoniadis, Dudas, Sagnotti '99]  
 [Angelantonj '99]  
 ...  
 [Raucci '22]  
 [Basile, Raucci, Thomée '22]  
 ...  
 [Mourad, Sagnotti '23]

**Example:** Sugimoto Model (  $USp(N)$  Type I with  $N \bar{D}8$  and  $N O9^+$  )

[Sugimoto '99]

Action:

$$S = \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-G} \left( \mathcal{R} - \frac{1}{2} (\partial\Phi)^2 \right) - T_9 \int d^{10}x \left( (N + 32) \sqrt{-G} e^{\frac{3}{2}\Phi} - (N - 32) A_{10} \right) + \dots$$

Dudas-Mourad solution - preserving 9d Poincaré invariance:

*See talk by Salvatore Raucci tomorrow!*

[Dudas, Mourad '00]

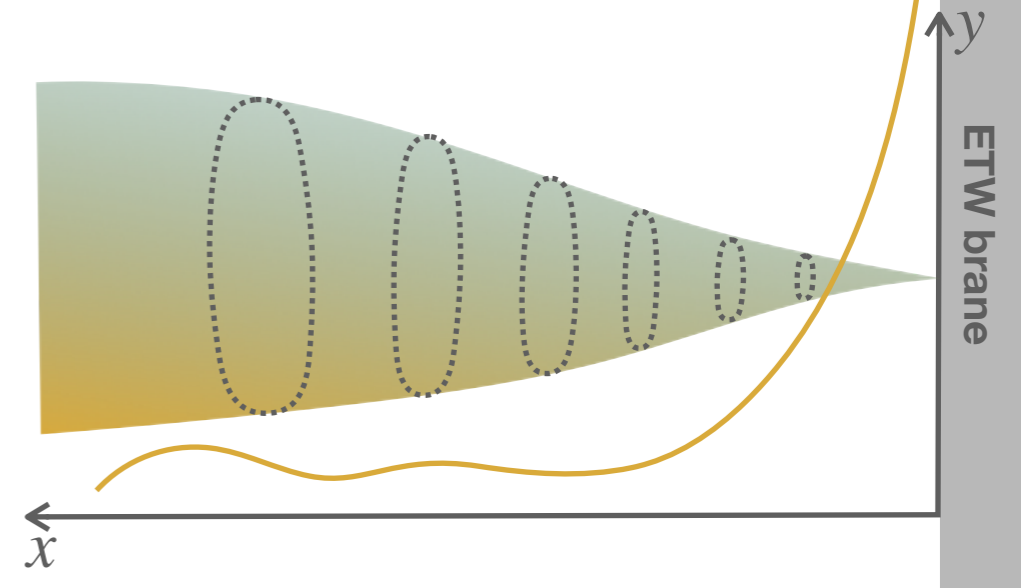
$$ds_E^2 = \sqrt{\alpha_E}^{1/9} e^{-\frac{\alpha_E}{8} y^2} \eta_{\mu\nu} dx^\mu dx^\nu + \sqrt{\alpha_E} y^{-1} e^{-\frac{3\phi_0}{2}} e^{-\frac{9\alpha_E}{8} y^2} dy^2, \quad \alpha_E = 64k^2 T_9$$

→ singularities at finite spacetime distance, spontaneous compactification to 9d

# Dynamical Cobordism : Scaling Relations

## Common features:

- Solution extends over finite spacetime distance  $\Delta$
- Ricci curvature singularity
- Scalar diverges close to the singularity  
Field distance  $D \rightarrow \infty$



**Interpretation:** [Buratti, Delgado, Uranga '21]

*Physical mechanism cutting off spacetime = cobordism defect of the initial theory*  
 $\rightarrow$  *End-of-the-World (ETW) brane*

[Buratti, Calderon-Infante, Delgado, Uranga '21]

## Cobordism Distance Conjecture:

- An infinite field distance limit is realized as running into a cobordism wall of nothing.
- In this limit one has the scaling relations

$$\Delta \sim e^{-\frac{1}{2}\delta D}, \quad \mathcal{R} \sim e^{\delta D}, \quad \delta > 0.$$

*Universal local description possible in terms of critical exponent  $\delta$*

[Angius, Calderon-Infante, Delgado, Huertas, Uranga '21]

*See talk by Roberta Angius later!*

# Application to a non-supersymmetric setup

[Blumenhagen, Cribiori, Kneissl, AM'22]

**Setup:** Gauge neutral, non-supersymmetric 9d object w/ brane-like dilaton coupling

**Physical realisation:** non-BPS  $\widehat{D8}$ -brane, non-SUSY stack of  $16 \times \bar{D}8 + O8^{++}$

[Blumenhagen, Font '00]

**Action:** 
$$S = \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-G} \left( \mathcal{R} - \frac{1}{2} (\partial\Phi)^2 \right) - T \int d^{10} \sqrt{-g} e^{\frac{5}{4}\Phi} \delta(r).$$

↑  
transverse  
direction

**Solution Ansatz:** 
$$ds^2 = e^{2\mathcal{A}(r,y)} ds_8^2 + e^{2\mathcal{B}(r,y)} (dr^2 + dy^2).$$

$$\mathcal{A} = A(r) + U(y)$$

$$\mathcal{B} = B(r) + V(y)$$

$$\Phi = \chi(r) + \psi(y)$$

## Solution I

[Blumenhagen, Font '00]

$$A(r) = \frac{1}{8} \log \left| \sin \left( 8K \left( r - \frac{R}{2} \right) \right) \right|$$

$$U(y) = -Ky$$

$$B(r) = \frac{1}{8} \log \left| \sin \left( 8K \left( r - \frac{R}{2} \right) \right) \right|$$

$$V(y) = 0$$

$$\chi(r) = -\frac{3}{2} \log \left| \tan \left( 4K \left( r - \frac{R}{2} \right) \right) \right| + \phi_0$$

$$\psi(y) = 0$$

# Application to a non-supersymmetric setup

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## Solutions $II^\pm$

[Blumenhagen, Font '00]

$$A(r) = \frac{1}{8} \log \left| \sin \left( 8K \left( r - \frac{R}{2} \right) \right) \right|$$

$$B(r) = \frac{\alpha^\pm}{8} \log \left| \sin \left( 8K \left( r - \frac{R}{2} \right) \right) \right| \mp 2 \log \left| \tan \left( 8K \left( r - \frac{R}{2} \right) \right) \right| + \phi_0$$

$$\chi(r) = \frac{\mu}{8} \log \left| \sin \left( 8K \left( r - \frac{R}{2} \right) \right) \right| \mp \frac{\alpha^\pm}{8} \log \left| \tan \left( 8K \left( r - \frac{R}{2} \right) \right) \right|$$

$$U(y) = \frac{1}{8} \log (\cosh(8Ky))$$

$$V(y) = \alpha^\pm U(y)$$

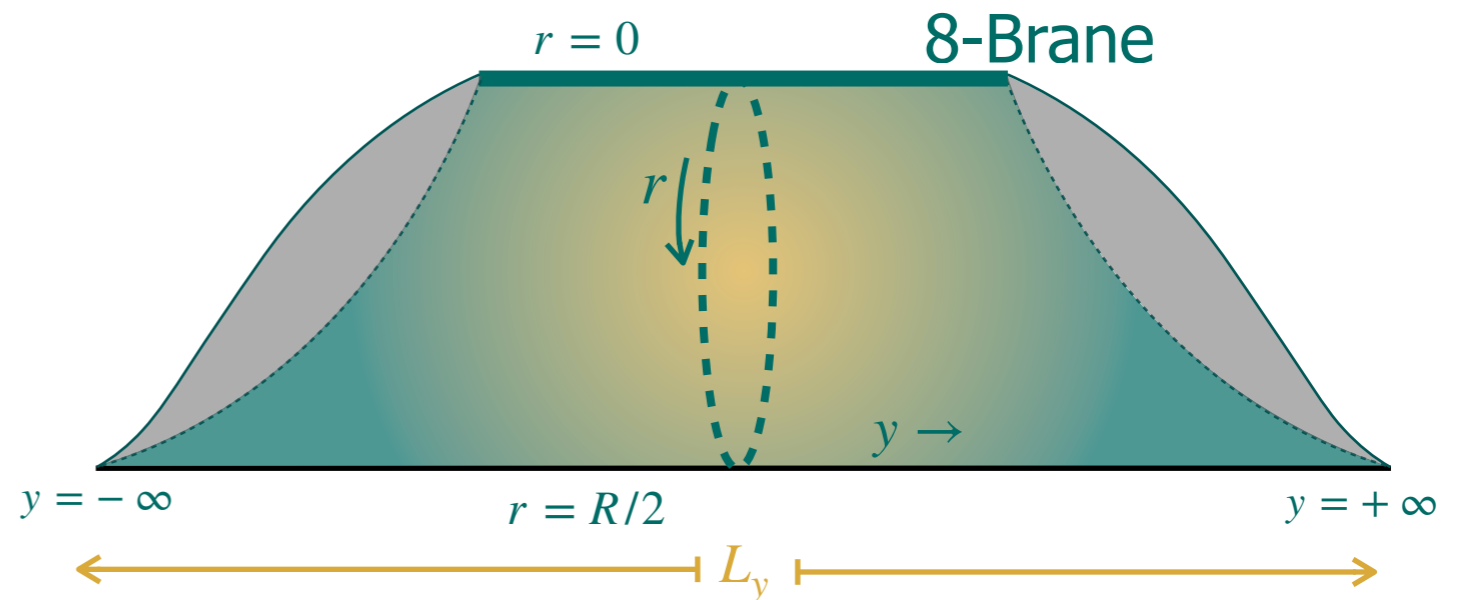
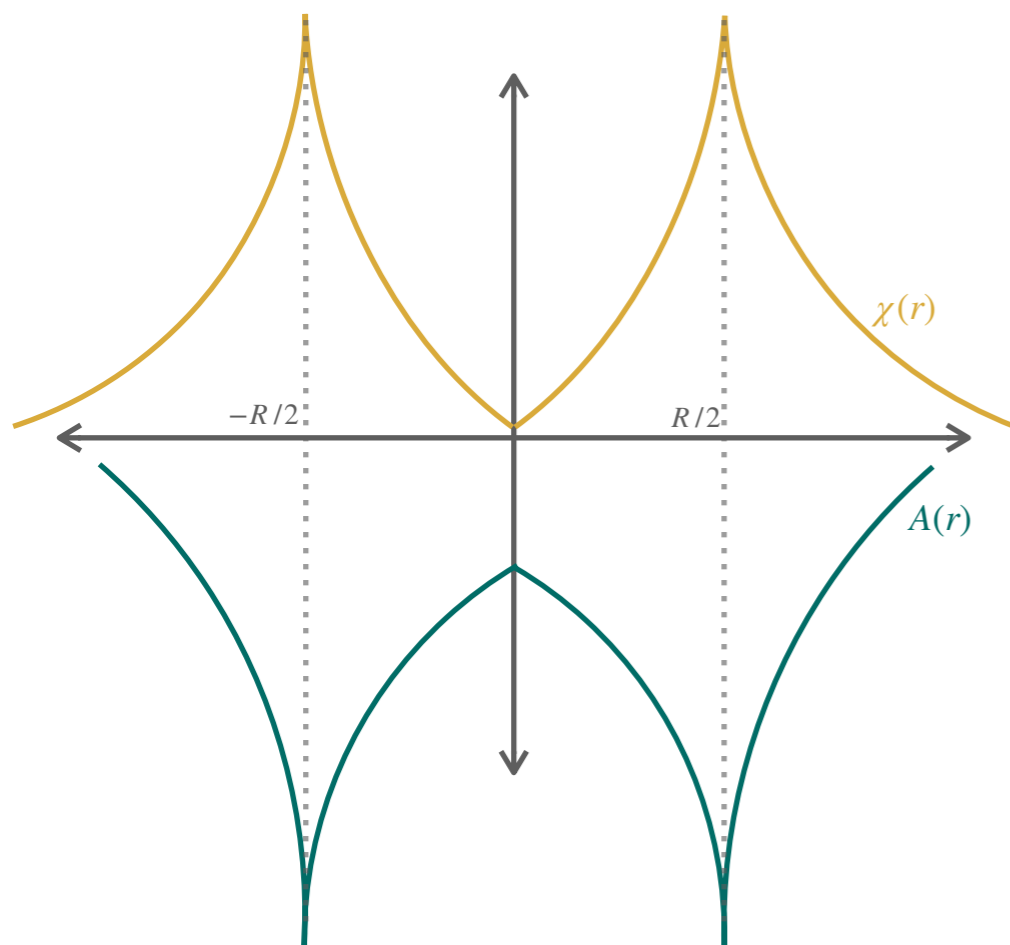
$$\psi(y) = -\frac{5\alpha^\pm}{4} U(y)$$

# Qualitative behaviour

r-direction spontaneously compactified on  $S^1$  with radius  $R$ :  $e^{\frac{5}{4}\phi_0} \sim \frac{1}{\lambda R}$   
 $\lambda R \sim \kappa_{10}^2 T$

Logarithmic singularities at  $r = \pm \frac{R}{2}$ , string coupling diverges

y-direction: infinite length in sols  $I, II^-$ , becomes finite interval in sol  $II^+$



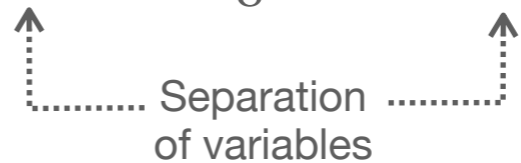
At  $(r, y) = (R/2, \pm \infty)$  :  $\left\{ \begin{array}{l} \Delta \sim L_y \sim e^{-\frac{5}{4}\sqrt{2}D} \\ \mathcal{R} \sim e^{\frac{5}{2}\sqrt{2}D} \end{array} \right\} \delta = \frac{5\sqrt{2}}{2}$



# ETW Defects

**Input:** 8-dimensional defect : log-singularity,  $S^1$  direction capped off  
 Poincaré symmetry along the brane preserved  
 2d transversal rotational symmetry broken

**Non-Isotropic Solution Ansatz:**  $ds^2 = e^{2\hat{\mathcal{A}}(\rho,\varphi)} ds_8^2 + e^{2\hat{\mathcal{B}}(\rho,\varphi)} (d\rho^2 + \rho^2 d\varphi^2)$ .


  
 Separation of variables

**Solutions:**  $\hat{A}(\rho) = \frac{1}{8} \log \left| \cosh(8\hat{K} \log(\frac{\rho}{\rho_0})) \right|$ ,  $\hat{B}(\rho) = -\log(\frac{\rho}{\rho_0}) + (\frac{\hat{\alpha}^2}{32} - \frac{7}{2})\hat{A}(\rho)$

$$\hat{\chi}(\rho) = \hat{\alpha}\hat{A}(\rho)$$

$$\hat{U}(\varphi) = \frac{1}{8} \log \left| \cos(8\hat{K}\varphi) \right|, \quad \hat{V}(\varphi) = -\left(\frac{\hat{\alpha}^2}{32} - \frac{9}{2}\right)\hat{U}(\varphi) + \frac{\hat{\alpha}}{16}\hat{\psi}(\varphi)$$

$$\hat{\psi}(\varphi) = \frac{\hat{\alpha}}{8} \log \left| \cos(8\hat{K}\varphi) \right| \pm 2 \log \left| \tan(4\hat{K}\varphi + \frac{\pi}{4}) \right|$$


  
*ETW*  $7^\pm$  solutions

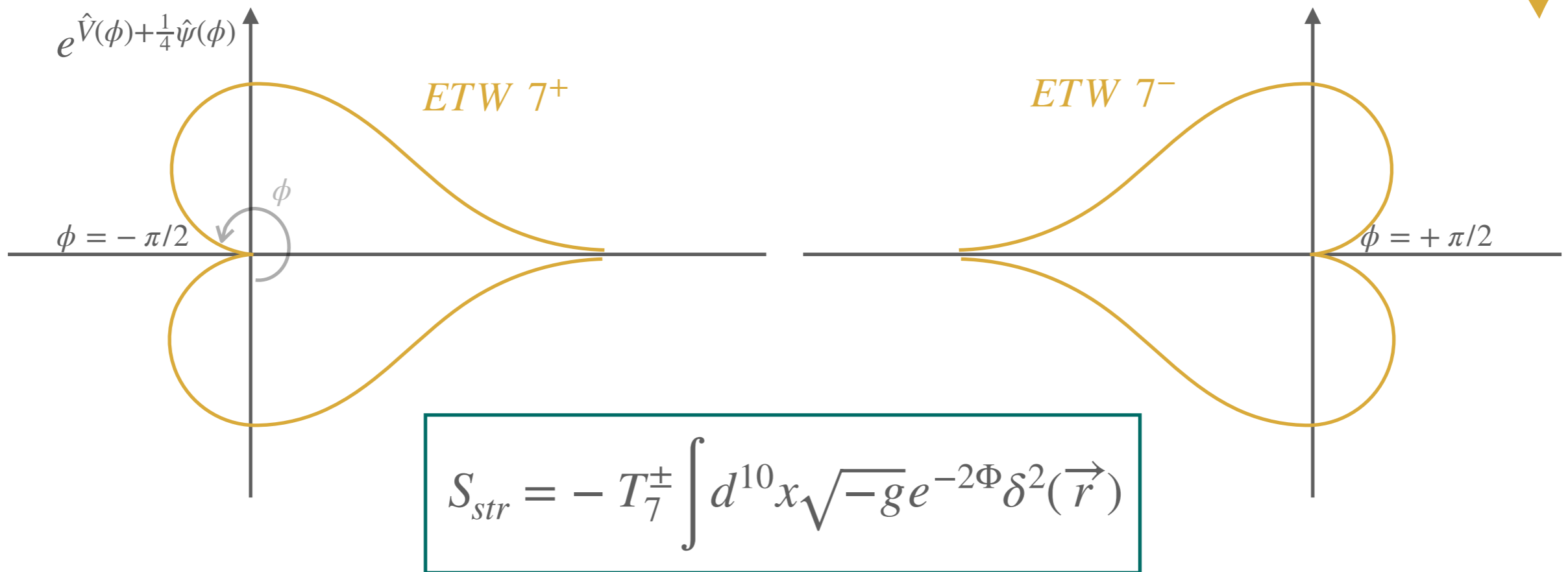
# Qualitative behavior

Logarithmic singularities at  $\rho = 0$ , string coupling diverges

For appropriate constant ( $\hat{\alpha} = \alpha^+$ ) same scaling as 9d defect

Dynamical Cobordism scaling satisfied:  $\Delta \sim e^{-\frac{5}{4}\sqrt{2}D}$ ,  $\mathcal{R} \sim e^{\frac{5}{2}\sqrt{2}D}$

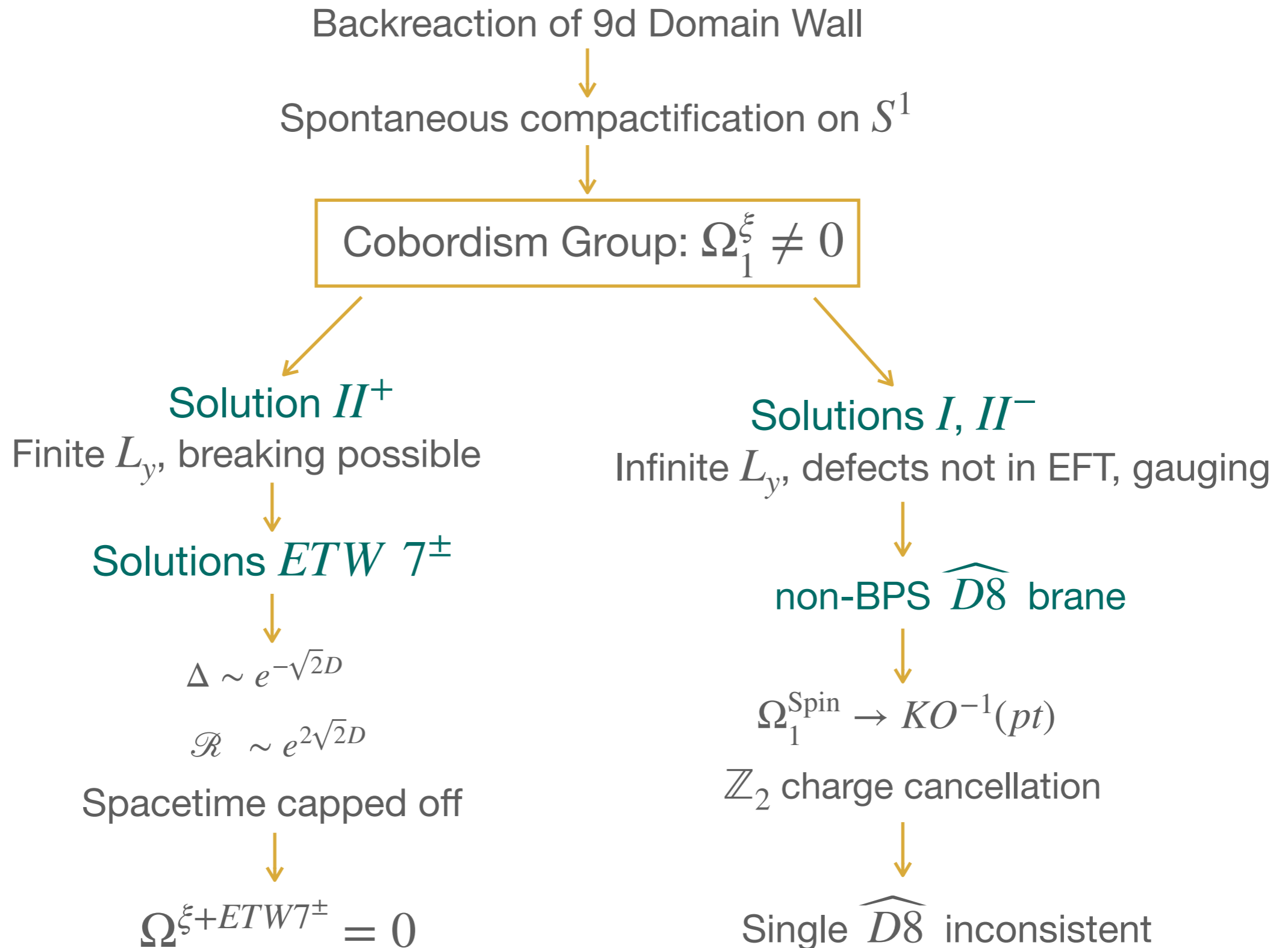
$$\delta = \frac{5\sqrt{2}}{2} \checkmark$$



$$S_{str} = -T_7^{\pm} \int d^{10}x \sqrt{-g} e^{-2\Phi} \delta^2(\vec{r})$$

with  $\kappa_{10}^2 T_7 = 2\pi$

# Cobordism Interpretation



# Generalisation and more evidence

[Blumenhagen, Kneissl, Wang '23]

**Setup:** Gauge neutral, non-supersymmetric codimension-1 object w/ arbitrary dilaton coupling in arbitrary dimensions  $D$

**Physical realisation:** Generalised BF model

**Action:** 
$$S = \frac{1}{2\kappa_D^2} \int d^D x \sqrt{-G} \left( \mathcal{R} - \frac{1}{2} (\partial\Phi)^2 \right) - T \int d^D x \sqrt{-G} e^{a\Phi} \delta(r) .$$

**ETWs:** neutral  $(d-2)$ -dimensional branes

# Generalisation and more evidence

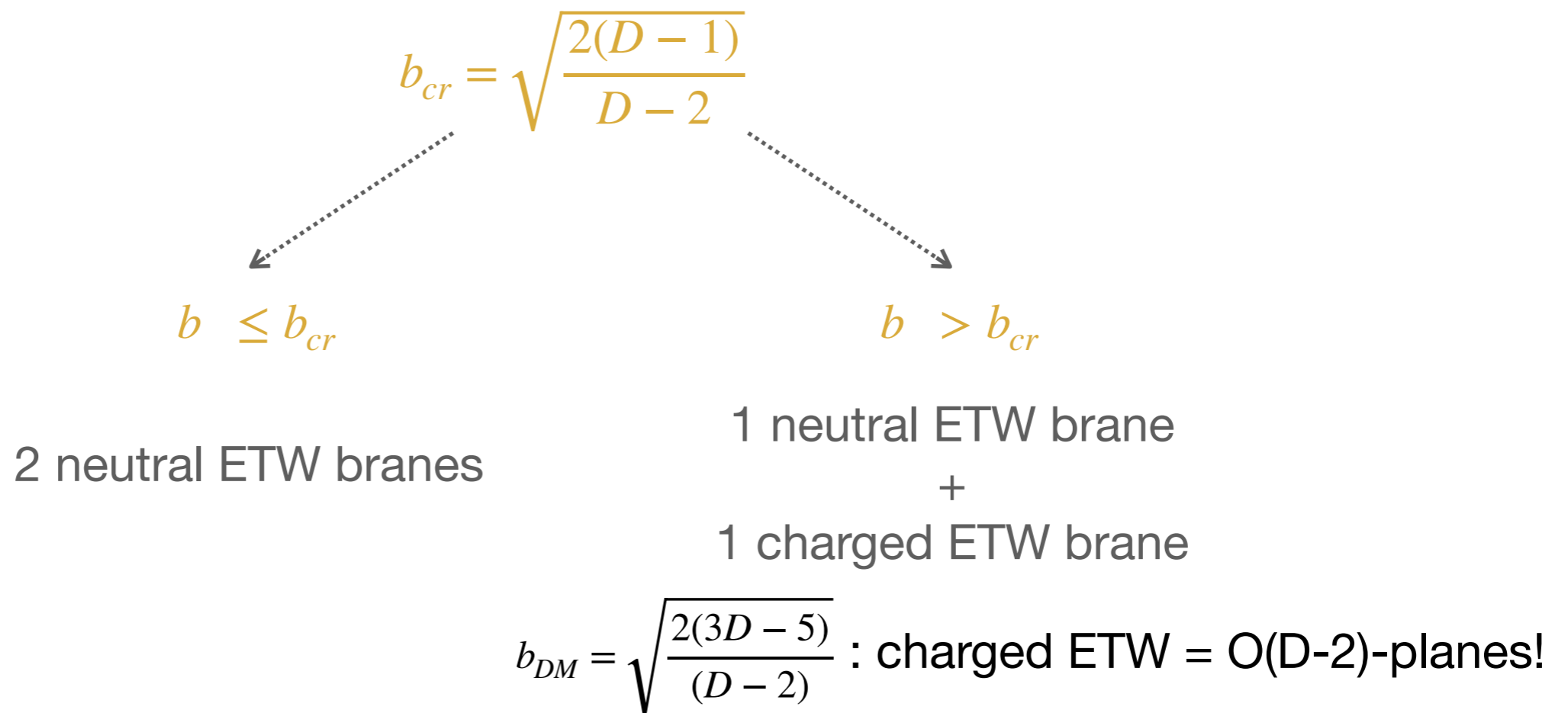
[Blumenhagen, Kneissl, Wang '23]

**Setup:** Gauge neutral, non-supersymmetric codimension-0 object  
w/ arbitrary dilaton coupling in arbitrary dimensions  $D$

**Physical realisation:** Generalised DM model,  $SO(16) \times SO(16)$ , IIA with Romans mass

**Action:** 
$$S = \frac{1}{2\kappa_D^2} \int d^D x \sqrt{-G} \left( \mathcal{R} - \frac{1}{2} (\partial\Phi)^2 \right) - T \int d^D x \sqrt{-G} e^{b\Phi} .$$

*Important difference: depending on value of  $b$ , charged ETWs also required!!*



# Summary and Outlook

- ETW-branes can be detected within the Dynamical Cobordism framework.
- Previously puzzling singularities in non-SUSY setups explained  
→ dynamical cobordism scalings verified.
- Explicit descriptions of cobordism-predicted ETW branes possible.  
→ possibly new objects?
- Independent verification of new ETW branes? Stability?
- Relation with other cobordism detected ETW-objects?
- Physical meaning of  $\delta$ , relation with Distance Conjecture?
- Universal Local Dynamical Cobordism description for higher codimension objects?

*Thank you for the attention!*

# Gauging Cobordism Charges

K-theory charges are gauged  
[Freed '00]

Gauged symmetries accompanied by tadpole conditions:  $0 = \int_M dF_{n-1} = \int_M J_n$

Idea: Cobordism and K-theory can be combined in single tadpole

In practice: add K-theory + cobordism contributions

$$0 = \int_M dF_{n-1} = \sum_{j \in \text{def}} \int_M Q_j \delta^{(n)}(\Delta_{10-n,j}) + \sum_{i \in \text{inv}} a_i \mu_{n,i}$$

↑
↑  
D-brane
Geometric  
contributions
contributions

Bottom-up way to construct tadpoles without string theoretical input!  
(up to constants!)

# An example:

Focus on interplay of  $K^{-6}(pt) = \mathbb{Z}$ ,  $\Omega_6(pt) = 2\mathbb{Z}$ :

Quantities entering tadpole:

D3-branes

Cobordism invariants

$$0 = \sum_i \int_M N_i \delta^{(6)}(\Delta_{4,i}) + a_1^{(6)} \frac{c_2(M)c_1(M)}{24} + a_2^{(6)} \frac{c_1^3(M)}{2}.$$

D3-tadpole for F-theory compactified on CY 4-fold  $Y$  over base  $B$ :

[Sethi, Vafa, Witten '96]

$$\int_B \sum_i \delta^{(6)}(\Delta_{4,i}) + \dots = \frac{\chi}{24} = \int_B \left( \frac{1}{2} c_2(B)c_1(B) + 15c_1^3(B) \right)$$

Perfectly reproduced for  $a_1^{(6)} = -12$ ,  $a_2^{(6)} = -30$