

# Light-ray operators in gauge and gravity theory: new infrared finite local observables

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based on 2012.01406 with A.Pokraka  
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- 2 Local momentum and angular momentum flow in gauge and gravity theory
- 3 The light-ray basis, Cordova-Shao algebra and angular momentum in QCD
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# Why light-ray operators?

Open questions from Ian Moulton: (talk 08/07)

- What other matrix elements of light-ray operators are well defined?  
Relation to IR finite S-matrices?
- What are the implications of asymptotic symmetry algebras?  
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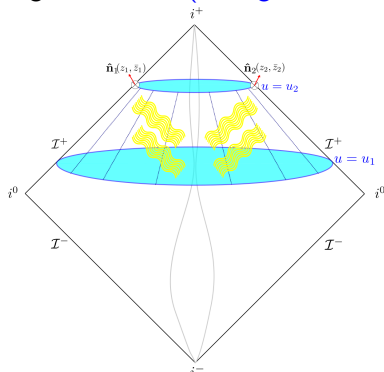
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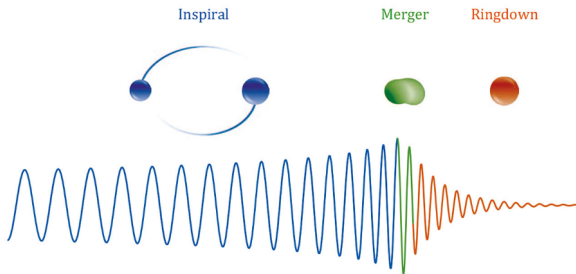
- **Light-ray operators** naturally arise by integrating Einstein equations at null infinity along the light-cone time ("charge densities")

- **Phenomenologically relevant for EM, QCD and gravity**: one can define a notion of localized detector collecting quanta of radiation in a direction  $\hat{n}$



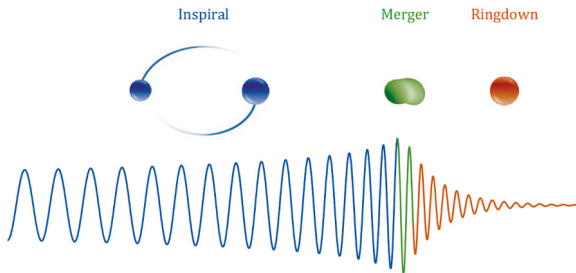
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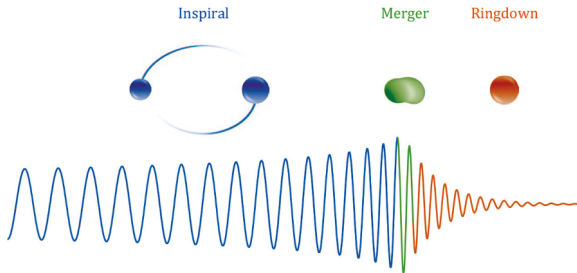
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# Gravitational waves: a theoretical lab for detector operators

- Recent interest in **GWs emitted in the merging of compact binary systems** (black holes, neutron stars, . . . ): we need accurate **waveform templates**



- The **waveform is the most natural local observable!** How about the gravitational **energy and angular momentum flow**?
- (Energy-energy) correlators are a smoking gun signature of quantum physics:** classical physics ( $\hbar \rightarrow 0$ ) implies **factorization of  $n$ -pt correlation functions**

# Momentum and angular momentum flow in gauge theory

- Framework: Electromagnetism or Einstein gravity in  $d = 4$



# Momentum and angular momentum flow in gauge theory

- Framework: **Electromagnetism or Einstein gravity in  $d = 4$**
- Let's start with QED. Using the standard **mode expansion of gauge fields**

$$\mathbb{A}_\mu(x) = \frac{1}{\sqrt{\hbar}} \sum_{\sigma=\pm} \int d\Phi(k) \left[ a_\sigma(k) \varepsilon_\mu^{\sigma*}(k) e^{-i\frac{k \cdot x}{\hbar}} + a_\sigma^\dagger(k) \varepsilon_\mu^\sigma(k) e^{i\frac{k \cdot x}{\hbar}} \right],$$

local **gauge-invariant momentum** [Sveshnikov, Tkachov; Sterman, Korchemsky] and **angular momentum** [RG, Ilderton] **flow**  $\mathcal{P}^\mu$ ,  $\mathcal{N}^{\mu\nu}$  can be defined from  $\mathcal{P}^\mu \sim \int d^3x \hat{\delta}^2(\Omega - \Omega_{\hat{n}}) T^{\mu 0}$ ,  $\mathcal{N}^{\mu\nu} \sim \int d^3x \hat{\delta}^2(\Omega - \Omega_{\hat{n}}) x^{[\mu} T^{\nu]0}$ ,

$$\mathcal{P}^\mu = \sum_{\sigma=\pm} \int d\Phi(k) \hat{\delta}^2(\Omega - \Omega_{\hat{n}}) k^\mu a_\sigma^\dagger(k) a_\sigma(k),$$

$$\mathcal{N}_{\text{QED}}^{\mu\nu} = \sum_{\sigma=\pm} \int d\Phi(k) \hat{\delta}^2(\Omega - \Omega_{\hat{n}}) \varepsilon_\sigma^\alpha(k) a_\sigma^\dagger(k) \left[ (\mathcal{J}_{\text{QED}})^{\mu\nu}_{\alpha\beta} \right] \varepsilon_\sigma^{*\beta}(k) a_\sigma(k),$$

$$(\mathcal{J}_{\text{QED}})^{\mu\nu}_{\alpha\beta} = -i\eta_{\alpha\beta} k^{[\mu} \frac{\overleftrightarrow{\partial}}{\partial k_{\nu]}} - i\delta_\alpha^{[\mu} \delta_\beta^{\nu]},$$

Non-trivial **gauge-invariance under  $\varepsilon_\sigma^\alpha(k) \rightarrow \varepsilon_\sigma^\alpha(k) + \xi k^\alpha!$**

# Momentum and angular momentum flow in gravity theory

- In **gravity**, the story is similar by expanding the **metric**  $g_{\mu\nu} = \eta_{\mu\nu} + \kappa\mathfrak{h}_{\mu\nu}$

$$\mathfrak{h}_{\mu\nu}(x) = \frac{1}{\sqrt{\hbar}} \sum_{\sigma=\pm} \int d\Phi(k) \left[ a_{\sigma}(k) \varepsilon_{\mu\nu}^{\sigma*}(k) e^{-i\frac{k\cdot x}{\hbar}} + a_{\sigma}^{\dagger}(k) \varepsilon_{\mu\nu}^{\sigma}(k) e^{i\frac{k\cdot x}{\hbar}} \right],$$

to define **local momentum and angular momentum flow** [RG, Ilderton]

$$\mathcal{P}^{\mu} = \sum_{\sigma=\pm} \int d\Phi(k) \hat{\delta}^2(\Omega - \Omega_{\hat{n}}) k^{\mu} a_{\sigma}^{\dagger}(k) a_{\sigma}(k),$$

$$\mathcal{N}_{\text{GR}}^{\mu\nu} = \sum_{\sigma=\pm} \int d\Phi(k) \hat{\delta}^2(\Omega - \Omega_{\hat{n}}) \varepsilon_{\sigma}^{\alpha\alpha'}(k) a_{\sigma}^{\dagger}(k) \left[ (\mathcal{J}_{\text{GR}})_{\alpha\alpha'\beta\beta'}^{\mu\nu} \right] \varepsilon_{\sigma}^{*\beta\beta'}(k) a_{\sigma}(k),$$

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The global  $\mathcal{J}_{\text{GR}}^{\mu\nu} = \int d\Omega \mathcal{N}_{\text{GR}}^{\mu\nu}$  for the classical two-body problem has received a lot of attention (zero energy gravitons are subtle!) [Damour; Manohar, Ridgway, Shen; Heissenberg, Russo, Veneziano], much less the local!

# Connection with asymptotic light-ray operators

- Projecting the local operators on the tetrad basis  $(n^\mu, \bar{n}^\mu, m^\mu, \bar{m}^\mu)$  at  $\mathcal{I}^+$  in Bondi-like flat null coordinates  $(u, r, z, \bar{z})$ , we obtain the light-ray basis [Korchemsky, Oderda, Sterman; Hofman, Maldacena; Cordova, Shao, RG, Pokraka]

$$\mathcal{E}(\hat{n}) = \int_{-\infty}^{+\infty} du \lim_{r \rightarrow \infty} r^2 T_{uu}(u, r, z_{\hat{n}}, \bar{z}_{\hat{n}}) \stackrel{\text{saddle-point}}{\sim} \bar{n}_\mu \mathcal{P}^\mu,$$

$$\mathcal{K}(\hat{n}) = \int_{-\infty}^{+\infty} du u \lim_{r \rightarrow \infty} r^2 T_{uu}(u, r, z_{\hat{n}}, \bar{z}_{\hat{n}}) \stackrel{\text{saddle-point}}{\sim} \bar{n}_\nu n_\mu \mathcal{N}^{\mu\nu},$$

$$\mathcal{N}_z(\hat{n}) = \int_{-\infty}^{+\infty} du \lim_{r \rightarrow \infty} r^2 T_{uz}(u, r, z_{\hat{n}}, \bar{z}_{\hat{n}}) \stackrel{\text{saddle-point}}{\sim} \bar{n}_\mu m_\nu \mathcal{N}^{\mu\nu},$$

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- These operators appear by integrating Einstein's equations near  $\mathcal{I}^+$ : natural interpretation of local measurements of energy and angular momentum!
- Light-ray algebra consistent with complexified Cordova-Shao algebra!

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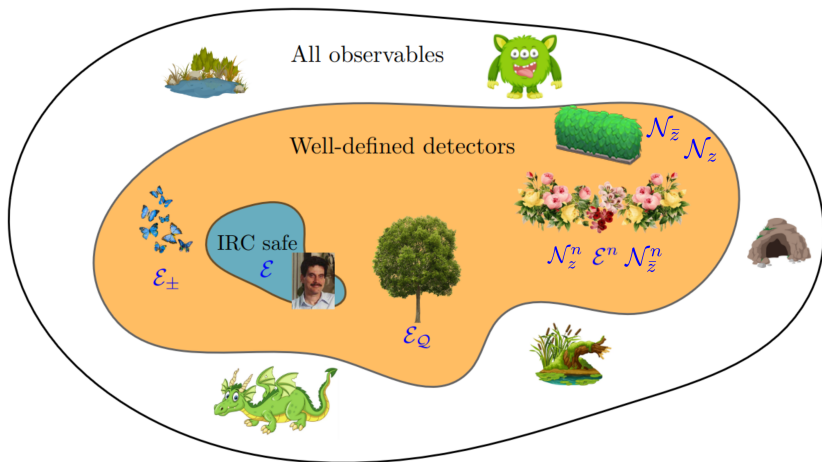
- **Angular momentum flow?** Subtle issue: **mixing between radiative (on-shell) and coulombic (potential) modes of the field** [Ashtekar,Bonga;RG,Pokraka]

$$\begin{aligned} \mathcal{N}_z \sim & \sum_{\sigma=\pm} \int d\Phi(k) \\ & \times \left\{ \underbrace{\hat{\delta}^2(z_{\hat{n}} - z_{\hat{k}}) a_\sigma^{a\dagger}(k) \left[ i \overset{\leftrightarrow}{\partial}_{z_n} \right] a_\sigma^a(k)}_{\sim \text{Orbital contribution}} + \underbrace{\sigma \left[ i \partial_{z_n} \delta^2(z_{\hat{n}} - z_{\hat{k}}) \right] a_\sigma^{a\dagger}(k) a_\sigma^a(k)}_{\sim \text{Spin contribution}} \right. \\ & \left. + 2g A_z^a(u = +\infty, z, \bar{z}) \underbrace{\left[ -if^{abc} a_\sigma^{b\dagger}(k) a_\sigma^c(k) \right]}_{\text{Density of color charge}} \right\} \end{aligned}$$

More to be understood! Potential implications for colliders?



# Expanding the space of detector observables?



# Compton scattering in EM and gravity: $\langle \mathcal{P}^\mu \rangle$ and $\langle \mathcal{N}^{\mu\nu} \rangle$

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- Classically, for a coherent state with real amplitude  $A_\pm$  and a phase  $\delta_\pm$

$$|\gamma\rangle = \mathcal{N}_\gamma \exp\left(\sum_{\sigma=\pm} \int d\Phi(k) \gamma^\sigma(k) a_\sigma^\dagger(k)\right) |0\rangle, \quad m^\mu \varepsilon_\mu^{*(\pm)} \gamma^{(\pm)} = \underbrace{A_\pm}_{\text{amplitude}} \underbrace{e^{i\delta_\pm}}_{\text{phase}}.$$

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- We find that  $\langle \mathcal{P}^\mu \rangle$  encodes the amplitude of the outgoing wave while  $\langle \mathcal{N}^{\mu\nu} \rangle$  is sensitive to both the amplitude and the phase

$$\langle \gamma | \mathcal{P}^\mu | \gamma \rangle = \sum_{\sigma=\pm} \int \frac{d\omega \omega}{4\pi} \omega n^\mu |A_\sigma(\omega, \omega \hat{n})|^2,$$

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- What about quantum effects?  $N$ -pt correlators of local observables?

# Gravitational event shapes: energy-energy correlators

- The action of the "shear-inclusive ANEC operator" at infinity [Wall]

$$\begin{aligned}\mathcal{E}(\hat{n}) \prod_{\alpha \in \text{species}} |X_\alpha\rangle &= \int_{-\infty}^{+\infty} du \left[ \lim_{r \rightarrow \infty} r^2 T_{uu}^{\text{matter}} + \frac{1}{16\pi G} N_{\zeta\zeta} N^{\zeta\zeta} \right] \prod_{\alpha \in \text{species}} |X_\alpha\rangle \\ &= \underbrace{\sum_{i=1}^n E_i \delta^2(\Omega_{\hat{k}_i} - \Omega_{\hat{n}})}_{w_{\mathcal{E}}(\hat{n})} \prod_{\alpha \in \text{species}} |X_\alpha\rangle.\end{aligned}$$

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- Gravitational energy event shapes are defined as ( $|\text{out}\rangle = S |\text{in}\rangle$ ) [RG, Pokraka]

$$\langle \text{out} | \mathcal{E}(\hat{n}) | \text{out} \rangle = \sum_{\mathbf{X}} \langle \text{in} | S^\dagger \mathcal{E}(\hat{n}) | \mathbf{X} \rangle \langle \mathbf{X} | S | \text{in} \rangle = \sum_{\mathbf{X}} w_{\mathcal{E}}(\hat{n}) |\langle \mathbf{X} | S | \text{in} \rangle|^2,$$

where  $\sum_{\mathbf{X}}$  is the phase space integration over the final states.

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where  $\sum_{\mathbf{X}}$  is the phase space integration over the final states.

- Upshot: we can study  $\langle \text{out} | \mathcal{E}(\hat{n}_1) \mathcal{E}(\hat{n}_2) | \text{out} \rangle$  not only in gauge theory but also in gravity: **power-suppression guarantees infrared-finiteness**



# Factorization of gravity EEC in the classical limit $\hbar \rightarrow 0$

- What can we learn from gravity EECs of relevance for GWs (inspiral phase)?

# Factorization of gravity EEC in the classical limit $\hbar \rightarrow 0$

- What can we learn from **gravity EECs** of **relevance for GWs (inspiral phase)**?
- **Classically** ( $\hbar \sim 0$ ), there is a **unique field** determined by the classical eoms: the **classical S-matrix** exponentiates [Cristofoli, RG, Moynihan, O'Connell, Ross, Sergola, White; DiVecchia, Heissenberg, Russo, Veneziano]

$$\begin{array}{c}
 k_1 \quad \dots \quad k_N \\
 \nearrow \quad \quad \searrow \\
 p_2 \quad \quad \quad p'_2 \\
 \longrightarrow \quad \quad \longrightarrow \\
 \text{---} \quad \text{---} \\
 \text{---} \quad \text{---} \\
 \text{---} \quad \text{---} \\
 p_1 \quad \quad \quad p'_1 \\
 \longrightarrow \quad \quad \longrightarrow \\
 \text{---} \quad \text{---} \\
 \text{---} \quad \text{---} \\
 \text{---} \quad \text{---} \\
 \tilde{\mathcal{S}}^{\text{cl}}
 \end{array}
 \sim e^{\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \tilde{\mathcal{K}}^{\text{cl}} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}} e^{\int d\Phi(k) \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \tilde{\mathcal{K}}_{5,\mathcal{R}}^{\text{cl}} \\ \text{---} \\ \text{---} \\ \text{---} \end{array}} a^\dagger(k) + h.c.$$

and because of the coherent state structure the **EEC factorizes** [RG, Pokraka]

$$\langle \text{out} | \mathcal{E}(\hat{\mathbf{n}}_1) \mathcal{E}(\hat{\mathbf{n}}_2) | \text{out} \rangle \Big|_{\hbar \rightarrow 0} = \langle \text{out} | \mathcal{E}(\hat{\mathbf{n}}_1) | \text{out} \rangle \Big|_{\hbar \rightarrow 0} \langle \text{out} | \mathcal{E}(\hat{\mathbf{n}}_2) | \text{out} \rangle \Big|_{\hbar \rightarrow 0}$$



# Summary and future directions

- We provide **novel gauge-invariant expressions for the momentum and angular momentum flow**, defining a **space of local observables** in gauge and gravity

Theory	Global radiative observables	Local radiative observables
QED	$\sigma, \mathbb{K}_{\text{QED}}^\mu, \mathbb{J}_{\text{QED}}^{\mu\nu}$	$\frac{d\sigma}{d\Omega}, \mathcal{P}_{\text{QED}}^\mu, \mathcal{N}_{\text{QED}}^{\mu\nu}$
GR	$\sigma, \mathbb{K}_{\text{GR}}^\mu, \mathbb{J}_{\text{GR}}^{\mu\nu}$	$\frac{d\sigma}{d\Omega}, \mathcal{P}_{\text{GR}}^\mu, \mathcal{N}_{\text{GR}}^{\mu\nu}$

- We made contact with the **light-ray basis in the BMS framework** (via the saddle-point method at  $r \rightarrow +\infty$ ), **verifying Cordova-Shao algebra** and highlighting some **subtleties for the angular momentum flow in QCD**
- We study the **properties of  $\langle \mathcal{P}^\mu \rangle$  and  $\langle \mathcal{N}^{\mu\nu} \rangle$** : **classical factorization of correlation functions**, relation with the **amplitude and phase of the waveform**, **singularity structure at high-energy** (better than global operators!)
- Lots to do**: Angular momentum in QCD? IR-finite S-matrix for gravity and related observables? Quantum effects for gravity correlators of local observables (probing non-coherence of GWs)? Collider OPE tools for gravity?