Energy Correlators at the Collider Frontier

Ian Moult



Exclusive Processes: S-Matrix

• Tremendous progress in the understanding of exclusive scattering processes: analytic structure, multi-loop perturbative data, amplituhedron, S-matrix bootstrap,...



Practical Outcome: Ability to accurately describe complicated SM scattering processes.

The High Multiplicity Regime

- A complementary regime: high multiplicity
 - Collisions with $E \gg m_{gap}$
 - Conformal Field Theories





• Good observables are correlations in fluxes at (null) infinity.

Motivation

- How can we characterize a theory using asymptotic data?
- Theoretical motivation:
 - What is the space of observables at null infinity?
 - How are they related to (C)FT data?
 - How do we constrain theories in the absence of S-matrix and/ or local ops (e.g. CFT coupled to gravity) [Maldacena, Zhiboedov]
- Phenomenological motivation:
 - Can we relate asymptotic measurements to parameters of the underlying theory? (couplings, transport coefficients,)
 - Can we identify universal features that can be computed to high precision?



• Wealth of collider data provides a practical testing ground.

Energy Correlators: Past

To make this idea more quantitative we define for any state a, an
"angular energy current" in the
$$e^+e^-$$
 CM frame:

$$j_a(\Omega) = \sum_{i=1}^{n_a} \eta_i \delta(\Omega - w_i)$$
(1)
where the sum is over the n_a massless particles in a, with energies
 $\{\eta_i\}$ and momentum directions $\{w_i\}$ (w_i stands for angles θ_i and w_i).

Energy Flux

• The expectation value of energy flux, $\langle \mathcal{E}(\vec{n}_1) \rangle$, at specific angles on the celestial sphere is calculable in perturbation theory. (Sterman, 1975)



Our ensembles will thus be specified in terms of sets of jet-related states. To make this idea more quantitative we define for any state \underline{a} , an "angular energy current" in the $e^{+}e^{-}$ CM frame:

$$\begin{array}{l} n_{a} \\ j_{a}(\Omega) = \sum & \eta_{i} \delta(\Omega - \omega_{i}) \\ i = 1 \end{array}$$
 (1)

where the sum is over the n_g massless particles in <u>a</u>, with energies $\{n_{\underline{i}}\}$ and momentum directions $\{w_{\underline{i}}\}$ $\{w_{\underline{i}}\}$ stands for angles $\theta_{\underline{i}}$ and $\alpha_{\underline{i}}\}$. Jetrelated states have the same j(fi). Each group of particles with collinear momenta may be described as a jet, and any set of jet-related states is characterized by the number of jets, as well as their energies and directions.

"Energy Flow becomes the focus of calculability"





Detector Operators

- Are now known to be part of a wide class of "detector operators".
- Significant recent progress in understanding the classification of detector operators, and their algebras and OPEs.

Sterman 1975:

To make this idea more quantitative we define for any state a, an "angular energy current" in the e^+e^- CM frame:

ANEC

 $j_a(\Omega) = \sum_{i=1}^{N_a} n_i \delta(\Omega - w_i)$ (1)

where the sum is over the n_a massless particles in <u>a</u>, with energies $\{T_{q}\}$ and momentum directions $\{w_{q}\}$ (w_{q} stands for angles θ_{q} and w_{q}).

[Korchemsky, Sterman] [Sveshnikov, Tkachov] [Hofman, Maldacena]

Maldacena]

$$\mathcal{E}(\vec{n}) = \lim_{r \to \infty} r^2 \int_{-\infty}^{\infty} dt \ n^i T_{0i}(t, r\vec{n})$$

Light Transform

[Kravchuk, Simmons Duffin]

$$\mathcal{E}(\vec{n}) = 2 \ \mathbf{L}[T](\infty,z)|_{z=(1,\vec{n})}$$

Ō

$$\mathbf{L}[\mathcal{O}](x, z) = \int_{-\infty}^{\infty} d\alpha (-\alpha)^{-\Delta - J} \mathcal{O}\left(x - \frac{z}{\alpha}, z\right)$$



$$\begin{bmatrix} \bigoplus_{i}, \bigoplus_{j} \end{bmatrix} = \sum D_{ijk} \bigoplus_{k}$$
$$\bigoplus_{i} (\theta_{1}) \bullet \bigoplus_{j} (\theta_{2}) = \sum (\theta_{1} - \theta_{2})^{\gamma} C_{ijk} \bigoplus_{k} (\theta_{1})$$

Detector Operators

• Now play a central role in many areas of theoretical physics.



ENERGY OPERATORS IN PARTICLE PHYSICS, QUANTUM FIELD THEORY AND GRAVITY

December 16-20, 2024

FOR

SIMONSCENTER Story Brook

Organized by: Thomas Hartman (Cornell), Zohar Komargodski (SCGP), Gregoire Mathys (EPFL), Ian Moult (Yale)

This sections will explore the record developments in the study of energy constraints pairs, the physics, Barant mills Theory QPT and gravity, In pairsic physics, constaints here: those of energy constraints are collider physics observables used for precision measurements of a standard Hodds. To TC, recording gravitation plan approach to the in particle, energy comparison are related to doublewase and causality in the ASC/TC compandences. These downer topics are all fuel together through energy constraints, in Falsy to an 4300 per standard bards and the together through energy constraints. We all those is advised and explore the standard bard approach and the or advised are energy than different fields and to able energy integrates and expendives.

For more information visit: scgp.stonybrook.edu



Energy Correlators

Correlators of energy flow operators (\$\mathcal{E}(\vec{n}_1) \cdots \mathcal{E}(\vec{n}_k)\$) characterize the final state in collider experiments in QCD (Ellis et al. 1977)



Energy Correlations in electron - Positron Annihilation: Testing QCD

C.Louie Basham (Washington U., Seattle), Lowell S. Brown (Washington U., Seattle), Stephen D. Ellis (Washington U., Seattle), Shetwin T. Love (Washington U., Seattle) Aug, 1978

13 pages Published in: Phys.Rev.Lett. 41 (1978) 1585 DOI: 10.1103/PhysRevLett.41.1585 Report number: RLO-1388-759 View in: OSTI Information Bridge Server, ADS Abstract Service





🖾 reference search 🛛 🕤 416 citations



Abstract

The possible role of gluons in hadronic processes as suggested by the guark-parton model and QCD is discussed and evaluated.

Introduction

The most dramatic and encouraging development in Strong Interaction physics in the past few years has been the emergence of a candidate field theory to describe these interactions. This situation is in sharp contrast to that which obtained ten years ago when, at least at one well known institution of higher learning, it was "taught" that field theory was irrelevant to Strong Interactiona¹¹.

Footnote and References

 Recollections of a Caltech graduate student, S. D. Ellis, unpublishable.

Energy Correlators: Beyond QCD

• Energy Correlators are interesting observables for characterizing generic quantum theories. (Hofman, Maldacena)

Conformal collider physics: Energy and charge correlations



• Well defined at weak coupling, strong coupling, in a CFT, coupled to gravity,

Energy Correlators: Simplicity in Perturbation Theory

• Physical observables can be beautiful: Belitsky, Hohenegger, Korchemsky, Sokatchev, Zhiboedov (2013) illustrated the simplicity of energy correlators in perturbation theory.



$$F(z;a) = aF_1(z) + a^2 \left[(1-z)F_2(z) + F_3(z) \right]$$

$$\begin{split} F_2(z) &= 4\sqrt{2} \left[L_2\left(-\sqrt{z}\right) - LL_2\left(\sqrt{z}\right) + \frac{\ln z}{2} \ln \left(\frac{1+\sqrt{z}}{1-\sqrt{z}}\right) \right] + (1+z) \left[2LL_2(z) + \ln^2(1-z) \right] + 2\ln(1-z) \ln \left(\frac{z}{1-z}\right) + z\frac{\pi^2}{3} \right] \\ F_3(z) &= \frac{1}{4} \left[(1-z)(1+2z) \left[\ln^2 \left(\frac{1+\sqrt{z}}{1-\sqrt{z}}\right) \ln \left(\frac{1-z}{z}\right) - 8LL_3\left(\frac{\sqrt{z}}{\sqrt{z-1}}\right) - 8LL_3\left(\frac{\sqrt{z}}{\sqrt{z+1}}\right) \right] - 4(z-4)LL_3(z) \\ &+ 6(3+3z-4z^2)LL_3\left(\frac{z}{z-1}\right) - 2z(1+4z)\zeta_3 + 2 \left[2(2z^2-z-2)\ln(1-z) + (3-4z)z\ln z \right] LL_2(z) \\ &+ \frac{1}{3} \ln^2(1-z) \left[4(3z^2-2z-1)\ln(1-z) + 3(3-4z)z\ln z \right] + \frac{\pi^2}{3} \left[2z^2 \ln z - (2z^2+z-2)\ln(1-z) \right] \right\}. \end{split}$$
(15)

- Since then, extended in many directions, by many authors:
 - Higher loop $\mathcal{N} = 4$: Henn, Sokatchev, Yan, Zhiboedov
 - QCD: Dixon, Shtabovenko, Yang, Zhu
 - Higher Point: Yan, Yang, Zhang

Energy Correlators at the Collider Frontier

 Transition from GeV→ TeV, and the experimental development of jet substructure, provides access to a new regime of QCD!



- Amazing dataset provides exciting opportunities:
 - Experiment Informing Theory: Measurement of correlators, lightray OPE, etc. to reveal interesting behavior.
 - 2 Theory Informing Experiment: New sophisticated theory techniques to better study the SM at colliders.

Energy Correlators at the Collider Frontier

• Spectacular recent progress bridging theory and experiment!



Energy Correlators: Present







Scaling Behavior



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Scaling Behavior in QFT

• Euclidean scaling behavior is well understood.

λ -point of Helium







$$\mathcal{O}(x)\mathcal{O}(0) = \\ \sum x^{\gamma_i} c_i \mathcal{O}_i$$

The OPE Limit of Lightray Operators

• Energy flow operators admit a Lorentzian OPE: "the lightray OPE"



- Scaling can be derived in generic (non-conformal) theories using factorization theorems. [Dixon, Moult, Zhu] See early work by [Konishi, Ukawa, Veneziano]
- Predicts universal scaling behavior in correlations of energy flux at energies $E\gg\Lambda_{\rm QCD}$.

Theory-Experiment Gap

Conformal collider physics: Energy and charge correlations

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Still true as of 2023...

Scaling Behavior in Jets

• Scaling measured inside jets by STAR, ALICE and CMS from 15 GeV to 1784 GeV:

An experimental realization of the detector OPE!



• Can we accurately extract anomalous exponents of different detectors?

The Spectrum of a Jet

• The light-ray OPE predicts that the N-point correlators develop an anomalous scaling that depends on N. [Maldacena, Hofman] [Chen, Moult Zhang, Zhu]



• Directly probes the spectrum of (twist-2) lightray operators from asymptotic energy flux.

Anomalous Scaling



• Uses scaling anomalous dimensions at three-loop order.

Beautiful quantitative test of QFT!

The Strong Coupling

• Use scaling to extract value of the strong coupling constant α_s at 4% accuracy.



• Very clear target for improved perturbative calculations. e.g. NNLO $2 \rightarrow 3$ hard functions, NP corrections, ... not yet included.

Strings 2024 Talk by David Simmons-Duffin

CMS determination of α_s from $\langle EEE \rangle / \langle EE \rangle$ [CMS '24]



• First study of lightray OPE in any system!

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Higher Point Functions in Energy Flux



Multipoint Correlators

• Higher-point correlators probe detailed aspects of the underlying microscopic interactions. e.g. CMB three-point functions allow to distinguish models of inflation.



• What is the structure of higher-point functions of energy flux?

Multipoint Correlators

• The only explicit results for correlators with N>2 were the remarkable strong coupling results of Hofman and Maldacena:

$$\begin{split} \langle \mathcal{E}(\vec{n}_1) \cdots \mathcal{E}(\vec{n}_n) \rangle &= \left(\frac{q}{4\pi}\right)^n \left[1 + \sum_{i < j} \frac{6\pi^2}{\lambda} [(\vec{n}_i . \vec{n}_j)^2 - \frac{1}{3}] + \right. \\ &\left. + \frac{\beta}{\lambda^{3/2}} [\sum_{i < j < k} (\vec{n}_i . \vec{n}_j) (\vec{n}_j . \vec{n}_k) (\vec{n}_i . \vec{n}_k) + \cdots] + o(\lambda^{-2}) \right] \end{split}$$

• The wealth of techniques developed to compute perturbative scattering amplitudes can be applied to multi-point correlators at weak coupling.



Correlators in Perturbation Theory

- Two approaches to calculate energy correlators:
 - 1 Light transforming N-point functions of stress tensors:

$$\langle 0|\mathcal{O}^{\dagger}T\cdots T\mathcal{O}|0\rangle \rightarrow \langle 0|\mathcal{O}^{\dagger}\mathcal{E}(\vec{n}_{1})\cdots \mathcal{E}(\vec{n}_{k})\mathcal{O}|0\rangle$$

Two Point NLO in $\mathcal{N} = 4$: [Belitsky, Hohenegger, Korchemsky, Sokatchev, Zhiboedov] Two Point NNLO in $\mathcal{N} = 4$: [Henn, Sokatchev, Yan, Zhiboedov] LO Charge-Charge Correlator in QCD: [Chicherin, Henn, Sokatchev, Yan]

2 Perturbative phase space integrals using (squared) form factors:



$$\frac{\langle \Psi | \mathcal{E}(\vec{n}_1) \cdots \mathcal{E}(\vec{n}_k) | \Psi \rangle}{\langle \Psi | \Psi \rangle} = \sum_{i_1, \dots, i_k} \int d\sigma \prod_{j=1}^k E_{i_j} \delta(\vec{n}_j - \vec{p}_{i_j}/p_{i_j}^0)$$

Two Point LO in QCD: [Basham, Ellis, Brown, Love] Two Point NLO in QCD: [Dixon, Luo, Shtabovenko, Yang, Zhu] Three Point Collinear LO in QCD: [Chen, Luo, Moult, Yang, Zhang, Zhu] Three Point General Angle LO in $\mathcal{N} = 4$: [Yan, Zhang] Three Point General Angle LO in QCD: [Yang, Zhang] Four Point Collinear LO in $\mathcal{N} = 4$: [Chicherin, Moult, Sokatchev, Yan, Zhu]

Correlators in Perturbation Theory

• For generic angles, the correlator depends on the cross ratios $\zeta_{ij} = \frac{1 - \cos \theta_{ij}}{2}$, and the source.



- In the collinear (OPE) limit, $\zeta_{ij} \rightarrow 0$, it becomes a function of 2(N-2) variables that is independent of the source.
- The LO contribution to the N-point function is given by a *finite* integral in (N-1) dimensional projective space of the *universal* splitting functions:



 This limit can be physically measured inside high energy jets at the LHC.

Three-Point Correlator at Weak Coupling

- First non-trivial correlator: tree level three-point correlator in the collinear limit $G(z, \bar{z})$. [Chen, Luo, Moult, Yang, Zhang, Zhu]
- Turns out to have an elegant perturbative structure. e.g. in $\mathcal{N}=4$

$$G_{\mathcal{N}=4}(z) = \frac{1+u+v}{2uv}(1+\zeta_2) - \frac{1+v}{2uv}\log(u) - \frac{1+u}{2uv}\log(v) - \frac{1+u}{2uv}\log(v) - (1+u+v)(\partial_u + \partial_v)\Phi(z) + \frac{(1+u^2+v^2)}{2uv}\Phi(z) + \frac{(z-\bar{z})^2(u+v+u^2+v^2+u^2v+uv^2)}{4u^2v^2}\Phi(z) + \frac{(u-1)(u+1)}{2uv^2}D_2^+(z) + \frac{(v-1)(v+1)}{2u^2}D_2^+(1-z) + \frac{(u-v)(u+v)}{2uv}D_2^+\left(\frac{z}{z-1}\right)$$

• where Φ and D_2^+ are

$$\begin{split} \Phi(z) &= \frac{2}{z - \bar{z}} \left(\operatorname{Li}_2(z) - \operatorname{Li}_2(\bar{z}) + \frac{1}{2} \left(\log(1 - z) - \log(1 - \bar{z}) \right) \log(z\bar{z}) \right) \\ D_2^+(z) &= \operatorname{Li}_2(1 - |z|^2) + \frac{1}{2} \log(|1 - z|^2) \log(|z|^2) \end{split}$$

• Provides important perturbative data for the development of the lightray OPE.

Shape Dependence of Non-Gaussianities

• Multipoint correlators can be directly measured in high energy jets: Simple analytic functions for the *actual measured observable!*

• Non-Gaussianities inside high energy jets at the LHC:

CMS Open Data, $R_L \in (0.3, 0.4)$





Four Point Correlator

[Chicherin, Moult, Sokatchev, Yan, Zhu]

- Simple structure makes energy correlators a nice playground for exploration of *physical observables* in perturbation theory.
- Four point correlator computed in $\mathcal{N} = 4$ SYM by direct integration in parameter space, using simple form of $1 \rightarrow 4$ splitting function.





Kai Yan

- Compact result expressed in terms of weight three polylogarithms: much structure still to be explored.
- Would be interesting to extend to QCD using known $1\to 4$ splitting functions. [Del Duca, Duhr, Haindl, Lazopoulos, Michel]
- Can one push to higher points or make general statements?.

The Four Point Correlator

• Intricate view of correlations of energy flow. Access to OPE limits, spinning operators, ...



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The Four Point Correlator

Experimental tour de force to enable precision measurements of higher point correlators. $p_T > 100 \text{ GeV}$ Dipole $0.1 < R_L < 0.2$

Thanks to Simon Rothman and Phil Harris!

The Four Point Correlator

Experimental tour de force to enable precision measurements of higher point correlators.



Thanks to Simon Rothman and Phil Harris!

Energy Correlators: The Future of the Collider Frontier



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LHC Targets:

- Measurements on more complicated states:
 - Imaging the Quark Gluon Plasma











Resolving the Scales of the QGP



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Quark Gluon Plasma

• Heavy ion collisions provide an example of an extremely complicated asymptotic state, where we do not understand the microscopic dynamics that created it.





• Many physical effects: medium induced radiation, medium response, wake,

Multipoint Correlators

• Can we map out the space of "shapes" of models of the QGP?



• Ability to understand energy correaltors at strong coupling, weak coupling, in different theories, will prove essential.

Imaging the Quark Gluon Plasma

• QGP scales cleanly imprinted in two-point correlation.



[Andres, Dominguez, Holguin, Kunnawalkam Elayavalli, Marquet, Moult]

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Imaging the Wake

• Higher point correlators allow the "shape" of the medium response to be imaged.



• See talk by Krishna, Ananya, Hannah, Arjun.

Quark Gluon Plasma

• Excited for forthcoming measurements and progress!



12th International Conference on Hard and Electromagnetic Probes of High-Energy Nuclear Collisions

Sep 22–27, 2024 DEJIMA MESSE NAGASAKI Asia/Tokyo timezone

Overview	Contribution List
Scientific Program	
limetable	🗮 21/340 🔎 correlator 🛞 👅 🗋 🔗
Call for Abstracts	330. Jets: Substructures and energy-energy correlator © 9/26/24.11:15 AM
Registration/Apply for foung Scientist Support	Plenary Session VI
Contribution List	313. A fast evaluation method for higher point energy correlators and a new probe for medium properties
Announcomont	▲ Ankita Budhraja (Nicher)

Quark Gluon Plasma

- Motivates understanding of asymptotic observables in thermal or large charge states.
- Two recent approaches:
 - Heavy half-BPS operators in $\mathcal{N} = 4$:

$$O_K(x) = \operatorname{tr}[\phi^K(x)]\,, \qquad \quad \phi(x) = \sum_{I=1}^6 Y^I X^I\,.$$

[Chicherin, Korchemsky, Sokatchev, Zhiboedov]



• Large charge states using semi-classics:

$$\langle \mathcal{E}(\mathbf{n}_1)\mathcal{E}(\mathbf{n}_2)\rangle \approx \left(\frac{E}{\Omega_{d-2}}\right)^2 \left(1 + O(1/\Delta_Q)\right) \qquad (Q \to \infty).$$

[Firat, Monin, Rattazzi, Walters]



- The top quark mass is one of the most important parameters of the SM. e.g. electroweak vacuum stability/criticality, electroweak fits, etc.
- Need simple observables with top mass sensitivity that can be computed from first principles field theory.







• Extract the mass ratio between the W and top quark from the shape of the three-point correlator. $3 p_{pp \to tX} + \frac{1}{T(c,0,\infty)}$



[Holguin, Moult, Pathak, Procura, Schofbeck, Schwarz] See also: [Xiao, Ye, Zhu]



• Motivates precision calculations of correlators on top decays.

• Initial investigations illustrate has minor sensitivity to experimental systematics, and global event: successfully isolates dynamics of top decay.



Figure 2. The expected uncertainties of m_t (in % of $m_t = 171$ GeV) using E3C and m_{jet} distributions, at $\mathcal{L} = 36\,\mathrm{h^{-1}}$. The statistical uncertainties and a breakdown of the systematic uncertainties are shown.

Contamination 175-----175. _____ $pp \rightarrow tX$, Pythia 8.3, MPI+Had VZi/Zw [GeV] Zw [GeV] $pp \rightarrow tX$, Pythia 8.3, MPI+Had Jet Energy Scale Variation Color Reconnection M CP5CR1 17 17 CP5CR2 Dowr 4 173 173 , Muu 172Pvthiz ĝ 400 425 450 475 500 525 550 575 600 400 425 450 475 500 525 550 575 600 PT. iet GeV pr. iet [GeV]

DGLAP evoluti

[Xiao, Ye, Zhu]

• Motivates precision calculations of correlators on top decays, and further experimental investigation.

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Beyond Energy Flux



The Space of Detectors

- Details of the hadronization process are encoded in the quantum numbers (charge, flavor, ...): By definition, energy flux is insensitive!
- What is the space of detectors over which we can gain theoretical control?



Factorization

- More general observables can be calculated by combining factorization into universal matrix elements, with the Renormalization Group.
- Tremendous recent progress in understanding renormalization group evolution of functions characterizing correlations in the hadronization process (beyond DGLAP).



• Enables the calculation of correlations on energy flux carried by hadrons of specific quantum numbers: e.g. $\langle \Psi | \mathcal{E}_+(\hat{n}_1) \cdots \mathcal{E}_-(\hat{n}_k) | \Psi \rangle$

Charged Energy Flux

- Two examples of experimental interest involving electromagnetic charge:
 - $\langle \Psi | \mathcal{E}_{\mathsf{charged}}(\hat{n}_1) \mathcal{E}_{\mathsf{charged}}(\hat{n}_2) | \Psi \rangle$
 - $\langle \Psi | \mathcal{E}_{+}(\hat{n}_{1}) \mathcal{E}_{-}(\hat{n}_{2}) | \Psi \rangle$, $\langle \Psi | \mathcal{E}_{+}(\hat{n}_{1}) \mathcal{E}_{+}(\hat{n}_{2}) | \Psi \rangle$



• How far can we push into the confinement transition? Experimental measurements will be crucial.

Revisiting Old Data with New Resolution

• Reanalysis of ALEPH data has measured the two-point energy correlator $\langle \mathcal{E}(n_1)\mathcal{E}(n_2) \rangle$ on tracks with spectacular resolution. 10 $\mathbb{E}^{\text{ALEPH e'e, fs} = 91.2 \, \text{GeVeV}_{\text{CV}} + 12 \, \text{GeVeV}_{\text{CV}}}$



• Combined with precision track based calculations, opens up a new playground for precision studies of QFT.

Revisiting Old Data with New Resolution

 Detector based formalism enables calculation at NNLL-collinear and N³LL back-to-back on tracks!



NEECs and the EIC: Will be Covered by Others...









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Towards a Non-Perturbative Understanding



Confinement

Figure: Wenging Fan



Non-Perturbative Structure

- Formulating collider physics in terms of energy correlators allows us to reduce it to the study of the four-point correlator: $\langle OTTO \rangle$
- CFTs:
 - Planar $\mathcal{N} = 4$ (CFT bootstrap + Integrability)
 - 3d Ising (CFT bootstrap, Fuzzy sphere)
- Gapped Theories:
 - Form factor + S-matrix bootstrap



Figure 1: The Gauge Theory Bootstrap builds a (strongly coupled) bridge between the weakly coupled QCD at high energies and a weakly coupled EFT of pions at low energy

 Provides significant motivation to push towards understanding of analyticity properties of the four point correlator in gapped theories, and its simplification in the lightray limit. July 2024 57 / 61

Summary

- Recent progress understanding correlation functions of detector operators from explicit perturbative calculations, and the light-ray OPE.
- Multi-point correlators of lightray operators can be directly measured at the LHC: How can we best use them?

 Provides the opportunity to use theoretically beautiful objects to learn about the real world.



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Open Questions: "Phenomenologist"

- General structure of N-point energy correlators in perturbation theory? Efficient IBP for finite integrals in Feynman parameter space.
- What is the "light transform of the correlahedron"?
- What other matrix elements of lightray operators are well defined? Relation to IR finite S-matrices?
- What are the implications of asymptotic symmetry algebras? [Cordova, Shao] [Korchemsky, Sokatchev, Zhiboedov]
- What is the structure of the lightray OPE in non-conformal theories?
- How to incorporate data from (multiparticle) S-matrix bootstrap?
- What other systems have interesting detector operators and how can we experimentally access them?
- Correlators in thermal states? transport? confining theories? [Csaki, Ismail]
- How to interpolate between weak and strong coupling? Integrability?
- How to relate colliider measurements of detector operators to other Lorentzian singularities (Pomeron,...)?

[Caron-Huot, Kologlu, Kravchuk, Meltzer, Simmons Duffin]

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Open Questions: "Formal Field Theorist"

• Open Questions from David Simmons Duffin "Strings 2024" Talk:

Some open questions

- How do you measure a general detector operator at a collider?
- · Can we measure detectors in a condensed matter system?
- Can we formulate EFT running and matching for detectors? (Could help organize understanding of confinement effects in (*E* ··· *E*) [Jaarsma, Li, Moult, Waalewijn, Zhu '23; Csaki, Ismail '24].)
- What does the Chew-Frautschi plot look like at finite λ and finite N?
- Can we find positivity/rigidity conditions for light-ray operators? Can we formulate bootstrap conditions?
- What behavior in the deep Regge limit is possible? Transparency vs. chaos? [Stanford '15; Murugan, Stanford, Witten '17; Caron-Huot, Gobeil, Zahree '20]
- · Are other Lorentzian singularities described by other types of operators?
- Does "factorization theorem" = OPE? [Chen '23]
- What is the general form of the light-ray OPE?
- · How are light-ray operators and conformal line defects related?
- Do light-ray operators participate in interesting algebras? [Casini, Teste, Torroba '17; Cordova, Shao '18; Korchemsky, Sokatchev, Zhiboedov '21; Korchemsky, Zhiboedov '21; Faulkner, Speranza '24]

• Conformal Colliders have met the Collider Frontier!

