Bayesian methods and uncertainty quantification in heavy-ion physics

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Overview

- Heavy-ion collisions and multi-stage physics models
- Model emulators
- Results of Bayesian parameter estimation from different studies
- Quantifying theoretical uncertainties: Model discrepancy •
- Challenges

Heavy-ion collision



The Large Hadron Collider (LHC) in Europe (27 Km circumference)



The Relativistic Heavy Ion Collider in US (3.8 Km circumference)



CMS collision events: first lead ion collisions - CERN Document Server



Many stages of heavy-ion collisions



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Physics goals



Quark Gluon Plasma (QGP) phase: Hints of "fluid" formation. Two decades of research

- Equation of state? $P(T, \mu)$, $\epsilon(T, \mu)$ Taken from lattice results for hydro simulation
- Transport properties of formed QGP: Coefficient of shear viscosity: η Coefficient of bulk viscosity: ζ First principle calculation have large uncertainties Needs to be inferred from experiments

shear viscosity η reaction to a change of shape

bulk viscosity ζ reaction to a change of volume





Physics goals



Challenges: models are multi-stage, uncertain, and expensive

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Measurements

Particle yields for kaons, pions and protons:

$$\frac{dN_{\pi}}{dy}, \frac{dN_{K}}{dy}, \frac{dN_{p}}{dy}$$

Mean transverse energy: $dE_T/d\eta$

Elliptic, triangular and quadrangular flows: $v_2^{ch}\{2\}, v_3^{ch}\{2\}, v_4^{ch}\{2\}$

Mean transverse momentum fluctuations: δp_T

 $< p_T >$

 $dN_{ch}/d\eta$



Mean transverse momenta of kaons, pions, protons: $\langle p_T \rangle \pi, \langle p_T \rangle K, \langle p_T \rangle p$





Slide adapted from D. Liyanage

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Simulation models



Slide adapted from D. Liyanage

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Slide adapted from D. Liyanage

JETSCAPE SIMS calibration

D. Everett *et al.* 2010.03928, 2011.01430

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Model parameters



Slide adapted from D. Liyanage

JETSCAPE SIMS calibration

D. Everett et al. 2010.03928, 2011.01430

.3 0.4	a _{high}	
.3 0.4	.3	0.4
	.3	0.4



Simulation models: typical numbers

- The model parameter θ is a vector of length ℓ (\approx 15).
- The multi-stage model is stochastic: For a given θ , model simulations $\mathbf{y}_{sim}(\theta)$ results in distributions of d (~100) observables of interest. luncertain
- Model is computationally expensive: ≈ 1000 CPU hours needed for a given parameter setting θ (\approx 3000 events for each θ). Many model observables still have large statistical uncertainty even after \approx 1000 CPU hours.
- Posterior inference requires model simulations at $> 10^6$ samples of parameter space. Inference is out of reach without emulation.
- Need for fast and accurate model surrogates to perform any Bayesian study.

expensive





Gaussian Process based model emulators

- We want emulators to "interpolate" in the ℓ (\approx 15) dimensional parameter space. •
 - Consider the mean values of the d (\approx 100) distributions for each θ_i : $\bar{y}_{sim}(\theta_i)$.
 - For $n \approx 1000$ training set $\{\theta_1, \dots, \theta_n\}$ (LHS), define a $n \times d$ matrix: $\Xi \equiv \{\bar{\mathbf{y}}_{sim}(\theta_1), \dots, \bar{\mathbf{y}}_{sim}(\theta_n)\}\$
 - Standardize dataset by removing mean and scaling to unit variance for all d-distributions: $\Xi \rightarrow \tilde{\Xi}$
 - Principal Component Analysis \rightarrow Reduce dimensionality of dataset 0
 - Transform by doing PCA and keep p < d principal components. In most applications, $p \ll d$ is sufficient to describe almost all the variance in the original dataset.
- Train *p* independent Gaussian process corresponding to the *p* reduced observables (means). Each Gaussian process is ℓ (\approx 15) dimensional.





Emulator: Challenges

- GP training does not take into account the inherent stochastic uncertainty.
- Principal Component Stochastic Kriging (PCSK) takes into account the variance to some extent.
- We need a way to reduce dimensionality of dataset consisting of distributions.
- Efficient emulator training through Transfer Learning.
- Current status: model emulators are fast, but not precise (Next slide ——)

BAND SURMISE package M. Plumlee, Ö. Sürer, S. Wild, M. Chan https://surmise.readthedocs.io

D. Liyanage, Ö. Sürer *et al.* 2302.14184 Hendrik Roch, Syed Afrid Jahan, Chun Shen: 2405.12019

D. Liyanage, et.al. (JETSCAPE Collaboration) 2201.07302





Comparison of different emulators

- For Au-Au collision for 3 different $\sqrt{S_{NN}}$
- Overall 544 experimental observables
 - Error in prediction (mean):

$$\mathcal{E} \equiv \sqrt{\left\langle \left(\frac{\text{prediction} - \text{truth}}{\text{truth}}\right)^2 \right\rangle}$$



PCGP, PCSK

BAND SURMISE package M. Plumlee, Ö. Sürer, S. Wild, M. Chan https://surmise.readthedocs.io



separated by black lines. All emulators are trained with the same 970 LHD points.

970 training sets. 30 validation sets



Comparison of different emulators Hendrik Roch, Syed Afrid Jahan, Chun Shen: 2405.12019

- For Au-Au collision for 3 different $\sqrt{S_{NN}}$
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Normalized residuals:
$$\mathcal{H} \equiv \ln\left(\sqrt{\left\langle \left(\frac{\text{prediction} - \text{truth}}{\text{prediction uncertainty}}\right)^2\right\rangle}\right)$$

For an accurate prediction, we expect the values of $\mathcal{E} \to 0$ and $\mathcal{H} \to 0$. In the case where $\mathcal{H} > 0$, the emulator and $n \to 0$. In the case share too small compared to the gives uncertainties that are too small compared to the χ actual error away from the true values; when $\mathcal{H} < 0$, the returned uncertainty estimates are too conservative.

Better metric for comparison can be KL divergence



lines separate different training sets. All emulators are trained with the same 970 LHD points.

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Best fit (MAP) output from the calibrated Models:

Scikit GP RBF kernel





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D. Liyanage et al., 2302.14184

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D. Everett et al. (JETSCAPE Collaboration), Phys. Rev. C 103, 054904 (2021)

D. <u>Liyanage</u> et al., <u>2302.14184</u>

Best fit (MAP) output from the calibrated Models:

- MAP predictions for VAH+PTMA are in better agreement with SIMS+14-moment model.
- Our aim Correct inference of physical parameters.

GRAD

SMASH

D. Everett et al. (JETSCAPE Collaboration), Phys. Rev. C 103, 054904 (2021)

D. Liyanage et al., 2302.14184

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JETSCAPE SIMS model parameters

≈15 model parameters

FIG. 10. The posterior for Grad (blue) and Chapman-Enskog (red) viscous corrections for select parameters related to the initial state, prehydrodynamic evolution, and switching temperature. The histograms on the diagonal are the marginal distributions for each parameter, D. Everett et al. (JETSCAPE Collaboration), Phys. Rev. C 103, 054904 (2021)

Two different viscous corrections: 14moment Grad, CE

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≈15 model parameters

D. Everett et al. (JETSCAPE Collaboration), Phys. Rev. C 103, 054904 (2021)

- Two different viscous corrections: 14moment Grad, CE
- Different posterior distributions of the parameters for the two different schemes.

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- Two different viscous corrections: 14moment Grad, CE
- Different posterior distributions of the parameters for the two different schemes.
- Correlation between parameters also differ.

Note: rotate red contours by 90 degrees for comparison with blue.

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VAH+PTMA model parameters D. Liyanage *et al.*, 2302.14184

Pb - Pb at 2.76 TeV ≈100 Observables ≈15 model parameters

• Posterior distributions of the common parameters differ substantially from the JETSCAPE SIMS model

Tension between different studies

90% credible intervals

JETSCAPE SIMS calibration

D. Everett et al. 2010.03928, 2011.01430

Three different models: Grad, CE, PTB

 η/s seems to be in agreement

 ζ/s varies for different studies

Viscous Anisotropic Hydrodynamics Model

M. McNelis et al. 2101.02827 D. Liyanage, O. Surer, et al. 2302.14184

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Accounting for theoretical uncertainties: Model discrepancy

- uncertainty.

• All models are approximations of reality and should only be used within their valid domains to fit experimental data. Extending a model beyond its intended scope leads to incorrect parameter values, making them mere fitting variables. Therefore, it is crucial to account for the theory's

• Consideration of theoretical uncertainties for the complex multi-stage heavy-ion models are beyond current theoretical capabilities. We need a statistical framework to model this uncertainty.

Accounting for theoretical uncertainties: Model discrepancy

- uncertainty.
- Possible framework: GP based model discrepancy by O'Hagan et. al.

$$y(x_i) = \eta(x_i, \theta) + \delta(x_i) + \delta(x_i) + y_i + Model + \delta(x_i) + \delta($$

Physical observation

Accounts for discrepancy between model and truth

Model $\delta(x_i)$ as a gaussian process. Choice of covariance kernel motivated from the physics.

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• Consideration of theoretical uncertainties for the complex multi-stage heavy-ion models are beyond current theoretical capabilities. We need a statistical framework to model this uncertainty.

Observation error

M. Kennedy, A. O'Hagan, Bayesian calibration of computer models, https://doi.org/10.1111/1467-9868.00294

D. Higdon, M. Kennedy, et. al., https://doi.org/10.1137/S1064827503426693

Work in progress...

Challenges

Expensive heavy-ion model simulations demands fast and accurate model emulators.

Quantifying theoretical uncertainties is a *necessity* for correct parameter estimation.

Thank You!

O. Soloveva, D. Fuseau, J. Aichelin, E. Bratkovskaya, 2011.03505

Large Uncertainties

Transport coefficients from different calculations

Valeriya Mykhaylova Thesis

QCD phase diagram

Foka, Panagiota et al - arXiv:1702.07233

Initial Energy Deposition (TRENTO)

Parametrization for energy deposition at proper time $\tau=0^+$

p = +1 p = 0Pb-Pb @ 2.76 TeV w = 0.4fm arXiv:1904.08290v1 p = -1

Slide adapted from D. Everett

	Parameter	Sym
re	educed thickness	p
	nucleon width	W
	energy normalization	N
	multiplicity fluctuation	σ_k
m	in. distance btw. nucleons	$d_{ m mi}$

 σ_k controls 'contrast'

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Pre-hydro (freestreaming)

Freestream massless particles: $f(t, \mathbf{x}; \mathbf{p}) = f(t_0, \mathbf{x} - \mathbf{v}\Delta t; \mathbf{p})$

Take initial momentum-distribution isotropic in transverse plane

$$T^{\mu\nu}(\tau_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} T^{\tau\tau}(\tau_0, \mathbf{x}_T - \mathbf{x}_h) = \frac{\tau_0}{\tau_h} \int \frac{d\phi}{2\pi} \hat{p}^{\mu} \hat{p}^{\nu} \hat{p}^{\mu} \hat$$

Slide adapted from D. Everett

Parameter	Sym
ref. proper time	$ au_R$
energy dependence	α

 $\Delta \tau = \tau_R \left(\frac{\langle \epsilon \rangle}{\epsilon_R} \right)^{\alpha}$

$-\hat{\mathbf{p}}_T\Delta \tau; \mathbf{p}_T$

Viscous Hydro

The viscosity of QGP:

$\tau_{\Pi}\dot{\Pi} + \Pi = -\zeta\theta + \dots$ $\tau_{\pi}\dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = 2\eta\sigma^{\mu\nu} + \dots$

Quantify transport properties : shear and bulk viscosities

Slide adapted from D. Everett

Parameter	Symbol
temperature of kink	T_{η}
shear at kink	$(^{\eta}/_{S})_{kink}$
shear low-T slope	$a_{\rm low}$
shear high-T slope	a_{high}
temperature of bulk peak	Tζ
bulk at peak	$(\zeta/_{S})_{\max}$
bulk width	Wζ
bulk skewness	λ
shear relax. time	b_{π}

Viscous Hydro

Viscosity parameterizations:

Parameter	Symbo
temperature of kink	T_{η}
shear at kink	$(^{\eta}/_{S})_{\mathrm{kin}}$
shear low-T slope	$a_{\rm low}$
shear high-T slope	$a_{ m high}$
temperature of bulk peak	Τ _ζ
bulk at peak	$(\zeta/s)_{\rm mat}$
bulk width	Wζ
bulk skewness	λ
shear relax. time	b_{π}

Emulator: challenges

 $\langle \text{emulator}/\text{model} \rangle$, $n_{\text{design}} = 2000$, $n_{\text{val}} = 200$

the experimental error of the data point (blue), the average statistical error of our model (green) and the range of model predictions given our prior range, defined as the standard deviation over the mean (orange). Ideally the

Trajectum G. Nijs *et al.* 2010.15134

